

Designing Matrix Models for Zeta Functions

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This talk is based on a collaboration with
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[arXiv:1807.07342 \[math-ph\]](#)

Matrix model for Riemann zeta via its local factors
and some ongoing work

Disclaimer: We are **not** trying to prove the Riemann hypothesis,
lest you think we've lost it!

Outline

Introduction & Motivation

Phase Space Description of the Unitary Matrix Model

Unitary Matrix Model for the Symmetric Zeta-function

UMM for the Local ζ -function and Attempts at a Synthesis

Zeta function: infinite sum and product

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 - ▶ extended to real $s > 1$ by Chebyshev.
 - ▶ analytically continued by Riemann to the complex s -plane as a meromorphic function.
- Useful in regularising infinities in Physics.

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The **symmetric zeta-function** $\xi(s) = \frac{1}{2} \pi^{-s/2} s(s-1) \Gamma\left(\frac{s}{2}\right) \zeta(s)$ is an **entire function** that satisfies $\xi(s) = \xi(1-s)$. Its zeroes are at the non-trivial zeroes of $\zeta(s)$, at $s = \gamma_m = \frac{1}{2} + it_m$.

Some related functions

n	γ_n
1	$\frac{1}{2} \pm i 14.1347 \dots$
2	$\frac{1}{2} \pm i 21.0220 \dots$
3	$\frac{1}{2} \pm i 25.0108 \dots$
4	$\frac{1}{2} \pm i 30.4248 \dots$
5	$\frac{1}{2} \pm i 32.9350 \dots$
6	$\frac{1}{2} \pm i 37.5861 \dots$
7	$\frac{1}{2} \pm i 40.9187 \dots$
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First few zeroes

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$$\mathcal{J}(x) = Li(x) - \sum_n Li(x^{\gamma_n})$$

Zeroes as the spectrum

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Dyson pointed out that this has the same behaviour as the two-point correlator of the eigenvalues of an ensemble of random hermitian matrices.

Numerical evidence and extension

Odlyzko confirmed this behaviour from his numerical computation of Riemann zeroes.

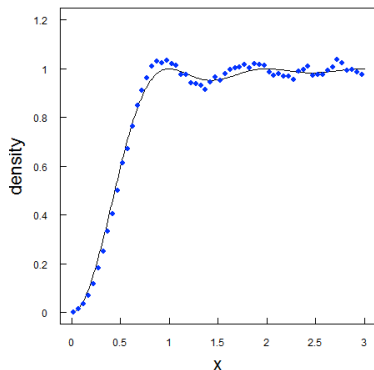


Figure : Blue dots describe the normalized spacings of the first 10^5 non-trivial zeros of the Riemann zeta function. The solid line describes the two-point correlation function of GUE of random matrices. (source: Wikipedia)

Numerical evidence and extension

Odlyzko confirmed this behaviour from his numerical computation of Riemann zeroes.

Rudnick-Sarnak extended it to higher correlators.

Özlük extended to the zeroes of Dirichlet L -functions:

$$L(s, \chi) = \sum_{n=1}^{\infty} \frac{\chi(n)}{n^s} = \prod_{p \in \text{primes}} \frac{\chi(p)}{(1 - p^{-s})}$$

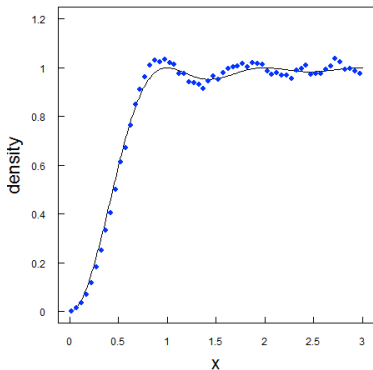


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Search for the Hamiltonian

Berry-Keating (and Connes) proposed the quantised form of the classical xp Hamiltonian : $H_{BK} = (xp + px) = -2i\hbar \left(x \frac{d}{dx} + \frac{1}{2} \right)$.

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They were motivated by the similarity of the fluctuating part of the **prime distribution function** and the **Gutzwiller trace formula** relating the fluctuating part of the energy eigenvalues and the periods of a **chaotic dynamical system**.

Riemann zeroes



Primes

~

Energy eigenvalues



Periods

(of primitive periodic orbits)

Intriguing similarities

The fluctuating part of the **distribution function**

$J(x) = \sum_m \Theta(x - t_m)$ has the form

$$J_{\text{fl}}(x) = -\frac{1}{\pi} \sum_p \sum_{n \in \mathbb{N}} \frac{1}{n} e^{-n \ln p / 2} \sin(xn \ln p)$$

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This is to be compared with the fluctuating part of the energy eigenvalues in the **Gutzwiller trace formula** relating to the periods of a **chaotic dynamical system**:

$$\varrho_{\text{fl}}(E) = \frac{1}{\pi} \sum_p \sum_{n \in \mathbb{N}} \frac{1}{n} e^{-n \lambda_p \tau_p / 2} \sin \left(n S_p(E) - \frac{\pi}{2} n \mu_p \right)$$

Some issues

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- Suggestion: Restrict the values of x and p .

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New ingredients: p -adic analysis, Hilbert space over p -adic numbers, in particular, wavelets as a basis, and operators on this space.

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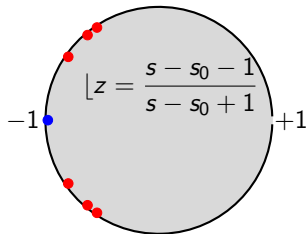
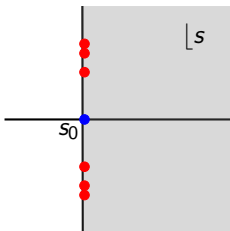
UMM for the Local ζ -function and Attempts at a Synthesis

Conformal map

The eigenvalues of large $N \times N$ unitary matrices gives a density $\rho(\theta) = \sum \delta(\theta - \theta_i)$ (distribution function) on the unit circle.

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The eigenvalues of large $N \times N$ unitary matrices gives a density $\rho(\theta) = \sum \delta(\theta - \theta_i)$ (distribution function) on the unit circle. Given a distribution on the line $\text{Re } s = s_0$, one can find a Gaussian Unitary Ensemble (GUE) such that its eigenvalue distribution is related to it.



$$s - s_0 = \frac{1 + z}{1 - z} = \frac{1 + e^{i\theta}}{1 - e^{i\theta}} = i \cot \frac{\theta}{2}$$

One-plaquette UMM

The partition function of the **one-plaquette model** is defined by:

$$\mathcal{Z} = \int \mathcal{D}U \exp \left[N \sum_{n=0}^{\infty} \frac{\beta_n}{n} \left(\text{Tr } U^n + \text{Tr } U^{\dagger n} \right) \right] = \int \prod_{i=1}^N \frac{d\theta_i}{2\pi} e^{-N^2 S_{\text{eff}}(\theta_i)}$$

$$\text{where, } S_{\text{eff}}(\theta_i) = - \sum_{n=1}^{\infty} \sum_{i=1}^N \frac{2\beta_n}{n} \cos(n\theta_i) - \frac{1}{2} \sum_{i \neq j} \ln \left(4 \sin^2 \frac{\theta_i - \theta_j}{2} \right)$$

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In the large N limit, $x = \frac{i}{N} \in [0, 1]$ and $\theta_i \rightarrow \theta(x)$

$$S[\theta] = - \sum_{n=1}^{\infty} \int_0^1 dx \frac{2\beta_n}{n} \cos n\theta(x) - \frac{1}{2} \int_0^1 dx \oint_0^1 dy \ln \left(4 \sin^2 \frac{\theta(x) - \theta(y)}{2} \right)$$

Saddle point

The **saddle point** of the action is determined by

$$\oint \frac{d\theta'}{2\pi} \rho(\theta') \cos\left(\frac{\theta - \theta'}{2}\right) = \sum_{n=1}^{\infty} 2\beta_n \sin n\theta \quad \left(= \frac{dV(\theta)}{d\theta}\right)$$

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Taylor expansion of the resolvent.

$$2\pi\rho(\theta) = 2\text{Re} [R(e^{i\theta})] - 1 = \lim_{\epsilon \rightarrow 0} \left[R((1 + \epsilon)e^{i\theta}) - R((1 - \epsilon)e^{i\theta}) \right]$$

UMM in terms of Irreps (Schematic)

The PF of a UMM can also be expanded in terms of the **irreducible representations (irreps)** of $U(N)$

$$\mathcal{Z} \sim \sum_{R \in \text{irreps}} \sum_{\vec{k}, \vec{\ell}} \alpha(\vec{\beta}, \vec{k}) \alpha(\vec{\beta}, \vec{\ell}) \chi_R(C(\vec{k})) \chi_R(C(\vec{\ell}))$$

(where $\chi_R(C(\vec{k}))$ is the **character** of the **conjugacy class** $C(\vec{k})$ of the **permutation group** $S_{K=\sum n k_n}$).

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(where $\chi_R(C(\vec{k}))$ is the **character** of the **conjugacy class** $C(\vec{k})$ of the **permutation group** $S_{K=\sum n k_n}$). The following have been used

$$\prod_n (\text{Tr } U^n)^{k_n} = \sum_R \chi_R(C(\vec{k})) \text{Tr}_R(U)$$

$$\int \mathcal{D}U \text{Tr}_R(U) \text{Tr}_{R'}(U^\dagger) = \delta_{RR'}$$

Young diagrams and momenta

Irreps can be labelled by the **number of boxes** in **Young diagrams**.
In the **large N** limit

$$\mathcal{Z} = \int \mathcal{D}h(x) \int d\vec{k} d\vec{\ell} \exp \left(-N^2 S_{\text{eff}}[h(x), \vec{k}, \vec{\ell}] \right)$$

where $u(h)dh \sim dx$ is another density function.

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The variables h are the **momenta conjugate** to the eigenvalues θ .

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There is a density $\Omega(\theta, h)$ in the phase space, such that

$$\int dh \Omega(\theta, h) = \rho(\theta) \quad \text{and} \quad \int d\theta \Omega(\theta, h) = u(h)$$

Expectation: Phase space description may lead to a Hamiltonian.

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- ▶ Write the partition function of the UMM in phase space:

$$Z \sim \int d\theta d\mathfrak{h} e^{-H(\theta, \mathfrak{h})}.$$

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- ▶ Construct a **unitary matrix model (UMM)** for which these zeroes are the eigenvalues: $Z \sim \int \mathcal{D}U e^{S(U)}$.
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$$Z \sim \int d\theta d\mathfrak{h} e^{-H(\theta, \mathfrak{h})}.$$
- ▶ Try to realise this as the **trace** of some **operator**: $Z \sim \text{Tr } \hat{\rho}$.

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- ▶ This determines the parameters of the one plaquette model:

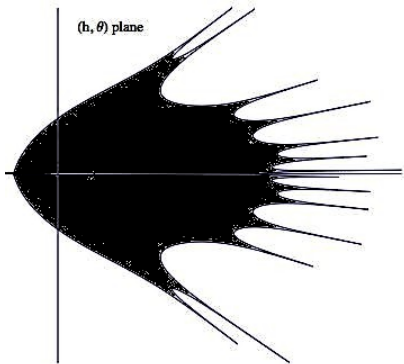
$$\beta_n = -\frac{1}{2n \ln 2} \lambda_n = \frac{1}{2 \ln 2} \oint_{C_1} \frac{ds}{2\pi i} \frac{s^{n-1}}{(s-1)^n + 1} \ln \xi(s)$$

in terms of the Keiper-Li numbers[‡]

$$\lambda_n = \frac{1}{(n-1)!} \frac{d^n}{ds^n} s^{n-1} \ln \xi(s) \Big|_{s=1} = \sum_i \left[1 - \left(1 - \frac{1}{\gamma_i} \right)^n \right]$$

Phase space density of the UMM of $\xi(s)$

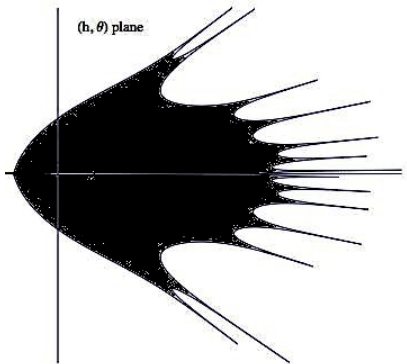
The density in the phase space is



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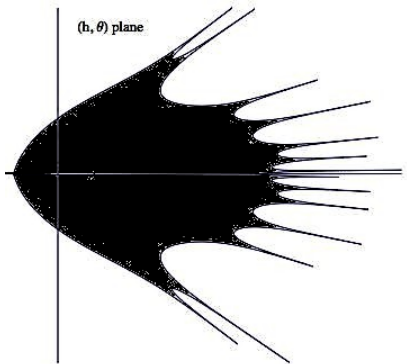
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The prime power counting function $J(x)$ jumps by $1/n$ at every p^n :

$$\begin{aligned} J(x) &= \sum_{p,n} \frac{1}{n} \Theta(x - p^n) \\ &= \langle J \rangle(x) + J_{\text{fl}}(x) \end{aligned}$$

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Turns out that the momentum density $h(x) \sim J_{\text{fl}}(x)$, the **fluctuating part** of the counting function.

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Singling out a prime: Leaving rationality

A p -adic number $\xi = p^N (\xi_0 + \xi_1 p + \xi_2 p^2 + \cdots)$, where N is an integer, $\xi_k = \{0, 1, \cdots p-1\}$ but $\xi_0 \neq 0$, and

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\mathbb{Q}_p 's are close relatives of the real numbers, although the notion of continuity and 'nearness' are very different, determined by divisibility wrt p . For example, \mathbb{Z}_p , the completion of integers \mathbb{Z} is compact in \mathbb{Q}_p .

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These poles can be brought on the unit circle on $z = \frac{s-1}{s+1}$ plane.

$$R_{<}(z) = 1 + \frac{z}{(1-z)^2} \frac{p^{-s(z)}}{1 - p^{-s(z)}}, \quad R_{>}(z) = -\frac{z}{(1-z)^2} \frac{p^{s(z)}}{1 - p^{s(z)}}$$

The resolvent above satisfies all the properties ($R_{<}(0) = 1$, $R_{>}(z \rightarrow \infty) = 0$ and $R_{<}(z) + R_{>}(1/z) = 1$).

★(Caveat)

A well-known measure

The following is easily computed:

$$\int_{p\mathbb{Z}_p} |h|_p^{s-1} dh = \frac{(1 - p^{-1})p^{-s}}{(1 - p^{-s})}, \quad p\mathbb{Z}_p = \left\{ h \in \mathbb{Q}_p : |h|_p < 1 \right\}$$

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$$\text{So } 2R_{<}(z) - 1 = p \int_{p\mathbb{Z}_p} dh \left(1 + \frac{2z}{(p-1)(1-z)^2} |h|_p^{\frac{1+z}{1-z}-1} \right)$$

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This is suggestive of a phase space density

$$\Omega_p(\theta, h) = p - \frac{p}{2(p-1)\sin^2\left(\frac{\theta}{2}\right)} |h|_p^{-i\cot\left(\frac{\theta}{2}\right)-1} \sim p - \frac{p^{-in\cot\left(\frac{\theta}{2}\right)}}{2(p-1)\sin^2\left(\frac{\theta}{2}\right)}$$

Vladimirov derivative and Kozyrev wavelets

The totally disconnected topology of \mathbb{Q}_p , makes differentiation difficult. [Vladimirov] defined derivative as an integral kernel:

$$\left(D_{(p)}^{\alpha} f\right)(x) = \frac{1}{\Gamma_p(-\alpha)} \int_{\mathbb{Q}_p} dx \frac{f(x) - f(y)}{|x - y|_p^{\alpha}}, \quad \alpha \in \mathbb{C}$$

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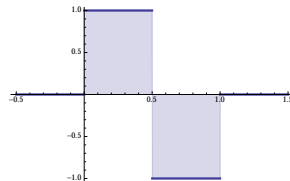
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Scaling and translation generate the affine group $t \rightarrow a t + b$, for $a > 0$ and $b \in \mathbb{R}$.

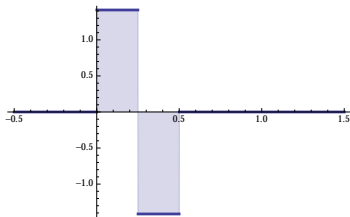
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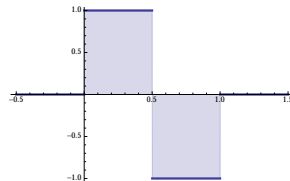


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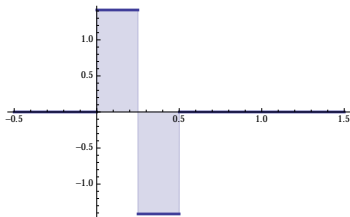


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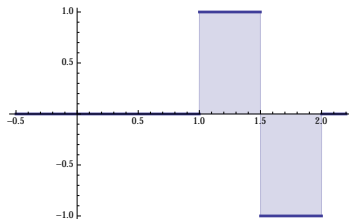
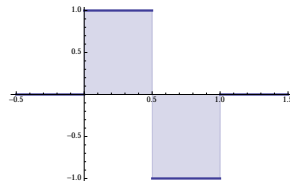


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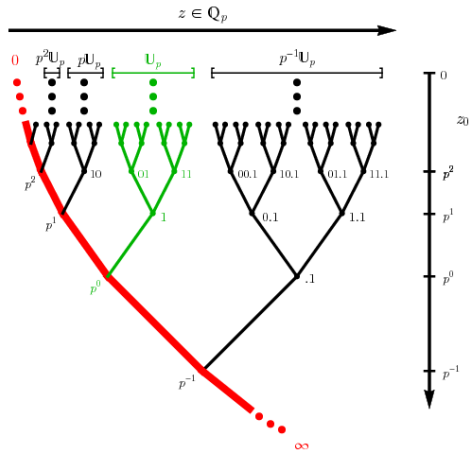
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(Fig. from Gubser et al)



Wavelets on \mathbb{Q}_p

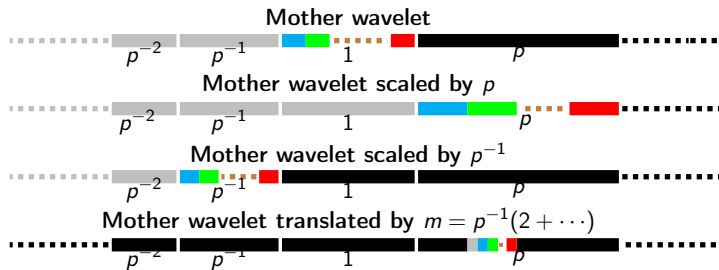


Figure : A schematic representation of the wavelets. The sets are ordered by the values $|\xi|_p = p^n$. (Colour code: grey = 1, black = 0, other colours correspond to primitive roots of unity.) [Dutta,-DG-Lala (2018)]

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The Hilbert space $\mathcal{H}_{(p)}$ of the quantum Hamiltonian is spanned by a subset of the Kozyrev wavelets (which are eigenfunctions of the Vladimirov derivative.)

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$$\sum_p \ln p \frac{dJ_p(\xi)}{d\xi} = \frac{d\psi(\xi)}{d\xi} = 1 - \underbrace{\sum_i \xi^{\gamma_i-1}}_{\text{non-trivial zeroes}} - \sum_n \xi^{2n-1}$$

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Keeping only the non-trivial zeroes γ_i

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Clearly μ has to be independent of i . The reflection symmetry of ζ -function implies that $\mu > 1$ and *assuming* Riemann hypothesis $\mu > \frac{1}{2}$.

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Leads to a **one-parameter family of Hamiltonians**

$$H_\mu \sim H - \mu P$$

Trivial zeroes and local zeta at infinity

Consider $\zeta_{\mathbb{A}}(s) = \pi^{-\frac{s}{2}} \Gamma\left(\frac{s}{2}\right) \prod_p \zeta_p(s) \equiv \zeta_{\infty}(s) \prod_p \zeta_p(s)$, the *adelic zeta function*. It involves \mathbb{Q}_p for all primes and \mathbb{R} , and satisfies $\zeta_{\mathbb{A}}(s) = \zeta_{\mathbb{A}}(1-s)$.

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Wigner functions for local models

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After integrating over **position** (**momentum**) Wigner function gives the **counting function** (resp. **eigenvalue density**) upto an infinite factor.

Wigner function in the large phase space

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Thank you!