

# Heterodyne Detection of Axion Dark Matter in an RF Cavity

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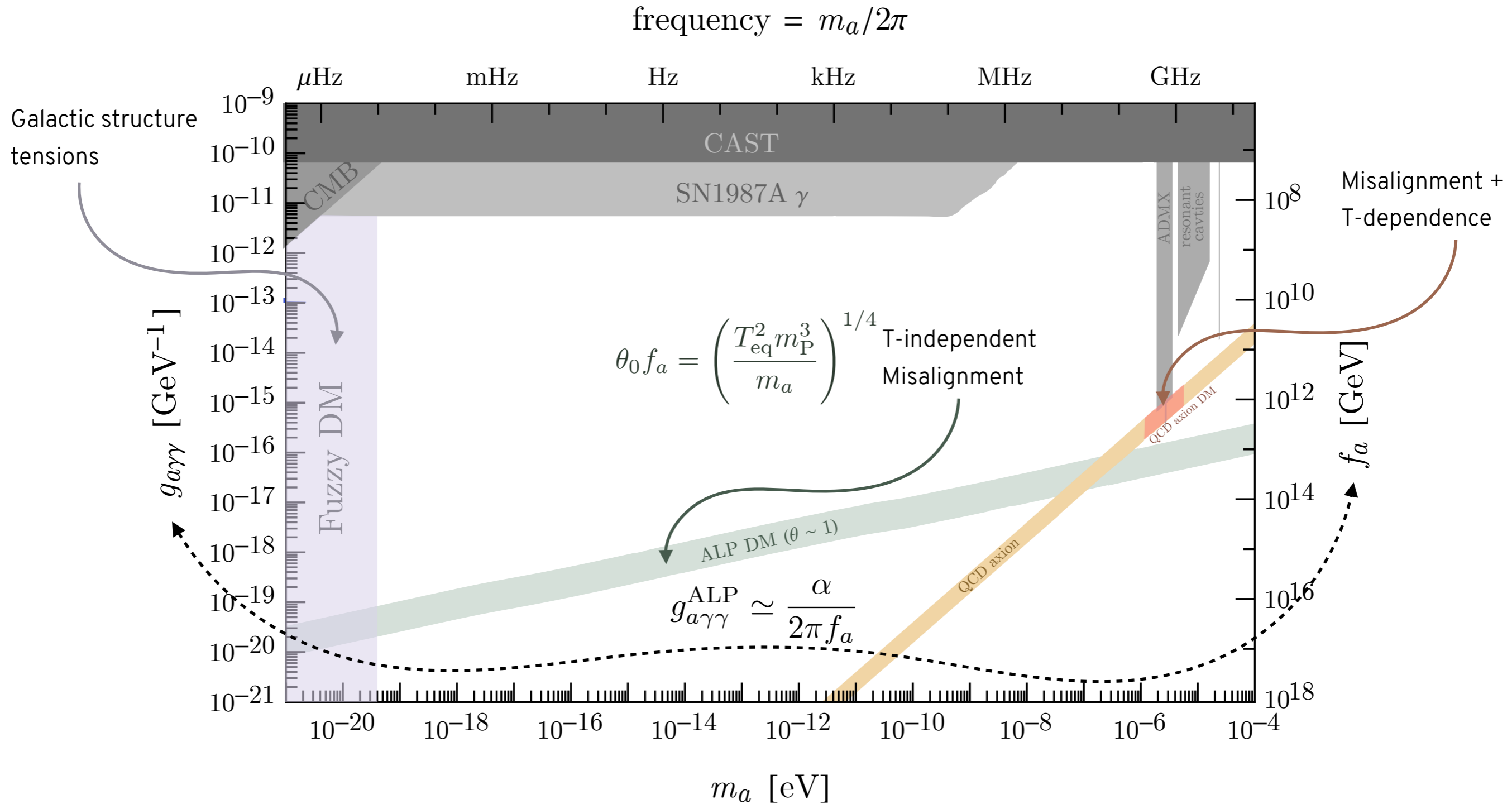
hep-ph/2007.15656

A. Berlin, R. T. D'Agnolo, SARE, K. Zhou

*JHEP* 07 (2020) 088, hep-ph/1912.11048

A. Berlin, R. T. D'Agnolo, SARE, P. Schuster, N. Toro, C.  
Nantista, J. Neilson, S. Tantawi, K. Zhou

# Axions as Dark Matter: Targets



# Resonant Axion Searches

Axion electrodynamics:  $\mathcal{L} \supset -\frac{g_{a\gamma\gamma}}{4} a F \tilde{F} = -g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$

$$\begin{aligned}\nabla \cdot \mathbf{E} &= \rho - g_{a\gamma\gamma} \mathbf{B} \cdot \nabla a \\ \nabla \times \mathbf{B} &= \partial_t \mathbf{E} + \mathbf{J} - g_{a\gamma\gamma} (\mathbf{E} \times \nabla a - \mathbf{B} \partial_t a)\end{aligned}$$

Maxwell's new and improved Equations

Axion dark matter:  $a(t) \simeq \frac{\sqrt{2\rho_{\text{DM}}}}{m_a} \cos(m_a t + \varphi)$

$$J_{\text{eff}}(t) \sim g_{a\gamma\gamma} B_0(t) \sqrt{\rho_{\text{DM}}} \cos m_a t \implies B_a(t) \propto J_{\text{eff}}(t)$$

# Resonant Axion Searches

Axion-induced magnetic field induces an E.M.F.:  $\mathcal{E}_a \sim V^{2/3} \partial_t B_a$

$$P_{\text{sig}}^{(r)} \sim \frac{\mathcal{E}_a^2}{R} \min\left(1, \frac{\tau_a}{\tau_r}\right) \sim \omega_{\text{sig}}^2 B_a^2 V \min\left(\frac{Q_r}{\omega_{\text{sig}}}, \frac{Q_a}{m_a}\right)$$

$$1/\tau_a \sim m_a \langle v^2 \rangle$$

$$1/\tau_r \sim \omega_{\text{sig}}/Q_r$$

$$Q_a \sim 1/\langle v^2 \rangle$$

**Maximise:**  $\omega_{\text{sig}}, B_a, V$



**Cavities, LC:**  $B_a \sim J_{\text{eff}} V^{1/3}$

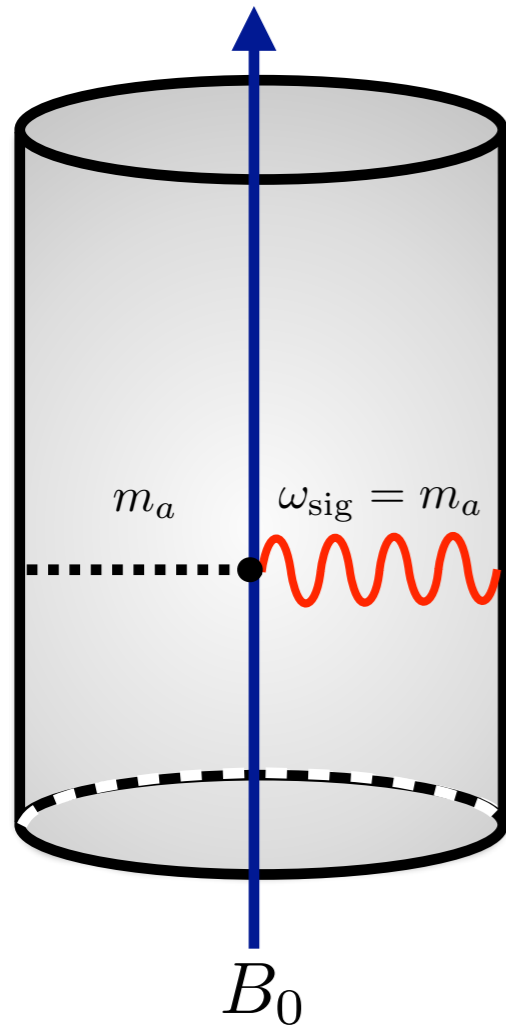
**Cavities:**  $\omega_{\text{sig}} \sim V^{-1/3}$

# Resonant Approaches

Static-field Haloscope:

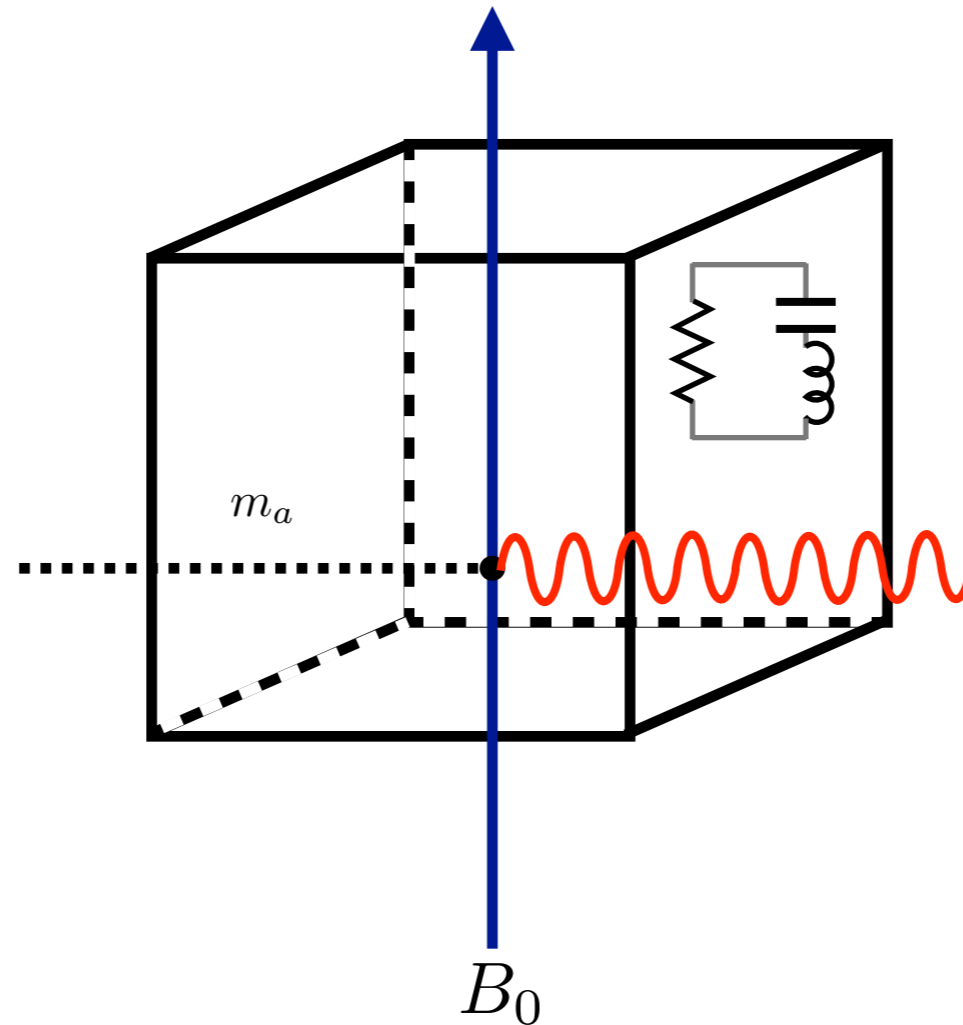
e.g. ADMX

$$\omega_{\text{sig}} = m_a \sim V^{-1/3}$$



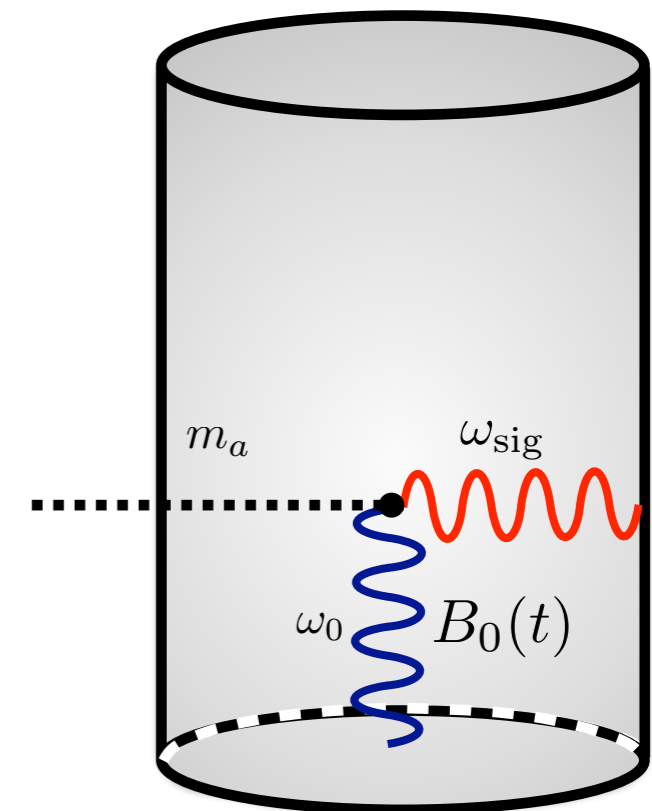
LC Resonator:

$$\omega_{\text{sig}} = m_a = \omega_{\text{LC}}$$



Heterodyne Resonator:

$$\omega_{\text{sig}} \sim \omega_0 \pm m_a \sim V^{-1/3}$$



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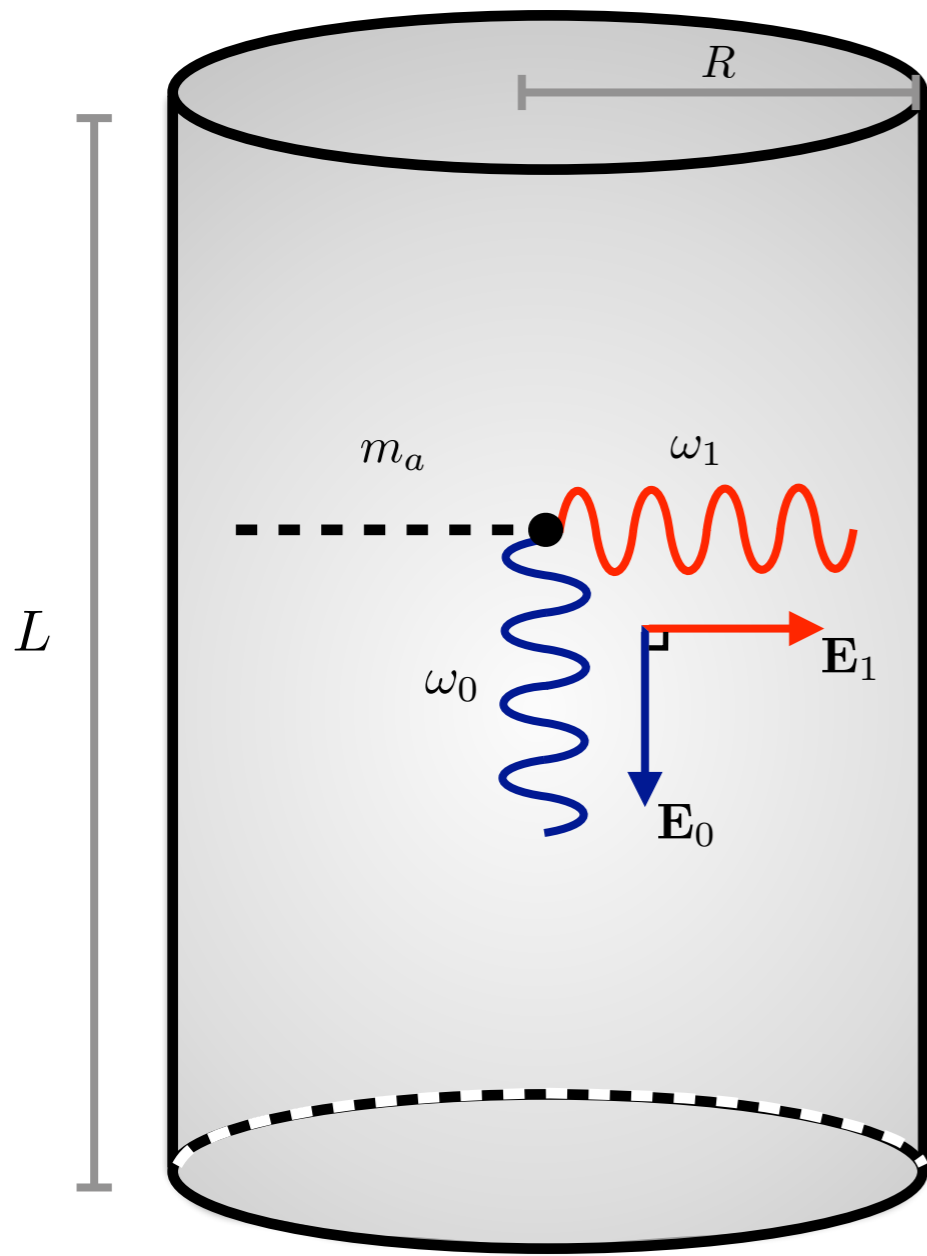
A. Berlin, R. T. D'Agnolo, SARE, P.

Schuster, N. Toro, C. Nantista, J.

Neilson, S. Tantawi, K. Zhou

Also: R. Lasenby hep-ph/1912.11467

# Allowed mode transitions



Cylindrical cavity modes:  $\text{TM}_{mnp}$   $\text{TE}_{mnp}$

Axion allowed transitions:  $\eta \propto \int_V \mathbf{E}_1^* \cdot \mathbf{B}_0 \neq 0$   
**MAXIMISE**

*spin-0*:  $m_0 = m_1$

*pseudoscalar*:  $p_0 + p_1 = \text{odd}$

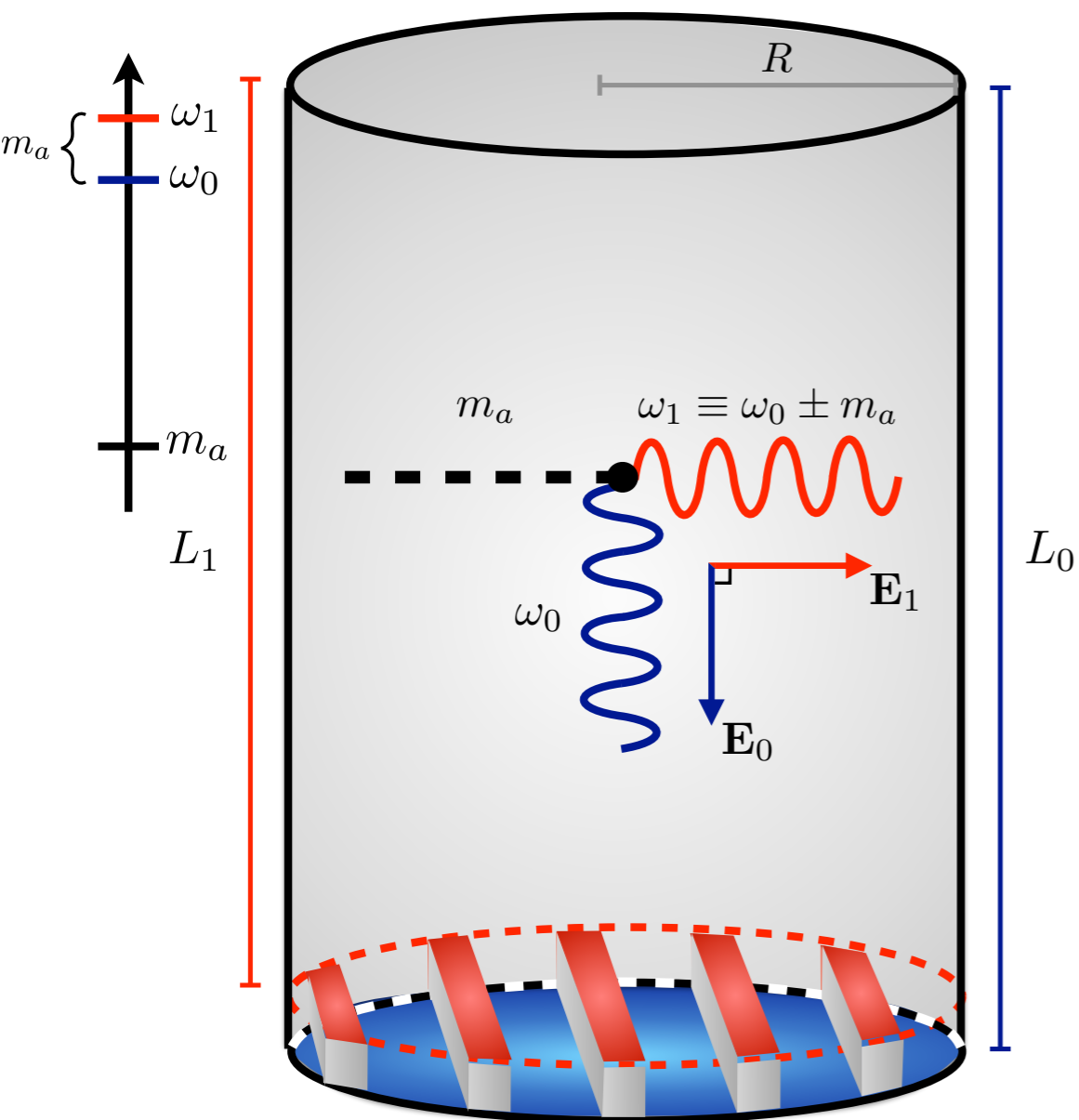
$\text{TM}_{mnp} \rightleftharpoons \text{TM}_{mnp}$

$\updownarrow$

$\text{TE}_{mnp} \rightleftharpoons \text{TE}_{mnp}$

*mode degeneracy*:  $\frac{L}{R} = \left( \frac{\pi(p_1^2 - p_0^2)}{x_{mn_0}^2 - x'_{mn_1}{}^2} \right)^{1/2}$

# Scanning Axion Resonant Frequency Conversion



Superconducting RF Cavity

$$\omega_0 = \omega_1 \sim \text{GHz}$$

$$Q_{\text{int}} \sim 10^9 \div 10^{13}$$

Tunability:

$$\delta\omega \lesssim \text{MHz} \quad \text{piezos}$$

$$\delta\omega \gtrsim \text{MHz} \quad \text{fins}$$

Degeneracy:

$$\frac{L}{R} = \left( \frac{\pi(p_1^2 - p_0^2)}{x_{mn_0}^2 - x'_{mn_1}{}^2} \right)^{1/2}$$

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Toro, C. Nantista, J. Neilson, S. Tantawi, K. Zhou

**Broadband:**

hep-ph/2007.15656

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# Axion Signal

Signal Power Spectral Density (PSD):

$$S_{\text{sig}}(\omega) = \frac{\omega_1}{Q_1} (g_{a\gamma\gamma} \eta_{10} B_0)^2 V \frac{\omega^2}{(\omega^2 - \omega_1^2)^2 + (\omega \omega_1 / Q_1)^2} \int \frac{d\omega'}{(2\pi)^2} (\omega' - \omega)^2 S_{b_0}(\omega') S_a(\omega - \omega')$$

Axion PSD:  $\langle a(t)^2 \rangle = \frac{1}{(2\pi)^2} \int d\omega S_a(\omega) = \frac{\rho_{\text{DM}}}{m_a^2}$

Background magnetic field PSD: To be discussed further...

$$S_{b_i}(\omega) = \pi^2 \left( \delta(\omega - \omega_i) + \delta(\omega + \omega_i) \right) + S_{b_i}^{(\text{phase})} + S_{b_i}^{(\text{mech})}$$

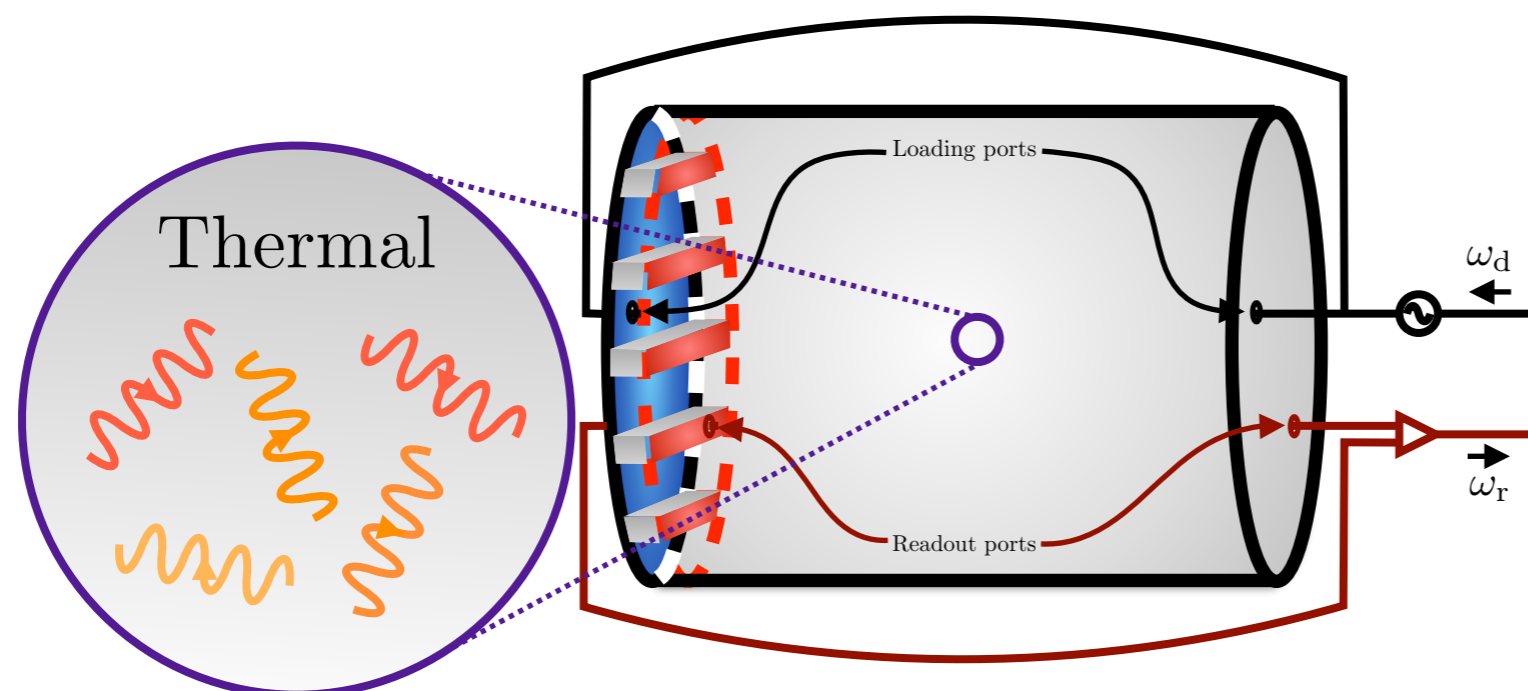
NB:  $B_i \equiv \sqrt{\frac{1}{V_{\text{cav}}} \int_{V_{\text{cav}}} |\mathbf{B}_i(x)|^2}$        $\mathbf{B}_i(x, t) = \mathbf{B}_i(x) b_i(t)$



# Standard Noise Sources: Thermal Noise

Power Spectral Density:

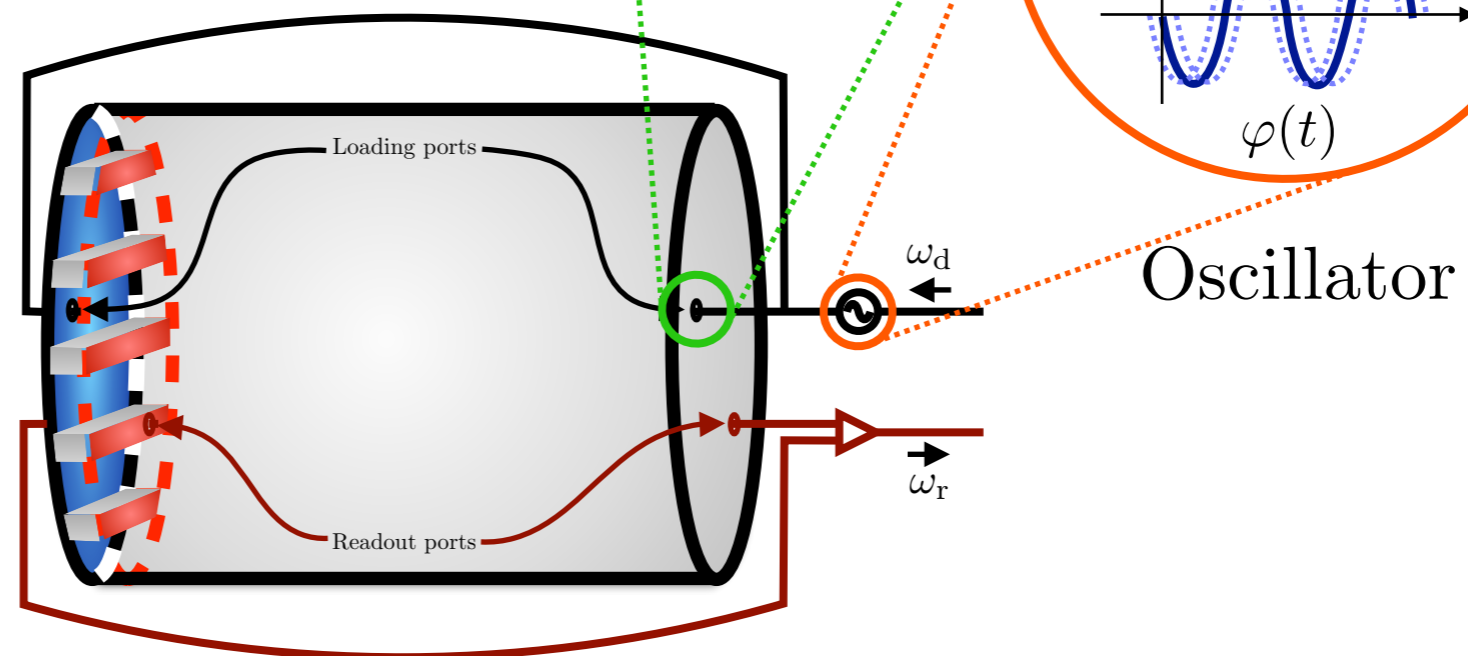
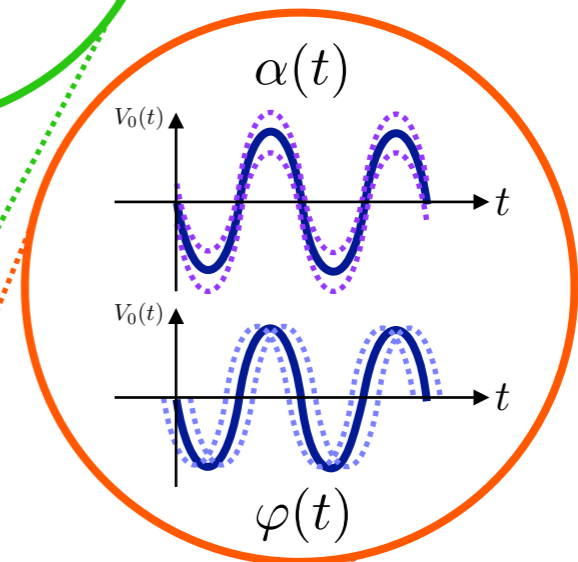
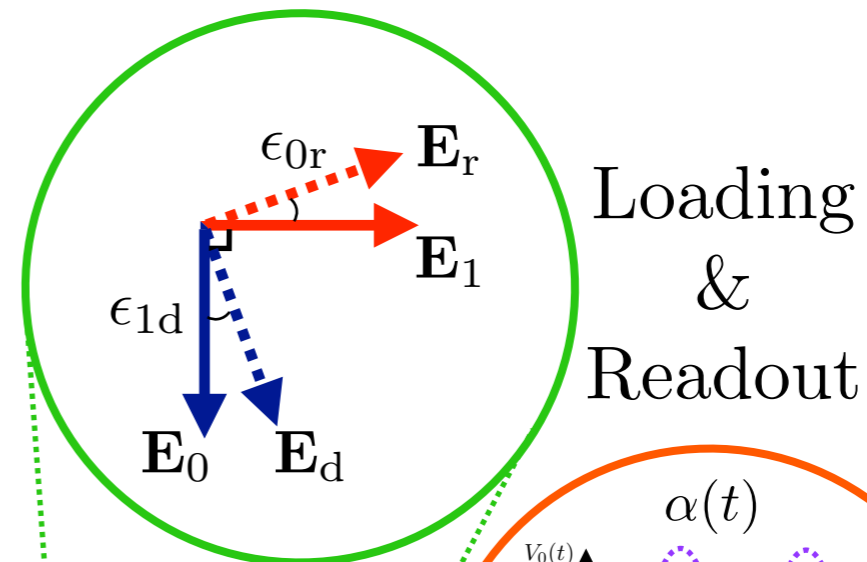
$$S_{\text{th}}(\omega) = \frac{Q_1}{Q_{\text{int}}} \frac{4\pi T (\omega \omega_1 / Q_1)^2}{(\omega^2 - \omega_1^2)^2 + (\omega \omega_1 / Q_1)^2}$$



# Non-standard Noise Sources: Phase Noise

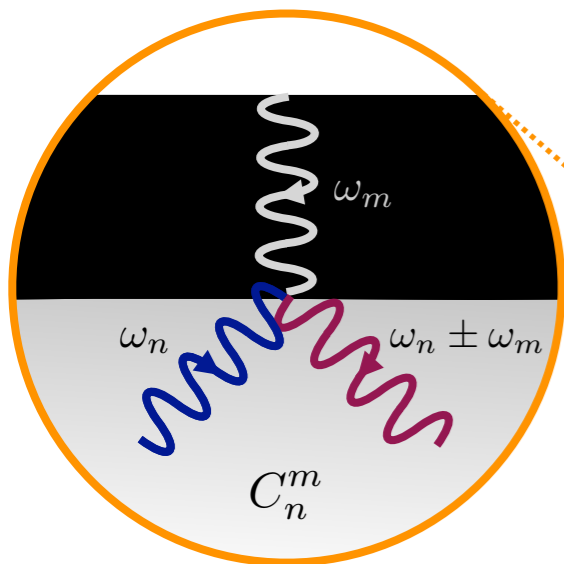
Power Spectral Density:

$$S_{\text{phase}}(\omega) \approx \frac{1}{2} \epsilon_{1d}^2 S_{\varphi}(\omega - \omega_0) \times \frac{(\omega \omega_1 / Q_1)^2}{(\omega^2 - \omega_1^2)^2 + (\omega \omega_1 / Q_1)^2} \frac{\omega_0 Q_1}{\omega_1 Q_0} P_{\text{in}}$$



# Non-standard Noise Sources: Wall Vibrations

Vibrations

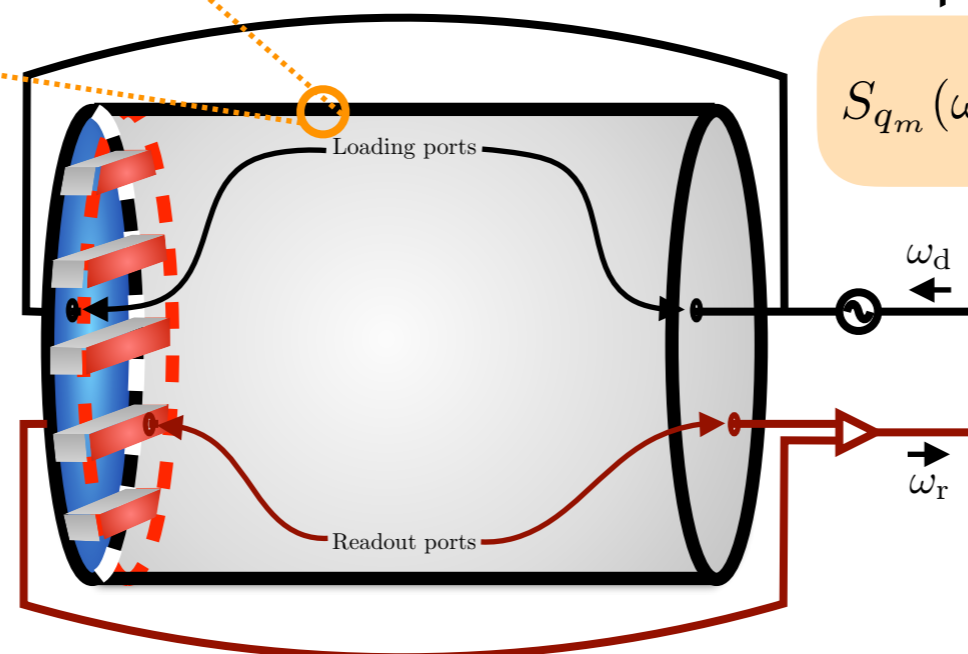


Power Spectral Density:

$$S_{\text{mech}}(\omega) \simeq \frac{\epsilon_{1d}^2}{4} \frac{\omega_0}{Q_0} P_{\text{in}} \times \sum_{n=0,1} \frac{(S_{q_m}(\omega - \omega_0)/V^{2/3}) (\omega_n/Q_n) \omega_n^4 \omega^2}{[(\omega^2 - \omega_n^2)^2 + (\omega \omega_n/Q_n)^2] [(\omega_0^2 - \omega_n^2)^2 + (\omega_0 \omega_n/Q_n)^2]}$$

Displacement PSD:

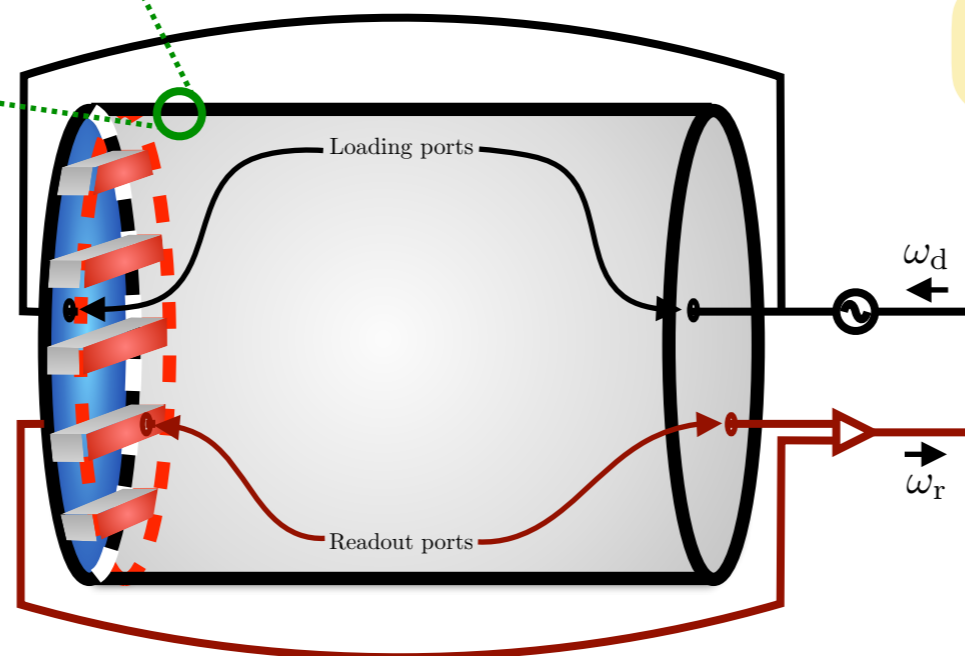
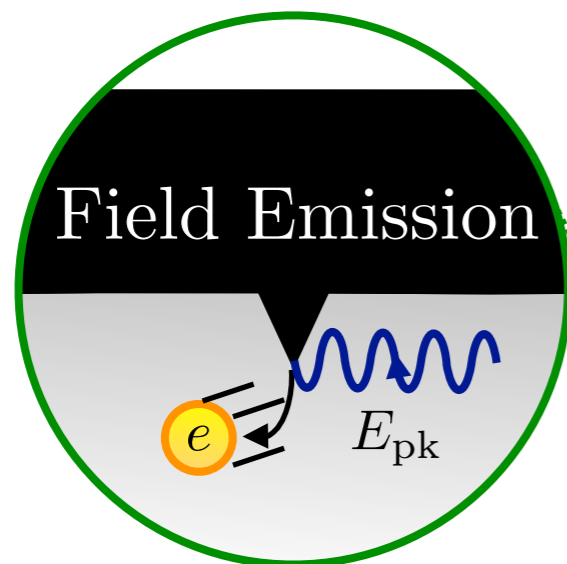
$$S_{q_m}(\omega) \simeq \frac{1}{M^2} \frac{S_{f_m}(\omega)}{(\omega^2 - \omega_m^2)^2 + (\omega_m \omega/Q_m)^2}$$



# Non-standard Noise Sources: Field Emission

Power Spectral Density:

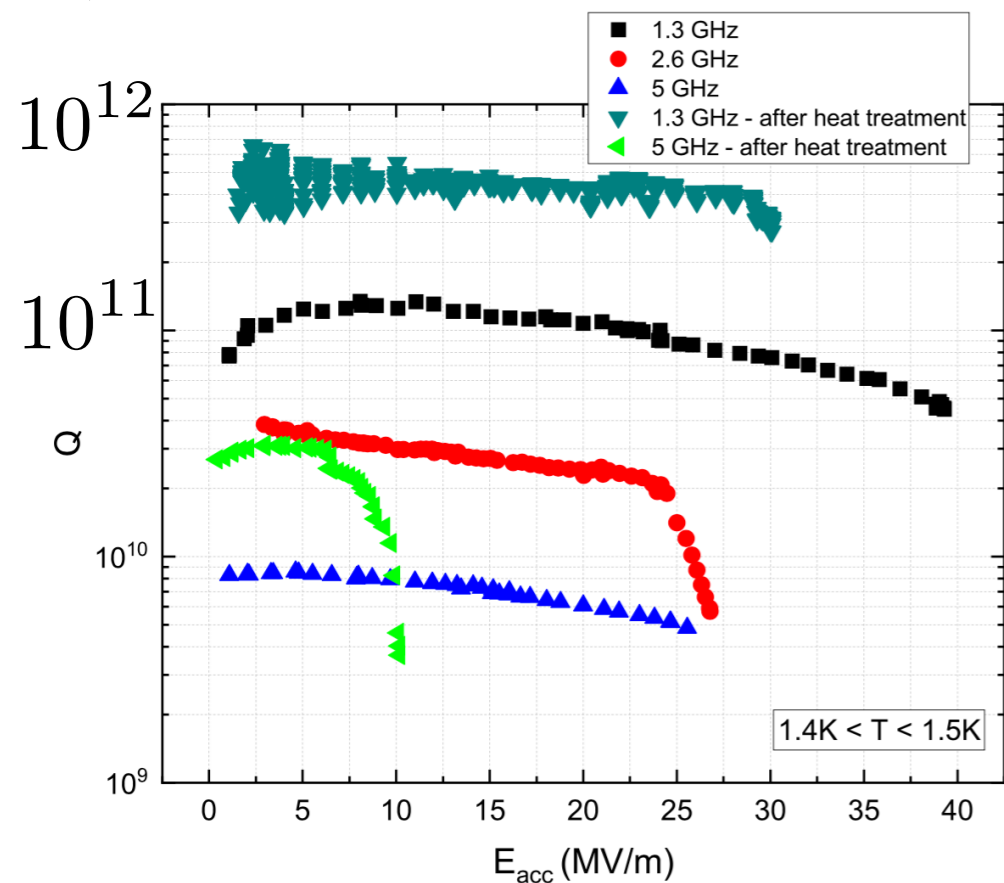
$$\frac{S(\omega_1)}{4\pi T} \sim \frac{P_{\text{tot}}}{0.1 \text{ W}} \times \begin{cases} 1 & \text{synchrotron} \\ 10^{-6} & \text{transition} \\ 10^{-5} & \text{Bremsstrahlung,} \end{cases}$$



Limits max B-field ~ 0.2T

# Experimental precedent

## Q-factor & B-field:



arXiv: 1810.03703 Romanenko et al.



gr-qc/0502054 Ballantini et al.

physics/0004031 Bernard, Gemme, Parodi, Picasso

Low-frequency  
seismic noise:

$$\Delta\omega/\omega \sim \delta \sim 10^{-10}$$

DarkSRF (2020)

# Signal to Noise

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Roughly:

$$(\text{SNR})^2 \simeq t_{\text{int}} \int_0^\infty d\omega \left( \frac{S_{\text{sig}}(\omega)}{S_{\text{noise}}(\omega)} \right)^2$$

Thermal noise dominated:

$$\text{SNR} \sim \frac{\rho_{\text{DM}} V}{m_a \omega_1} (g_{a\gamma\gamma} \eta_{10} B_0)^2 \left( \frac{Q_a Q_{\text{int}} t_e}{T} \right)^{1/2}$$

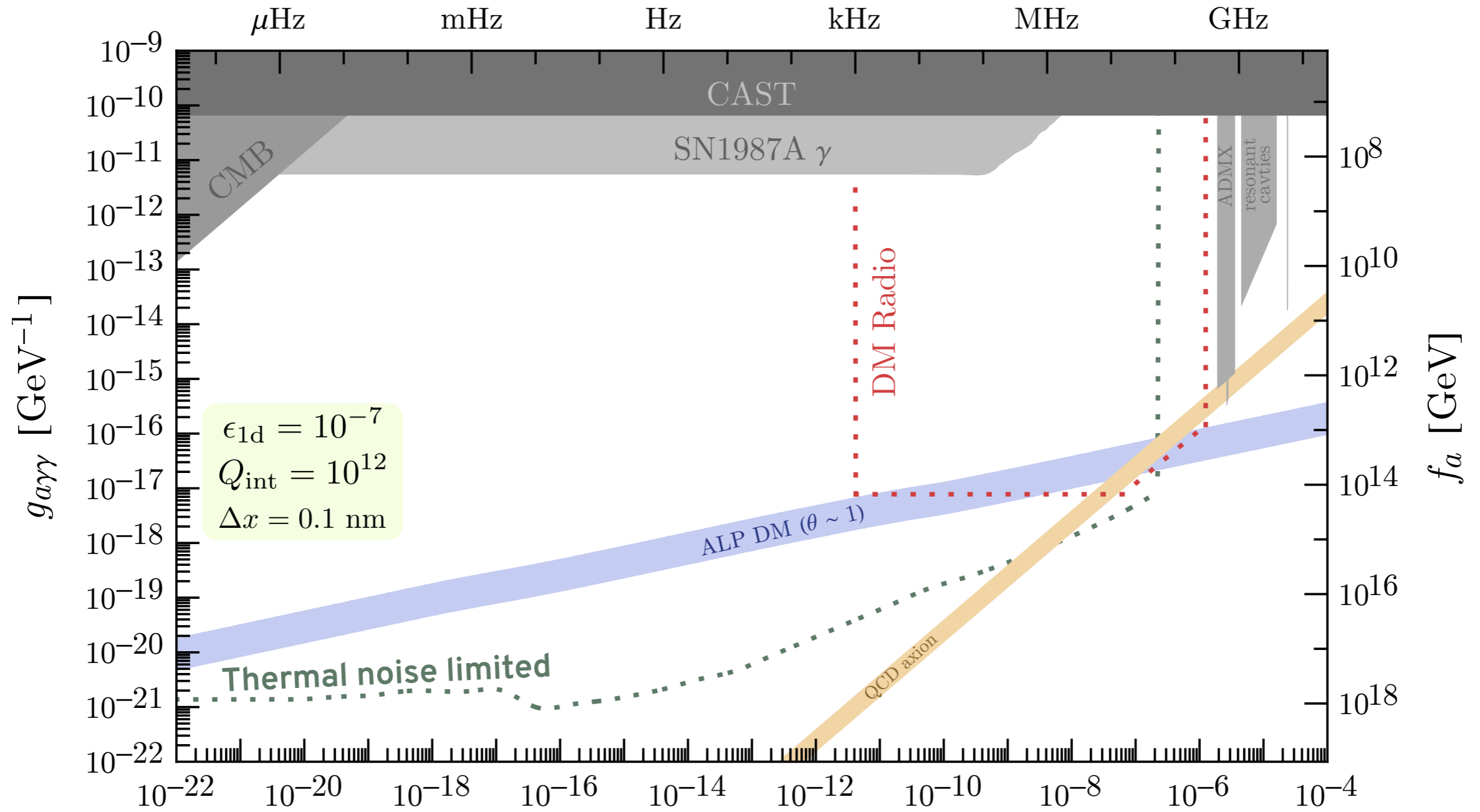
Comparison with LC resonator:

$$\frac{\text{SNR}}{\text{SNR}^{\text{LC}}} \sim \frac{\omega_0 \pm m_a}{m_a} \left( \frac{Q_{\text{int}}}{Q_{\text{LC}}} \right)^{1/2} \left( \frac{T_{\text{LC}}}{T} \right)^{1/2} \left( \frac{B_0}{B_{\text{LC}}} \right)^2$$

# Resonant Axion Resonant Frequency Conversion

**B = 0.2 T, T = 2K,  $\omega_0 = 1$  GHz**

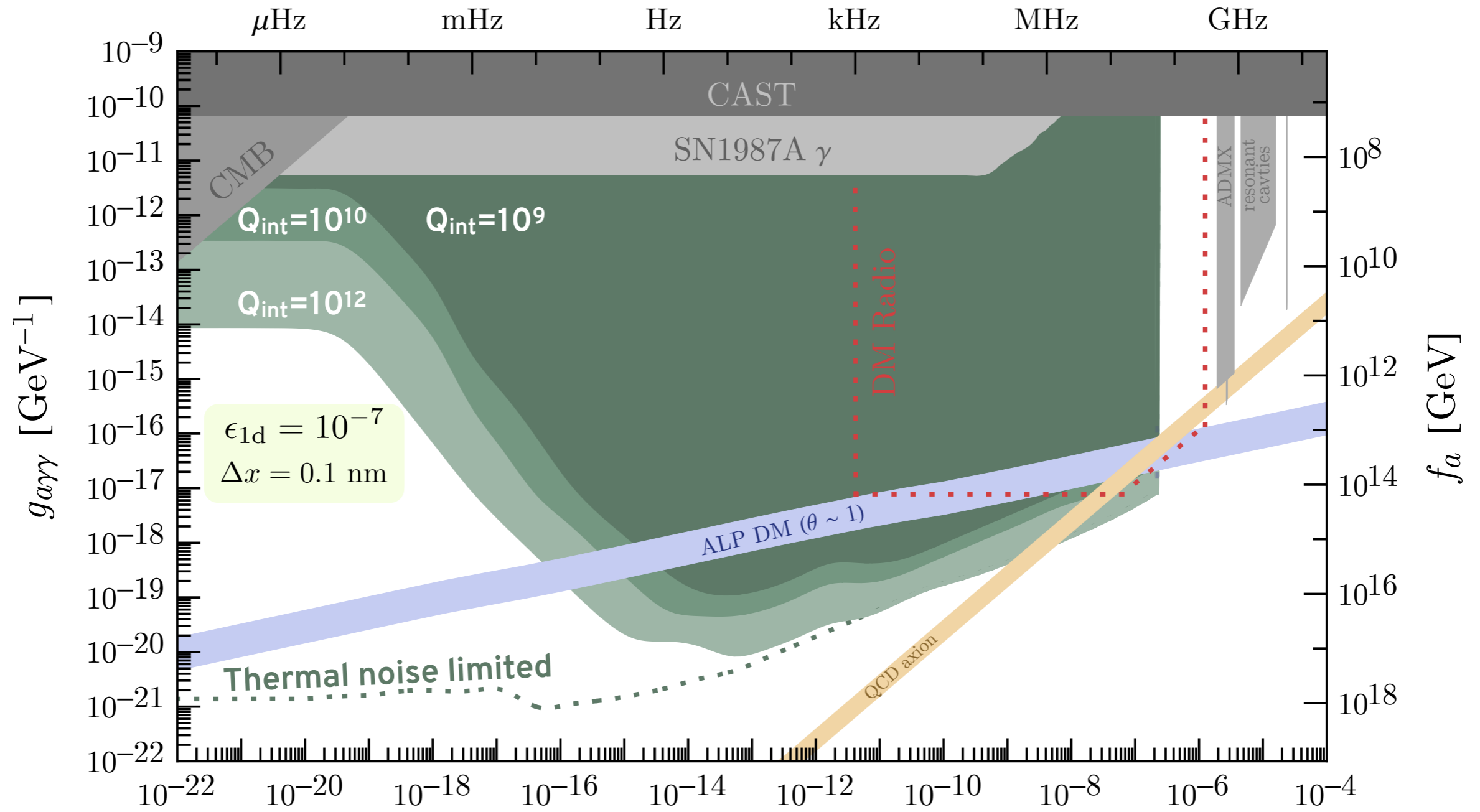
frequency =  $m_a/2\pi$



# Resonant parameter variations: $Q$ -factor

**B = 0.2 T, T = 2K,  $\omega_0 = 1$  GHz**

frequency =  $m_a/2\pi$

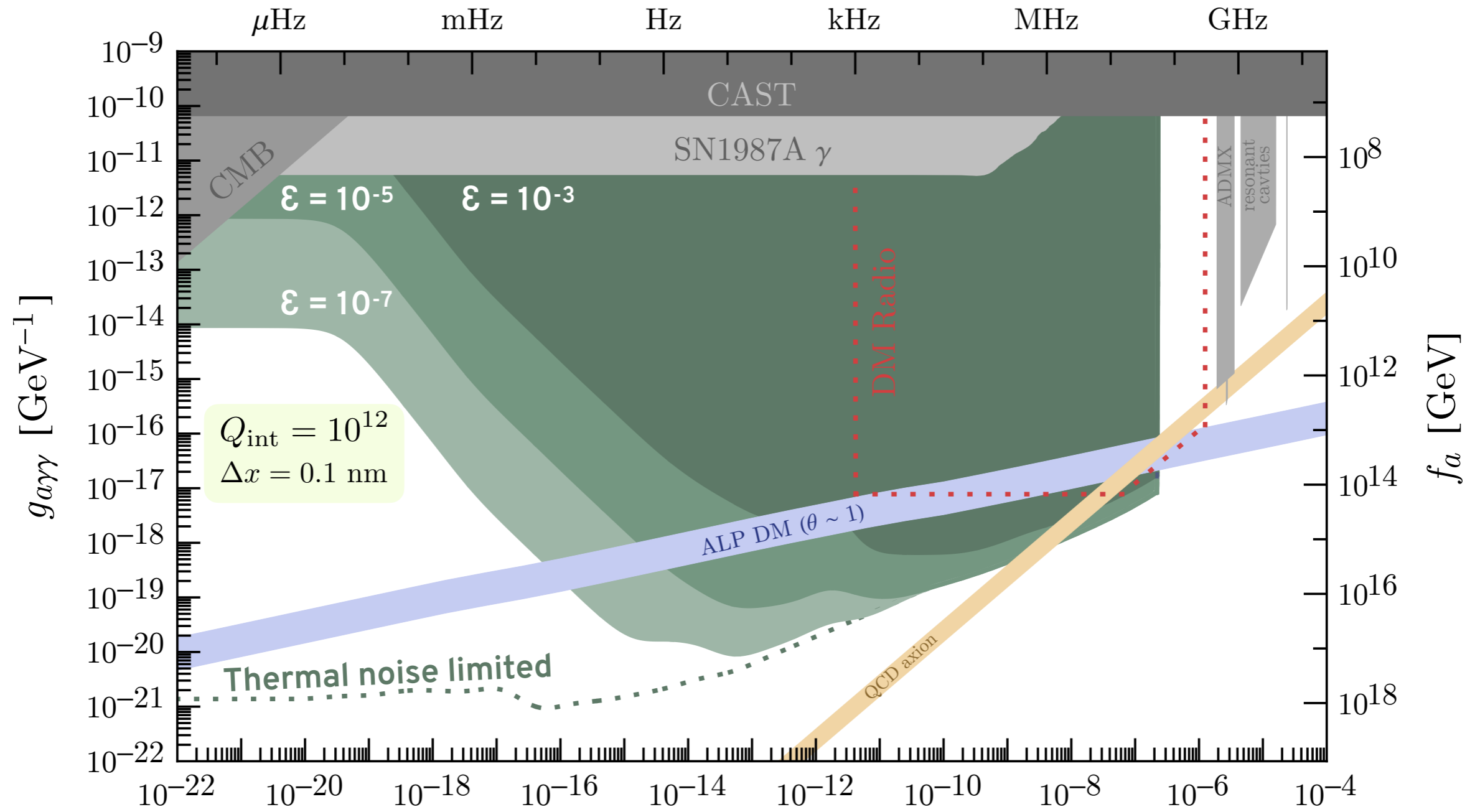




# Resonant parameter variations: *mode rejection*

**B = 0.2 T, T = 2K,  $\omega_0 = 1$  GHz**

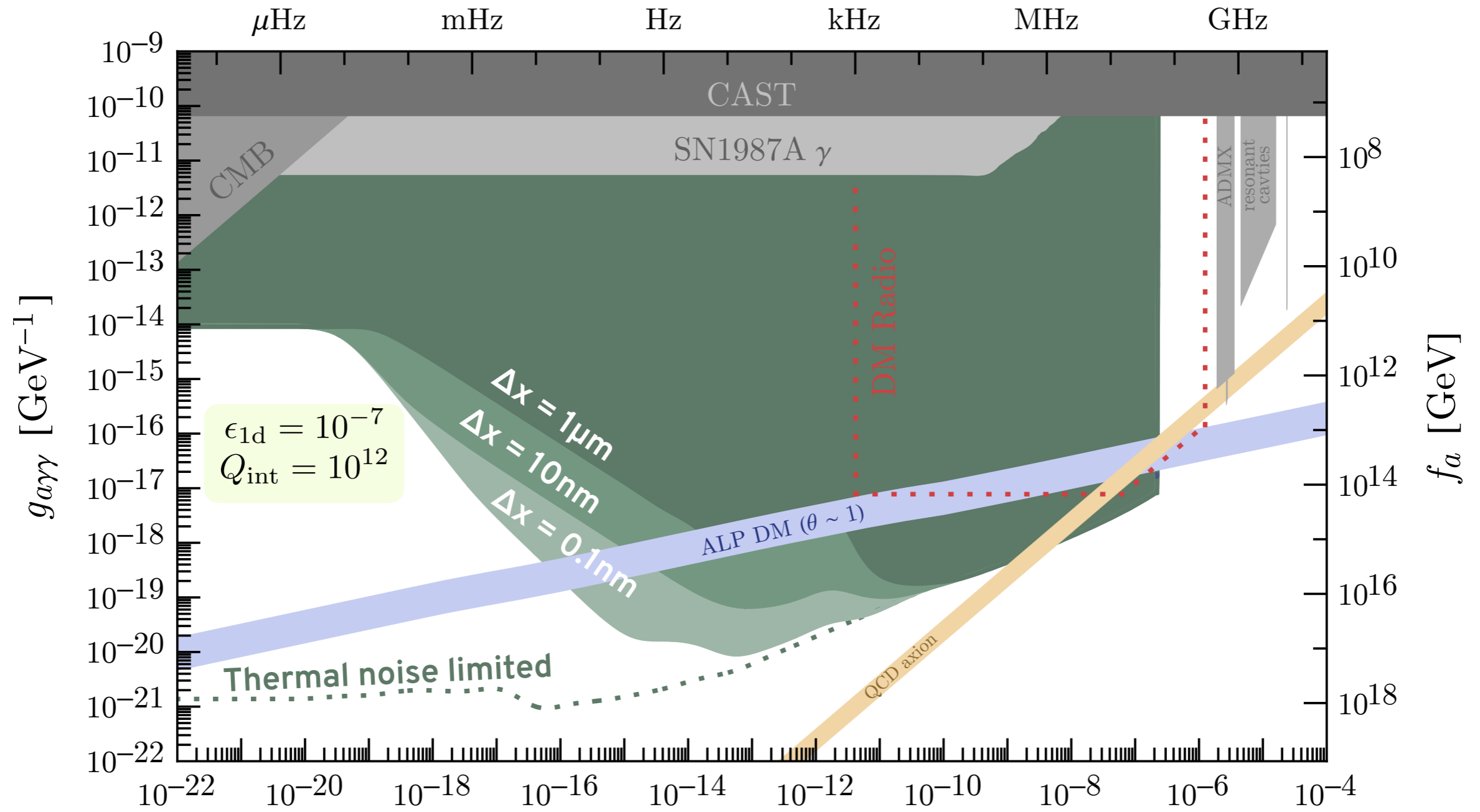
frequency =  $m_a/2\pi$



# Resonant parameter variations: *mode rejection*

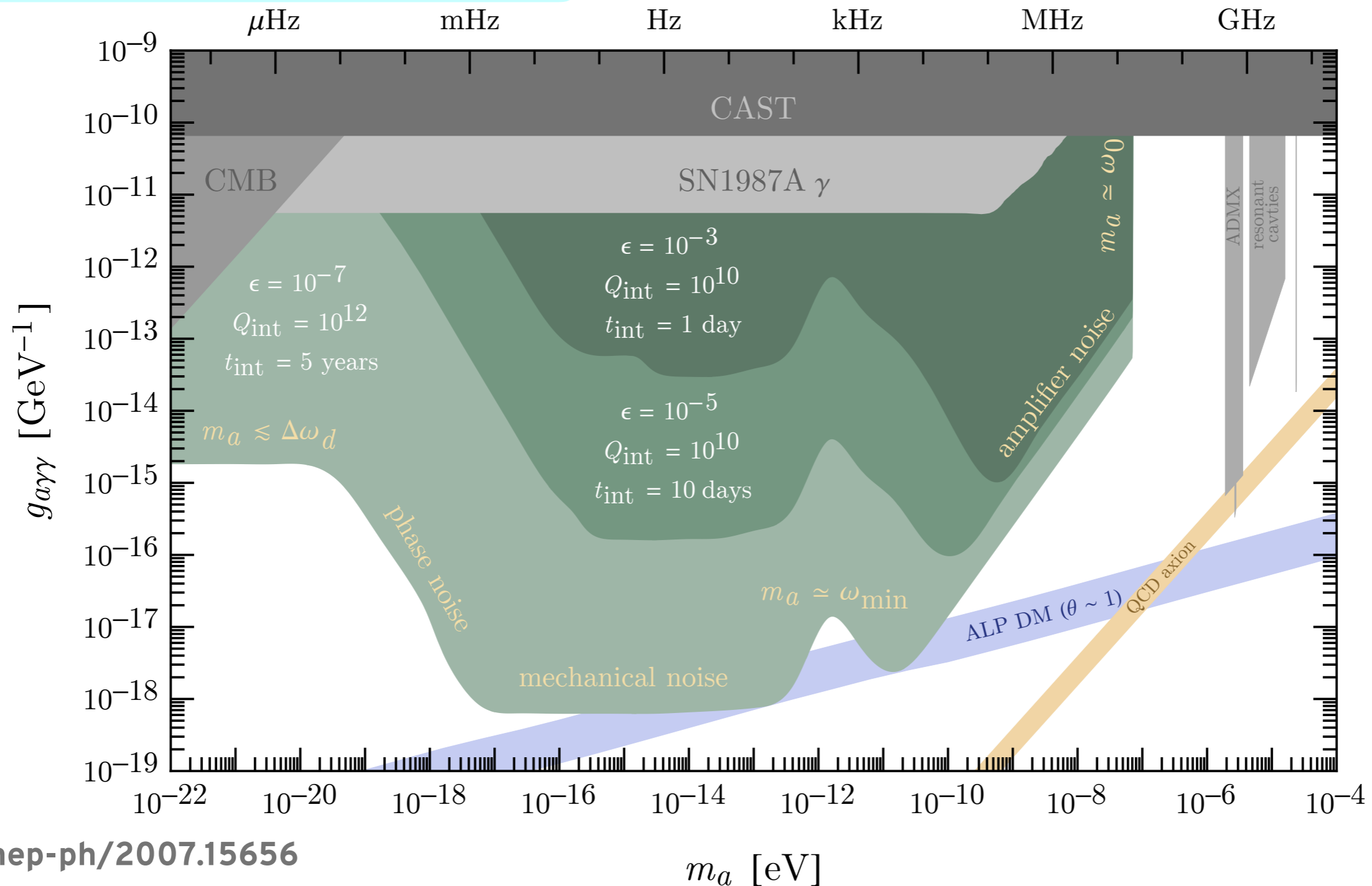
**B = 0.2 T, T = 2K,  $\omega_0 = 1$  GHz**

frequency =  $m_a/2\pi$



# Broadband Axion Resonant Frequency Conversion

**B = 0.2 T, T = 2K,  $\omega_0 = 100$  MHz** frequency =  $m_a/2\pi$



hep-ph/2007.15656

$m_a$  [eV]

# Outlook

## Radio-Frequency **up-conversion** approach

$$\omega_{\text{sig}} = \omega_0 \pm m_a$$

## Parametric gain for small axion masses vs. LC Resonator

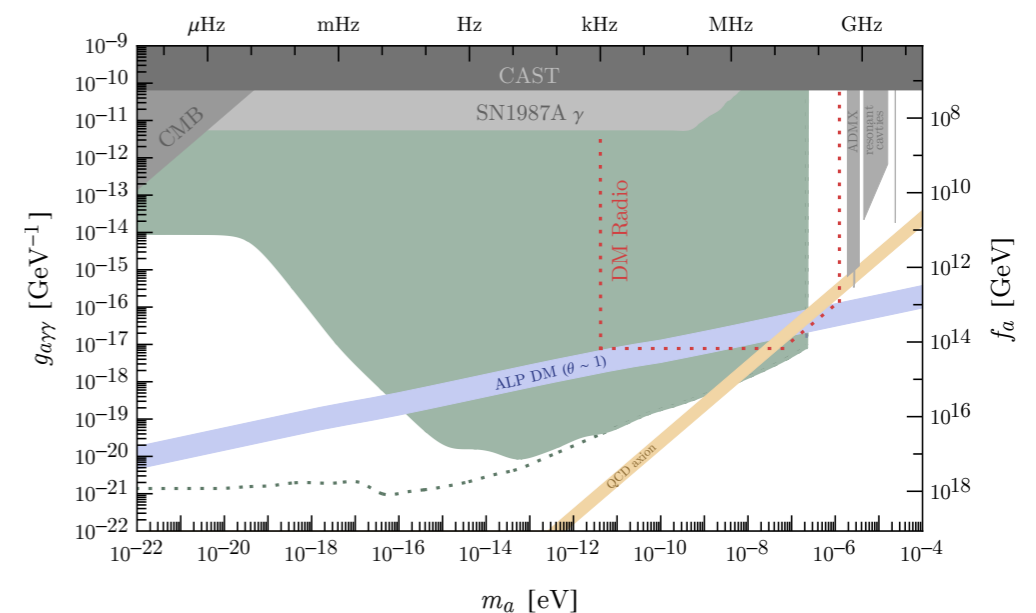
$$\frac{\text{SNR}}{\text{SNR}^{\text{LC}}} \sim \frac{\omega_0 \pm m_a}{m_a} \left( \frac{Q_{\text{int}}}{Q_{\text{LC}}} \right)^{1/2} \left( \frac{T_{\text{LC}}}{T} \right)^{1/2} \left( \frac{B_0}{B_{\text{LC}}} \right)^2$$

frequency =  $m_a/2\pi$

SLAC group seeking internal funding

In discussions for Physics Beyond Colliders @ CERN

Snowmass LOI CF2 & AF5



# Outlook

