

**Recent results from LIGO and Virgo searches of** continuous gravitational waves Cristiano Palomba (INFN Roma) on behalf of the LIGO/Virgo/Kagra Collaborations

**Talk outline** 

Sources of continuous gravitational waves (CWs)

Searches for CWs

Some recent search results



Future prospects

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# Sources of Gravitational Waves (GW)

Modeled

**Un-modeled** 





Transient

sistent

### **CW** signals

Persistent, semi-periodic signals, with duration much larger than the typical observation times (months-years)

CW are emitted by spinning neutron stars, if asymmetric with respect to their rotation axis: constant amplitude, small spin-down

Asymmetry can be due to:
elastic stresses in the star body
inner magnetic field not aligned to the rotation axis
matter accretion from a companion star
excitation of long-lasting oscillations (like r-modes).

See e.g. review by Glampedakis & Gualtieri, https://arxiv.org/abs/1709.07049

 We KNOW that potential sources of CW exist: more than 3,000 NS are observed through their EM emission, up to 1 billion are expected to exist in the Galaxy

• We DO NOT KNOW the amplitude of the emitted signals

$$h_0 \cong 10^{-27} \left( \frac{I_{zz}}{10^{38} \text{kg} \cdot \text{m}^2} \right) \left( \frac{10 \text{kpc}}{d} \right) \left( \frac{f}{100 \text{Hz}} \right)^2 \left( \frac{\varepsilon}{10^{-6}} \right)$$

Expected signal maximum frequency below ~2kHz.

The source we are searching for are in the Galaxy, d<O(10kpc)</li>

• The ellipticity is largely unknown;  $\epsilon_{max} \sim 10^{-6}$  for standard NSs, but some exotic EOS foresee  $\epsilon_{max} \sim 10^{-4}$  or even more.

• What are the typical values?

# CW-like signals are also expected from newborn highly deformed neutron stars, like magnetars





# The very strong spin-down makes these signals last hours-days in the detector sensitivity band $\epsilon_{\rm B} = k \left( E_{\rm B,\phi}/W \right) \approx 4 \times 10^{-4} \left( k/4 \right) B_{\rm int,16}^2 R_6^4 M_{1,4}^{-2} \quad \text{ellipticity} \sim B^2$

Magnetars formed in core collapse supernovae have an estimated event rate of 1 per year within ~10 Mpc

CW-like signals can be also associated to Dark Matter (DM)

Three channels are being explored at the moment:

ultra-light boson "clouds" around spinning black holes (BHs)

direct interaction of ultra-light DM particles with ITF mirrors

sub-solar mass binary BHs

Ultra-light boson "clouds" should form around spinning BHs through superradiance, and then dissipate through GW emission [Arvanitaki+, PRD91, 084011 (2015)]

Signal mainly depends on the boson and BH mass

M<sub>RH</sub> [solar masses]



M<sub>RH</sub> [solar masses]

DM ultra-light candidates could directly interact with ITF components through different mechanisms

E.g. dark photons (DP) could couple to protons and/or neutrons of the mirrors producing a stochastic oscillatory force [Pierce+, PRD99, 075002 (2019)]



 $\langle h_{\rm total}^2 \rangle = \langle h_D^2 \rangle + \langle h_C^2 \rangle$ 

 e particle/DP coupling constant (normalized to EM coupling constant)

$$\sqrt{\langle h_C^2 \rangle} \simeq 6.58 \times 10^{-26} \left( \frac{\epsilon}{10^{-23}} \right)$$

Frequency spread related to the Maxwell-Boltzman velocity distribution of DPs

# Searches for CWs

Intrinsic signal frequency changes due to spin-down (up)

$$f_0(t) = f_0 + \dot{f}_0(t - t_0) + \frac{\ddot{f}_0}{2}(t - t_0)^2 + \dots$$

Received signal frequency is affected by Doppler modulation (with a further modulation if the source is in a binary system)

There is a sidereal phase and amplitude modulation due to the detector nonuniform response function







Small relativistic effects (namely, Einstein and Shapiro delay) may play a role

Pulsar glitches are likely to produce GW signal phase discontinuities

Signal frequency random fluctuations can also be present



In general, signals are expected to be deeply buried into the noise BUT

→ signal duration very long respect to typical observation times → Signal-to-noise ratio increases with time

→ signals have very specific time-frequency behavior → This helps also in rejecting noise artifacts Robustness with respect to un-modeled signal th would be tifacts oise 0 00 00

T<sub>cob</sub>: search coherence time

Based on matched filter.  $h_{\min} \approx [10 - 25] \sqrt{\frac{S_n(f)}{T_{obs}}}$ E.g. pulsars for which accurate Sensitivity ephemeris are available Allows for a small Targeted search mismatch between the observation time GW and the EM signal Hierarchical methods. Narrow-band Follow-up of the most search interesting candidates. **Computationally bound**  $h_{\min} \approx \frac{\Lambda}{N^{1/4}} \sqrt{\frac{S_n(f)}{T_{coh}}}$ Wide-band Based on semi-coherent methods. ('directed') E.g. SNR remnants search  $h_{\min} \approx 30 \sqrt{\frac{S_n(f)}{T}}$ All-sky search NS in binary Binary system orbital parameters must be systems taken into account. (e.g. Sco X-1) **Computationally bound** S<sub>n</sub>(f): detector noise power spectrum **Computational load** 

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### Some recent search results

Based on O3 data (several analyses still on-going)

See e.g. Sieniavska & Bejger, Universe 2019, 5, 217; doi:10.3390/universe5110217 for a (relatively) recent review of methods and past results



**Figure 1.** O3a noise PSD for H1, L1, and V1 shown in red, green, and purple. The H1 and L1 PSDs are calculated during a time period of optimal performance for the detector, while the Virgo PSD is averaged over the run. The vertical dashed lines indicate the searched frequency region for each of the five pulsars.

# ApJL 902 L21, 2020 O3a covers April 1<sup>st</sup> –October 1<sup>st</sup> 2019

Full O3 targeted and narrow-band searches on-going GWs constraints on the equatorial ellipticity of millisecond pulsars Targets: 5 radio pulsars, O3a data For the first time, a constraint on the fraction of spin-down energy due to GWs emission has been obtained for a millisecond pulsar

Pulsar name (J2000)	$h_0^{\rm sd}$ (10 <sup>-26</sup> )	Analysis method	$C_{21}^{95\%}$ (10 <sup>-26</sup> )	$C_{22}^{95\%}$ (10 <sup>-27</sup> )	$h_0^{95\%}$ (10 <sup>-26</sup> )	$Q_{22}^{95\%}$ (10 <sup>32</sup> kg m <sup>2</sup> )	$\epsilon^{95\%}$	$h_0^{95\%}/h_0^{\rm sd}$
			Y	oung pulsars <sup>a</sup>		- 200		<u>2</u>
J0534+2200 <sup>b</sup> (Crab)	140	Bayesian	12.7(7.9)	6.3(5.6)	1.5(1.2)	6.6(5.7)	$8.6(7.4) \times 10^{-6}$	0.010(0.009)
		$\mathcal{F}/\mathcal{G}$ -statistic	8.9(6.2)	7.9(7.1)	1.9(1.5)	7.9(6.3)	$10(8.1) \times 10^{-6}$	0.014(0.011)
		5n-vector	15.9(12.4)		3.0(2.9)	12.6(12.1)	$16.3(15.7) \times 10^{-6}$	0.021(0.021)
J0835–4510 (Vela)	330	Bayesian	1100(980)	120(84)	22(17)	91(73)	$12.0(9.5) \times 10^{-5}$	0.067(0.052)
		$\mathcal{F}/\mathcal{G}$ -statistic	1470(1370)	116(48)	23(12)	96(50)	$12.4(6.4) \times 10^{-5}$	0.070(0.036)
		5n-vector	1700(1400)		24(24)	100(102)	$13.0(13.2) \times 10^{-5}$	0.073(0.073)
			R	ecycled pulsars	S			5
J0437-4715	0.79	Bayesian	2.2	4.1	0.78	0.0074	$9.5 \times 10^{-9}$	0.99
		$\mathcal{F}/\mathcal{G}$ -statistic	2.1	7.2	0.86	0.0082	$11.0 \times 10^{-9}$	1.1
		5n-vector	•••					•••
J0711—6830	1.2	Bayesian	2.6	3.5	0.82	0.0064	$8.3 \times 10^{-9}$	0.68
		$\mathcal{F}/\mathcal{G}$ -statistic	2.4	9.4	0.98	0.0059	$7.7 \times 10^{-9}$	0.82
		5n-vector	2.9	••••	0.91	0.0053	$7.2 \times 10^{-9}$	0.76
J0737-3039A	0.62	Bayesian	5.9	3.3	0.69	0.80	$1.0 \times 10^{-6}$	1.1
		$\mathcal{F}/\mathcal{G}$ -statistic	3.0	1.2	0.99	1.10	$1.4 \times 10^{-6}$ 1	.3 1.6
		5n-vector				1111	100	100

### All-sky search for CWs from unknown NSs in binary systems

#### PRD 103 064017, 2021

O3a data

Frequency range: [50, 300] Hz

Orbital period range of [3, 45] days and projected semi-major axis of [2, 40] I-s

**GPU-accelerated** pipeline



#### Constraints from the energetic young X-ray pulsar PSR J0537-6910

ApJL 913 L27, 2021

PSR 0537-6910 is also known as the "big glitcher"

Ephemerides from NICER data

Full O3 LIGO/Virgo data

Search at both once and twice f<sub>rot</sub>=62 Hz

For the first time for this pulsar, we have been able to constrain the fraction of spin-down energy due to GW (less than 14%)



Figure 4. Posterior probability distribution for ellipticity and  $h_0$  for the analyses with unrestricted and restricted priors on the pulsar orientation. The 95% credible upper limits are shown as vertical colored lines, while the spin-down limit is given by the vertical dashed black line. 15

# Constraints on GW emission due to r-modes in the glitching pulsar PSR J0537-6910

#### Accepted in ApJ (2104.144170)



Frequency range: [86, 97] Hz to deal with EOS uncertainty

Ephemeris obtained. from NICER data Full O3 search, motivated by the measured interglitch braking index which suggests that r-modes could be active  $\alpha = \sqrt{\frac{5}{8\pi}} \frac{c^5}{G} \frac{h_0}{(2\pi f)^3} \frac{d}{MR^3 \tilde{J}}$ 





 $\approx 0.017 \left(\frac{90 \text{ Hz}}{f}\right)^3 \left(\frac{h_0}{10^{-26}}\right)^3$ 

#### CW search from young supernova remnants

Submitted to ApJ (2105.11641)

15 SNR, O3a data, 3 algorithms

#### Frequency range: [10, 2000] Hz

$$h_0^{
m age} = 2.27 imes 10^{-24} \left(rac{1 \, 
m kpc}{D}
ight) \left(rac{1 \, 
m kyr}{t_{
m age}}
ight)$$





Source	Minimum $t_{age}$ (kyr)	D (kpc)	$T_{\rm coh}$ (hours)	f (Hz)	$\dot{f}~({ m Hz/s})$
G1.9 + 0.3	0.10	8.5	1.0	[31.56, 121.7]	$\left[-3.858 \times 10^{-8}, 3.858 \times 10^{-8}\right]$
$G15.9 {+} 0.2$	0.54	8.5	1.0	[44.03, 657.1]	$\left[-3.858\times10^{-8}, 3.858\times10^{-8}\right]$
G18.9–1.1	4.4	2	1.9	[31.02, 1511]	$\left[-1.507\times10^{-8}, 1.507\times10^{-8}\right]$
G39.2 - 0.3	3.0	6.2	2.8	[62.02, 459.2]	$\left[-1.968  imes 10^{-8}, 1.968  imes 10^{-8} ight]$
G65.7 + 1.2	20	1.5	4.7	[35.10, 1128]	$\left[-3.149\times10^{-9}, 3.149\times10^{-9}\right]$
G93.3 + 6.9	5.0	1.7	1.9	$\left[ 30.00, 1668 \right]$	$\left[-1.335  imes 10^{-8}, 1.335  imes 10^{-8} ight]$
G111.7 - 2.1	0.30	3.3	1.0	$\left[25.71, 365.1 ight]$	$\left[-3.858  imes 10^{-8}, 3.858  imes 10^{-8} ight]$
G189.1+3.0	3.0	1.5	1.4	$\left[26.13,2000\right]$	$\left[-1.968\times10^{-8}, 1.968\times10^{-8}\right]$
G266.2 - 1.2	0.69	0.2	1.0	[18.36, 839.6]	$\left[-3.858 \times 10^{-8}, 3.858 \times 10^{-8}\right]$
G291.0-0.1	1.2	3.5	1.0	[31.97, 1460]	$\left[-3.858 imes 10^{-8}, 3.858 imes 10^{-8} ight]$
G330.2 + 1.0	1.0	5	1.1	[36.57, 1039]	$\left[-3.858\times10^{-8}, 3.858\times10^{-8}\right]$
G347.3 - 0.5	1.6	0.9	1.0	[21.74, 1947]	$\left[-3.858\times10^{-8}, 3.858\times10^{-8}\right]$
G350.1 - 0.3	0.60	4.5	1.0	[31.96, 730.1]	$\left[-3.858\times10^{-8}, 3.858\times10^{-8}\right]$
G353.6-0.7	27	3.2	10	$\left[77.86, 318.3 ight]$	$\left[-2.295 \times 10^{-9}, 2.295 \times 10^{-9} ight]$
G354.4 + 0.0	0.10	5	1.0	[25.72, 121.7]	$\left[-3.858 imes10^{-8}, 3.858 imes10^{-8} ight.$

 $h_{0,min} \sim 7.7 \ 10^{-26} \ for \ G65.7+1.2$   $\epsilon_{min} \sim 6 \ 10^{-8} \ for \ G266.2-1.2/Vela \ Jr.$ Search sensitive to a possible spin wandering

#### Constraints on dark photon dark matter using data from LIGO's and Virgo's third observing run

PRL, submitted (2105.1385)

Full O3 search, 2 pipelines

Explored DP mass range:  $10^{-14} - 10^{-11}$  eV/c<sup>2</sup>

Improvement of O(100) w.r.t. dark matter direct detection experiments in the range  $[2-4] 10^{-13} \text{ eV/c}^2$ 



**Multi-messenger approach** EM observations are crucial for: - more sensitive searches of known sources - updated ephemeris restricted parameter space - searches of new potentially promising sources possible confirmation/identification of all-sky search candidates estimate distance to the source (which would allow to extract physical information from the measured strain)

The connection among GWs searches and DM will likely strengthen

#### **Future prospects**

#### We are currently completing O3 analyses

- Three crucial ingredients to increase chance of detecting CWs are:
- More sensitive, robust and fast searches

**New sources** 

More sensitive detectors  $2G \rightarrow 2.5G \rightarrow 3G$  (Einstein Telescope Cosmic Explorer)



https://dcc.ligo.org/public/0094/P1200087/058/ObservingScenarios.pdf

# Thank you for the attention!