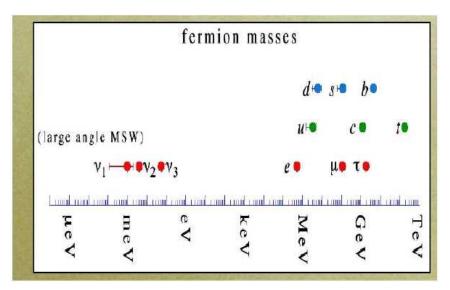
Flavour Physics in SUSY extensions of the SM in the LHC era

P. Paradisi



Third Workshop on Theory, Phenomenology and Experiments in Heavy Flavour Physics 2010 Capri, Italy 5-7 July, 2010

The fermion mass puzzle



SM flavour puzzle

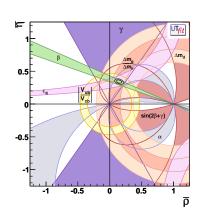
Smallness and Hierarchy

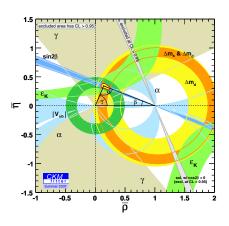
$$\begin{split} Y_t \sim 1, \quad Y_c \sim 10^{-2}, \quad Y_u \sim 10^{-5} \\ Y_b \sim 10^{-2}, \quad Y_s \sim 10^{-3}, \quad Y_d \sim 10^{-4} \\ Y_\tau \sim 10^{-2}, \quad Y_\mu \sim 10^{-3}, \quad Y_e \sim 10^{-6} \\ |V_{us}| \sim 0.2, \quad |V_{cb}| \sim 0.04, \quad |V_{ub}| \sim 0.004, \quad \delta_{\rm KM} \sim 1 \end{split}$$

- For comparison: $g_s \sim 1$, $g \sim 0.6$, $g' \sim 0.3$, $\lambda \sim 1$
- The SM flavor parameters have structure: smallness and hierarchy
- Why? = The SM flavor puzzle

Nir

SM success



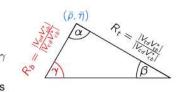


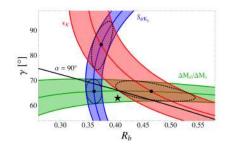
Very likely, flavour and CP violation in FC processes are dominated by the CKM mechanism (Nir)

• Recent theoretical improvements in $\epsilon_{\it K}$ expose some tensions in the UT analysis [Lunghi & Soni,

Buras & Guadagnoli]

- Look at ε_K, S_{ψKS} (sin 2β), ΔM_d/ΔM_s in the R_b-γ plane
- R_b , γ can be obtained from tree-level processes

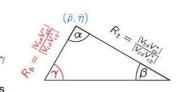


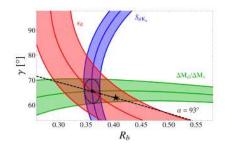


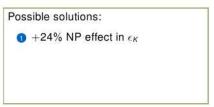
• Recent theoretical improvements in ϵ_K expose some tensions in the UT analysis [Lunghl & Soni,

Buras & Guadagnoli]

- Look at ε_K, S_{ψKS} (sin 2β), ΔM_d/ΔM_S in the R_b-γ plane
- R_b , γ can be obtained from tree-level processes

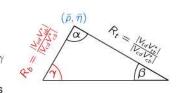


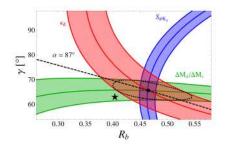




 Recent theoretical improvements in ε_K expose some tensions in the UT analysis (Lunghi & Soni, Buras & Guadagnoli)

- Look at ε_K, S_{ψKS} (sin 2β), ΔM_d/ΔM_s in the R_b-γ plane
- R_b, γ can be obtained from tree-level processes





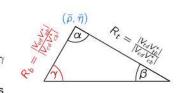
Possible solutions:

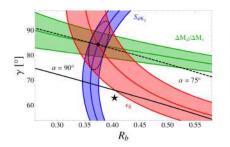
- $_{f 0}$ +24% NP effect in ϵ_K
- $oldsymbol{0} = -6.5^{\circ}$ NP phase in B_d mixing

• Recent theoretical improvements in ϵ_K expose some tensions in the UT analysis [Lunghi & Soni,

Buras & Guadagnoli]

- Look at ϵ_K , $S_{\psi K_S}$ (sin 2β), $\Delta M_\sigma/\Delta M_S$ in the R_b - γ plane
- R_b , γ can be obtained from tree-level processes





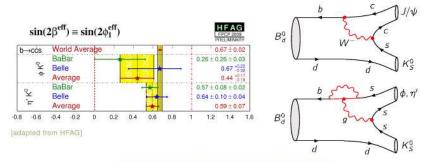
Possible solutions:

- $_{\odot}$ +24% NP effect in ϵ_{K}
 - $2 6.5^{\circ}$ NP phase in B_d mixing
 - 3 -22% NP effect in $\Delta M_d/\Delta M_s$ (requiring $\alpha \sim 74^{\circ}$)

$\sin 2\beta_{eff}$ tensions

- In the SM, mixing-induced CP asymmetries in $B_d \to \psi K_S$, ϕK_S , $\eta' K_S$ all $\approx \sin 2\beta$
- $B_d \to \psi K_S$ dominated by tree level, ϕK_S and $\eta' K_S$ are loop-induced

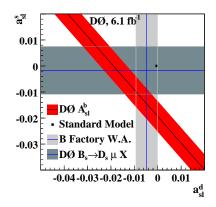
Data indicate
$$S_{\phi K_S} < S_{\eta' K_S} < S_{\psi K_S}$$



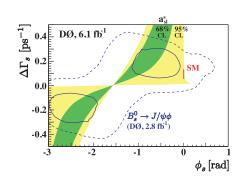
New physics in the decay amplitudes?

Can only be resolved at SuperB

CPV in B_s mixing



$$S_{\psi\phi}=\sin(2|eta_{
m S}|-2\phi_{
m B_{
m S}})\;,$$



$$A_{\mathsf{SL}}^q \equiv rac{\Gamma(ar{B}_q
ightarrow I^+ X) - \Gamma(B_q
ightarrow I^- X)}{\Gamma(ar{B}_q
ightarrow I^+ X) + \Gamma(B_q
ightarrow I^- X)}$$

New Physics in the B_s mixing phase?

Why CP violation?

Motivation:

- Baryogenesis requires extra sources of CPV
- $\bullet \ \, \text{The QCD} \, \overline{\theta}\text{-term} \, \mathcal{L}_{\mathit{CP}} = \overline{\theta} \tfrac{\alpha_s}{8\pi} \mathsf{G} \tilde{\mathsf{G}} \, \text{is a CPV source beyond the CKM}$
- Most UV completion of the SM have many CPV sources

NP search strategies

Where to look for New Physics at the low energy?

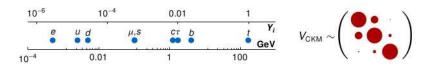
- Processes very suppressed or even forbidden in the SM
 - FCNC processes ($\mu \to {\bf e}\gamma, \, \tau \to \mu\gamma, \, {\it B}^0_{s,d} \to \mu^+\mu^-, \, {\it K} \to \pi \nu \bar{\nu}$)
 - CPV effects in the electron/neutron EDMs, de,n...
 - FCNC & CPV in B_{s,d} decay/mixing & D mixing amplitudes
- Processes predicted with high precision in the SM
 - EWPO as $\Delta \rho$, $(g-2)_{\mu}$
 - LU in $R_M^{e/\mu}$ = $\Gamma(K(\pi) \to e\nu)/\Gamma(K(\pi) \to \mu\nu)$

Flavour Matrix

		FLAVOUR COUPLING						
		$b \rightarrow s \left[-\lambda^2 \text{ in SM}\right]$	$b \to d \left[-\lambda^3 \text{ in SM} \right]$	$s \to d \ \left[\text{-} \lambda^5 \text{in SM} \right]$				
ELECTROWEAK STRUCTURE	ΔF=2 box	$\begin{array}{c} \Delta M^{}_{Bs} \\ A^{}_{CP}(B^{}_s {\rightarrow} \psi \phi), \;\; \epsilon^{}_{Bs} \end{array}$	$\begin{array}{l} \Delta M_{Bd} \\ A_{CP}(B_d{\to}\psi K), \ \ \epsilon_{Bd} \end{array}$	$\varepsilon_{_{\rm K}}$				
	ΔF=1 4-quark ops.	$A_{CP}(B_d \rightarrow \phi K)$	$A_{CP}(B_s \rightarrow \phi K)$					
	gluon penguin	$\begin{split} &A_{CP}(B_d{\to}\phi K)\\ &[\Gamma,\Delta\Gamma_{CP}](B\to X_s\gamma) \end{split}$	$[\Gamma, \Delta \Gamma_{CP}](B \to \! \rho/\pi \gamma)$	$\Gamma(K_L \rightarrow \pi^0 I^* I)$				
	γ penguin	$\begin{split} & [\Gamma, \Delta \Gamma_{\text{CP}}] (B \to & X_s \gamma) \\ & [\Gamma, \Delta \Gamma_{\text{CP}}] (B \to & X_s \Gamma \Gamma) \\ & A_{\text{FB}} (B \to & X_s \Gamma \Gamma) \end{split}$	$\begin{split} & [\Gamma, \Delta\Gamma_{CP}](B \to \! \rho/\pi \; \gamma) \\ & [\Gamma, \Delta\Gamma_{CP}](B \to \! \rho/\pi \; \Gamma\Gamma) \\ & A_{FB}(B \to \! \rho/\pi \; \Gamma\Gamma) \end{split}$	$\Gamma(K_L {\to} \pi^0 l^* l^*)$				
	Z ⁰ penguin	$\begin{split} & [\Gamma, \Delta \Gamma_{\text{CP}}] (B \rightarrow X_s f f) \\ & A_{\text{FB}} (B \rightarrow X_s f T) \\ & \Gamma(B_s \rightarrow \mu \mu) \end{split}$	$\begin{split} & [\Gamma, \Delta \Gamma_{\rm CP}] (B \rightarrow \! \rho/\pi \ l' l') \\ & A_{\rm FB} (B \rightarrow \! \rho/\pi \ l' l') \\ & \Gamma (B_d \! \rightarrow \! \mu \mu) \end{split}$	$\Gamma(K^* \rightarrow \pi^* \nu \nu)$ $\Gamma(K_L \rightarrow \pi^0 \nu \nu)$ $\Gamma(K_L \rightarrow \pi^0 l^* l^*)$				
	H ⁰ penguin	$\Gamma(B_s \rightarrow \mu\mu)$	$\Gamma(B_d \rightarrow \mu\mu)$					

SM vs. SUSY flavour problems

Flavour violation is highly non-generic already in the SM!



The two problems should be related!

Minimal Flavour Violation (MFV)

- Yukawa couplings are the only sources of flavour violation
- Effective theory
- Pragmatic approach
- Pessimistic phenomenology

Flavour Models

- Flavour structure of Yukawa couplings and soft terms generated by spontaneous breaking of a flavour symmetry
- Ambitious approach
- Diverse phenomenology

Minimal Flavour Violation

SM without Yukawa interactions: SU(3)⁵ global flavour symmetry

$$SU(3)_u \otimes SU(3)_d \otimes SU(3)_Q \otimes SU(3)_e \otimes SU(3)_L$$

- Yukawa interactions break this symmetry
- Proposal for any New Physics model:

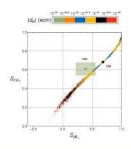
Yukawa structures as the only sources of flavour violation

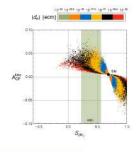
 \Downarrow

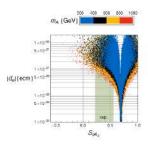
Minimal Flavour Violation
MFV allows for new "flavour blind" CPV phases!

Altmannshofer, Buras and P.P., '08

Flavor blind MSSM ≈ MFV + CPV







- ► CP violating △F = 0 and △F = 1 dipole amplitudes can be strongly modified
- S_{φKS} and S_{η'KS} can simultaneously be brought in agreement with the data
- sizeable and correlated effects in
 A^{bsn}_{CP} ≈ 1% − 6%
- ► lower bounds on the electron and neutron EDMs at the level of d_{e,n} ≥ 10⁻²⁶ ecm
- ► large and correlated effects in the CP asymmetries in B → K* µ⁺ µ⁻ (WA, Ball, Bharucha, Buras, Straub, Wick)

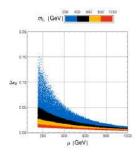
- the leading NP contributions to ΔF = 2 amplitudes are not sensitive to the new phases of the FBMSSM
- CP violation in meson mixing is SM like
- lackbox i.e. small effects in $S_{\psi\,\phi},\,S_{\psi\,\kappa_S}$ and $\epsilon_{\it K}$
- lacksquare in particular: $0.03 < S_{\psi\,\phi} < 0.05$

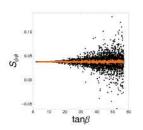
A combined study of all these observables and their correlations constitutes a very powerful test of the FBMSSM

Phenomenology of the flavor blind MSSM

Kaon mixing

- ► The mixing amplitude M^K₁₂ has no sensitivity to the new flavor blind phases
- Still, €_K ∝ Im(M^K₁₂) can get a positive NP contribution up to 15%
- ▶ But only for a very light SUSY spectrum: μ , $m_{T_c} \simeq 200 \text{GeV}$





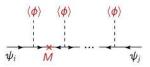
- B_d and B_s mixing
- Leading NP contributions to M^{d,s}₁₂ are insensitive to the new phases of a FBMSSM.
 (at least for moderate tan β...)
- ▶ For large $\tan \beta$, the constraint from $b \rightarrow s\gamma$ does not allow for sizeable effects
- \blacktriangleright $S_{\psi K_S}$ and $S_{\psi \phi}$ are SM like $(S_{\psi \phi} \simeq 0.03 0.05)$

SUSY flavour models

Main idea: hierarchies in Yukawa couplings generated by spontaneous breakdown of flavour symmetry (horizontal symmetry, family symmetry)

- · Generalization of the Froggat-Nielsen mechanism
- Yukawa hierarchies explained by different powers of small ε:

$$\Rightarrow Y_{ij} \propto \left(\frac{\langle \phi \rangle}{M}\right)^{(a_i+b_j)} = \epsilon^{(a_i+b_j)}$$



 Possible to relate Yukawa matrices and sfermion mass matrices/trilinear couplings

SUSY flavour models can explain the origin of the hierarchies in the Yukawa couplings and solve the SUSY flavour problem

 Many different viable models exist, with abelian or non-abelian flavour symmetries

Abelian vs. non-Abelian flavour models

Abelian vs. Non-abelian

- In most non-abelian models, 1st & 2nd generatio sfermions are approximately degenerate
- In abelian models, sfermions of different generations need not be degenerate
 - O(1) 1-2 mass splitting leads to O(λ) (δ_μ^{LL})₁₂ in the SCKM basis

Chirality structure of flavour violating terms

- Different flavour symmetries lead to different patterns of flavour violation
- Mass insertions: $M_{\tilde{d}}^2 = \operatorname{diag}(\tilde{m}^2) + \tilde{m}^2 \begin{pmatrix} \delta_d^{LL} & \delta_d^{LR} \\ \delta_d^{RL} & \delta_d^{RR} \end{pmatrix}$
- δ^{LL}, δ^{RR}, δ^{LR} fixed by the flavour symmetry (up to O(1) factors)

Examples of flavour models

4 representative flavour models with different chirality structures in the \tilde{d} sector:

AC model

U(1)

Large, O(1) RR mass insertions

AKM model

SU(3)

Only CKM-like RR mass insertions

RVV model [Ross.

SU(3)

CKM-like LL & BB mass insertions

δLL model

 $(S_3)^3$

Only CKM-like LL mass insertions

$$\delta_{\sigma}^{LL} \sim \begin{pmatrix} \cdot & 0 & 0 \\ 0 & \cdot & \lambda^2 \\ 0 & \lambda^2 & \cdot \end{pmatrix}$$

$$\delta_d^{LL} \sim \begin{pmatrix} \cdot & 0 & 0 \\ 0 & \cdot & 0 \\ 0 & 0 & \cdot \end{pmatrix}$$

$$\delta_d^{LL} \sim \begin{pmatrix} \cdot & 0 & 0 \\ 0 & \cdot & \lambda^2 \\ 0 & \lambda^2 & \cdot \end{pmatrix} \quad \delta_d^{LL} \sim \begin{pmatrix} \cdot & 0 & 0 \\ 0 & \cdot & 0 \\ 0 & 0 & \cdot \end{pmatrix} \quad \delta_d^{LL} \sim \begin{pmatrix} \cdot & \lambda^3 & \lambda^2 \\ \lambda^3 & \cdot & \lambda \\ \lambda^2 & \lambda & \cdot \end{pmatrix} \delta_d^{LL} \sim \begin{pmatrix} \cdot & \lambda^5 & \lambda^3 \\ \lambda^5 & \cdot & \lambda^2 \\ \lambda^3 & \lambda^2 & \cdot \end{pmatrix}$$

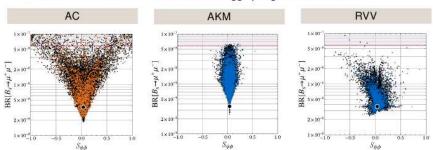
$$\delta_d^{LL} \sim \begin{pmatrix} \cdot & \lambda^5 & \lambda^3 \\ \lambda^5 & \cdot & \lambda^2 \\ \lambda^3 & \lambda^2 & \cdot \end{pmatrix}$$

$$\delta_d^{RR} \sim \begin{pmatrix} \cdot & 0 & 0 \\ 0 & \cdot & 1 \\ 0 & 1 & \cdot \end{pmatrix}$$

$$\delta_d^{RR} \sim \begin{pmatrix} \cdot & 0 & 0 \\ 0 & \cdot & 1 \\ 0 & 1 & \cdot \end{pmatrix} \quad \delta_d^{RR} \sim \begin{pmatrix} \cdot & \lambda^3 & \lambda^3 \\ \lambda^3 & \cdot & \lambda^2 \\ \lambda^3 & \lambda^2 & \cdot \end{pmatrix} \\ \delta_d^{RR} \sim \begin{pmatrix} \cdot & \lambda^3 & \lambda^2 \\ \lambda^3 & \cdot & \lambda \\ \lambda^2 & \lambda & \cdot \end{pmatrix} \\ \delta_d^{RR} \sim \begin{pmatrix} \cdot & 0 & 0 \\ 0 & \cdot & 0 \\ 0 & 0 & \cdot \end{pmatrix}$$

$Br(B_{s} o \mu^{+}\mu^{-})$ vs. $S_{\psi\phi}$

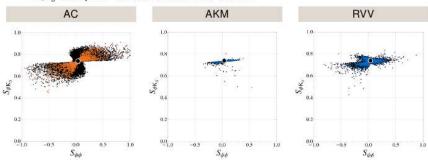
- Both observables can deviate significantly from the SM in all 3 models
- large $S_{\psi\phi} \Rightarrow$ large BR($B_s \rightarrow \mu^+\mu^-$) in the AC and AKM models
- Correlation arises from dominance of Higgs penguin contributions



- Orange points: UT tension solved through contribution to $\Delta M_d/\Delta M_s$
- Blue points: UT tension solved through contribution to ϵ_K
- Scan ranges: $m_0 < 2$ TeV, $M_{1/2} < 1$ TeV, $|A_0| < 3m_0$, $5 < \tan \beta < 55$, O(1) parameters varied within $[\frac{1}{5}, 2]$

$S_{\phi K_S}$ vs. $S_{\psi \phi}$

- In the AC model, both $S_{\phi K_S}$ and $S_{\psi \phi}$ can have large effects, but a simultaneous enhancement of $S_{\psi \phi}$ and suppression of $S_{\phi K_S}$ (as indicated by the data) is impossible
- S_{φKS} nearly SM-like in AKM and RVV models

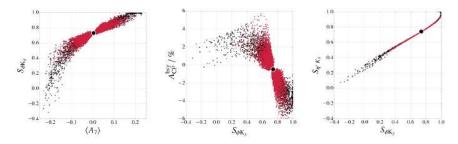


- Orange points: UT tension solved through contribution to $\Delta M_d/\Delta M_s$
- Blue points: UT tension solved through contribution to ϵ_K
- Scan ranges: $m_0 < 2$ TeV, $M_{1/2} < 1$ TeV, $|A_0| < 3m_0$, $5 < \tan \beta < 55$, O(1) parameters varied within $[\frac{1}{2}, 2]$

Model with purely left-handed currents

Pattern of NP effects in the δ LL model:

- No large effects in S_{ψφ}
- Large, correlated effects in $S_{\phi K_S}$, $S_{\eta' K_S}$, $A_{\sf CP}(b o s \gamma)$, $\langle A_{7,8}
 angle$
- ⟨A_{7,8}⟩: T-odd CP asymmetries in B → K*ℓ⁺ℓ⁻



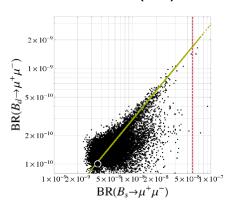
• Scan ranges: $m_0<2$ TeV, $M_{1/2}<1$ TeV, $|A_0|<3m_0,5<\tan\beta<55,$ O(1) parameters varied within $[\frac{1}{2},2]$

$Br(B_s \to \mu^+ \mu^-)$ vs. $Br(B_d \to \mu^+ \mu^-)$

Abelian (AC)

2×10^{-9} BR($B_d^+\mu^+$) 2×10^{-10} 1×10^{-10} $1 \times 10^{-9} \times 10^{-9}$ $5 \times 10^{-9} 1 \times 10^{-8} \times 10^{-8}$ $5 \times 10^{-8} 1 \times 10^{-7}$ $BR(B_s \rightarrow \mu^+ \mu^-)$

Non abelian (RVV)



$$Br(B_s o \mu^+\mu^-)/Br(B_d o \mu^+\mu^-) = |V_{ts}/V_{td}|^2$$
 in MFV models

CPV in D-physics

CPV in $D^0 - \bar{D}^0 \sim ((V_{cb}V_{ub})/(V_{cs}V_{us})) \sim 10^{-3}$ in the SM

$$\mathbf{S_f} = 2\Delta Y_f = \frac{1}{\Gamma_D} \left(\hat{\Gamma}_{\bar{D}^0 \to f} - \hat{\Gamma}_{D^0 \to f} \right)$$

$$\eta_f^{\text{CP}} \mathbf{S}_f = \mathbf{x} \left(\left| \frac{\mathbf{q}}{\mathbf{p}} \right| + \left| \frac{\mathbf{p}}{\mathbf{q}} \right| \right) \sin \phi - \mathbf{y} \left(\left| \frac{\mathbf{q}}{\mathbf{p}} \right| - \left| \frac{\mathbf{p}}{\mathbf{q}} \right| \right) \cos \phi$$

$$\textbf{a}_{\mathrm{SL}} = \frac{\Gamma(D^0 \to K^+\ell^-\nu) - \Gamma(\bar{D}^0 \to K^-\ell^+\nu)}{\Gamma(D^0 \to K^+\ell^-\nu) + \Gamma(\bar{D}^0 \to K^-\ell^+\nu)} = \frac{|q|^4 - |p|^4}{|q|^4 + |p|^4}$$

CPV in D-physics vs. neutron EDM in SUSY

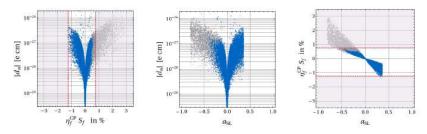


FIG. 3: Correlations between d_n and S_f (left), d_n and a_{SL} (middle) and a_{SL} and S_f (right) in SUSY alignment models. Gray points satisfy the constraints (8)-(10) while blue points further satisfy the constraint (11) from ϕ . Dashed lines stand for the allowed range (18) for S_f .

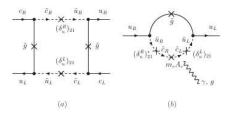


FIG. 2: Examples of relevant Feynman diagrams contributing (a) to $D^0 - \bar{D}^0$ mixing and (b) to the up quark (C)EDM in SUSY alignment models.

"DNA-Flavour Test"

	GMSSM	AC	RVV2	AKM	δ LL	FBMSSM	
$S_{\phi K_S}$	***	***		В	***	***	CunorD
$A_{\sf CP}\left({\sf B} o {\sf X}_{\sf S}\gamma ight)$	***		10	10	***	***	
$B ightarrow K^{(*)} u ar{ u}$	••	1	10	10	100	10	Super
$ au ightarrow \mu \gamma$	***	***	***	ш	***	***	
$D^0 - \bar{D}^0$	***	***		10		и	0
$A_{7,8}(B \to K^* \mu^+ \mu^-)$	***			10.	***	***	vs.
$A_9(B \rightarrow K^* \mu^+ \mu^-)$	***	0		10	88	11	ruch
$S_{\psi\phi}$	***	***	***	***	10	11	THCP
$B_s \rightarrow \mu^+ \mu^-$	***	***	***	***	***	***	
ϵ_{K}	***		***	***	- 11	10	
$K^+ o \pi^+ u ar{ u}$	***	0		10	100	16	
$K_L \rightarrow \pi^0 \nu \bar{\nu}$	***				100	11	
$\mu ightarrow oldsymbol{e} \gamma$	***	***	***	***	***	***	
$\mu + N \rightarrow e + N$	***	***	***	***	***	***	i
d _n	***	***	***	***		***	
d _θ	***	***	***		100	***	
$(g-2)_{\mu}$	***	***	***		***	***	

G. Isidori – Flavour Physics now and in the LHC era

Flavour physics in the LHC era

LHC [high p_T]

A unique effort toward the high-energy frontier



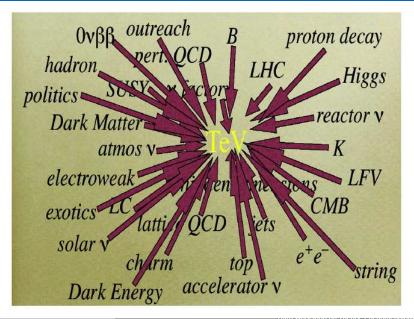
to determine the energy scale of NP

Flavour physics



A collective effort toward the high-intensity frontier [to determine the flavour structure of NP]

Murayama's view



Masiero's view

DM - FLAVOR A MAJOR for DISCOVERY LEAP AHEAD IS NEEDED and/or FUND. TH. NEW RECONSTRUCTION. PHYSICS AT THE ELW SCALE "LOW ENERGY" DARK MATTER PRECISION PHYSICS $m_y n_y \sigma_y \dots$ FCNC, CP \neq , (g-2), ($\beta\beta$)_{0vv} LINKED TO COSMOLOGICAL EVOLUTION

Possible interplay with dynamical DE