

From PT Quantum Mechanics to PT Holography

Karl Landsteiner



Based on:

- **SciPost Phys. 9 (2020) no.3, 032,**
[arXiv:1912:06647]
with *Daniel Areán, Ignacio Salazar-Landea*
- **e-Print 2203.02524 [hep-th]**
with *Sergio Morales-Tejera*

QCD@work, Lecce 28. June, 2022

Resources

Some references:

- Seminal article:

Bender, Boettcher: Phys.Rev.Lett. 80 (1998) 5243-5246

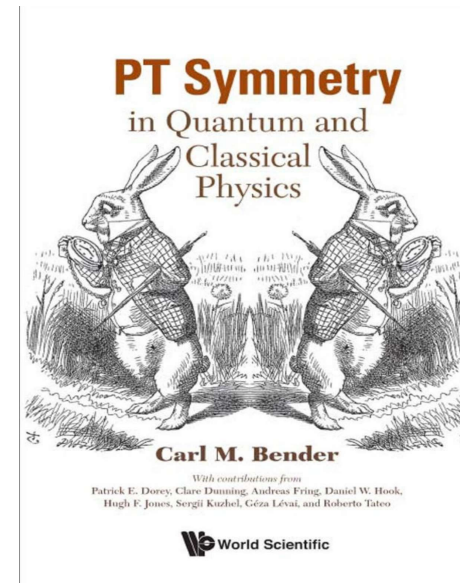
- Book:

Bender: “PT symmetry in quantum and classical physics” (2019)

- Time dependence:

- **Mostafazadeh:** Entropy 22 (2020) 4, 471 (**Review**)

- **Fring:** e-Print: 2201.05140 [quant-ph] (**Lecture notes**)

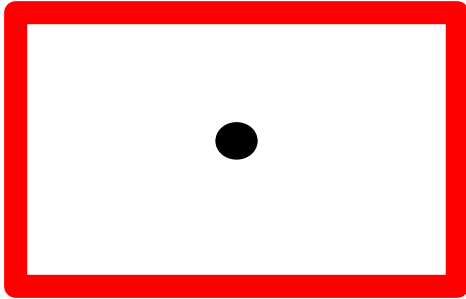


Outline

- PT Quantum Mechanics
- Holographic model
 - Static properties, PT phase transition
 - Dynamics: Hermitian vs. Non-Hermitian
- Conclusions

PT Quantum Mechanics

Particle in a Box



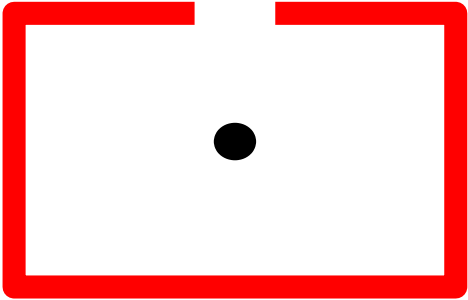
$$\Psi(t) = e^{-iEt}$$

$$H = E$$

$$p = |\Psi|^2 = 1$$

PT Quantum Mechanics

Particle in a **leaky** Box



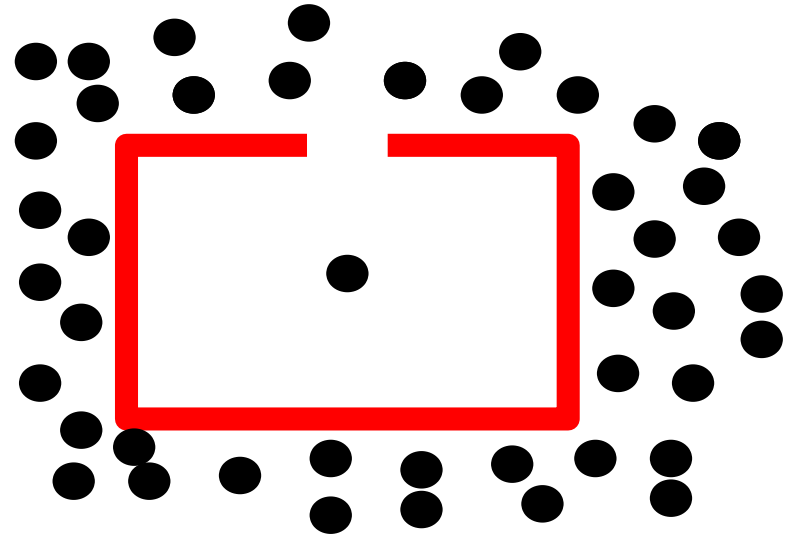
$$p(t) = e^{-2\Gamma t}$$

$$\text{“}\Psi(t)\text{”} = e^{-i(E-i\Gamma)t}$$

$$\text{“}H\text{”} = E - i\Gamma$$

PT Quantum Mechanics

Particle in a **leaky** Box in a reservoir



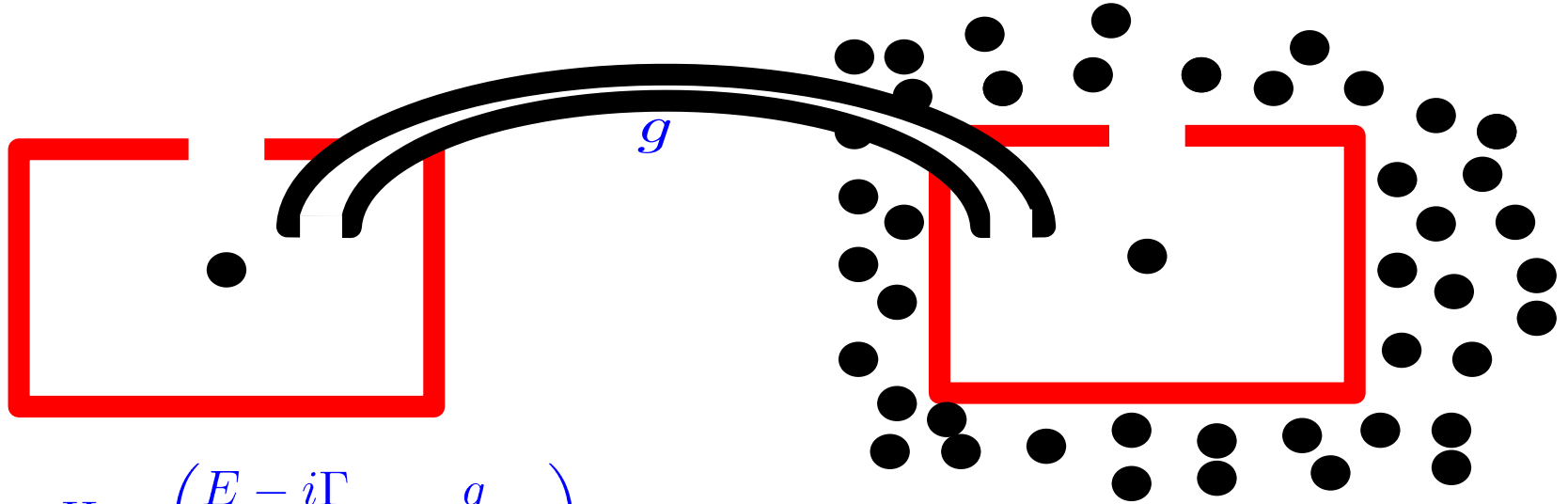
$$n(t) = e^{2\Gamma t}$$

$$\text{“}\Psi(t)\text{”} = e^{-i(E+i\Gamma)t}$$

$$\text{“}H\text{”} = E + i\Gamma$$

PT Quantum Mechanics

Particle in a **leaky** Box coupled to a particle in a Box immersed in a reservoir



$$H = \begin{pmatrix} E - i\Gamma & g \\ g & E + i\Gamma \end{pmatrix} \quad H^\dagger \neq H$$

$$\epsilon_{\pm} = E \pm \sqrt{g^2 - \Gamma^2} \quad \text{Real if } |g| \geq |\Gamma|$$

PT Hamiltonian = Open system with exact balance between source and sink

PT Quantum Mechanics

Time reversal: $\mathcal{T}(H) = H^* \quad i \rightarrow -i$

Parity: exchange the two boxes $\mathcal{P} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

$$\left. \begin{array}{l} \mathcal{T}(H) = H^* \quad i \rightarrow -i \\ \mathcal{P} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \end{array} \right\} \begin{array}{l} [\mathcal{PT}, H] = 0 \\ (\mathcal{PT})^2 = 1 \end{array}$$

$$(\mathcal{PT})^2 |n\rangle = \lambda_n^* \lambda_n |n\rangle \Rightarrow \lambda_n = e^{i\phi}$$

$$\epsilon_n^* = \epsilon_n$$

PT Symmetric phase $|g| \geq |\Gamma|$

$$\mathcal{PT}|n, +\rangle = |n, -\rangle$$

Or: $\epsilon_{n,+} = \epsilon_{n,-}^*$

PT broken phase $|g| < |\Gamma|$

PT Quantum Mechanics

Symmetry perspective: $h = E\mathbf{1}_2 + \vec{g} \cdot \vec{\sigma}$

Hermitian 2x2

Any two theories are related by SU(2) $D = e^{i\vec{\alpha} \cdot \frac{\vec{\sigma}}{2}}$

$$\vec{g} = \begin{pmatrix} g' \\ 0 \\ 0 \end{pmatrix} \quad D^\dagger h D = \begin{pmatrix} E - g' \sin(\alpha) & g' \cos(\alpha) \\ g' \cos(\alpha) & E + g' \sin(\alpha) \end{pmatrix}$$

Generate a non-Hermitian Matrix by setting $\alpha \rightarrow i\hat{\alpha}$

Equivalently: transform couplings: $\vec{g} \rightarrow O(\alpha) \cdot \vec{g}$

PT Quantum Mechanics

Symmetry perspective: $h = E\mathbf{1}_2 + \vec{g} \cdot \vec{\sigma}$

Complexify transformation: $\eta = e^{\hat{\alpha} \cdot \frac{\sigma_2}{2}}$ Hermitian

$$\vec{g} = \begin{pmatrix} g' \\ 0 \\ 0 \end{pmatrix} \quad H = \eta^{-1} h \eta = \begin{pmatrix} E - ig' \sinh(\hat{\alpha}) & g' \cosh(\hat{\alpha}) \\ g' \cosh(\hat{\alpha}) & E + g'i \sinh(\hat{\alpha}) \end{pmatrix}$$

$$g = g' \cosh(\hat{\alpha}) \quad , \quad \Gamma = g' \sinh(\hat{\alpha})$$

$$\epsilon_{\pm} = E \pm \sqrt{g'^2 (\cosh(\hat{\alpha})^2 - \sinh^2(\hat{\alpha}))} = E \pm g'$$

PT Quantum Mechanics

Pseudo-Hermitian Hamiltonians

$$H = \eta^{-1} h \eta, \quad h^\dagger = h, \quad \eta^\dagger = \eta$$

Schrödinger Equation: $i\partial_t|\psi\rangle = H|\psi\rangle$

Metric: η^2

Unitary time evolution: $\frac{d}{dt}\langle\Phi|\eta^2|\Psi\rangle = i\langle\Phi|H^\dagger\eta^2|\Psi\rangle - i\langle\Phi|\eta^2 H|\Psi\rangle = 0$

PT symmetric phase = pseudo-Hermitian Hamiltonian

PT Quantum Field Theory

$$\left. \begin{array}{l} \text{Dirac mass term } \mathcal{O}_1 = \bar{\psi}\psi \\ \text{Axial mass term } \mathcal{O}_2 = i\bar{\psi}\gamma_5\psi \end{array} \right\} \mathcal{O} = \mathcal{O}_1 + i\mathcal{O}_2 \text{ Has charge 2 under axial U(1) } e^{i\alpha\gamma_5}$$

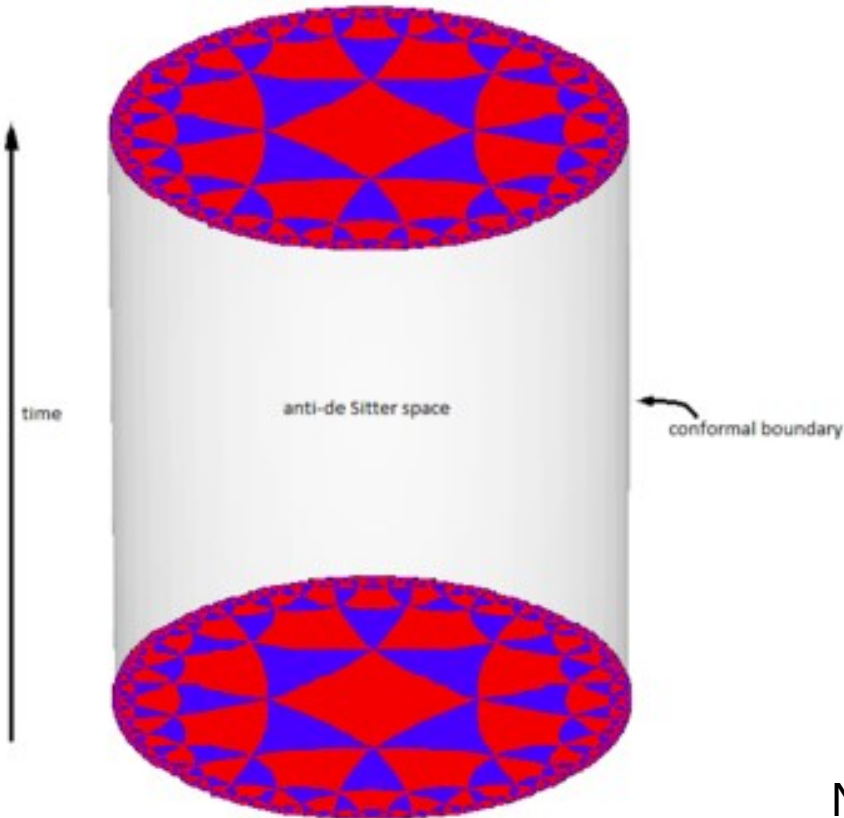
$$\text{Non-hermitian mass term } M_0\bar{\psi}\psi + M_5\bar{\psi}\gamma_5\psi \quad (Me^{2\hat{\alpha}}\bar{\mathcal{O}} + Me^{-2\hat{\alpha}}\mathcal{O})$$

$$\text{Spectrum } \omega = \sqrt{\vec{k}^2 + M_0^2 - M_5^2}$$

$$\text{Real for } |M_0| > |M_5|$$

Holography (Gauge/Gravity Duality)

Gravity in asymptotically AdS = QFT



Holographic Dictionary

Gravity	Quantum Field Theory
<i>Metric</i>	<i>Energy Momentum Tensor</i>
<i>Gauge field</i>	<i>Conserved current</i>
<i>Scalar field</i>	<i>Scalar operator</i>

Notice: Symmetry in QFT = gauge principle in AdS

PT-Holography

$$S = \int d^4x \sqrt{-g} \left[R - 2\Lambda - |D\phi|^2 - m^2|\phi|^2 - \frac{v}{2}|\phi|^4 - \frac{1}{4}F^2 \right]$$

PT = Non-Hermitian boundary conditions at $r = \infty$

$$\begin{aligned}\phi &\approx \frac{\phi_0}{r} + \frac{\langle \bar{\mathcal{O}} \rangle}{r^2} + \dots \\ \bar{\phi} &\approx \frac{\bar{\phi}_0}{r} + \frac{\langle \mathcal{O} \rangle}{r^2} + \dots\end{aligned}$$

Hermitian

$$\phi_0 = M'$$

$$\bar{\phi}_0 = M'$$



Non-Hermitian

$$\phi_0 = M' e^{\hat{\alpha}} = (1 - \xi)M$$

$$\bar{\phi}_0 = M' e^{-\hat{\alpha}} = (1 + \xi)M$$

PT-Holography

$$S = \int d^4x \sqrt{-g} \left[R - 2\Lambda - |D\phi|^2 - m^2|\phi|^2 - \frac{v}{2}|\phi|^4 - \frac{1}{4}F^2 \right]$$

How does PT work here: $\mathcal{T} : A \rightarrow -A$, $i \rightarrow -i$, $ds^2 \rightarrow ds^2$

$$\phi_R \rightarrow \phi_R \quad , \quad \phi_I \rightarrow -\phi_I$$

$$\mathcal{P} : A \rightarrow -A \quad , \quad ds^2 \rightarrow ds^2$$

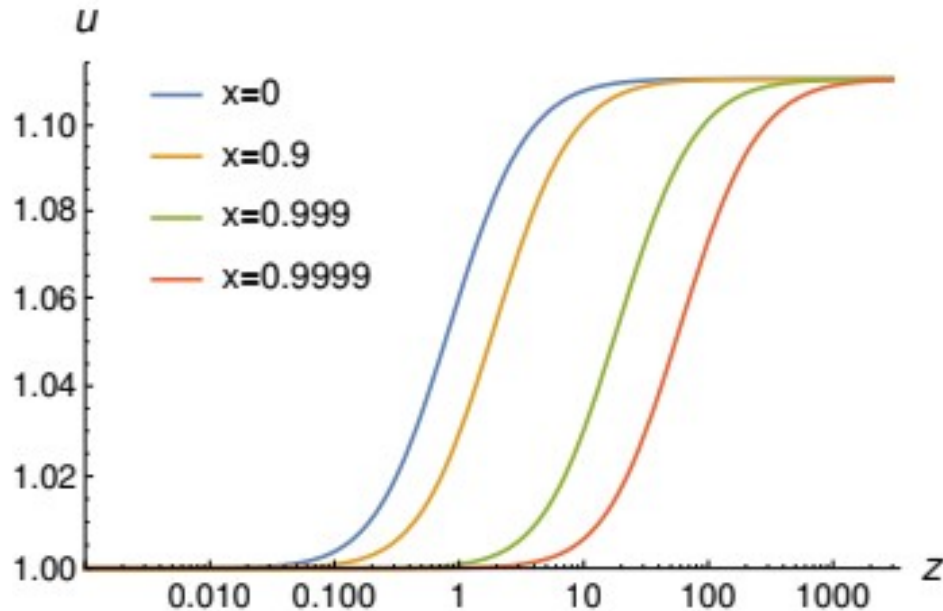
$$\phi_R \rightarrow \phi_R \quad , \quad \phi_I \rightarrow -\phi_I$$

PT leaves boundary conditions invariant: PT Hamiltonian

PT-Holography

T=0:

IR Solution:
$$\phi_{IR} = \sqrt{\frac{-m^2}{v} \frac{1 + \xi}{1 - \xi}}$$



- We find usual domain wall solutions
- Domain wall moves towards IR with ξ approaching 1
- For $|\xi| > 1$ no real solution!
PT phase transition
- At $\xi=1$ exceptional line: formally pure AdS
- Background is equivalent to hermitian solution with bcs
$$\phi_0 = \sqrt{1 - \xi^2 M'} \quad , \quad \bar{\phi}_0 = \sqrt{1 - \xi^2 M'}$$

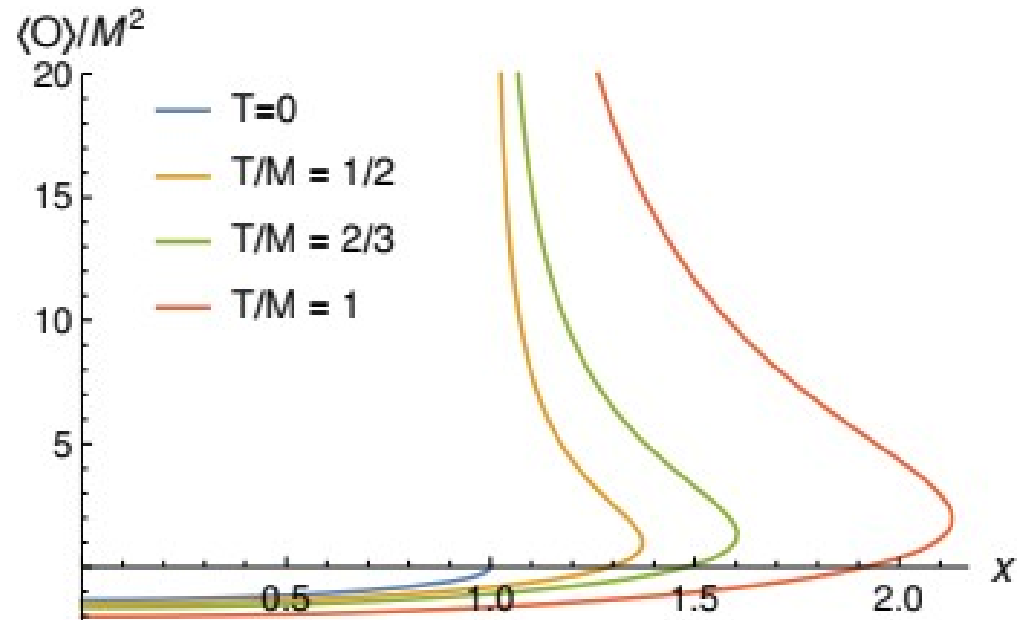
PT-Holography

$T > 0$:

Finite Temperature = black hole

$$u = u_0(r - r_H) + \dots$$

Surprisingly we find
real solutions for $|\xi| > 1$!



PT-Holography

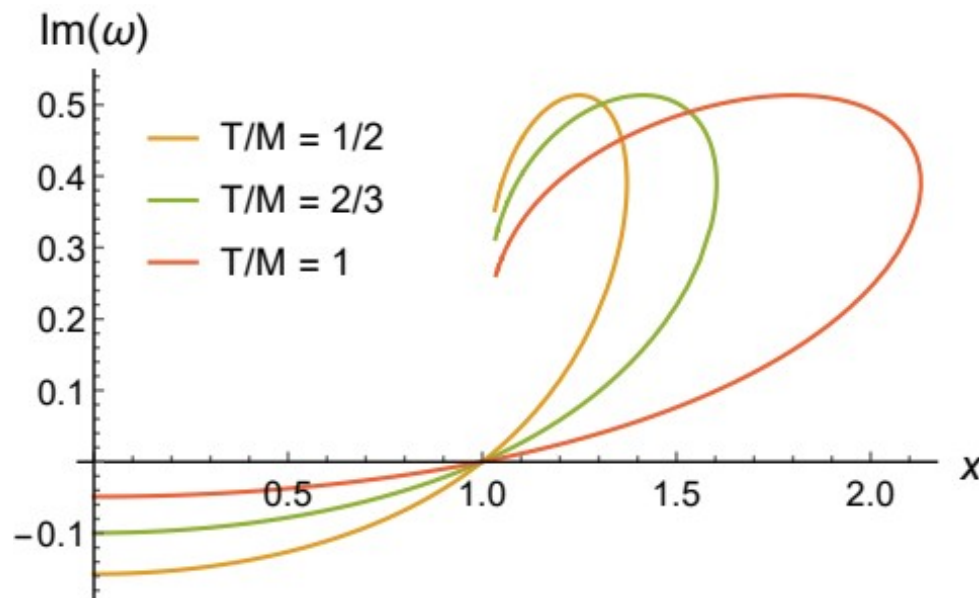
What is going on? Does finite temperature stabilize the system beyond the PT “quantum” phase transition even for $|\xi| > 1$?



NO!

Quasi-normal mode spectrum shows **instability** precisely **at $x=1$** (“pseudo”-diffusion mode),
No stable solution found

$$\omega = -i\Gamma(x)$$

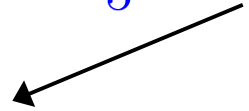


PT-Holography

QFT homework:

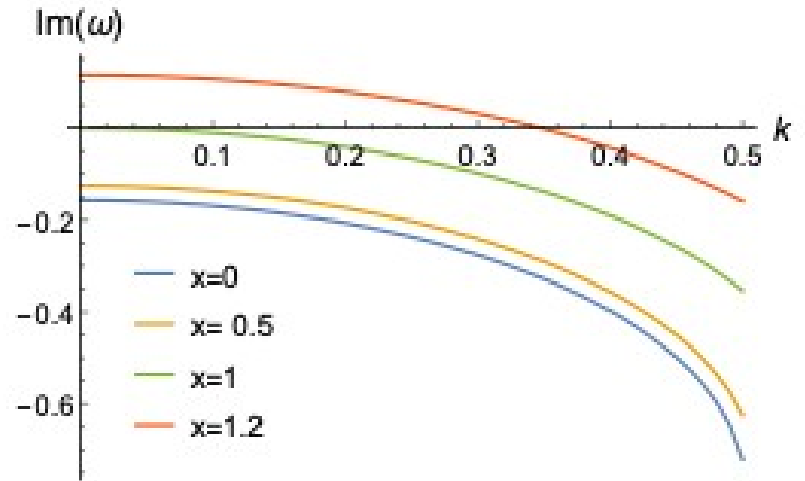
$$E = \sqrt{\vec{k}^2 + M^2 - M_5^2 + \delta M_T^2}$$

Thermal mass



Instability in (axial) density fluctuations

$$\langle \rho_5 \rho_5 \rangle \approx \frac{i\sigma k^2}{\omega + iDk^2 + i\Gamma}$$



PT-Holography

Time dependent non-hermitian coupling: $\frac{\partial}{\partial t} \langle \Psi | \Phi \rangle_\eta = 2 \langle \Psi | \eta^{-1} \partial_t \eta | \Phi \rangle$

[Mostafazadeh]
[Fring]

Non-unitary time evolution, what to do?

- 2 Options:
- Just live with it
 - Repair unitarity in some way

How to repair? With a “gauge” field !

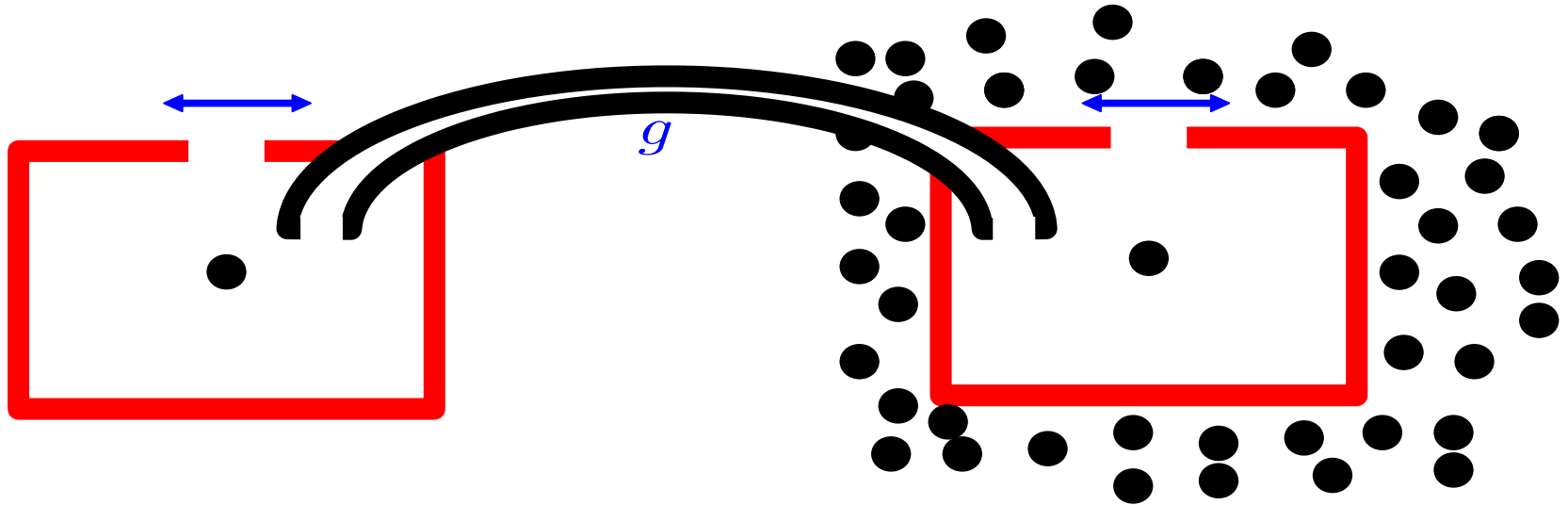
$$i(\partial_t - iA_t)|\Psi\rangle = H|\Psi\rangle$$

$$H = \eta^{-1}(t)h\eta(t) \quad , \quad A_t = i\eta^{-1}\partial_t\eta$$

[Chernodub, Millington]

PT Quenches

PT Hamiltonian = Open system with exact balance between source and sink



- Non-unitary: change openings simultaneously and in the same way
- Unitary: change openings in the same way but also do something more, corresponding to the gauge field (forward vs backward hopping ...)

$$h = g\sigma_x \quad , \quad \eta = e^{\hat{\alpha}\sigma_y} \quad , \quad A_t = i\partial_t\hat{\alpha}\sigma_y$$

PT Quenches

In Holography this can be done explicitly, time evolution can be calculated in both cases!

Time dependent boundary conditions:

$$\phi = [1 - \xi(t)]M$$
$$\bar{\phi} = [1 + \xi(t)]M$$

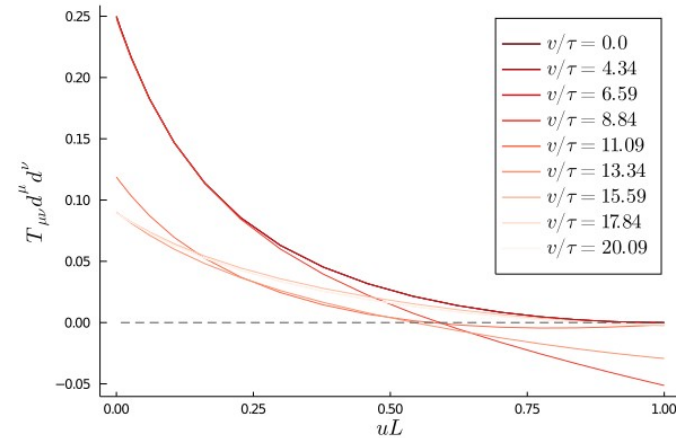
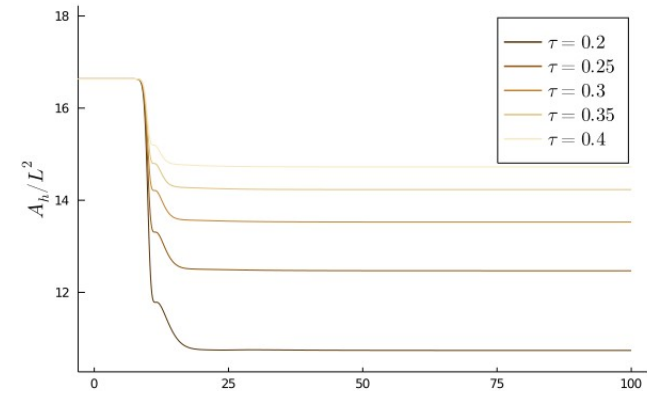
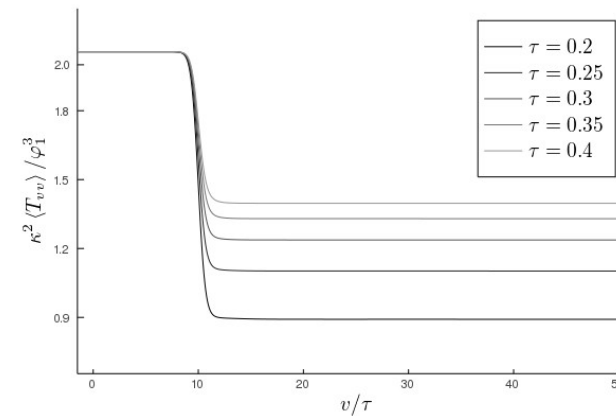
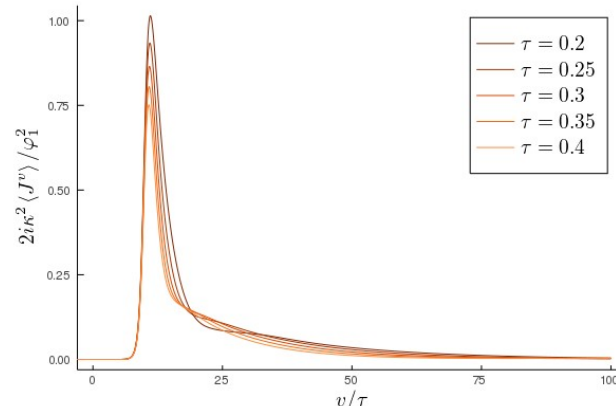
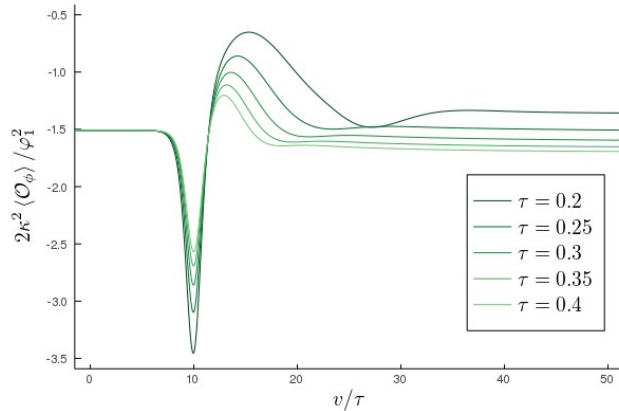
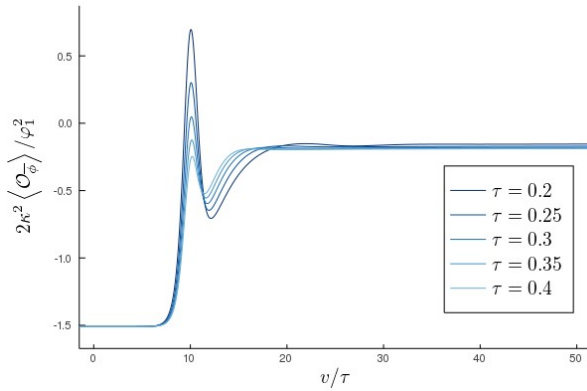
For unitary evolution introduce gauge field:

$$A_t = i \frac{\partial_t \xi}{1 - \xi^2}$$

Note: in the unitary time evolution formalism we can not reach the exceptional point or go into the PT broken regime! Singularity in the gauge field.

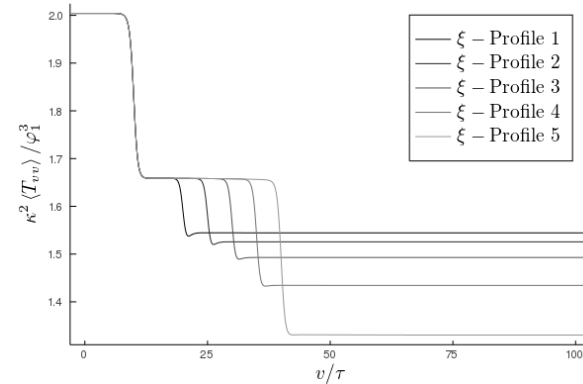
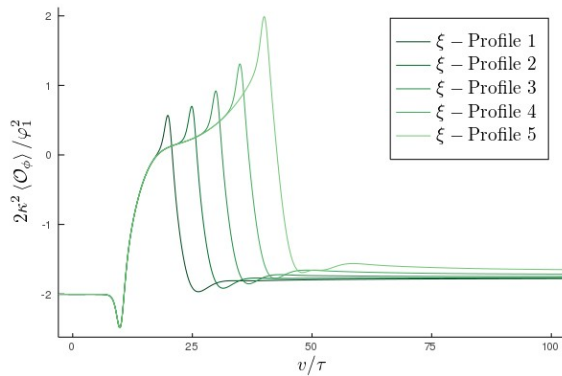
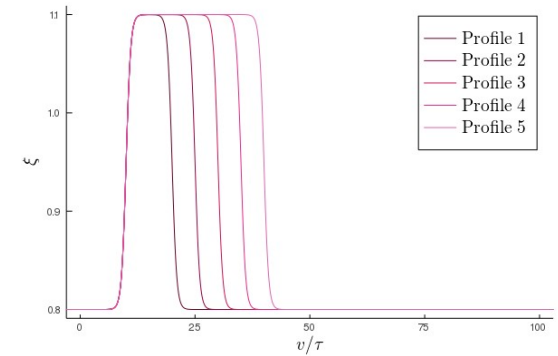
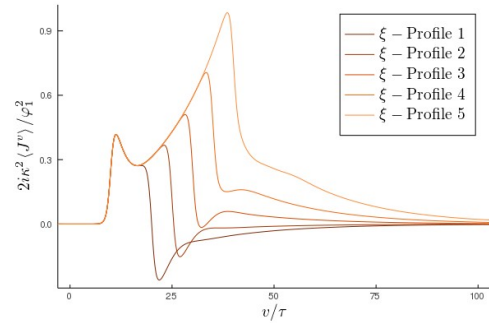
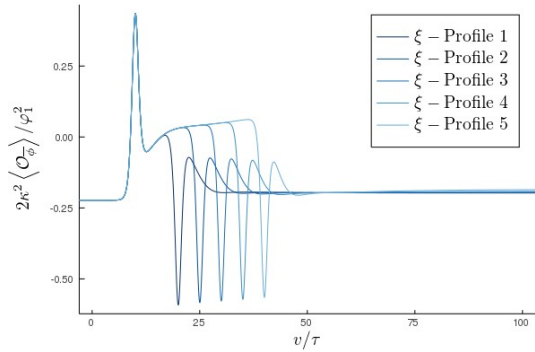
PT Holographic quench

Non-unitary evolution I: $\xi = 0 \rightarrow \xi = 0.8$



PT Holographic quench

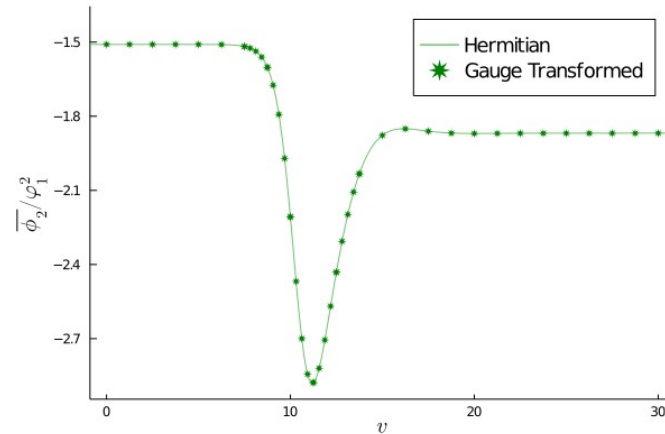
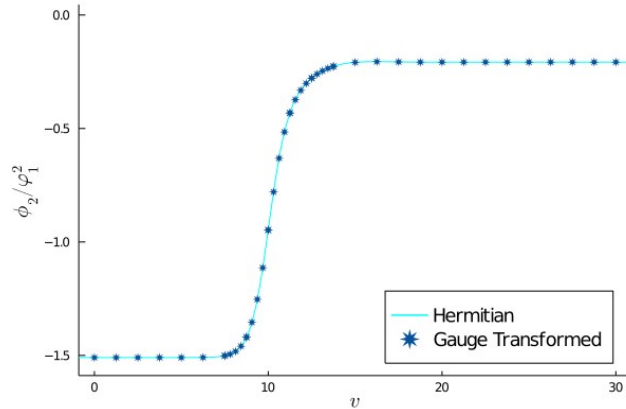
Non-unitary evolution III: $\xi = 0.8 \rightarrow \xi = 0.8$ With excursion into PT broken regime!



PT Holographic quench

Unitary evolution: scalars as before but also gauge field $A_t(r = \infty) = i \frac{\partial_t \xi}{1 - \xi^2}$

Compare to Hermitian description: $\phi = \sqrt{1 - \xi(t)^2} \phi(0)$



Summary

- We have constructed the first example of PT holographic model
- PT phase transition exists at $\xi=1$ just as in PT QM
- Predictions for PT quantum field theory
- PT-quantum quenches show very interesting phenomenology
- Unitary with gauge field equivalent to Hermitian theory
- If \exists holographic material \rightarrow violate energy conditions in Laboratory?

THANK YOU!