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PARTICLE PHYSICS 粒子物

Semiconductor particle detectors

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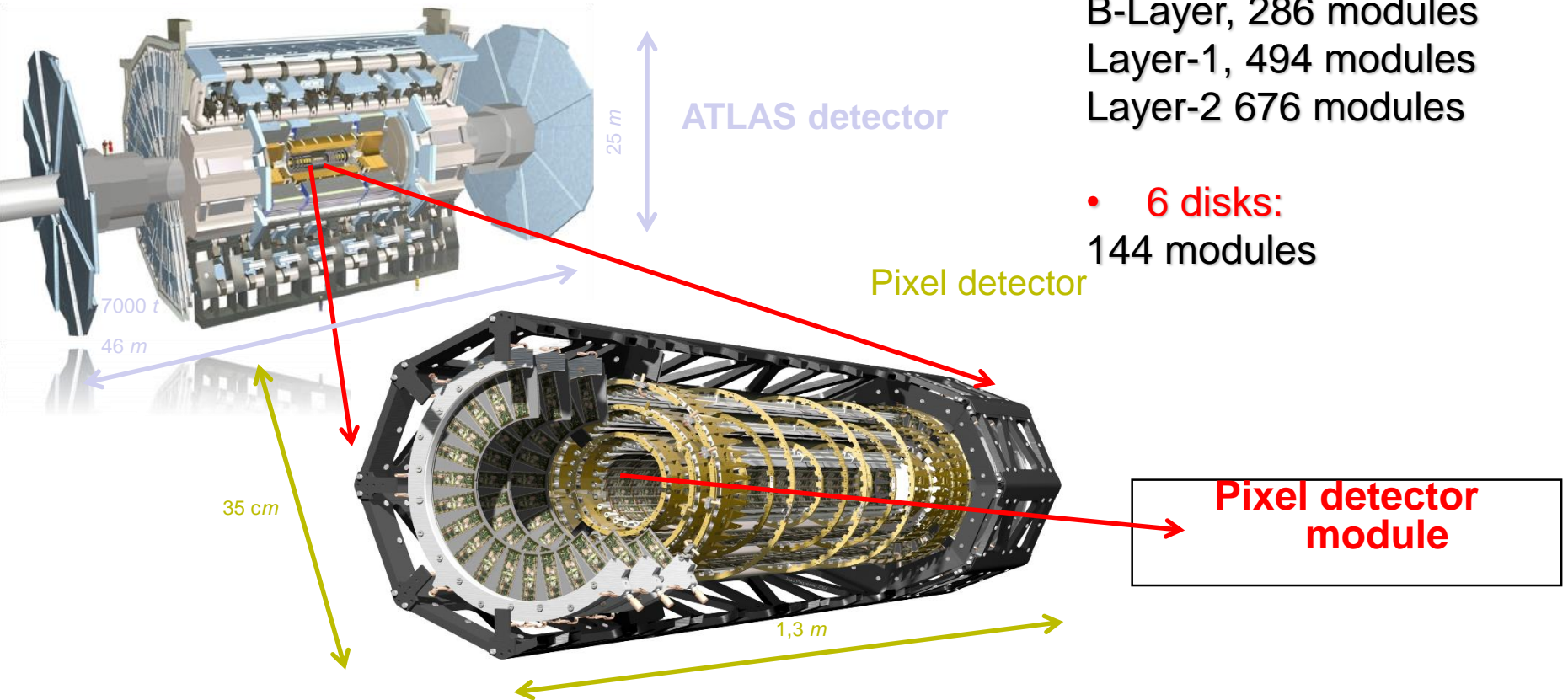


Detectors in high-energy physics (1)

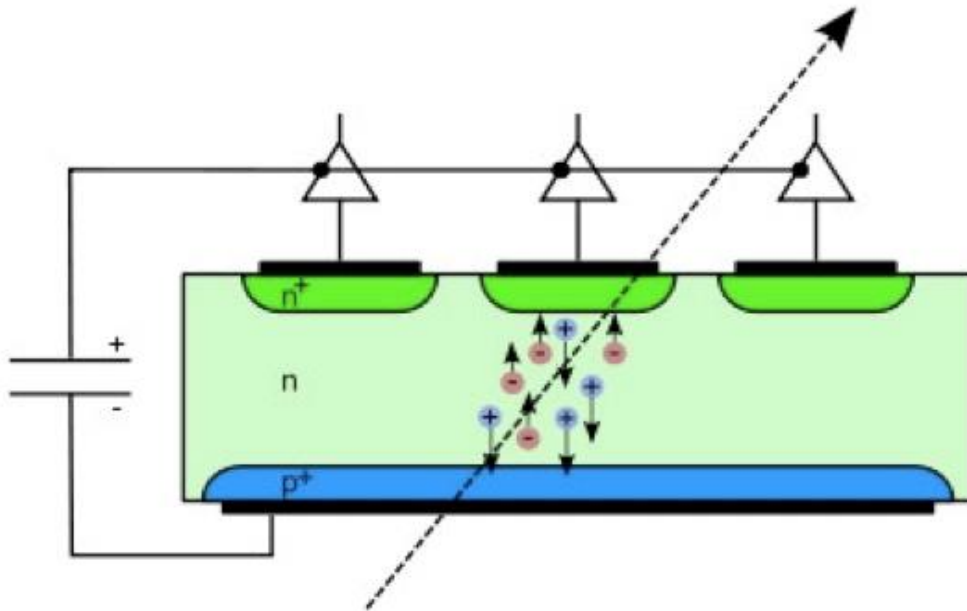
- **particle tracking**: identify that a particle crosses a given surface at a given time
- **energy spectroscopy**: identify the energy of a particle



Detectors in high-energy physic (2)



What a detector does



- the incoming particle generates free carriers inside a semiconductor
- the electric field inside the device separates the charges and creates a current at the terminals
- the electrical signal is amplified by an electronic read-out



Outline

- **Semiconductors**
- p-n junction
- Main figures of merits for detectors



Semiconductors (1)

						VIIIA	
						2	He 4.003
		IIIA	IVA	VA	VIA	VIIA	
		5	6	7	8	9	10
		B 10.811	C 12.011	N 14.007	O 15.999	F 18.998	Ne 20.183
		13	14	15	16	17	18
		Al 26.982	Si 28.086	P 30.974	S 32.064	Cl 35.453	Ar 39.948
IB	IIB						
29	30	31	32	33	34	35	36
Cu 63.54	Zn 65.37	Ga 69.72	Ge 72.59	As 74.922	Se 78.96	Br 79.909	Kr 83.80
47	48	49	50	51	52	53	54
Ag 107.870	Cd 112.40	In 114.82	Sn 118.69	Sb 121.75	Te 127.60	I 126.904	Xe 131.30
79	80	81	82	83	84	85	86
Au 196.967	Hg 200.59	Tl 204.37	Pb 207.19	Bi 208.980	Po (210)	At (210)	Rn (222)

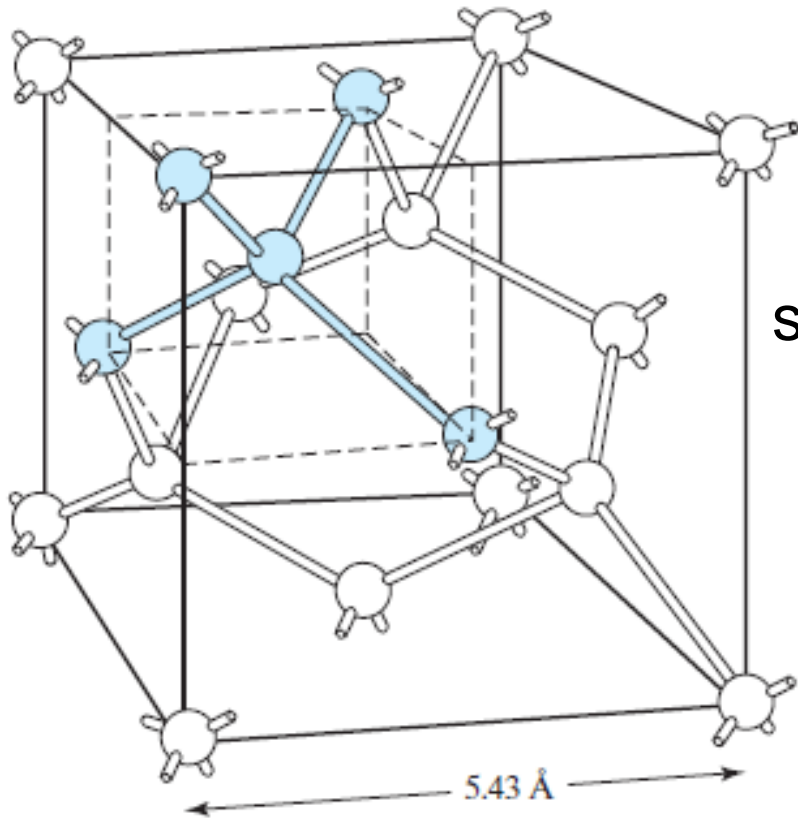
- group IV: Si, Ge
- III-V compounds (GaAs, InP, InAs, GaN...)
- Alloys ($\text{Si}_{1-x}\text{Ge}_x$, $\text{In}_{1-x}\text{Ga}_x\text{As}$)
- CdTeSe, ZnTeSe, CdZnTe

- behave as conductive or insulating materials by adding impurities a/o changing the device bias



Semiconductors (2)

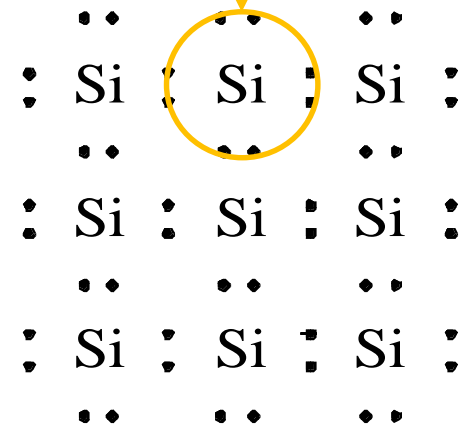
- crystal structure of silicon



simplified 2D view

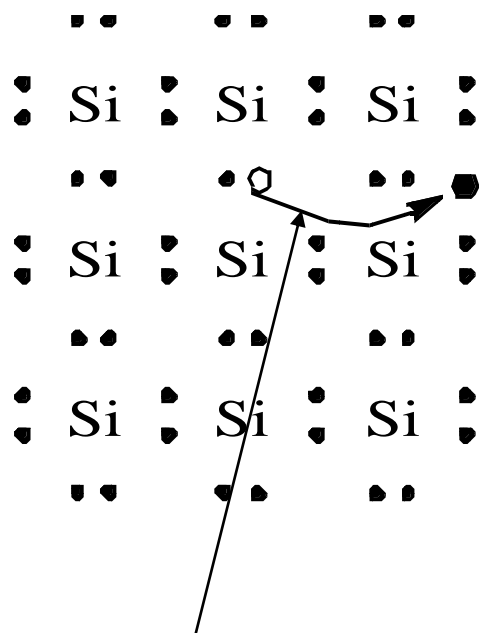


8 electrons in
the outer shell

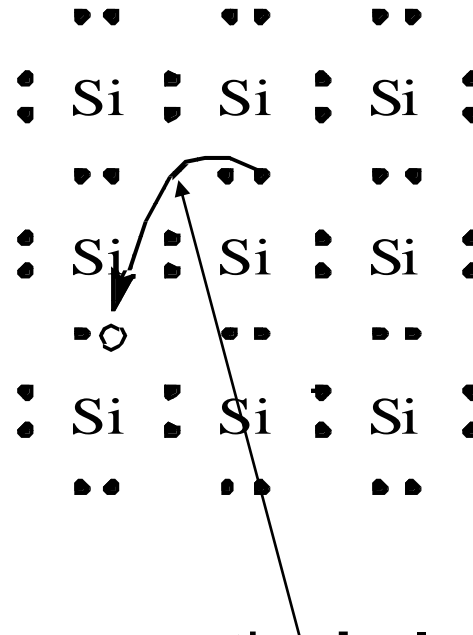




Free carriers: electrons and holes



one **electron** can break the bond and become a **free carrier** (negative charge)

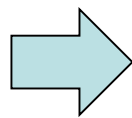
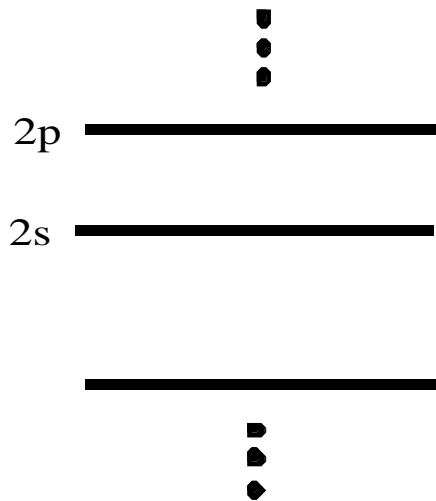


the **hole** that is left can move too: free carrier with positive charge

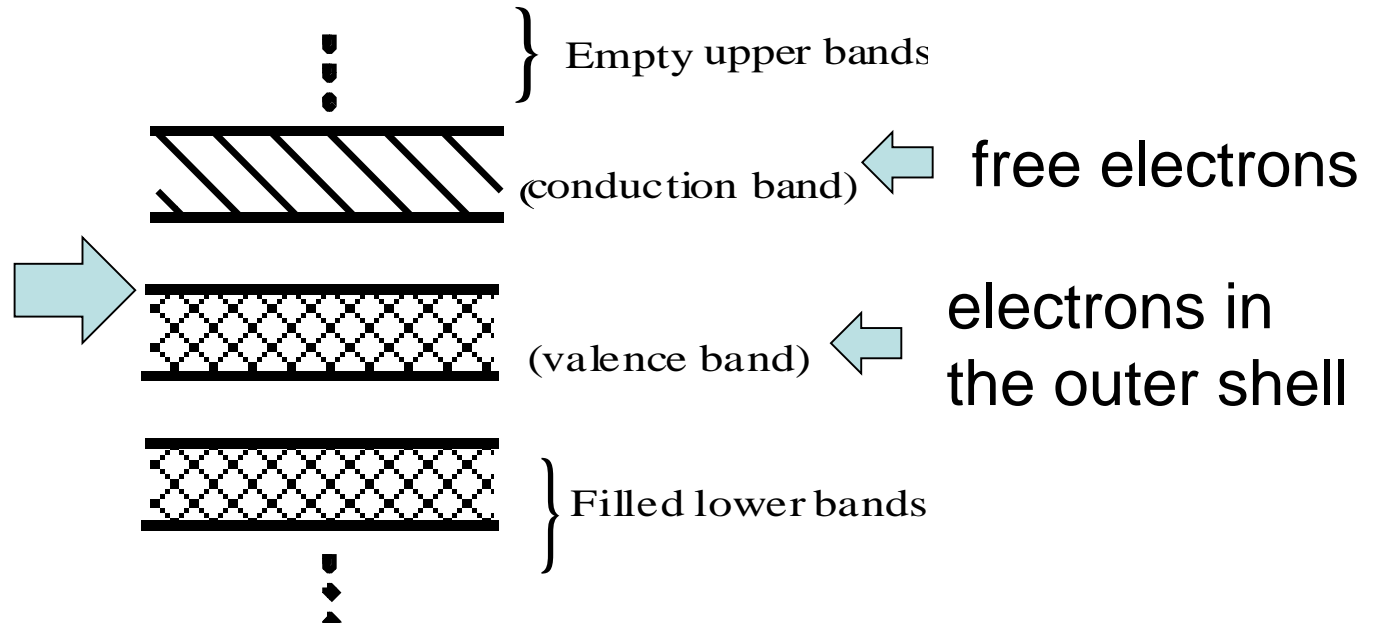


Energy bands

single atom:
energy levels

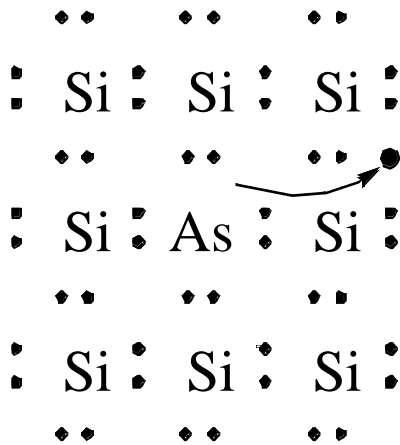


crystal:
energy bands





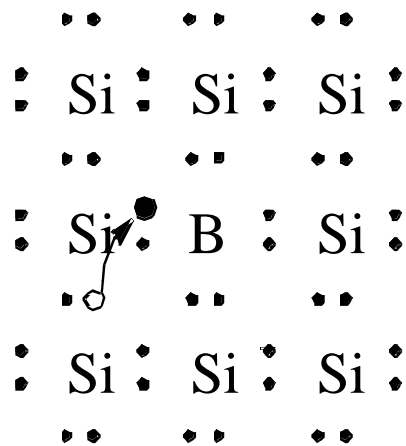
Doping (1)



doping with donors: As
has 5 electrons in the outer
shell → one is not needed
for bonding and becomes a
free electron



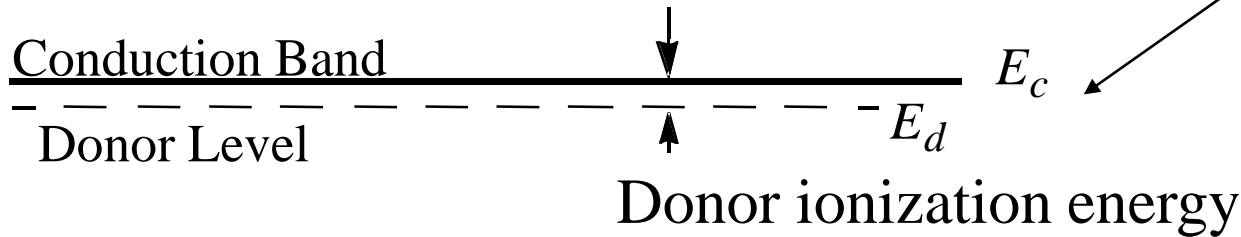
Doping (2)



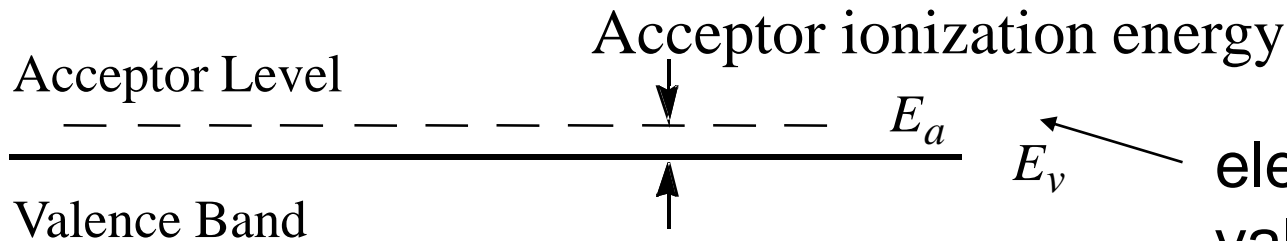
doping with acceptors: B has 3 electrons in the outer shell and this results in a **hole** that can move in the crystal



Doping (3)



electron from the donor easily jumps to the conduction band



electron in the valence band easily captured by the acceptor → a hole is left in the valence band



Doping (4)

It can be shown that

$$np = n_i^2$$

with

$$n_i = \sqrt{N_c N_v} e^{-E_g/2kT}$$

free electron
concentration

hole
concentration

intrinsic density ($\sim 10^{10} \text{cm}^{-3}$
in Si@room temperature)

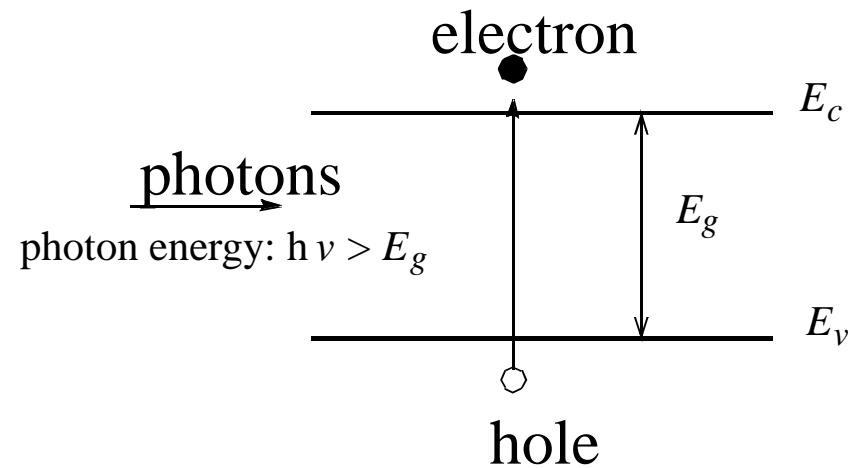
Charge neutrality: $n + N_a = p + N_d$

Example: $N_a = 10^{16} \text{cm}^{-3}$, $N_d = 0 \rightarrow p = 10^{16} \text{cm}^{-3}$, $n = 10^4 \text{cm}^{-3}$

Note: density of atoms in Si is $5 \cdot 10^{22} \text{cm}^{-3}$



e-h generation by radiation



- the photon promotes an electron in the conduction band
- a hole is left in the valence band

- if $h\nu \gg E_g$: number of pairs $\sim h\nu / E_g$
- The same happens with many types of particles, not only photons
- To measure the generated e-h pairs we need to separate electrons and holes \rightarrow **p-n junction**



Outline

- Semiconductors
- **p-n junction**
- Main figures of merits for detectors



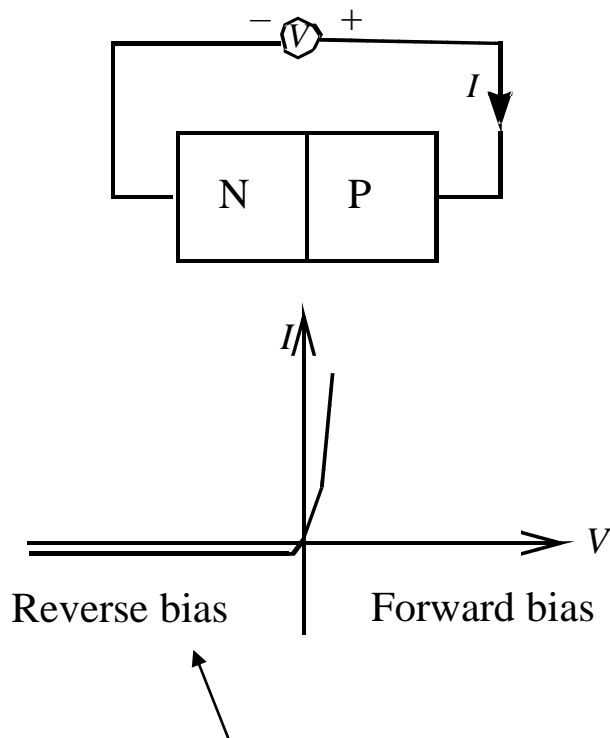
Structure

N-region:

- donor doping;
- plenty of free electrons
- balanced by fixed positive charge

P-region:

- acceptor doping;
- plenty of free holes
- balanced by fixed negative charge

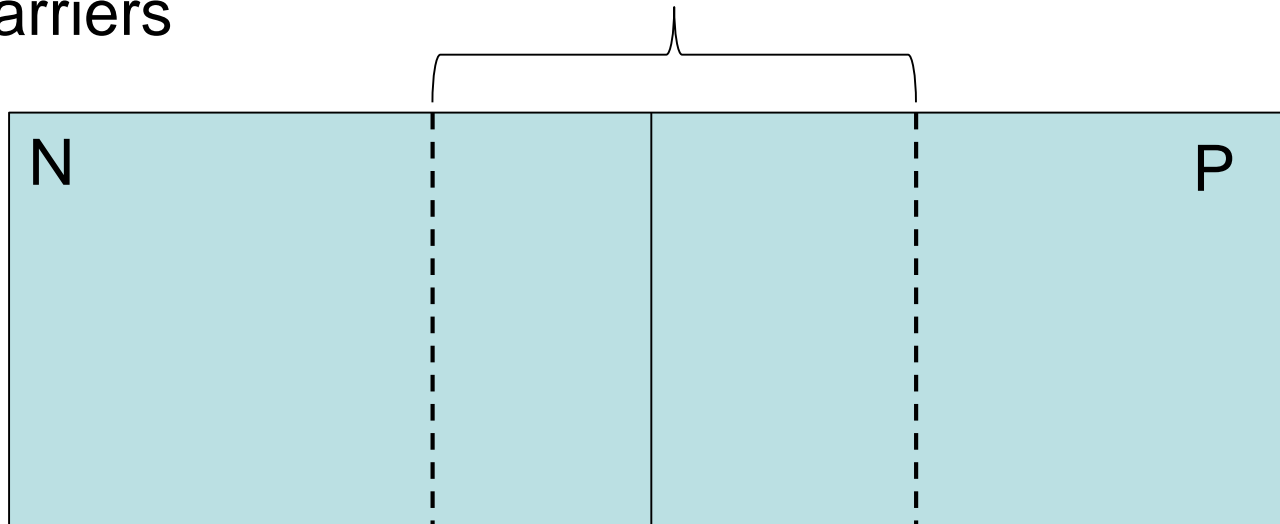


region where the detectors operate



Charge profile (1)

at the border between the N and P region we have very few free carriers



neutral N-region:

- donor doping;
- plenty of free electrons
- balanced by fixed positive charge

depletion region:

only fixed charges

neutral P-region:

- acceptor doping;
- plenty of free holes
- balanced by fixed negative charge



Charge profile (2)

net charge concentration:

$$\rho = q(p - n + N_D - N_A)$$

N			P
$n=N_D$	$n \sim 0$	$n \sim 0$	$p=N_A$
$p \sim 0$	$p \sim 0$	$p \sim 0$	$n \sim 0$
$\rho=0$	$\rho=qN_D$	$\rho=-qN_A$	$\rho=0$

neutral N-region:

- donor doping;
- plenty of free electrons
- balanced by fixed positive charge

depletion region:

only fixed charges

neutral P-region:

- acceptor doping;
- plenty of free holes
- balanced by fixed negative charge



Poisson equation

Gauss's Law:

$$\epsilon_S \mathbf{E}(x + \Delta x)A - \epsilon_S \mathbf{E}(x)A = \rho \Delta x A$$

ϵ_S : permittivity ($\sim 12\epsilon_0$ for Si)

ρ : charge density (C/cm³)

$$\frac{\epsilon_S \mathbf{E}(x + \Delta x) - \epsilon_S \mathbf{E}(x)}{\Delta x} = \rho$$

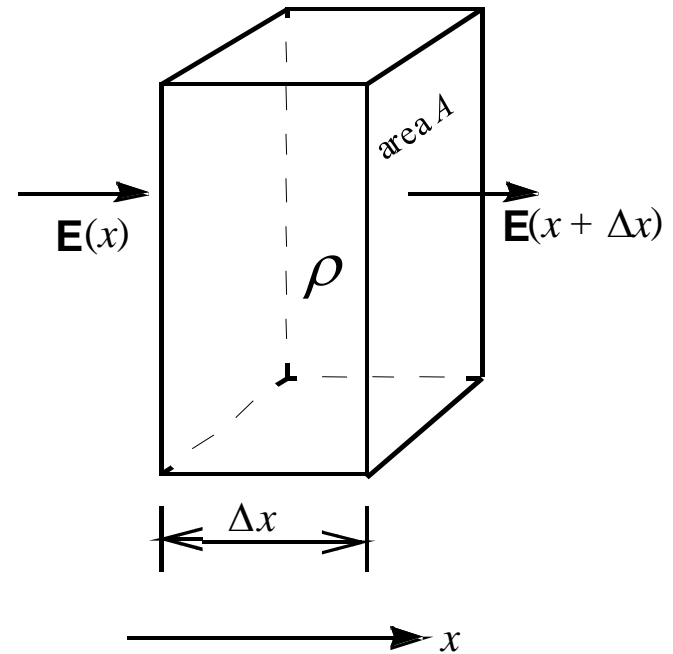


$$\frac{d\mathbf{E}}{dx} = \frac{\rho}{\epsilon_S}$$



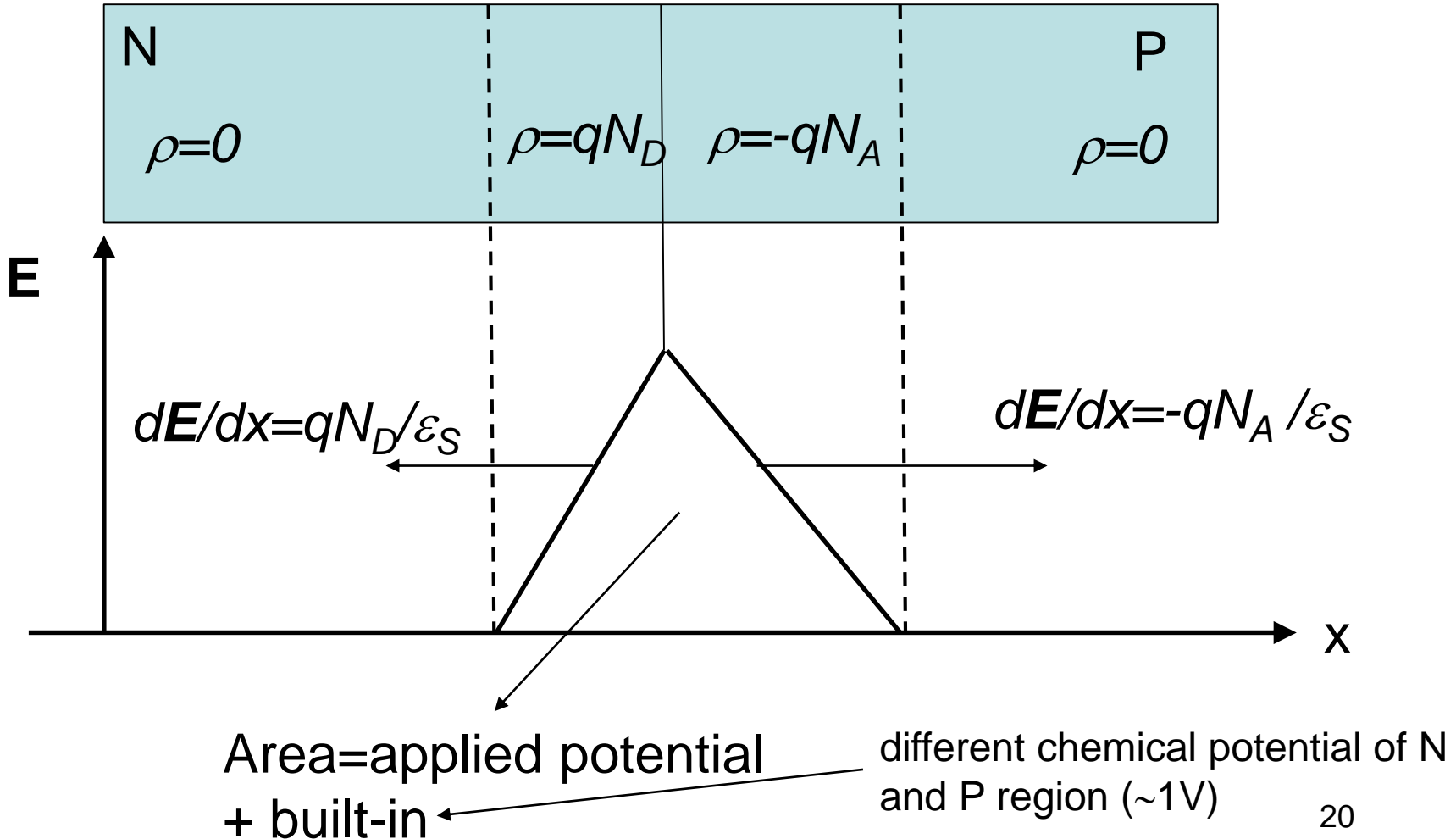
$$\frac{d^2V}{dx^2} = -\frac{d\mathbf{E}}{dx} = -\frac{\rho}{\epsilon_S}$$

Poisson's equation



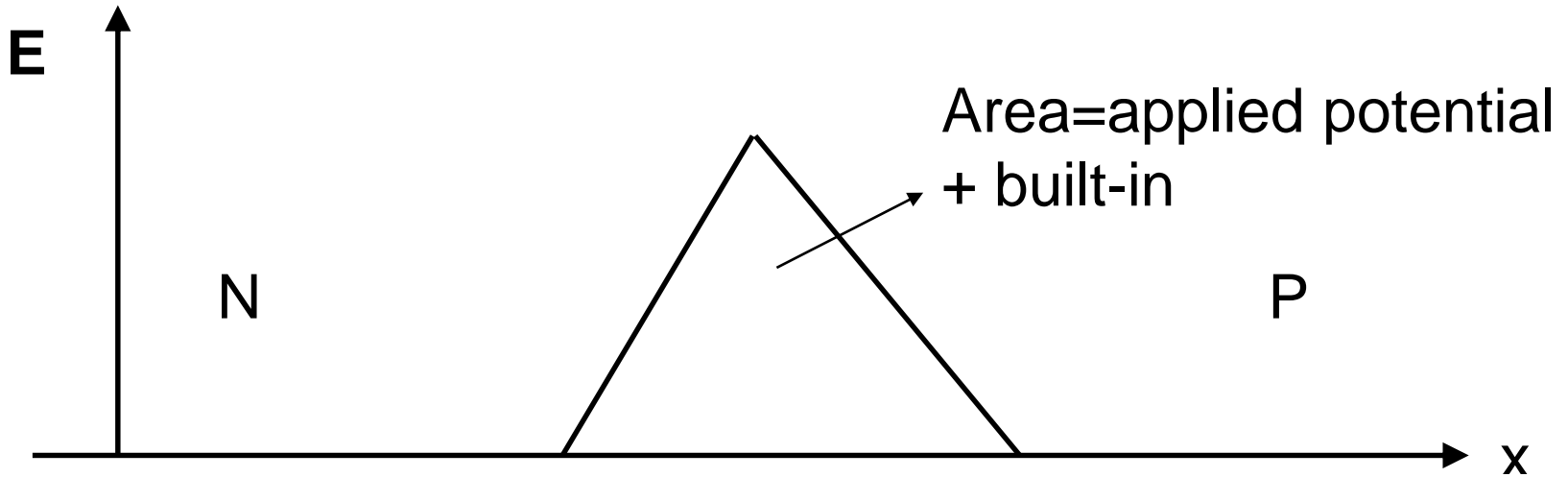


Electric field profile (1)

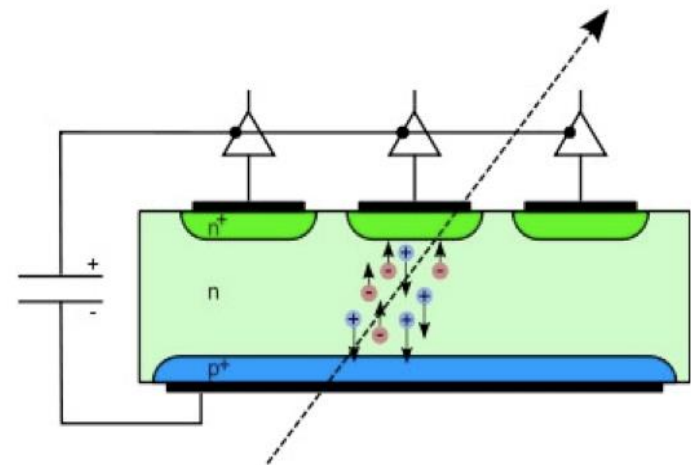




Electric field profile (2)

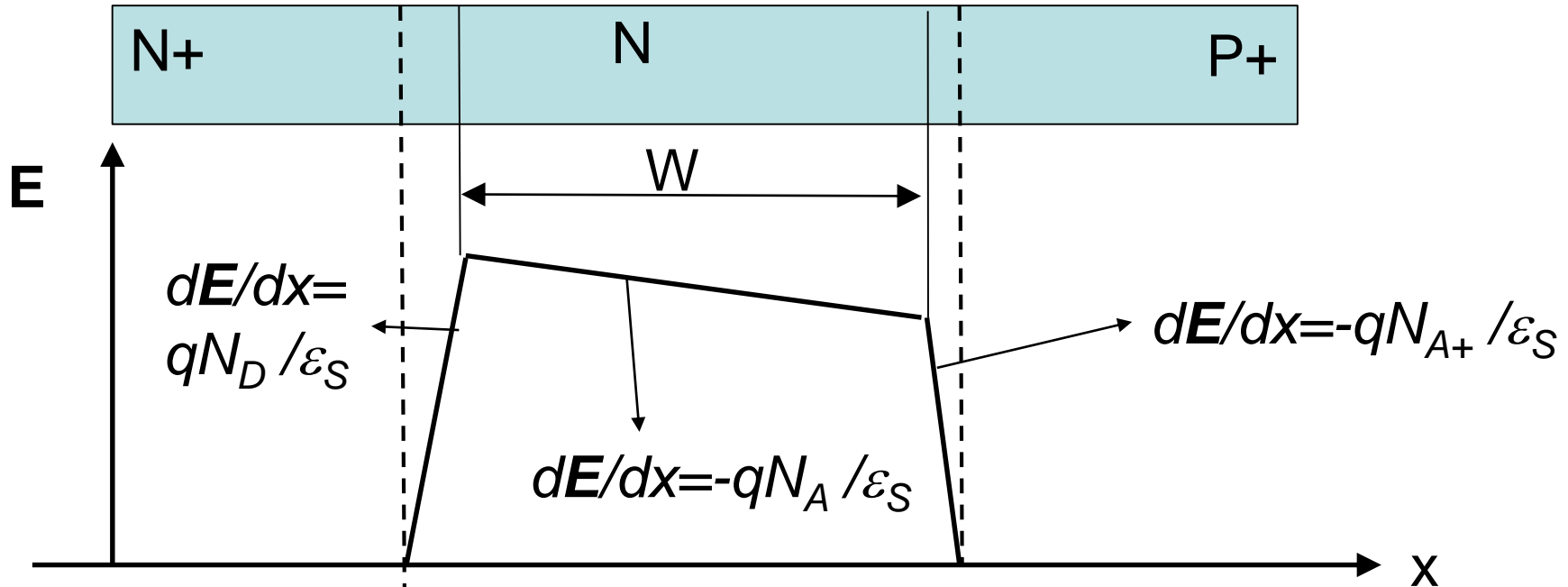


- positive electric field \rightarrow electrons (holes) generated by radiation go to the N (P) region





Electric field profile (3)

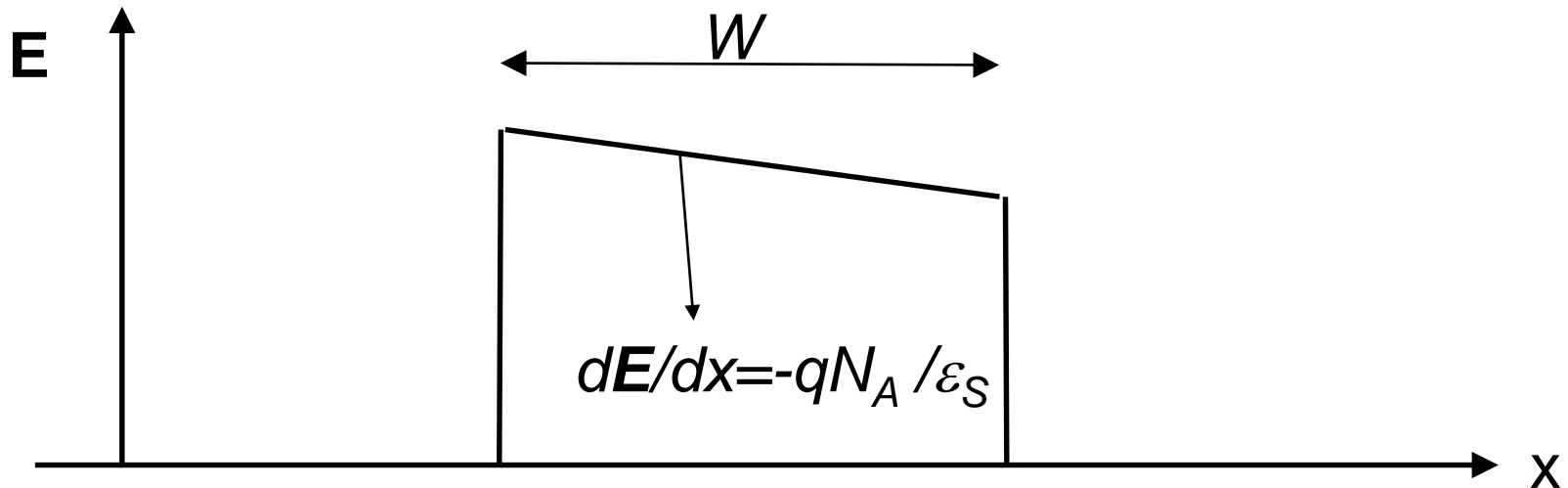


- insertion of a low-doped region: trapezoidal electric field profile \rightarrow same “area” with lower peak electric field



Electric field profile (4)

- doping in the P+ and N+ regions is large so that:



- this however requires: $V + V_{BI} > \frac{qN_A}{\epsilon_S} W^2$ to fully deplete the low doping region



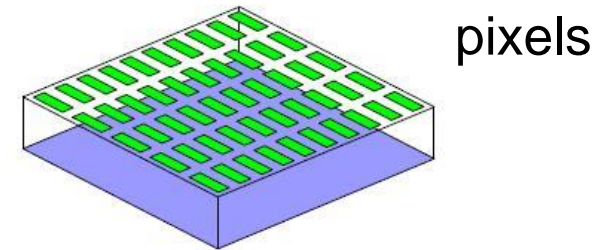
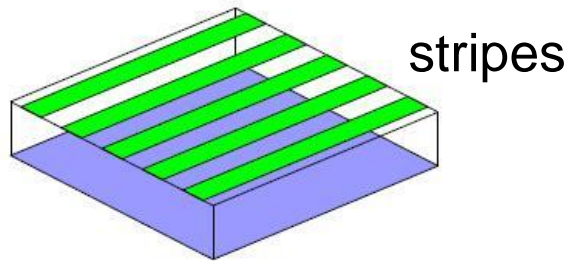
Outline

- Semiconductors
- p-n junction
- **Main figures of merits for detectors**

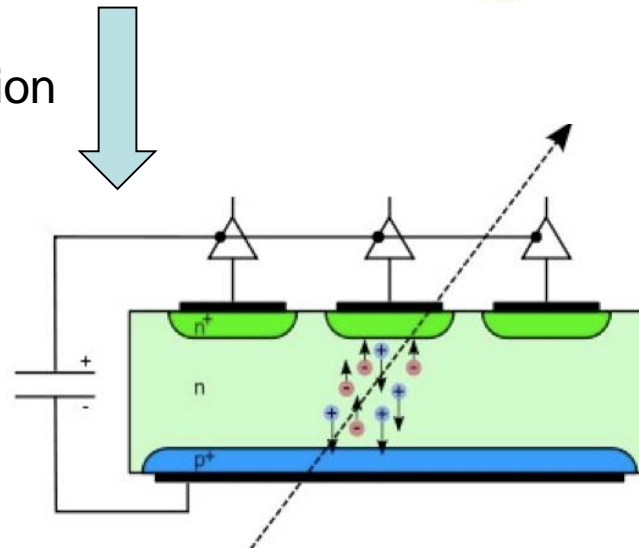


Spatial resolution

- the structure needs to be split in stripes or pixels

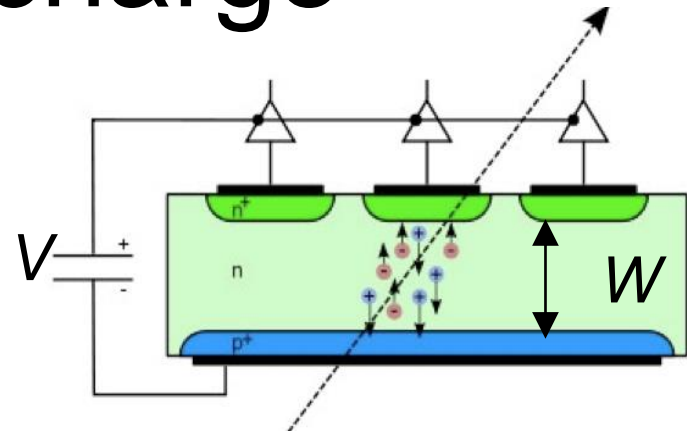


cross section



Collected charge

- the longer the particle travels in the detector the larger the induced charge
- **collected charge proportional to $W \cdot A$** where A is the pixel area.
- large W is required,
- but the condition for depletion



$$V + V_{BI} > \frac{qN_A}{\epsilon_S} W^2$$

implies using **high bias** and **very low doping** levels (high purity crystals)



Dark current

- there are other mechanisms creating e-h pairs;
- e-h creation is assisted by defects;
- simple model:

$$I_{dark} \sim \frac{qWAn_i}{\tau}$$

depletion width → W
area → A
 $n_i = \sqrt{N_c N_v} e^{-E_g / 2kT}$
generation time → τ

- a particle can be detected only if the induced current is larger than the dark current;
- this requires **high purity** material and **low temperature**

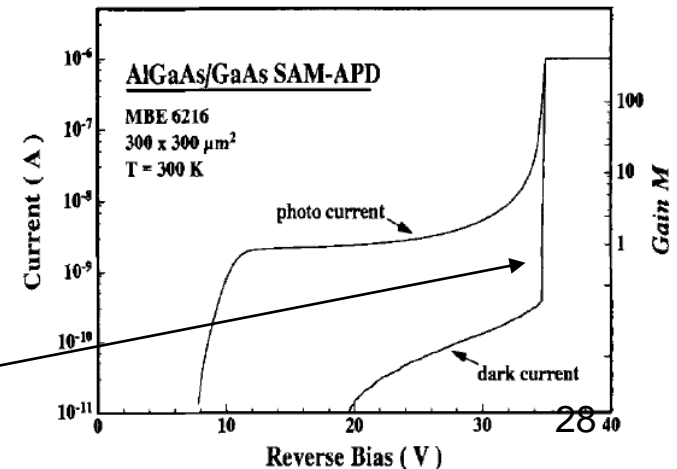
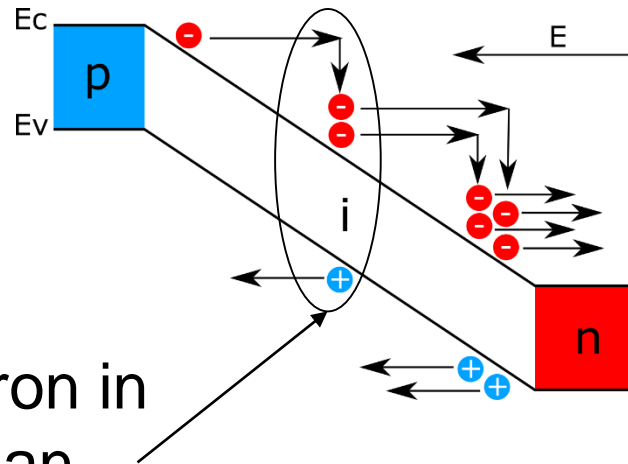


Breakdown voltage (1)

high electric field results in unwanted e-h pair generation by **impact ionization**

a free electron hits an electron in the valence band obtaining an additional free electron + a hole

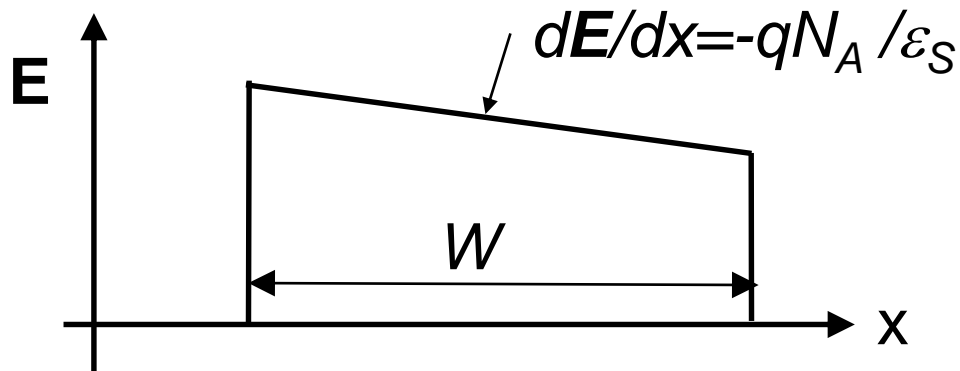
the same process can be initiated by a hole → the dark current increases significantly





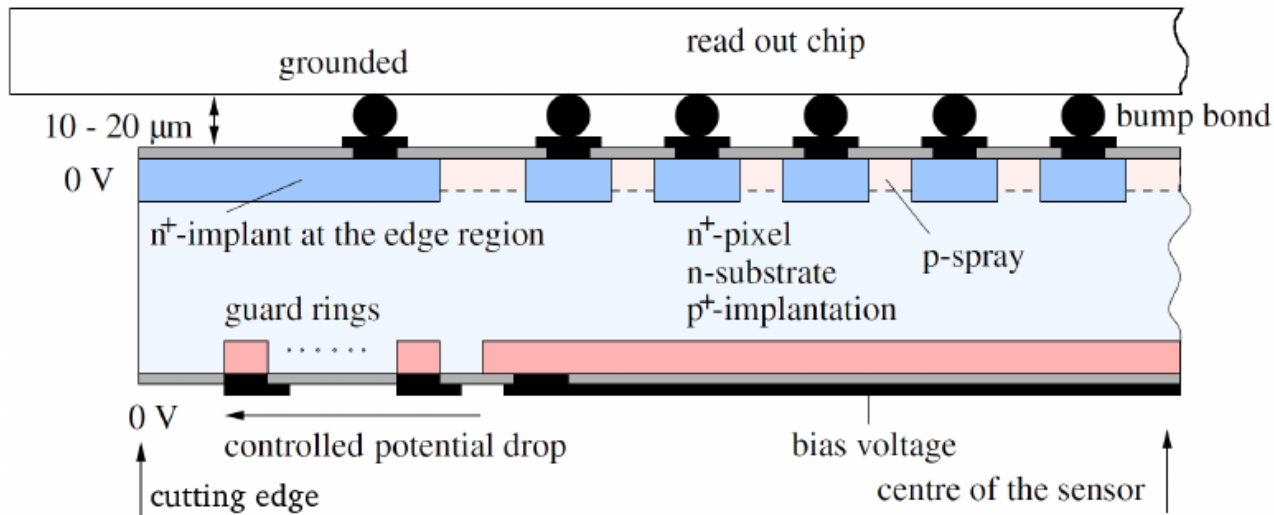
Breakdown voltage (2)

- reduce peak electric field as much as possible → **low doping**

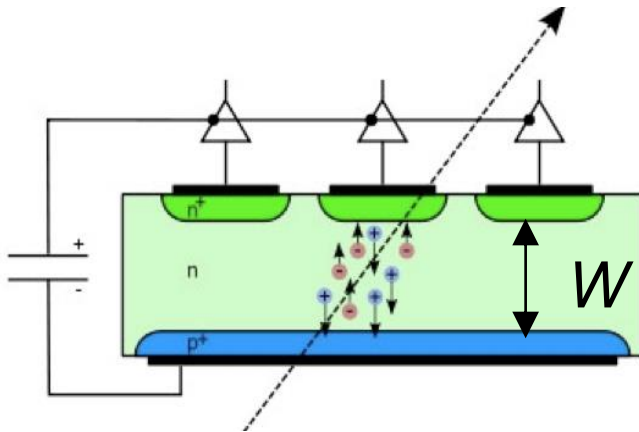


Breakdown voltage (3)

- avoid breakdown at the borders → guard rings



Speed



- generated e-h pairs need to reach the high-doping regions

- the duration of the current pulse induced by a particle is in the order of

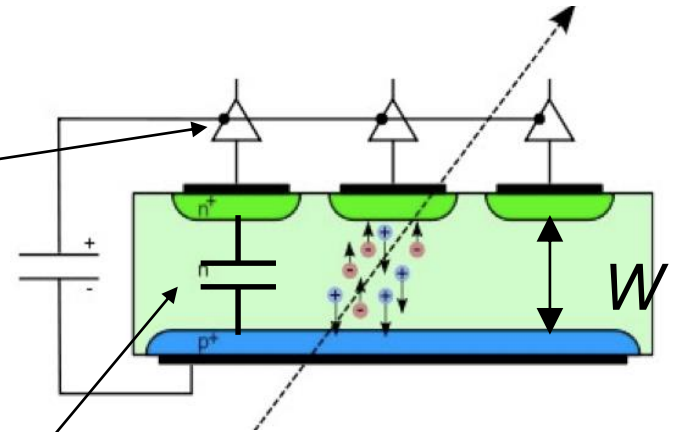
$$T_{pulse} = \frac{W}{\min(v_e, v_h)}$$

carrier velocity $\sim 10^7$ cm/s



Capacitance

- The electronic read-out converts the generated charge into a signal
- it adds noise (unwanted random signal) during the process
- this noise is amplified by the detector capacitance
- C should be as small as possible



$C \sim \epsilon_S \frac{A}{W}$ ← area of the pixel