Light meson spectroscopy and gluonium searches in η_c and $\Upsilon(1S)$ decays at BaBar

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Outline:

Study of $\Upsilon(1S)$ radiative decays to $\gamma\pi^+\pi^-$ and γK^+K^-

Light meson spectroscopy in three-body η_c decays

8th Workshop on Theory, Phenomenology and Experiments in Flavour Physics,
Anacapri, Island of Capri (Napoli), Italy, June 11–13, 2022.

Introduction: Gluonium physics in radiative $\Upsilon(1S)$ decays

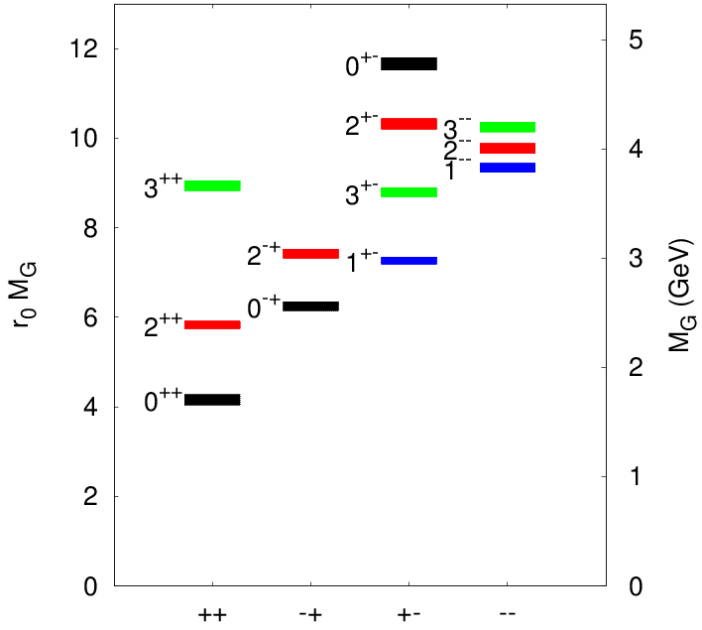
- The search for gluonium states is still a hot topic for QCD.
- Glueball spectrum from Lattice QCD.

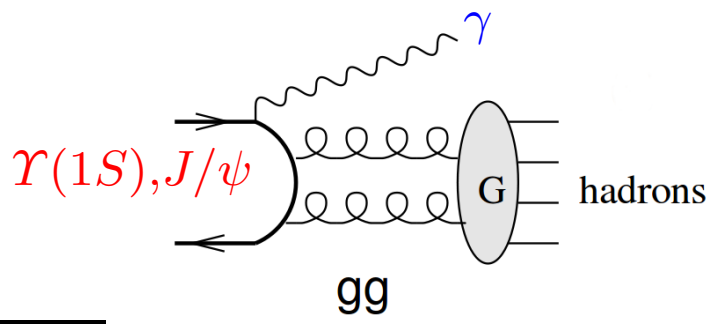
(Y. Chen et al. Phys. Rev. D **73**, *014516 (2006))*

- The $J^{PC} = 0^{++}$ glueball is expected in the mass region between 1.5 and 2.0 GeV.
- Intensive work done in the past using $\bar{p}p$ annihilations, central production, J/ψ radiative decays, ...
- Difficult problem, also because gluonium states could easily mix with standard $q\bar{q}$ mesons.
- Proposed scalar gluonium candidates are:

$f_0(500)$, $f_0(1370)$, $f_0(1500)$, $f_0(1710)$, $f_0(1865)$

- A similar work as for J/ψ decays could be done in $\Upsilon(1S)$ radiative decays.
- Challenging: radiative $\Upsilon(1S)$ decays branching fractions expected to be suppressed by a factor 25 with respect to the corresponding J/ψ branching fractions.





Gluonium searches in charmonium decays

□ Charmonium decays can be used to obtain new information on light meson spectroscopy.

□ The η_c resonance is strongly coupled to scalar mesons.

□ New information can be obtained on the scalar mesons properties and the identification of the scalar glueball.

□ The η' meson is supposed to contain a significant gluonic contribution.

(S. D. Bass and P. Moskal, Rev.Mod.Phys.91, 015003 (2019))

□ It is of interest to compare η_c decays to:

$\eta_c \to \eta$ (scalar-meson)
$\eta_c \to \eta'$ (scalar-meson)

where the scalar meson decays to $\pi^+\pi^-$ or K^+K^-.

□ Compare with results from J/ψ radiative decays.

$\Upsilon(1S)$ Analysis Strategy

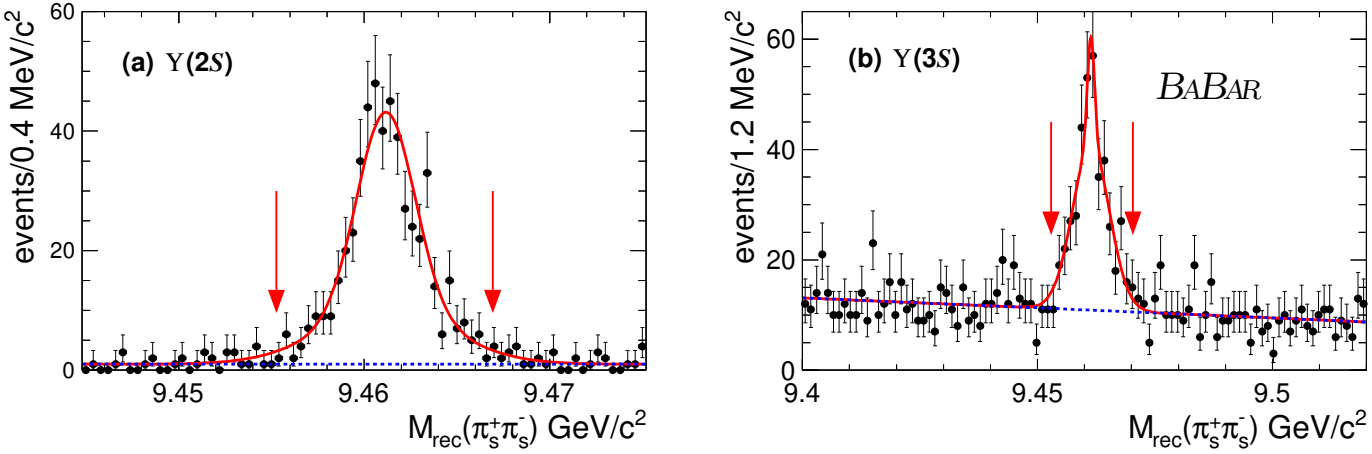
□ In the present analysis we make use of $\Upsilon(2S)$ and $\Upsilon(3S)$ decays with integrated luminosities of 13.6 and 28.0 fb^{-1} *(J.P. Lees et.al., Phys.Rev. D97 (2018) no.11, 112006.)*

□ We reconstruct the decay chains:

$$\Upsilon(2S)/\Upsilon(3S) \to \pi_s^+\pi_s^-\Upsilon(1S) \quad \begin{array}{l} \to \gamma\pi^+\pi^- \\ \to \gamma K^+K^- \end{array}$$

□ Require momentum balance and compute the recoiling mass.

$$M^2_{\rm rec}(\pi_s^+\pi_s^-) = |p_{e+} + p_{e-} - p_{\pi_s^+} - p_{\pi_s^-}|^2,$$

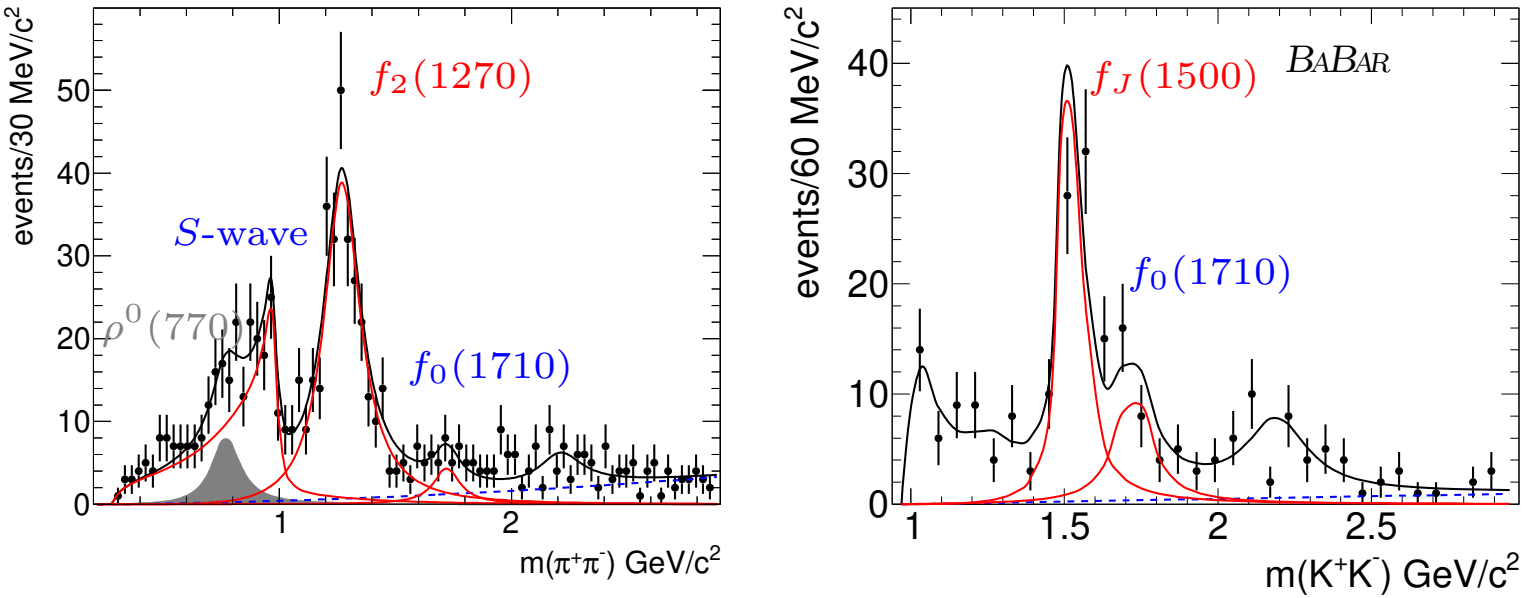


□ Require: $|M^2_{\rm rec}(\pi_s^+\pi_s^-) - m(\Upsilon(1S))_f| < 2.5\sigma$

□ Select the $\Upsilon(1S)$: $9.1\ {\rm GeV}/c^2 < m(\gamma h^+h^-) < 9.6\ {\rm GeV}/c^2$

$\Upsilon(2S)/\Upsilon(3S)$ Combined $\pi^+\pi^-$ and K^+K^- mass spectra

- Observe rich resonance production.



- The total S-wave is described by a coherent sum of $f_0(500)$ and $f_0(980)$ as:

$$S\text{-wave} = \mid BW_{f_0(500)}(m) + cBW_{f_0(980)}(m)e^{i\phi} \mid^2 .$$

- The $f_0(980)$ is described by a coupled channel Breit-Wigner.

- For $f_0(500)$ we obtain:

$$m(f_0(500)) = 0.856 \pm 0.086 \text{ GeV}/c^2, \Gamma(f_0(500)) = 1.279 \pm 0.324 \text{ GeV}, \ \phi = 2.41 \pm 0.43 \ rad$$

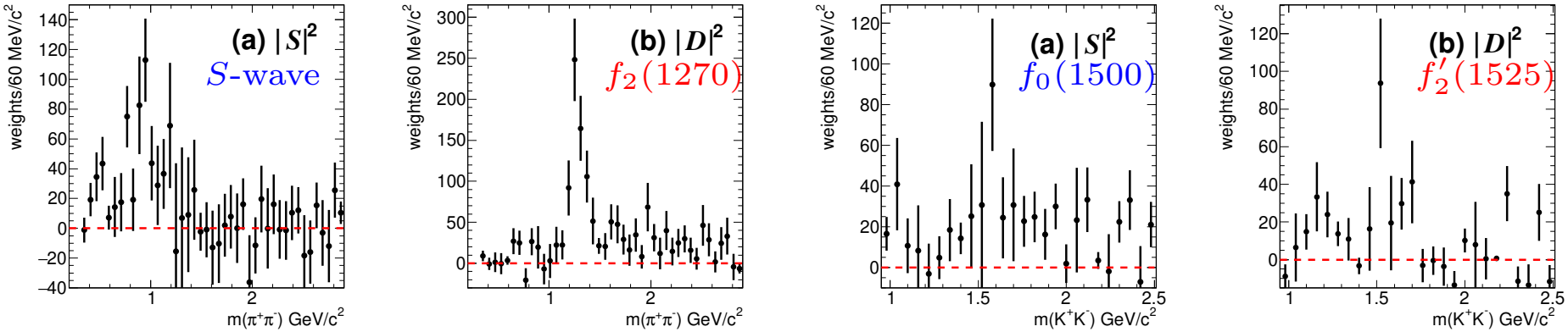
- Contaminations from a $\rho^0(770) \to \pi^+\pi^-$ in the $\Upsilon(3S)$ data to $\gamma\pi^+\pi^-$.

A Simple Partial Wave Analysis.

□ In a simplified procedure, the Y_L^0 moments are related to the S and D waves by the system of equations:

$$\sqrt{4\pi}\langle Y_0^0\rangle = S^2 + D^2,$$
$$\sqrt{4\pi}\langle Y_2^0\rangle = 2SD\cos\phi_{SD} + 0.639D^2,$$
$$\sqrt{4\pi}\langle Y_4^0\rangle = 0.857D^2,$$

□ The system can be solved directly for S and D waves:



□ We obtain an estimate of the S-wave$\to\pi^+\pi^-$ yield

$$N(S-\text{wave}) = 629 \pm 128,$$

□ The K^+K^- mass spectrum shows evidence for both $f_0(1500)$ and $f_2'(1525)$.

Measured $\Upsilon(1S)\to\gamma R$ branching fractions

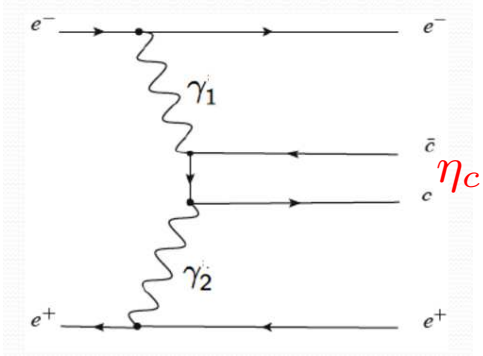
□ We label with $f_J(1500)$ the total enhancement in the 1500 MeV mass region.

Resonance	$\mathcal{B}(10^{-5})$ (*BABAR*)	CLEO (Phys.Rev. D73 (2006) 032001)
$\pi\pi$ S-wave	$4.63 \pm 0.56 \pm 0.48$	$(f_0(980))$ $1.8^{+0.8}_{-0.7} \pm 0.1$
$f_2(1270)$	$10.15 \pm 0.59\ ^{+0.54}_{-0.43}$	$10.2 \pm 0.8 \pm 0.7$
$f_0(1710)\to\pi\pi$	$0.79 \pm 0.26 \pm 0.17$	
$f_J(1500)\to K\bar{K}$	$3.97 \pm 0.52 \pm 0.55$	$3.7^{+0.9}_{-0.7} \pm 0.8$
$f_2'(1525)$	$2.13 \pm 0.28 \pm 0.72$	
$f_0(1500)\to K\bar{K}$	$2.08 \pm 0.27 \pm 0.65$	
$f_0(1710)\to K\bar{K}$	$2.02 \pm 0.51 \pm 0.35$	$0.76 \pm 0.32 \pm 0.08$

□ Theoretical expectations for gluonium assignments of $f_0(1710)$ *(R. Zhu, JHEP* **1509***, 166 (2015))*, $f_0(1500)\to K\bar{K}$, *(X. G. He et al.,Phys. Rev. D* **66***, 074015 (2002))* and $f_0(1370)$ *(R. Zhu, JHEP* **1509***, 166 (2015))* are in the range of our meaurements.

η_c production in two-photon interactions

□ In two-photon interactions we select events in which the e^+ and e^- beam particles are scattered at small angles and remain undetected.



□ Only resonances with $J^{PC} = 0^{\pm+}, 2^{\pm+}, 4^{\pm+}$ can be produced.

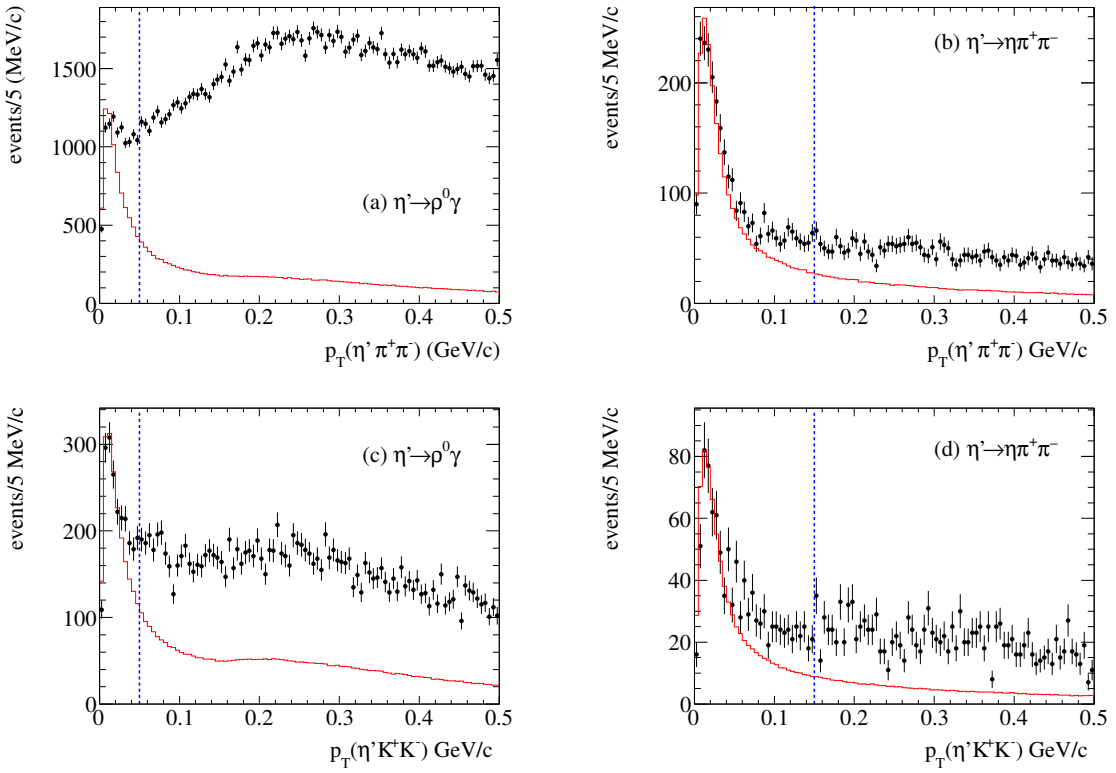
□ We have studied the following final states.

- □ $\eta_c \to \eta' \pi^+ \pi^-$, with $\eta' \to \rho^0 \gamma$ and $\eta' \to \eta \pi^+ \pi^-$ $(\eta \to \gamma\gamma)$.
- □ $\eta_c \to \eta' K^+ K^-$.
- □ $\eta_c \to \eta \pi^+ \pi^-$ with $\eta \to \gamma\gamma$ and $\eta \to \pi^+ \pi^- \pi^0$.

□ Dataset: 519 fb^{-1} recorded with the *BABAR* detector at center-of-mass energies at and near the $\Upsilon(nS)$ $(n = 2, 3, 4)$ resonances.

Two-photon signals for $\eta' h^+ h^-$

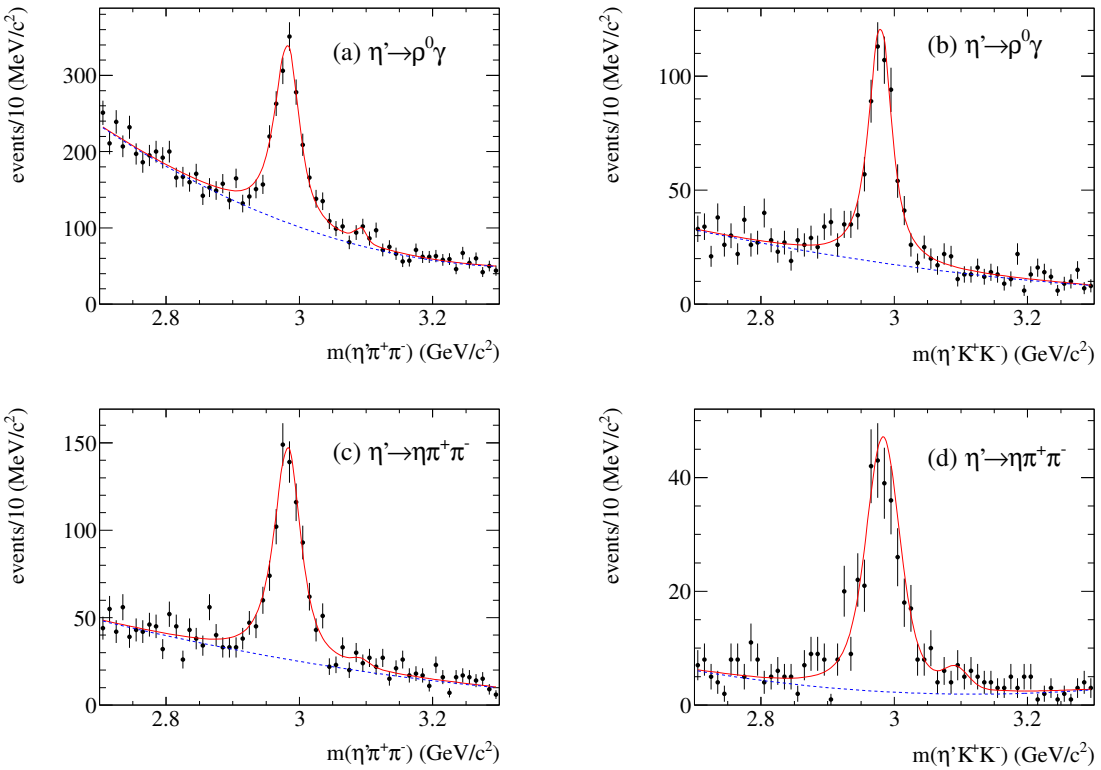
- Removed Initial-State-Radiation $e^+e^- \to \gamma_{ISR}(\eta' h^+ h^-)$ and $\gamma\gamma \to 4h$ events.
- Two-photon events are isolated from the balance of the transverse momentum p_T with respect to the e^+e^- direction.
- p_T distributions for $m(\eta' h^+ h^-) > 2.7$ GeV (charmonium region).



- In red are signal MC simulations normalized to the threshold peak.
- p_T selection optimized on the η_c signal.

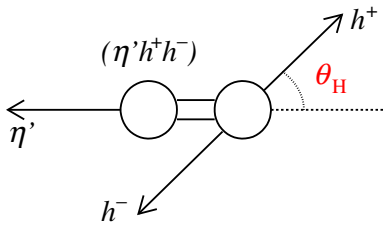
η_c signals for $\eta_c \to \eta' h^+ h^-$

- Mass resolution modeled as the sum of a Gaussian and a Crystal Ball function.
- η_c described by a Breit-Wigner convolved with the experimental resolution.
- Binned χ^2 fits. η_c parameters fixed to PDG values.
- When left free in the fit, parameters consistent with PDG averages.
- Background described by a 2^{nd} order polynomial.



- Included a J/ψ contribution from residual ISR background.
- **First observation of $\eta_c \to \eta' K^+ K^-$.**

Efficiency



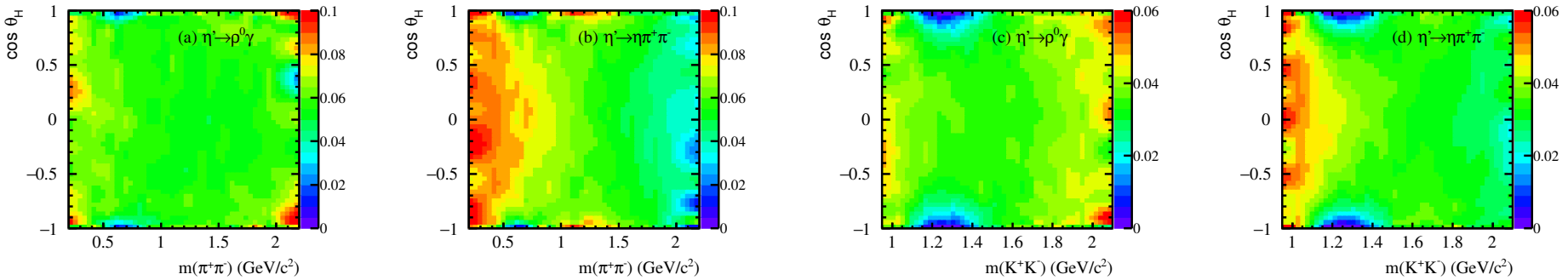
- We compute the ratio of branching fractions as:

$$\mathcal{R} = \frac{\mathcal{B}(\eta_c \to \eta' K^+ K^-)}{\mathcal{B}(\eta_c \to \eta' \pi^+ \pi^-)} = \frac{N_{\eta' K^+ K^-}}{N_{\eta' \pi^+ \pi^-}} \frac{\epsilon_{\eta' \pi^+ \pi^-}}{\epsilon_{\eta' K^+ K^-}}$$

for each η' decay mode.

- $\epsilon_{\eta' K^+ K^-}$ and $\epsilon_{\eta' \pi^+ \pi^-}$ are the corresponding weighted efficiencies.
- The efficiency is projected on the helicity angle $cos\theta_H$ vs. $m(h^+h^-)$.
- Fitted using Legendre polynomials in $cos\theta_H$ and interpolated along $m(h^+h^-)$.

$\eta' \pi^+ \pi^-$ $\qquad\qquad$ $\eta' K^+ K^-$



- Depletions close to $|cos\theta_H| \approx 1$ due to not reconstructed low momentum particles.
- For a given η' decay mode we apply the same selections to the numerator and denominator, except for particle identification.

Branching fractions

- Each event is weighted by the inverse of the efficiency.
- Background subtraction performed by giving positive weights in the η_c signal region and negative weights in the η_c sidebands.

Final state	yield
$\eta_c \to \eta' \pi^+ \pi^-$ $(\eta' \to \rho^0 \gamma)$	$1160 \pm 57 \pm 47$
$\eta_c \to \eta' K^+ K^-$	$473 \pm 29 \pm 3$
$\eta_c \to \eta' \pi^+ \pi^-$ $(\eta' \to \pi^+ \pi^- \eta)$	$619 \pm 35 \pm 11$
$\eta_c \to \eta' K^+ K^-$	$249 \pm 20 \pm 11$

- We obtain the following values of the branching fractions

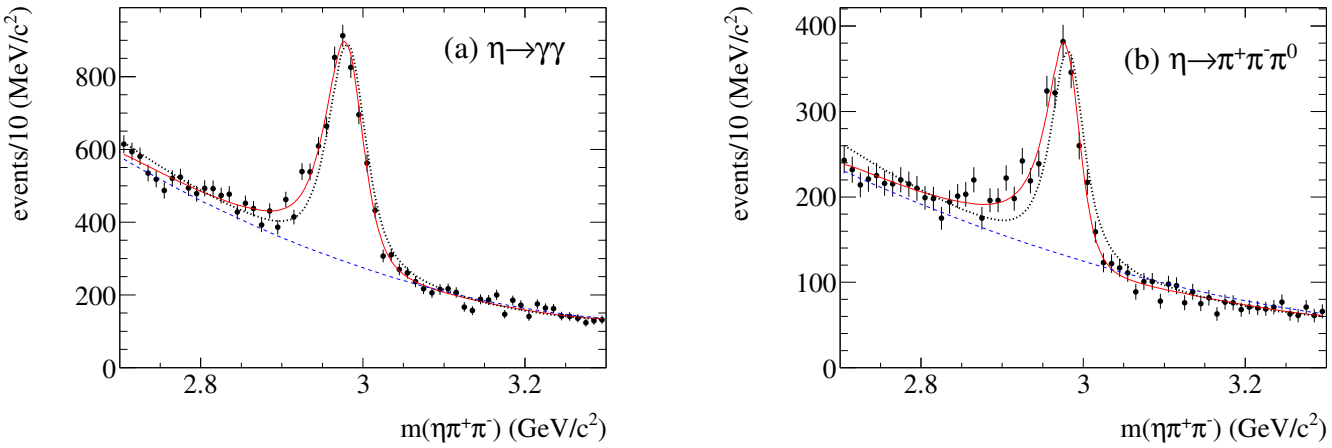
$\mathcal{R}_1(\rho^0\gamma) = 0.629 \pm 0.049_{\rm stat} \pm 0.035_{\rm sys}$ and $\mathcal{R}_2(\eta\pi^+\pi^-) = 0.672 \pm 0.066_{\rm stat} \pm 0.078_{\rm sys}$

- Average:

$$\mathcal{R} = \frac{\mathcal{B}(\eta_c \to \eta' K^+ K^-)}{\mathcal{B}(\eta_c \to \eta' \pi^+ \pi^-)} = 0.644 \pm 0.039_{\rm stat} \pm 0.032_{\rm sys}.$$

Study of $\eta_c \to \eta\pi^+\pi^-$

□ $\eta\pi^+\pi^-$ mass spectra for $\eta\to\gamma\gamma$ and $\eta\to\pi^+\pi^-\pi^0$.



□ The fits to the $\eta\pi^+\pi^-$ mass spectra return η_c masses shifted down by ≈ 10 MeV.

□ Fix the η_c parameters and introduce interference between the η_c and the two-photon background

$$f(m) = |A_{nres}|^2 + |A_{res}|^2 + c \cdot 2Re(A_{nres}A^*_{res})$$

where $|A_{nres}|^2$ is described by a 2^{nd} order polynomial, $A_{res} = \alpha \cdot BW(m) \cdot exp(i\phi)$.

□ We obtain $\chi^2/ndf = 77/54(160/55)$ and $\chi^2/ndf = 46/54(139/55)$ for the two η decay modes, with and without interference.

□ **No evidence for interference effects is observed for other η_c to η or η' decay modes.**

Dalitz plot analysis.

□ We perform a Dalitz plot analysis of $\eta_c \to \eta' \pi^+ \pi^-$, $\eta_c \to \eta' K^+ K^-$ and $\eta_c \to \eta \pi^+ \pi^-$ in the η_c mass region using unbinned maximum likelihood fits.

□ Efficiencies, normalization integrals and purities are computed separately for the two contributing final states.

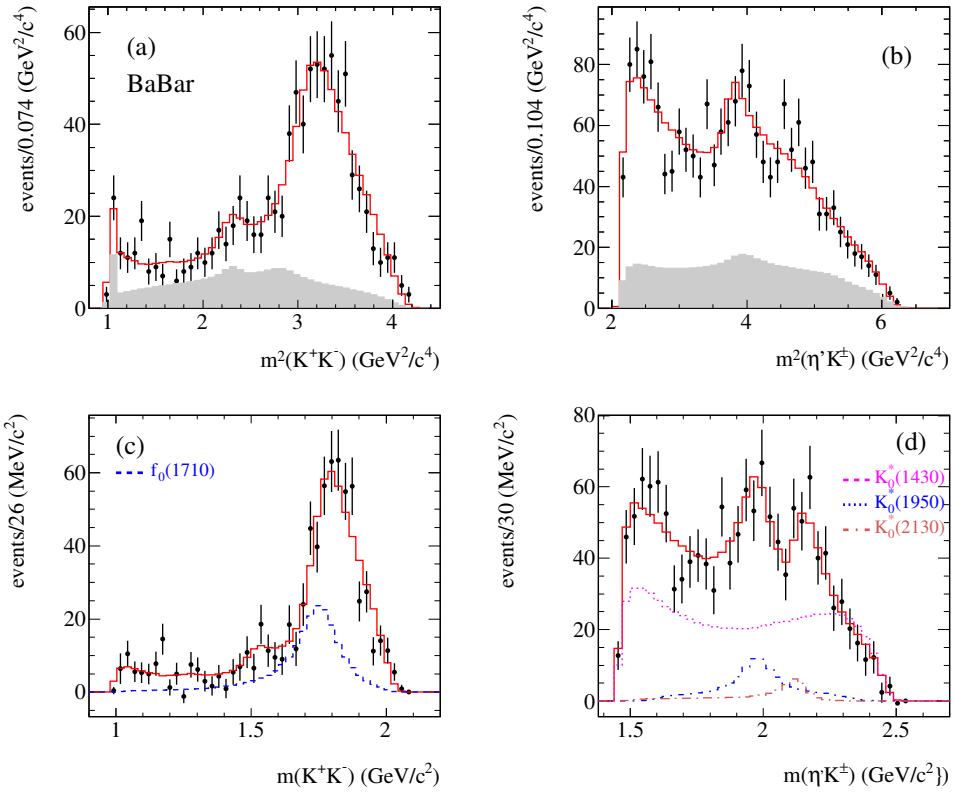
□ Backgrounds evaluated from the η_c sidebands.

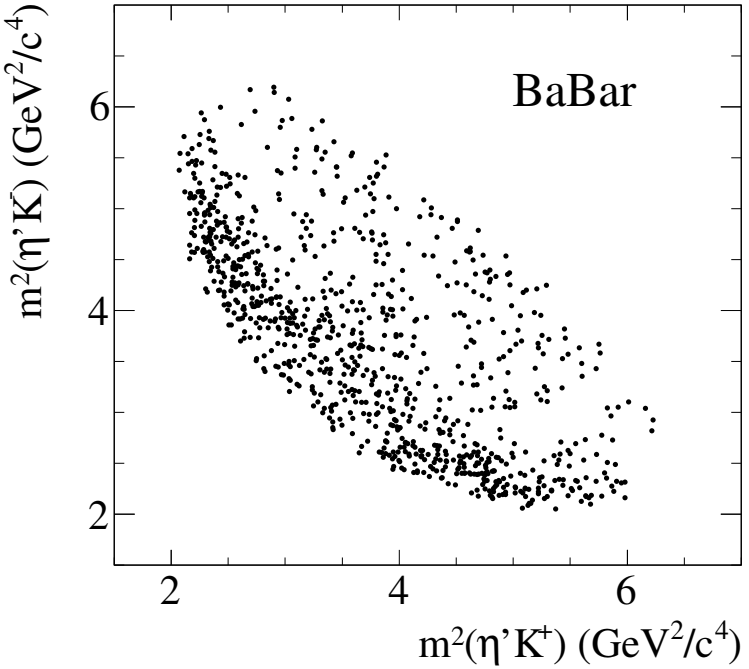
□ We scan the likelihood adding resonances one by one and testing the likelihood increase.

Final state	Decay mode	Yield	Fraction (%)	Purity
$\eta_c \to \eta' K^+ K^-$	$\eta' \to \rho^0 \gamma$	656	0.705	69.7 ± 1.7
	$\eta' \to \pi^+ \pi^- \eta$	274	0.295	85.7 ± 2.0
$\eta_c \to \eta' \pi^+ \pi^-$	$\eta' \to \rho^0 \gamma$	2239	0.717	51.8 ± 1.1
	$\eta' \to \pi^+ \pi^- \eta$	883	0.283	69.0 ± 1.6
$\eta_c \to \eta \pi^+ \pi^-$	$\eta \to \gamma\gamma$	6512	0.700	58.0 ± 0.6
	$\eta \to \pi^+ \pi^- \pi^0$	2791	0.300	52.7 ± 1.0

Dalitz plot analysis of $\eta_c \to \eta' K^+ K^-$

- Combined $\eta' K^+ K^-$ Dalitz plot.
- Mass projections in quadratic and background subtracted linear mass scales.
- Shaded is the fitted background.





- Presence of $f_0(1710)$, $K_0^*(1430)$, $K_0^*(1950)$ and possibly a $K_0^*(2130)$ resonance.

Model for the $K_0^*(1430)$ resonance.

□ We describe the $K_0^*(1430)$ resonance using a coupled channel Breit Wigner with couplings to the $K\pi$ and $K\eta'$ final states.

$$BW(m) = \frac{1}{m_0^2 - m^2 - i(\rho_1(m) g_{\pi K}^2 + \rho_2(m) g_{\eta' K}^2)},$$

where $\rho_i(m) = 2P/m$ and $\rho_2(m)$ becomes imaginary below the $K\eta'$ threshold.

□ We obtain the $K_0^*(1430)$ parameters from a fit to the $K\pi$ $\mathcal{S}$-wave from $\eta_c \to K\bar{K}\pi$.

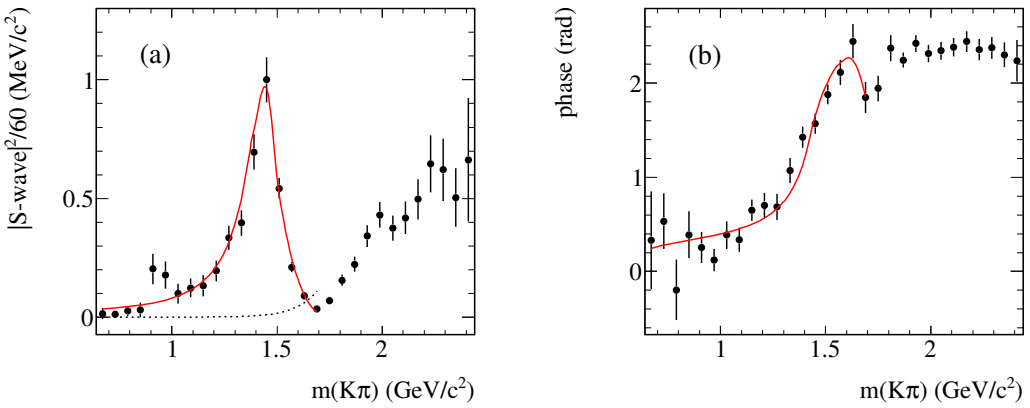
$$\mathcal{S}\text{-wave}(m) = B(m) + c \cdot BW_{K\pi}(m) e^{i\phi},$$

where $B(m) = \rho_1(m) e^{-\alpha m}$.

□ Combined $K_S^0 K\pi$ and $K^+K^-\pi^0$ data from BaBar.

(J.P. Lees et.al., Phys. Rev. D 93, 012005 (2016).)

$m_0 = 1447 \pm 8$ MeV,
$g_{K\pi}^2 = 0.414 \pm 0.026$ GeV$^2/c^4$,
$g_{K\eta'}^2 = 0.197 \pm 0.105$ GeV$^2/c^4$,
$\frac{g_{K\eta'}^2}{g_{K\pi}^2} = 0.476 \pm 0.254.$



Results from the Dalitz plot analysis of $\eta_c \to \eta' K^+ K^-$

- Using the above $K_0^*(1430)$ parameters we obtain the following fractions and relative phases.

Final state	fraction (%)	phase (rad)
$f_0(1710)\eta'$	$29.5 \pm 4.7 \pm 1.6$	0.
$K_0^*(1430)^+ K^-$	$53.9 \pm 7.2 \pm 2.0$	$0.61 \pm 0.13 \pm 0.45$
$K_0^*(1950)^+ K^-$	$2.4 \pm 1.2 \pm 0.4$	$0.46 \pm 0.29 \pm 0.50$
$f_0(1500)\eta'$	$0.8 \pm 1.0 \pm 0.3$	$0.32 \pm 0.54 \pm 0.10$
$f_0(980)\eta'$	$4.7 \pm 2.7 \pm 0.4$	$-0.74 \pm 0.55 \pm 0.05$
$f_2(1270)\eta'$	$2.9 \pm 1.5 \pm 0.1$	$2.9 \pm 0.38 \pm 0.09$
sum	94.3 ± 9.3	

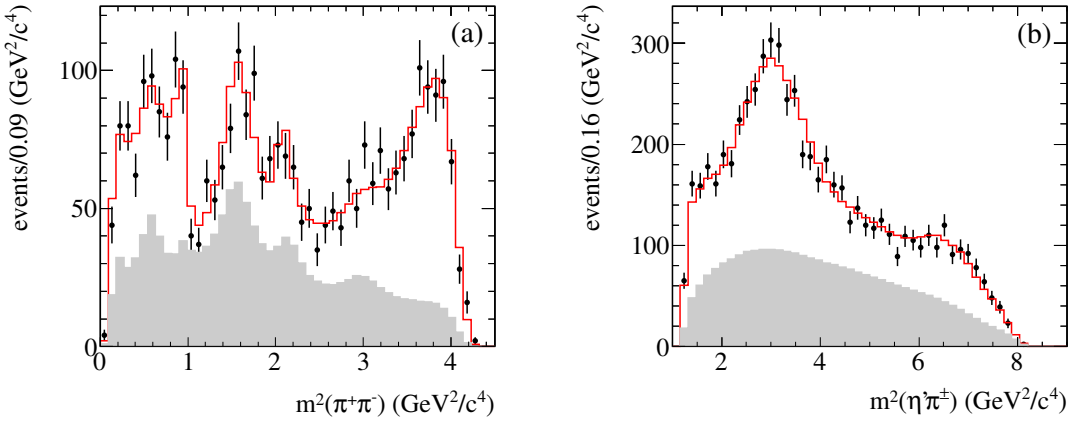
- Leaving free the $g_{K\eta'}$ parameter, we obtain $g^2_{K\eta'} = 0.113 \pm 0.279$ GeV$^2/c^4$.
- We also leave free other resonances parameters and obtain:

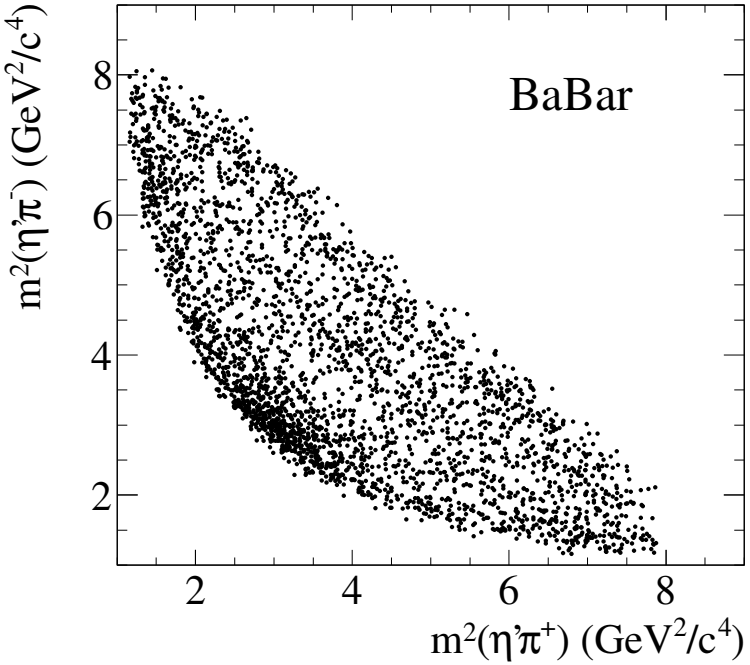
Resonance	Mass (MeV/c^2)	Γ (MeV)	significance (nσ)
$f_0(1710)$	$1757 \pm 24 \pm 9$	$175 \pm 23 \pm 4$	8.2
$K_0^*(1950)$	$1979 \pm 26 \pm 3$	$144 \pm 44 \pm 21$	4.5
$K_0^*(2130)$	$2128 \pm 31 \pm 9$	$95 \pm 42 \pm 76$	3.1

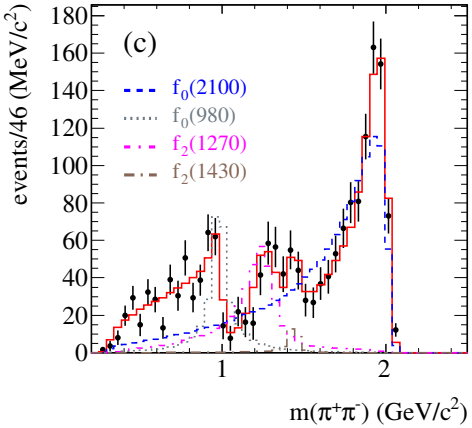
- We obtain $\chi^2/ndf = 281/262 = 1.1$ on the $(m(K^+K^-), cos\theta_H)$ plane, and $ndf = N_{cells} - N_{par}$.
- Non-resonant contribution consistent with zero.

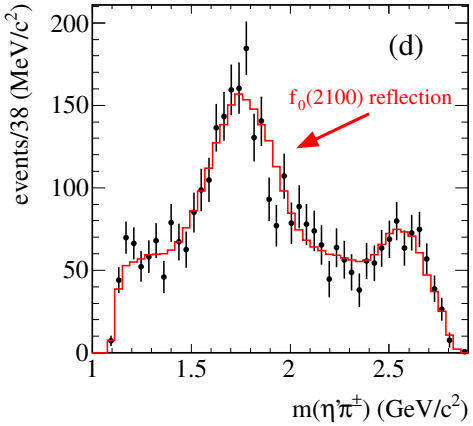
Dalitz plot analysis of $\eta_c \to \eta' \pi^+ \pi^-$

- Combined $\eta' \pi^+ \pi^-$ Dalitz plot.
- Mass projections in quadratic and background subtracted linear mass scales.









- **Decay dominated by the $f_0(2100) \to \pi^+ \pi^-$ resonance.**

Results from the Dalitz plot analysis of $\eta_c \to \eta' \pi^+ \pi^-$

Final state	fraction (%)	phase (rad)
$f_0(2100)\eta'$	$74.9 \pm 7.5 \pm 3.6$	0.
$f_0(500)\eta'$	$4.3 \pm 2.3 \pm 0.7$	$-5.89 \pm 0.24 \pm 0.10$
$f_0(980)\eta'$	$16.1 \pm 2.4 \pm 0.5$	$-5.31 \pm 0.16 \pm 0.04$
$f_2(1270)\eta'$	$22.1 \pm 2.9 \pm 2.4$	$-3.60 \pm 0.16 \pm 0.03$
$f_2(1430)\eta'$	$1.9 \pm 0.7 \pm 0.1$	$-2.45 \pm 0.32 \pm 0.11$
$a_2(1710)\pi$	$3.2 \pm 1.9 \pm 0.5$	$-0.75 \pm 0.27 \pm 0.11$
$a_0(1950)\pi$	$2.5 \pm 1.1 \pm 0.1$	$-0.02 \pm 0.32 \pm 0.06$
$f_2(1800)\eta'$	$5.3 \pm 2.2 \pm 1.4$	$0.67 \pm 0.24 \pm 0.08$
sum	$130.5 \pm 9.5 \pm 4.7$	

- Non-resonant contribution consistent with zero. $\chi^2/ndf = 409/386 = 1.1$
- We leave free the parameters of the $f_0(500)$, $f_2(1430)$, and $f_0(2100)$.

Resonance	Mass (MeV/c^2)	Γ (MeV)	significance (nσ)
$f_0(500)$	953 ± 90	335 ± 81	
$f_2(1430)$	$1440 \pm 11 \pm 3$	$46 \pm 15 \pm 5$	4.8
$f_0(2100)$	$2116 \pm 27 \pm 17$	$289 \pm 34 \pm 15$	10

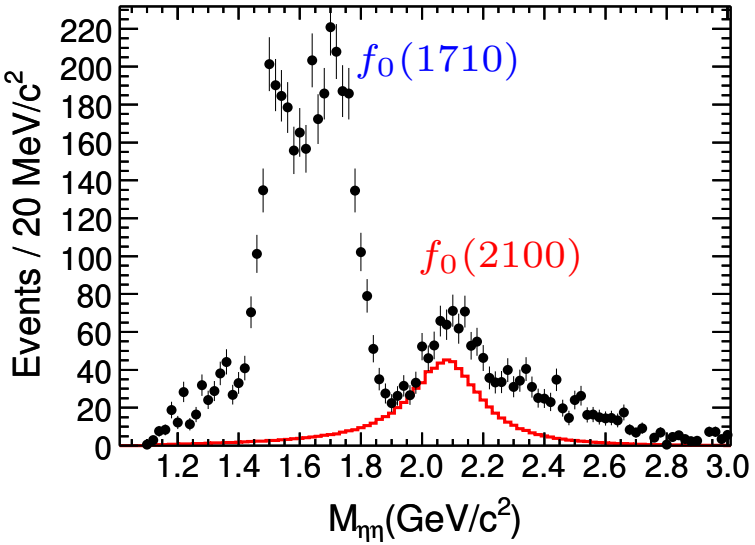
- A weak $J^{PC} = 2^{++}$ signal in the 1.4-1.5 GeV mass region was seen in πp, pp interactions and J/ψ radiative decays.

$f_0(2100)$ resonance: comparison with BESIII results

□ The $f_0(2100)$ parameters are consistent with BESIII values measured in $J/\psi \to \gamma\eta\eta$

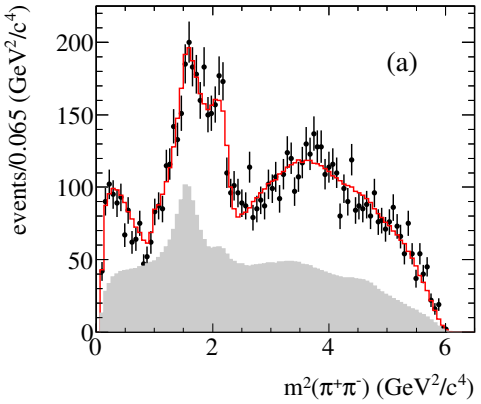
(M. Ablikim *et al.*, Phys. Rev. D **87**, 092009 (2013).)

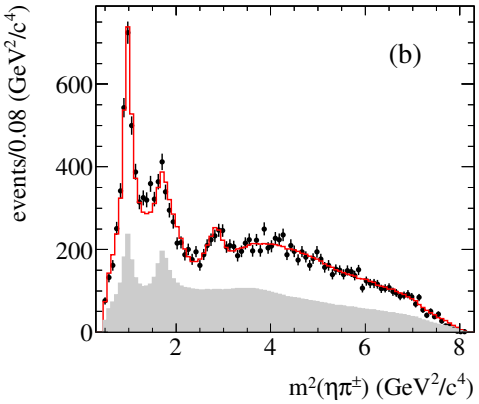
$$m(f_0(2100)) = 2081 \pm 13^{+24}_{-36} \text{ MeV}, \Gamma(f_0(2100)) = 273^{+27+70}_{-24-23} \text{ MeV}$$

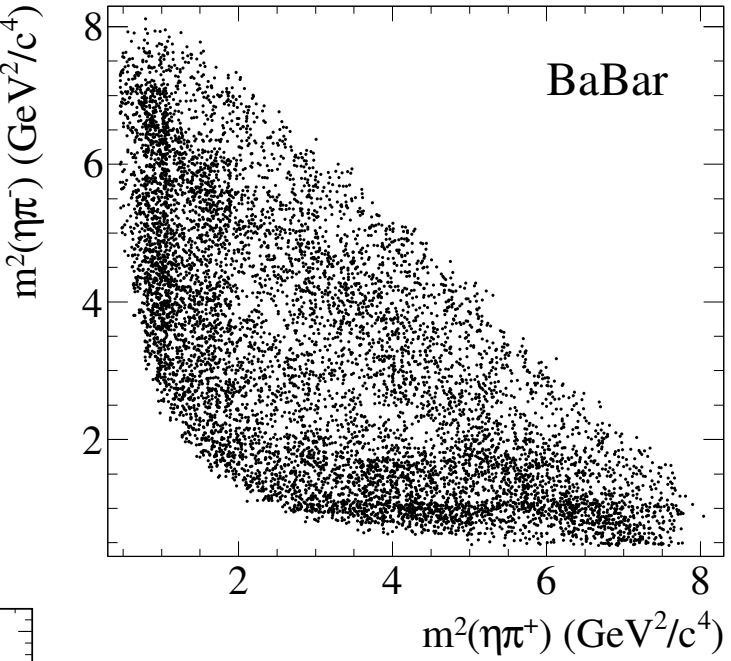


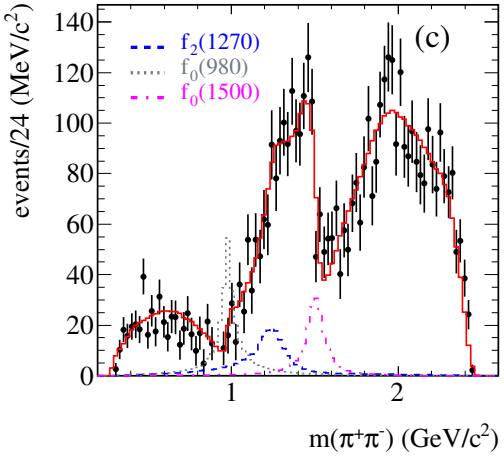
Dalitz plot analysis of $\eta_c \to \eta\pi^+\pi^-$

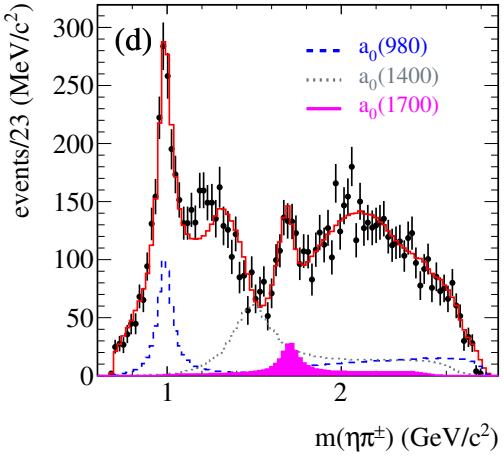
- Combined $\eta\pi^+\pi^-$ Dalitz plot.
- Mass projections in quadratic and background subtracted linear mass scales.











- Complex resonant structures both in signal and background regions.
- A new $a_0(1700) \to \eta\pi$ resonance.

Just confirmed by BESIII as $a_0(1710)^+ \to K^0_S K^+$

(M. Ablikim et al. arXiv:2204.09614.)

Results from the Dalitz plot analysis of $\eta_c \to \eta\pi^+\pi^-$

Final state	fraction (%)	phase (rad)
$a_0(980)\pi$	$12.3 \pm 1.2 \pm 0.9$	0.
$a_2(1310)\pi$	$2.5 \pm 0.7 \pm 0.6$	$-1.04 \pm 0.13 \pm 0.20$
$f_0(500)\eta$	$4.3 \pm 1.3 \pm 0.7$	$0.54 \pm 0.14 \pm 0.20$
$f_2(1270)\eta$	$4.6 \pm 0.9 \pm 0.4$	$-1.15 \pm 0.11 \pm 0.05$
$f_0(980)\eta$	$5.7 \pm 1.3 \pm 1.0$	$-2.41 \pm 0.09 \pm 0.04$
$f_0(1500)\eta$	$4.2 \pm 0.7 \pm 0.6$	$2.32 \pm 0.13 \pm 0.05$
$a_0(1450)\pi$	$15.0 \pm 2.4 \pm 2.1$	$2.60 \pm 0.09 \pm 0.11$
$a_0(1700)\pi$	$3.5 \pm 0.8 \pm 0.6$	$1.39 \pm 0.15 \pm 0.12$
$f_2(1950)\eta$	$4.2 \pm 1.0 \pm 0.6$	$-1.59 \pm 0.15 \pm 0.20$
sum	$56.3 \pm 3.7 \pm 2.9$	
NR	$172.7 \pm 8.0 \pm 10.0$	$1.67 \pm 0.07 \pm 0.03$

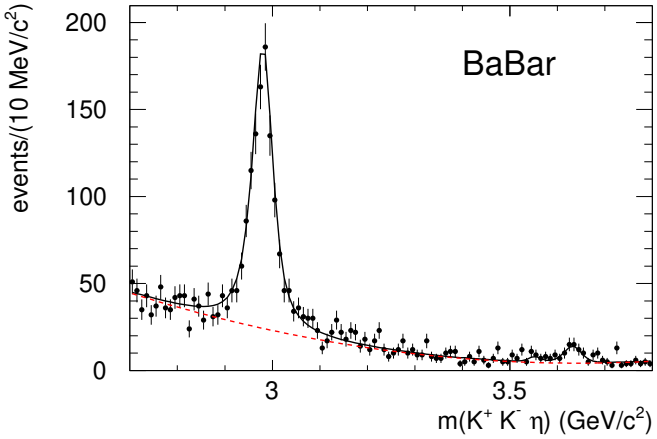
- χ^2 on the $(m(\pi^+\pi^-), \cos\theta_H)$ plane gives $\chi^2/ndf = 419/382 = 1.1$.
- A new $a_0(1700)$ resonance is observed in the $\eta\pi^\pm$ mass spectrum, with fitted parameters:

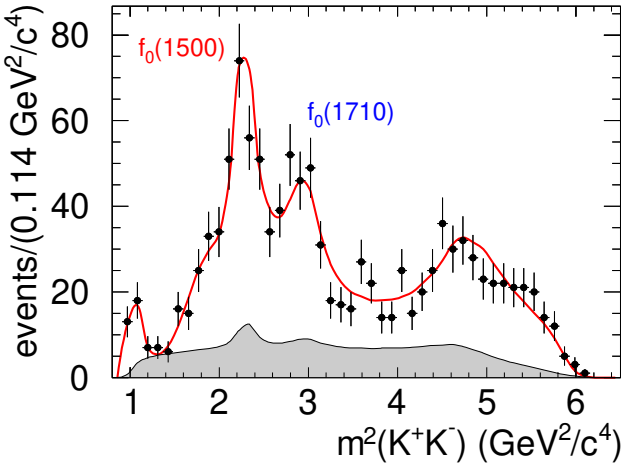
$$m = 1704 \pm 5 \pm 2\,\text{MeV},\ \Gamma = 110 \pm 15 \pm 11\,\text{MeV},\ n\sigma = 8$$

- Large non-resonant contribution, possibly related to the η_c interference with the two-photon background.

Results from $\eta_c \to \eta K^+ K^-$

- BaBar Dalitz plot analysis of $\eta_c \to \eta K^+ K^-$ with $\eta \to \gamma\gamma$ and $\eta \to \pi^+\pi^-\pi^0$ (*J.P. Lees et.al., Phys. Rev. D89, 112004 (2014)*).
- $\eta_c \to \eta K^+ K^-$ signal and $K^+ K^-$ squared mass projection.





- Signals of $f_0(1500)$ and $f_0(1710)$.

Summary on the study of η_c decays.

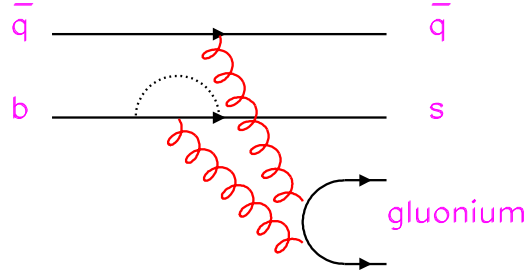
- First observation of $\eta_c \to \eta' K^+ K^-$ and measurement of its branching fraction.
- Observation of a new $a_0(1700) \to \eta\pi$ resonance.
- Evidence for $f_2(1430) \to \pi^+\pi^-$.
- Comparison between the fractional contributions in η_c decays to gluonium candidates for final states involving η or η' meson. *Not corrected for unseen decay modes.*

Final state	$f_0(1500)(\%)$	$f_0(1710)(\%)$	$f_0(2100)(\%)$
$\eta K^+ K^-$	$23.7 \pm 7.0 \pm 1.8$	$8.9 \pm 0.2 \pm 0.4$	
$\eta\pi^+\pi^-$	$4.2 \pm 0.7 \pm 0.9$		0.
$\eta' K^+ K^-$	$0.8 \pm 1.0 \pm 0.3$	$29.5 \pm 4.7 \pm 1.6$	
$\eta'\pi^+\pi^-$	0.3 ± 0.2		$74.9 \pm 7.5 \pm 3.5$

- We observe an enhanced contributions of $f_0(1710)$ and $f_0(2100)$ in η_c decays to η'.
- This effect may point to an enhanced gluonium content in the $f_0(1710)$ and $f_0(2100)$ resonances.
- A similar conclusion is drawn in the study of J/ψ radiative decays by BESIII.
- **The results obtained in this analysis allow to add the $f_0(2100)$ resonance in the list of the candidates for the scalar glueball.**

What Next?

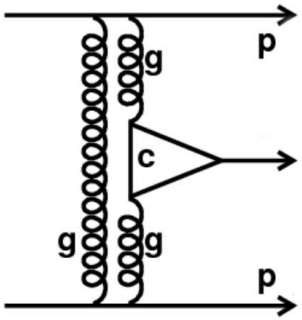
□ B decays with strangeness in the final state have been suggested as possible final states where to search for gluonium (H. Fritzsch,Phys. Lett. B **415** (1997), 83-89)



□ Large samples of three-body and four-body charmless B-decays are being collected by LHCb

□ Data are also being collected on "Central Exclusive Production" at LHCb.

□ Double-pomeron-exchange processes are a Glue Laboratory.



□ The spectroscopy of states having exotic J^{PC} quantum number $(0^{--}, 0^{+-}\ 1^{-+}, 2^{+-}, ...)$ still needs to be explored.

□ The COMPASS/AMBER experiment at CERN is actively working in this field.

□ New results will come from e^+e^- experiments, BESIII and BELLEII.

□ The search for gluonium and hybrids is one of the main parts of the program of the study of $\bar{p}p$ annihilations of PANDA experiment at FAIR.