Detecting Axion Dark Matter through the Radio Signal from Omega Centauri

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Astroparticle Physics, SISSA



JW Wang, XJ Bi, RM Yao, PF Yin, arXiv:2101.02585, PRD

JW Wang, XJ Bi, PF Yin, arXiv:2109.00877



Barolo, September 9th, 2021

- Introduction and motivation
- The properties of Omega Centauri
- The properties of compact stars
- Detecting sensitivity on SKA 1 and LOFAR
- Summary

Dark matter in the Universe

 The astrophysical and cosmological observations have provided compelling evidences of the existence of dark matter (DM).







Planck 2018 [1807.06209]



Cold DM (~26%) $\Omega_c h^2 = 0.11933 \pm 0.00091$ Baryons (~5%) $\Omega_b h^2 = 0.02242 \pm 0.00014$ Dark energy (~69%) $\Omega_{\Lambda} = 0.6889 \pm 0.0056$ • Strong CP problem: the QCD Lagrangian can contain a CP violating term

$$\mathcal{L}_{\theta} = \bar{\theta} \frac{g_s^2}{32\pi^2} \varepsilon^{\mu\nu\alpha\beta} F^a_{\mu\nu} F^a_{\alpha\beta}$$

Current limit from neutron EDM:

$$\bar{\theta} < 10^{-10}$$

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- Current limit from neutron EDM: $\bar{\theta} < 10^{-10}$
- Introduce a global U(1) PQ symmetry, promote $\overline{\theta}$ to a dynamical quantity instead of a constant parameter.
- The goldstone boson associated with U(1) PQ is called axion, with is a good DM candidate.

The coupling between axion and electromagnetic sector:

- -

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Axion-photon conversion in the external magnetic field

The coupling between axion and electromagnetic sector:

$$\mathcal{L}_{a\gamma\gamma} = -(1/4)g_{a\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu} = g_{a\gamma}a\boldsymbol{E}\cdot\boldsymbol{B} \qquad \longrightarrow \qquad \begin{array}{c} \mathcal{A} \qquad \boldsymbol{\gamma} \\ \boldsymbol{\gamma}_{\text{virtual}} \\ \boldsymbol{B} \\ \boldsymbol{X} \end{array}$$

Axion-photon conversion in the external magnetic field

• Axion helioscope [1705.02290, 1401.3233 ---]



The coupling between axion and electromagnetic sector:

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- light shining through a wall [1004.1313, 1302.5647 ---]



The coupling between axion and electromagnetic sector:

Axion-photon conversion in the external magnetic field

- Axion helioscope [1705.02290, 1401.3233 ---]
- light shining through a wall [1004.1313, 1302.5647 ---]
- Neutron stars [0711.1264, 1804.03145, 1803.08230, 2008.01877, 2011.05378, 2004.06486 ----]
- Magnetic white dwarf stars [2101.02585]



From Benjamin Safdi's talk

Single NS/MWD results

• The axion-photon mixing equations

$$\begin{bmatrix} -i\frac{d}{dr} + \frac{1}{2k} \begin{pmatrix} m_a^2 - \xi \,\omega_p^2 & -\Delta_B \\ -\Delta_B & 0 \end{pmatrix} \end{bmatrix} \begin{pmatrix} \tilde{A}_y \\ \tilde{a} \end{pmatrix} = 0 \quad \text{with} \quad \begin{cases} \xi = \frac{\sin^2 \tilde{\theta}}{1 - \frac{\omega_p^2}{\omega^2} \cos^2 \tilde{\theta}}, \quad \Delta_B = Bg_{a\gamma} \omega \frac{\xi}{\sin \tilde{\theta}} \end{cases}$$

• Resonant conversion condition $m_a = \omega_p$. When $\theta = \pi/2$

$$P_{ay} = \frac{1}{2v_c^2} g_{ay}^2 B^2(r_c) L, \text{ where } L = \begin{cases} \sqrt{\frac{2\pi k(r_c)r_c}{3m_A^2(r_c)}} & \text{for NS} \\ \sqrt{\frac{2\pi v_c H_{cor}}{m_A(r_c)}} & \text{for MWDs} \end{cases}$$

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• Resonant conversion condition $\overline{m_a = \omega_p}$. When $\theta = \pi/2$

 10^{3}

ADM

 10^{-5}

10⁴

CAST



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 10^{-6}

 10^{-4}

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• The DM profile:
$$\rho_{\rm NFW} = \rho_s \left(\frac{r}{r_s}\right)^{-1} \left(1 + \frac{r}{r_s}\right)^{-2}$$
 with $\begin{array}{l} \rho_s = 7650.59 \ M_{\odot} {\rm pc}^{-34} \\ r_s = 1.63 \ {\rm pc} \end{array}$

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- The King model and position dependent MB distribution.
- N-body simulations of GCs by using CMC code [K. Kremer, etc, 2020]

12531 NSs and 102990 MWDs with $B_0 > 0.1 MG$

• Dipole magnetic field configuration, set	$n_{e0} = 10^{10} \text{ cm}^{-3}$	and	$T_{\rm cor} = 10^6 {\rm K}$	
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- For NSs, the distribution of P_{NS} , α_{NS} , and B_0 are derived by the NS evolution models

$$P(t)P'(t) = \begin{cases} \frac{2}{3} \frac{P_{\rm in}^2}{\tau_{\rm in}} & \text{for active NSs} \\ \frac{2}{3} \frac{P_{\rm in}^2}{\tau_{\rm in}} \sin^2 \alpha(t) & \text{for dead NSs} \end{cases}, \qquad B_0/P^2 > 0.34 \times 10^{12} \text{ G s}^{-2} \\ \frac{d}{dt} \log \sin\alpha(t) = -\frac{2}{3} \frac{\cos^2 \alpha_{\rm in}}{\tau_{\rm in}}, \qquad \tau_{\rm in} = \frac{I}{\pi \mu^2 f_0^2} \approx 8904 \left(\frac{10^{12} \text{ G}}{B_0}\right)^2 \left(\frac{P_{\rm in}}{0.01}\right)^2 \text{ yr} \\ = -B_0 \left[\frac{1}{\tau_{\rm ohm}} + \left(\frac{B_0}{B_{\rm in}}\right)^2 \frac{1}{\tau_{\rm ambip}}\right], \qquad \tau_{\rm ohm} \sim \frac{1.8 \times 10^9}{Q_{\rm imp}} \text{ yr}, \qquad \tau_{\rm ambip} \sim 3 \times 10^9 \left(\frac{t_{\rm E}}{10^6 \text{ yr}}\right)^{-1/3} \left(\frac{B_{\rm in}}{10^{12} \text{ G}}\right)^{-2} \text{ yr} \end{cases}$$

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The sensitivity of SKA1 and LOFAR

Technical parameters of the SKA1 and LOFAR [13, 16]						
Channel		Range [GHz]	Resolution [kHz]	FoV [arcmin]		
SKA1 LOW		$0.05 \sim 0.35$	1.0	327		
SKA1 MID	B1	$0.35 {\sim} 1.05$	3.9	109		
	B2	$0.95 {\sim} 1.76$	3.9	60		
	$B3^*$	$1.65 {\sim} 3.05$	9.7	42		
	$B4^*$	$2.80 \sim 5.18$	9.7	42		
	B5a	$4.6 \sim 8.5$	9.7	12.5		
	B5b	$8.3 \sim 15.3$	9.7	6.7		
LOFAR LBA		$0.03 {\sim} 0.08$	195	470.9		
LOFAR HBA		$0.11 \sim 0.24$	195	94.8		

• The minimal detectable flux density:

$$S_{\min} = \frac{\text{SEFD}}{\eta_s \sqrt{n_{\text{pol}} \mathcal{B} t_{\text{obs}}}} , \qquad \text{SEFD} = \frac{2k_B}{A_{\text{eff}}/T_{\text{sys}}}$$

• The total flux density from Omega Centauri:

$$S_{a\gamma}^{\text{total}} = \left(\sum_{N_{\text{NS}}} S_{a\gamma}^{\text{NS}} + \sum_{N_{\text{WD}}} S_{a\gamma}^{\text{WD}}\right) \times \Theta\left(\frac{\text{FoV}}{2} - P_{\text{CS}}\right)$$

Result of pure NSs



• The largest detectable axion mass:

$$m_a^{\rm max} \approx 69.2 \times \left| 3\cos^2 \theta - 1 \right|^{1/2} \sqrt{\frac{B_0}{10^{14} \text{ G}} \frac{1 \text{ s}}{P}} \ \mu \text{eV}$$

Pure MWDs results



• The statistical fluctuation is quiet small (law of large numbers)

Combined results: NS+WMDs



• For NS model 1, the contribution of NS is dominant

• For NS model 2, the contribution of MWDs is relatively larger in the region of $m_a \lesssim 2 \,\mu {
m eV}$

Summary

- Axion is one of the most compelling DM candidates, and can also solve the strong CP problem in an elegant way;
- The axion-photon conversion provides a possible way to detect the axion dark matter, especially around objects with strong magnetic field;
- Compared with single CS, the Omega Centauri is a much more effective target. The constrains on $g_{a\gamma}$ can up to $10^{-14} \sim 10^{-15} \text{ GeV}^{-1}$

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