

Pairs of massless quarks on the lattice from staggered fermions

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Numerically investigated by Ph. de Forcrand & collaborators (next talk).

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 - ▶ Describe massless fermions in staggered framework
 - ▶ Exact zero-modes, index theorem, η' mass...

Usual overlap construction

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Overlap Dirac operator [Neuberger, PLB 1998]:

$$D_{ov} = \frac{\rho}{a} \left(1 + \frac{D_W - M}{\sqrt{(D_W - M)^\dagger (D_W - M)}} \right)$$

has exact zero-modes, index theorem, axial anomaly etc.

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- ▶ has purely imaginary spectrum
- ▶ \Rightarrow would-be zero-modes do not become real negative eigenvalue modes for $D_{st} - M$
- ▶ \Rightarrow Overlap construction based on $D_{st} - M$ fails:
e.g. no exact zero-modes or index theorem

Reason to still hope for a staggered overlap construction

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Consider free field staggered Dirac operator in flavor space rep
[Gliozzi, Kluberg-Stern et al.]:

$$D_{st} = (\gamma_\mu \otimes 1) \nabla_\mu + \frac{1}{2} (\gamma_5 \otimes \gamma_5 \gamma_\nu) \Delta_\nu$$

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Then $\tilde{\gamma}_5 = \gamma_5$ and

$$i(\gamma_5 \otimes 1) D_{st} = (\tilde{\gamma}_\mu \otimes 1) \nabla_\mu + \frac{1}{2} (1 \otimes \tilde{\gamma}_\nu) \Delta_\nu \equiv D_{stW}$$

– has structure of Wilson-Dirac operator

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⇒ Can expect the spectral flow of $H_{st}(m)$ in the interacting case determined by gauge field topology.

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\Rightarrow Cannot simply replace $\gamma_5 \rightarrow \Gamma_5$ in overlap construction.

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Go back and start the staggered overlap construction from $D_{st} - M$, but with

$$M = \frac{\rho}{a} \Gamma_{55} \Gamma_5$$

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$$\Gamma_{55} \chi(x) = (-1)^{n_1+n_2+n_3+n_4} \chi(x) \quad , \quad x = an \in a\mathbf{Z}^4$$

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$\Rightarrow M$ is a “flavor-chiral” mass term

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\Rightarrow get 2 physical fermion flavours!

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Take the staggered overlap Dirac operator to be

$$D_{sov} = \frac{\rho}{a} \left(1 + \frac{D_{st} - M}{\sqrt{(D_{st} - M)^\dagger (D_{st} - M)}} \right) \quad M = \frac{\rho}{a} \Gamma_{55} \Gamma_5$$

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The “ $\Gamma_5^2 \neq 1$ problem” is hereby solved.

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\therefore The exact flavored chiral symm of D_{st} has been turned into an exact *unflavored* GW chiral symm of D_{sov} .

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Its spectral flow is already known to be determined by gauge field topology, giving $2Q$.

[D.A., PRL (2010)]

Theoretical justification for D_{SOV} (continued)

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Note that the staggered overlap operator,

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where the hermitian operator \tilde{H}_{st} is

$$\tilde{H}_{st} = \Gamma_{55} (D_{st} - M) = \Gamma_{55} D_{st} - \frac{\rho}{a} \Gamma_5$$

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Theorem:

$$\tilde{H}_{st}(m) = \Gamma_{55} D_{st} - m\Gamma_5 \quad \text{and} \quad H_{st}(m) = iD_{st} - m\Gamma_5$$

have the same eigenvalue spectrum.

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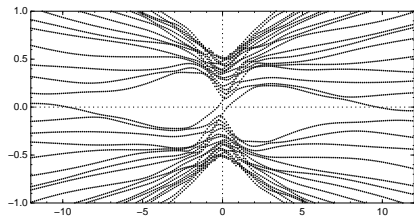
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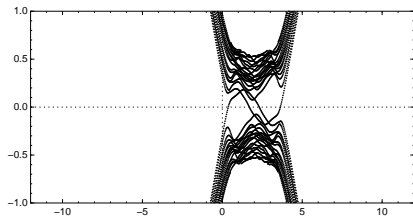
$\Rightarrow D_{sov}$ has exact zero-modes satisfying 2-flavor index theorem.

Illustration: spectral flow in a $U(1)$ background with $Q = 1$ in 2 dimensions

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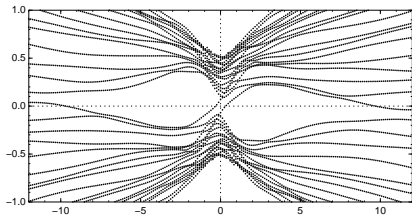


staggered

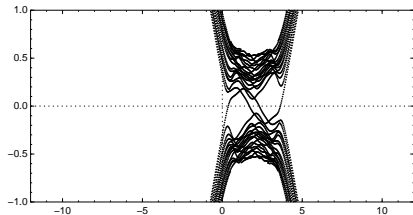


Wilson

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staggered



Wilson

Results in 4 dimensions will be shown in Ph. de Forcrand's talk.

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- ▶ Again 2 physical (almost massless) flavors
- ▶ Gives truncation of staggered overlap, just like in Wilson case [Neuberger, PRD (1998)].
- ▶ Completely unrelated to a previous proposal for staggered domain wall fermions by G. Fleming and P. Vranas (PRD, 2002).

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- ▶ A good option for calculating the η' mass in Lattice QCD?