

Axion quality from gauge flavour symmetries

Based in part on arXiv:2102.05055 (L. Darmé, EN) and on work in progress (Grilli Di Cortona, L. Darmé, C. Smarra)

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- Non-pert. (instanton related) QCD effects generate a potential $V_{\text{QCD}}(a) = -(m_\pi f_\pi)^2 \cos(a/v_a)$ that drives $\langle a/v_a \rangle \rightarrow 0$ at the minimum

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The generating functional $Z \sim \int [DA_\mu D\Phi] D\psi D\bar{\psi} \exp(iS)$

is not invariant under a PQ transformation. $U(1)_{PQ}$ cannot be "imposed"

- In benchmark axion models, Φ is a complex scalar, and a gauge singlet.

Renormalizable terms $\mu^3\Phi$, $\mu^2\Phi^2$, $\mu\Phi^3$, $\lambda\Phi^4$ do not break gauge or Lorentz and cannot be forbidden. However, they would destroy $U(1)_{PQ}$ and the axion solution!

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- moreover: Non-pt. quantum gravity effects do not respect global symmt.

Controlled solutions [Euclid. wormholes] do generate : $e^{-S_{wh}} M_P^3 \Phi + h.c.$

Safe suppression requires $S_{wh} > 190$, while typically $S_{wh} \sim \text{Log}(M_P/v_a) \sim 15$

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- Quality: Effective ~~PQ~~ opts.: ~~PQ~~ vacuum eng. density $< 10^{-10} V_{QCD}(a)$

For $f_a \sim 10^{10} \text{ GeV}$ and effective scale M_P , this implies Eff. Opt. Dim. ≥ 10

[Barr & Seckel '92, Kamionkowski & March-Russel '92, Holman et al. '92, Ghigna et al. '92]

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- **Non-Abelian $SU(N)_L \times SU(N)_R$** , $a(x) \in Y_{n \times n}$ "Orbital mode" $Y = U \hat{Y} V^\dagger e^{ia/v_a}$
 $N > 4$, \mathcal{L}_{ren} has an automatic rephasing symm. $V = V(Y Y^\dagger)$ $Y \rightarrow e^{i\xi} Y$.
1st ~~PQ~~ opt. $\Lambda^{4-N} \det Y$ $\dim = N$. This requires again $N \geq 10$
(Fong, EN '14 [in $SU(3) \times SU(3)$], Di Luzio, Ubaldi, EN '17)

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Gauge invariants are constructed with Kronecker δ and Levi-Civita ε

δ -invariants involve $Y^\dagger Y$:

They are obviously **all Hermitian** \Rightarrow accidental U(1): $Y \rightarrow e^{i\xi} Y$

ε -invariants (non-Hermitian): there is none $\varepsilon_{\alpha\beta\dots\sigma} Y_{\alpha i} Y_{\beta j} \dots Y_{\sigma r} = 0$.

Some $SU(n)$ index (i, j, \dots, r) must coincide: ε opts. vanish symmetrically

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Note: for a $Y_{n \times n}$ square matrix $\varepsilon_{\alpha\beta\dots\sigma} \varepsilon_{ij\dots r} Y_{\alpha i} Y_{\beta j} \dots Y_{\sigma r} \propto \det Y \neq 0$

Such automatic exact U(1) symmetries are peculiar of local `rectangular' symmetries

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- Two mixed invariants $I_\varepsilon = \varepsilon_{\alpha\beta\gamma} \varepsilon_{ij} Y_{ai} Y_{\beta j} Z_\gamma \neq 0$ preserves $U(1)_\varepsilon$ ~~$U(1)_\delta$~~
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Operators for which $\langle O \rangle = 0$ do not break the symmetries of the minimum,
the vacuum can enjoy a larger symmetry than the Lagrangian.

Scalar bosons associated with these symmetries remain massless [Georgi & Pais '75]

The axion can remain protected even if $V(Y, Z)$ breaks $U(1)_{PQ}$

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The non-Abelian local G_F thus is a flavour symmetry!

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We are currently studying flavour groups that we would never have considered, had it not been for the axion !