

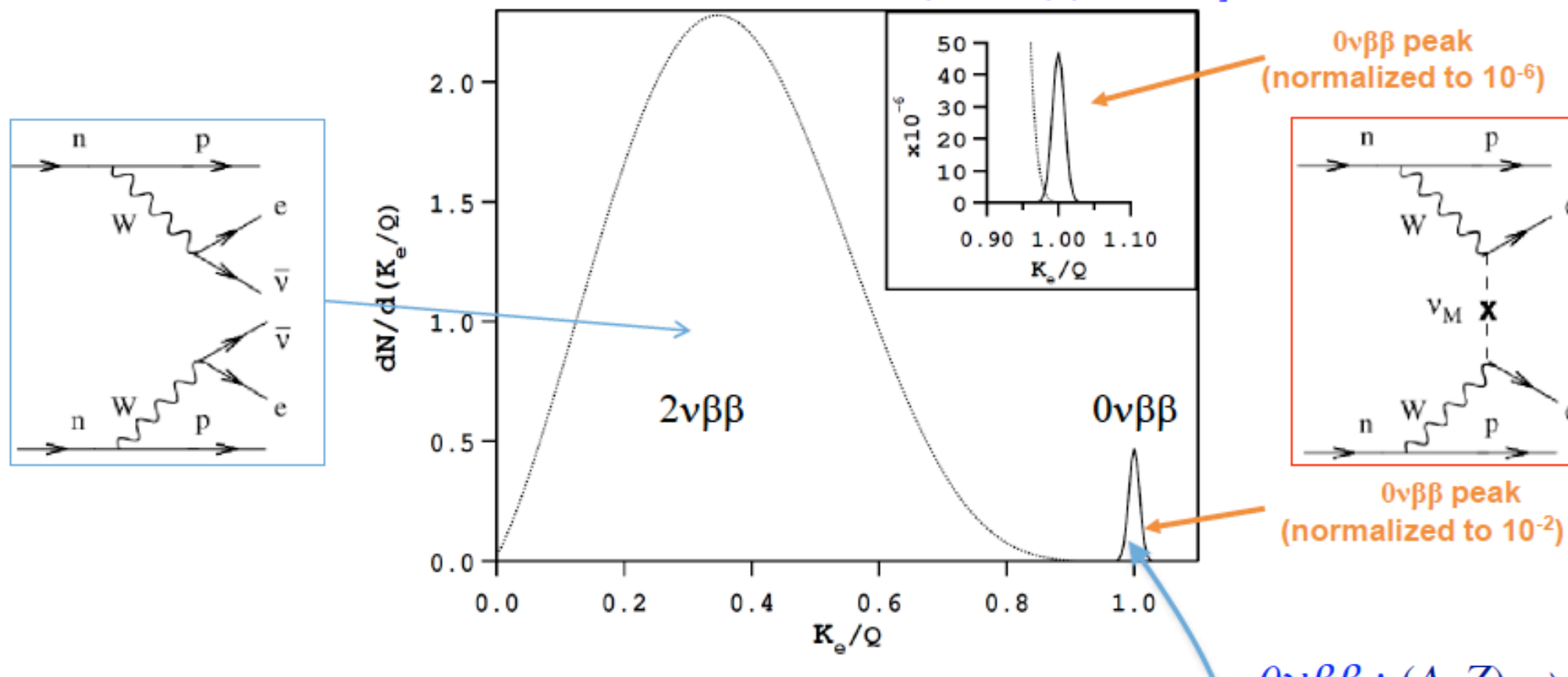
Investigation of neutrino nature with LEGEND

G.Salamanna

$0\nu 2\beta$ decays

$$\Delta L=2$$

[arXiv:hep-ph/0611243]



- Two beta decays at the same time
- Only a few isotopes able to undergo 2β

• among which **Ge76**

$0\nu\beta\beta : (A, Z) \rightarrow (A, Z+2) + 2e^-$
new physics, $T_{1/2} > 10^{25}$ yrs

Only feasible way at present to determine ν nature:

$\nu = \bar{\nu} \Rightarrow$ Majorana particle

$$(T_{1/2}^{0\nu})^{-1} = G^{0\nu}(Q_{\beta\beta}, Z) |M^{0\nu}|^2 \left(\frac{\langle m_{ee} \rangle}{m_e} \right)^2$$

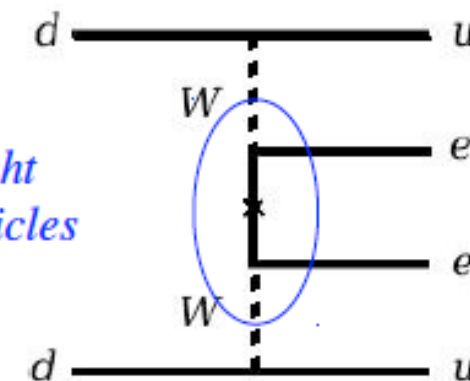
phase space factor

nuclear matrix element

effective Majorana neutrino mass

$\langle m_{ee} \rangle = \left| \sum_i U_{ei}^2 m_i \right|$

exchange of light Majorana particles



- Has also important cosmological repercussions

Challenges of double- β decay experiments

$$\text{Sensitivity on } T_{1/2} \propto \varepsilon \cdot A \cdot \sqrt{\frac{M \cdot T}{b \cdot \Delta E}}$$



ε detection efficiency

A isotopic abundance

M active mass

T exposure

b background rate

ΔE energy resolution

$T_{1/2}^{0\nu} > 10^{25}$ years !!

→ Need:

- high target mass
- high exposure
- low background rate
- good energy resolution

Possible **backgrounds** from:

External:

- γ from Th and U chain
- neutrons
- μ from cosmic rays (prompt and delayed)

Internal:

- cosmogenic ^{60}Co ($T_{1/2}=5.3$ yr)
- cosmogenic ^{68}Ge ($T_{1/2}=271$ d)
- Radioactive surface contaminations

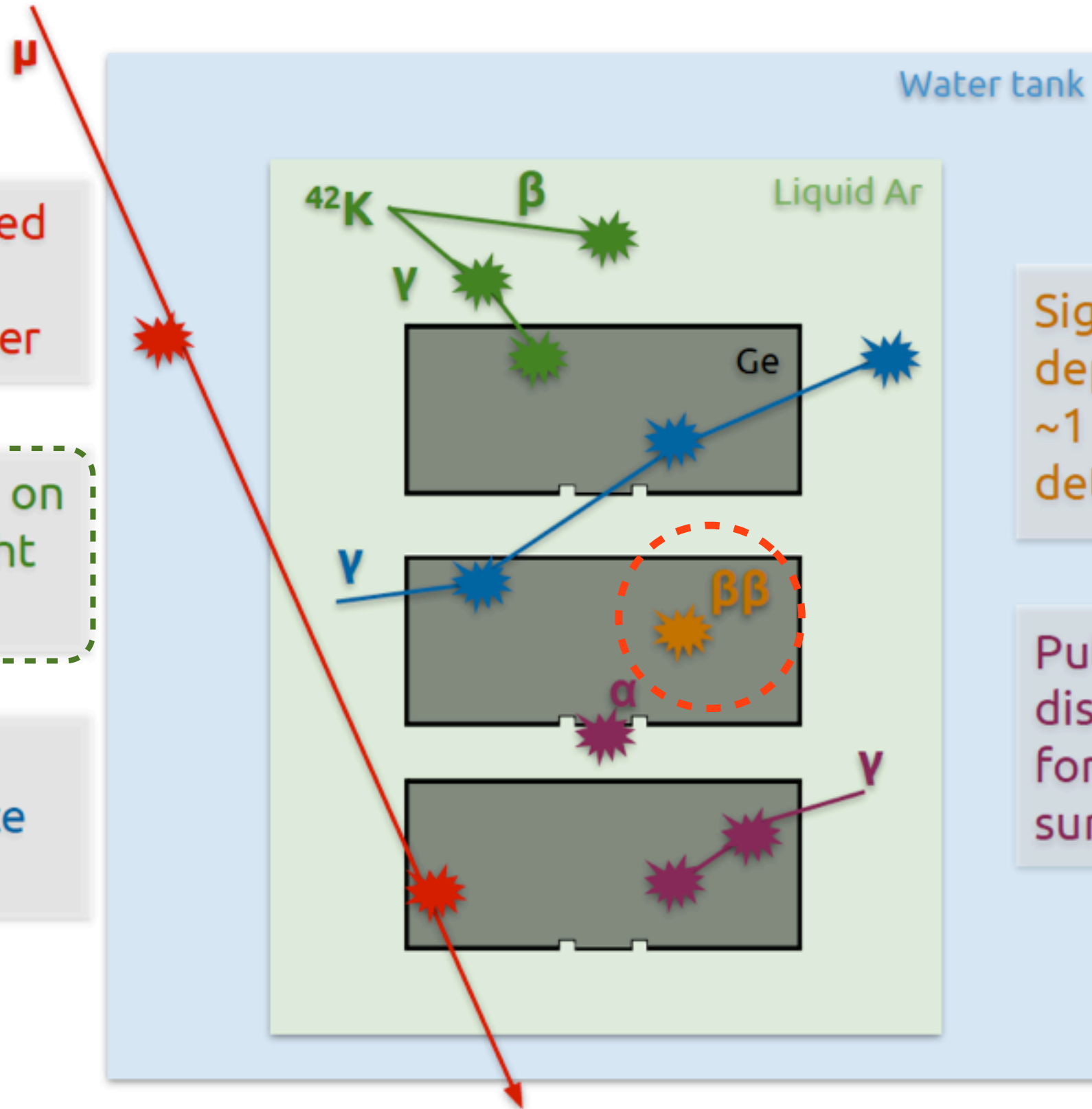
- + **Ge**: source + detector: high ε
- + has best ΔE on the market
- expensive to get to high A for Ge76
- many bkg around $Q(\text{bb, Ge})$

BACKGROUND REDUCTION CONCEPT

Muon veto based on Cherenkov photons in water

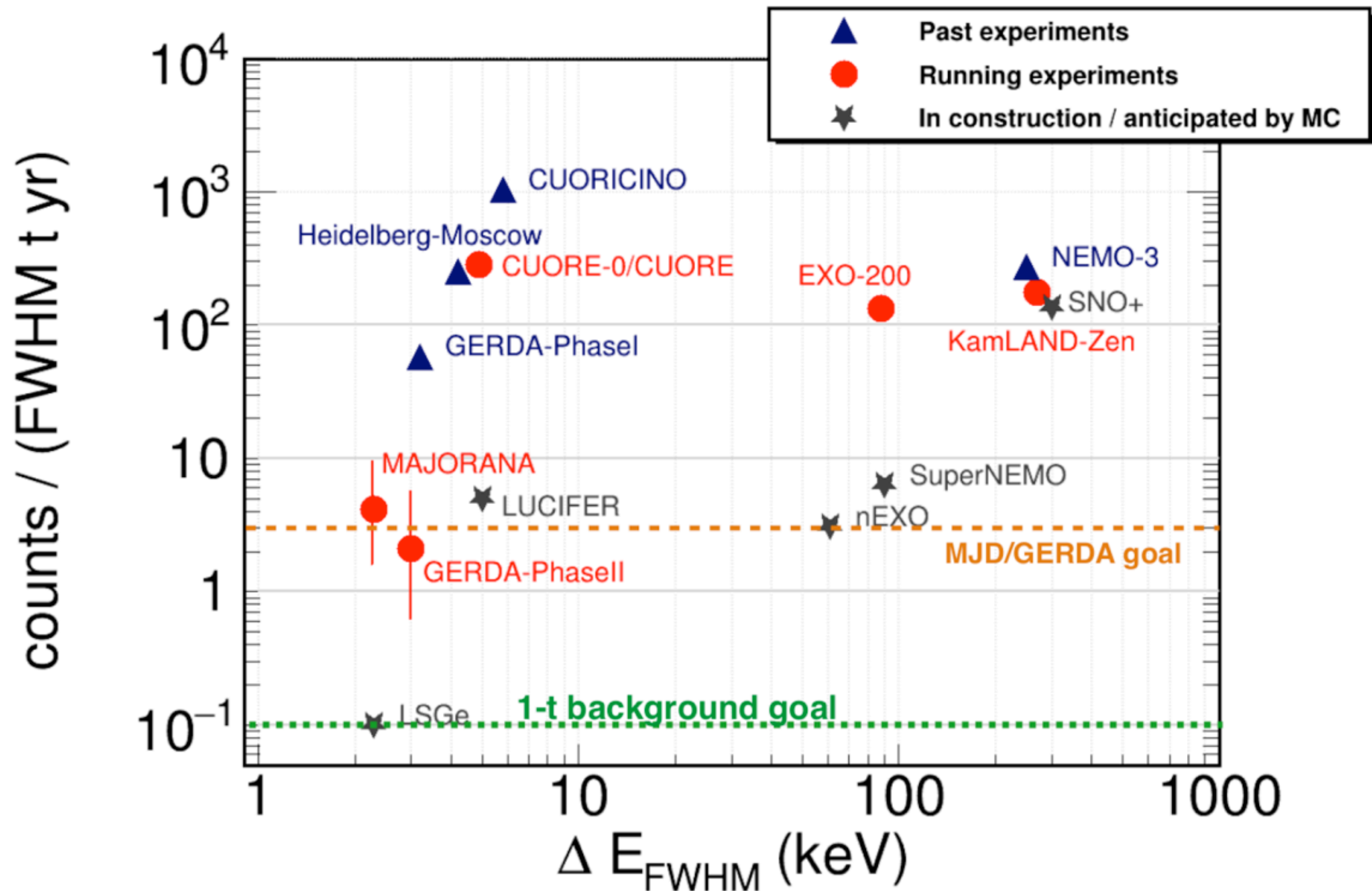
LAr veto based on scintillation light from γ s and β s

Detector anti-coincidence of the array



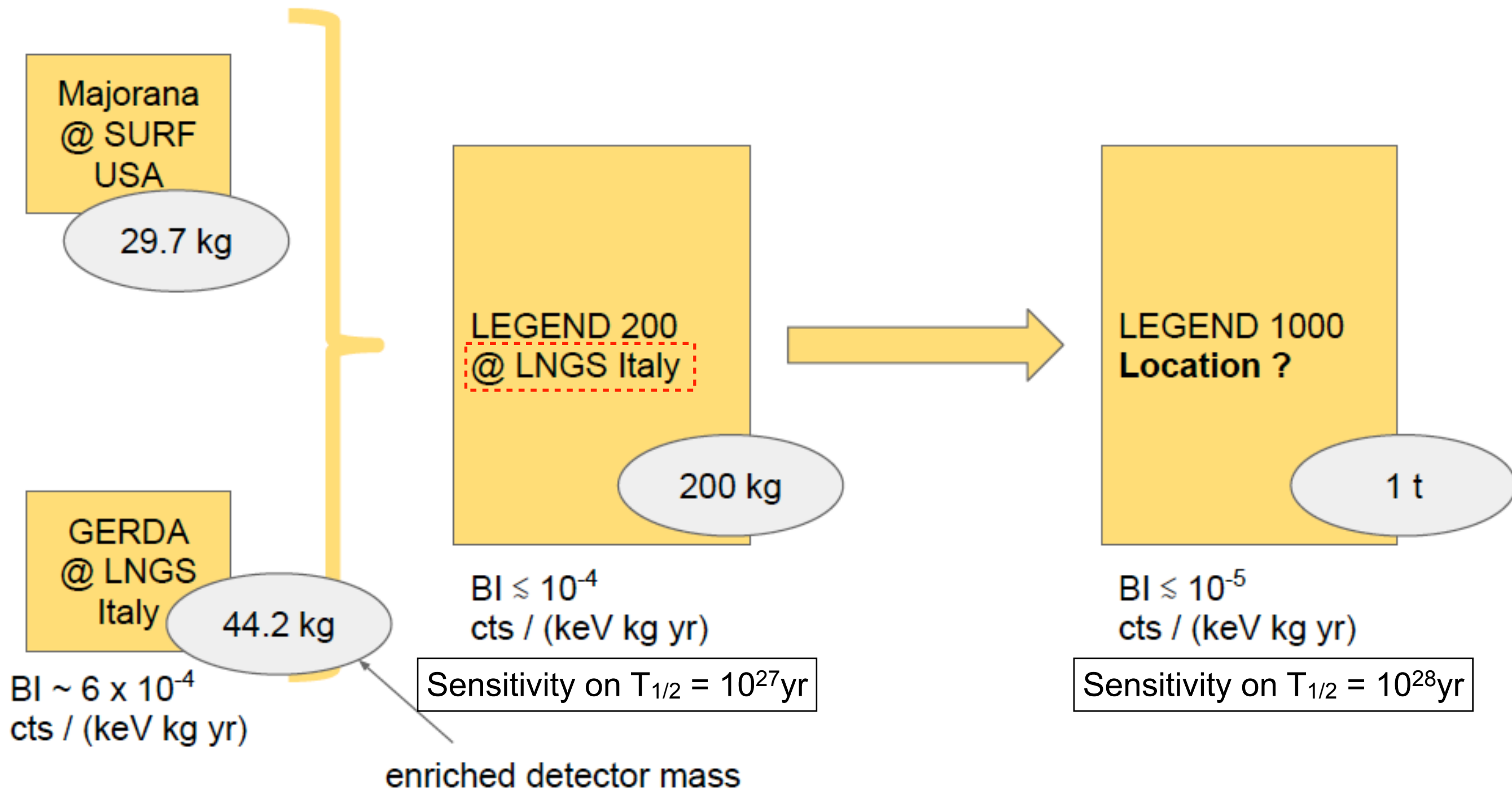
Signal: single energy deposition within $\sim 1 \text{ mm}^3$ of the detector volume

Pulse shape discrimination (PSD) for multi-site and surface events



Ge76 experiments aggressive in keeping background low

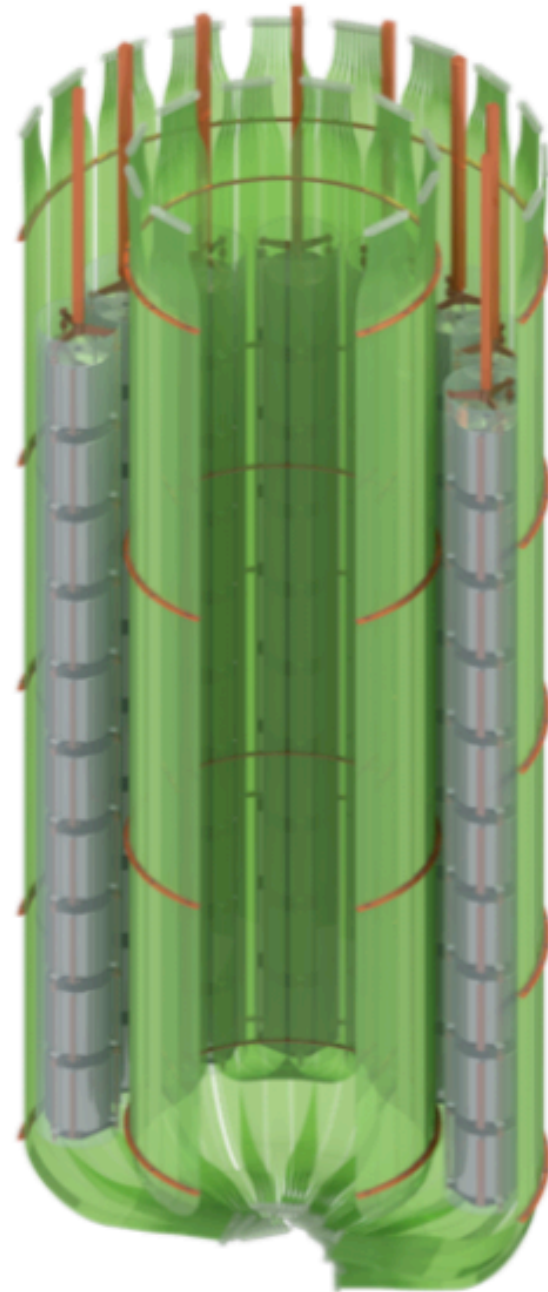
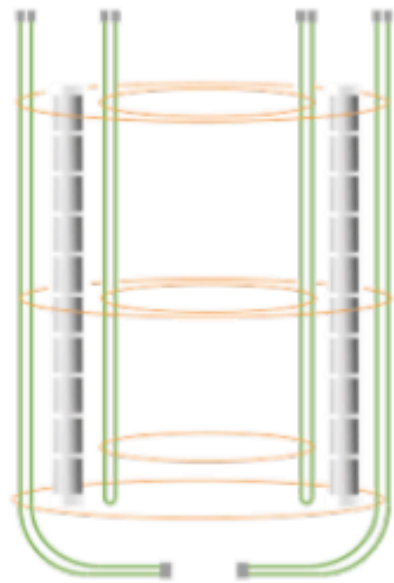
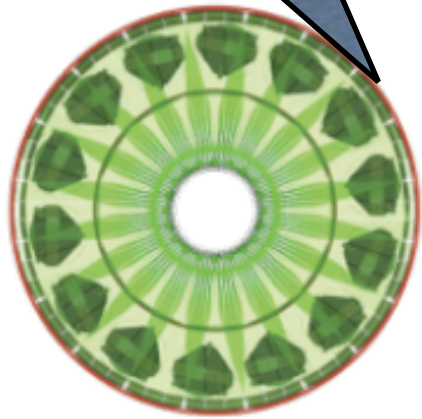
The Large Enriched Germanium Experiment for Neutrinoless $\beta\beta$ decay - LEGEND



Planned local activities on LEGEND LAr veto

LEGEND-200 LAr veto

LEGEND-200
prelim. design



SiPM top & bottom readout for outer curtain
Removable segments to install central strings

- An active veto is one of main resources to reach such low bkg
- LAr veto improved in coverage and read-out wrt Gerda
- Veto efficiency is then introduced in simulations to estimate expected BI close to signal region
- important to response know well:
 - attenuation length at desired purity
 - light yield
 - fast vs slow component
 - reflectivity of immersed materials
- room in LEGEND-1000 for R&D to improve veto capability

LEGEND-200: Activities/ I

- Study optical properties of pure samples of LAr for LEGEND-200
 - cryogenic conditions: Ar is liquid at $T=87$ K
 - expected $\lambda_{\text{att}} \sim 50\text{-}100$ cm, but large discrepancies in available measurements such that more studies are necessary
- ✓ coherent with local activities on liquid scintillators (e.g. JUNO) and should be able to profit from dedicated cryo equipment
 - ⊙ + integrations to keep cold and monitor larger volumes: seeking advise of cryo experts at LNGS on set-up design
 - ⊙ + best set-up: there are various ways to make a λ_{att} measurement; we need to consider systematic light collection effects and space/funds constraints
- [People](#): Diego T and Giuseppe S + LNGS LEGEND people
- Tune Geant-4 LAr optical response (+smart speed optimization on specific bkg processes?) to model full LAr veto
 - [People](#): Andrea A

LEGEND-200: Activities/2

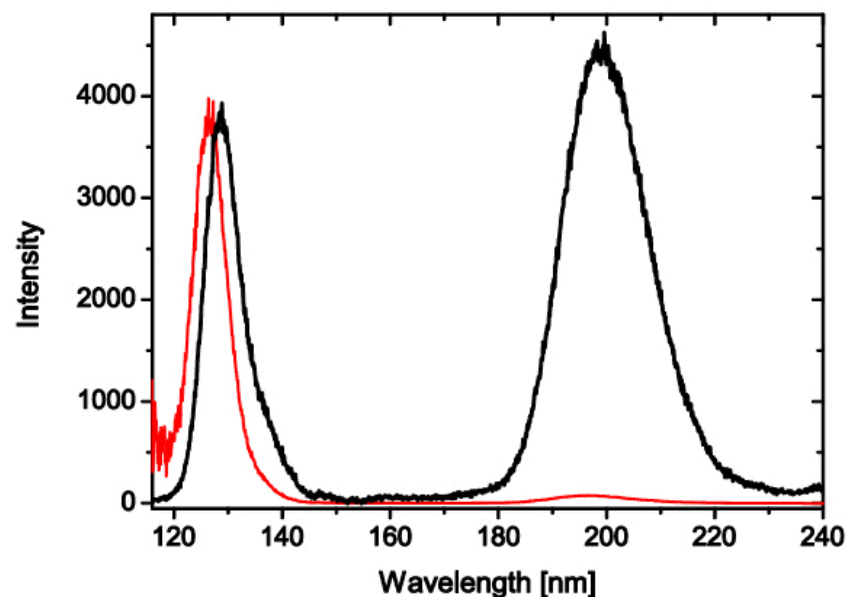
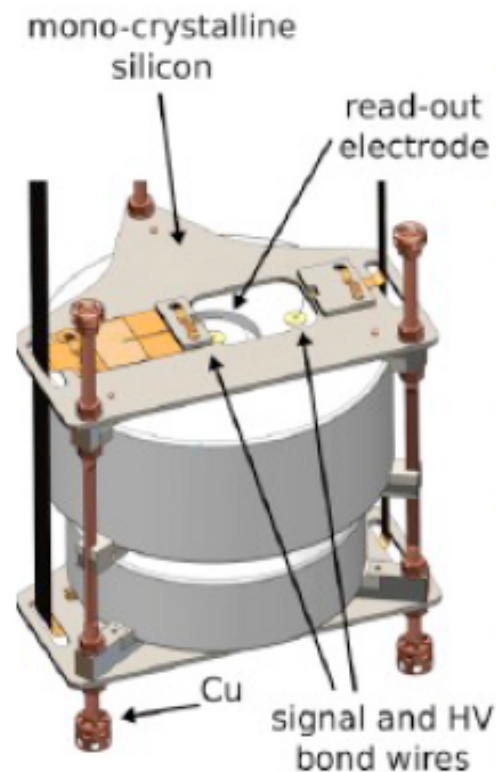
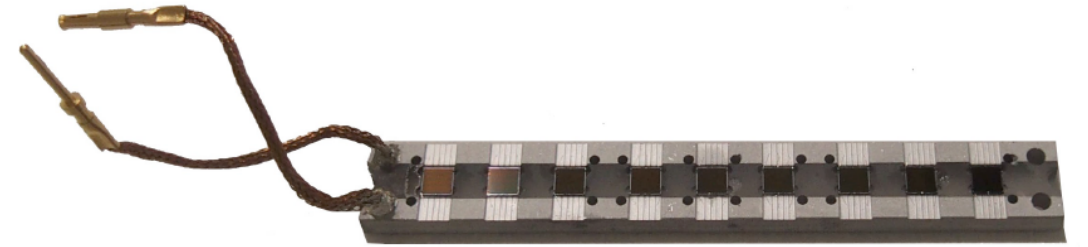


Fig. 5. Emission spectrum of unpurified liquid argon (85K) with (thin line) and without (thick line) sensitivity correction of our setup. The broad emission structure around 197 nm may be attributed to an oxygen impurity [12, 24].

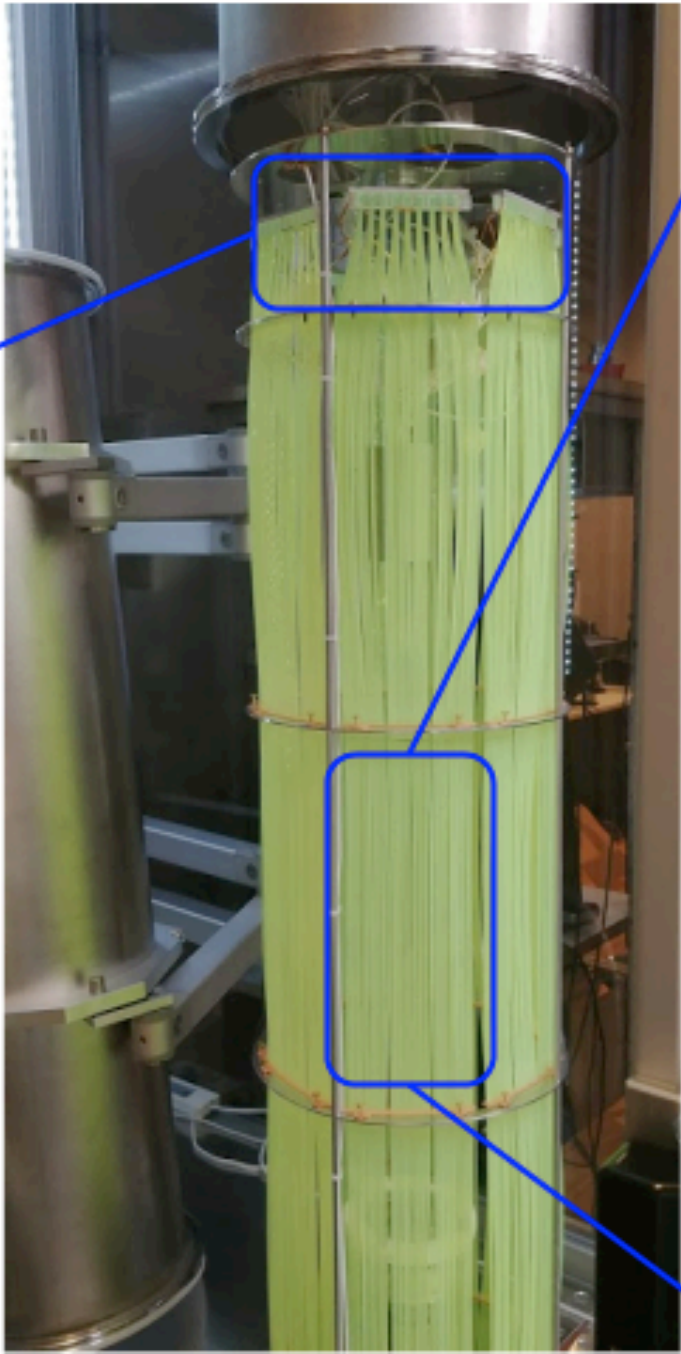
- Materials immersed in LAr (detectors, support structures) reflect scintillation photons
- need to know reflectivity down to VUV ($\lambda_{\text{LAr}} = 128 \text{ nm}$) to model overall light collection in veto
- many old measurements of reflectivity, but it heavily depends on surface
 - e.g. polishing of Ge detectors + incident angle will change reflectivity
- Scouting to do campaign of dedicated measurements with vacuum optical set-up; e.g. at synchrotron radiation from 128 to ~500 nm with Ge, Si, Cu samples
- [People](#): Enrico B, Giuseppe S

Activities/3

Electronics + anti-muon trigger



- 9 SiPM = 81 fiber signals in parallel are collected and amplified by Cremat CR-112
 - Overall 52 channels plugged to 8-channel NIM boards for FE (FADC for digitization) together with Ge detectors
- a. Can “packaging” be improved (change card output to board and group more effectively)?
- b. Study possibility to introduce “fast-trigger” directly on analog signals using topology to veto $Ar40 + n \rightarrow Ar41 + \gamma$ (n from cosmogenic muons close to Ge detectors)
- [People](#): Diego + Giuseppe + TUMunich people for simulations



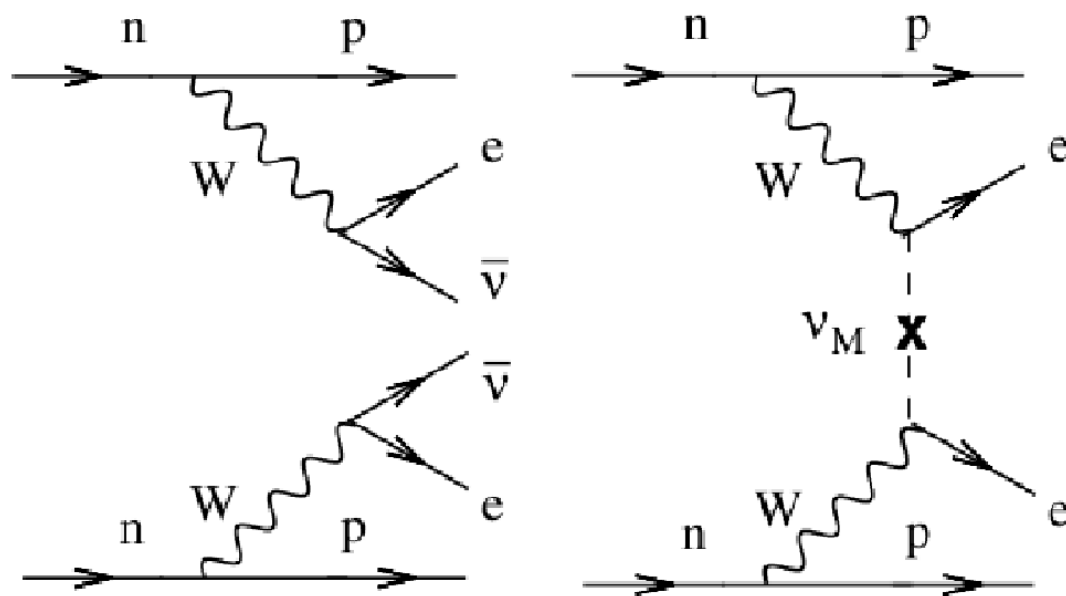
LEGEND-1000

- This opens door to LEGEND-1000
- In R&D phase, with chance to study solutions to attain VERY challenging background levels (really the issue!)
- Identifying some ideas that can be pursued:
 - Screen off incoming radioactivity in double way: “mechanically” with low-rad plastic shrouds + “electrically” against K42 ions that are created close to, and get drifted towards, Ge p-n junction
 - better light yield, ΔE and PSD by doping Ar with materials like Xe *in the right concentrations*
 - shifts wavelength efficiently, suppresses slow and fast component differently + has different peaks with changing intensity based on bkg particle mass
 - small needed concentrations of Xe (10 to 1000 ppm) at 87K a challenge

Back up

$0\nu 2\beta$ decays

$$\Delta L=2$$



- Two beta decays at the same time
- Only a few isotopes able to undergo 2β
 - those with peculiar energy level arrangements such that the emission is more convenient than the unstable isotope

$$0\nu\beta\beta : (A, Z) \rightarrow (A, Z+2) + 2e^-$$

new physics, $T_{1/2} > 10^{25}$ yrs

$$2\nu\beta\beta : (A, Z) \rightarrow (A, Z+2) + 2e^- + 2\nu_e$$

2nd order process, observed, $T_{1/2} \sim 10^{19}$ - 10^{24} yrs

^{76}Ge : $T_{1/2} \sim 10^{21}$ yrs

$$\left(T_{1/2}^{0\nu}\right)^{-1} = G^{0\nu}(Q_{\beta\beta}, Z) |M^{0\nu}|^2 \left(\frac{\langle m_{ee} \rangle}{m_e}\right)^2$$

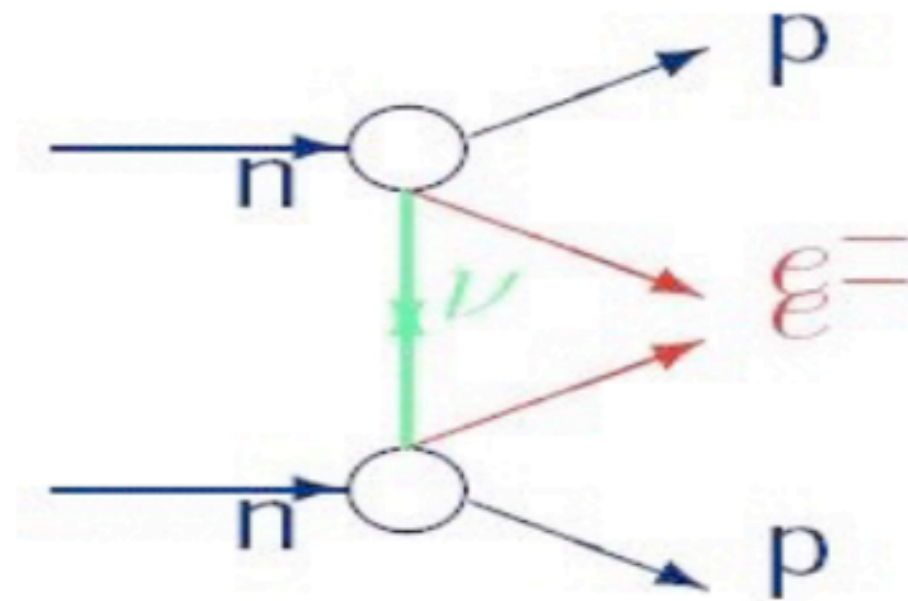
phase space factor
nuclear matrix element

$$\langle m_{ee} \rangle = \left| \sum_i U_{ei}^2 m_i \right|$$

effective Majorana
neutrino mass

Neutrinoless double beta decay

Neutrinoless double beta decay, $(A, Z) \rightarrow (A, Z+2) + 2 e^-$, tests the nature of neutrinos. It violates L by 2 units.



The half-life time depends on neutrino properties

$$\left[T_{0\nu}^{1/2}(0^+ \rightarrow 0^+) \right]^{-1} \propto |M_F - g_A^2 M_{GT}|^2 |\langle m \rangle|^2$$

$$|\langle m \rangle| \equiv |m_1 |U_{e1}|^2 + m_2 |U_{e2}|^2 e^{i\alpha_{21}} + m_3 |U_{e3}|^2 e^{i\alpha_{31}}|,$$

Mixing angles (known)

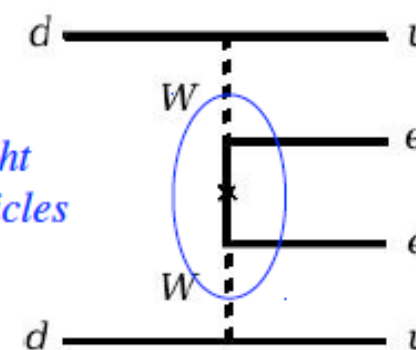
CPV phases (unknown)

Why look for it...

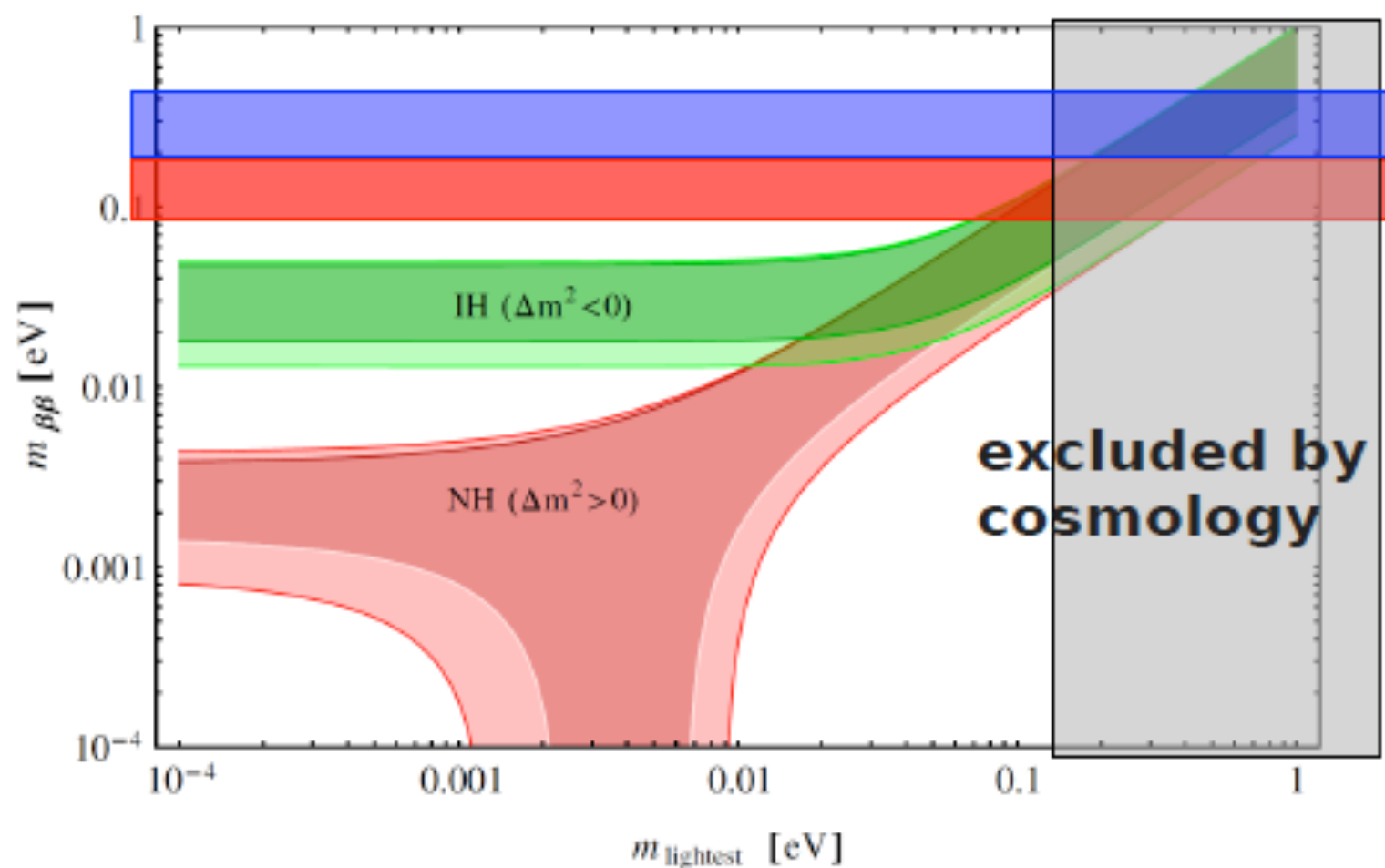
- ◆ Only way to determine if neutrino is its own antiparticle:

$$\nu = \bar{\nu} \Rightarrow \text{Majorana particle}$$

exchange of light
Majorana particles



S. Dell'Oro, S. Marcocci, F. Vissani, PRD 90 (2014)



- A weighted average of neutrino masses enters the decay rate (decay half life)

• **NB: experiments measure $T_{1/2}^{0\nu}$**

- Limits on m_{ee} from above, can also exclude IH

- because one is still looking only at electron flavour, therefore a mix of mass eigenstates, with their MH arrangement in $\langle m_{ee} \rangle$

- nuclear matrix element uncertainties are the biggest spoiler in the conversion

$$\left(T_{1/2}^{0\nu}\right)^{-1} = G^{0\nu}(Q_{\beta\beta}, Z) |M^{0\nu}|^2 \left(\frac{\langle m_{ee} \rangle}{m_e}\right)^2$$

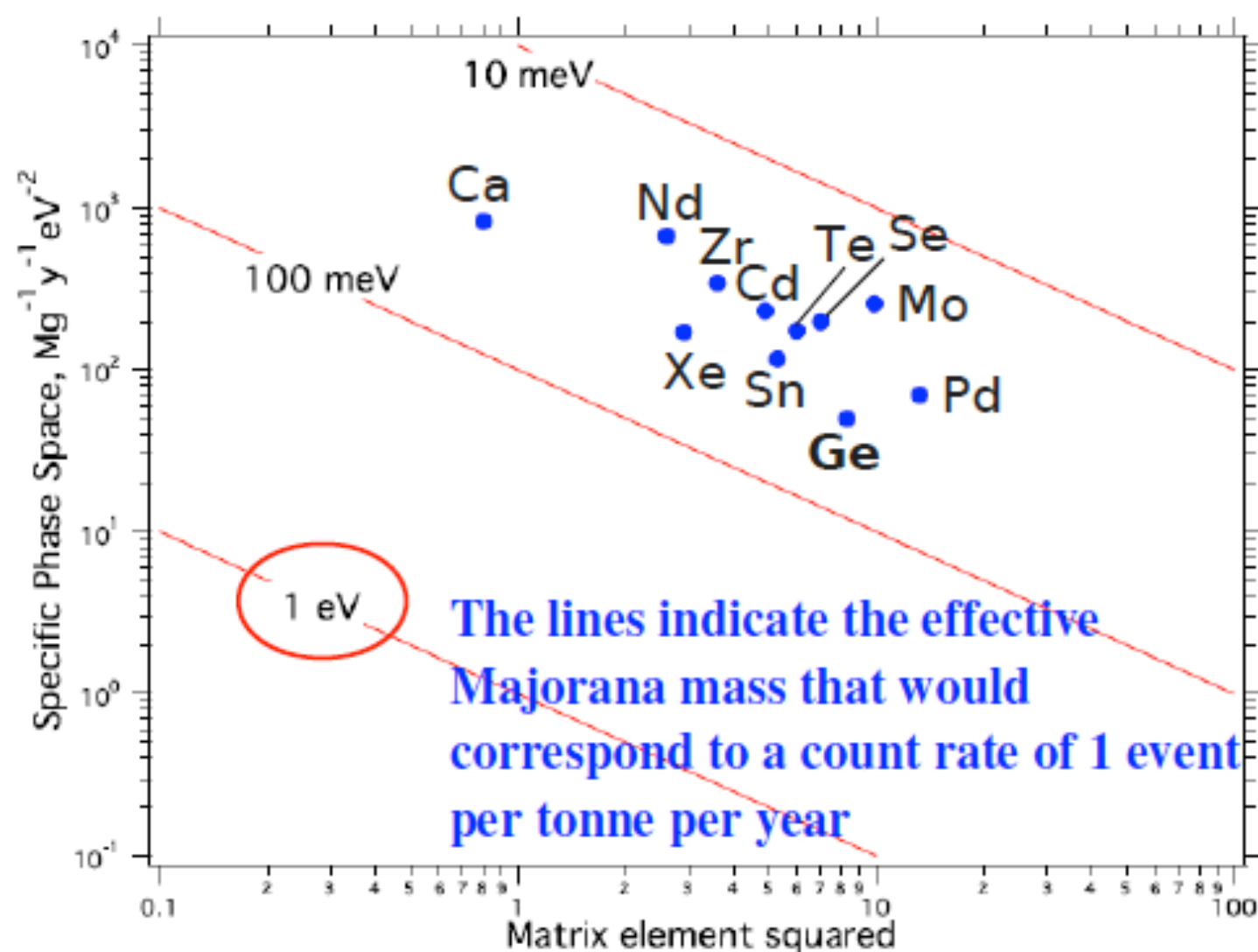
phase space factor

nuclear matrix element

$\langle m_{ee} \rangle = \left| \sum_i U_{ei}^2 m_i \right|$
effective Majorana
neutrino mass

Comparing different isotopes

R.G.H. Robertson arXiv:1301.1323



$$\left(T_{1/2}^{0\nu}\right)^{-1} = G^{0\nu}(Q_{\beta\beta}, Z) \left|M^{0\nu}\right|^2 \left(\frac{\langle m_{ee} \rangle}{m_e}\right)^2$$

\uparrow nuclear matrix element
 \uparrow phase space factor

$\langle m_{ee} \rangle = \left| \sum_i U_{ei}^2 m_i \right|$
 effective Majorana neutrino mass

No theoretical preference

- Phase Space and NME inversely correlated. Tend to compensate.
- Theoretical uncertainties very large

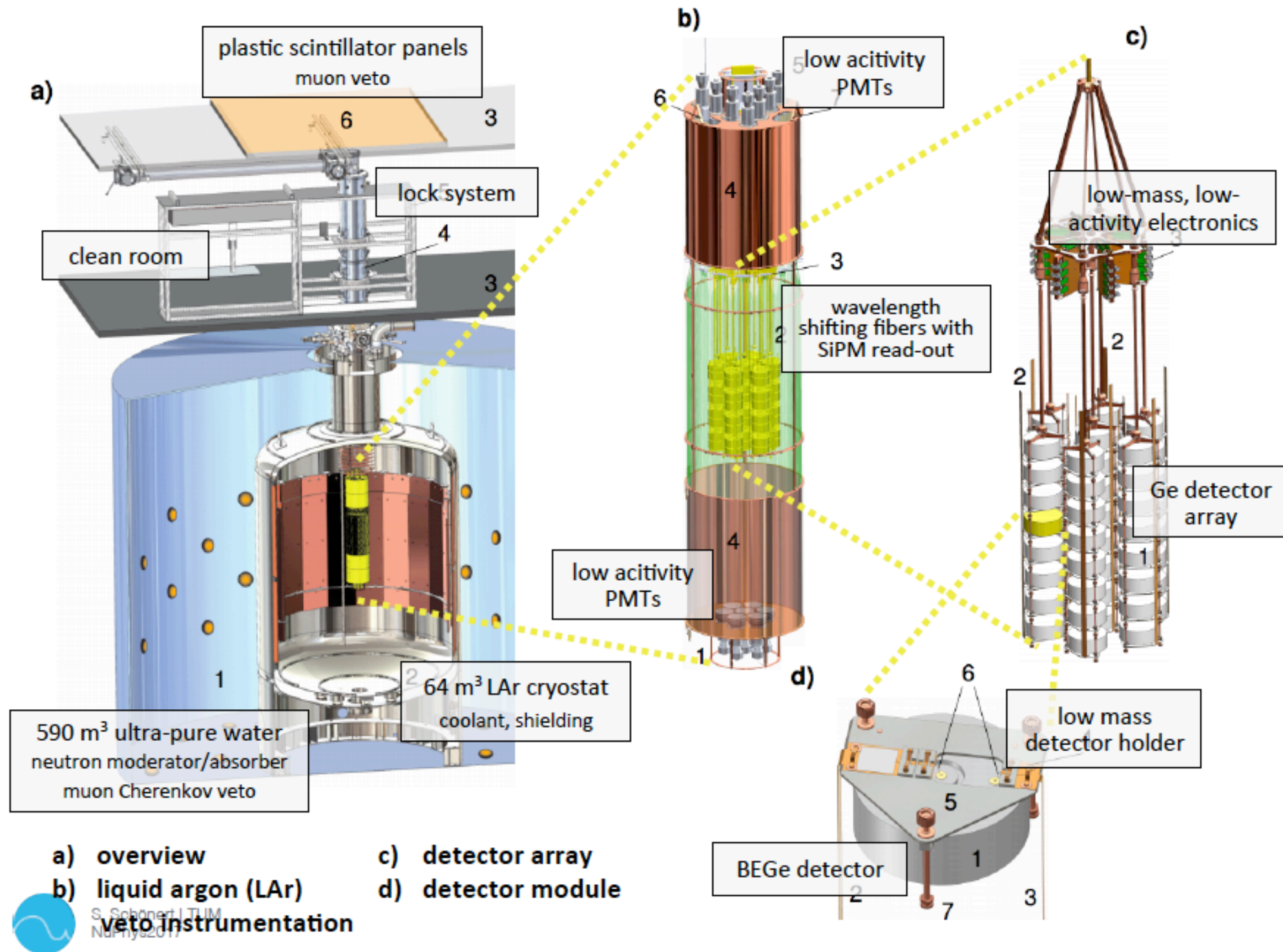
Experimental/practical criteria

- Enrichment cost
- Energy resolution
 - ◆ Narrow peak for discovery
- Background index
 - ◆ Ultraclean components
 - ◆ Avoid surfaces
 - ◆ Especially in a vacuum
- Scalability
 - ◆ Liquids, gases, large crystals

$0\nu\beta\beta$ decay

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Gerda/LEGEND



- high-purity germanium (HPGe) detectors enriched in ^{76}Ge to (86–88)%: source + detector
- detectors mounted in low-mass holders (to minimize radioactive bkg)
- **embedded in liquid argon (LAr): cryogenic coolant and absorber against external radiation**
- ultrapure water tank: buffer around cryostat as additional absorber + Cherenkov mu veto

Synoptic comparison (*not most up to date*)

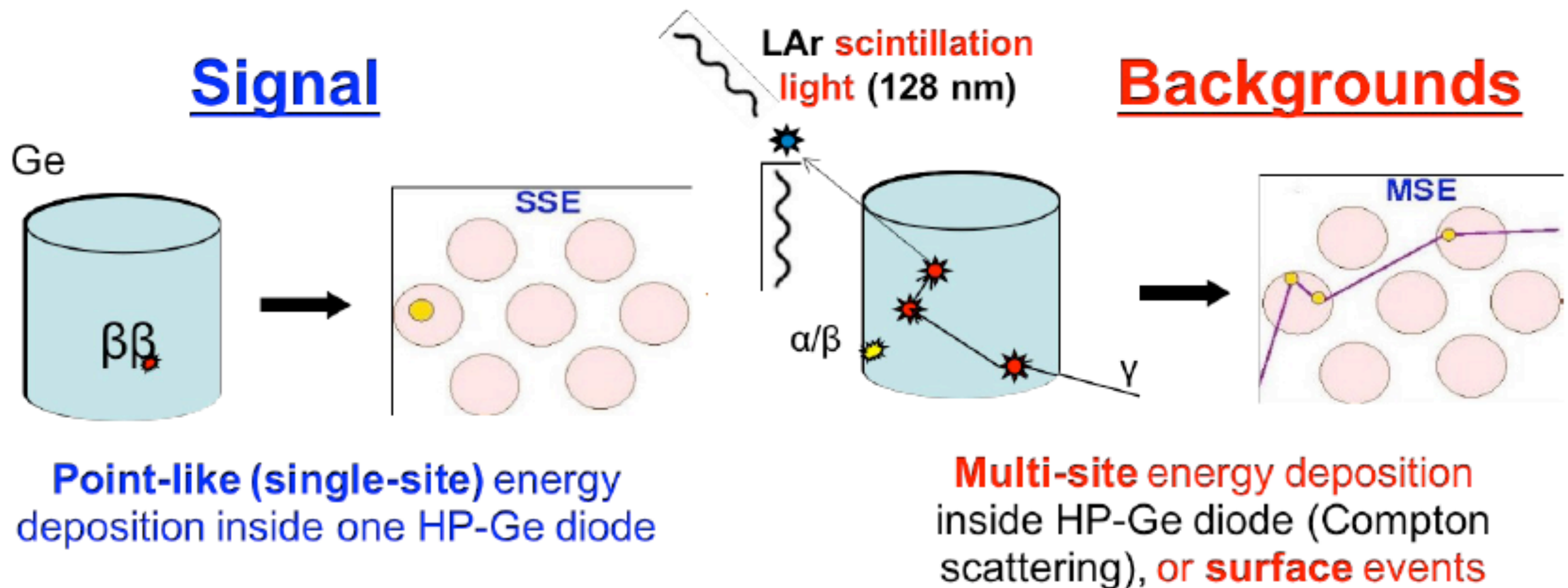
Isotope	Experiment	Exposure (kg yr)	$T_{1/2}^{0\nu\beta\beta}$ average sensitivity (10^{25} yr)	$T_{1/2}^{0\nu\beta\beta}$ (10^{25} yr) 90%CL	$\langle m_\nu \rangle$ (meV) Range from NME*	Reference
^{76}Ge	GERDA	46.7	5.8	>8.0	<120-270	L. Pandola for GERDA Collab, TAUP 2017
	Majorana Demonstrator	10	>2.1	>1.9	<240-520	C.E. Aalseth, arXiv:1710.11608v1
^{130}Te	CUORE	86.3	0.7	>1.5	<140-400	C. Alduino, et al., arXiv:1710.07988v1
^{136}Xe	EXO-200	177.6	3.7	>1.8	<147-398	Albert et al. arXiv: 1707.08707 (2017)
	KamLAND- ZEN	504**	4.9	>11 (run 2)	<60-161	Gando et al., PRL 117 (2016) 082503

Note that the range of “viable” NME is chosen by the experiments and uncertainties related to g_A are not included. ** All Xe. Fiducial Xe is more like ~150 kg yr

To achieve higher sensitivity, the next generation of experiments will be at the ton-scale.

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Active background reduction tools



- Anti-coincidence with the muon veto (MV)
- Anti-coincidence between detectors (cuts multi-site) (AC)
- Active veto using LAr scintillation (LAr Veto)
- Pulse shape discrimination (PSD)

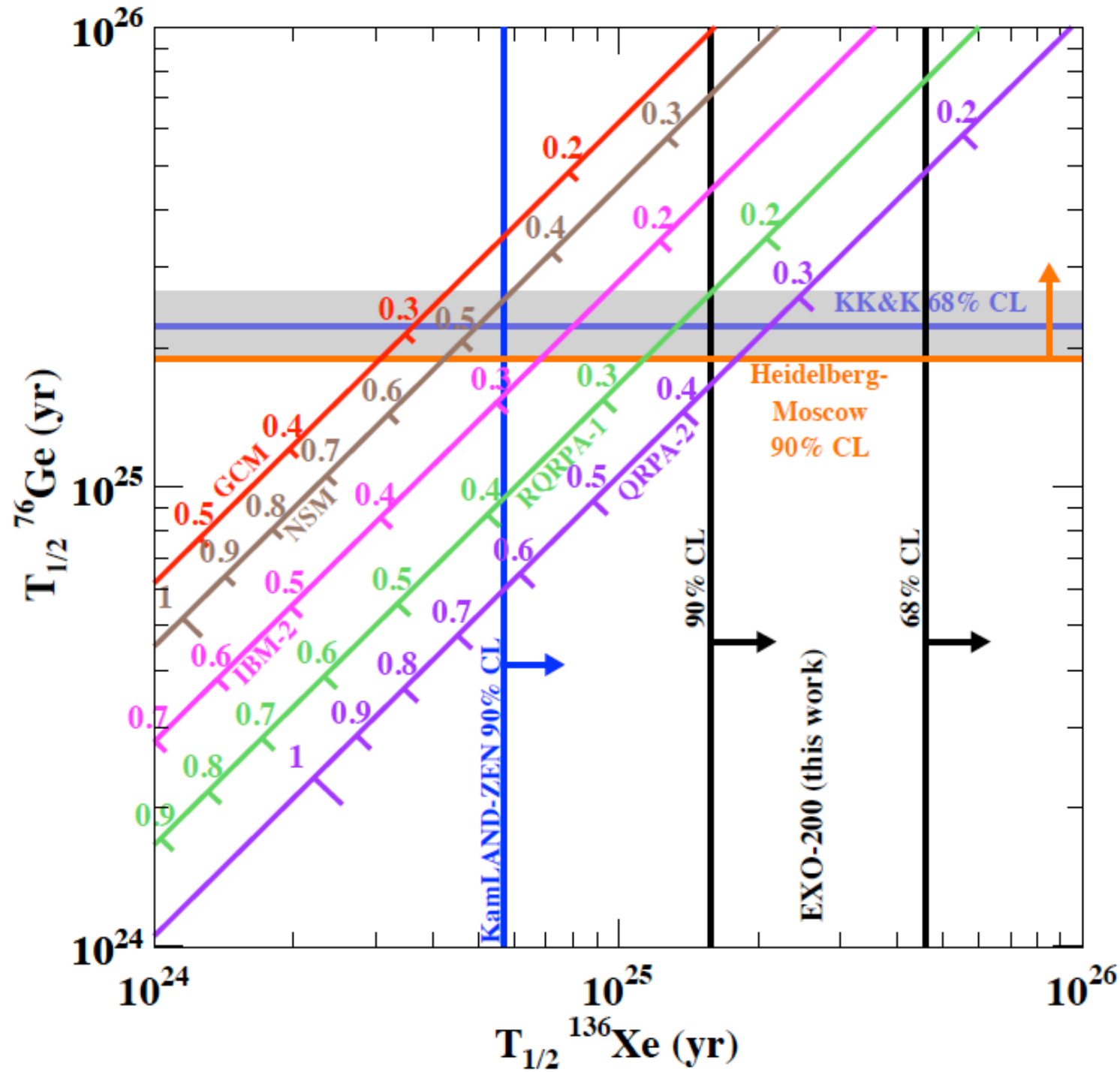
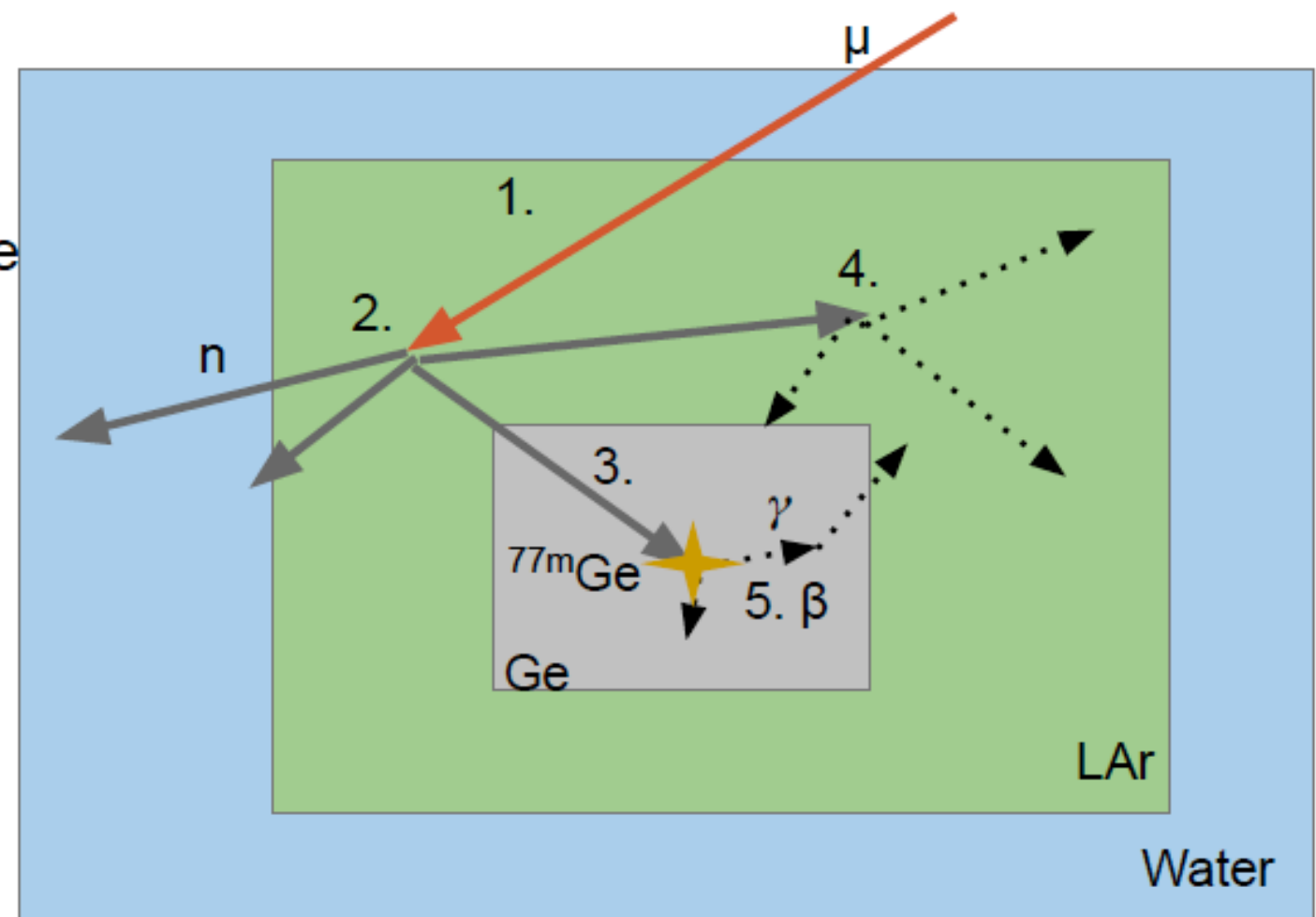


FIG. 6: Relation between the $T_{1/2}^{0\nu\beta\beta}$ in ^{76}Ge and ^{136}Xe for different matrix element calculations (GCM [20], NSM [21], IBM-2 [22], RQRPA-1 [23] and QRPA-2 [5]). For each matrix element $\langle m \rangle_{\beta\beta}$ is also shown (eV). The claim [4] is represented by the grey band, along with the best limit for ^{76}Ge [19]. The result reported here is shown along with that from [7].

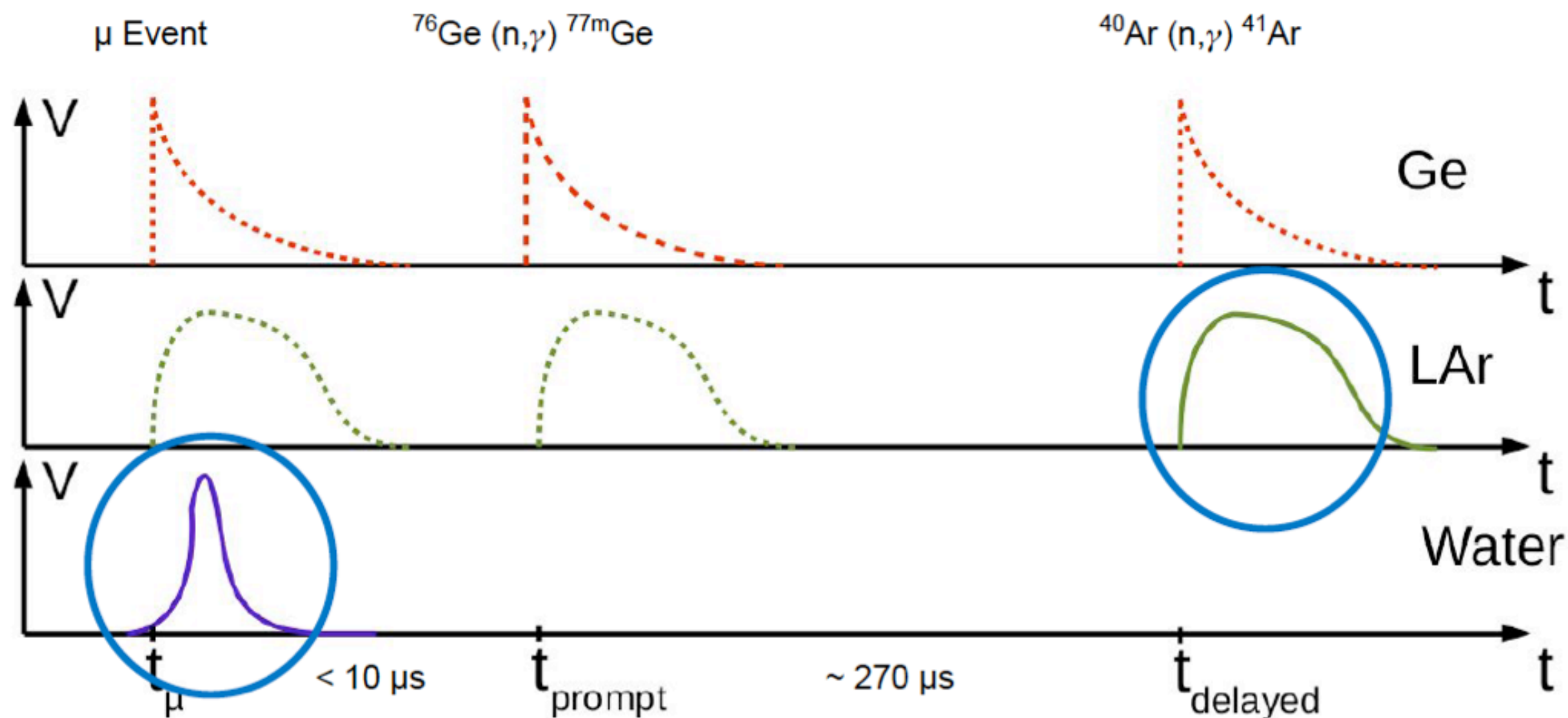
Topology of μ induced events with ^{77m}Ge production

1. A cosmic muon passes through the setup creating signals in water and LAr channels
2. Neutrons are produced; usually $n > 1$
3. Neutron capture on ^{76}Ge producing ^{77m}Ge & prompt gammas
4. Neutrons of step 2. can be captured on ^{40}Ar , producing gammas, $\tau \sim 271 \mu\text{s}$
5. β decay of ^{77m}Ge : $T_{1/2} \sim 53 \text{ s}$



Signal structure of $^{77\text{m}}\text{Ge}$ production events

- Possible to tag $^{77\text{m}}\text{Ge}$ production by muon veto and delayed LAr signals only
- Delayed signals can reduce dead time compared to prompt signals

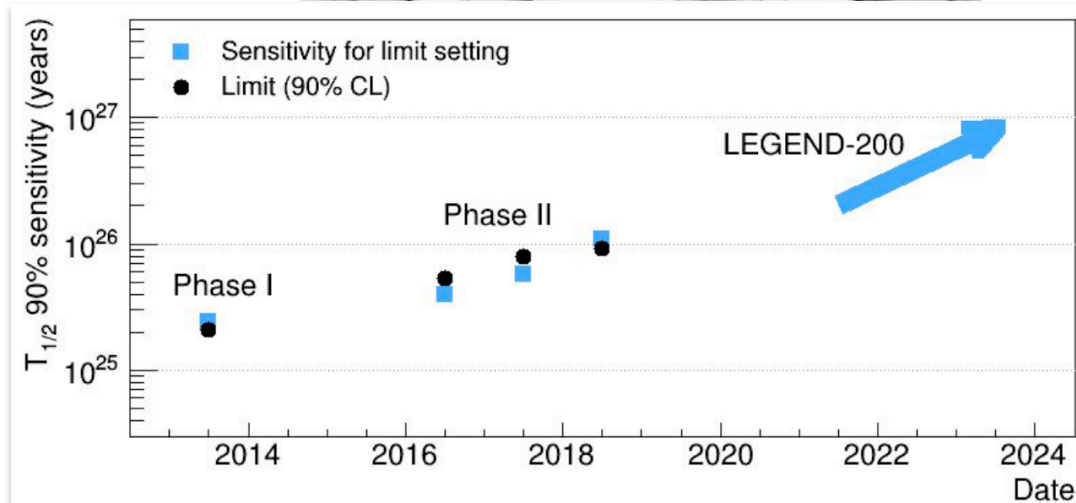


LEGEND program



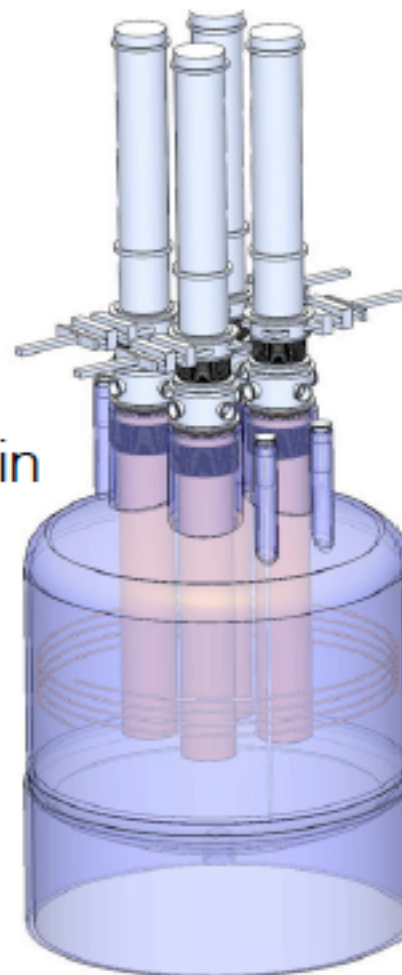
LEGEND-200 (first phase):

- up to 200 kg of detectors
- BI ~ 0.6 cts/(FWHM \cdot t \cdot yr) $\rightarrow \sim 1/5$ of Gerda!
- use existing GERDA infrastructure at LNGS
- design exposure: 1 t \cdot yr
- Sensitivity 10^{27} yr

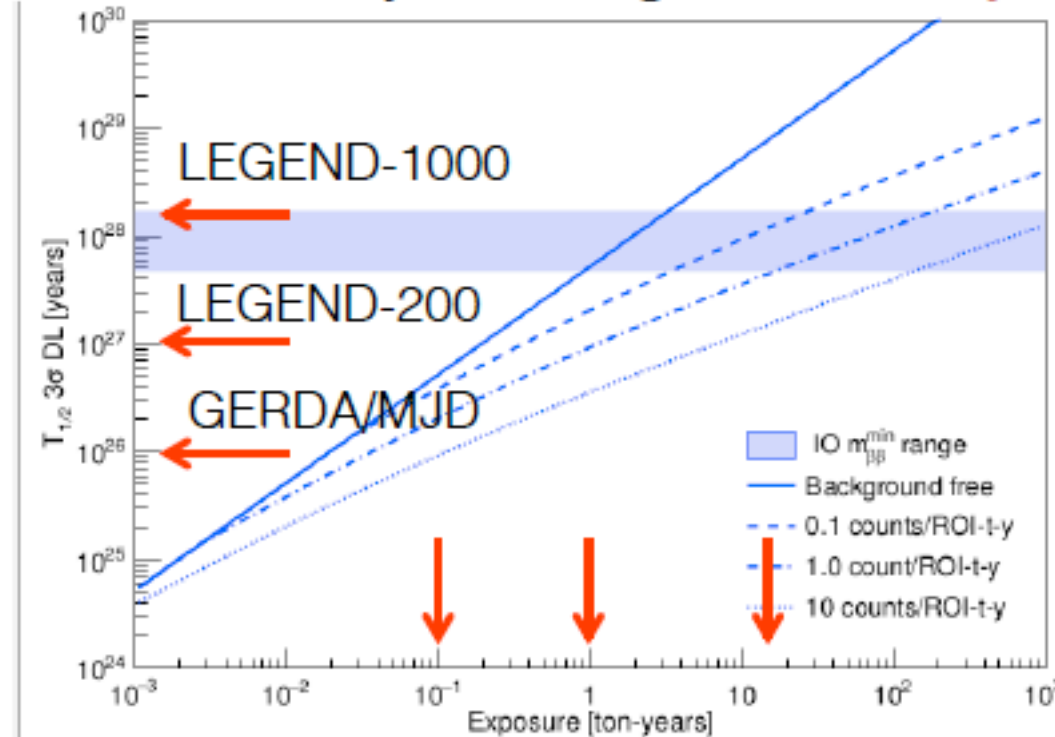


LEGEND-1000 (second phase):

- 1000 kg of detectors (deployed in stages)
- BI < 0.1 cts/(FWHM \cdot t \cdot yr)
- Location tbd
- Design exposure 12 t \cdot yr
- $1.2 \cdot 10^{28}$ yr



Sensitivity for 3σ signal discovery

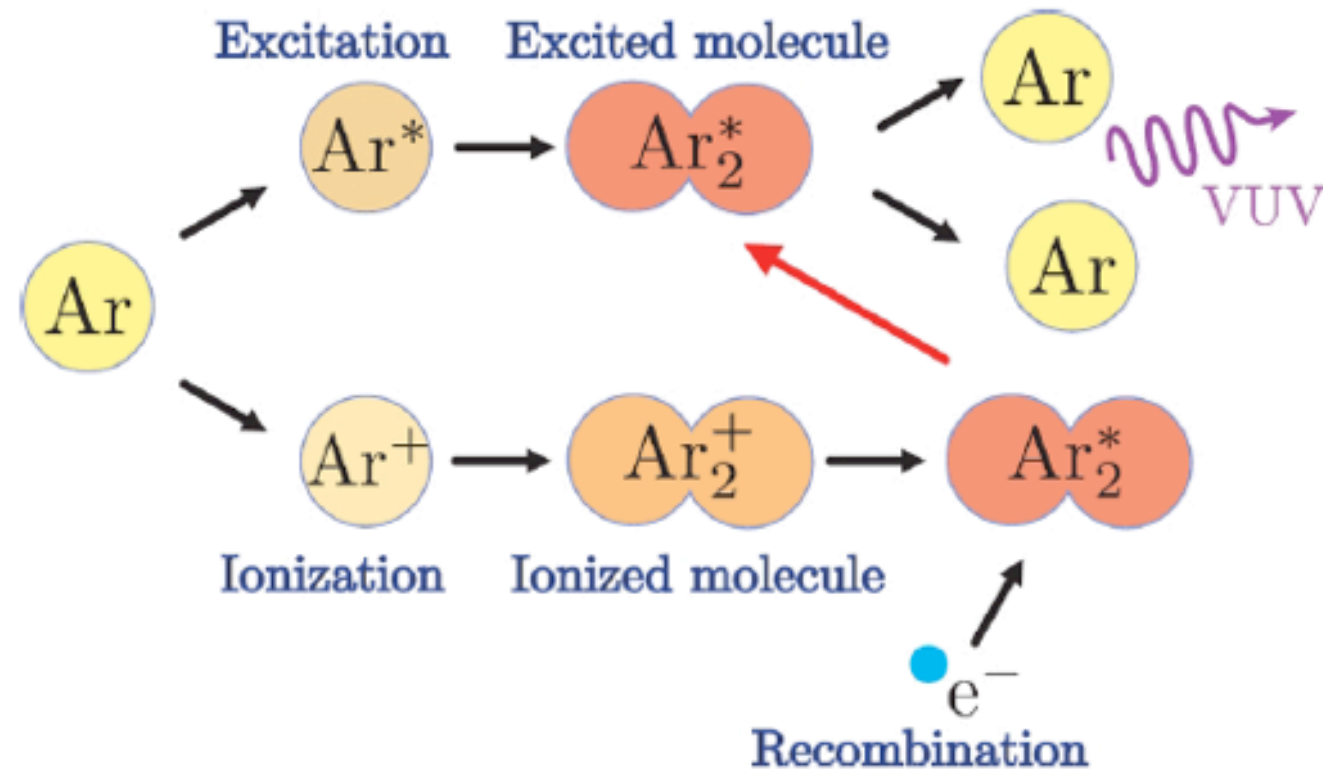


17 meV discovery sensitivity
for "worst case" NME of 3.5



S. Schönert | TUM
NuPhys2017

Argon scintillation mechanism



[arXiv:0708.2621v1]

- Ionizing radiation leads to excited or ionized argon atoms
→ Forming molecules with ground state argon atoms
- Ar_2^+ recombine with free electrons into excited states
- Excited states are created in:
 - singlets (allowed, $\tau \sim 4 - 7$ ns)
 - triplets (forbidden, $\tau \sim 1.0 - 1.7$ μ s)
 - singlet to triplet production ratio is 0.3 for electrons
- decay by emission of photons, $\lambda = 128$ nm
- Contaminations lead to reduction of triplet lifetime and extinction of scintillation light

Consequences of impurities in LAr

O₂ : Free electrons are bound by Oxygen molecules $e^- + O_2 \rightarrow O_2^-$

Non-radiative collision $Ar_2^* + O_2 \rightarrow 2Ar + O_2$

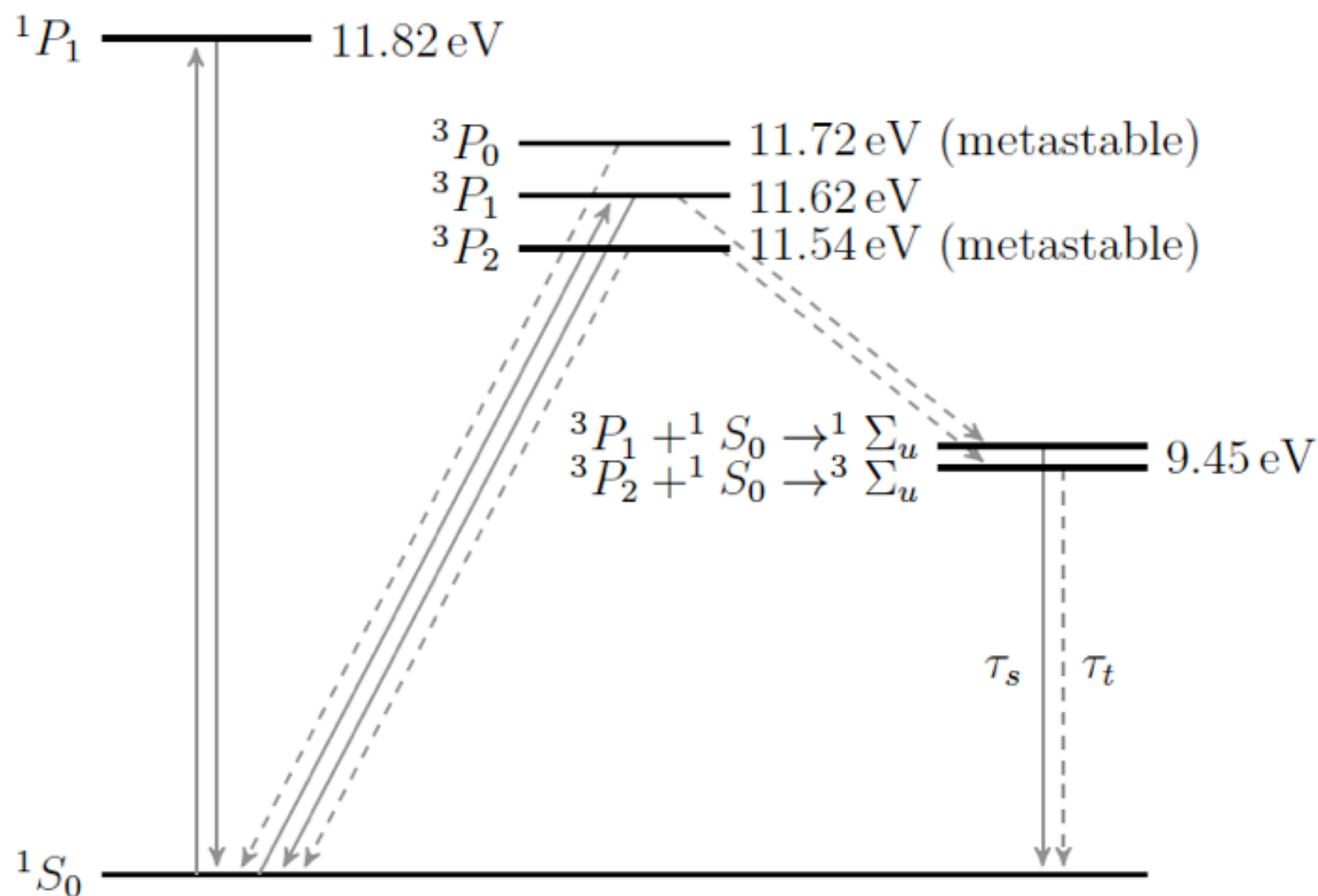
Absorption of scintillation photons, emission of Oxygen resonance line at 557 nm

N₂ : Non-radiative collision $Ar_2^* + N_2 \rightarrow 2Ar + N_2$

Xe : Emission of the Xenon resonance line around 149 nm,
emission of Xenon excimer at 175 nm

α_{att}	Impurity	concentration	Reference
50 cm	O ₂ , H ₂ O, CH ₄ Xe, Kr Hg N ₂	100 ppb 1 ppb 10 ppb 1 ppm	arXiv:1611.02481
110 cm	Xe	0.1 ppm	arXiv:1511.07725
66 ± 3 cm 118 ± 10 cm	Xe	3%	interpreted as scattering NIMPRS A: V. 384, 380-386
1790 ± 160 (m?) 30 ± 3 m	N ₂ N ₂	37 ppb 2 ppb	arXiv:1306.4605

Argon scintillation mechanism



- Electron configuration of argon: $[\text{Ne}](3s)^2(3p)^6 \rightarrow [\text{Ne}](3s)^2(3p)^5(4s)$
- Radiation from 1P_1 and 3P_1 to 1S_0 ground state reabsorbed in argon
- Atom in 3P_1 and 3P_2 state can form molecule with 1S_0 argon atom
- γ 's from Ar_2^* have not enough energy to be absorbed by 1S_0 argon atoms
- Peaks from singlet and triplet excimer decay are not resolved