

Perspectives in Fundamental Physics

Discussing the future of Particle Physics

Joint Seminar INFN, Dipartimento di Fisica

and Scuola Normale Superiore Pisa

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The theory perspective

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- **having important questions to pursue**

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- **creating opportunities to answer them**

The vision for particle physics builds on

- **having important questions to pursue**
- **creating opportunities to answer them**
- **being able to constantly add to our knowledge, while seeking those answers**

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- **Data driven:**
 - DM
 - Neutrino masses
 - Matter vs antimatter asymmetry
 - Dark energy
 - ...

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- **Data driven:**

- DM
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- **Theory driven:**

- The hierarchy problem and naturalness
- The flavour problem (origin of fermion families, mass/mixing pattern)
- Origin of inflation
- Quantum gravity
- ...

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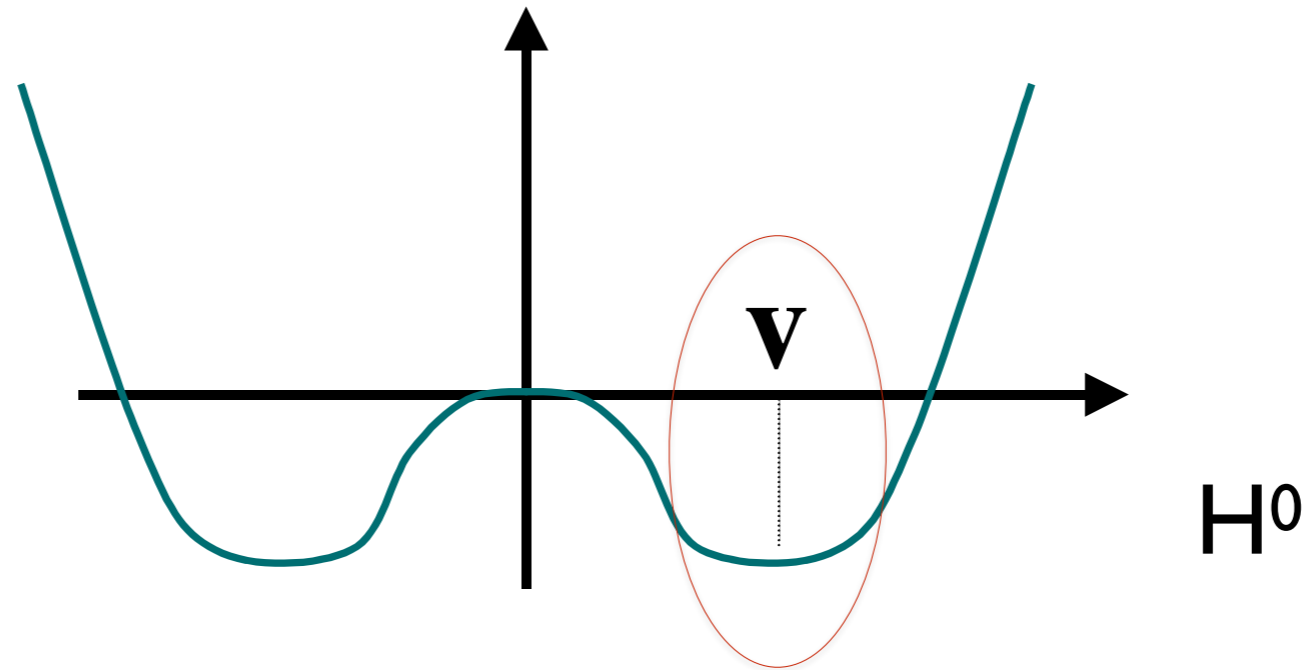
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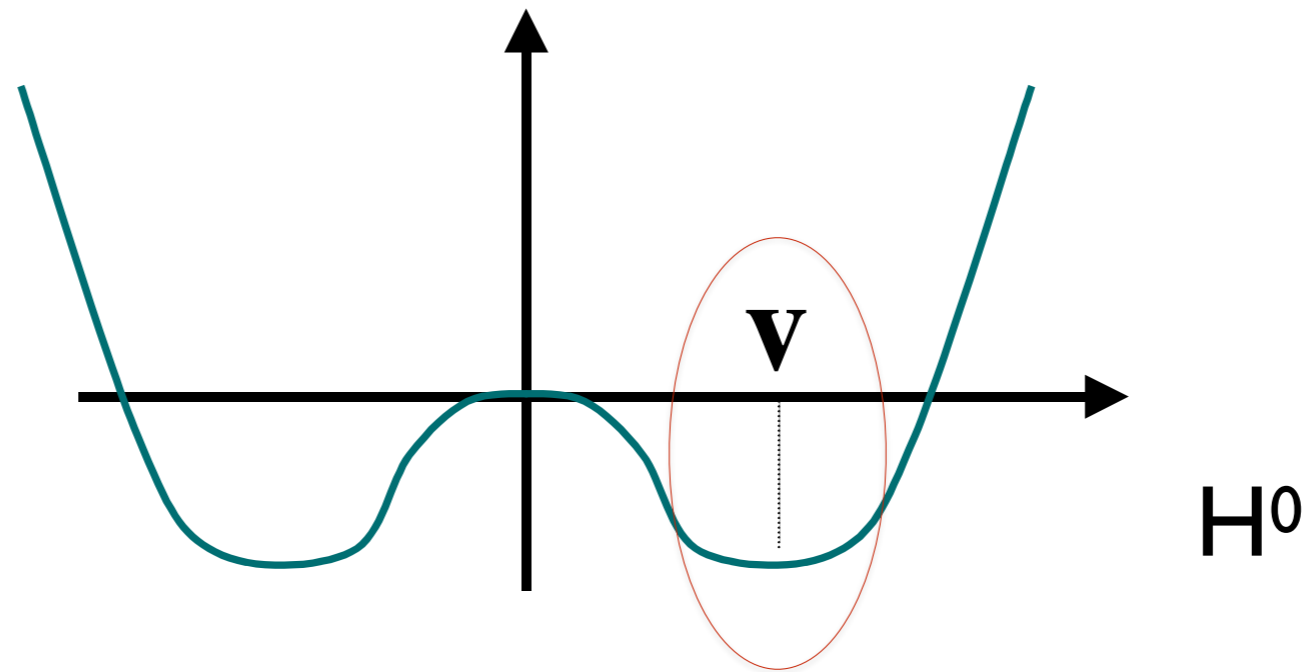
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One question, however, has emerged in stronger and stronger terms from the LHC, and appears to single out a unique well defined direction....



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Who ordered that ?

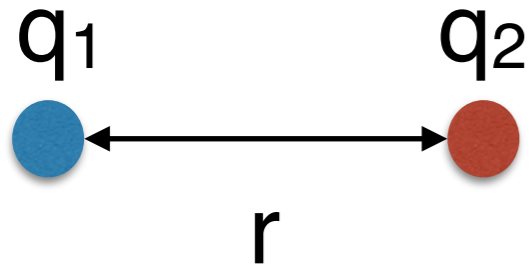


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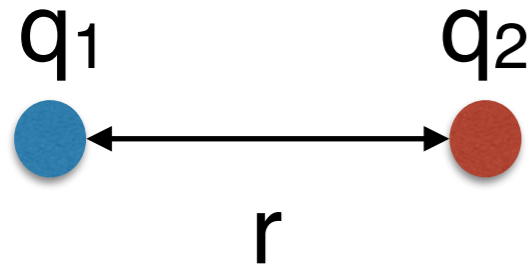
We must learn to appreciate the depth and the value of this question, which is set to define the future of collider physics

Electromagnetic vs Higgs dynamics



$$V(r) = + \frac{q_1 \times q_2}{r^1}$$

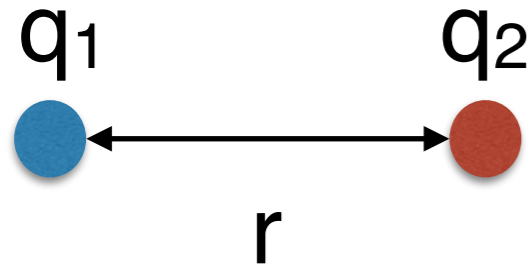
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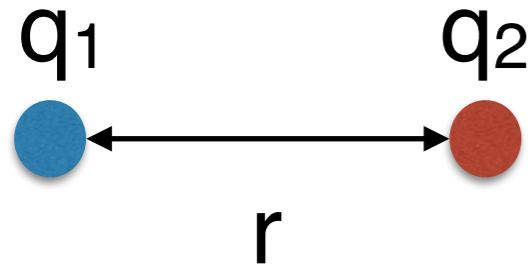
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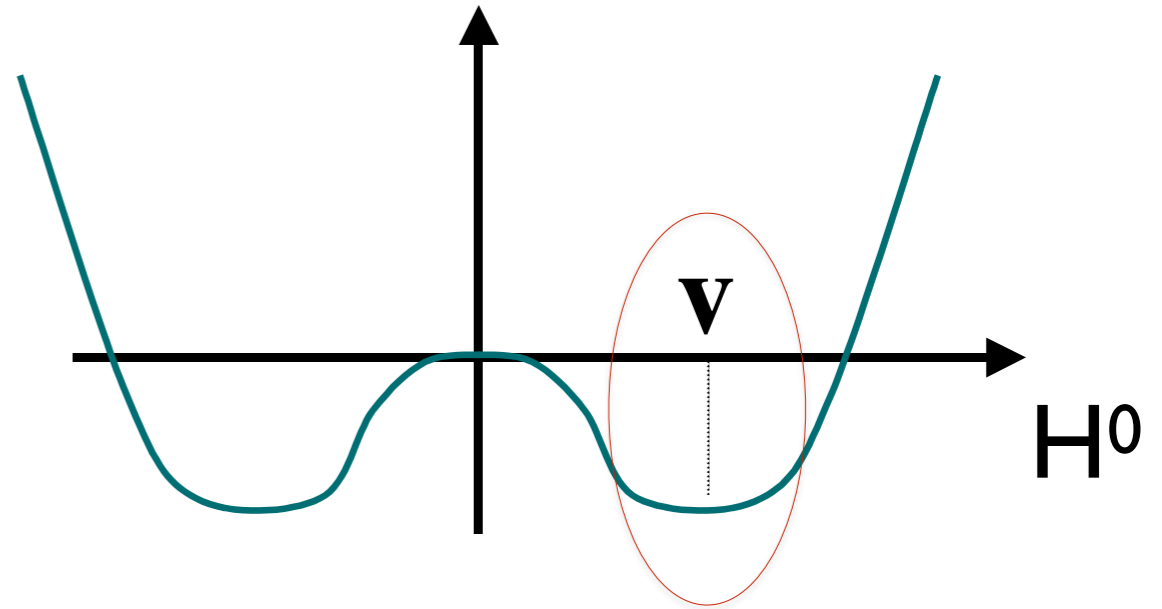


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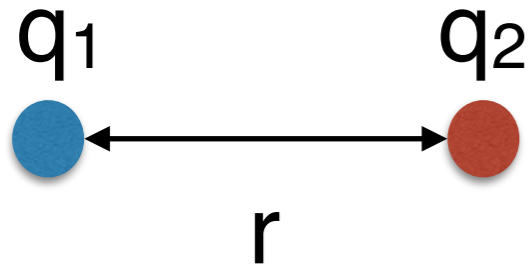
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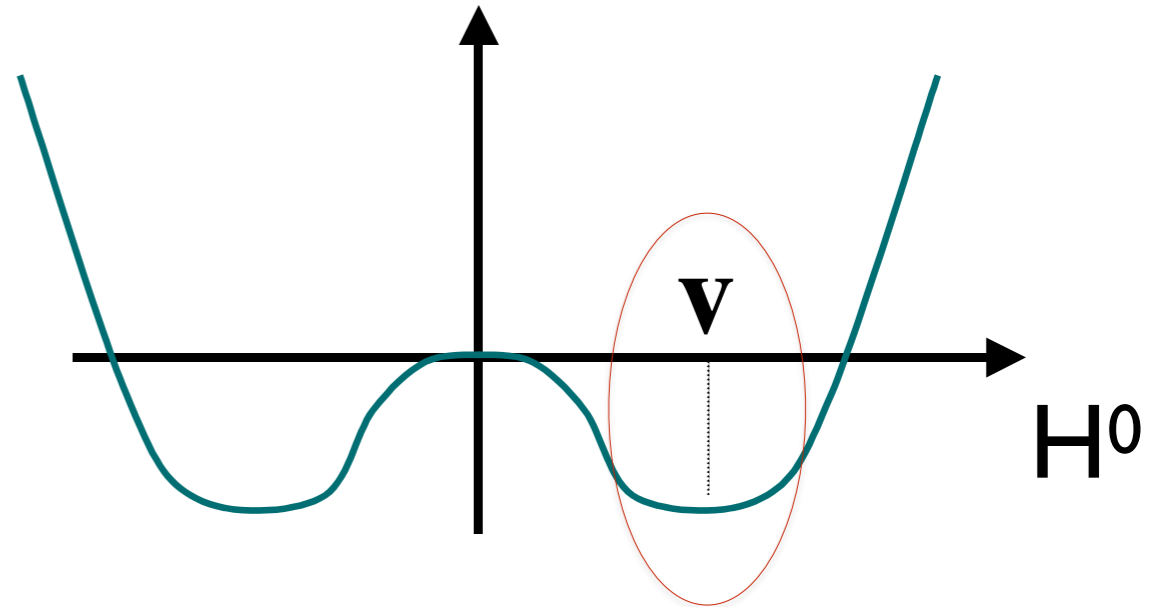


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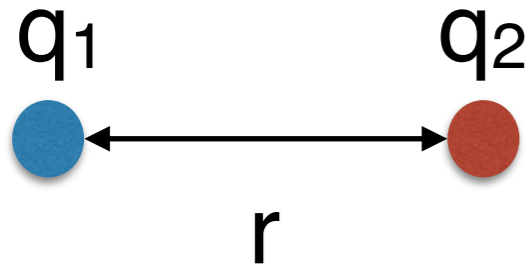


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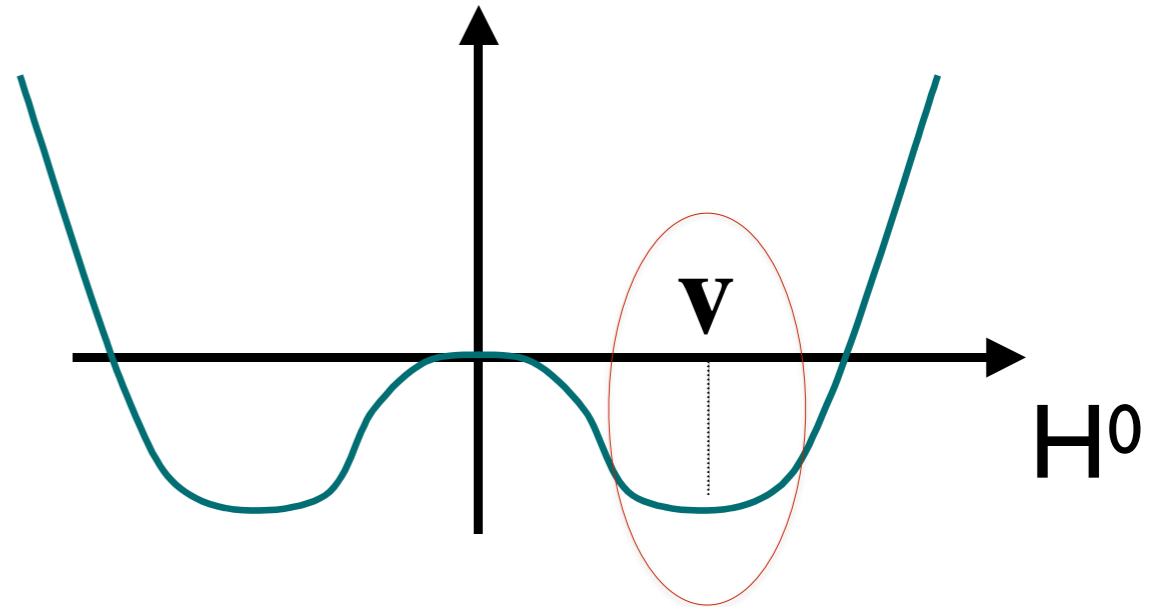


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any function of $|H|^2$ would be
ok wrt known symmetries

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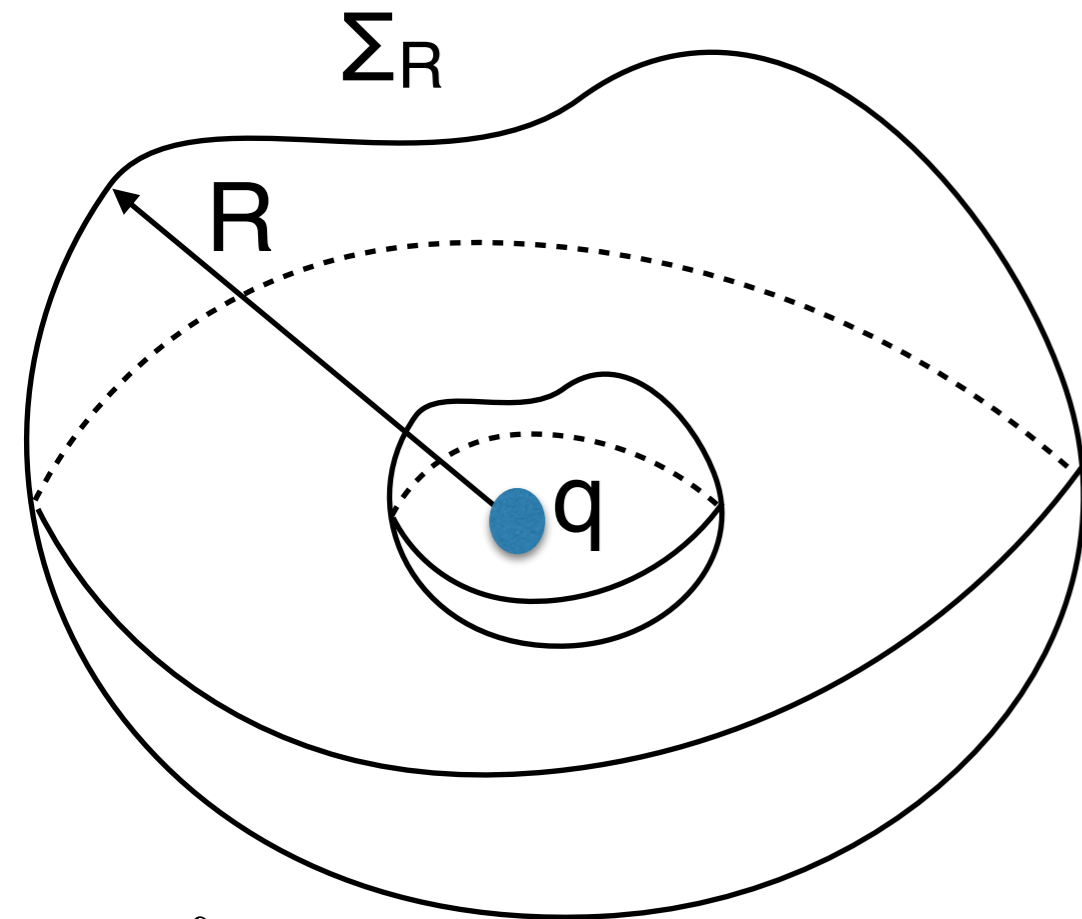
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- For superconductivity, this came later, with the identification of e^-e^- Cooper pairs as the underlying order parameter, and BCS theory. In particle physics, we still don't know whether the Higgs is built out of some sort of Cooper pairs (composite Higgs) or whether it is elementary, and in either case we have no clue as to what is the dynamics that generates the Higgs potential. With Cooper pairs it turned out to be just EM and phonon interactions. With the Higgs, none of the SM interactions can do this, and **we must look beyond.**

examples of possible scenarios

- **BCS-like**: the Higgs is a composite object
- **Supersymmetry**: the Higgs is a fundamental field and
 - $\lambda^2 \sim g^2 + g'^2$, it is not arbitrary (MSSM, w/out susy breaking, has one parameter less than SM!)
 - potential is fixed by susy & gauge symmetry
 - EW symmetry breaking (and thus m_H and λ) determined by the parameters of SUSY breaking
- ...

Decoupling of high-frequency modes

E&M



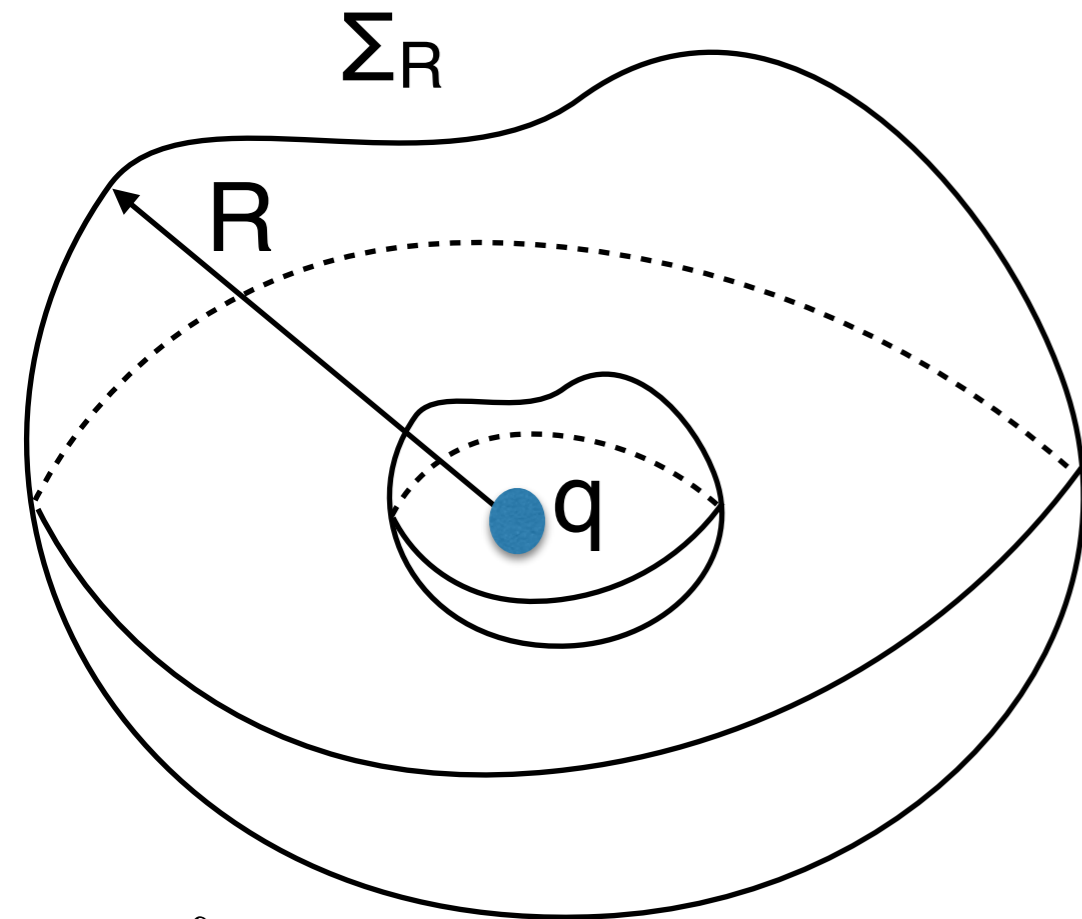
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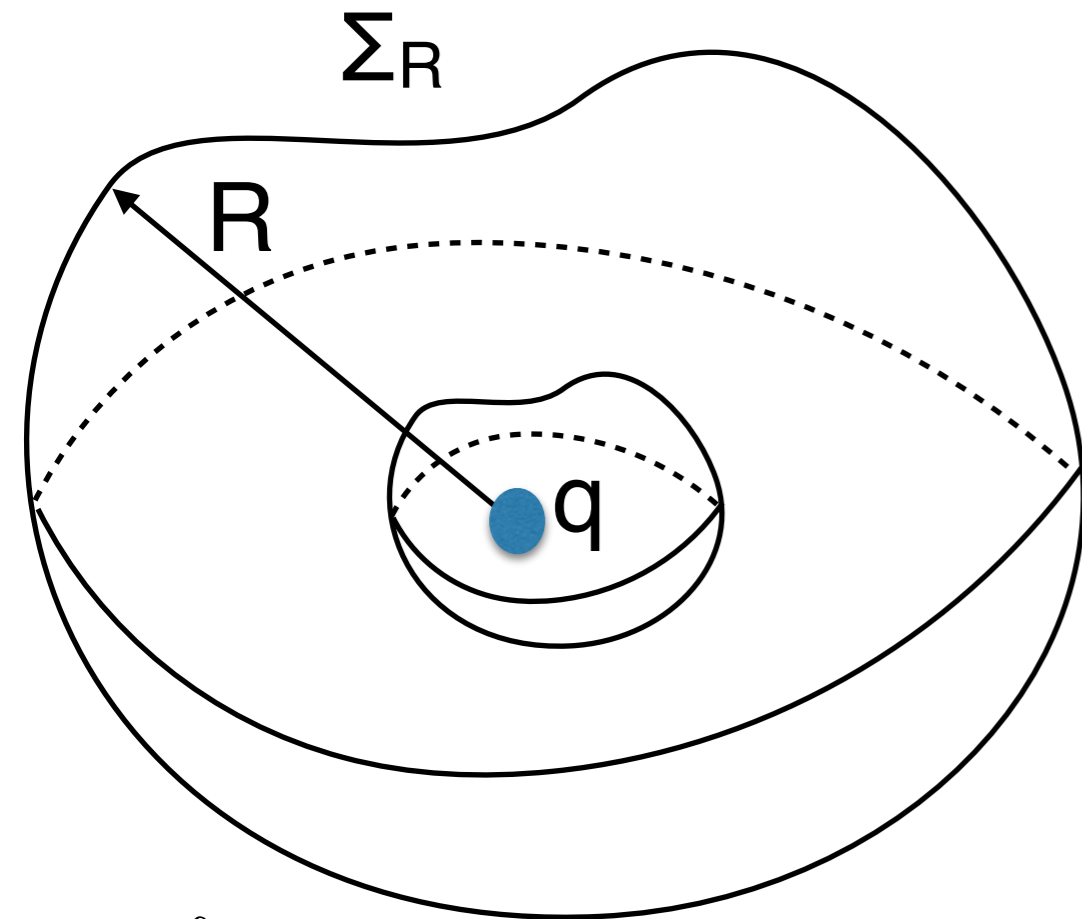


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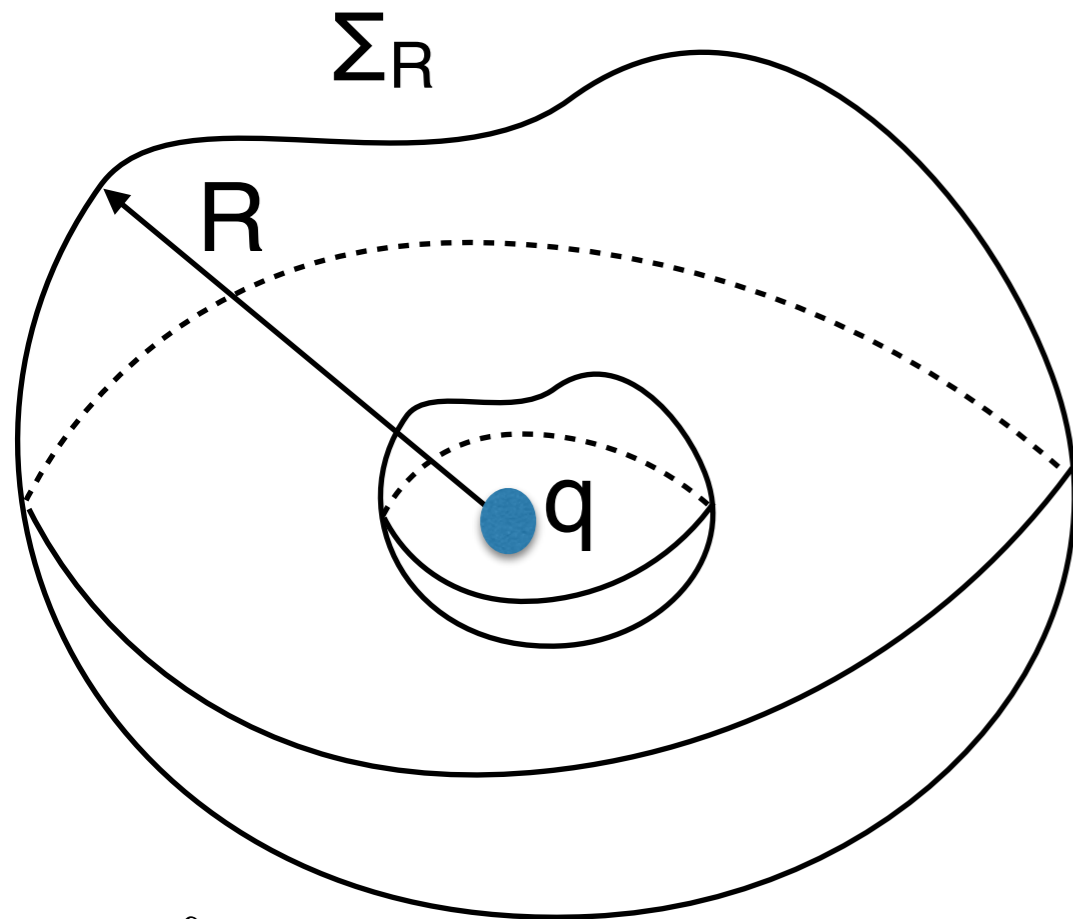
$$\mu^2_{\text{ren}} = \mu^2 + \frac{g^2}{16\pi^2} \ln \frac{\Lambda}{\mu} + \frac{y_t^2}{16\pi^2} \ln \frac{\Lambda}{\mu}$$

The diagram shows a horizontal dashed line representing a propagator. On the left, a solid black dot is connected to the line by a vertical dashed line, with μ^2_{ren} written below. This is followed by an equals sign, then a dashed line with μ^2 below it. A plus sign follows, then a dashed line entering a dashed circle labeled W, H with g^2 below it. Another plus sign follows, then a dashed line entering a solid circle labeled t with $-y_t^2$ below it. The line continues to the right.

$$\Delta\mu^2 \sim (c_B m_B^2 - c_F m_F^2) \times (\Lambda / v)^2$$

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$$\lambda_{\text{ren}} = \lambda - y_t^4 + \lambda^4$$

$$\Rightarrow \frac{d\lambda}{d \log \mu} \propto \lambda^4 - y_t^4 \propto a m_H^4 - b m_t^4$$

high-energy modes can change size and sign of both μ^2 and λ , dramatically altering the stability and dynamics \Rightarrow **hierarchy problem**

bottom line

- The Higgs dynamics is sensitive to all that happens at any scale larger than the Higgs mass !!! A very **unnatural fine tuning** is required to protect the Higgs dynamics from the dynamics at high energy
- This issue goes under the name of **hierarchy problem**
- Solutions to the hierarchy problem require the introduction of new symmetries (typically leading to the existence of new particles), which decouple the high-energy modes and allow the Higgs and its dynamics to be defined at the “natural” scale defined by the measured parameters v and m_H

⇒ **naturalness**

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 - Lack of experimental evidence so far for a straightforward answer to naturalness, forces us to review our biases, and to take a closer look even at the most basic assumptions about Higgs properties
- ➡ *the Higgs discovery does not close the book, it opens a whole new chapter of exploration, based on precise measurements of its properties, which can only rely on a future generation of colliders*

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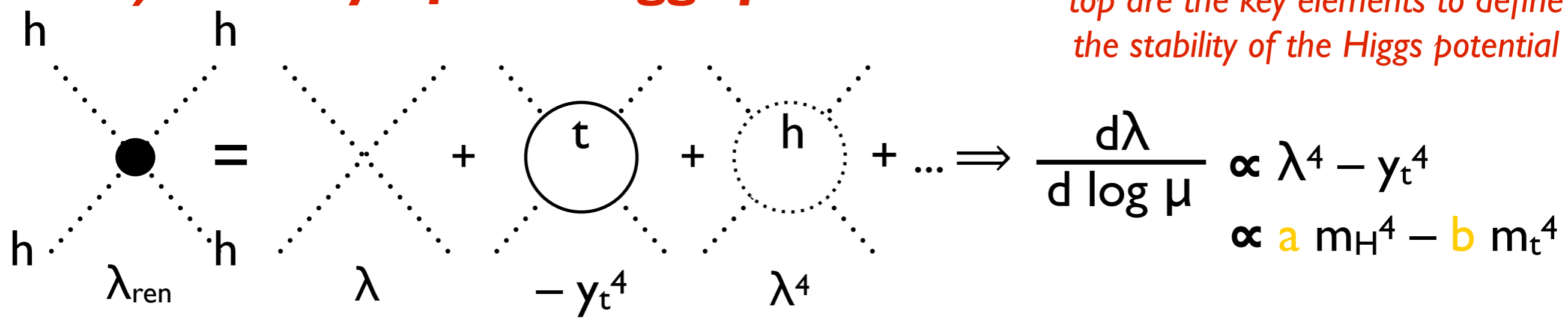
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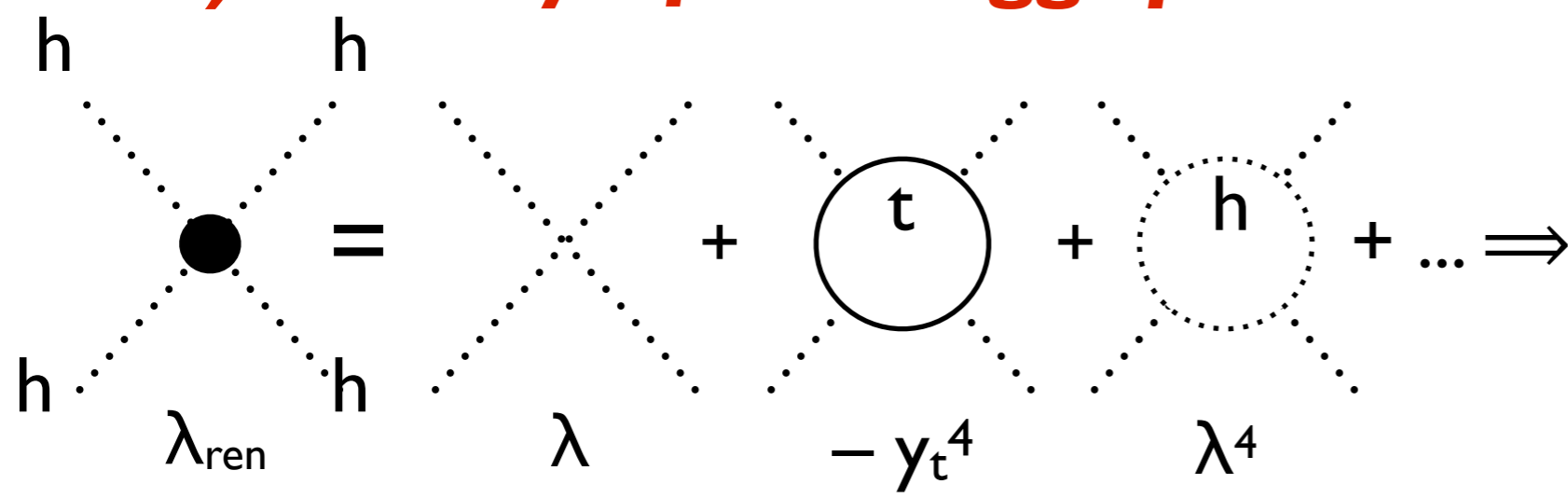
(meta)Stability of the Higgs potential

Higgs selfcoupling and coupling to the top are the key elements to define the stability of the Higgs potential



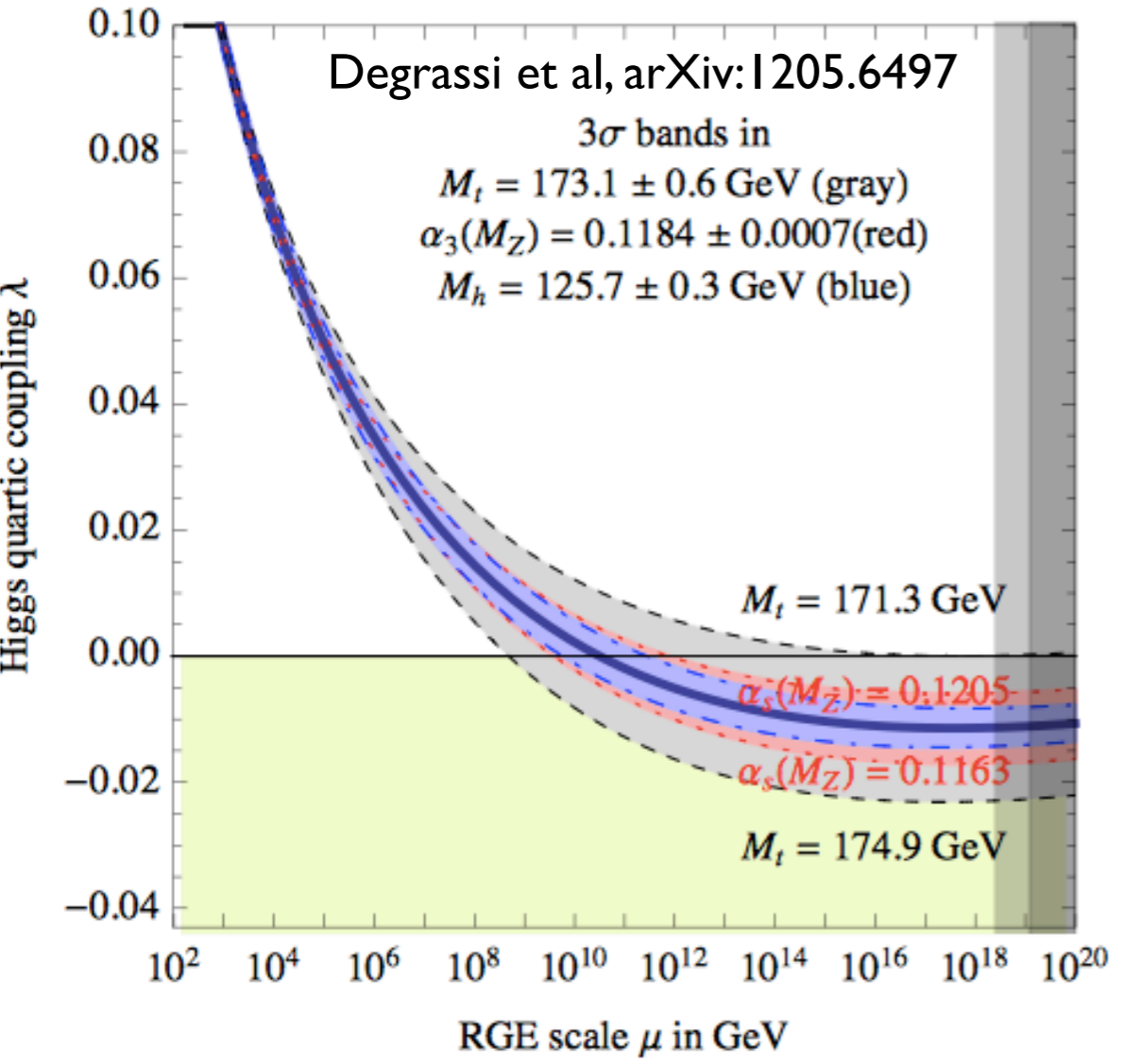
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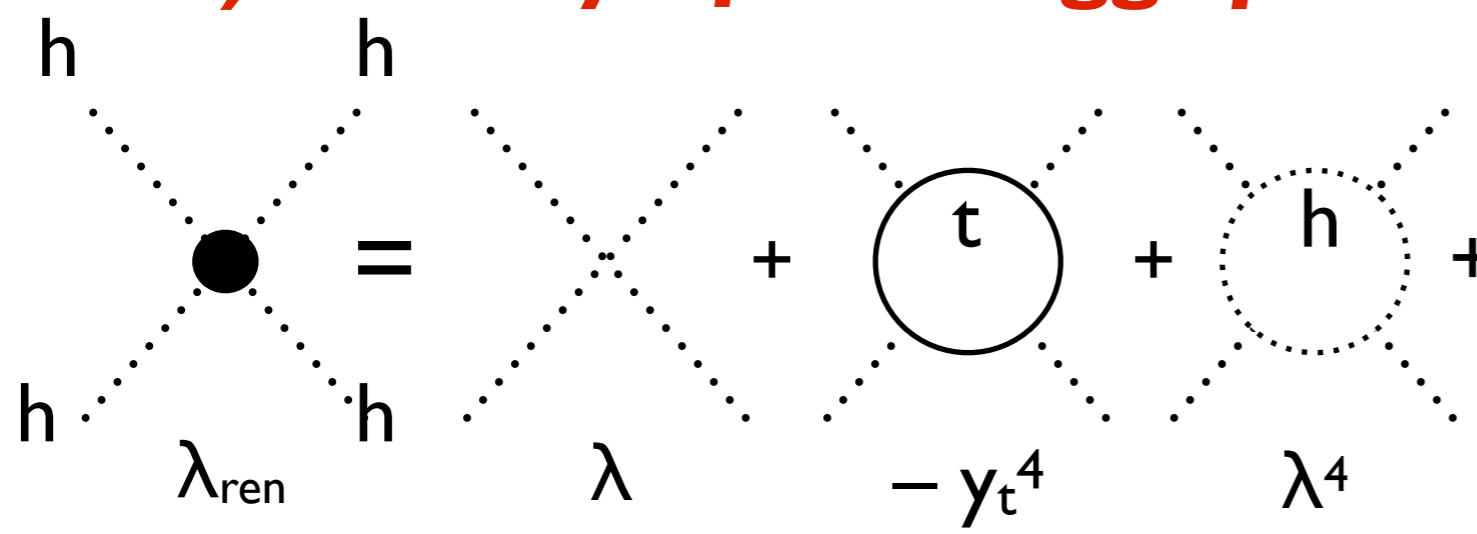
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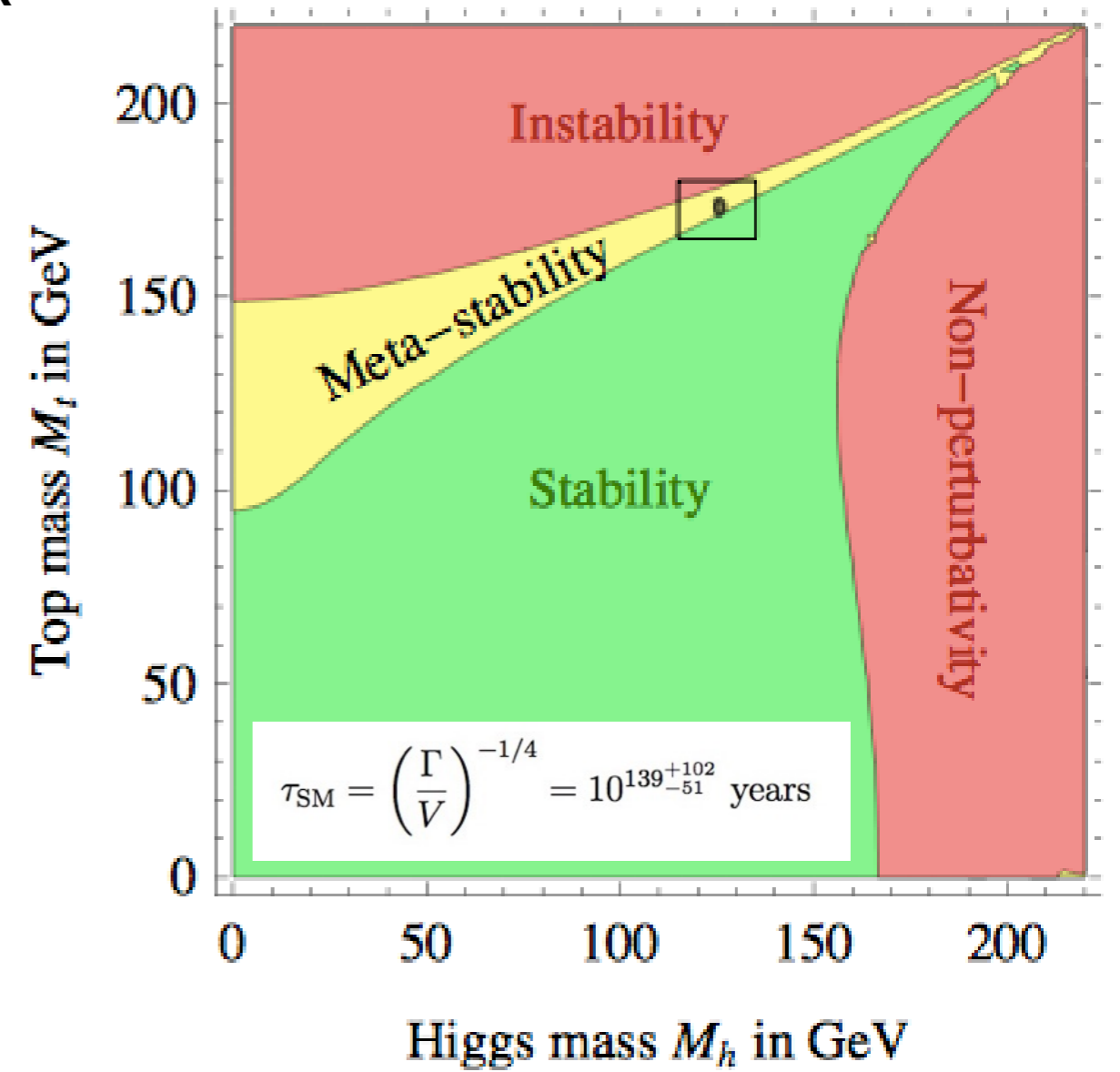
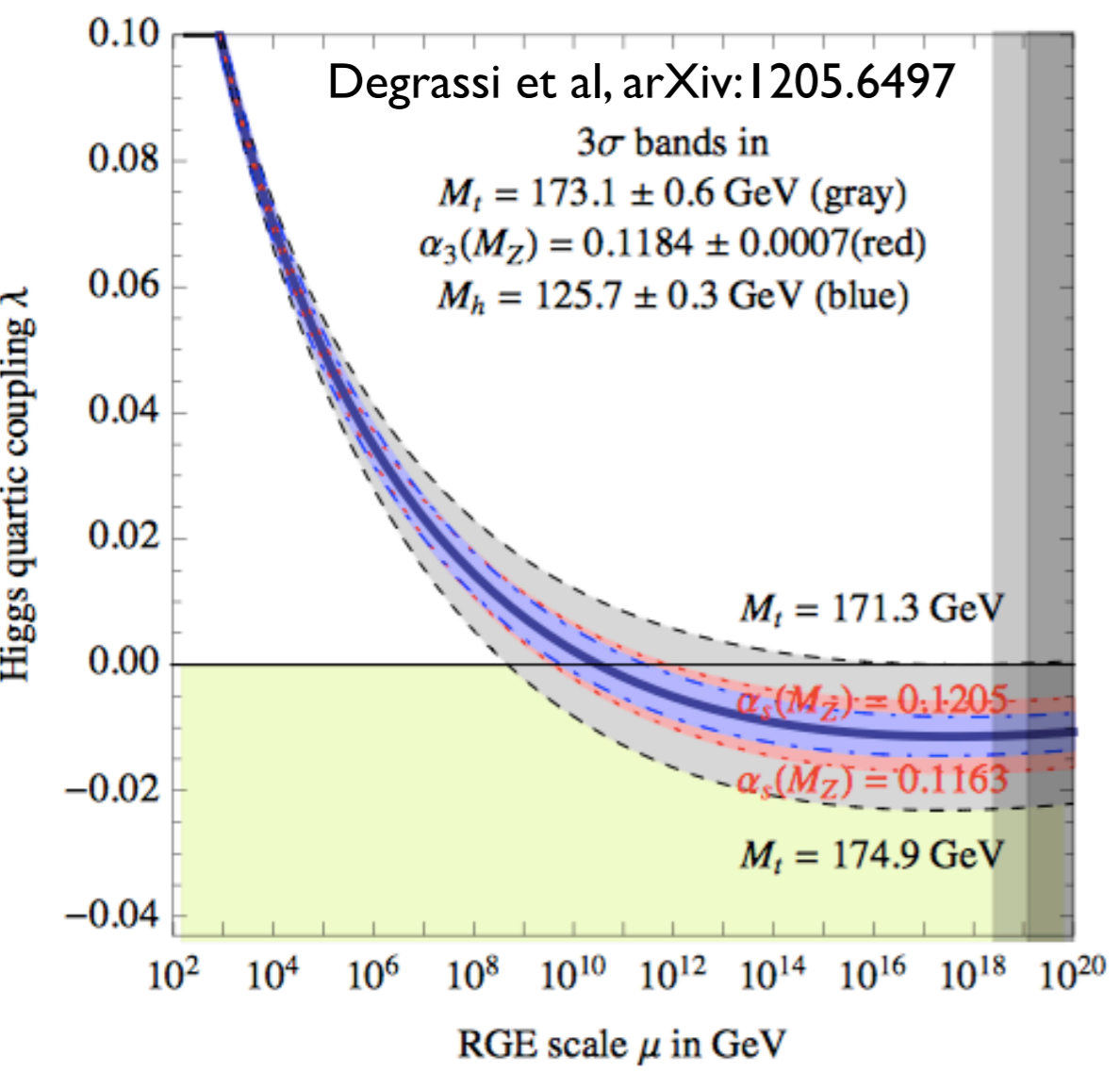
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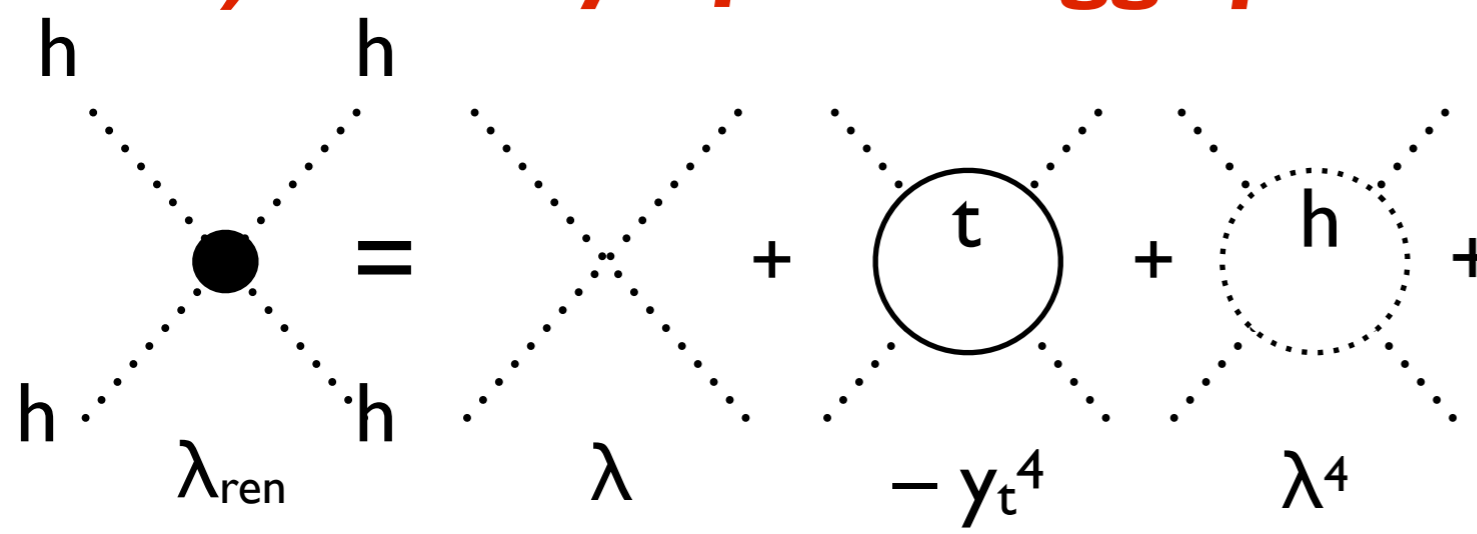
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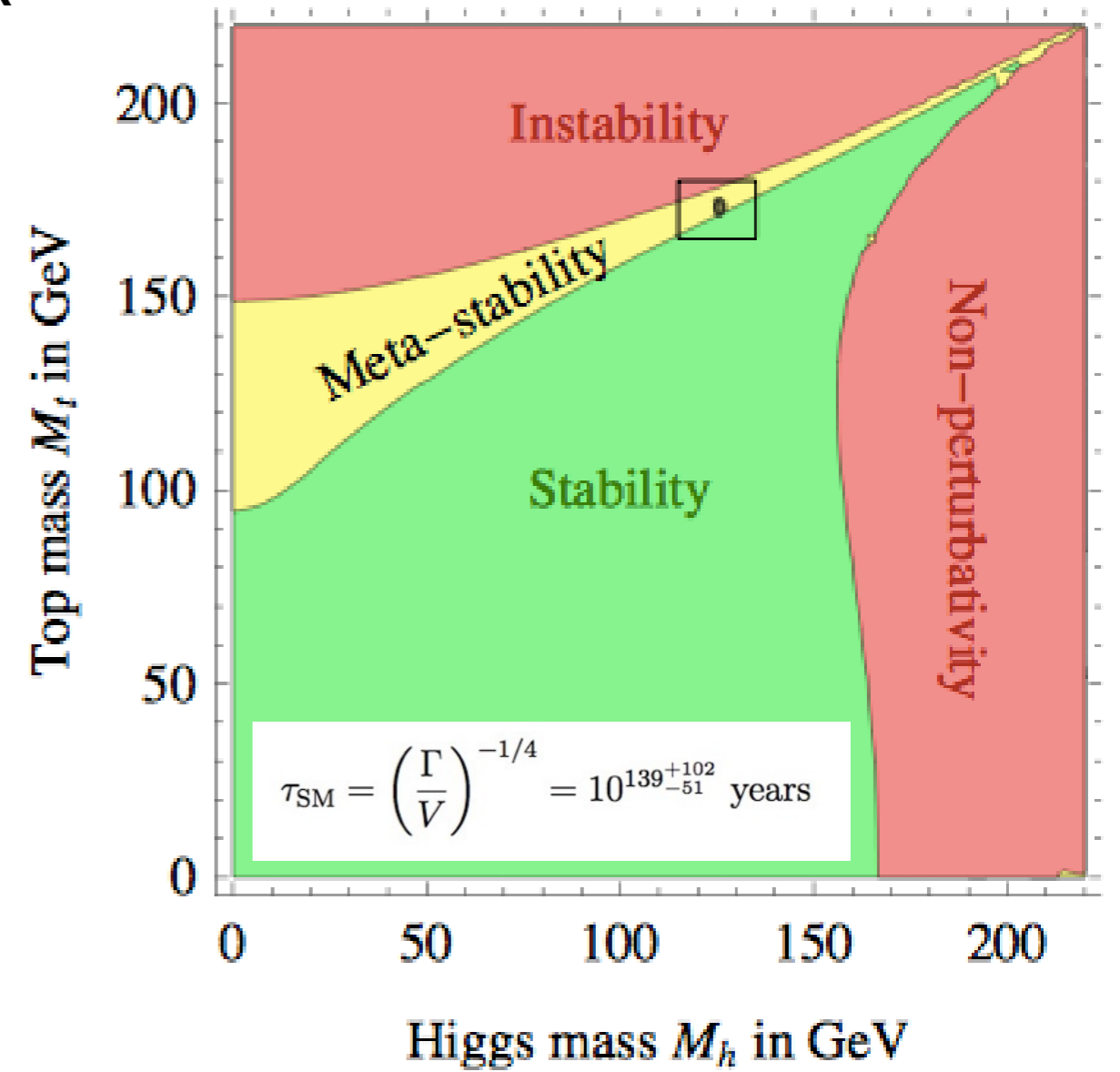
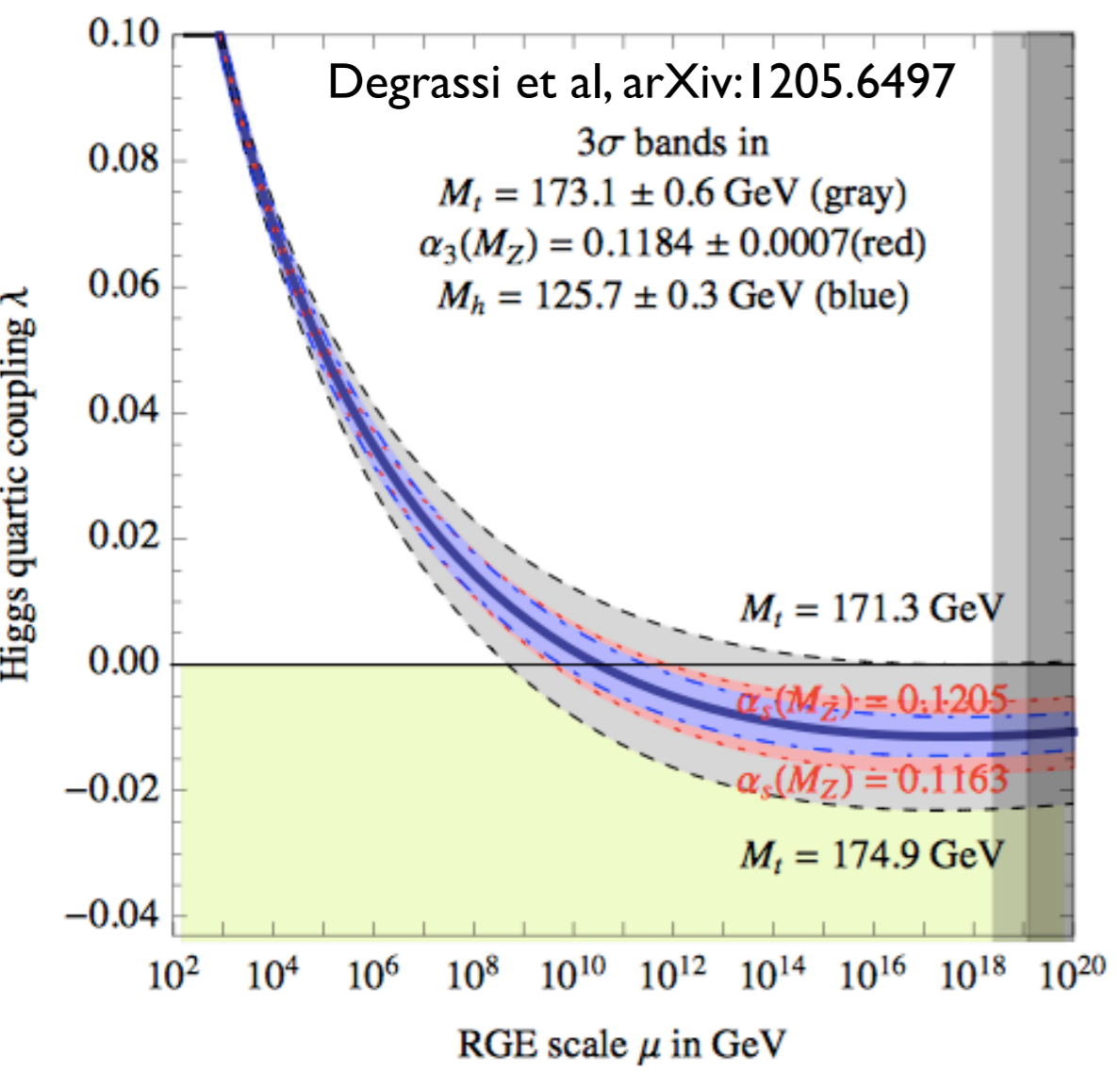
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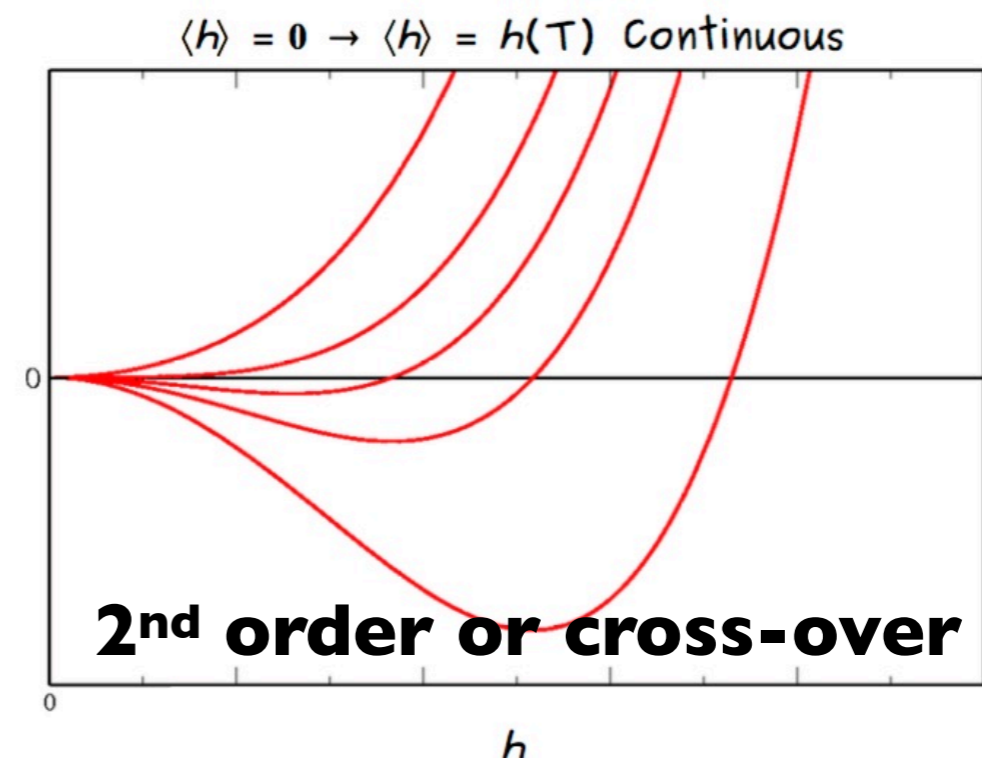
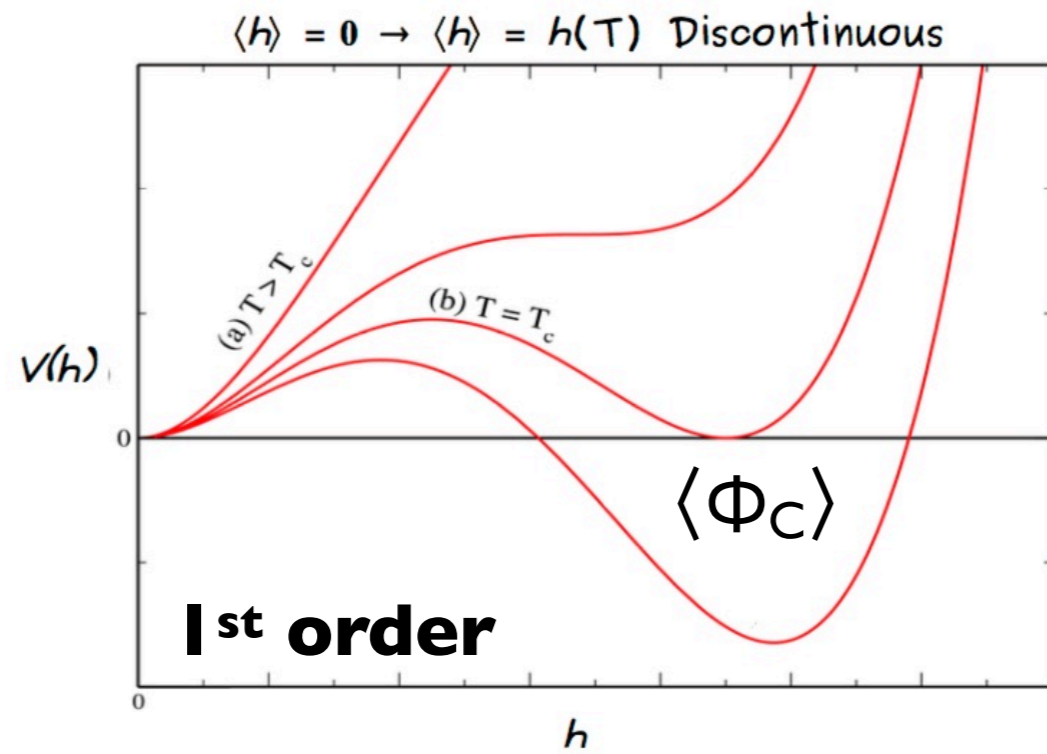


Not an issue of concern for the human race.... but the closeness of m_{top} to the critical value where the Higgs selfcoupling becomes 0 at M_{Planck} (namely 171.3 GeV) might be telling us something fundamental about the origin of EWSB ... incidentally, $y_{top}=1$ (!?)

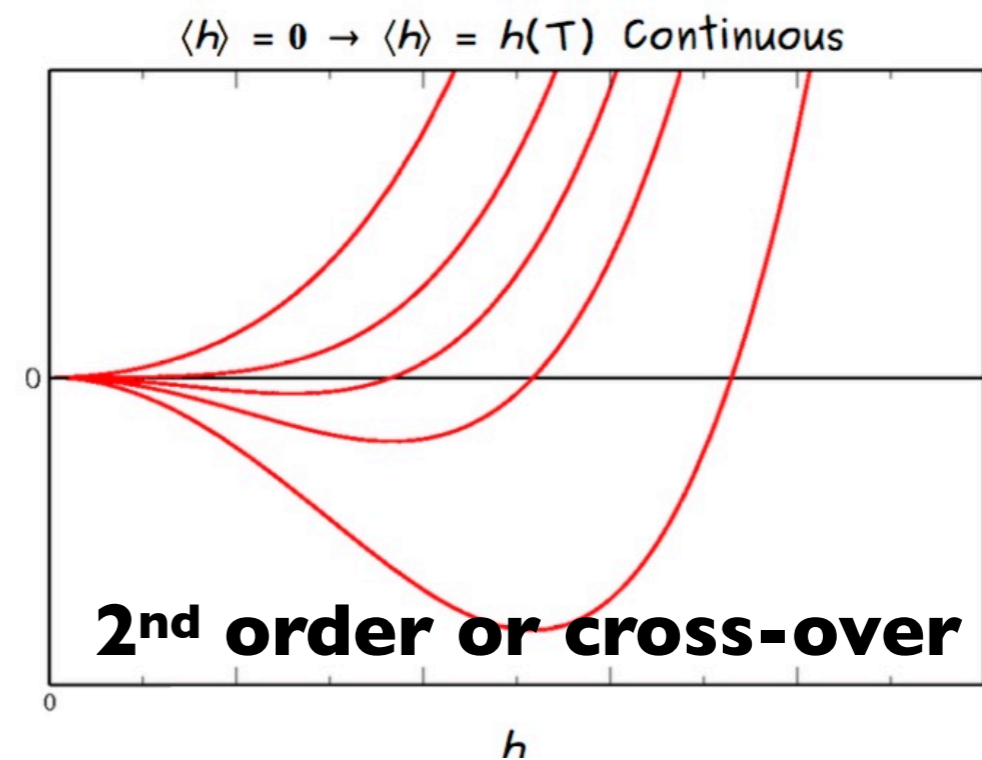
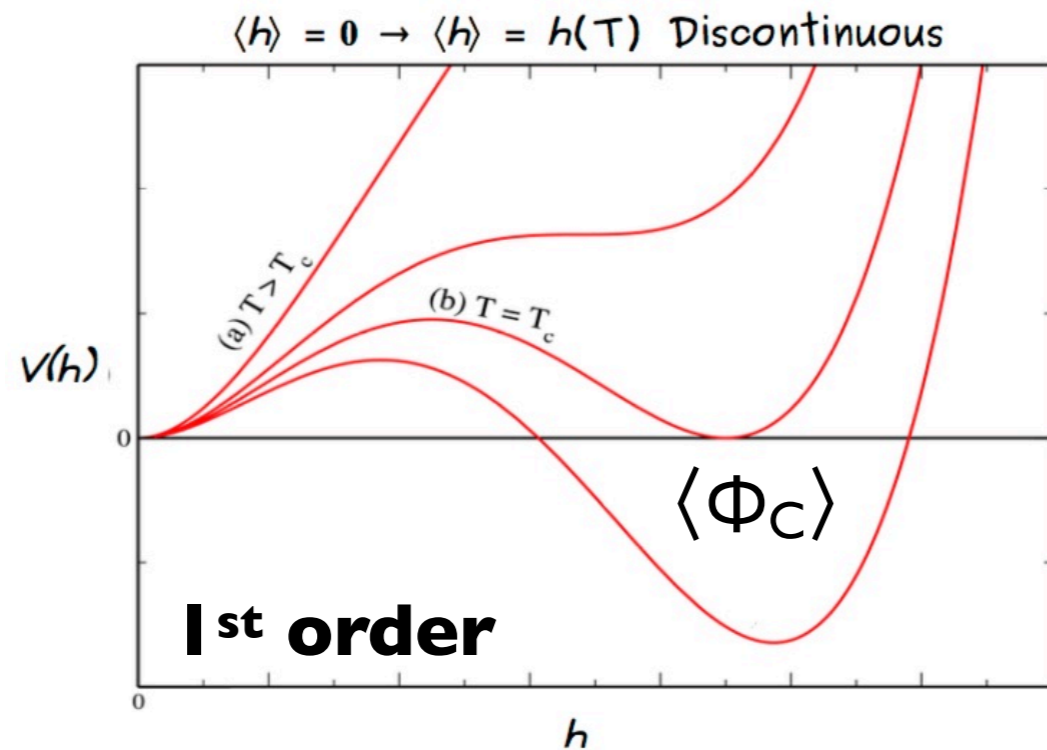
Other important open issues on the Higgs sector

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 - what's the order of the phase transition?
 - are the conditions realized to allow EW baryogenesis?

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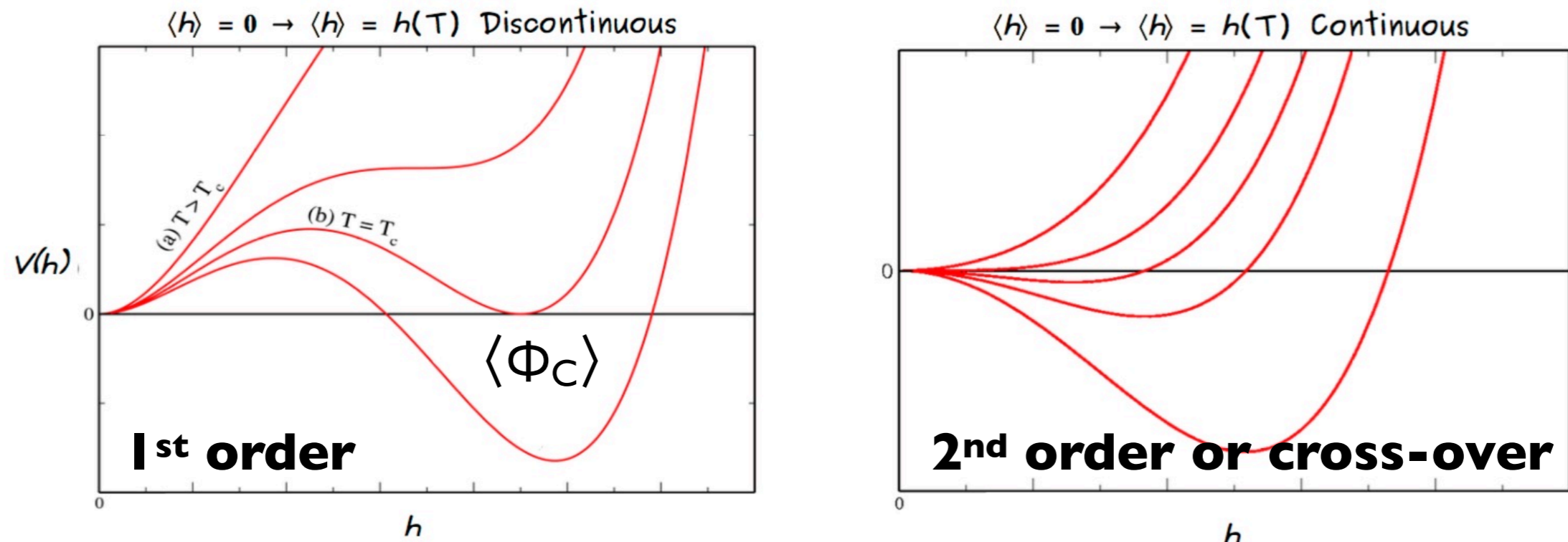
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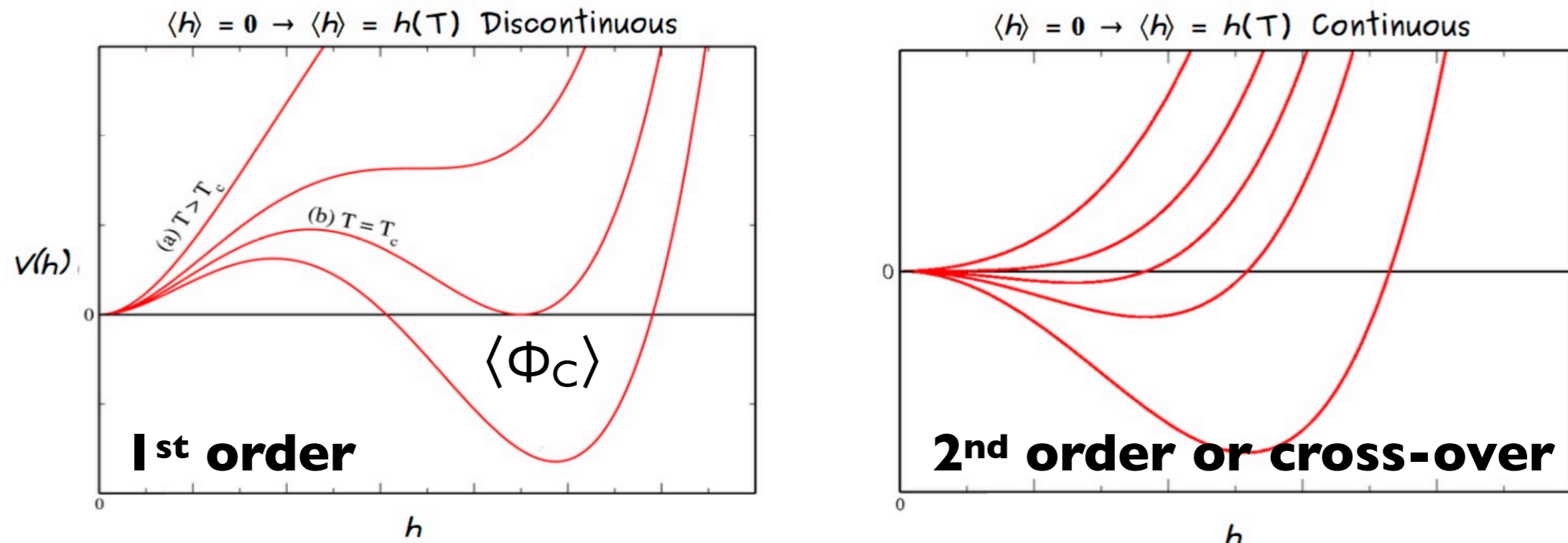
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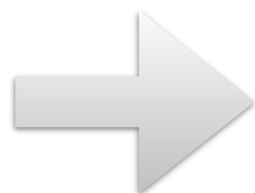


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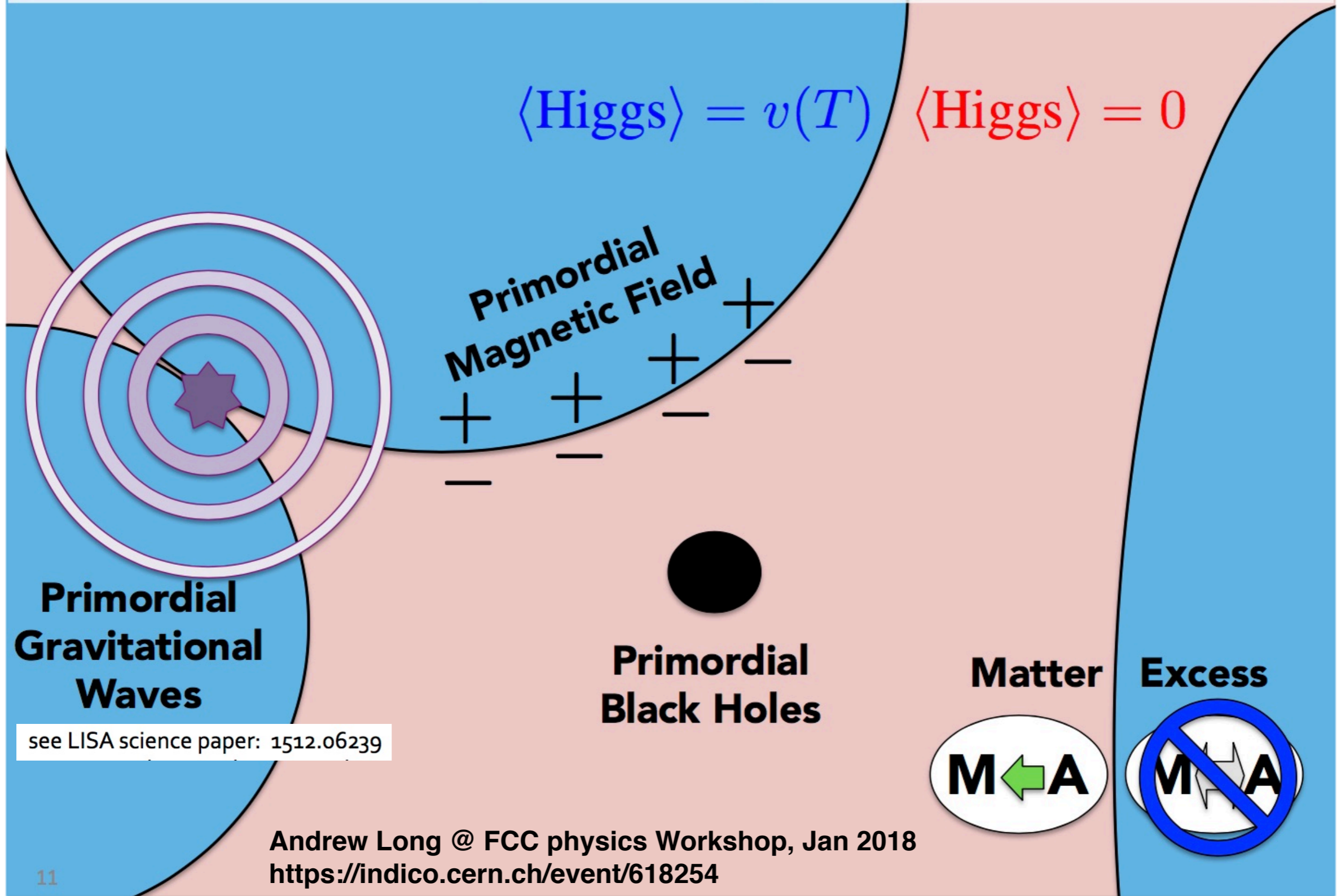
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- Probe higher-order terms of the Higgs potential (selfcouplings)
- Probe the existence of other particles coupled to the Higgs

1st Order EWPT has profound implications for cosmology



Andrew Long @ FCC physics Workshop, Jan 2018
<https://indico.cern.ch/event/618254>

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- Is there a relation between Higgs and Dark Matter?
- etc.etc.

The only way we know how to address these questions is by directly studying the properties of the Higgs boson, produced in a collider

What are we talking about when we're talking future colliders: at CERN...



pp @ 14 TeV, 3ab^{-1}

✓ Approved
2026-37

What are we talking about when we're talking future colliders: at CERN...



pp @ 14 TeV, 3ab⁻¹

**✓ Approved
2026-37**



e⁺e⁻ @ 380 GeV, 1.5 & ~3 TeV

**CDR 2012+
update '16**

CDR: Conceptual Design Report

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CDR: Conceptual Design Report



CDR (end '18)

100km tunnel

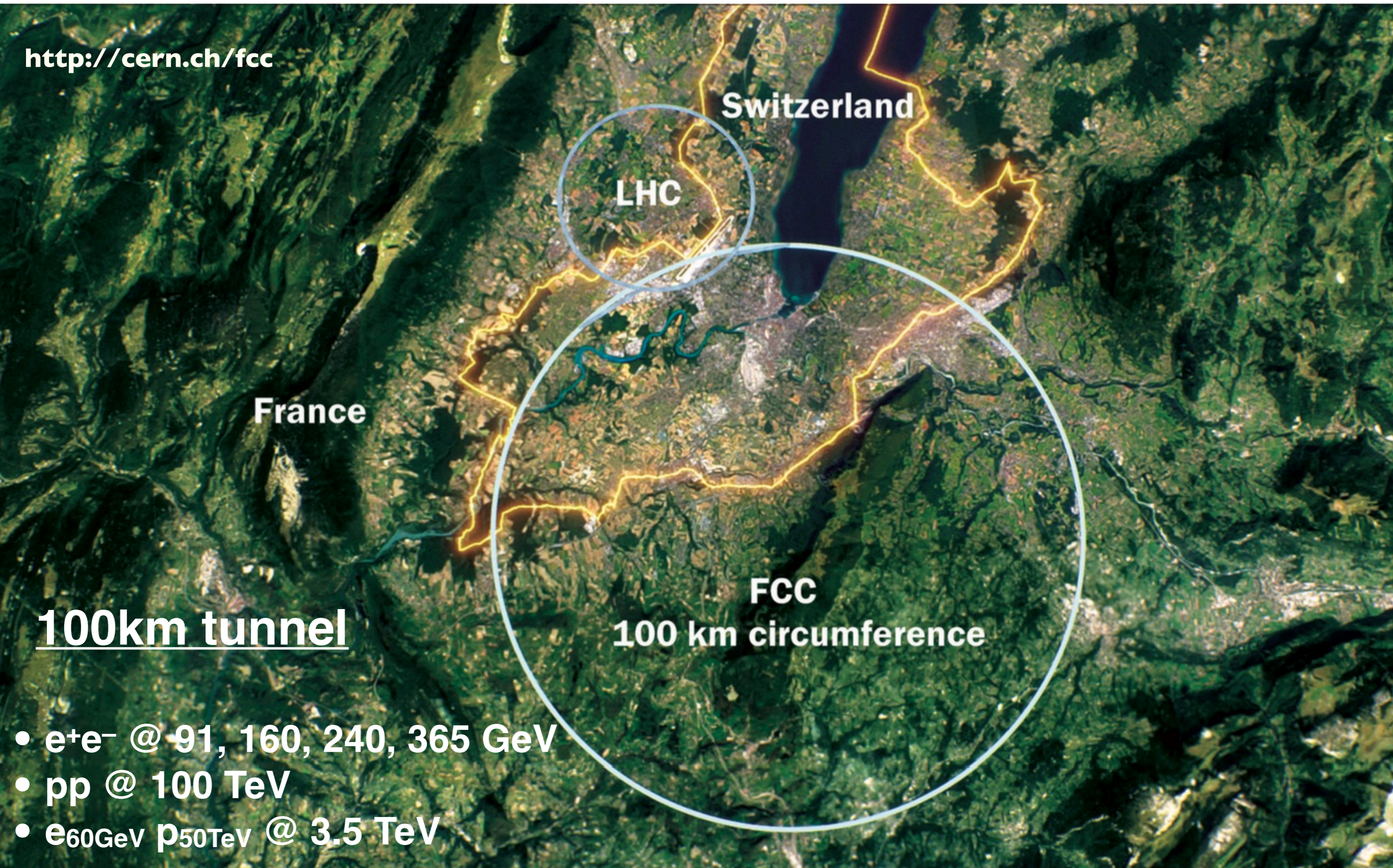
- pp @ 100 TeV
- e⁺e⁻ @ 91, 160, 240, 365 GeV
- e_{60GeV} p_{50TeV} @ 3.5 TeV

LHC tunnel: HE-LHC

- pp @ 27 TeV, 15ab⁻¹

Future Circular Collider

<http://cern.ch/fcc>



... and in the rest of the world:



e^+e^- @ 250, 350, 500 GeV

TDR 2012,
decision postponed
to end 2020

TDR: Technical Design Report



CDR (Fall '18)

100km tunnel

- e^+e^- @ 91, 240 GeV (but possibly 160 & 350)
- Future possible pp @ ~ 70 TeV and $e_{60\text{GeV}} p_{35\text{TeV}}$

Key question for the future steps of LHC and beyond:
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These two scenarios are a priori equally likely, but they impact in different ways the future of HEP, and thus the assessment of the physics potential of possible future facilities

Readiness to address both scenarios is the best hedge for the field:

- *extended energy/mass reach*
- *sensitivity (to elusive signatures)*
- *precision*

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 - *E.g. match the mass scales for new physics that could be exposed via indirect precision measurements in the EW and Higgs sector*
- Provide firm Yes/No answers to questions like:
 - is there a TeV-scale solution to the hierarchy problem?
 - is DM a thermal WIMP?
 - could the cosmological EW phase transition have been 1st order?
 - could baryogenesis have taken place during the EW phase transition?
 - could neutrino masses have their origin at the TeV scale?
 - ...

Some examples

=>

see Franco's talk for more details about the e^+e^- physics potential.

Here I **focus on pp@100 TeV** and its complementarity/synergy with ee

Higgs couplings (κ fit): HL-LHC \rightarrow FCC-ee \rightarrow hh

	HL-LHC(§)	FCC-ee	FCC-hh
$\delta\Gamma_H / \Gamma_H$ (%)	SM(§§)	1.3	tbd
$\delta g_{HZZ} / g_{HZZ}$ (%)	1.5	0.17	tbd
$\delta g_{HWW} / g_{HWW}$ (%)	1.7	0.43	tbd
$\delta g_{Hbb} / g_{Hbb}$ (%)	3.7	0.61	tbd
$\delta g_{Hcc} / g_{Hcc}$ (%)	~ 70	1.21	tbd
$\delta g_{Hgg} / g_{Hgg}$ (%)	2.5 (gg \rightarrow H)	1.01	tbd
$\delta g_{H\tau\tau} / g_{H\tau\tau}$ (%)	1.9	0.74	tbd
$\delta g_{H\mu\mu} / g_{H\mu\mu}$ (%)	4.3	9.0	0.65 (*)
$\delta g_{H\gamma\gamma} / g_{H\gamma\gamma}$ (%)	1.8	3.9	0.4 (*)
$\delta g_{Htt} / g_{Htt}$ (%)	3.4	—	0.95 (**)
$\delta g_{HZ\gamma} / g_{HZ\gamma}$ (%)	9.8	—	0.9 (*)
$\delta g_{HHH} / g_{HHH}$ (%)	50	~ 30 (indirect)	6.5
BR_{exo} (95%CL)	$BR_{\text{inv}} < 2.5\%$	< 1%	$BR_{\text{inv}} < 0.025\%$

§ M. Cepeda, S. Gori, P. J. Ilten, M. Kado, and F. Riva, (conveners), et al, Higgs Physics at the HL-LHC and HE-LHC, [arXiv:1902.00134](https://arxiv.org/abs/1902.00134)

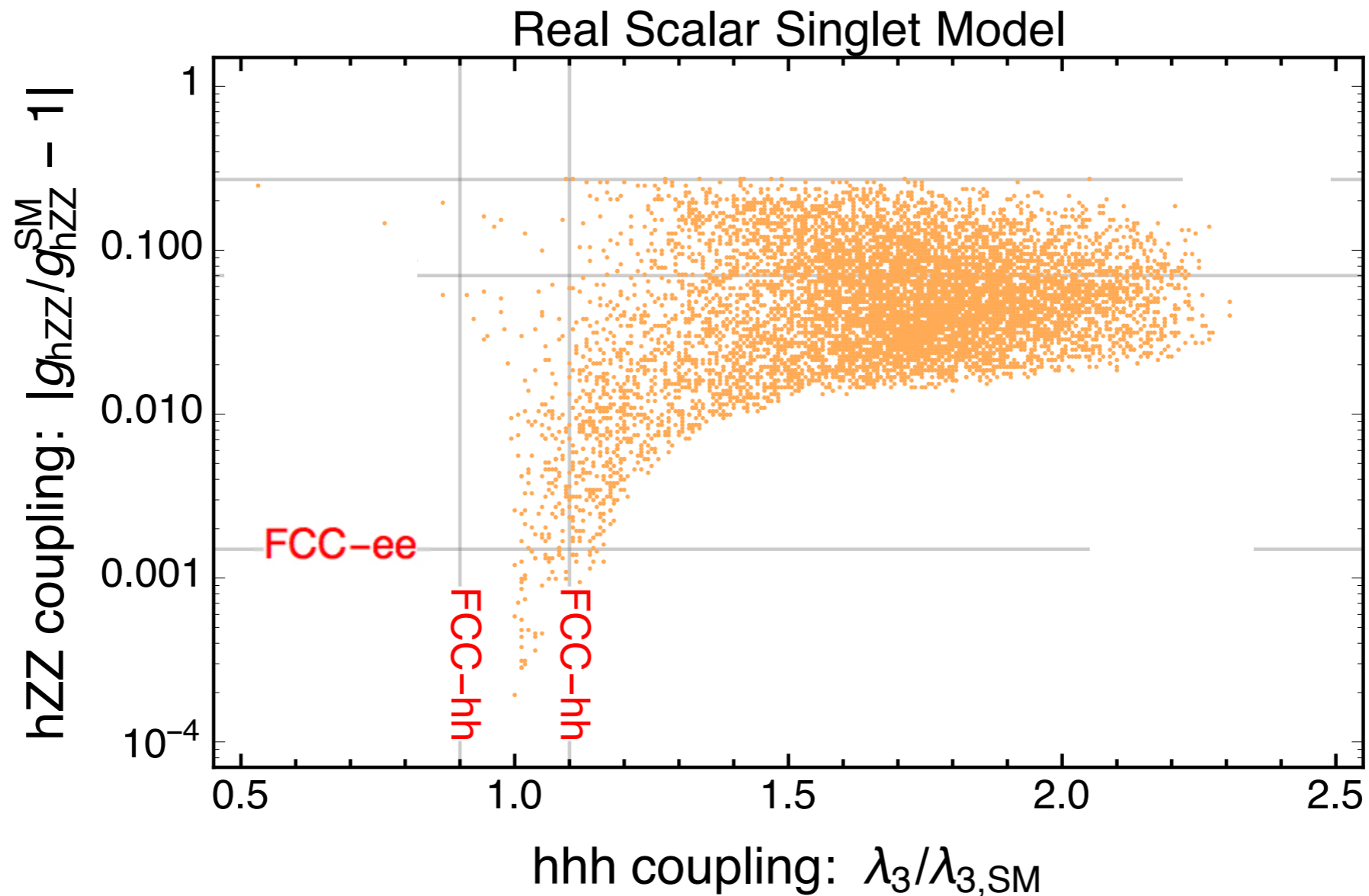
§§ SM width assumed in the global fit. Will be measured to $\sim 20\%$ (68%CL) via off-shell $H \rightarrow 4l$, to $\sim 5\%$ (95%CL) from global fit of Higgs production cross sections.

* From BR ratios wrt $B(H \rightarrow 4\text{lept})$ @ FCC-ee

** From $pp \rightarrow ttH$ / $pp \rightarrow ttZ$, using $B(H \rightarrow bb)$ and ttZ EW coupling @ FCC-ee

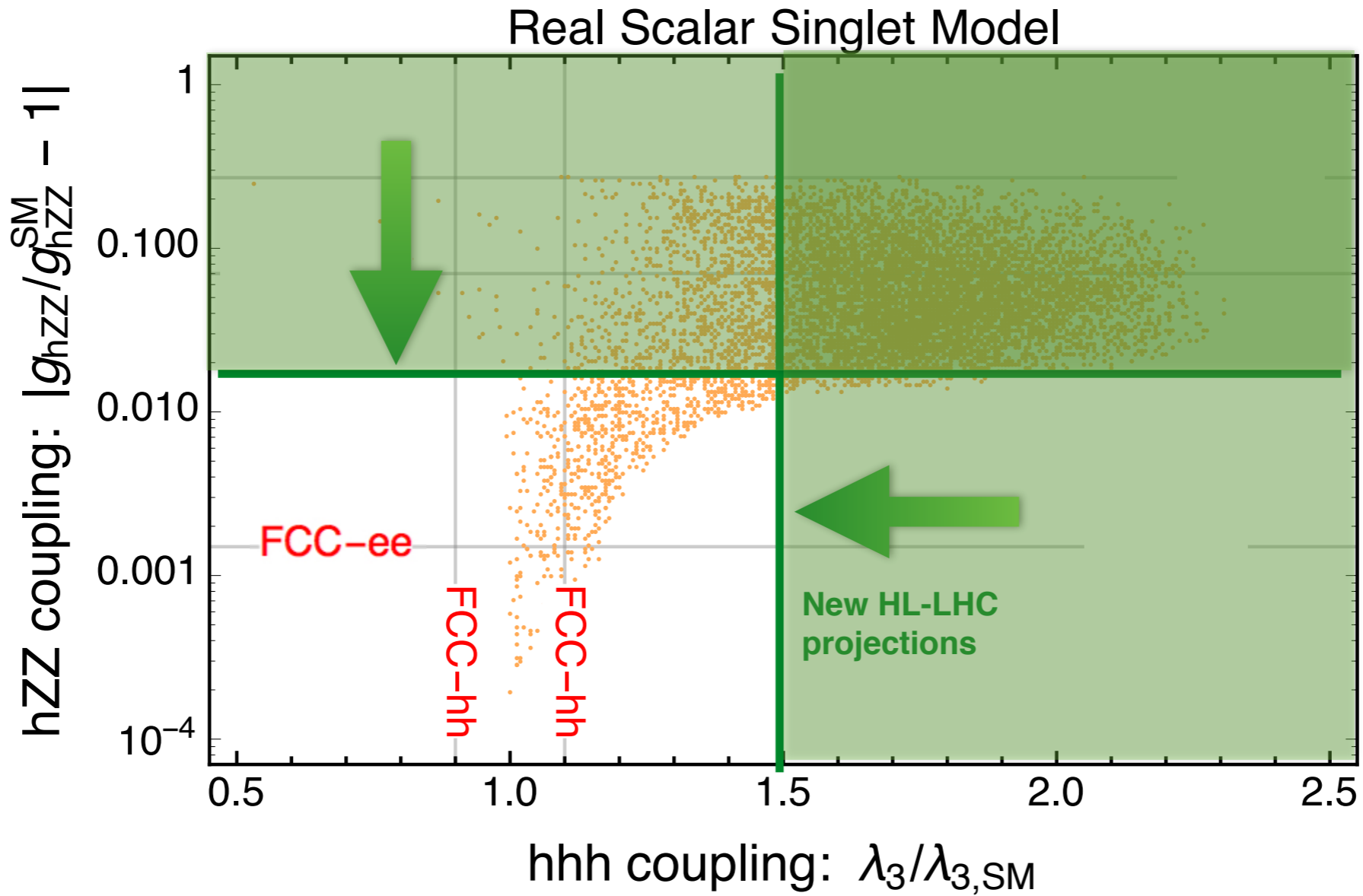
Constraints on models with 1st order phase transition: after the HL-LHC

$$V(H, S) = -\mu^2 (H^\dagger H) + \lambda (H^\dagger H)^2 + \frac{a_1}{2} (H^\dagger H) S + \frac{a_2}{2} (H^\dagger H) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$$



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 \end{aligned}$$



Bringing the HL-LHC sensitivity to the $\pm 50\%$ level, makes a big dent in this class of BSM models!

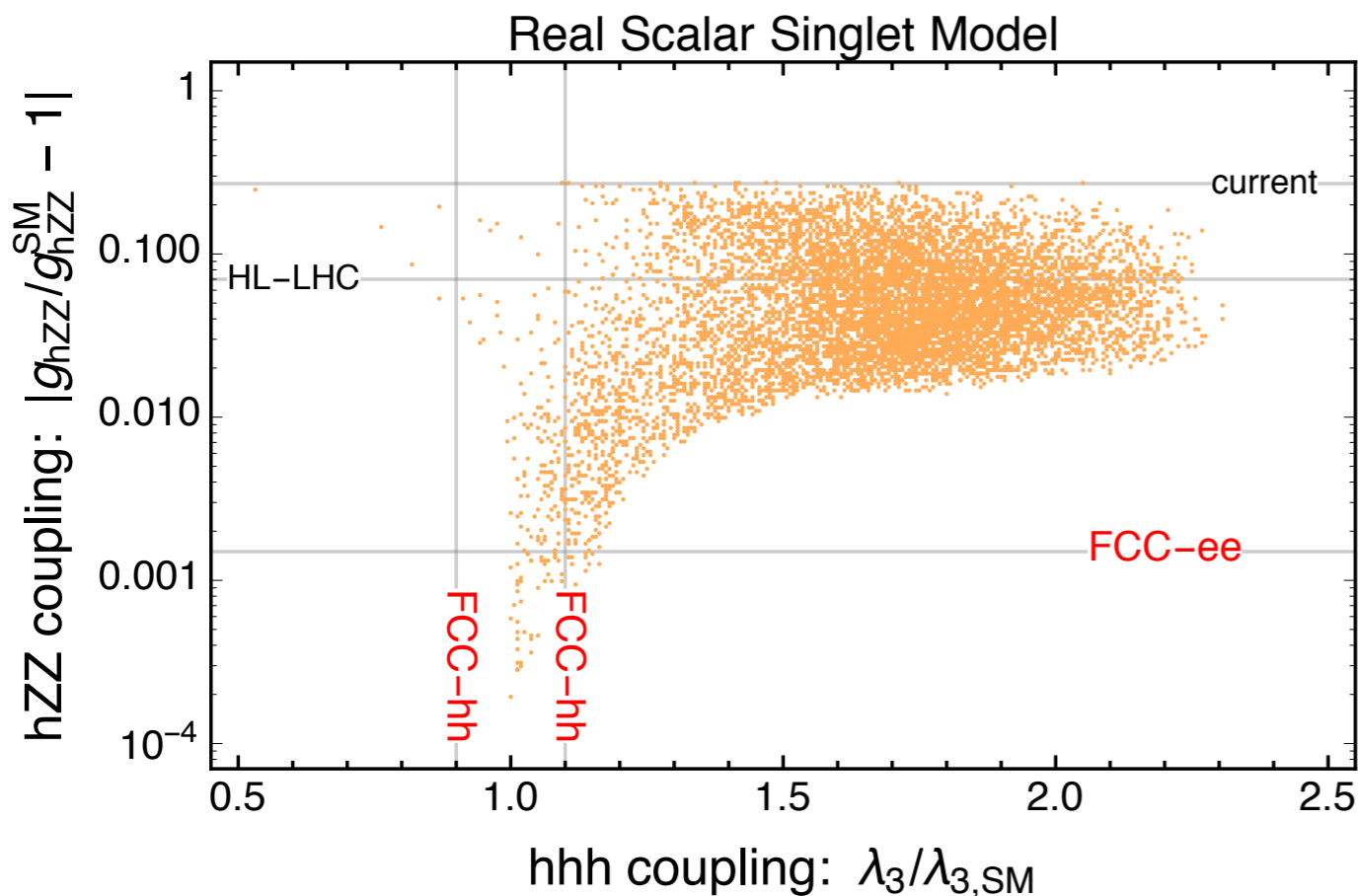
Constraints on models with 1st order phase transition: after the FCC

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Combined constraints from precision Higgs measurements at FCC-ee and FCC-hh

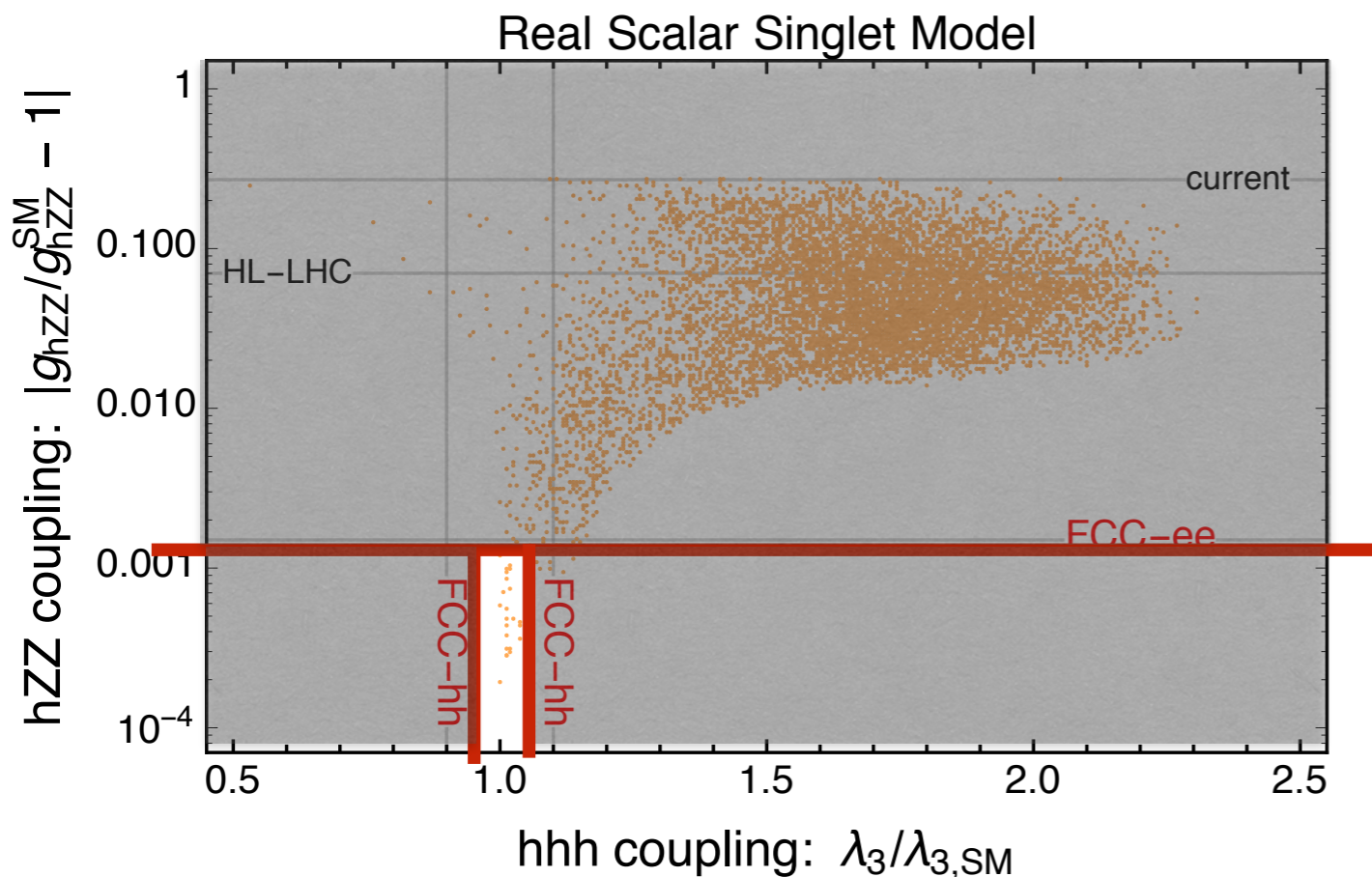


Parameter space scan for a singlet model extension of the Standard Model. The points indicate a first order phase transition.

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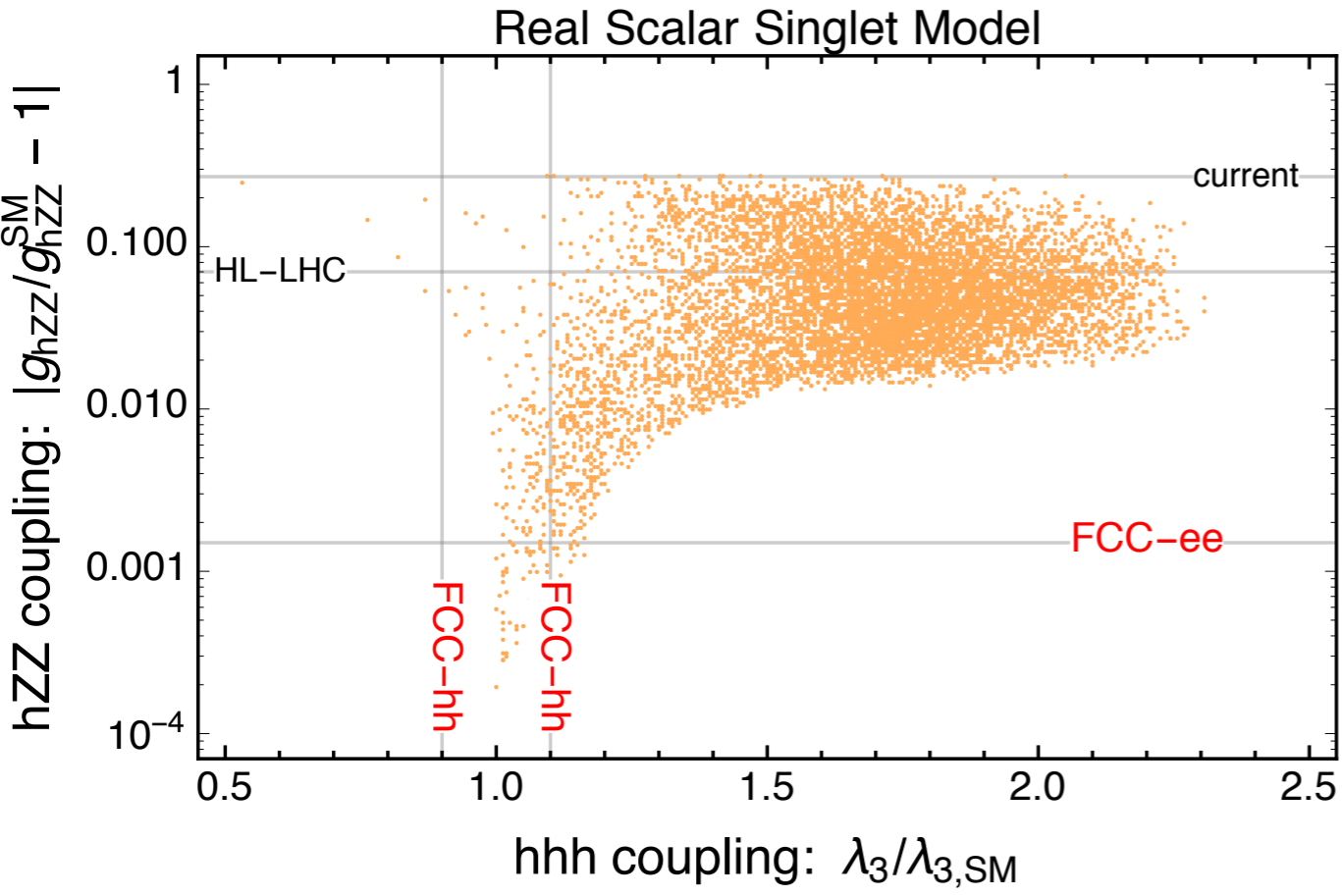


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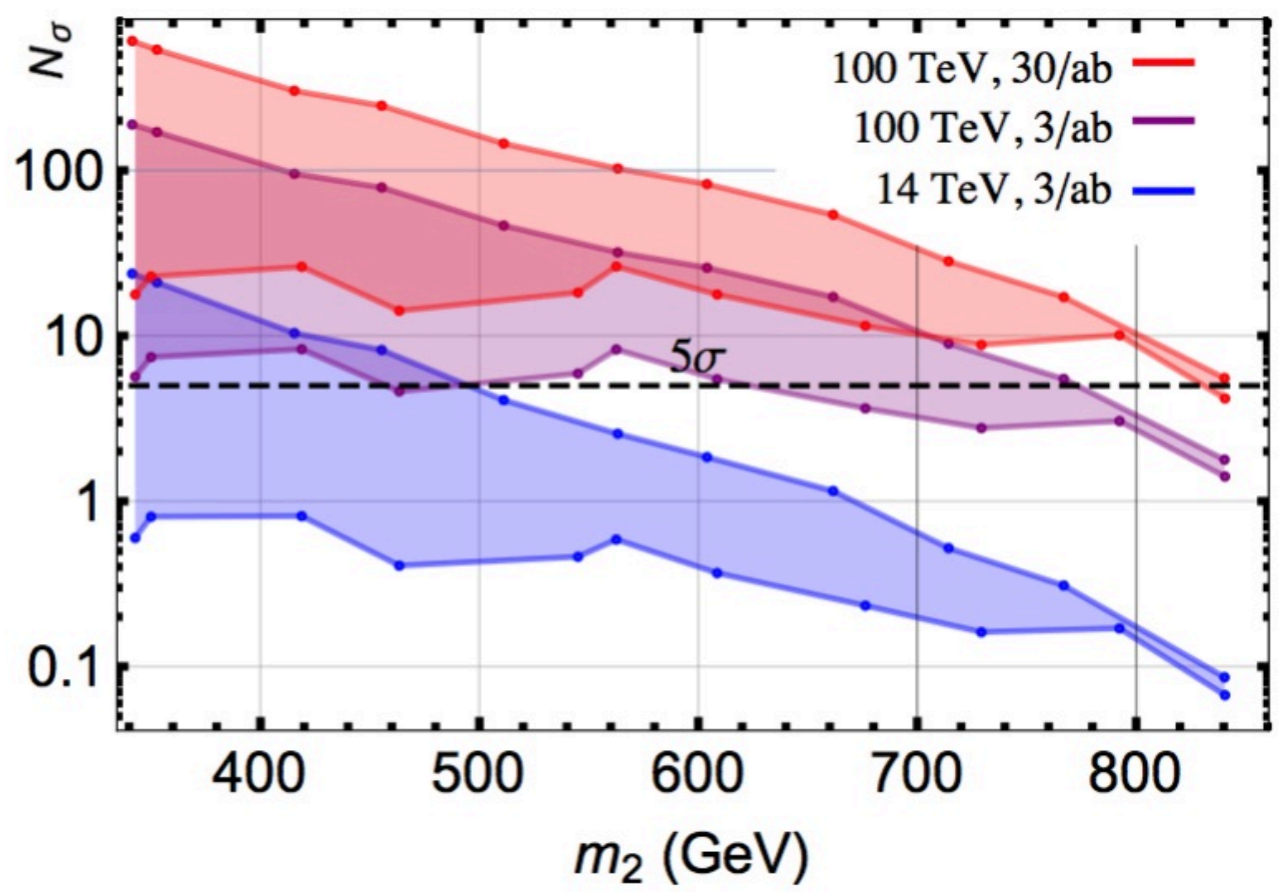
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Direct detection of extra Higgs states at FCC-hh



$h_2 \rightarrow h_1 h_1 \quad (b\bar{b}\gamma\gamma + 4\tau)$
 $(h_2 \sim S, \quad h_1 \sim H)$

On the interplay of precision and kinematic reach in probing new physics indirectly

$$L = L_{SM} + \frac{1}{\Lambda^2} \sum_k \mathcal{O}_k + \dots$$

$$O = | \langle f | L | i \rangle |^2 = O_{SM} [1 + O(\mu^2 / \Lambda^2) + \dots]$$

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$$\delta O \sim \left(\frac{v}{\Lambda}\right)^2 \sim 6\% \left(\frac{\text{TeV}}{\Lambda}\right)^2 \Rightarrow \text{precision probes large } \Lambda$$

e.g. $\delta O = 1\% \Rightarrow \Lambda \sim 2.5 \text{ TeV}$

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Complementarity between super-precise measurements at ee collider and large-Q studies at 100 TeV

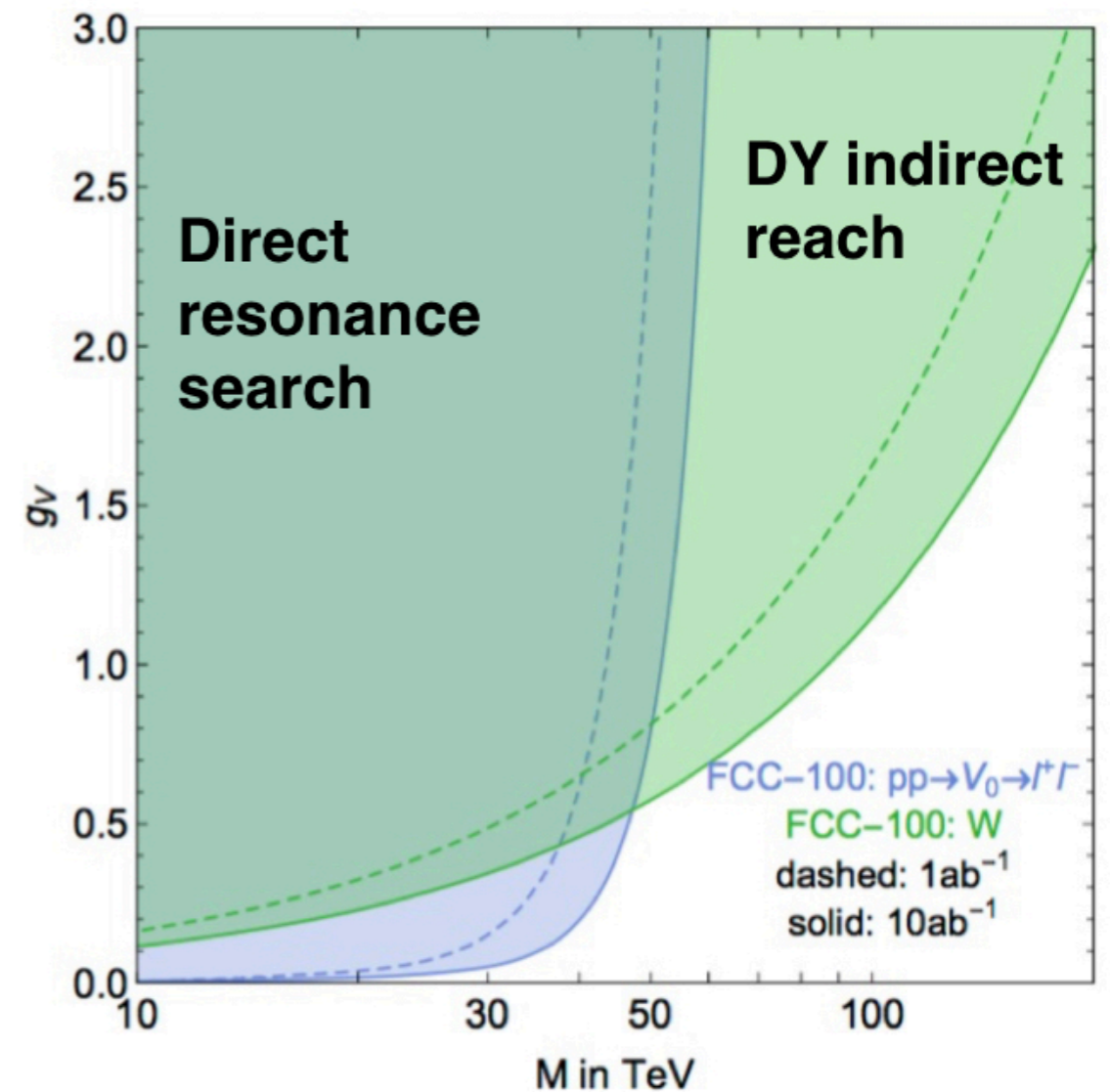
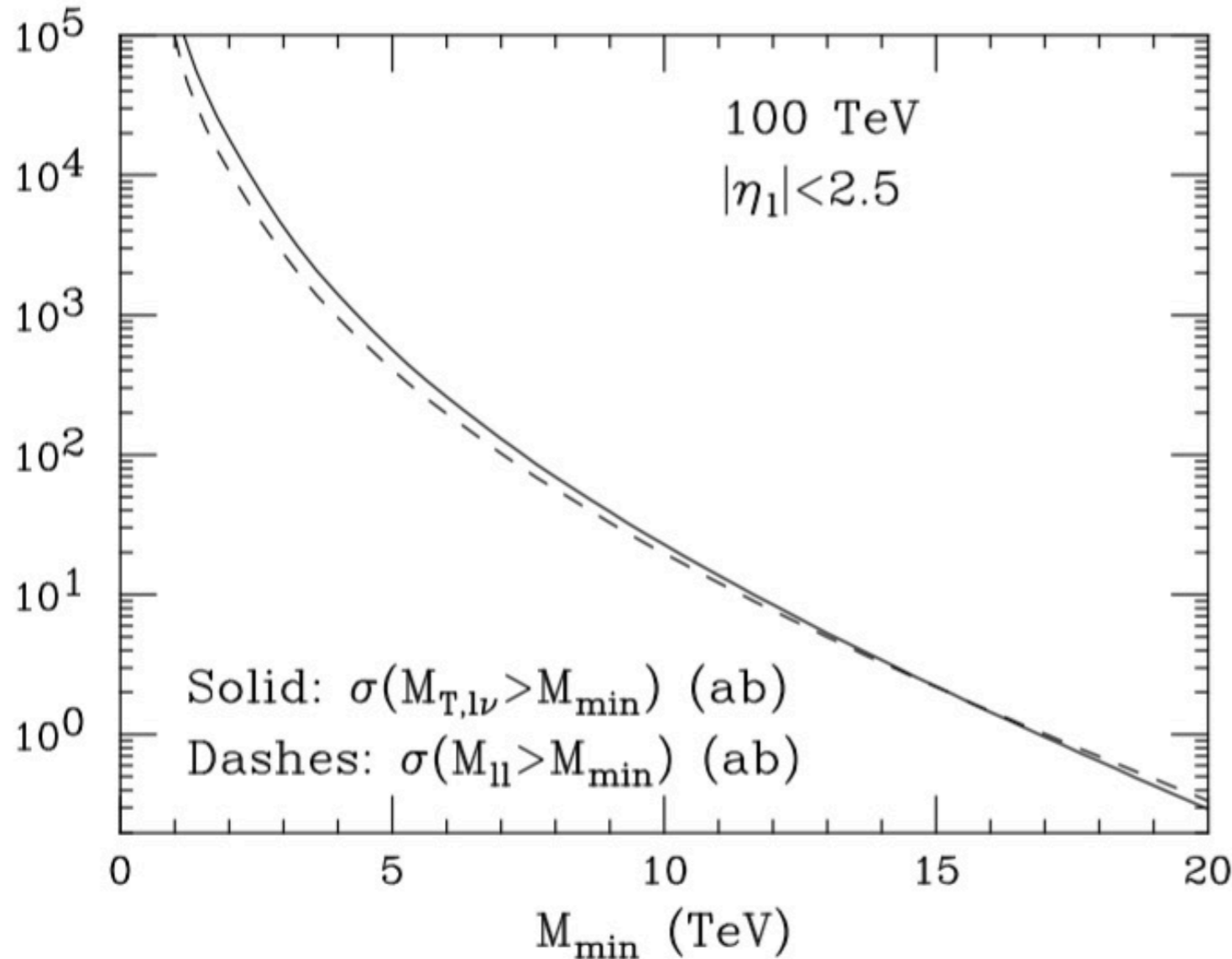
Example: high mass DY

Constraints on Higher-dim op's

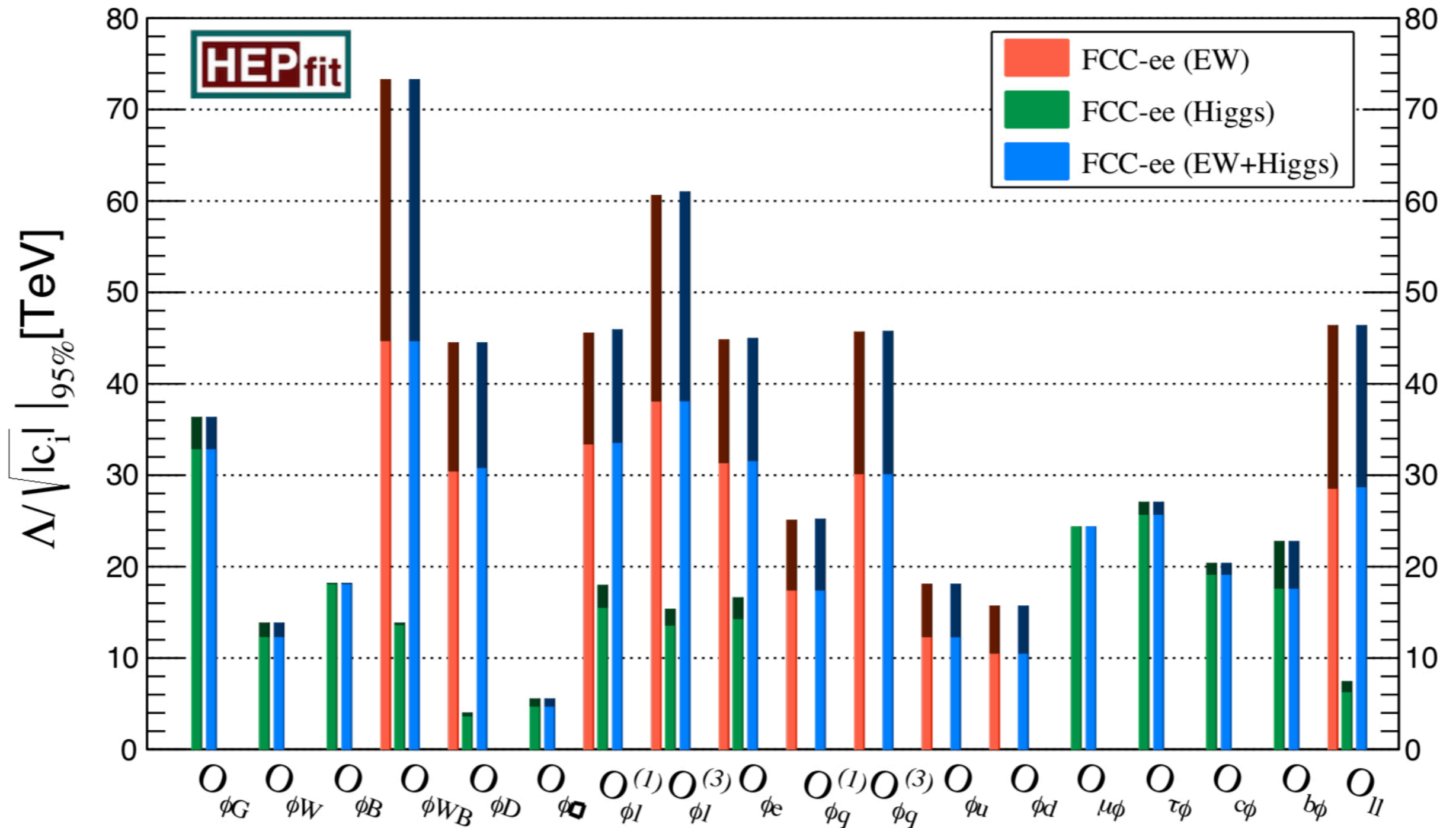
$$\hat{W} = -\frac{W}{4m_W^2} (D_\rho W_{\mu\nu}^a)^2, \quad \hat{Y} = -\frac{Y}{4m_W^2} (\partial_\rho B_{\mu\nu})^2$$

	LEP	LHC 13	FCC 100	ILC	TLEP	CEPC	ILC 500	CLIC 1	CLIC 3	
luminosity	$2 \times 10^7 Z$	0.3/ab	3/ab	10/ab	$10^9 Z$	$10^{12} Z$	$10^{10} Z$	3/ab	1/ab	1/ab
$W \times 10^4$	[-19, 3]	± 0.7	± 0.45	± 0.02	± 4.2	± 1.2	± 3.6	± 0.3	± 0.5	± 0.15
$Y \times 10^4$	[-17, 4]	± 2.3	± 1.2	± 0.06	± 1.8	± 1.5	± 3.1	± 0.2	$\sim \pm 0.5$	$\sim \pm 0.15$

$W / 4m_W^2 < 1 / (100 \text{ TeV})^2$



Global EFT fits to EW and H observables at FCC-ee



Constraints on the coefficients of various EFT op's from a global fit of (i) EW observables, (ii) Higgs couplings and (iii) EW+Higgs combined. Darker shades of each color indicate the results neglecting all SM theory uncertainties. 30

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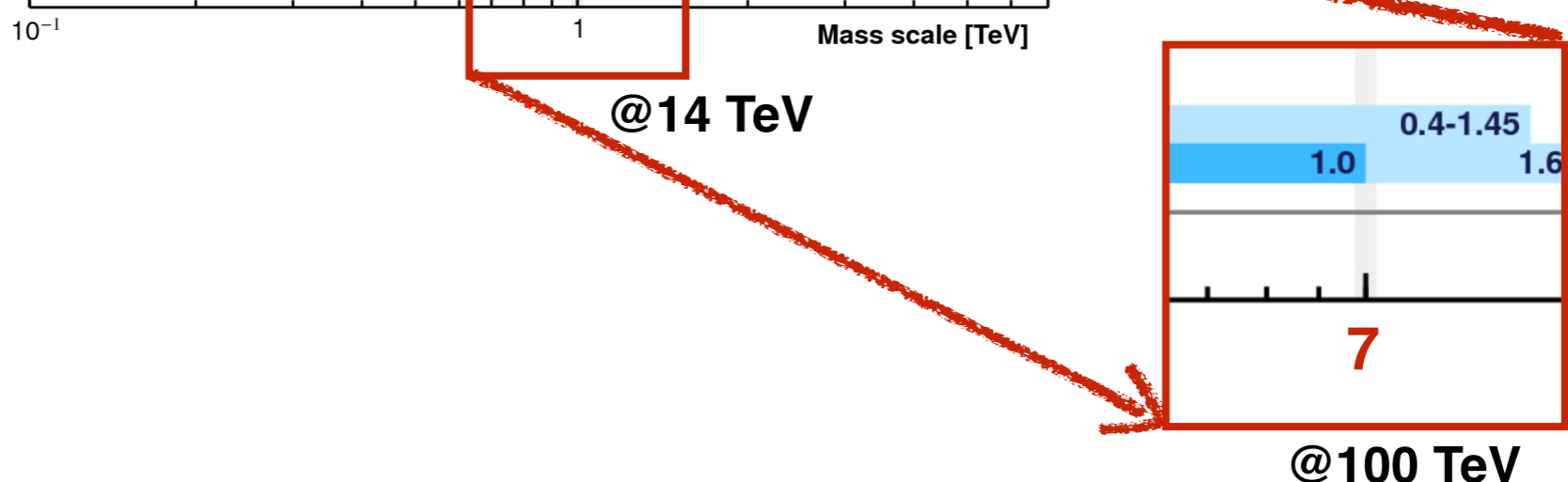
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 - *$m(e) \Rightarrow$ just a parameter; $m(p) \Rightarrow$ just QCD dynamics; Higgs couplings \Rightarrow ???*

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 - *$m(e) \Rightarrow$ just a parameter; $m(p) \Rightarrow$ just QCD dynamics; Higgs couplings \Rightarrow ???*
- ... but who knows how important a given measurement can become, to assess the validity of a future theory?
 - the day some BSM signal is found somewhere, the available precision measurements, will be crucial to establish the nature of the signal, whether they agree or deviate from the SM

Direct discovery potential at 100 TeV

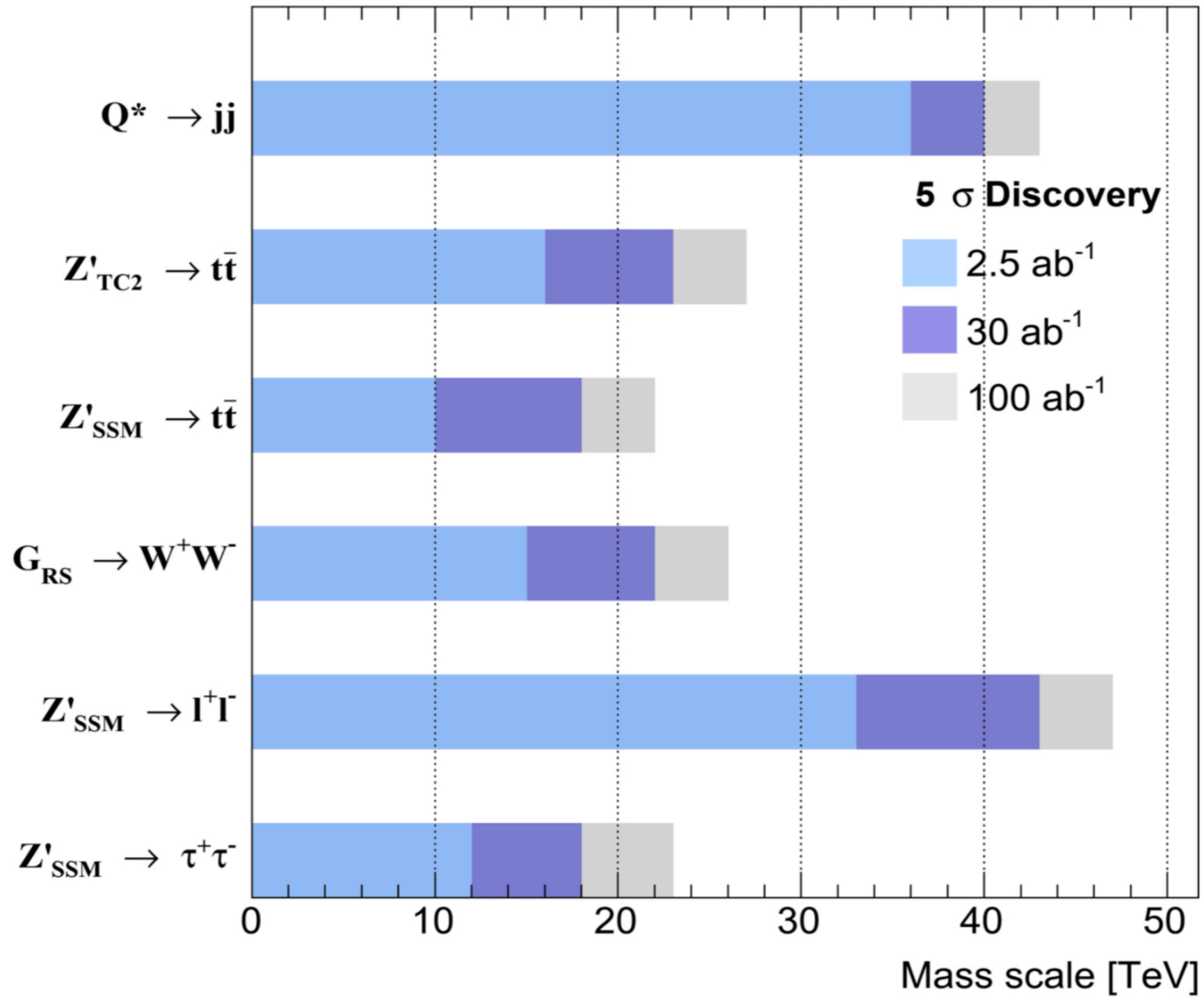
Model	Signature	$\int \mathcal{L} dt \text{ [fb}^{-1}\text{]}$	Mass limit	Reference		
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 e, μ mono-jet	2-6 jets 1-3 jets E_T^{miss} 36.1	\tilde{q} [2x, 8x Degen.] \tilde{q} [1x, 8x Degen.] 0.43 0.71 0.9 1.55	$m(\tilde{\chi}_1^0) < 100 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1712.02332 1711.03301
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0 e, μ	2-6 jets E_T^{miss} 36.1	\tilde{g} \tilde{g} Forbiden 0.95-1.6 2.0	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{\chi}_1^0) = 900 \text{ GeV}$	1712.02332 1712.02332
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	3 e, μ $ee, \mu\mu$	4 jets 2 jets E_T^{miss} 36.1	\tilde{g} \tilde{g} 1.2 1.85	$m(\tilde{\chi}_1^0) < 800 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 50 \text{ GeV}$	1706.03731 1805.11381
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$	0 e, μ 3 e, μ	7-11 jets 4 jets E_T^{miss} 36.1	\tilde{g} \tilde{g} 0.98 1.8	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200 \text{ GeV}$	1708.02794 1706.03731
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$	0-1 e, μ 3 e, μ	3 b 4 jets E_T^{miss} 79.8 36.1	\tilde{g} \tilde{g} 1.25 2.25	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300 \text{ GeV}$	ATLAS-CONF-2018-041 1706.03731
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/t\tilde{\chi}_1^\pm$	Multiple Multiple Multiple	36.1 36.1 36.1	\tilde{b}_1 Forbiden 0.9 \tilde{b}_1 Forbiden 0.58-0.82 \tilde{b}_1 Forbiden 0.7	$m(\tilde{\chi}_1^0) = 300 \text{ GeV}, \text{BR}(b\tilde{\chi}_1^0) = 1$ $m(\tilde{\chi}_1^0) = 300 \text{ GeV}, \text{BR}(b\tilde{\chi}_1^0) = \text{BR}(t\tilde{\chi}_1^\pm) = 0.5$ $m(\tilde{\chi}_1^0) = 200 \text{ GeV}, m(\tilde{\chi}_1^\pm) = 300 \text{ GeV}, \text{BR}(t\tilde{\chi}_1^\pm) = 1$	1708.09266, 1711.03301 1708.09266 1706.03731
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$	0 e, μ	6 b E_T^{miss} 139	\tilde{b}_1 Forbiden \tilde{b}_1 0.23-0.48 0.23-1.35	$\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130 \text{ GeV}, m(\tilde{\chi}_1^0) = 100 \text{ GeV}$ $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130 \text{ GeV}, m(\tilde{\chi}_1^0) = 0 \text{ GeV}$	SUSY-2018-31 SUSY-2018-31
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$	0-2 e, μ	0-2 jets/1-2 b E_T^{miss} 36.1	\tilde{t}_1 1.0	$m(\tilde{\chi}_1^0) = 1 \text{ GeV}$	1506.08616, 1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1$, Well-Tempered LSP	Multiple	36.1	\tilde{t}_1 0.48-0.84	$m(\tilde{\chi}_1^0) = 150 \text{ GeV}, m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}, \tilde{t}_1 \approx \tilde{t}_L$	1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$	1 $\tau + 1 e, \mu, \tau$	2 jets/1 b E_T^{miss} 36.1	\tilde{t}_1 1.16	$m(\tilde{\tau}_1) = 800 \text{ GeV}$	1803.10178
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$	0 e, μ	2 c E_T^{miss} 36.1	\tilde{t}_1 0.46 0.85 \tilde{t}_1 0.43	$m(\tilde{\chi}_1^0) = 0 \text{ GeV}$ $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 50 \text{ GeV}$ $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1805.01649 1805.01649 1711.03301
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$	0 e, μ	mono-jet E_T^{miss} 36.1	\tilde{t}_2 0.32-0.88	$m(\tilde{\chi}_1^0) = 0 \text{ GeV}, m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 180 \text{ GeV}$	1706.03986
EW direct	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via WZ	2-3 e, μ $ee, \mu\mu$	≥ 1 E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.6 $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.17	$m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 10 \text{ GeV}$	1403.5294, 1806.02293 1712.08119
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via WW	2 e, μ	E_T^{miss} 139	$\tilde{\chi}_1^\pm$ 0.42	$m(\tilde{\chi}_1^0) = 0$	ATLAS-CONF-2019-008
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via Wh	0-1 e, μ	2 b E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.68	$m(\tilde{\chi}_1^0) = 0$	1812.09432
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via $\tilde{\ell}_L/\tilde{\nu}$	2 e, μ	E_T^{miss} 139	$\tilde{\chi}_1^\pm$ 1.0	$m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	ATLAS-CONF-2019-008
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}_1\nu(\tau\tilde{\nu}), \tilde{\chi}_2^0 \rightarrow \tilde{\tau}_1\tau(\nu\tilde{\nu})$	2 τ	E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.76 $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.22	$m(\tilde{\chi}_1^0) = 0, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 100 \text{ GeV}, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	1708.07875 1708.07875
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ 2 e, μ	0 jets ≥ 1 E_T^{miss} 139 36.1	$\tilde{\ell}$ 0.7 $\tilde{\ell}$ 0.18	$m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\ell}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	ATLAS-CONF-2019-008 1712.08119
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 e, μ 4 e, μ	$\geq 3 b$ 0 jets E_T^{miss} 36.1 36.1	\tilde{H} 0.13-0.23 \tilde{H} 0.3 0.29-0.88	$\text{BR}(\tilde{\chi}_1^0 \rightarrow h\tilde{G}) = 1$ $\text{BR}(\tilde{\chi}_1^0 \rightarrow Z\tilde{G}) = 1$	1806.04030 1804.03602
Long-lived particles	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet E_T^{miss} 36.1	$\tilde{\chi}_1^\pm$ 0.46 $\tilde{\chi}_1^\pm$ 0.15	Pure Wino Pure Higgsino	1712.02118 ATL-PHYS-PUB-2017-019
	Stable \tilde{g} R-hadron	Multiple	36.1	\tilde{g} 2.0		1902.01636, 1808.04095
	Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$	Multiple	36.1	\tilde{g} [$\tau(\tilde{g}) = 10 \text{ ns}, 0.2 \text{ ns}$] 2.05 2.4	$m(\tilde{\chi}_1^0) = 100 \text{ GeV}$	1710.04901, 1808.04095
RPV	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\tau\mu/\mu\tau$	$e\mu, e\tau, \mu\tau$	3.2	$\tilde{\nu}_\tau$ 1.9	$\lambda'_{311} = 0.11, \lambda'_{132/133/233} = 0.07$	1607.08079
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\ell\nu\nu$	4 e, μ	0 jets E_T^{miss} 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ [$\lambda'_{333} \neq 0, \lambda'_{12k} \neq 0$] 0.82 1.33	$m(\tilde{\chi}_1^0) = 100 \text{ GeV}$	1804.03602
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq$	4-5 large- R jets	36.1	\tilde{g} [$m(\tilde{\chi}_1^0) = 200 \text{ GeV}, 1100 \text{ GeV}$] 1.3 1.9 \tilde{g} [$\lambda'_{112} = 2e-4, 2e-5$] 1.05 2.0	Large λ'_{112}	1804.03568
	$\tilde{u}, \tilde{t} \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs$	Multiple	36.1	\tilde{g} [$\lambda'_{323} = 2e-4, 1e-2$] 0.55 1.05	$m(\tilde{\chi}_1^0) = 200 \text{ GeV}, \text{bino-like}$	ATLAS-CONF-2018-003
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$	2 jets + 2 b	36.7	\tilde{t}_1 [qq, bs] 0.42 0.61	$m(\tilde{\chi}_1^0) = 200 \text{ GeV}, \text{bino-like}$	ATLAS-CONF-2018-003
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$	2 e, μ 1 μ	2 b DV 36.1 136	\tilde{t}_1 0.4-1.45 \tilde{t}_1 [$1e-10 < \lambda'_{23k} < 1e-8, 3e-10 < \lambda'_{23k} < 3e-9$] 1.0 1.6	$\text{BR}(\tilde{t}_1 \rightarrow q\mu) = 100\%, \text{cos}\theta > 20\%$	1710.07171 1710.05544 ATLAS-CONF-2019-006	



*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

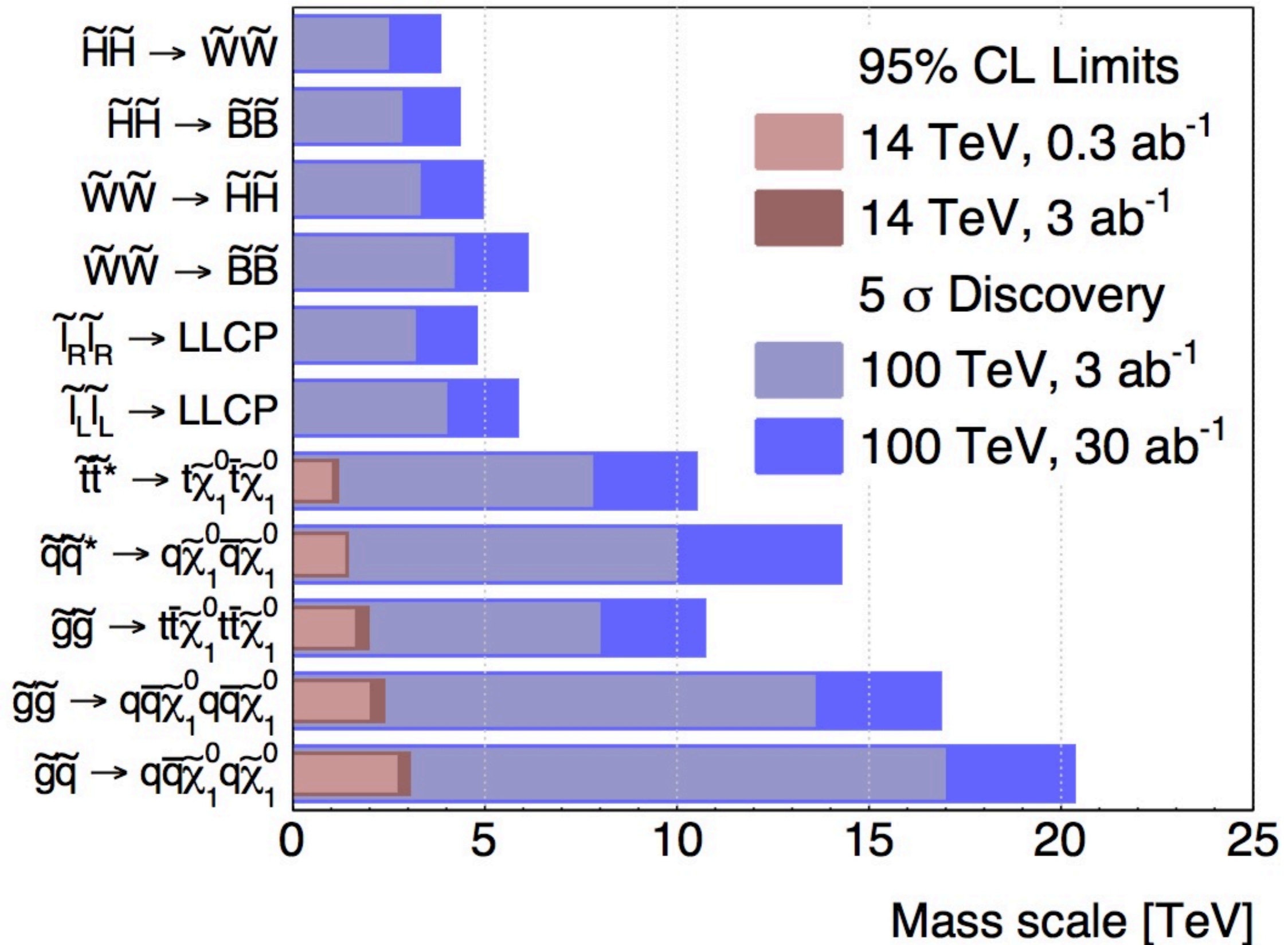
Direct discovery reach: s-channel resonances

FCC-hh Simulation (Delphes), $\sqrt{s} = 100$ TeV



FCC-hh reach ~ 6 x HL-LHC reach

SUSY reach at 100 TeV



Dark Matter

- DM could be explained by BSM models that would leave no signature at any future collider (e.g. axions).
- More in general, no experiment can guarantee an answer to the question "what is DM?"
- Scenarios in which DM is a WIMP are however compelling and theoretically justified
- **We would like to understand whether a future collider can answer more specific questions, such as:**
 - do WIMPS contribute to DM?
 - can WIMPS, detectable in direct and indirect (DM annihilation) experiments, be discovered at future colliders? Is there sensitivity to the explicit detection of DM-SM mediators?
 - what are the opportunities w.r.t. new DM scenarios (e.g. interacting DM, asymmetric DM,)?

WIMP DM theoretical constraints

For particles held in equilibrium by pair creation and annihilation processes, ($\chi \chi \leftrightarrow \text{SM}$)

$$\Omega_{\text{DM}} h^2 \sim \frac{10^9 \text{GeV}^{-1}}{M_{\text{pl}}} \frac{1}{\langle \sigma v \rangle}$$

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$$\Omega_{\text{DM}} h^2 \sim 0.12 \times \left(\frac{M_{\text{DM}}}{2 \text{TeV}} \right)^2 \left(\frac{0.3}{g_{\text{eff}}} \right)^4$$

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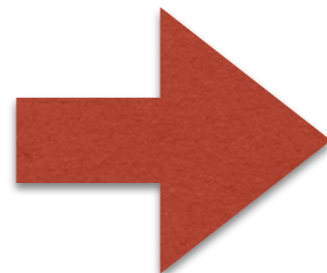
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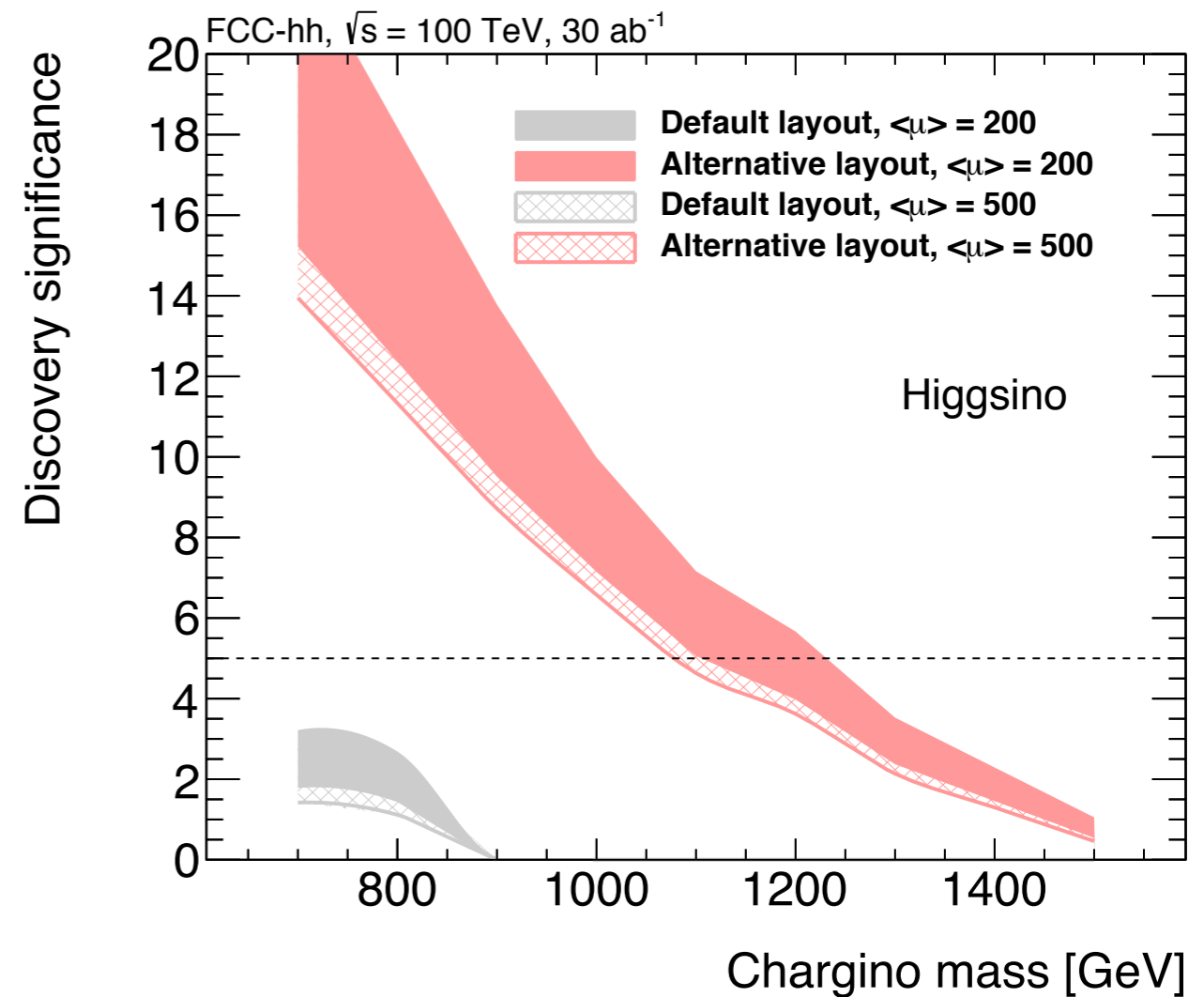
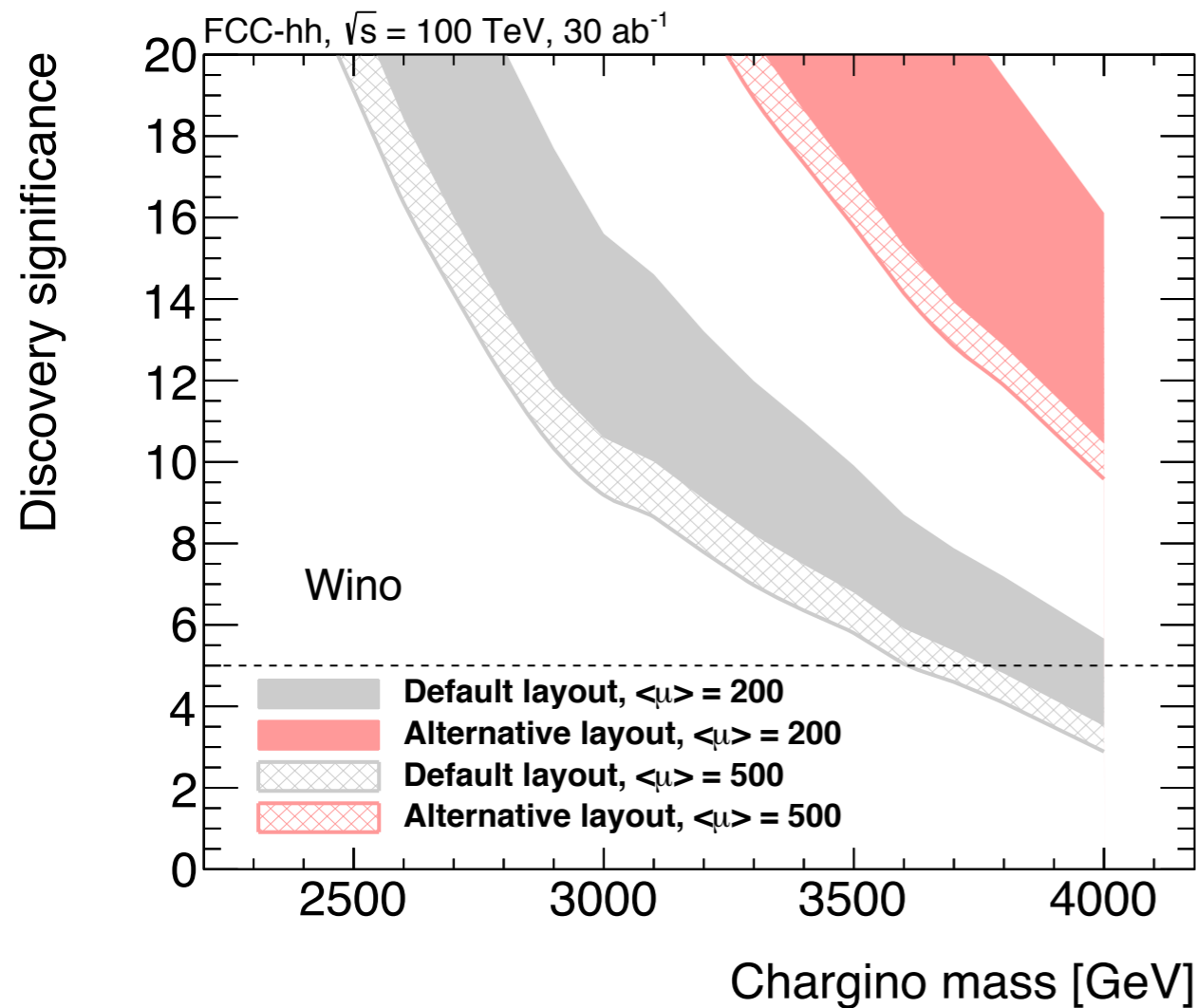
$$\Omega_{\text{DM}} h^2 \sim 0.12 \times \left(\frac{M_{\text{DM}}}{2 \text{TeV}} \right)^2 \left(\frac{0.3}{g_{\text{eff}}} \right)^4$$

$$\Omega_{\text{wimp}} h^2 \lesssim 0.12$$



$$M_{\text{wimp}} \lesssim 2 \text{TeV} \left(\frac{g}{0.3} \right)^2$$

Disappearing charged track analyses (at ~full pileup)



New detector performance studies

K. Terashi, R. Sawada, M. Saito, and S. Asai, *Search for WIMPs with disappearing track signatures at the FCC-hh*, (Oct, 2018) . <https://cds.cern.ch/record/2642474>.

**Prospects to discover/exclude WIMP DM:
coverage beyond the upper limit of the thermal
WIMP mass range for both higgsinos and winos !!**

$$M_{wimp} \lesssim 2 \text{ TeV} \left(\frac{g}{0.3} \right)^2 \Rightarrow m_{wino} < 3.5 \text{ TeV} , \quad m_{higgsino} < 1 \text{ TeV}$$

additional opportunities

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- Dedicated detectors for long-lifetime particles (like FASER, MATHUSLA)

Final remarks

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- Unique among the proposed projects for future colliders, the FCC builds on the tried and tested format forged by the LEP-LHC experience, integrating the well-established and complementary qualities of circular e^+e^- and pp colliders within a largely common, and partly existing, infrastructure.
- The sequence of FCC-ee and FCC-hh provides the most complete, detailed and accurate picture of Higgs properties achievable with the currently planned facilities, and gives direct access to the largest mass scales allowed by foreseeable technology