

Nucleon and $N^*(1535)$ Distribution Amplitudes

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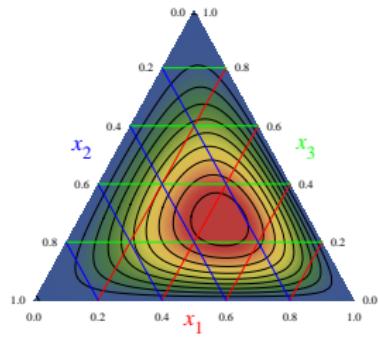


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¹V. M. Braun, S. Collins, M. Göckeler, T. R. Hemmert, R. Horsley, T. Kaltenbrunner, A. Lenz, Y. Nakamura, D. Pleiter, P. E. L. Rakow, J. Rohrwild, A. Schäfer, G. Schierholz, H. Stüben, N. Warkentin, P. Wein, J. M. Zanotti
[QCDSF collaboration]

Motivation

- ▶ **Wanted:** Hadron wave function
- ▶ Experiment:
 - ▶ Form factor measurements at JLab (especially with the 12 GeV upgrade), FAIR, ...
 - ▶ Light-Cone Sum Rules relate form factor data to distribution amplitudes
- ▶ Theory:
 - ▶ Direct calculation of distribution amplitudes from Lattice QCD



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Outline

Nucleon Distribution Amplitudes
in the Continuum ...
... and on the Lattice

Parity Separation

Results

Normalization Constants
Distribution Amplitudes

Conclusions and Outlook



Distribution Amplitudes in the Continuum

- ▶ In leading twist

- ▶ with transverse momentum components integrated out
- ▶ the nucleon wave function can be written as

$$|N, \uparrow\rangle = f_N \int \frac{[dx]\varphi(x_i)}{2\sqrt{24x_1x_2x_3}} \{ |u^\uparrow(x_1)u^\downarrow(x_2)d^\uparrow(x_3)\rangle - |u^\uparrow(x_1)d^\downarrow(x_2)u^\uparrow(x_3)\rangle \}$$

- ▶ where

- ▶ x_i : longitudinal momentum fractions
- ▶ $\int [dx] = \int_0^1 dx_1 dx_2 dx_3 \delta(1 - x_1 - x_2 - x_3)$
- ▶ f_N : Leading-twist normalization constant, "Wave function at the Origin"
- ▶ $\varphi(x_i)$: Nucleon Distribution Amplitude

- ▶ in next-to-leading twist

- ▶ λ_1, λ_2 : Next-to-leading twist normalization constants



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Distribution Amplitudes on the Lattice

- ▶ On the lattice, calculate correlators of the form

$$\langle \mathcal{O}(x)_{\alpha\beta\gamma} \bar{\mathcal{N}}(y)_\tau \rangle$$

- ▶ where

- ▶ $\bar{\mathcal{N}}$ is a smeared nucleon interpolator
- ▶ \mathcal{O} is a local three-quark operator with up to two derivatives

- ▶ to determine moments of the distribution amplitude

$$\varphi^{lmn} = \int [dx] x_1^l x_2^m x_3^n \varphi(x_1, x_2, x_3)$$

- ▶ Limited to first and second moments ($l + m + n \leq 2$) due to
 - ▶ Mixing with lower dimensional operators for higher moments
 - ▶ Need for higher momenta → more noise



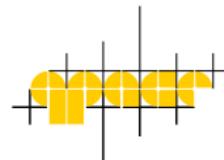
Overview of lattices

$N_f = 2$ clover Wilson fermions

κ	m_π / MeV	Size	# Configurations
$\beta = 5.29, a = 0.0753 \text{ fm}$			
0.13590	627	$24^3 \times 48$	901
0.13620	407	$24^3 \times 48$	850
0.13632	282	$32^3 \times 64$	578 + more in progress
0.13632	271	$40^3 \times 64$	in progress
0.13640	170	$40^3 \times 64$	coming soon
$\beta = 5.40, a = 0.0672 \text{ fm}$			
0.13610	648	$24^3 \times 48$	687
0.13625	558	$24^3 \times 48$	1180
0.13640	451	$24^3 \times 48$	1037
0.13660	233	$48^3 \times 64$	coming soon

[QCDSF collaboration]

New configurations are being generated on



Parity separation

$$32^3 \times 64$$

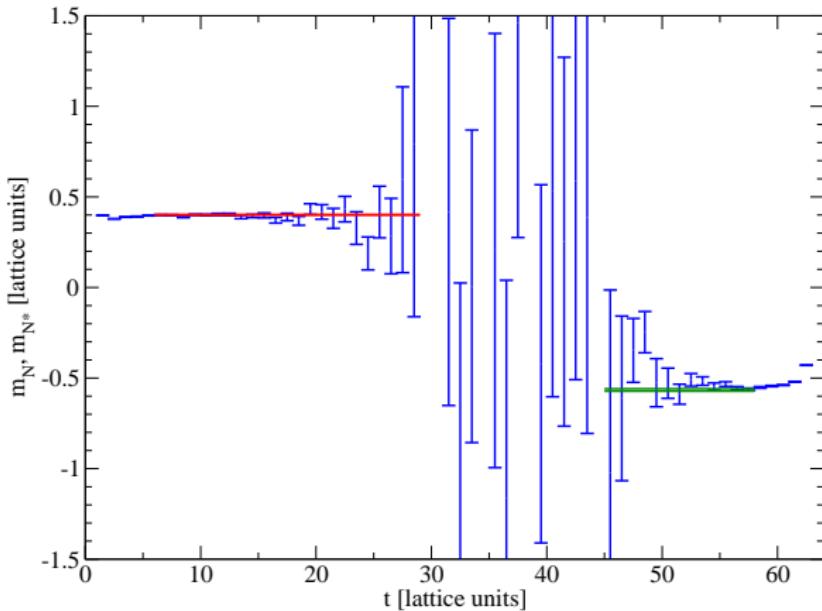
$$a = 0.0753 \text{ fm}$$
$$a^{-1} = 2.620 \text{ GeV}$$

$$m_\pi = 282(2) \text{ MeV}$$

$$m_\pi L = 3.44$$

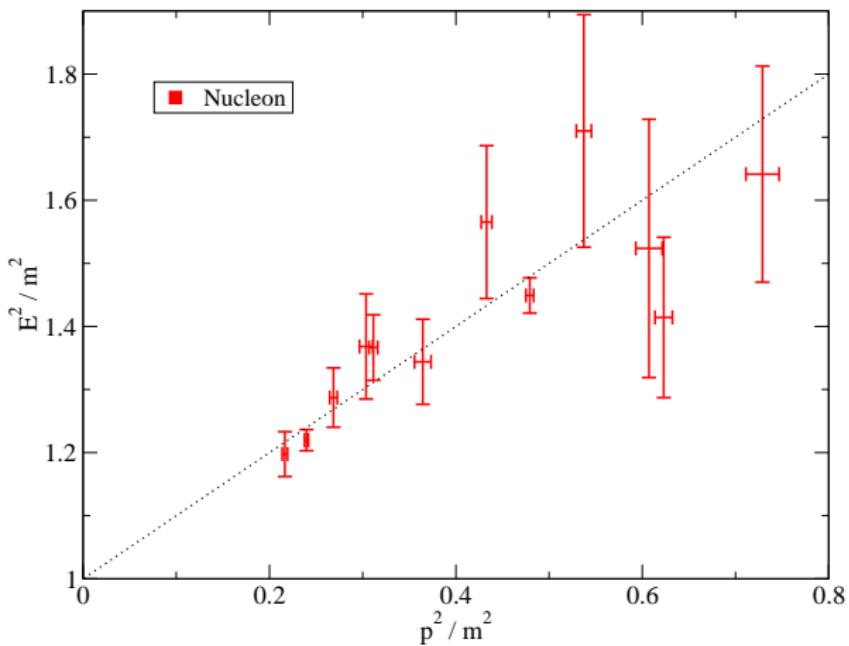
$$m_N = 1051(5) \text{ MeV}$$

$$m_{N^*} = 1482(17) \text{ MeV}$$



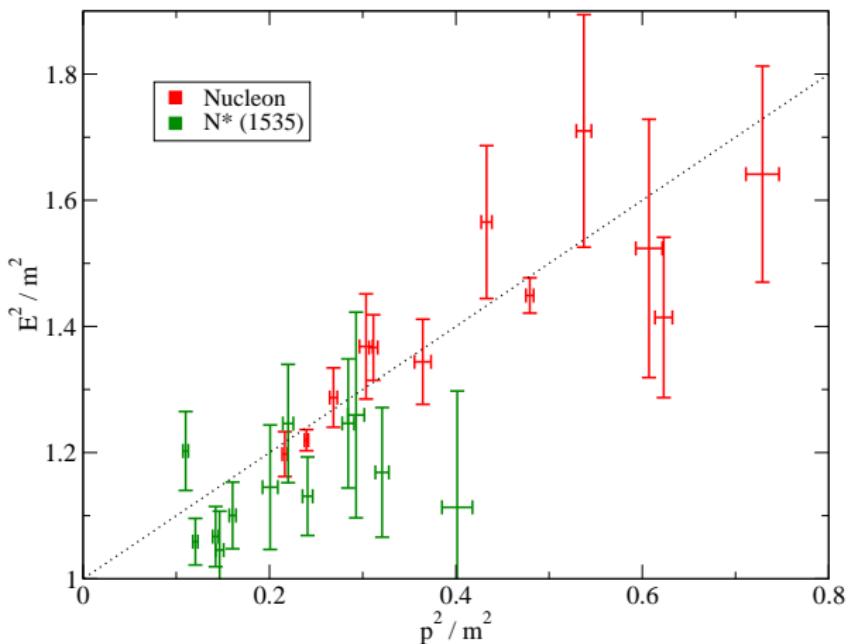
The generalized Lee-Leinweber parity projector, $\frac{1}{2}\Gamma (1 + \frac{m}{E}\gamma_4)$, where Γ is a suitable product of γ_i matrices, works well for all operator-momentum combinations that are used here

Dispersion Relation



- ▶ According to Einstein, $\frac{E^2}{m^2} = 1 + \frac{p^2}{m^2}$ (dotted line)
- ▶ All results preliminary, statistical errors only

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Wave functions at the origin

$$32^3 \times 64$$

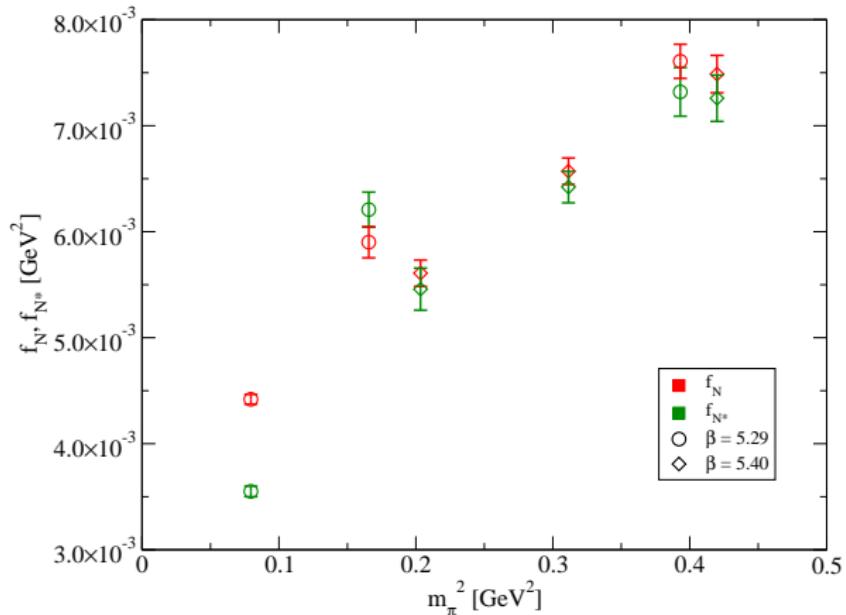
$$a = 0.0753 \text{ fm}$$

$$m_\pi = 282(2) \text{ MeV}$$

$$m_\pi L = 3.44$$

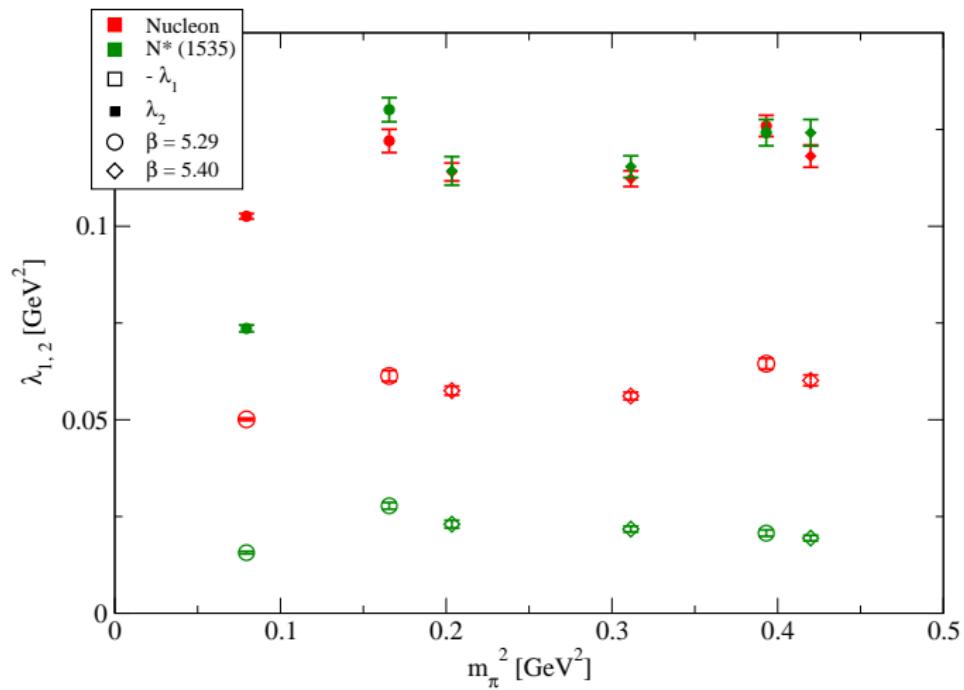
$$f_N = 4.42(5) \cdot 10^{-3} \text{ GeV}^2$$

$$f_{N^*} = 3.55(5) \cdot 10^{-3} \text{ GeV}^2$$



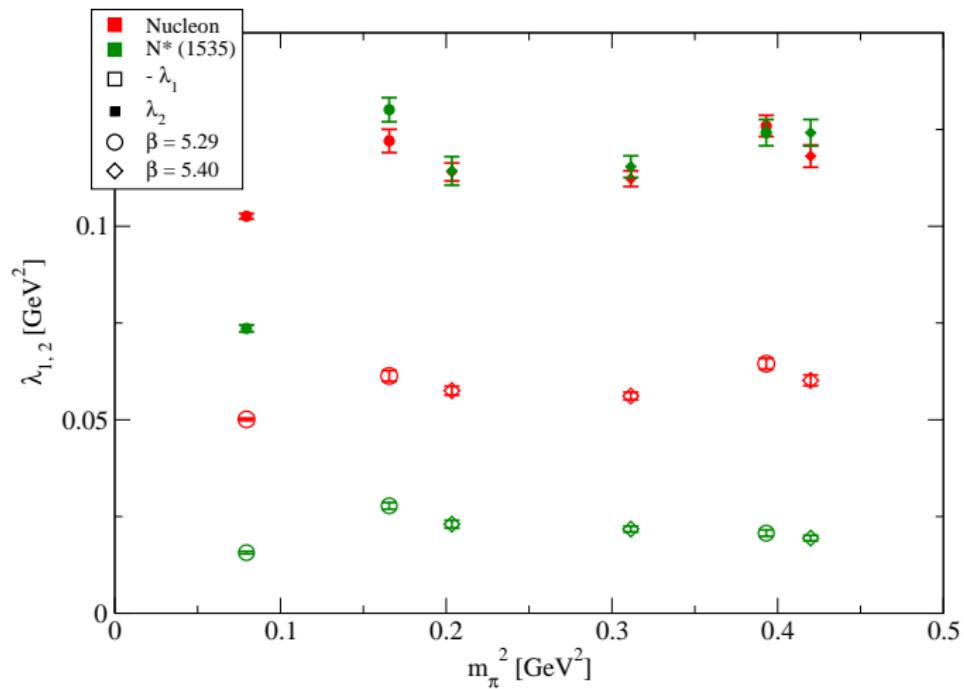
All results preliminary, statistical errors only, no chiral extrapolation

Next-to-leading twist Normalization Constants



All results preliminary, statistical errors only, no chiral extrapolation

Next-to-leading twist Normalization Constants

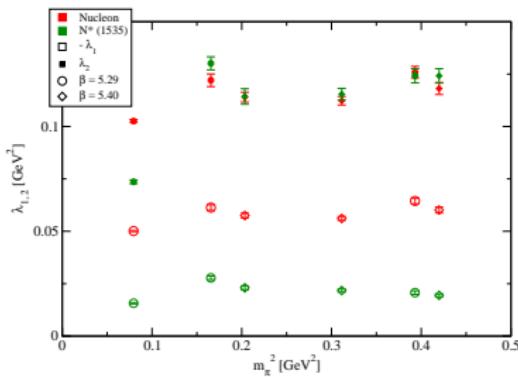


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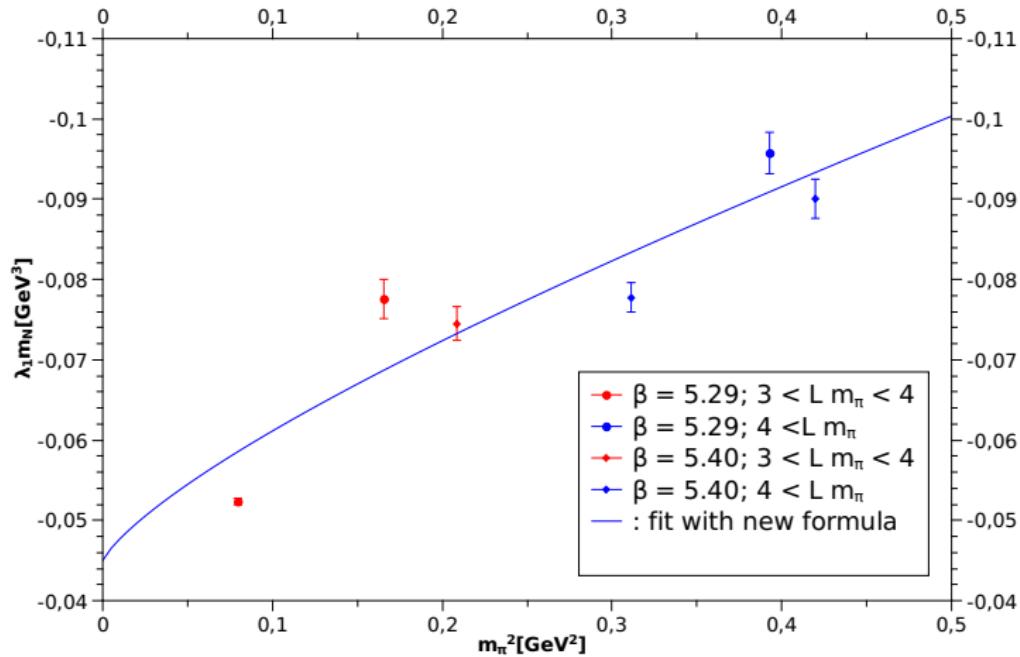
Leading-One-Loop Baryon χ PT for λ_1

$$\begin{aligned} \lambda_1 = & \frac{4}{m_N} \left[\alpha \left(1 + \frac{3}{64\pi^2 f^2} (1 + 3g_A^2) m_\pi^2 \ln \frac{\lambda^2}{m_\pi^2} - \frac{3g_A^2}{32\pi^2 f^2} m_\pi^2 + \right. \right. \\ & + \frac{1}{64\pi^2 f^2} \frac{m_\pi^3}{M} \frac{(12g_A + 18g_A^2) - (3g_A + 6g_A^2) \frac{m_\pi^2}{M^2}}{\sqrt{1 - \frac{m_\pi^2}{4M^2}}} \arccos \frac{m_\pi}{2M} + \\ & \left. \left. + \frac{1}{64\pi^2 f^2} \frac{m_\pi^4}{M^2} (3g_A + 6g_A^2) \ln \frac{m_\pi^2}{M^2} \right) + \epsilon^{(r)}(\lambda) m_\pi^2 \right] \end{aligned}$$

- ▶ (Wein, Hemmert, Schäfer, in preparation)
- ▶ $\alpha, \epsilon^{(r)}$: free parameters
- ▶ $f = \sqrt{2}f_\pi$
- ▶ g_A : axial coupling constant
- ▶ λ : renormalization scale
- ▶ $M = 938\text{MeV}$



χ PT-fit to λ_1 of the Nucleon



(thanks to Philipp Wein)

All results preliminary, statistical errors only, chiral extrapolation!

Wave functions and Distribution amplitudes

- ▶ Expand wave function in multiplicatively renormalizable terms (Braun, Manashov, Rohrwild):

$$\begin{aligned}\varphi(x_i; \mu^2) = & 120x_1x_2x_3 \left\{ 1 + \textcolor{red}{c_{10}}(x_1 - 2x_2 + x_3)L^{\frac{8}{3\beta_0}} \right. \\ & + \textcolor{red}{c_{11}}(x_1 - x_3)L^{\frac{20}{9\beta_0}} + \textcolor{red}{c_{20}} \left[1 + 7(x_2 - 2x_1x_3 - 2x_2^2) \right] L^{\frac{14}{3\beta_0}} \\ & + \textcolor{red}{c_{21}}(1 - 4x_2)(x_1 - x_3)L^{\frac{40}{9\beta_0}} \\ & \left. + \textcolor{red}{c_{22}}[3 - 9x_2 + 8x_2^2 - 12x_1x_3]L^{\frac{32}{9\beta_0}} + \dots \right\}\end{aligned}$$

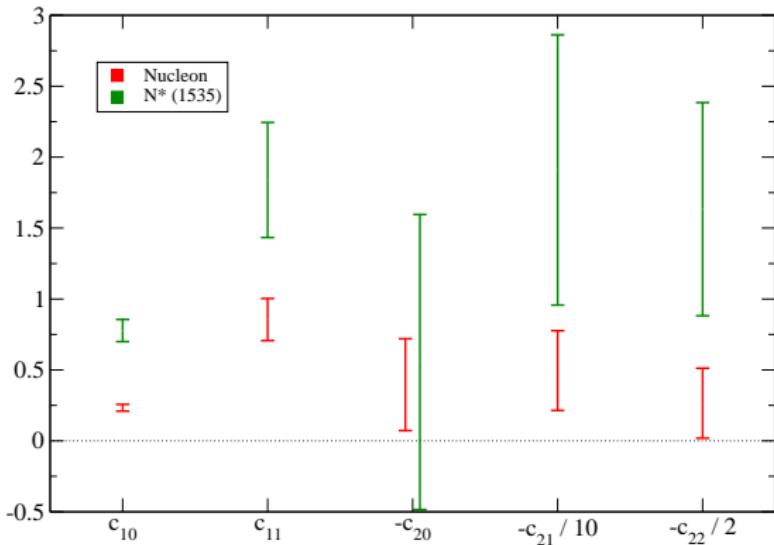
- ▶ where $L \equiv \alpha_s(\mu)/\alpha_s(\mu_0)$
- ▶ $\textcolor{red}{c_{ij}}$: “shape parameters” are obtained from a constrained fit to φ^{lmn} with $l + m + n \leq i$ such that $\varphi^{lmn} = \varphi^{(l+1)mn} + \varphi^{l(m+1)n} + \varphi^{lm(n+1)}$ is fulfilled



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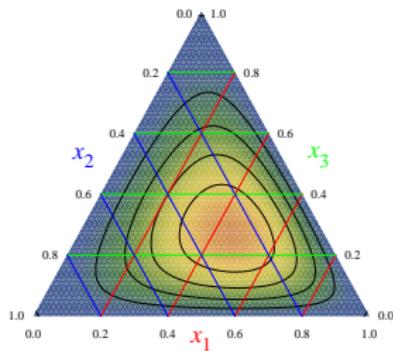
Shape parameters

$32^3 \times 64$
 $a = 0.0753 \text{ fm}$
 $m_\pi = 282(2) \text{ MeV}$
 $m_\pi L = 3.44$
 578×4
configurations



- ▶ c_{1j} are significantly different from zero, both for the nucleon and $N^*(1535)$
- ▶ All results preliminary, statistical errors only, no chiral extrapolation

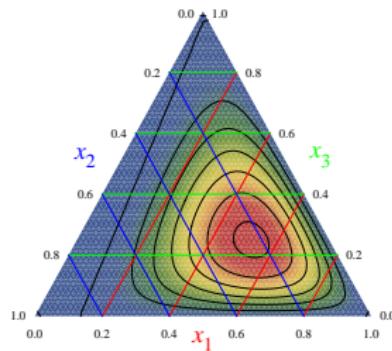
Shape parameters (II)



Nucleon

$$c_{10}^N = 0.233 \pm 0.024 ,$$

$$c_{11}^N = 0.855 \pm 0.145 ,$$



$N^*(1535)$

$$c_{10}^{N^*} = 0.778 \pm 0.078$$

$$c_{11}^{N^*} = 1.84 \pm 0.41$$

- ▶ The distribution amplitude of the $N^*(1535)$ is much more asymmetric than that of the nucleon
- ▶ All results preliminary, statistical errors only, no chiral extrapolation

Conclusions and Outlook

- ▶ $f_N, f_{N^*}, \lambda_1, \lambda_2$ and the first two moments of the nucleon and $N^*(1535)$ distribution amplitudes have been calculated at several pion masses, the lowest being $m_\pi = 282\text{MeV}$
- ▶ Finer and larger volume lattices are still being investigated
- ▶ Chiral Perturbation Theory calculations for more quantities in progress
- ▶ Next big step: extension of this study to all octet and decuplet baryons

