



# Schroedinger equation

Wave-particles wave function governed by

$$i\hbar \frac{\partial}{\partial t} \psi(t, \mathbf{x}) = \hat{H} \psi(t, \mathbf{x})$$

where  $\hat{H}$  is the "Hamiltonian" (K+V) and

$$\vec{p} \rightarrow -i\hbar \left( \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) = -i\hbar \vec{\nabla}, \quad K = \frac{p^2}{2m} = \dots$$

$$E \rightarrow i\hbar \frac{\partial}{\partial t}$$

Example of a free particle (i.e.  $\hat{H} = \hat{K}$ ) possible solution in one spatial dimension:

$$\psi(t,x) = \exp\left(\frac{\imath}{\hbar}\left(px - \frac{p^2}{2m}t\right)\right)$$

(verify as an exercise)



### Time independent

#### Schroedinger equation

(From here on assume always 1D, i.e.  $x \to x$ )

Under the general assumption

$$\psi(t,x) = \phi(t)\chi(x)$$

the Schroedinger equation can be split into:

$$i\hbar \frac{\partial}{\partial t} \left[\phi(t)\chi(x)\right] = -\frac{\hbar^2}{2m} \frac{\partial}{\partial x^2} \left[\phi(t)\chi(x)\right] + V(t,x)\phi(t)\chi(x)$$

if time independent potential V(t,x) = V(x)

$$\chi(x) i\hbar \frac{\partial}{\partial t} \phi(t) = \phi(t) \left( -\frac{\hbar^2}{2m} \frac{\partial}{\partial x^2} \chi(x) + V(x) \chi(x) \right)$$



### Time independent

Schroedinger equation

$$\frac{1}{\phi(t)} i\hbar \frac{\partial}{\partial t} \phi(t) = \frac{1}{\chi(x)} \left( -\frac{\hbar^2}{2m} \frac{\partial}{\partial x^2} \chi(x) + V(x) \chi(x) \right)$$
$$A(t) = B(x) \quad \Rightarrow A(t) = E = B(x)$$

Which means the time independent Schoredinger equation satisfies the eigenvalue problem:

$$\begin{split} -\frac{\hbar^2}{2m}\frac{\partial}{\partial x^2}\chi(x) & + & V(x)\chi(x) = E\chi(x) \\ E \to \text{"Eigenvalue"} & \chi(x) \to \text{"Eigenfunction"} \end{split}$$

Eigenvalue problem 
$$\hat{H}\chi(x) = E\chi(x)$$



#### Axiomatic formulation

• The state of a quantum mechanical system is completely specified by its wave function

$$\psi(t, \mathbf{x}), P \propto |\psi(t, \mathbf{x})|^2, \quad \int d\mathbf{x} \ dt \ |\psi(t, \mathbf{x})|^2 = 1$$

The wavefunction must be continuous and finite.

**2** To every observable in classical mechanics there corresponds a linear, Hermitian  $(H^{\dagger}=H)$  operator in quantum mechanics (ex:  $E \to -\imath \hbar \vec{\nabla}$ ).



### Axiomatic formulation

**3** Measurements of observables  $\hat{O}$  only outputs corresponding eigenvalues  $o_i$ , where by definition

$$\hat{O}\psi_i = o_i\psi_i$$

Note that a general state doesn't need to be an eigenstate, but can always be written as a superposition

$$\psi = \sum_{i} c_{i} \psi_{i}.$$

The wavefunction immediately "collapses" into the corresponding eigenstate  $\Psi_{\it i}$ 

4 If a system is in a state  $\psi$ , then the average value of the observable corresponding to  $\hat{O}$  is given by

$$<{\it O}>=\int_{-\infty}^{\infty}\psi^{*}\hat{\it O}\psi\;d au$$



#### Axiomatic formulation

**6** The wavefunction of a system evolves in time according to the time-dependent Schroedinger equation

$$\hat{H}\psi(\mathbf{r},t)=i\hbar\frac{\partial\psi}{\partial t}$$

(which must be accepted as a postulate)

The total wavefunction of a system of identical (fermions)bosons must be (anti)symmetric with respect to the interchange of all coordinates of one with those of another. Electronic spin must be included in this set of coordinates.



## Klein Gordon equation

Bosons

What is a generic relativistic equation for a scalar field? Start from Schroedinger:

$$\hat{H}\phi = i\hbar \frac{\partial \phi}{\partial t} \equiv \hat{E}\phi,$$

where now use the relativistic relation:

$$E = \sqrt{p^2c^2 + m^2c^4}$$

and, to avoid  $\sqrt{\ldots}$ , "square" it

Klein-Gordon equation 
$$\left(\Box + \frac{\mathit{m}^2\mathit{c}^2}{\hbar^2}\right) \phi = 0$$

where 
$$\Box = \partial_{\mu}\partial^{\mu}$$
, and  $\partial_{\mu} = \frac{\partial}{\partial x_{\mu}}$ 



### Dirac equation

Fermions

Dirac, 1928: is there a way to "linearize" KG equation? Introduce a set of 4  $\times$  4 matrices  $\gamma_{\mu}$  such that

$$\{\gamma_{\mu}, \gamma_{\nu}\} = 2g_{\mu\nu}$$

where

$$g_{\mu\nu} = diag(1, -1, -1, -1)$$

Dirac equation

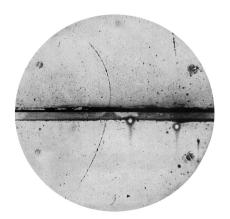
$$(\imath \gamma^{\mu} \partial_{\mu} - m) \psi = 0$$

Predicts antimatter!



### Positron discovery

Anderson, 1932



Cloud chamber photograph by C. D. Anderson of the first positron ever identified. A 6 mm lead plate separates the upper and lower halves of the chamber. The deflection and direction of the particle's ion trail indicate the particle is a positron.