





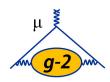
E.Bottalico

Muse General Meeting – 25 October 2019

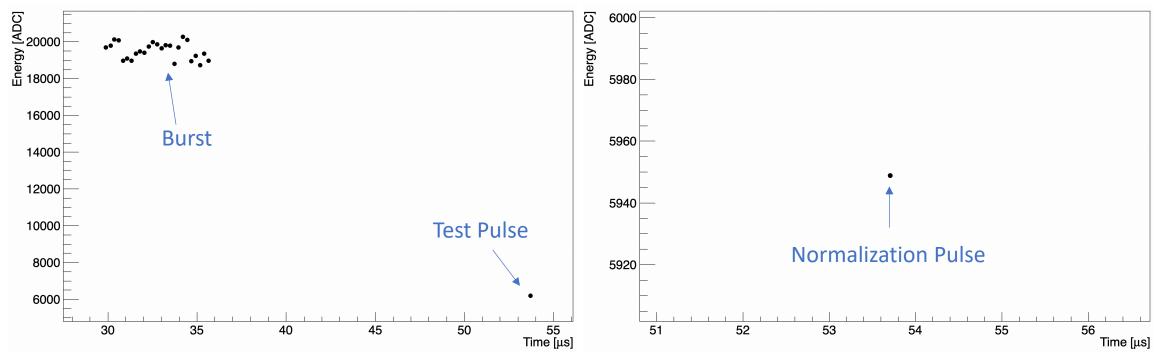




### LTDP Study

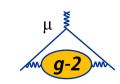


- Mimic the splash with laser in Burst mode.
- Measure the gain sag using a probe, Test Pulse (T), normalized to a Normalization Pulse (N).

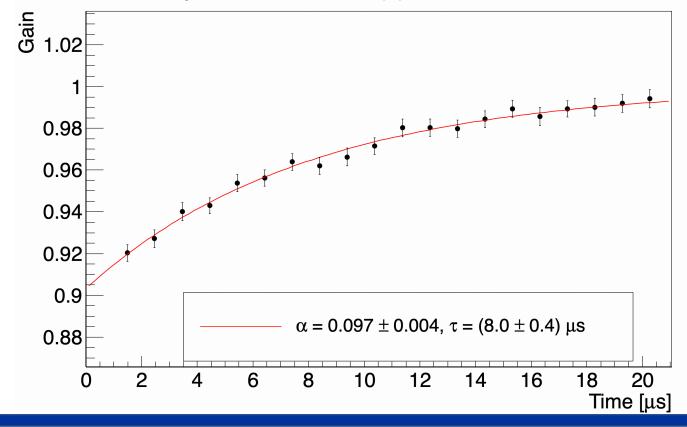




## LTDP Study

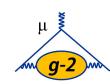


- We built the gain function computing for each point  $G(t) = \frac{\langle T \rangle}{\langle N \rangle}$  where t is the delay between the last pulse in the burst and the Test pulse.
- The fit function is an exponential:  $G(t) = 1 \alpha e^{-\frac{c}{\tau}}$

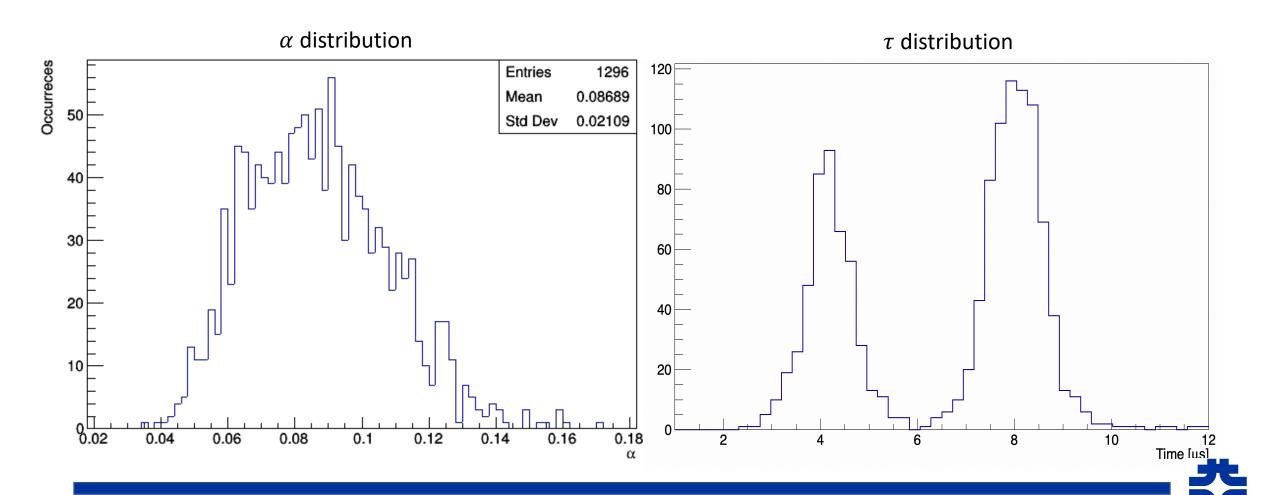




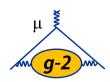
# LTDP Study



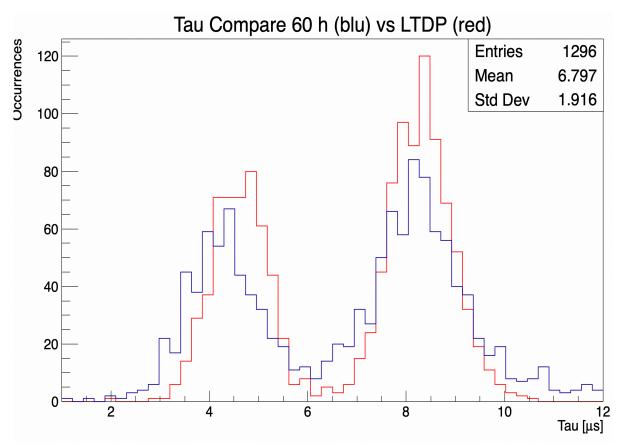
• Distribution of the  $\tau$  and  $\alpha$  parameters for each crystal.

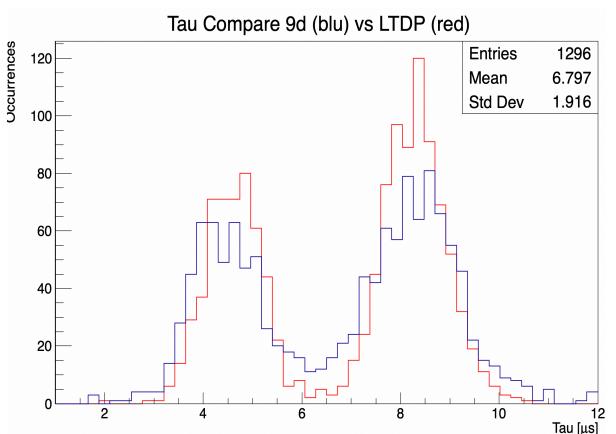






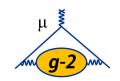
These are the comparison between the  $\tau$  distributions obtained by LTDP and the 3 datasets (60h, 9d and Endgame).



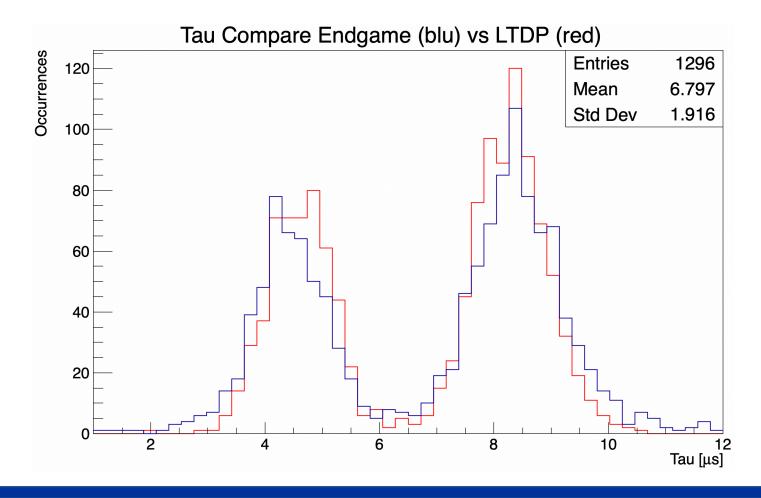






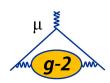


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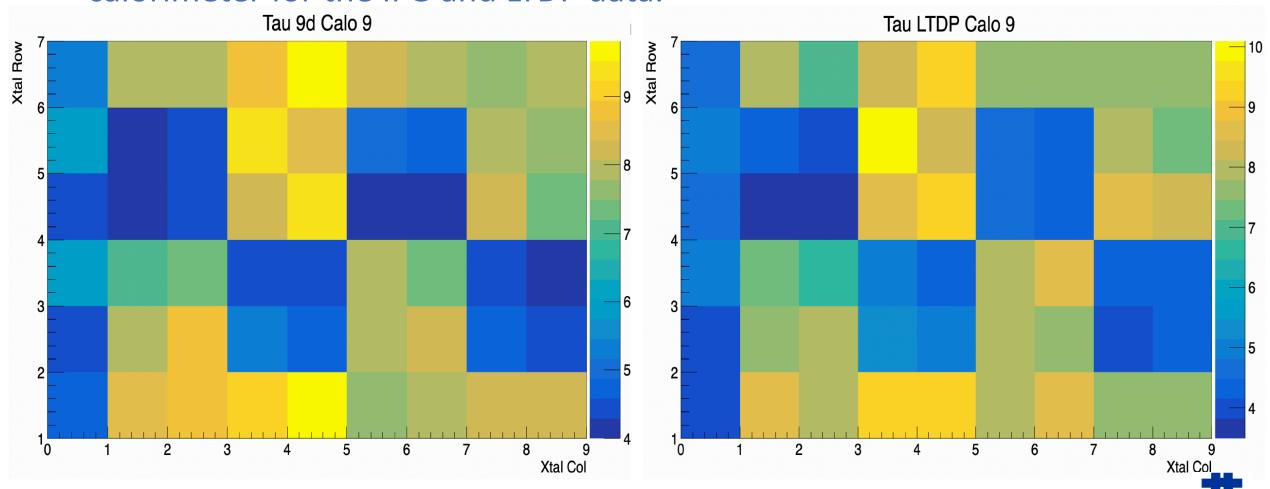








The distributions of the  $\tau$  parameters have the same pattern in the calorimeter for the IFG and LTDP data.

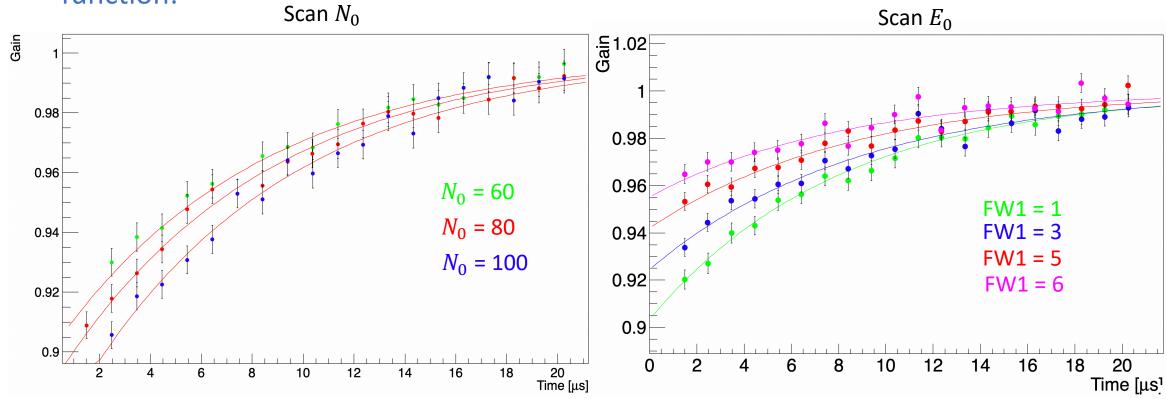




# **Energy Study**

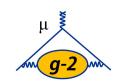
μ **§** 

• To study the energy dependence of  $\alpha$  we change the number of pulses in the Burst  $(N_0)$  and the energy of the single pulse in the burst  $(E_0)$ , by varying the Filter Wheel (FW) position. In the following plots are shown different gain function for different combination  $N_0$  and  $E_0$ , fitted with the same exponential function:





# **Energy Study**



We compute an analytical function to describe the gain sag due to the Burst mode:

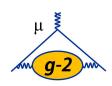
$$G(t) = 1 - (1 - e^{-p_0 E_0}) \underbrace{\frac{1 - e^{-\frac{N_0 \delta}{\tau}}}{1 - e^{-\frac{\delta}{\tau}}} e^{\frac{t}{\tau}}}_{F(N_0, \delta, \tau)}$$

We put  $(1 - e^{-p_0 E_0}) \cdot F(N_0, \delta, \tau)$  instead of  $\alpha$ , where  $p_0$  is the gain sag in function of energy and  $\delta$  the distance between 2 pulses in the burst (120 ns). In this way we can characterize the gain sag in function of the parameters  $p_0$  and  $\tau$ .



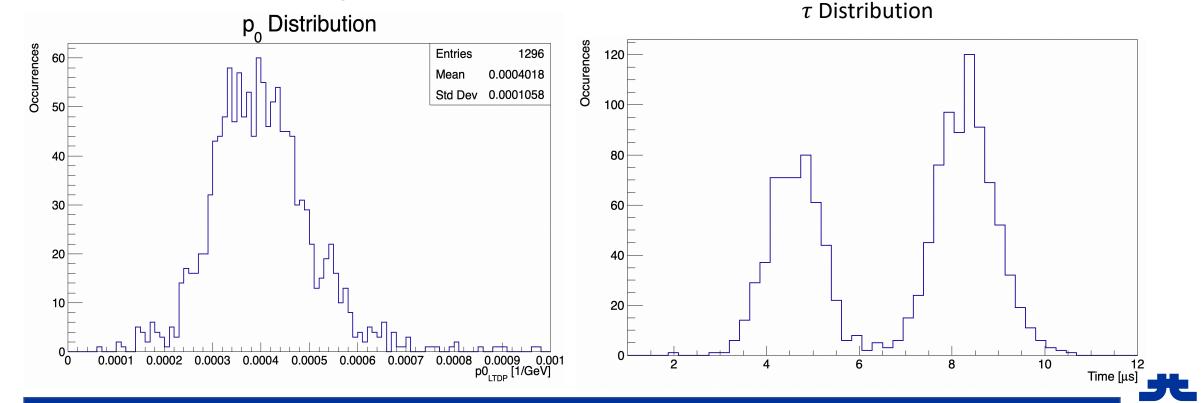


### Fit Mode



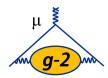
We perform a combined fit of the gain building a 3-dimensional histogram, for each SIPM, where each bin of coordinates  $(E_0, N_0, t)$  is weighted with the value of the gain in that bin. We perform the fit with the function of  $G(t) = 1 - (1 - e^{-p_0 E_0}) F(\delta, N_0, \tau) e^{-\frac{t}{\tau}}$  shown before,

perform the fit with the function of  $G(t) = 1 - (1 - e^{-p_0 E_0}) F(\delta, N_0, \tau) e^{-\tau}$  shown before with free parameters  $\tau$  and  $p_0$ . These are the distributions of the parameters:



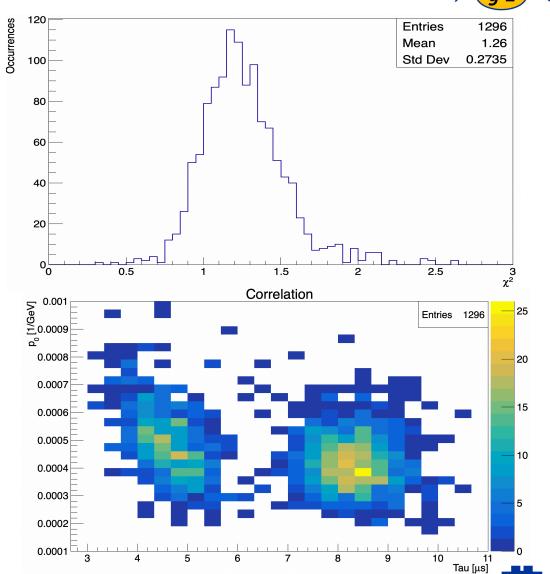


### Fit Mode



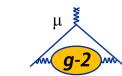
Distribution of the  $\chi^2$  of the 1296 fits

Scatter plot  $p_0$  -  $\tau$  to show the correlation between the parameters (p=-0.36)





# Definition $\alpha_{LTDP}$



The IFG curves are fitted with  $G_{IFG}(t)=1-lpha_{IFG}e^{-t/ au_{IFG}}$ . To compare the lpha parameter, we define  $lpha_{LTDP}$  modeling the splash with a Gaussian. Convoluting the Gaussian distribution  $E_0$ .

 $\frac{1}{\sqrt{2\pi}\sigma}e^{-\frac{t^2}{\delta t^2}}$  (where  $\delta t$  is the time spread of the Splash) with the SiPM response  $p_0e^{-\frac{t}{\tau}}$  we obtain:

$$\frac{E_0(p_0)}{\sqrt{2}}e^{\frac{\delta t^2 - 4\tau \cdot t}{4\tau^2}}(t=0) \to \frac{E_0(p_0)}{\sqrt{2}}e^{\frac{\delta t^2}{4\tau^2}}$$

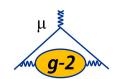
In such way we can approximate  $G(t)=1-(1-e^{-p_0E_0})F(\delta,N_0,\tau)e^{-\frac{t}{\tau}}$  as  $G(t)\sim 1-p_0E_0e^{-\frac{t}{\tau}}$  and assuming  $\alpha_{LTDP}$  as:

$$\alpha_{LTDP} = \frac{E_{splash} \cdot p_0}{\sqrt{2}} e^{\frac{\delta t^2}{4\tau^2}}$$

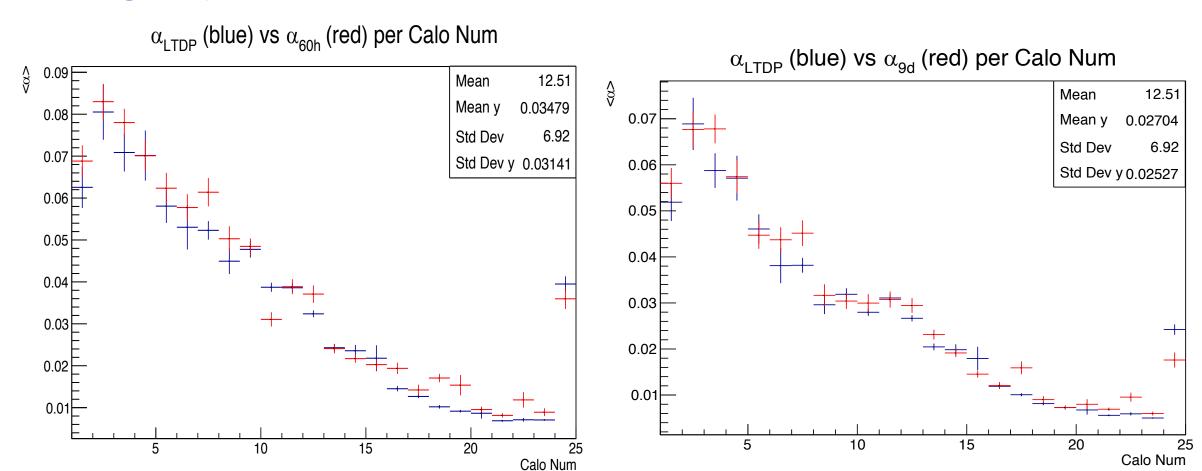
Where  $E_{splash \ meas}$  is the splash energy measured.





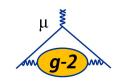


Comparison between  $\alpha_{LTDP}$  and  $\alpha_{IFG}$  for the 3 dataset (60h, 9d and Endgame):



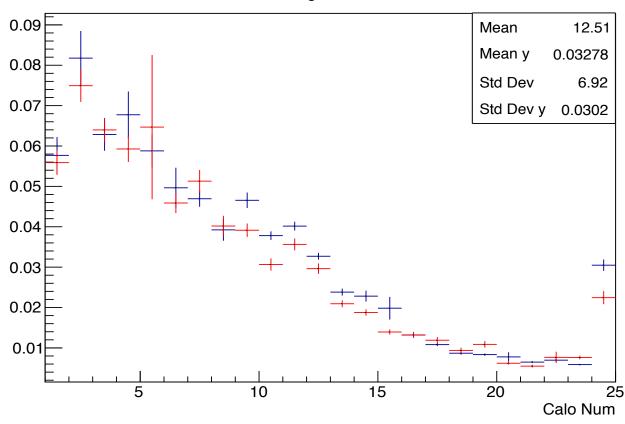






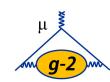
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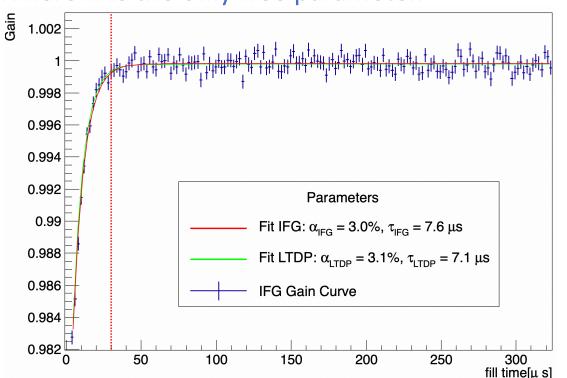




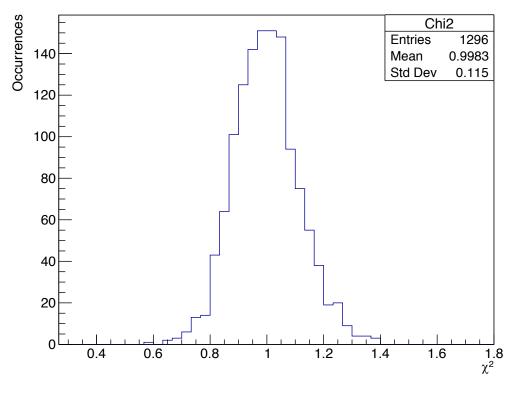
The next step is to build the LTDP correction functions to compare them with the IFG functions. In these plots are shown the IFG gain curves (blue dots), the IFG fit with  $G_{IFG}(t)$  (red curve) and the LTDP correction function (green curve) define in this way:

$$G_{LTDP}(t) = N \left( 1 - \alpha_{LTDP} \cdot e^{-\frac{t}{\tau_{LTDP}}} \right)$$

#### Where N is the only free parameter.

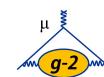


#### $\chi^2$ Distribution









Finally to study quantitatively the goodness of the LTDP corrections we apply the  $G_{LTDP}$  function (for the 60h dataset) instead of  $G_{IFG}$  to correct positron energy  $E_{\mathcal{D}}$ .

$$E_p' = \frac{E_p}{G_{LTDP}(\Delta t)}$$

The fit result are:

• LTDP Correction:  $\frac{\chi^2}{NDF}$ : 0.954134 with prob 0.973296 R: -42.707 ± 1.33528

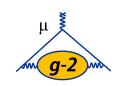
• IFG Correction:  $\frac{\chi^2}{NDF}$ : 0.954346 with prob 0.972728 R: -42.6989  $\pm$  1.3354

$$\delta R = 8.1 \text{ ppb}$$





#### Conclusion



- Using the Long Term Double Pulse we can describe the gain sag of the SIPM as a function of energy by 2 parameters  $(p_0, \tau)$ .
- By assuming a convolution model for the splash we can compute the gain sag as  $G_{LTDP}(t) = N\left(1 \alpha_{LTDP} \cdot e^{-\frac{t}{\tau_{LTDP}}}\right)$  and find a good agreement with the IFG.
- The effect of these corrections are inside the 20 ppb final score for gain corrections.
- This is a perfect indipendet check for the IFG corrections.