

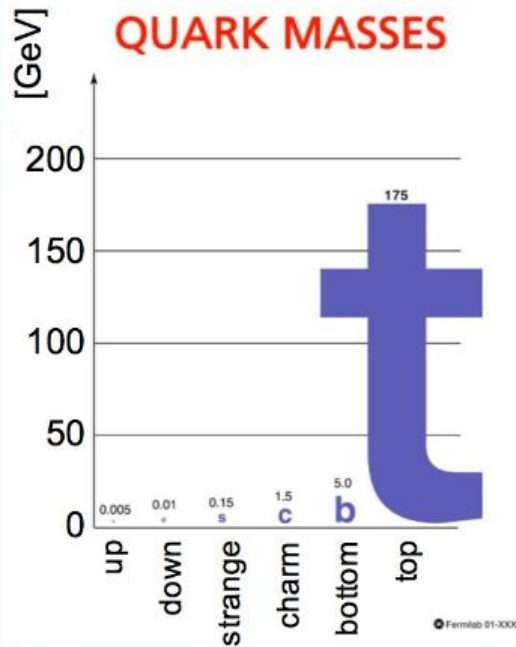


Overview on the top quark: between SM and BSM

COMPOSE-IT 28/01/2020

A.O.M. Iorio

The top quark: why so important?

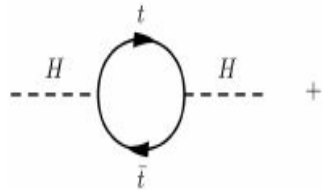


Heaviest particle known: top quark mass appears in most loops and corrections to SM and BSM

- **Strongest coupling with higgs:**

- possibly related to EWK Symmetry Breaking
- Leading contribution to Higgs mass corrections

$$m_f = \frac{\lambda_f v}{\sqrt{2}}$$



- **Only quark to decay before hadronisation**

- Properties can be directly measured

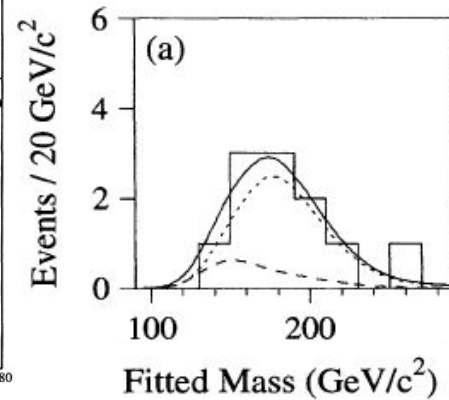
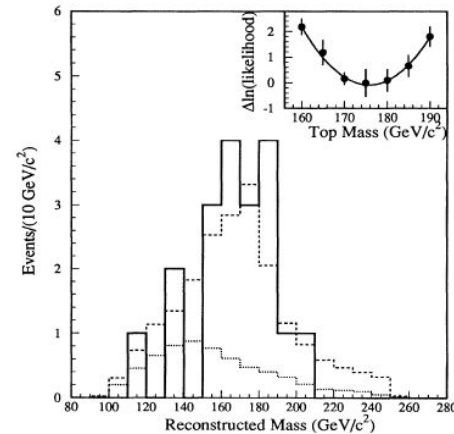
- **For new physics searches:**

- Top BSM partners usually have the lowest mass
- Least explored couplings

On top of all quarks since 25 years now!

We have gone a long way since first observation at TeVatron in '95!

→ Observation of pair ('95) and single (2009!) production



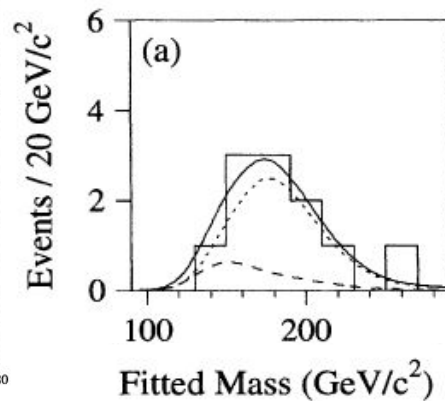
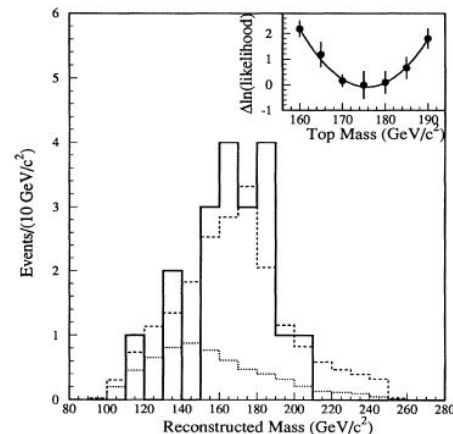
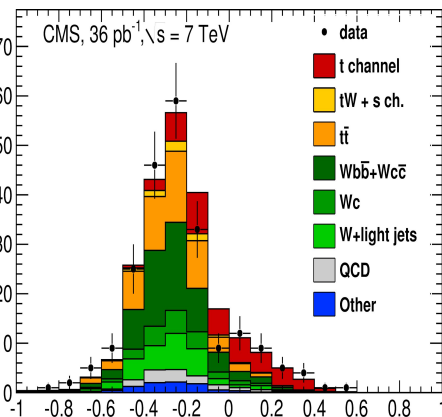
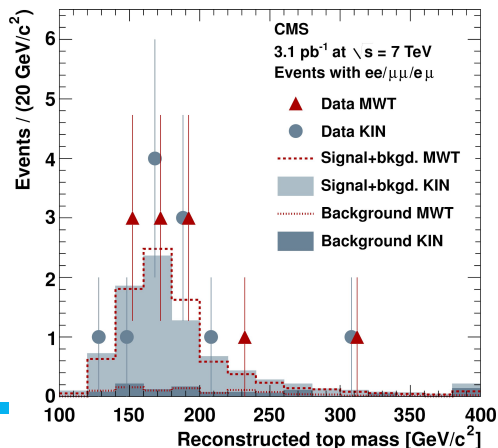
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What now? What does it mean living the top life nowadays?



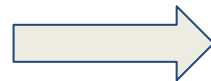
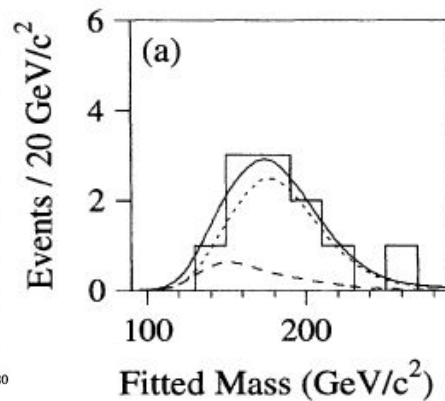
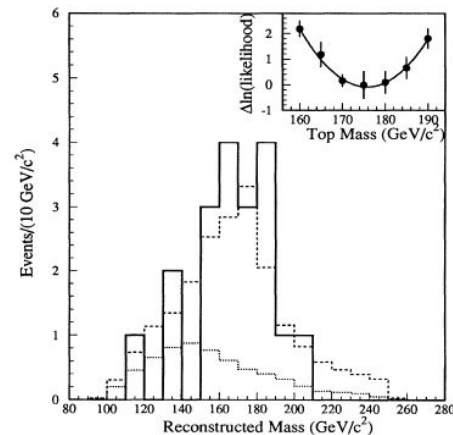
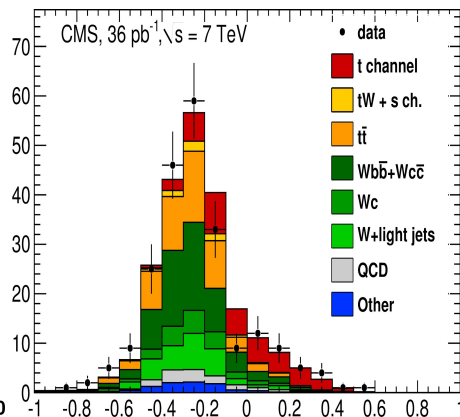
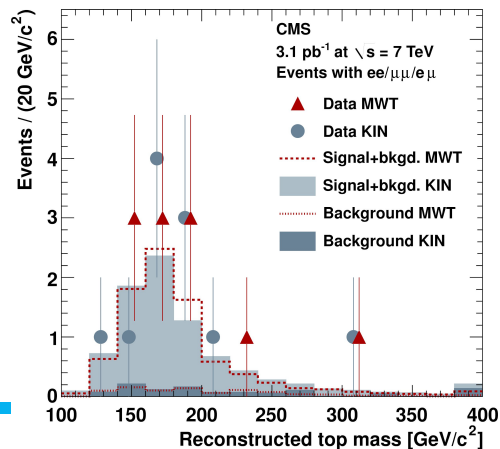
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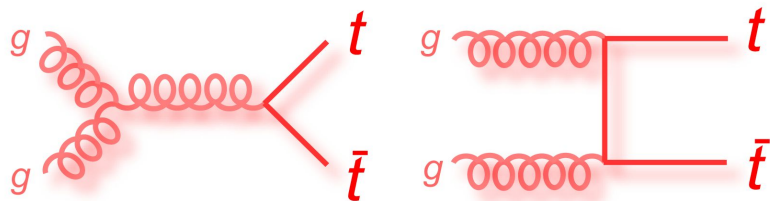


What now?

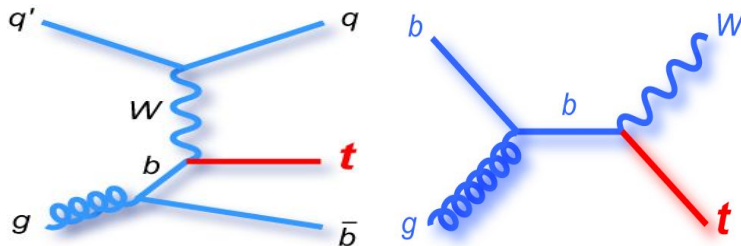
The “classics”: pair and single production

A great laboratory for top@ high precision:

- **Coupling properties:** inclusive and differential xSection
- Top quark-related **parameters:** mass, width, CKM elements etc.
- **Modeling of QCD** (perturbative and non-pert.) and **PDF**



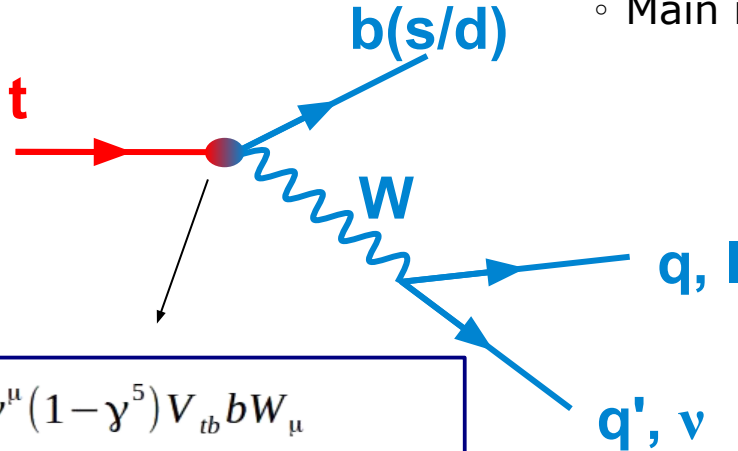
- **tt pairs** via strong interaction:
 - dominant at the LHC and Tevatron
 - depends on α_s
 - sensitive to pdf
- **σ at LO $\propto (\alpha_s/m_{\text{top}})^2$**



- **single-top quarks:**
 - weak charged current interactions
 - t-, s-channel and W-associated
 - tWb vertex in production
 - Sensitive to V_{tb}
- **σ at LO $\propto (\alpha \cdot |V_{tb}|)^2$**

Top-quark decays vademecum

- Main mechanism is **electroweak: no hadronisation**



W decay:

Branching Ratio

$W \rightarrow lv \text{ (any)}$	0.32
$W \rightarrow qq' \text{ (any)}$	0.68

$$\frac{-ig}{2\sqrt{2}} \bar{t} \gamma^\mu (1 - \gamma^5) V_{tb} b W_\mu$$

- Electroweak **tWb vertex:**

- V-A: **polarization** of the products and defined **W-helicity**
- CKM matrix element $|V_{tb}| \sim 1 \rightarrow \Gamma(t \rightarrow b) \gg \Gamma(t \rightarrow s, d)$

$1/m_{toq}$
production

<

$1/\Gamma_t$
lifetime

<

$1/\Lambda$
hadronization

<

m_t/Λ^2
spin decorrelation

What's the kind of measurements we do with it?

... our results are usually a mix of three:

X: High precision

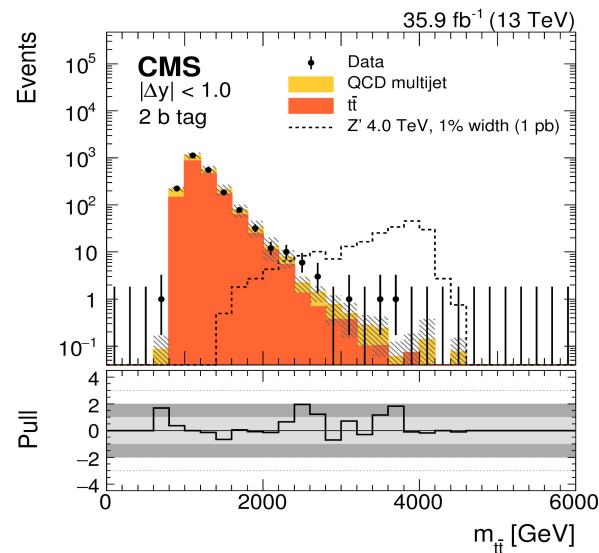
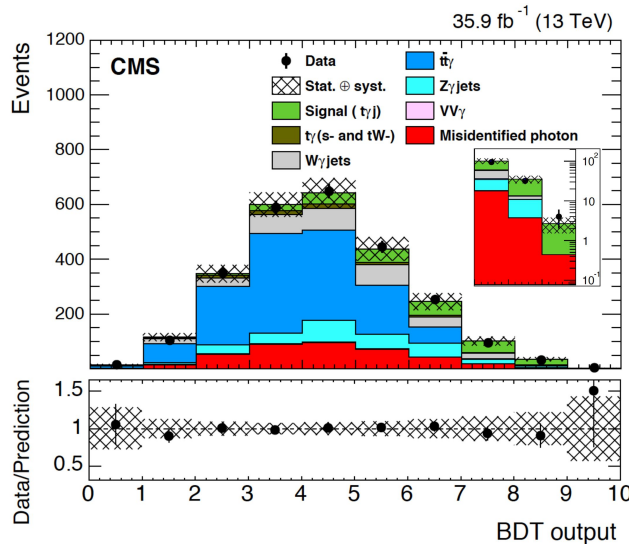
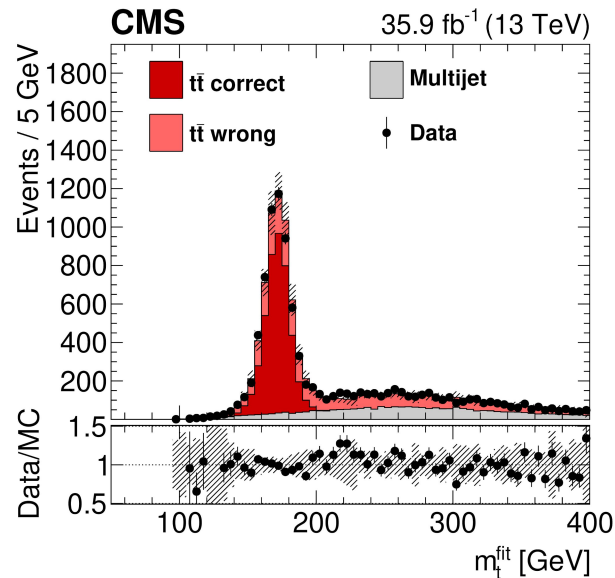
→ SM Measurements!

Y: Rare SM processes

→ Exploration!

Z: New physics signals

→ Searches!



→ Often as time goes one turns into the others

A further binding thread: EFT in top quarks

Effective lagrangians: taylor expansion of the unknown new physics lagrangian, truncated at $X=\dots 6$

$$L = L_{SM}^{(4)} + \sum_i \frac{c_i^{(5)}}{\Lambda_i} O_i^{(5)} + \sum_i \frac{c_i^{(6)}}{\Lambda_i^2} O_i^{(6)} + \dots$$

→ define a reduced number of operators for top quarks

→ See also <https://arxiv.org/pdf/1802.07237.pdf>

Operators

Wilson coefficients

A further binding thread: EFT in top quarks

Effective lagrangians: taylor expansion of the unknown new physics lagrangian, truncated at X=...6

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→ define a reduced number of operators for top quarks

→ See also <https://arxiv.org/pdf/1802.07237.pdf>

Two quark operators:

$$\begin{aligned} \dagger O_{u\varphi}^{(ij)} &= \bar{q}_i u_j \tilde{\varphi} (\varphi^\dagger \varphi), \\ O_{\varphi q}^{1(ij)} &= (\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{q}_i \gamma^\mu q_j), \\ O_{\varphi q}^{3(ij)} &= (\varphi^\dagger \overleftrightarrow{D}_\mu^I \varphi) (\bar{q}_i \gamma^\mu \tau^I q_j), \\ O_{\varphi u}^{(ij)} &= (\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{u}_i \gamma^\mu u_j), \\ \dagger O_{\varphi ud}^{(ij)} &= (\tilde{\varphi}^\dagger i D_\mu \varphi) (\bar{u}_i \gamma^\mu d_j), \\ \dagger O_{uW}^{(ij)} &= (\bar{q}_i \sigma^{\mu\nu} \tau^I u_j) \tilde{\varphi} W_{\mu\nu}^I, \\ \dagger O_{dW}^{(ij)} &= (\bar{q}_i \sigma^{\mu\nu} \tau^I d_j) \varphi W_{\mu\nu}^I, \\ \dagger O_{uB}^{(ij)} &= (\bar{q}_i \sigma^{\mu\nu} u_j) \tilde{\varphi} B_{\mu\nu}, \\ \dagger O_{uG}^{(ij)} &= (\bar{q}_i \sigma^{\mu\nu} T^A u_j) \tilde{\varphi} G_{\mu\nu}^A, \end{aligned}$$

Four quark operators:

$$\begin{aligned} O_{qq}^{1(ijkl)} &= (\bar{q}_i \gamma^\mu q_j) (\bar{q}_k \gamma_\mu q_l), \\ O_{qq}^{3(ijkl)} &= (\bar{q}_i \gamma^\mu \tau^I q_j) (\bar{q}_k \gamma_\mu \tau^I q_l), \\ O_{qu}^{1(ijkl)} &= (\bar{q}_i \gamma^\mu q_j) (\bar{u}_k \gamma_\mu u_l), \\ O_{qu}^{8(ijkl)} &= (\bar{q}_i \gamma^\mu T^A q_j) (\bar{u}_k \gamma_\mu T^A u_l), \\ O_{qd}^{1(ijkl)} &= (\bar{q}_i \gamma^\mu q_j) (\bar{d}_k \gamma_\mu d_l), \\ O_{qd}^{8(ijkl)} &= (\bar{q}_i \gamma^\mu T^A q_j) (\bar{d}_k \gamma_\mu T^A d_l), \\ O_{uu}^{(ijkl)} &= (\bar{u}_i \gamma^\mu u_j) (\bar{u}_k \gamma_\mu u_l), \\ O_{ud}^{1(ijkl)} &= (\bar{u}_i \gamma^\mu u_j) (\bar{d}_k \gamma_\mu d_l), \\ O_{ud}^{8(ijkl)} &= (\bar{u}_i \gamma^\mu T^A u_j) (\bar{d}_k \gamma_\mu T^A d_l), \\ \dagger O_{quqd}^{1(ijkl)} &= (\bar{q}_i u_j) \varepsilon (\bar{q}_k d_l), \\ \dagger O_{quqd}^{8(ijkl)} &= (\bar{q}_i T^A u_j) \varepsilon (\bar{q}_k T^A d_l), \end{aligned}$$

Two quark, two leptons operators:

$$\begin{aligned} O_{lq}^{1(ijkl)} &= (\bar{l}_i \gamma^\mu l_j) (\bar{q}_k \gamma_\mu q_l), \\ O_{lq}^{3(ijkl)} &= (\bar{l}_i \gamma^\mu \tau^I l_j) (\bar{q}_k \gamma_\mu \tau^I q_l), \\ O_{lu}^{(ijkl)} &= (\bar{l}_i \gamma^\mu l_j) (\bar{u}_k \gamma_\mu u_l), \\ O_{eq}^{(ijkl)} &= (\bar{e}_i \gamma^\mu e_j) (\bar{q}_k \gamma_\mu q_l), \\ O_{eu}^{(ijkl)} &= (\bar{e}_i \gamma^\mu e_j) (\bar{u}_k \gamma_\mu u_l), \\ \dagger O_{lequ}^{1(ijkl)} &= (\bar{l}_i e_j) \varepsilon (\bar{q}_k u_l), \\ \dagger O_{lequ}^{3(ijkl)} &= (\bar{l}_i \sigma^{\mu\nu} e_j) \varepsilon (\bar{q}_k \sigma_{\mu\nu} u_l), \\ \dagger O_{ledq}^{(ijkl)} &= (\bar{l}_i e_j) (\bar{d}_k q_l), \end{aligned}$$

Some EFT example contributions:

$$O_{\varphi Q}^{(3)} = i\frac{1}{2}y_t^2 \left(\varphi^\dagger \overleftrightarrow{D}_\mu^I \varphi \right) (\bar{Q}\gamma^\mu \tau^I Q)$$

$$O_{\varphi Q}^{(1)} = i\frac{1}{2}y_t^2 \left(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi \right) (\bar{Q}\gamma^\mu Q)$$

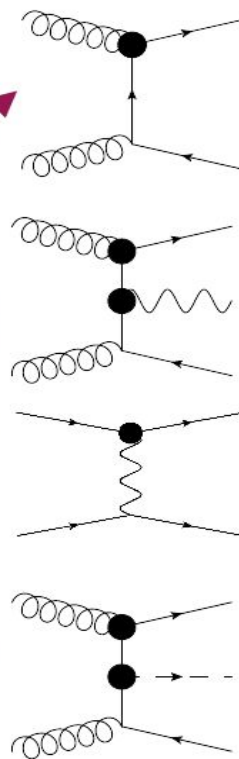
$$O_{\varphi t} = i\frac{1}{2}y_t^2 \left(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi \right) (\bar{t}\gamma^\mu t)$$

$$O_{tW} = y_t g_w (\bar{Q}\sigma^{\mu\nu} \tau^I t) \tilde{\varphi} W_{\mu\nu}^I$$

$$O_{tB} = y_t g_Y (\bar{Q}\sigma^{\mu\nu} t) \tilde{\varphi} B_{\mu\nu}$$

$$O_{tG} = y_t g_s (\bar{Q}\sigma^{\mu\nu} T^A t) \tilde{\varphi} G_{\mu\nu}^A$$

$$O_{t\phi} = y_t^3 (\phi^\dagger \phi) (\bar{Q}t) \tilde{\phi}$$



top quark pair

top pair + Z

single top quark

top pair + scalar

Cartoon taken from [this talk](#) of Eleni Vryonidou

EFT in top quarks in practice

In top quark sector means:

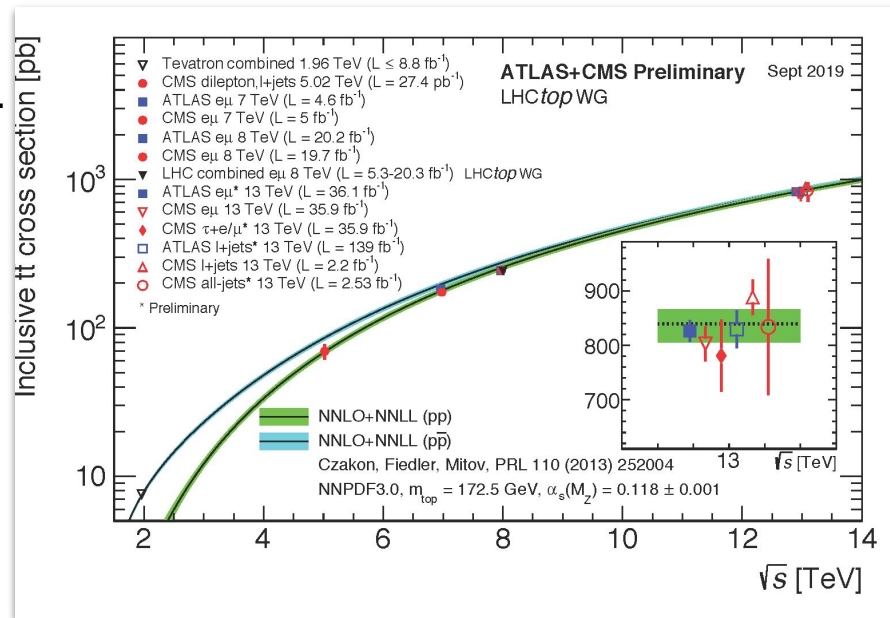
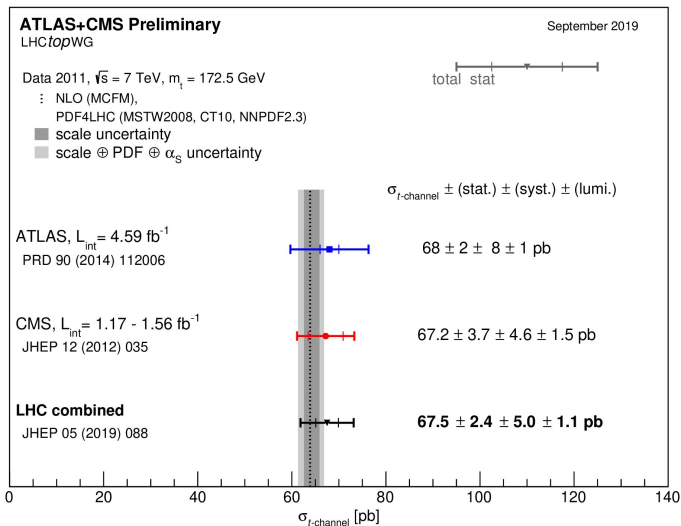
- Re-interpret existing unfolded measurements.
 - Design analyses to be explicitly sensitive EFT
- Effort in the [LHCTopWG](#) across Atlas and CMS to define common benchmarks and language
- Effort also across communities (W/Z measurements, top quark, Higgs)

In the following we will highlight: high impact measurements that have possible EFT implications / are designed to do so.

Top quark cross sections

Event counting - inclusive or differential:

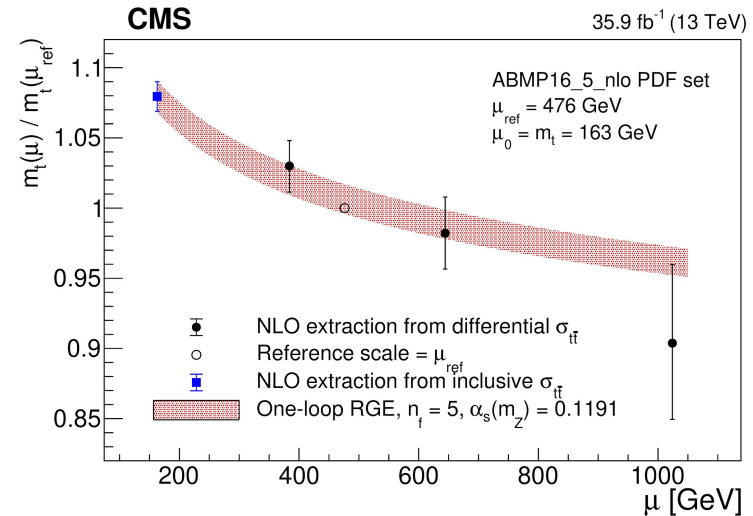
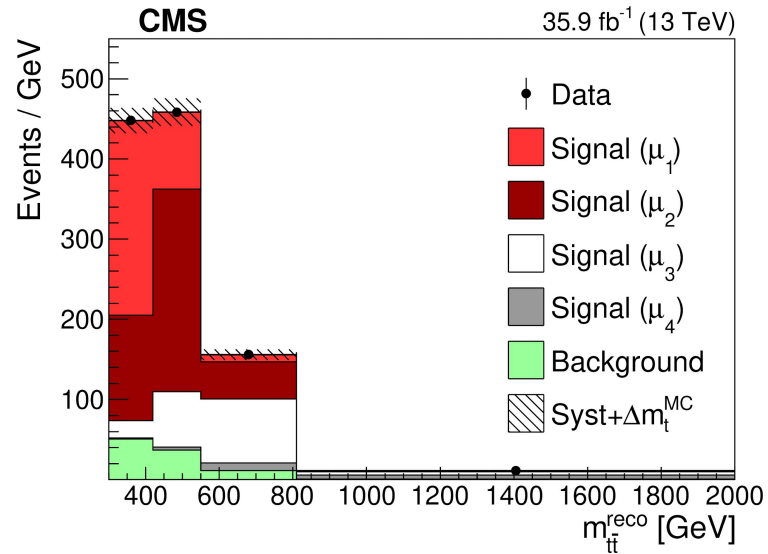
- Allows to access production properties, e.g. m_t , V_{tb} , α_S , etc.
- Inclusive cross section: first measurements, usually soon becoming systematics limited
- Allow to easily check the \sqrt{s} dependency across years



Differential spectra vs SM: top quark mass

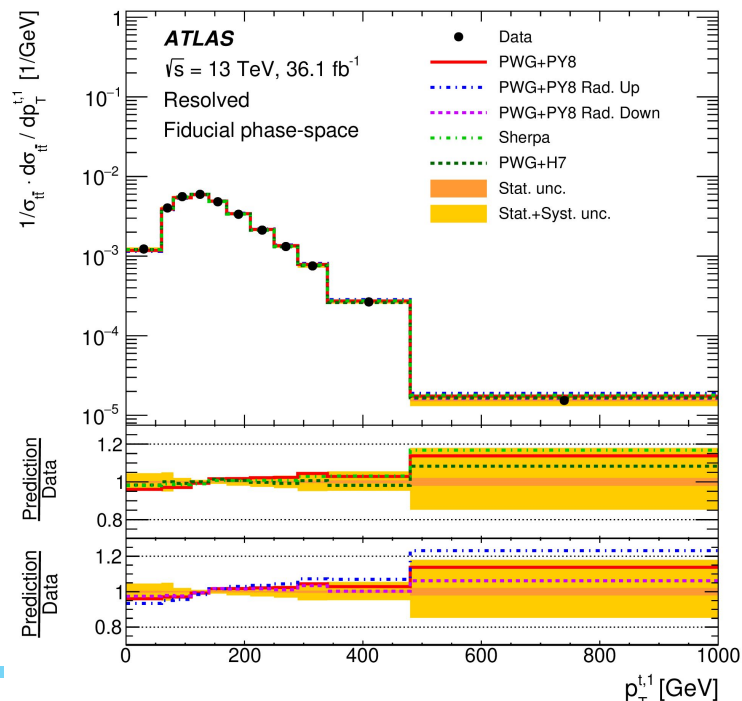
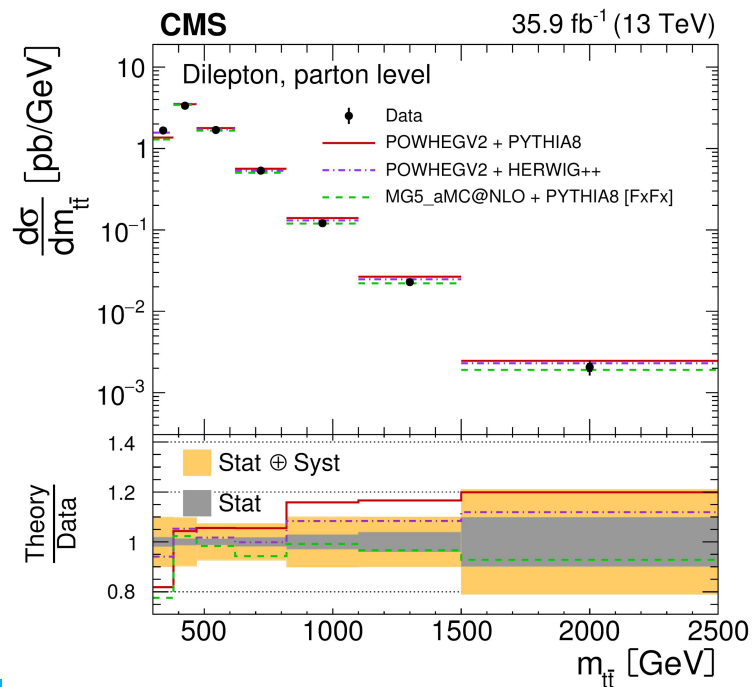
- **Running top quark mass**

- Dileptonic event selection: opposite-sign muon/electron
- Cross section in bins of reconstructed top quark mass
- SUT signal split as function of generator-level mass



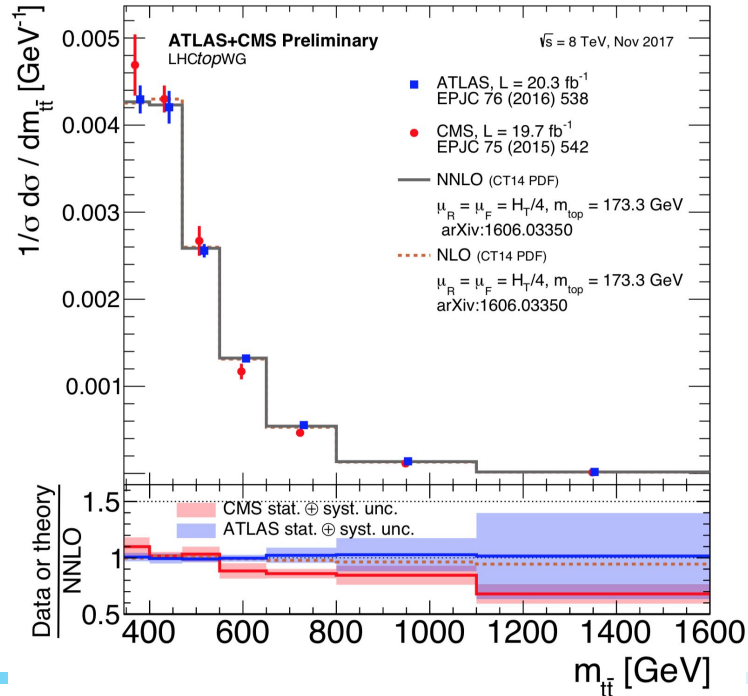
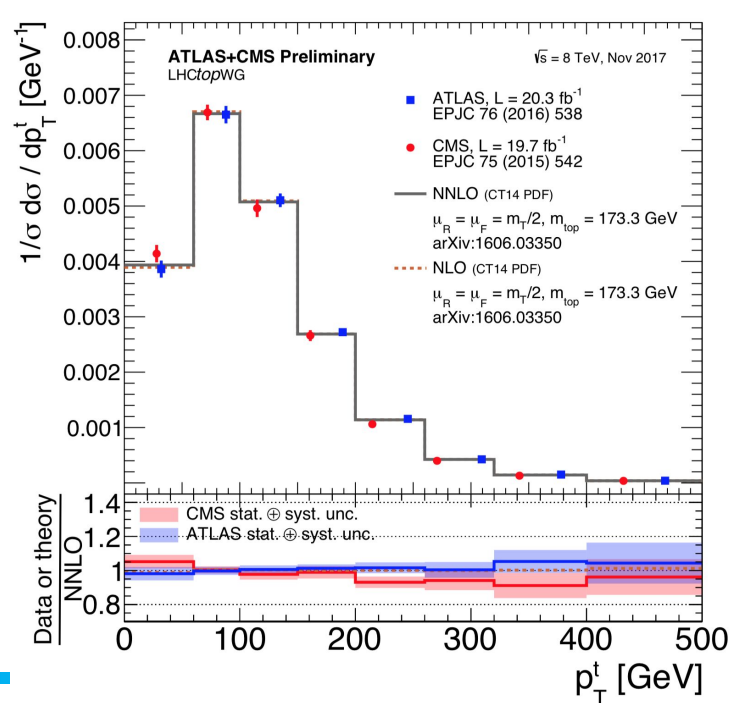
mtt spectrum vs QCD modelling

- **Pt spectra of the top quark pairs:** discrepancy coming from the transverse momentum spectra of the top quarks: prediction harder spectrum than data! Still under investigation where this effect comes from.
→ Work ongoing with full Run-II !



An “historical” view on differential mtt

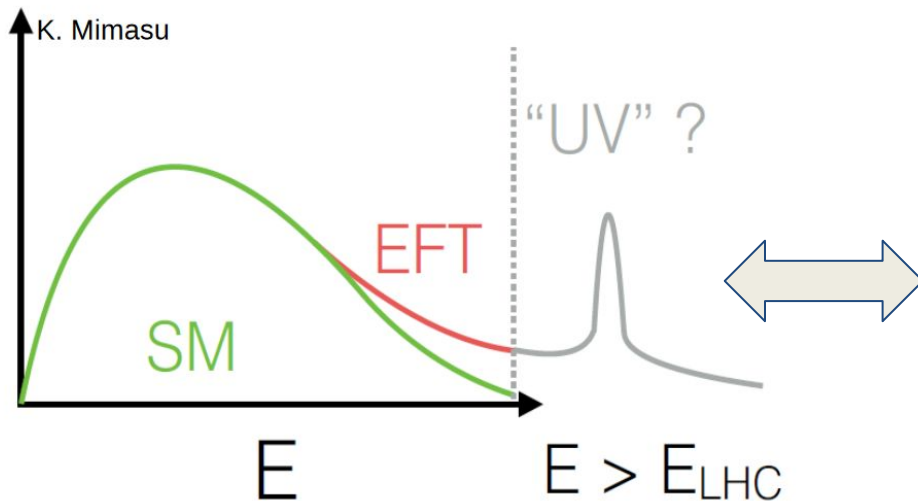
Discrepancies seen since long time: in Run-I data at 7-8 TeV discrepancy was already present
 → Not unambiguous interpretation.



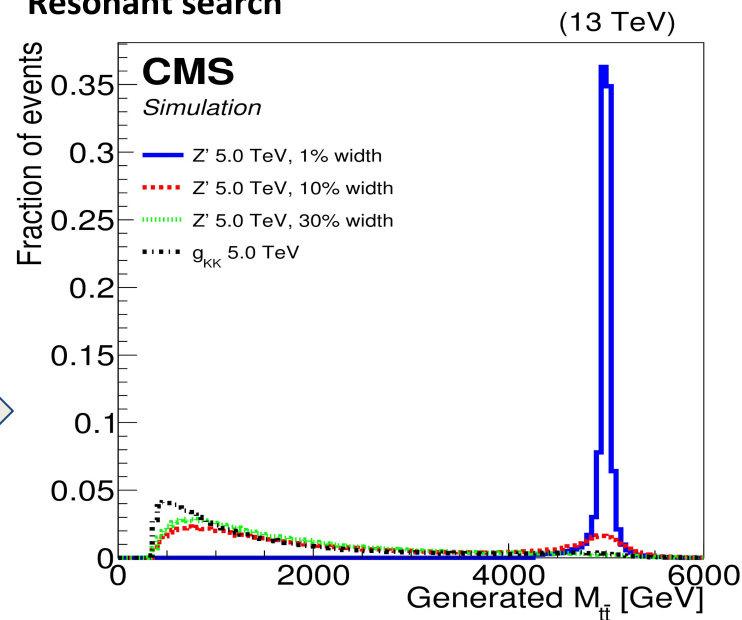
mtt beyond the SM, from EFT to resonances

Differential spectra at high-energy become key for **small anomaly detection** or traditional **bump-hunting**

EFT - type analysis



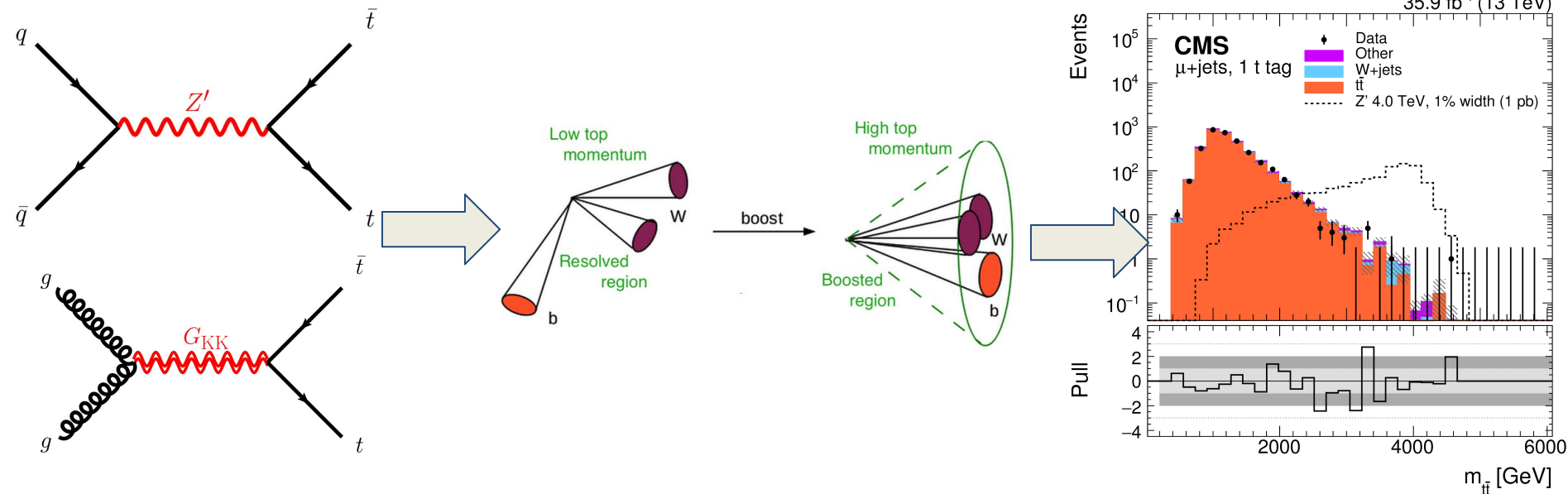
Resonant search



Nota bene: high-end tails will be extremely important for EFT analyses!

At high mtt: new physics searches

Search for a resonance \rightarrow top quark pairs: boosted final states!

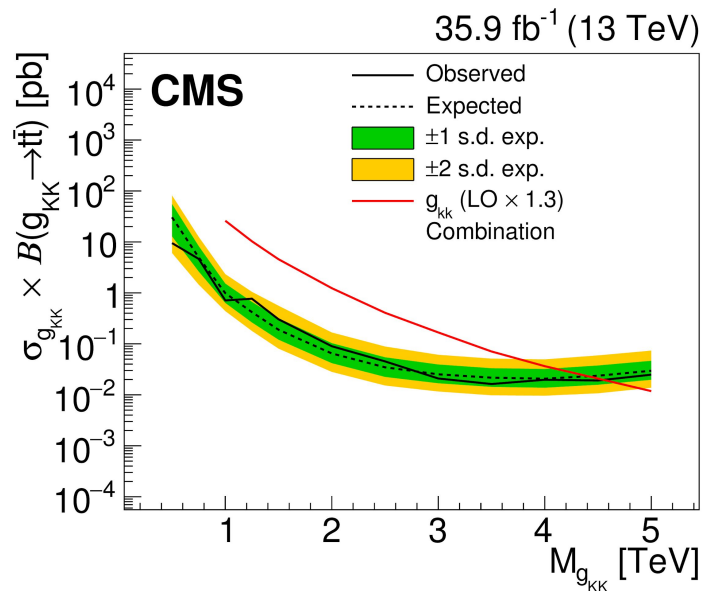
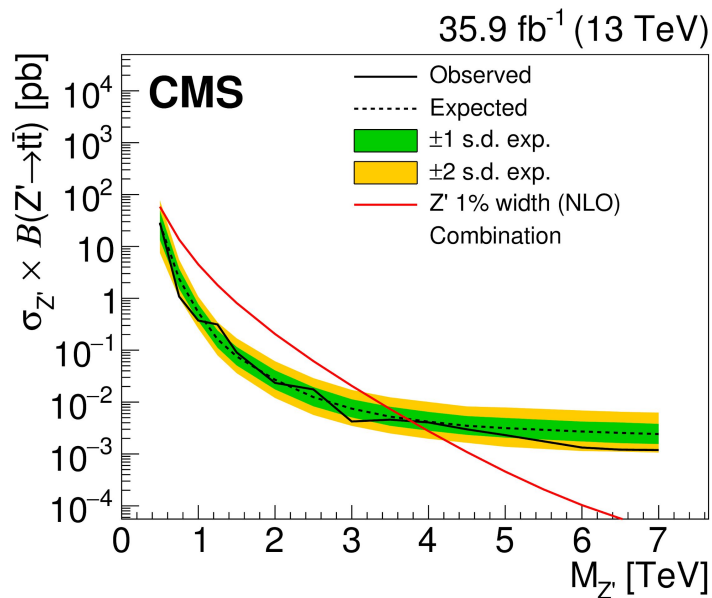


Boosted technologies: allow for probes of highest masses collecting top decay products in one fat jet

At high mtt: results and going forward

Limits go up: the mass shape does degenerate!

drawing from low x of the Parton Density Function \rightarrow searches will become “EFT-like”



Probing couplings in tt: angular properties

Angular properties for ttbar pairs :

1) Define appropriate basis for tt events:

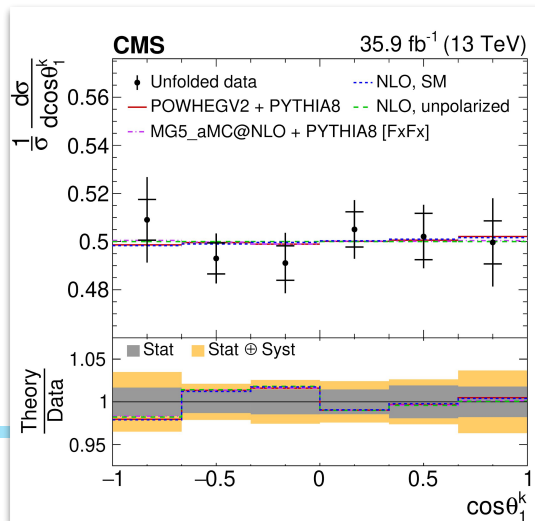
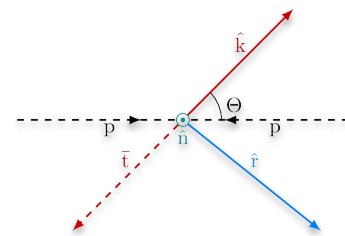
\mathbf{k} = top direction; \mathbf{n} = normal to scattering plane; \mathbf{r} orthogonal to other two

2) Decompose top polarisation $\mathbf{B}_1, \mathbf{B}_2$ and correlation matrix C in new basis

3) Define decay products: their distribution depends on top spin axis:

$$\frac{1}{\sigma} \frac{d^4\sigma}{d\Omega_1 d\Omega_2} = \frac{1}{(4\pi)^2} \left(1 + \mathbf{B}_1 \cdot \hat{\ell}_1 + \mathbf{B}_2 \cdot \hat{\ell}_2 - \hat{\ell}_1 \cdot C \cdot \hat{\ell}_2 \right)$$

4) One can write the products in 3) as angles in frame 1) and measure them

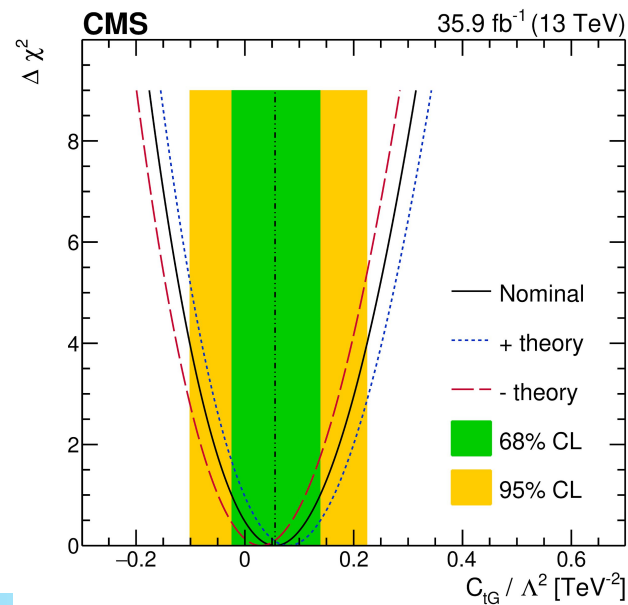
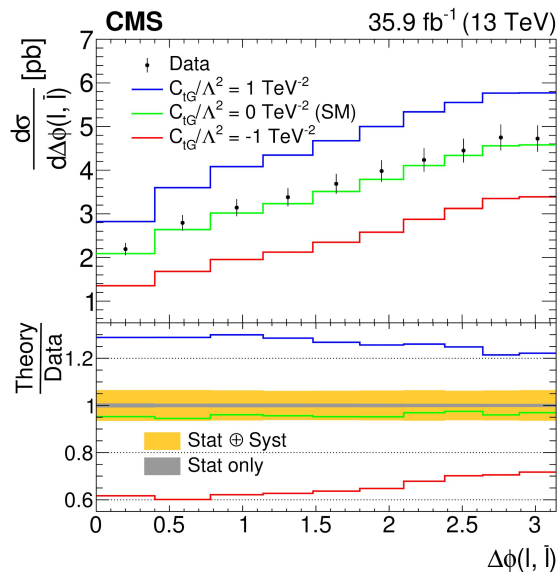
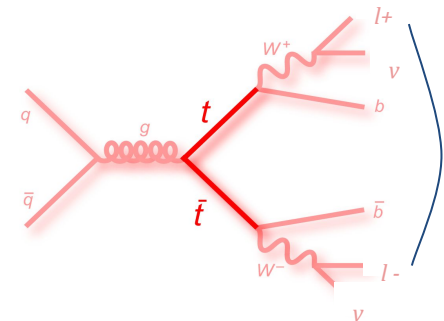


Probing EFT in tt: direct measurement

Angular properties: operators affecting the spin of tt, e.g. chromomagnetic dipole momentum

$$O_{tG} = y_t g_s (\bar{Q} \sigma^{\mu\nu} T^A t) \tilde{\varphi} G_{\mu\nu}^A$$

Angle between leptons: sensitive to operator O_{tG}



Single top quark cross section

All single top quarks: cross section depends on $|V_{tb}|^2$

→ sensitive to models that deviate from CKM = unitary

→ Sensitive to enhancements to V_{td}/V_{ts} from BSM.

Combination at LHC for Run-I:

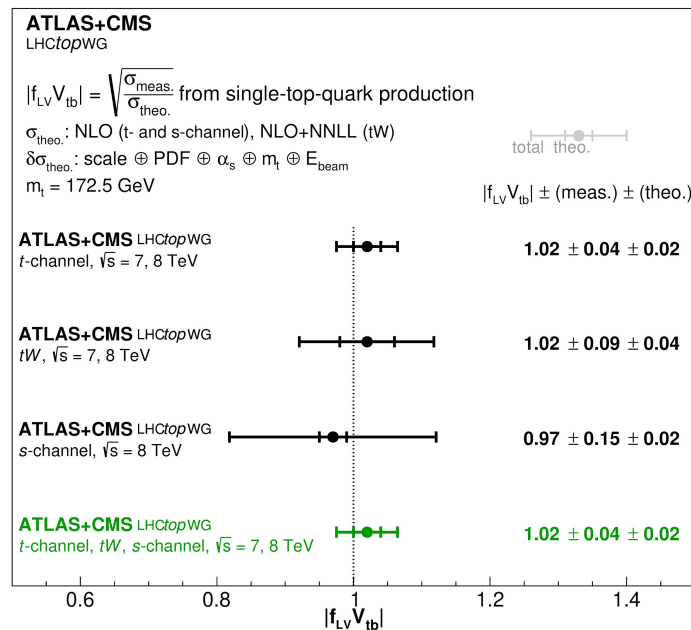
[10.1007/JHEP05\(2019\)088](https://arxiv.org/abs/1007/JHEP05(2019)088)

New measurement at 13 TeV under different assumptions: $|V_{tb}| = 1.00 \pm 0.03$

[TOP-17-012 Web page](#) / [CDS record](#)

Note: this measurement is also constraint on left-handed type operators:

$$\mathcal{L}_{Wtb}^{\text{dim-6}} = -\frac{g}{\sqrt{2}} \bar{b}(x) \gamma^\mu P_L t(x) W_\mu(x) \left(1 + \frac{C_{\varphi Q}^{(3)} y_t^2 v^2}{2\Lambda^2} \right) + \frac{2g v y_t C_{tW}}{\Lambda^2} \bar{b}(x) \sigma^{\mu\nu} P_R t(x) \partial_\nu W_\mu(x) + \text{h. c.},$$



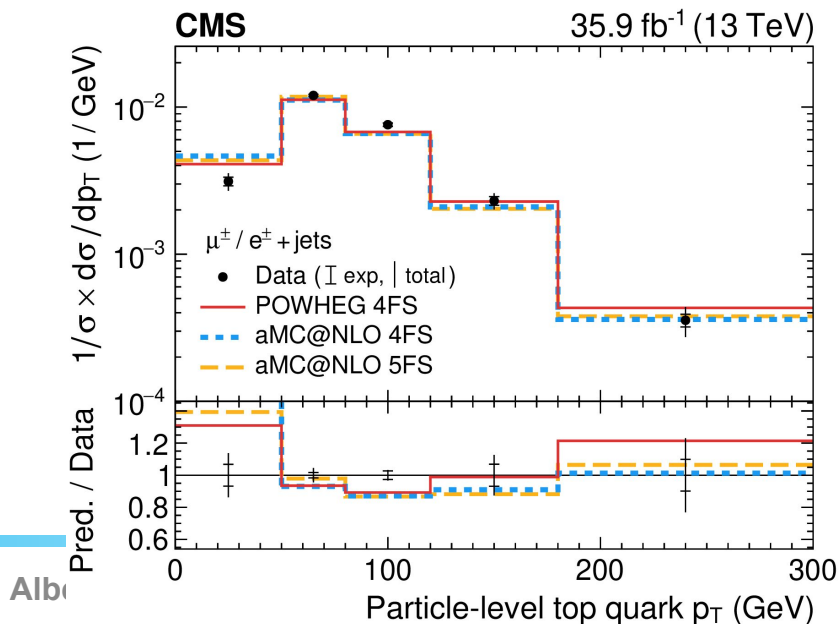
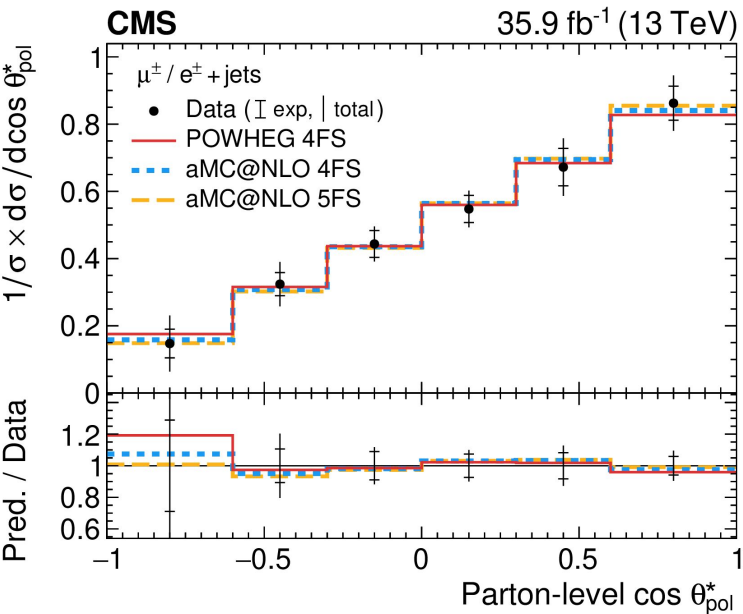
Going differential: spectrum and polarisation

Single top quark: different degrees of polarisation w.r.t. $t\bar{t}$ events: electroweak production.

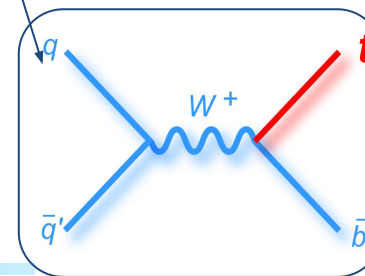
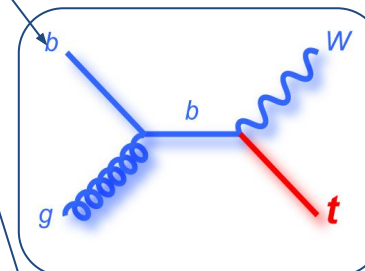
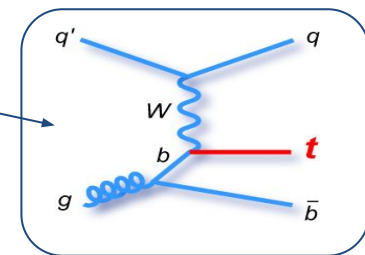
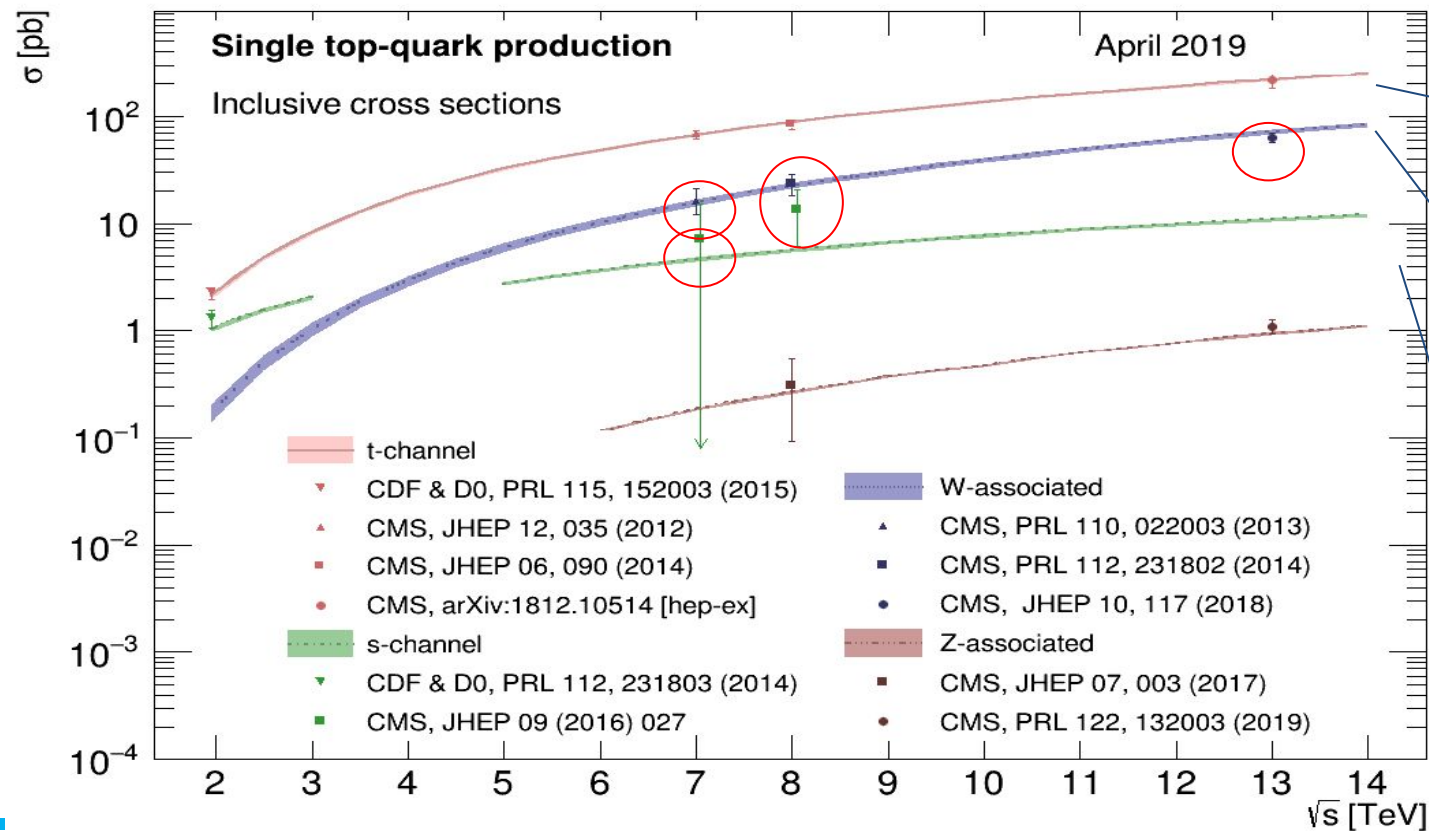
→ $\cos \theta^*$: lepton in the top rest frame vs spin axis, taken from the recoil jet

→ p_T -spectrum entirely different from $t\bar{t}$

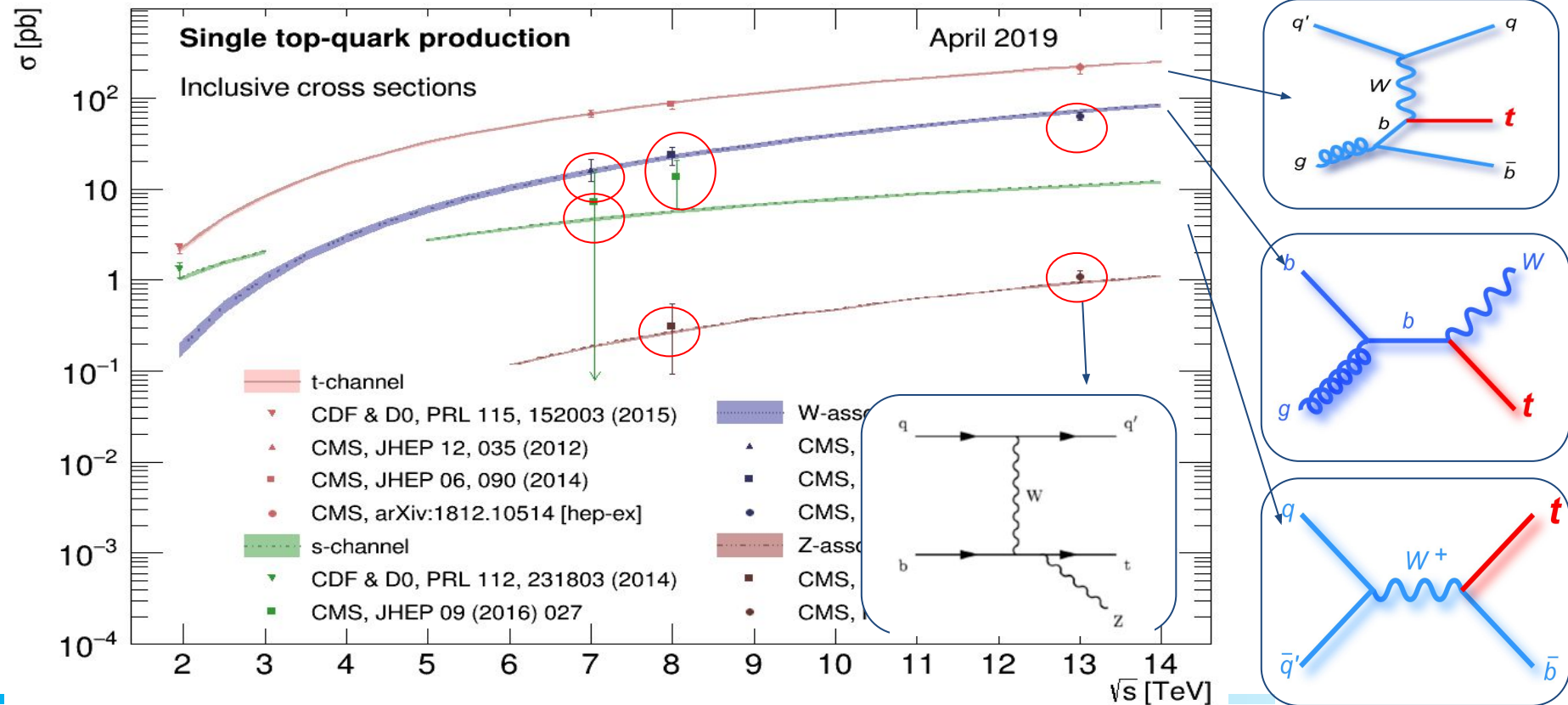
[TOP-17-023 Web page](#) / [arxiv:1907.08330](#)



Lower cross-section channels



Lower cross-section channels

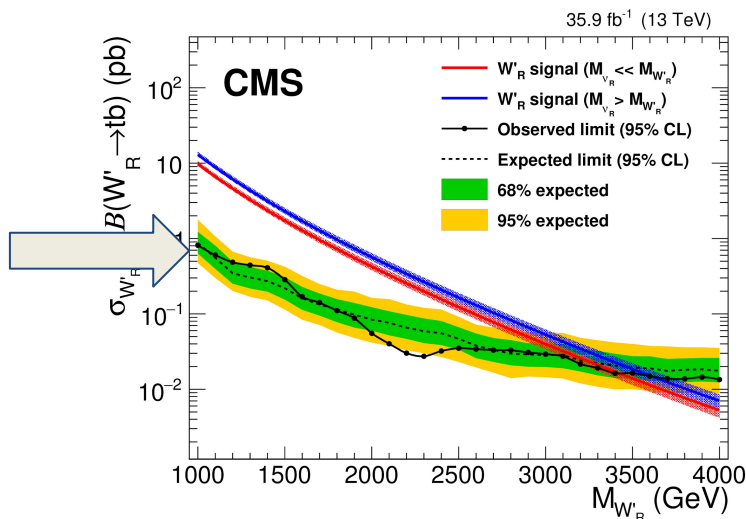
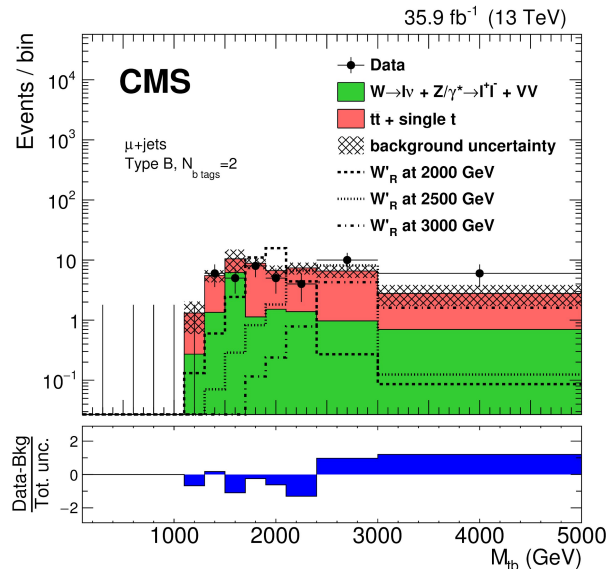
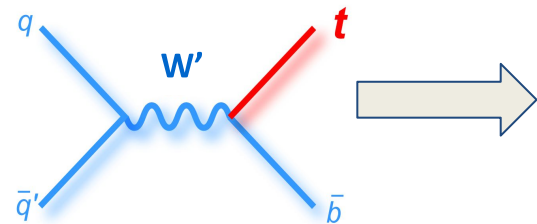


Single-top quark at high-pt / BSM

Single-top s-channel: in the SM it is strongly suppressed due to pdf effects and to the W being virtual.

→ In case of a new resonance, W' , this can be searched!

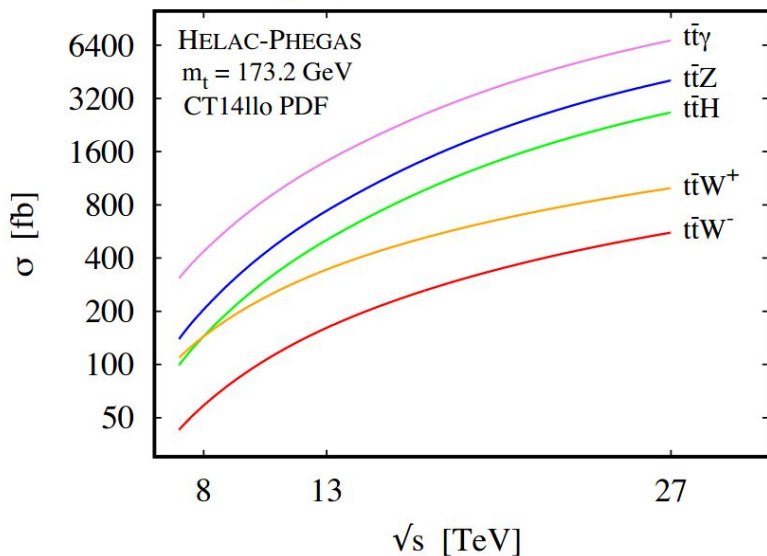
→ events with one lepton, one neutrino, one or two b-tagged jets:



Different Right-handed vs Left handed: can be inferred from angular quantities!

New goldmine of Run-II: $t\bar{t}X/tXj$ productions!

$t\bar{t}\gamma$, $t\bar{t}Z$, $t\bar{t}H$, $t\bar{t}W^+$ & $t\bar{t}W^-$ @ LHC

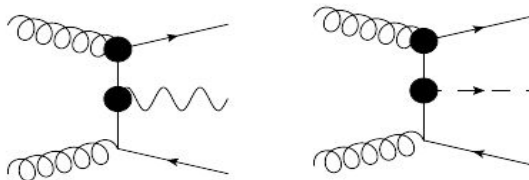


Plus 4-top production, single top+X channels etc

Evidence already!

Solid observation, several possibilities open:

- Can use precision measurements to constrain top operators
- Similar to what done for “classic” $t\bar{t}$ and t productions!

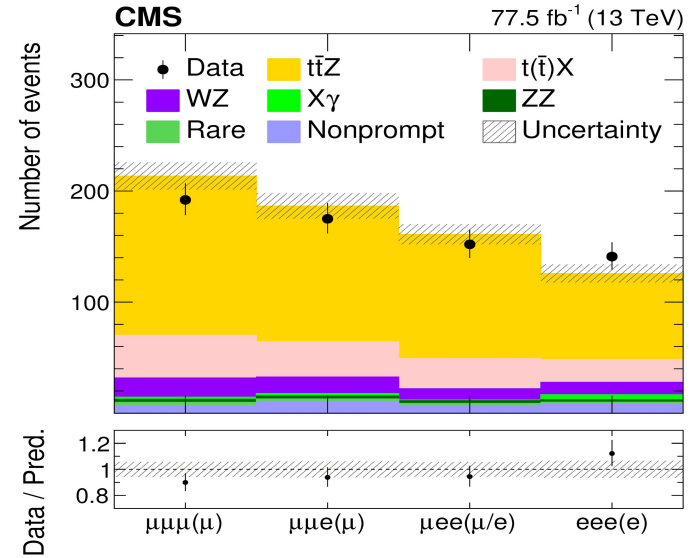
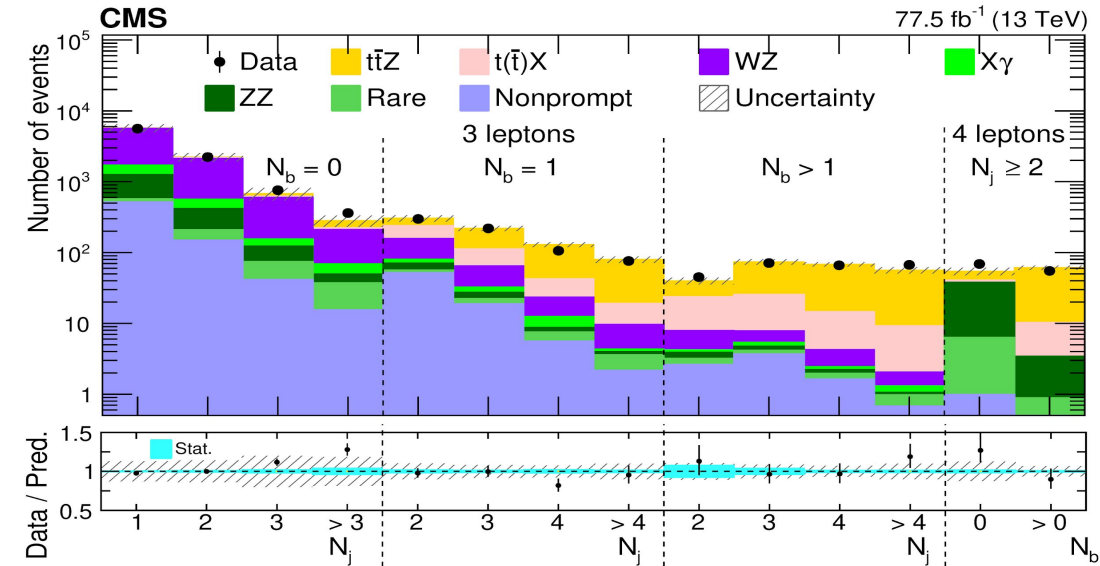


ttZ: from new channel to EFT probes

Very clean sample:

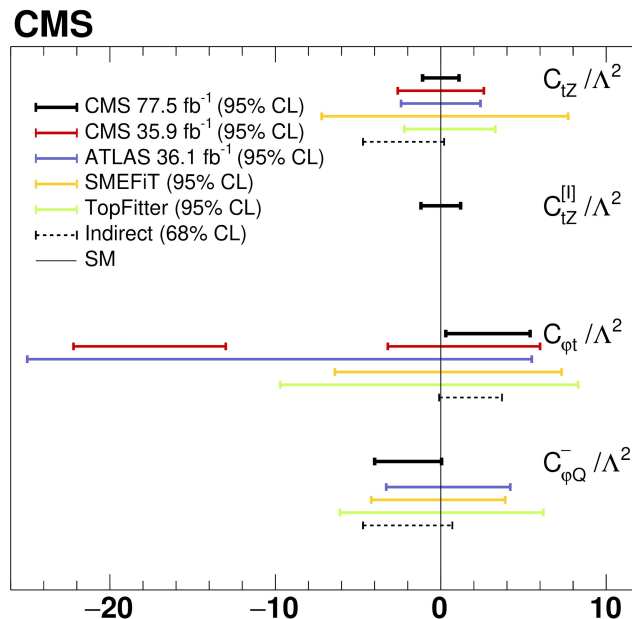
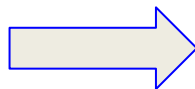
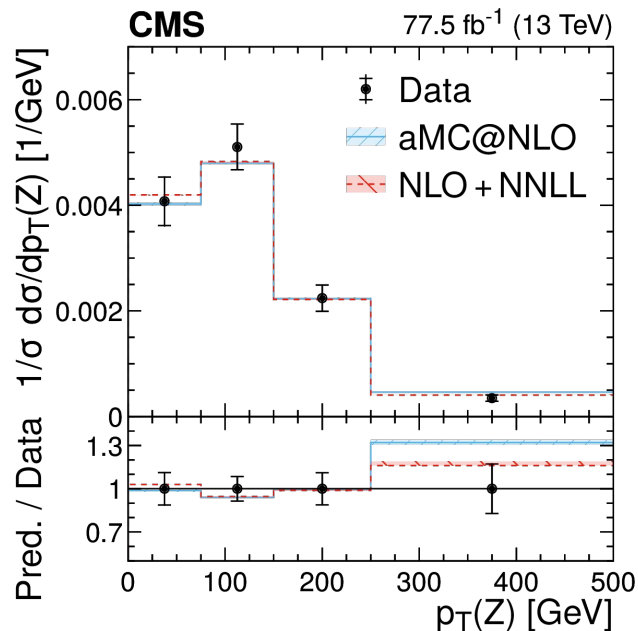
→ possible to have a lepton-antilepton pair + ≥ 1 b jet

→ Signal enriched regions: differential distributions to probe EFT!



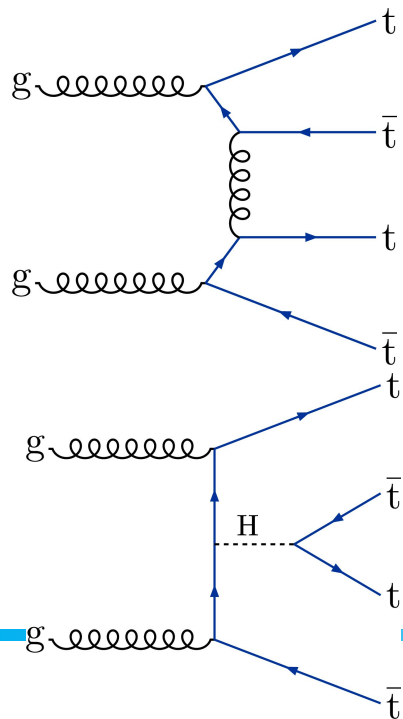
ttZ: from new channel to EFT probes

Contributions from EFT Can be seen in the Z angle and momentum:



4-top quark production

Very rare process: sensitive to:
 → 6-dimension operator involving top-higgs
 → top-higgs coupling y_t



Complex analysis:

≥ 2 leptons
 high HT >300 GeV
 ≥ 2 jets
 ≥ 2 bjets

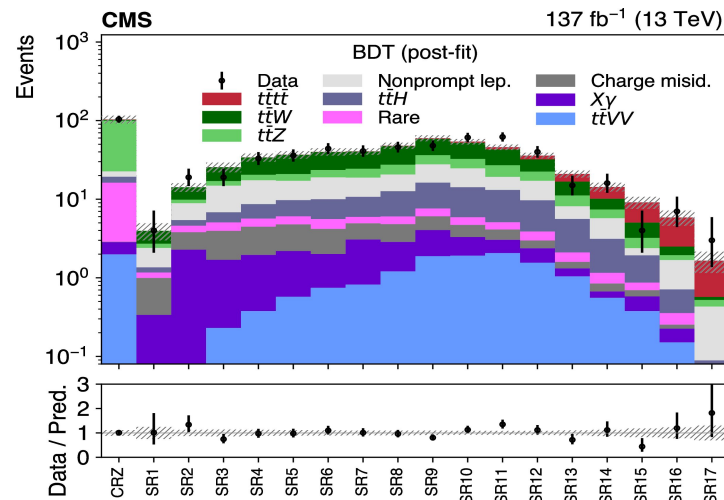
Multiple signal regions

depending on n leptons, jets, b

N_ℓ	N_b	N_{jets}	Region
2	2	≤ 5	CRW
		6	SR1
		7	SR2
		≥ 8	SR3
	3	5	SR4
		6	SR5
		≥ 7	SR6
≥ 3	≥ 4	≥ 8	SR7
		≥ 5	SR8
	2	5	SR9
		6	SR10
		≥ 7	SR11
	≥ 3	4	SR12
		5	SR13
		≥ 6	SR14
	Inverted resonance veto		CRZ

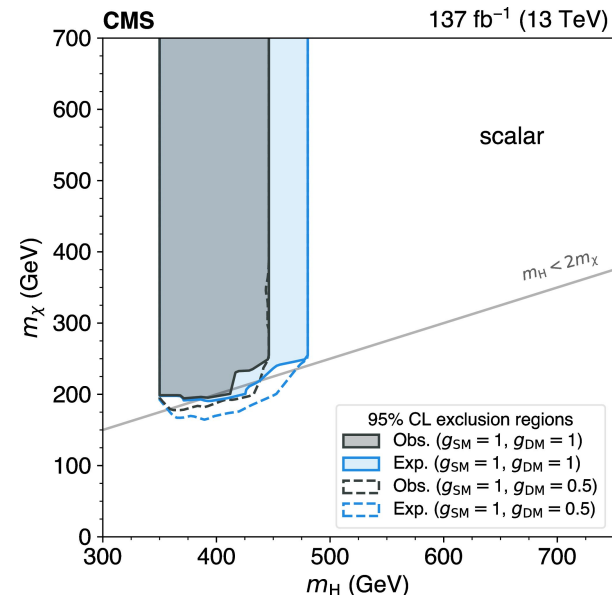
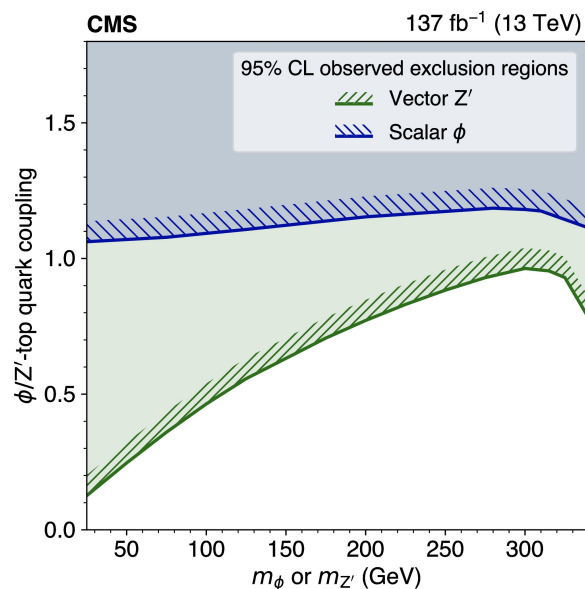
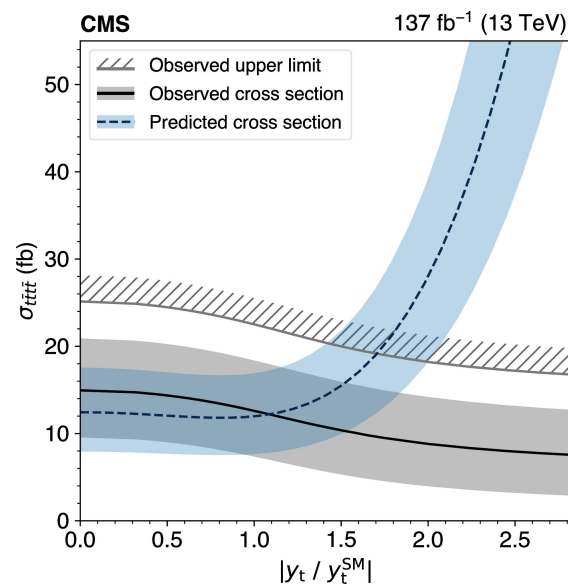
Simultaneous fit:

→ BDT in all channels
 → Cut and count analysis for cross-check



4-top quark production

Depending on the mediator: can be interpreted as EFT **or** new physics direct search (extra Z' /DM mediators)

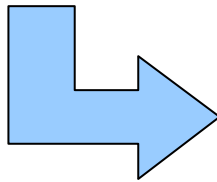


Complex BSM final states → Vector Like Quarks

Vector-Like quarks: new fermions appearing in several composite models, coupled mostly to 3rd family

SM		singlets	doublets	triplets
$\begin{pmatrix} u \\ d \end{pmatrix} \quad \begin{pmatrix} c \\ s \end{pmatrix} \quad \begin{pmatrix} t \\ b \end{pmatrix}$		T B	$\begin{pmatrix} X \\ T \end{pmatrix} \quad \begin{pmatrix} T \\ B \end{pmatrix} \quad \begin{pmatrix} B \\ Y \end{pmatrix}$	$\begin{pmatrix} X \\ T \\ B \end{pmatrix} \quad \begin{pmatrix} T \\ B \\ Y \end{pmatrix}$
$SU(2)_L$	2	1	2	3
$U(1)_Y$	$q_L = 1/6$ $U_R = 2/3$ $d_R = -1/3$	$2/3 \quad -1/3$	$1/6 \quad 7/6 \quad -5/6$	$2/3 \quad -1/3$

- Multiple representations
- $Q = I_3 + Y/2$



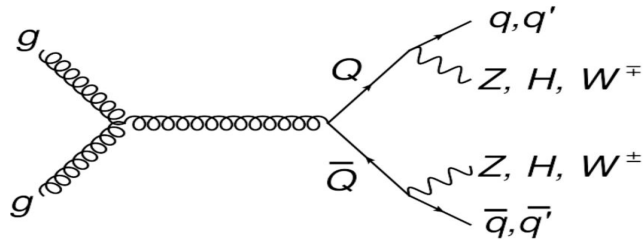
Appear in many SM extensions:

- **Little higgs, composite higgs**
- Warped **extra-dimensions**, Kaluza-Klein etc
- Some non-minimal **SUSY scenarios**

VLQ production

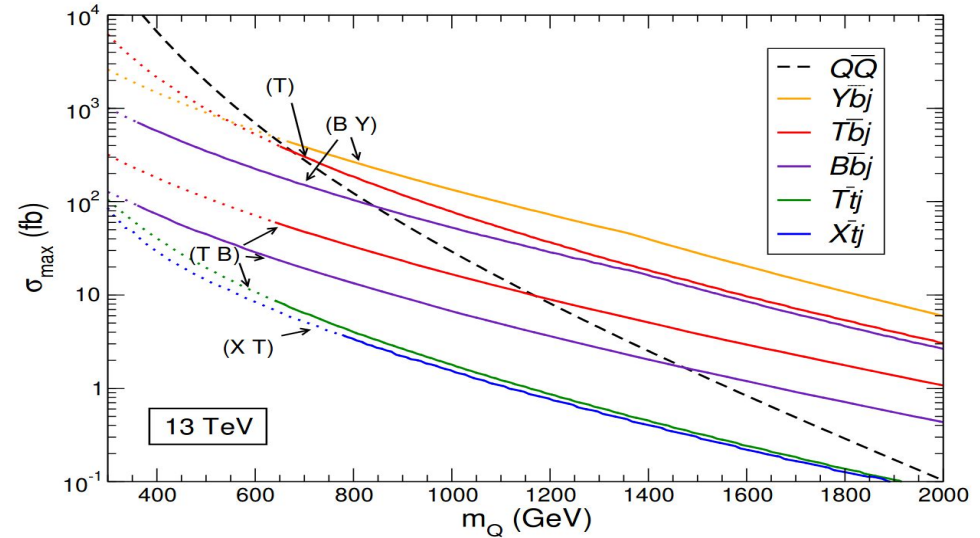
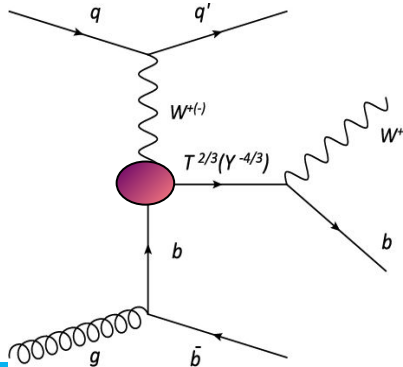
Pair-wise stong production: $Q\bar{Q}$

- depends only on the Q mass



Single production: qQ

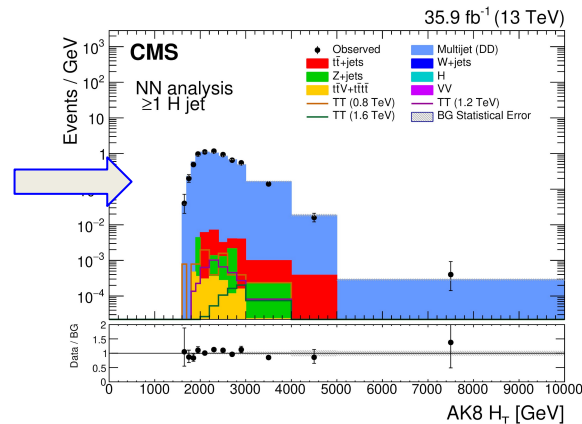
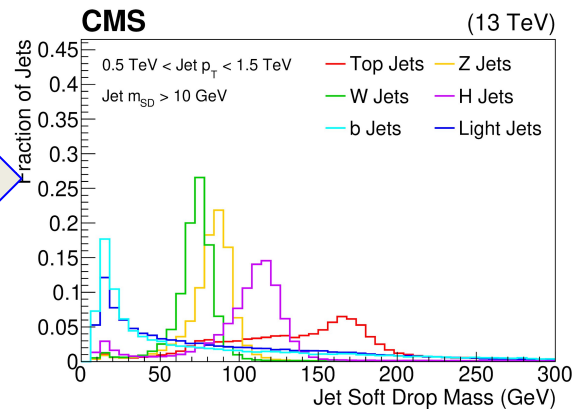
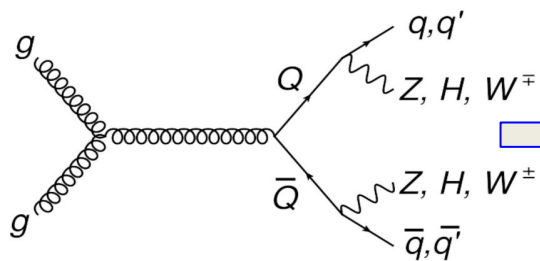
- depends on the Q mass and on the Qq coupling



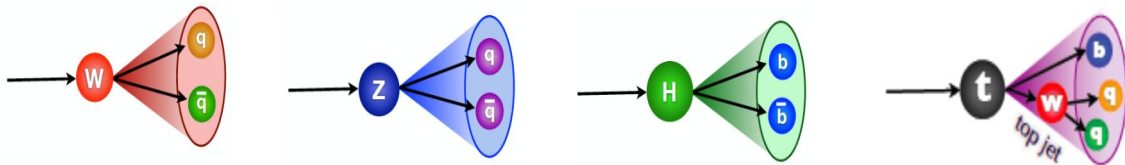
Both of them can reproduce the topologies we do measure in EFTs!

→ In particular single top quark ones!

VLQ pairs: all-hadronic



Inclusive search for $TT \rightarrow$ any combination of four: b, W, Z, H, t
 \rightarrow a-specific search looking at high energy spectrum
 \rightarrow Machine learning algorithm to separate different jets

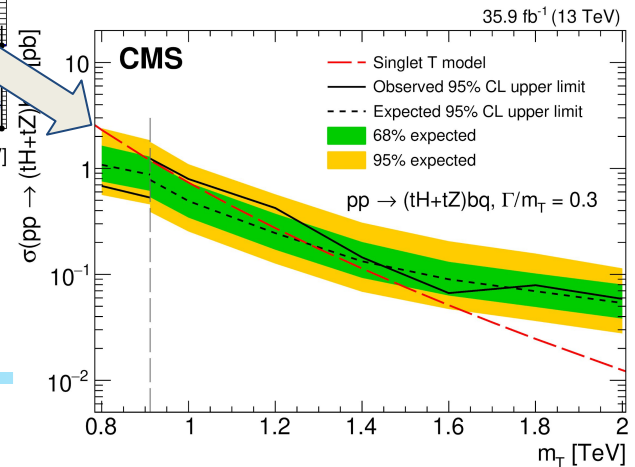
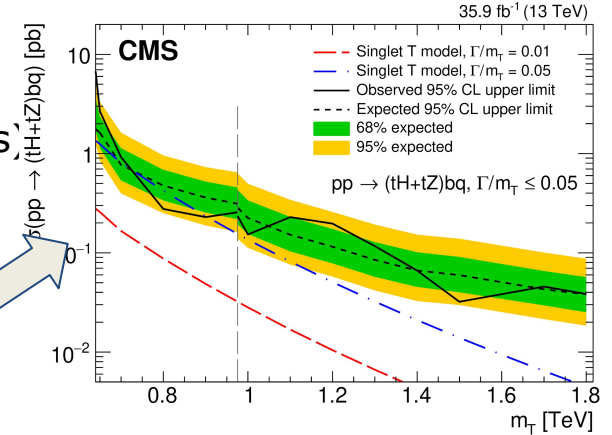
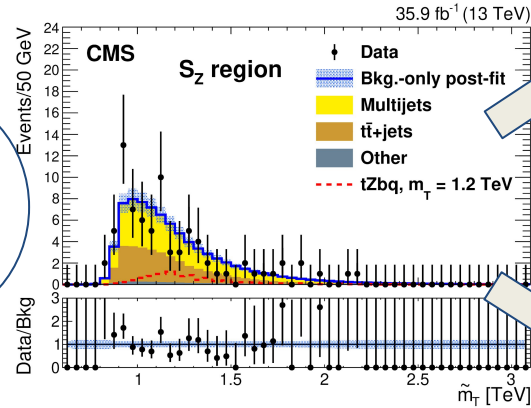
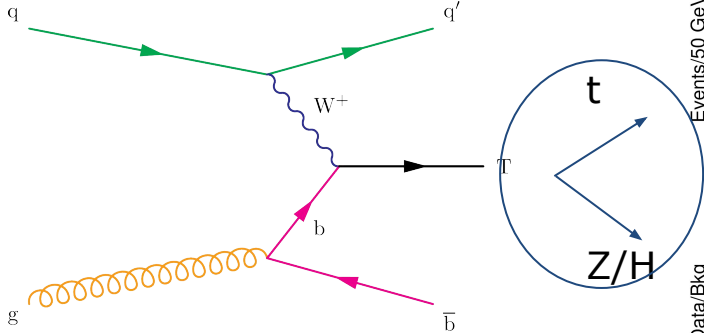


Single VLQs: $T \rightarrow tH/Z$

More specific search (in this case): one top + H/Z

→ Possible enhancement to tH/tZ in EFT

→ two analyses: resolved (5 narrow jets) and boosted (2 fat jets)



Also in this case: very different behavior depending on width

[B2G-18-003 Web page](#) / [JHEP01\(2020\)036](#)

Some final considerations

EFT @ top quark:

- Common framework for interpreting top quark measurements as potential hints for new physics
- Still plenty of work, not all of the interpretations are model independent!
- “Rare” channels also start from there!

Differential and inclusive spectra:

- Measurements of two-fold nature: comparison with SM and with EFT benchmarks.

New physics scenarios:

- High mass resonant searches : can have interplay with SM and EFT extensions!
- In general using different technologies and problems, but having a common language would benefit greatly looking forward Run-III and HL-LHC → Still work in progress!