Semileptonic D decays and **Weak Annihilation**

Paolo Gambino Università di Torino



In collaboration with Jernej Kamenik

Paolo Gambino Capri 5/7/2010

The Unitarity Triangle





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3

~1.60

UT_{fit}

6 σ

5

3

2

1.4^{__0}

B_k

1.2

1

$|V_{ub}|$ in the kinetic scheme -GGOU

PG,Giordano,Ossola,Uraltsev

Good consistency & small th error. OPE in a scheme with Wilsonian IR cutoff ~IGeV, all subleading I/m_b and $O(\alpha_s^2\beta_0)$ terms consistently included, careful treatment of high q² tail.

Inputs from global fit to the moments

+6.3-7.0% total error

very strong dependence on m_b , twice larger than in total rate



A global comparison arXiv:0907.5386



- Not all observables are equally clean. eg high q² tail is sensitive to WA
- More inclusive measurements, less dependence on mb
- Spectra can help discriminate models. Vary the cuts.
- Theory errors are partly parametric: mb dependence is crucial

HFAG 2010 Average |V_{ub}|x10³ GGOU 4.27(16)_{ex}+15₋₂₁ **2009** preliminary Belle Multivariate analysis only E_I>IGeV

 $|V_{ub}| = (4.45 \pm 0.26^{+0.13}_{-0.22}) \times 10^{-3}$

2.1 σ from excl, 2.5 σ from UTFit probably a bit less after fit upgrade

Includes about 90% of the rate: really <u>inclusive</u> measurement, no need for SF. Crucial input m_b, needs to be confirmed!

I.8σ from B→πIV (MILC-FNAL) 2.5σ from UTFit (because of sin2β)

NEW PHYSICS? eg LR models Chen, Nam

Weak Annihilation

Spectator dependent non-pert contribution localized at max E_1 (or max q^2) Bigi, Uraltsev 1993

In principle affects B^+ only but WA mixes with Darwin operator at O(1). Isosinglet component can be as large as isotriplet

Difficult to study on the lattice, can be constrained experimentally Rosner et al [Cleo coll]

WA may pollute all present inclusive determinations of V_{ub} and more severely the less inclusive ones (E₁ endpoint, high q²)





WA in the OPE

WA appears naturally in the OPE for total rates as a combination of 4-fermion dim-6 op

$$\mathcal{O}_1^{Qq} = \bar{Q}\gamma_\mu (1-\gamma_5)q\,\bar{q}\gamma^\mu (1-\gamma_5)Q$$
$$\mathcal{O}_2^{Qq} = \bar{Q}(1-\gamma_5)q\,\bar{q}(1-\gamma_5)Q$$

Its B meson exp value vanishes in the limit of factorization, but it is enhanced by the coefficient

 $B_{\rm WA}(\mu_{\rm WA}) = \langle B|O_1 - O_2|B\rangle$

Because of the mixing with Darwin op, the WA m.e. depends on μ_{WA}

$$\delta\Gamma \sim \left[C_{\rm WA} B_{\rm WA}(\mu_{\rm WA}) - \left(8\ln\frac{m_b^2}{\mu_{\rm WA}^2} - \frac{77}{6} \right) \frac{\rho_D^3}{m_b^3} + \mathcal{O}(\alpha_s) \right] \qquad C_{\rm WA} = 32\pi^2/m_b^3$$



New Cleo-c results

arXiv:0912.4232

Recently determined experimentally

B(D⁺ → Xeν) = (16.13 ± 0.20 ± 0.33)%
B(D⁰ → Xeν) = (6.46 ± 0.17 ± 0.13)%

Similar results for muons
Very recently results also for D_s decays

B(D_s → Xeν) = (6.52±, 0.39 ± 0.15)%

D_s and D₀ rates differ significantly $\Gamma(D^+ \to Xe^+\nu)/\Gamma(D^0 \to Xe^+\nu) = 0.985(28)$ $\Gamma(D_s^+ \to Xe^+\nu)/\Gamma(D^0 \to Xe^+\nu) = 0.828(57)$

Valence WA Cabibbo suppressed in D⁺, absent in D⁰, is it a sign of WA? Bigi, Mannel, Turczyk, Uraltsev 0911.3322

Ligeti,Luke,Manohar 1003.1351

New Cleo-c results (II) arXiv:0912.4232

 Cleo also measured the electron spectra for p>0.2GeV in the lab

- We have extrapolated them to p=0, computed their first moments, and boosted to the D rest frame
- We cannot compute spectra, but moments should follow OPE.



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New Cleo-c results (III)

$$\begin{split} \langle E_{\ell} \rangle_{exp}^{D^0} &= 0.459(3) \text{GeV} \,, \\ \langle E_{\ell} \rangle_{exp}^{D^+} &= 0.455(1) \text{GeV} \,, \\ \langle E_{\ell} \rangle_{exp}^{D_s} &= 0.456(11) \text{GeV} \,, \end{split}$$

$$\begin{split} \langle E_\ell^2 \rangle_{exp}^{D^0} &= 0.240(2) {\rm GeV}^2 \,, \\ \langle E_\ell^2 \rangle_{exp}^{D^+} &= 0.236(1) {\rm GeV}^2 \,, \\ \langle E_\ell^2 \rangle_{exp}^{D_s} &= 0.239(12) {\rm GeV}^2 \,, \end{split}$$

JF Kamenik, PG, arXiv:1004.0114

$$\langle (E_{\ell} - \langle E_{\ell} \rangle)^2 \rangle_{exp}^{D^0} = 0.029(2) \text{GeV}^2 ,$$
$$\langle (E_{\ell} - \langle E_{\ell} \rangle)^2 \rangle_{exp}^{D^+} = 0.029(1) \text{GeV}^2 ,$$
$$\langle (E_{\ell} - \langle E_{\ell} \rangle)^2 \rangle_{exp}^{D_s} = 0.031(12) \text{GeV}^2$$

No evidence for spectator effects! Is there really evidence for WA?

SU(3) violation in charm

• Decay constants on the lattice: $f_D=260(10)MeV$ vs $f_{Ds}=217(10)MeV$ Bazavov et al

- Hyperfine splittings $\Delta_{D_q}^{hf} = 3(m_{D_q^*}^2 m_{D_q}^2)/4 \approx \mu_G^2$ $\Delta_{D^+}^{hf} = 0.409(1) \text{GeV}^2, \quad \Delta_{D^0}^{hf} = 0.413(1) \text{GeV}^2, \quad \Delta_{D_q}^{hf} = 0.440(2) \text{GeV}^2$
 - SU(3) violation can be as large as 20%. Widths get much larger power corrections than moments and this might partially explain the observed width difference without WA

OPE expansion for charm

Pole mass scheme, $m_c=1.6$ GeV, $\mu_{WA}=0.8$ GeV

$$\begin{split} & \Gamma = \Gamma_0 \left[1 - 0.72 \,\alpha_s - 0.29 \,\alpha_s^2 \beta_0 - 0.60 \,\mu_G^2 - 0.20 \,\mu_\pi^2 + 0.42 \,\rho_D^3 + 0.38 \,\rho_{LS}^3 + 80 B_{\rm WA}^{(0)} \right] \,, \\ & \langle E \rangle = \langle E \rangle_0 \left[1 - 0.03 \,\alpha_s - 0.03 \,\alpha_s^2 \beta_0 - 0.07 \,\mu_G^2 + 0.20 \,\mu_\pi^2 + 1.4 \,\rho_D^3 + 0.29 \,\rho_{LS}^3 + 135 \bar{B}_{\rm WA}^{(1)} \right] \\ & \langle E^2 \rangle = \langle E^2 \rangle_0 \left[1 - 0.07 \,\alpha_s - 0.05 \,\alpha_s^2 \beta_0 - 0.14 \,\mu_G^2 + 0.52 \,\mu_\pi^2 + 3.5 \,\rho_D^3 + 0.66 \,\rho_{LS}^3 + 204 \bar{B}_{\rm WA}^{(2)} \right] \\ & \sigma_E^2 = (\sigma_E^2)_0 \left[1 - 0.09 \,\alpha_s - 0.05 \,\alpha_s^2 \beta_0 - 0.14 \,\mu_G^2 + 1.7 \,\mu_\pi^2 + 9.4 \,\rho_D^3 + 1.4 \,\rho_{LS}^3 + 641 \bar{B}_{\rm WA}^{(\sigma)} \right] \,, \end{split}$$

Poor convergence for rates. Moments have smaller pert and power corrections, but are still quite sensitive to WA

OPE expansion for charm

$$\begin{aligned} & \text{Kinetic scheme, } \mu = 0.5 \text{GeV, } m_c = 1.4 \text{GeV, } \mu_{\text{WA}} = 0.8 \text{GeV} \\ & \Gamma_{kin} = 1.2(3)10^{-13} \text{GeV} \left\{ 1 + 0.23 \,\alpha_s + 0.18 \,\alpha_s^2 \beta_0 - 0.79 \,\mu_G^2 - 0.26 \mu_\pi^2 + 1.45 \,\rho_D^3 + 0.56 \rho_{LS}^3 + 120 B_{\text{WA}}^{(0)} \right\} \\ & \langle E_\ell \rangle_{kin} = 0.415(21) \text{GeV} \left\{ 1 + 0.03 \,\alpha_s + 0.02 \,\alpha_s^2 \beta_0 - 0.09 \,\mu_G^2 + 0.26 \mu_\pi^2 + 2.7 \rho_D^3 + 0.44 \rho_{LS}^3 + 203 \bar{B}_{\text{WA}}^{(1)} \right\} , \\ & \langle E_\ell^2 \rangle_{kin} = 0.192(20) \text{GeV}^2 \left\{ 1 + 0.001 \,\alpha_s + 0.02 \,\alpha_s^2 \beta_0 - 0.18 \,\mu_G^2 + 0.68 \mu_\pi^2 + 6.6 \rho_D^3 + 0.99 \rho_{LS}^3 + 307 \bar{B}_{\text{WA}}^{(2)} \right\} , \\ & \sigma_{E,kin}^2 = 0.019(2) \text{GeV}^2 \left\{ 1 - 0.53 \,\alpha_s - 0.17 \,\alpha_s^2 \beta_0 - 0.18 \mu_G^2 + 2.2 \mu_\pi^2 + 17 \rho_D^3 + 2.1 \rho_{LS}^3 + 961 \bar{B}_{\text{WA}}^{(\sigma)} \right\} , \end{aligned}$$

In kinetic scheme, better convergence for low $0.5 < \mu < 0.8$ GeV. D-expectation values related to those in B by heavy quark symmetry Use B->X_c lv moments fit for m_c and OPE parameters with enlarged errors

$$\Gamma_{\rm WA}^{(n)}(D^0) \propto \cos^2 \theta_c B_{\rm WA}^{(n),s}(D^0) + \sin^2 \theta_c B_{\rm WA}^{(n),d}(D^0) , \Gamma_{\rm WA}^{(n)}(D^+) \propto \cos^2 \theta_c B_{\rm WA}^{(n),s}(D^+) + \sin^2 \theta_c B_{\rm WA}^{(n),d}(D^+) , \Gamma_{\rm WA}^{(n)}(D_s) \propto \cos^2 \theta_c B_{\rm WA}^{(n),s}(D_s) + \sin^2 \theta_c B_{\rm WA}^{(n),d}(D_s) ,$$

Cabibbo suppressed contributions generally negligible We can extract the singlet and "valence" contributions

WA dilution

Expect perturbative and nonperturbative smearing of WA. It is both α_s and $1/m_c$ suppressed.

The exclusive upper threshold is $2\pi \text{ or } \eta$

The smearing is <u>model</u> <u>dependent</u> and generally dilutes the WA contribution to higher E_1 moments, although it is generally a modest effect



Results

JF Kamenik, PG 1004.0114

Allowing for 20% SU(3) violation in the OPE parameters

 $\Delta B_{\rm WA}^{(0),s} \equiv B_{\rm WA}^{(0),s}(D_s) - B_{\rm WA}^{(0),s}(D^0) = -0.0014(12)(5) \,{\rm GeV}^3$

 $\Delta \bar{B}_{WA}^{(1),s} = 0.0000(3)(1) \text{GeV}^3$

 $\Delta \bar{B}_{WA}^{(2),s} = 0.0000(4)(2) \text{GeV}^3$

In worst dilution scenario the moments alone give (linearly adding errors)

Singlet component



Valence component always compatible with zero

$$\Delta B_{\rm WA}^{(0),s} = 0.0000(12)(3)$$

$$B_{\rm WA}^s = -0.0003(25) {\rm GeV}^3$$

equivalent to 30% error on rate

In principle, one can even probe dilution in data by comparing WA contributions to different moments resolution is now insufficient

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Implications for Vub

A factor 2.5 in going to B neglecting running and power corrections

$$B_{\mathrm{WA}}^{bq}(\mu_{\mathrm{WA}}) = \frac{m_B f_B^2}{m_D f_D^2} B_{\mathrm{WA}}^{cq}(\mu_{\mathrm{WA}})$$

Singlet component $|B_{wA}^b(\mu_{wA} = 0.8 \text{GeV})| \lesssim 0.006 \text{GeV}^3$

Valence component $-0.004 \,\mathrm{GeV}^3 \lesssim \Delta B^b_{wA}(\mu_{wA} = 0.8 \mathrm{GeV}) \lesssim 0.002 \,\mathrm{GeV}^3$

In GGOU V_{ub} analysis we had max B_{WA}=0.02. Factor 3 improvement. **Max 1% effect on V_{ub}** for most inclusive analyses B⁰ and B⁺ inclusive widths should not differ more than about 1%

Compatible results from widths only in Ligeti, Luke, Manohar, 1003.1351 but moments are cleaner than rates and our errors more conservative

Conclusions

- OPE seems to work surprisingly well for D decays
- We find no evidence for WA, mild indication in the widths
- WA uncertainty in inclusive $|V_{ub}|$ is reduced by factor 3
- Additional tests possible at Cleo and BES-III