Electron Polarimetry at EIC

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- Polarimetry at EIC
- Electron Polarimetry
 - ESR Compton
 - RCS Compton
 - Mott Polarimeters



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EIC will provide unique challenges for electron polarimetry

- → 10 ns between electron/hadron bunches at high luminosity configuration (~40 ns at higher CM configuration)
- \rightarrow Intense beams (0.26 to 2.5 A)
 - → Large synchrotron radiation for electron beams result in large effects at detectors

Requirements:

- → Bunch-by-bunch measurement of polarization
- → Simultaneous measurement of both P_L and P_T
- → Measurement fast enough to achieve 1% statistics for each bunch
- → Systematics *dP/P* = 1% or better

Table 1.1: Maximum luminosity parameters.

hadron	electron		
104.9			
275	10		
1160			
6.9	17.2		
1.0	2.5		
11.3	20.0		
1.0	1.3		
80	45		
7.2	5.6		
0.228/0.210	0.08/0.06		
0.119	0.211		
0.119	0.152		
0.012	0.072		
0.012	0.1		
2.9/2.0	-		
-	9.0		
6	0.7		
0.9	4		
Luminosity $[10^{34} \mathrm{cm}^{-2} \mathrm{s}^{-1}]$ 1.0			
	hadron 104 275 116 6.9 1.0 11.3 1.0 80 7.2 0.228/0.210 0.119 0.119 0.119 0.119 0.119 0.119 0.012 2.9/2.0 - - 6 0.9 1.0		



EIC Electron Polarimeter Map





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Polarization Time Dependence - electrons

- Electrons injected into the storage ring at full polarization (85%)
- Sokolov-Ternov effect (self-polarization) will re-orient spins to be anti-parallel to main dipole field → electrons will have different lifetime depending on polarization
- Bunches must be replaced relatively often to keep average polarization high
- Bunch-by-bunch polarization measurement required



Compton polarimeter will be upstream of upstream of detector IP

At Compton interaction point, electrons have both longitudinal and transverse (horizontal) components

- → Longitudinal polarization measured via asymmetry as a function of backscattered photon/scattered electron energy
- \rightarrow Transverse polarization from left-right asymmetry

Beam energy	PL	P _T
5 GeV	96.5%	26.1%
10 GeV	86.4%	50.4%
18 GeV	58.1%	81.4%

Polarization Components at Compton

Beam polarization will be fully longitudinal at detector IP, but accurate measurement of absolute polarization will require *simultaneous* measurement of P_L and P_T at Compton polarimeter

EIC Compton will provide first high precision measurement of P_1 and P_7 at the same time



Polarization Measurement via Compton Polarimetry

Compton longitudinal and transverse analyzing powers

$$A_{\text{long}} = \frac{2\pi r_o^2 a}{(d\sigma/d\rho)} (1 - \rho(1+a)) \left[1 - \frac{1}{(1 - \rho(1-a))^2} \right]$$

$$A_{\rm T} = \frac{2\pi r_o^2 a}{(d\sigma/d\rho)} \cos\phi \left[\rho(1-a)\frac{\sqrt{4a\rho(1-\rho)}}{(1-\rho(1-a))}\right]$$



Compton polarimetry – lessons from previous devices

- Longitudinal polarimetry
 - Electron detector needs sufficient segmentation to allow self-calibration "on-the-fly"
 - Photon detector integrating technique provides most robust results – perhaps not practical at EIC? → lower the threshold as much as possible
- Transverse polarimetry
 - Remove η -y calibration issue use highly segmented detectors at all times
 - Calorimeter resolution \rightarrow integrate over all energy?
 - Beam size/trajectory important build in sufficient beam diagnostics
- Common to both
 - Birefringence of vacuum windows can impact laser polarization → use back-reflected light (optical reversibility theorems)





Compton Placement and Integration



Figure courtesy Ciprian Gal (Miss. State U.)

- \rightarrow Laser IP in field-free area space to insert laser in beamline
- \rightarrow Photon detector 29 m from laser/beam IP
- → Quad after dipole (Q5EF) horizontally defocusing facilitates use of electron detector
- \rightarrow Synchrotron from D3EF4 may impact electron detector



The needed time t_D to achieve an accuracy $\Delta P_e/P_e$ is then

Compton Laser System Requirem

$$t_D^{-1} = \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < \frac{A_l^2}{1 - P_e^2 P_{\gamma}^2 A_l^2} > \simeq \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 P_{\gamma}^2 P_{\gamma}^2 \sigma_t < A_l^2 > 2 \mathcal{L} \left(\frac{\Delta P_e}{P_e}\right)^2 P_e^2 P_{\gamma}^2 P_$$

8	Configuration	Beam energy [GeV]	Unpol Xsec[barn]	Tot Unpol Xsec[barn]	Apeak [not used]	<a^2></a^2>	L	1/t(1%)	t[s]	t[min]	/
9	laser:532nm, photon long	18	0.432	0.432	0.310	2.07E-02	1.81E+05	1.17E-01	9		0.14
10	laser:532nm, photon trans	18	0.432	0.432	0.210	3.62E-03	1.81E+05	2.05E-02	49		0.81
11	laser:532nm, electron	18	0.301	0.432	0.320	4.57E-02	1.81E+05	1.80E-01	6		0.09
12											
13	laser:532nm, photon long	10	0.503	0.503	0.270	1.54E-02	1.55E+05	8.69E-02	12		0.19
14	laser:532nm, photon trans	10	0.503	0.503	0.170	2.15E-03	1.55E+05	1.21E-02	83		1.38
15	laser:532nm, electron	10	0.340	0.503	0.270	3.05E-02	1.55E+05	1.17E-01	9		0.14
16											
17	laser:532nm, photon long	5	0.569	0.569	0.160	5.82E-03	1.37E+05	3.29E-02	30		0.51
18	laser:532nm, photon trans	5	0.569	0.569	0.110	1.63E-03	1.37E+05	9.19E-03	109		1.81
19	laser:532nm, electron	5	0.323	0.569	0.160	1.14E-02	1.37E+05	3.65E-02	27		0.46

Ciprian Gal

Laser power constraint: sufficient power to provide ~ 1 backscattered photon/bunch-laser crossing → Want to make "single photon" measurements – not integrating

532 nm laser with ~5 W average power at same frequency as EIC electron bunches sufficient

StReBulting measurement times (for differential measurement, dP/P=1%) as noted above – easily meets beam lifetime constraints



Compton Polarimeter Laser System

Average of 1 backscattered photon/bunch crossing will allow Compton measurements on the ~1 minute time scale → can be achieved with pulsed laser system that provides about 5 W average power at 532 nm



JLab injector laser system

Polarization in vacuum set using "back-reflection" technique

→ Requires remotely insertable mirror (in vacuum) Proposed laser system based on similar system used in JLab injector and LERF

- Gain-switched diode seed laser variable frequency, few to 10 ps pulses @ 1064 nm
 - → Variable frequency allows optimal use at different bunch frequencies (100 MHz vs 25 MHz)
- 2. Fiber amplifier \rightarrow average power 10-20 W
- 3. Optional: Frequency doubling system (LBO or PPLN)
- 4. Insertable in-vacuum mirror for laser polarization setup



Prototype system under development (C. Gal, Mississippi State U.) errerson Lab

Electron Detector Size and Segmentation

- Electron detector (horizontal) size determined by spectrum at 18 GeV (spectrum has largest horizontal spread)
 - Need to capture zero-crossing to endpoint → detector should cover at least 60 mm
- Segmentation dictated by spectrum at 5 GeV (smallest spread)
 - Scales ~ energy \rightarrow 17 mm
 - Need at least 30 bins, so a strip pitch of about 550 μm would be sufficient
- At 18 GeV, zero-crossing about 3 cm from beam
 - 5 GeV \rightarrow 8-10 mm this might be challenging



Asymmetry at electron detector @18 GeV



Transverse Polarization Measureme

- At Compton location significant transverse beam
- → Unfortunately, this transverse polarization is in the horizontal direction
- \rightarrow Same coordinate as momentum-analyzing dipole

In the absence of the dipole, the transversely polarized electrons would result in a left-right asymmetry

- The "scattered electron cone" is much smaller than the photons
- → Left-right asymmetry is spread over much smaller distance (µm vs mm)

The large dispersion induced by the dipole makes measurement of the left-right asymmetry impossible

Electron detector can only be used for measurements of P_L Stony Brook University



100% transversely polarized beam

18GeV eDet(bQ9) polXsec



Electron detector considerations

Not clear if electron detector can live in vacuum directly – may need to be housed in a structure similar to Roman Pot

- → Preliminary wakefield calculations (alternate configuration) suggest power deposited manageable
- ightarrow This needs to be updated for latest EIC layout



Electron detector out of direct synchrotron fan, but single-bounce can deposit power on detector

- ightarrow Studies by Mike Sullivan (for different configuration) suggested large power deposition
- → Updated studies with GEANT4 for latest layout suggests that synchrotron backgrounds may not be a problem work in progress





Polarization Measurement with Photon Detector

Photon detector needs 2 components to measure both longitudinal and transverse polarization

- Calorimeter \rightarrow asymmetry vs. photon energy (P_L) Ο
- Position sensitive detector \rightarrow left-right asymmetry (P_T) Ο



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0.2 0.3	0.4 0.5 0 am E = 18 pc	.6 0.7 0 photon E / ma	.8 0.9 ιx photon Ε ρ	1
	٨			
	$- \Lambda$			
		-		

Beam energy	PL	P _T
5 GeV	96.5%	26.1%
10 GeV	86.4%	50.4%
18 GeV	58.1%	81.4%

Transverse size of detectors determined by backscattered photon cone at low energy

- \rightarrow +/- 2 cm adequate at 5 GeV
- → Longitudinal measurement requires good energy resolution from ~0 (as low as possible) to 3 GeV
- \rightarrow Fast time response also needed (10 ns bunch spacing)
- \rightarrow PbWO4 a possible candidate (slow component may be an issue)

Position sensitive detector segmentation determined by highest energy \rightarrow 18 GeV

 \rightarrow More investigation needed, but segmentation on the order of 100-400 μ m should work 14



Position Sensitive Detectors

- Requirements for position sensitive detectors
 - Radiation hard
 - Fast response (needed for bunch-by-bunch measurements)
 - High granularity (down to 25 µm pitch)

Size determined by 5 GeV hit distributions, segmentation by 18 GeV distributions

Diamond strip detectors have been used successfully at JLab in Compton polarimeters

- → No performance degradation after 10 Mrad dose during Q-Weak experiment @ JLab
- → Intrinsic time response is fast, but small signals require significant amplification custom electronics/ASIC will be required



500 μ m pCVD diamond w/TOTEM electronics



Electron Polarimetry Systematics	Beam energy	PL	P _T
	5 GeV	96.5%	26.1%
State of the art for Compton polarimetry:	10 GeV	86.4%	50.4%
Longitudinal	18 GeV	58.1%	81.4%

Longitudinal:

SLD @ SLAC: dP/P=0.5% \rightarrow Electron detector in multi-photon mode Q-Weak in Hall C @ JLab: dP/P=0.59% \rightarrow Electron detector, counting mode CREX in Hall A @ JLab: dP/P=0.44% \rightarrow Photon detector, integrating mode

Transverse:

TPOL @ HERA: dP/P=1.87% → Photon detector in counting mode

Total polarization extraction will rely on two quasi-independent measurements While 0.5% for P_L is plausible, P_T is less certain \rightarrow 1%?

At 18 GeV this results in dP/P=0.86% at 18 GeV



Rapid Cycling Synchrotron (RCS) Compton Polarimeter

RCS properties

- RCS accelerates electron bunches from 0.4 GeV to full beam energy (5-18 GeV)
- Bunch frequency \rightarrow 2 Hz
- Bunch charge → up to 28 nA
- Ramping time = 100 ms



Polarimetry challenges

- Analyzing power often depends on beam energy
- Low average current
- Bunch lifetime is short

Compton polarimeter can also be used for measurement of polarization in RCS

- → Measurements will be averaged over several bunches can tag accelerating bunches to get information on bunches at fixed energy
- → Requires measurement in multiphoton mode (many backscattered photons/electron bunch)



Highest precision transverse Compton polarimeter operated in single photon mode (HERA)

 \rightarrow RCS requires position sensitive measurement in **multi-photon** mode

Need highly segmented detector sensitive to signal size (not just counts above threshold)

 \rightarrow LEP polarimeter operated in this fashion, although with relatively low precision





Differential measurement of asymmetry vs. position at detector allows us to incorporate offsets in the fit

Example using Toy MC for integrating mode asymmetry vs. y assuming 0.1 mm segmentation (240 bunches)
→ Sufficient position resolution would allow determination of arbitrary offset in spectrum

→ Requires detector operating in integrating mode (~10,000 photons/bunch) with signal proportional to number of photons in each channel





Rate and measurement time estimates

$$t^{-1} = \mathcal{L}\sigma \left(\frac{\Delta P}{P}\right)^2 P^2 A_{method}^2$$

Average analyzing power: $A^2_{method} = \langle A \rangle^2 \rightarrow$ Average value of asymmetry over acceptance

Energy-weighted: $A_{method}^2 = \left(\frac{\langle EA \rangle}{\langle E \rangle}\right)^2 \rightarrow$ Energy deposited in detector for each helicity state

Differential: $A^2_{method} = \langle A^2 \rangle \rightarrow$ Measurement of asymmetry bin-by-bin vs. energy, etc.

Assuming 80% polarization, $\langle P_{laser} \rangle = 6$ mW, 300 μ m beam spot size..., \rightarrow time for 1% measurement

E _{beam}	A _{avg}	T _{avg}	A _{energy}	T _{energy}	A _{diff}	T _{diff}
5	4.51%	243 s	5.78%	148 s	5.48%	164 s
10	7.79%	92 s	10.15%	54 s	9.56%	61 s
18	11.29%	51 s	14.91%	29 s	13.96%	33 s



Summary

- Electron polarimetry at EIC has challenging requirements
 - Bunch-by-bunch polarization measurement with short times between bunches (as low as 10 ns)
 - -High precision: dP/P=1% (or better)
 - Simultaneous measurement of longitudinal and transverse components
- EIC Compton polarimeter in storage ring must meet all these requirements
 - Simultaneous detection of the backscattered photons and scattered electrons will allow high precision for both longitudinal and transverse polarization
 - -Fast detectors required due to bunch structure
- RCS Compton polarimeter needed to provide information on electron polarization during acceleration
 - -Will operate in multi-photon mode (several thousand photons/bunch crossing)
 - -Less stringent requirements for absolute precision



EXTRA



ESR Compton Detector Technology



JLab Hall C diamond detector

Several choices feasible for position sensitive detectors

- \rightarrow Diamond strip detectors are baseline choice
 - Radiation hard
 - Fast time response
 - Compatible with segmentation requirements
 - ASIC under development for LHC diamond detectors compatible with EIC timing requirements

Tungsten-powder calorimeter



Photon calorimeter more challenging

- → Timing requirements suggest lower resolution calorimeter must be used
- → OK for transverse measurement, but reduces precision on longitudinal



Position Sensitive Detectors

- Requirements for position sensitive detectors
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 - Fast response (needed for bunch-by-bunch measurements)
 - High granularity (down to 25 µm pitch)

Size determined by 5 GeV hit distributions, segmentation by 18 GeV distributions

Diamond strip detectors have been used successfully at JLab in Compton polarimeters

- → No performance degradation after 10 Mrad dose during Q-Weak experiment @ JLab
- → Intrinsic time response is fast, but small signals require significant amplification custom electronics/ASIC will be required



500 μ m pCVD diamond w/TOTEM electronics



Photons will not clear beamline magnet apertures in some cases
→ Quad inner aperture: R= 6 cm
→ Quad outer radius = 25 cm

Depending on final layout, will need to modify one or more quads to allow clear aperture for backscattered photons

If backscattered photons traverse ironfree region – coils can likely/hopefully be modified to accommodate → For 1st option, one quad may require a hole in the iron – but this should not have large impact on quad performance



Quad cross-section



Relatively straightforward to prepare/determine laser polarization before entering beamline

- → Stress on entrance window can introduce significant birefringence
- →Nearly impossible to measure directly without significant instrumentation in vacuum

Measurements at JLab suggest these effects can't be ignored

State 1: DOCP in exit line



JLab laser polarization measurements through 2 vacuum windows → Tightening bolts on flanges, vacuum stress has significant impact



Laser Polarization – Optical reversibility theorems

Propagation of light through the vacuum window to the IP can be described by matrix, $M_E \rightarrow Light$ propagating in opposite direction described by transpose matrix, $(M_E)^T \rightarrow If$ input polarization (ϵ_1) linear, polarization at cavity (ϵ_2) circular only if polarization of reflected light (ϵ_4) linear and orthogonal to input*



 $\varepsilon_4 = (M_F)^T M_F \varepsilon_1$

Laser polarization at a mirror (inside vacuum) can be set/determined by monitoring the back-reflected light in a single photodiode

→ Used this technique at JLab to constrain laser polarization to ~0.1%



*J. Opt. Soc. Am. A/Vol. 10, No. 10/October 1993



EIC will make use of two Mott polarimeters to measure the electron polarization from the source

- 1. Low voltage Mott polarimeter
- → Measure polarization at 20 keV immediately after photocathode
- 2. High voltage Mott polarimeter
- → Measure at 300 keV, in the beamline, before electron bunching
- → Requires spin rotator to change electron from longitudinal to transverse spin





Detector segmentation driven by requirement to be able to extract polarization (fit asymmetry) without any corrections due to detector resolution (see SLD Compton)

→ Studies with toy Monte Carlo suggest that about 30 bins (strips) between asymmetry zero crossing and endpoint results in corrections <0.1%





Luminosity

Luminosity for CW laser colliding with electron beam at non-zero crossing angle:

$$\mathcal{L} = \frac{(1 + \cos \alpha_c)}{\sqrt{2\pi}} \frac{I_e}{e} \frac{P_L \lambda}{hc^2} \frac{1}{\sqrt{\sigma_e^2 + \sigma_\gamma^2}} \frac{1}{\sin \alpha_c}$$

Pulsed laser:

$$\mathcal{L} = f_{coll} N_{\gamma} N_e \frac{\cos\left(\alpha_c/2\right)}{2\pi} \frac{1}{\sqrt{\sigma_{x,\gamma}^2 + \sigma_{x,e}^2}} \frac{1}{\sqrt{(\sigma_{y,\gamma}^2 + \sigma_{y,e}^2)\cos^2\left(\alpha_c/2\right) + (\sigma_{z,\gamma}^2 + \sigma_{z,e}^2)\sin^2\left(\alpha_c/2\right)}}$$

 $N_{\gamma(e)}$ = number of photons (electrons) per bunch

Assumes beam sizes constant over region of overlap (ignores "hourglass effect")

Beam size at interaction point with laser dictates luminosity (for given beam current and laser/electron beam crossing angle)

