

# Sigma term and strange content of the nucleon

[Preliminary results with a subset of the Budapest-Marseille-Wuppertal ensembles]

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## Definitions

$$\sigma_{\pi N} = \hat{m} \langle N(p) | (\bar{u}u + \bar{d}d)(0) | N(p) \rangle$$

$$\sigma_{ksN} = m_s \langle N(p) | (\bar{s}s)(0) | N(p) \rangle$$

$$y = \frac{2 \langle N | \bar{s}s | N \rangle}{\langle N | \bar{u}u + \bar{d}d | N \rangle}$$

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## Feynman-Hellman theorem

$$\frac{\partial E(\lambda)}{\partial \lambda} = \langle \lambda | \frac{\partial H(\lambda)}{\partial \lambda} | \lambda \rangle$$

If applied to QFT

$$\hat{m} \frac{\partial M_N}{\partial \hat{m}} = \hat{m} \langle N | \frac{\partial H_{QCD}}{\partial \hat{m}} | N \rangle = \sigma_{\pi N}$$

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## Important for

- Hadron spectrum
- The strangeness content of the nucleon ( $\rightarrow$  Detection of dark matter).
- The quark mass ratio  $m_s/\hat{m}$
- $\pi - N$  and  $K - N$  scattering amplitudes

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# Simulation details

- Gauge Action: Symanzik  $\mathcal{O}(a^2)$  improved.
- Fermion Action: 2 + 1 Wilson fermions. Tree level  $\mathcal{O}(a)$  improved.
- 6 levels of stout smearing.
- Algorithm: HMC (for 2 degenerate quarks), RHMC (for the s quark).
- Volume: Spatial size L big so that data can be corrected with (small) finite volume effects. ( $M_\pi L \geq 4$ ).
- Three values of the lattice spacing. ( $a \approx 0.124fm, 0.083fm, 0.063fm$ ).
- Pion masses range:  $190MeV \lesssim M_\pi \lesssim 460MeV$ .

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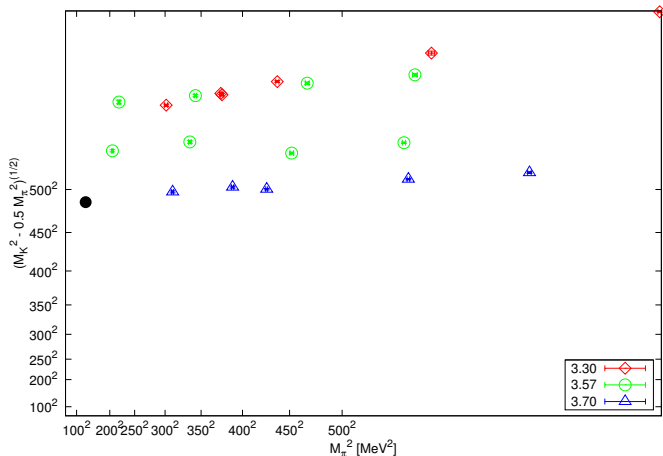
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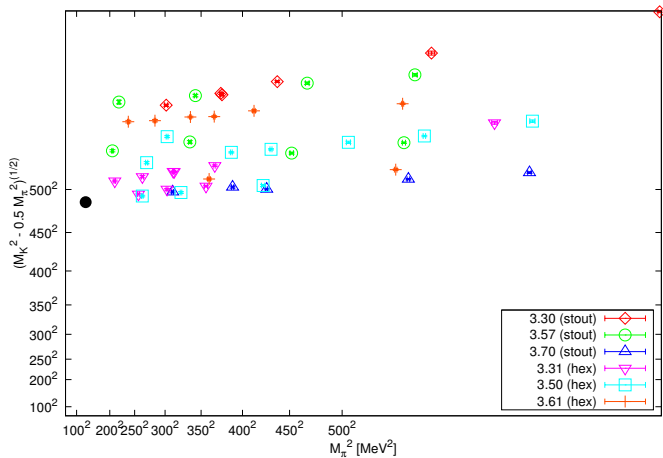
Figure: Dataset used in [S. Dürr et al Science 322 (2008)], [S. Dürr et al Phys.Rev.D 81 (2010)]



BMW also has ensembles generated with 2 steps of HEX smearing. [T. Kurth and C. Hoelbling talks].

- Three values of the lattice spacing ( $a \approx 0.11\text{fm}, 0.085\text{fm}, 0.07\text{fm}$ )

Figure: Landscape



# Chiral extrapolation.

Having data  $190\text{MeV} < M_\pi < 460\text{MeV}$  we need to extrapolate our data to the physical point.

# Chiral extrapolation.

## Taylor based extrapolations

The nucleon mass has an expansion around a regular point

$$(M_{\pi}^{\text{exp}})^2 = \frac{1}{2}[(M_{\pi}^{\Phi})^2 + (M_{\pi}^{\text{max}})^2]; M_{ks}^2 = M_K^2 - \frac{1}{2}M_{\pi}^2$$

$$M_N = \sum_{i=1}^{N_{\pi}} \alpha'_i \left[ M_{\pi}^2 - (M_{\pi}^{\text{exp}})^2 \right]^i + \sum_{i=1}^{N_{\pi}} \beta'_i \left[ M_{ks}^2 - (M_{ks}^{\Phi})^2 \right]^j$$

In particular at the physical point

$$M_N^{\Phi} = \sum_{i=1}^{N_{\pi}} \alpha'_i \left[ (M_{\pi}^{\Phi})^2 - (M_{\pi}^{\text{exp}})^2 \right]^i + \sum_{i=1}^{N_{\pi}} \beta'_i \left[ (M_{ks}^{\Phi})^2 - (M_{ks}^{\text{exp}})^2 \right]^j$$

# Chiral extrapolation.

## Taylor based extrapolations

Analytical expressions in quark masses.

$$(aM_N) = (aM_N^\Phi) \left\{ 1 + \sum_{i=1}^{N_\pi} \alpha_i \left[ \left( \frac{aM_\pi}{aM_N^\Phi} \right)^2 - \left( \frac{M_\pi^\Phi}{M_N^\Phi} \right)^2 \right]^i + \sum_{j=1}^{N_K} \beta_j \left[ \left( \frac{aM_{ks}}{aM_N^\Phi} \right)^2 - \left( \frac{M_{ks}^\Phi}{M_N^\Phi} \right)^2 \right]^j \right\}$$

# Chiral extrapolation.

## Taylor based extrapolations

Pade expressions to control higher order terms

$$(aM_N) = (aM_N^\Phi) \left\{ 1 + \sum_{i=1}^{N_\pi} \alpha_i \left[ \left( \frac{aM_\pi}{aM_N^\Phi} \right)^2 - \left( \frac{M_\pi^\Phi}{M_N^\Phi} \right)^2 \right]^i + \sum_{j=1}^{N_K} \beta_j \left[ \left( \frac{aM_{ks}}{aM_N^\Phi} \right)^2 - \left( \frac{M_{ks}^\Phi}{M_N^\Phi} \right)^2 \right]^j \right\}^{-1}$$

# Chiral extrapolation.

## Three flavour BChPT

For a better control of the chiral extrapolation, three flavour covariant ChPT formulas have been computed [\[See S. Dürr poster\]](#).



# Continuum extrapolation

- Formally the two different actions receive different cutoff corrections.
- Although our actions formally have  $\mathcal{O}(a\alpha_s)$  corrections, data suggests cutoff effects quadratic in  $a$  [BMW. Phys.Rev.D. 79 (2009)].

## Parametrisation of cutoff effects

Each of the “chiral LEC” get an  $a$  dependence.

$$p \longrightarrow p(1 + \mathcal{O}(a^n))$$

or

$$p \longrightarrow p(1 + \mathcal{O}(\alpha_s a^n))$$

## No cutoff effects

Data consistent with no cutoff effects!.

# Finite volume

In all our ensembles the bound  $M_\pi L \gtrsim 4$  is maintained.

## Parametrisation of finite volume effects

We parametrise the finite volume effects as

$$\propto e^{-M_\pi L}$$

## No finite volume effects

Data consistent with no finite volume effects.

# Excited states

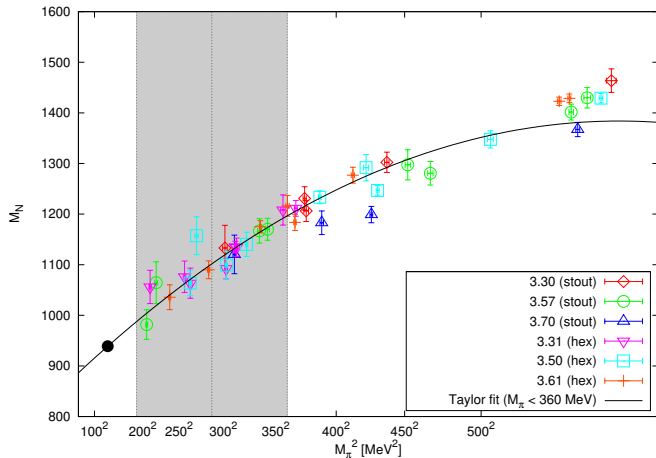
For each value of beta we use 2-3 values for the fitting interval strongly dominated by the ground state.

## Estimate excited states contributions

Use the total of 144 possible combinations of fitting intervals.

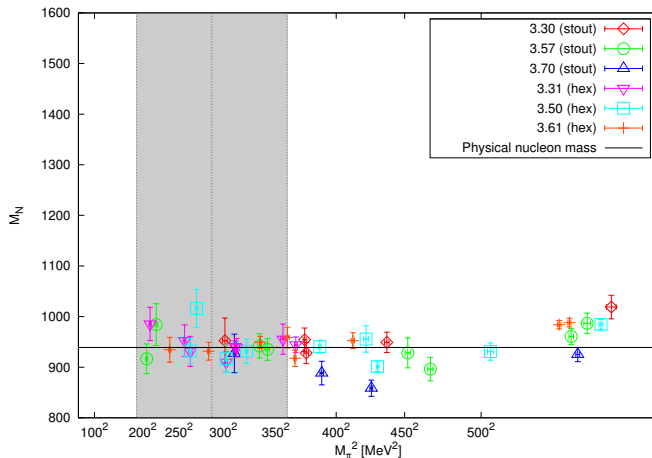
# A typical fit

Figure: A Taylor fit with  $M_\pi < 360\text{MeV}$ .  $\chi^2/dof = 17.3/17$



# A typical fit

Figure: Residuals after the fit



## Systematics *a la* BMW

- 2 formulas to extrapolate to the physical point: Taylor, Padé.
- 2 pion mass cuts:  $M_\pi < 360\text{MeV}$  and  $M_\pi < 460\text{MeV}$ .
- No cutoff effects or finite volume corrections (consistent with our data).
- 144 fitting intervals to estimate excited states contributions.

This amounts a total of  $144 \times 2 \times 2 = 576$  analysis.

## Central idea of the analysis.

Use all of them, and let data decide.

- use each of the 576 values, weighted by the quality of the fit, to construct a distribution of values.
- The final result is the median.
- The systematic error are given by the 16th and 84th percentiles.
- All the process is bootstrapped and the deviation over bootstraps samples taken as the statistical error.
- The total error is the sum by quadratures of the systematic and the statistical error.

## Following the procedure

- $\sigma_{\pi N} = 49(10)_{\text{stat}}(11)_{\text{sys}}$  MeV
- $\sigma_{K^*N} = 49(37)_{\text{stat}}(26)_{\text{sys}}$  MeV
- $y = 0.08(7)_{\text{stat}}(4)_{\text{sys}}$



Figure: Distribution of values

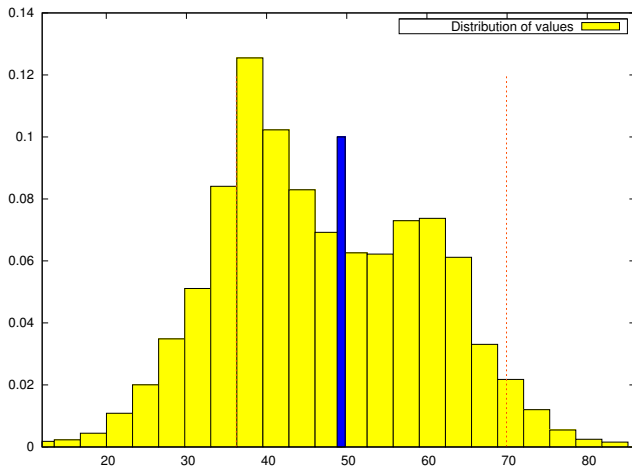
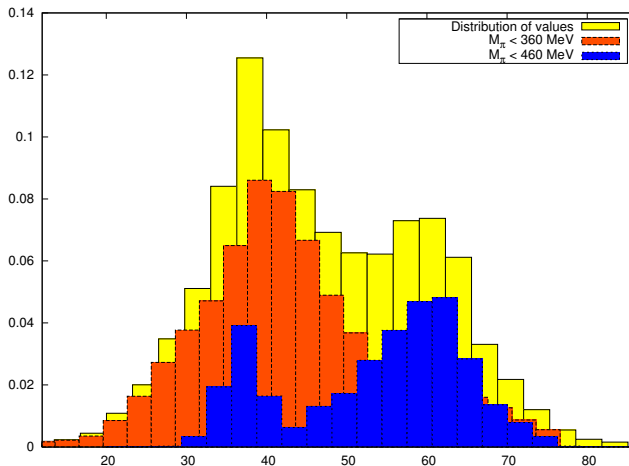


Figure: Different  $M_\pi$  cuts is the main source of systematic error



$M_\pi = 134$  MeV just by chance!!

## Imaginary world

Repeat the analysis in an imaginary world with a heavier  $m_{ud}$ .

- $M_\pi = 200$  MeV
- $M_K = 505$  MeV
- $M_N = 1000$  MeV

Figure: Distribution of values in imaginary world.

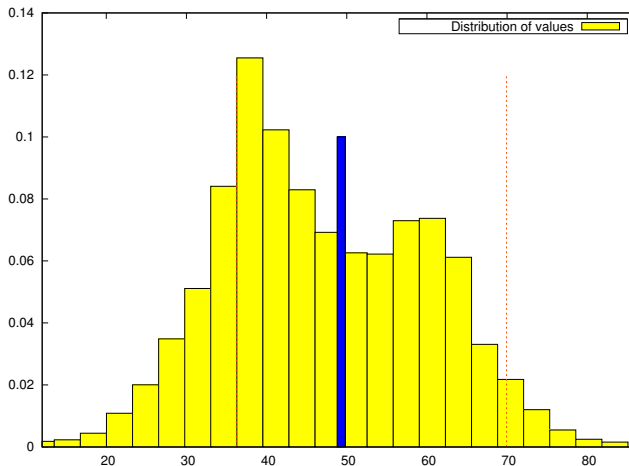
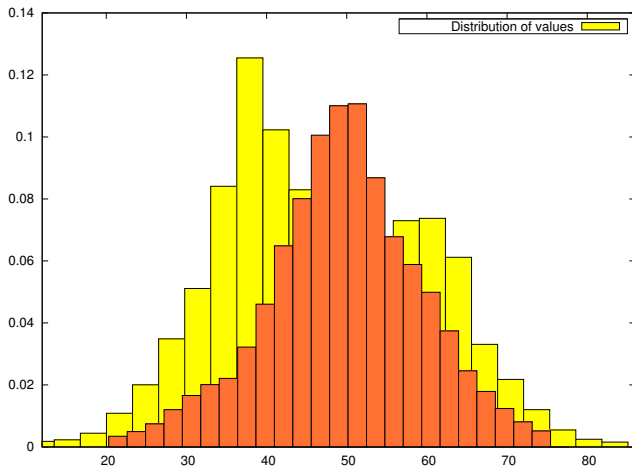


Figure: Distribution of values in imaginary world.



## Sigma term in imaginary world

$$\sigma_{\pi N} = XX(7)_{\text{stat}}(6)_{\text{sys}}$$

Both systematic and statistical errors reduced roughly by a 70%.

# Conclusions

- Using a sub-sample of the BMW ensemble:
  - ▶ 2 (similar) actions.
  - ▶ 6 values of  $\beta$ .
  - ▶ Large volumes ( $M_\pi L \gtrsim 4$ ).
  - ▶ Data in the range  $190\text{MeV} \lesssim M_\pi \lesssim 460\text{MeV}$ .
- We have presented preliminary results for the nucleon sigma term, obtaining  $\sigma_{\pi N} = 49 \pm 15 \text{ MeV}$ .
- For a better control over chiral extrapolation, and for a better sensitivity to the strange content, still need to work in the three flavour CBChPT [see S.Dürr poster].
- Computing sigma terms of all the octet.

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