EDMs and bounds on new physics

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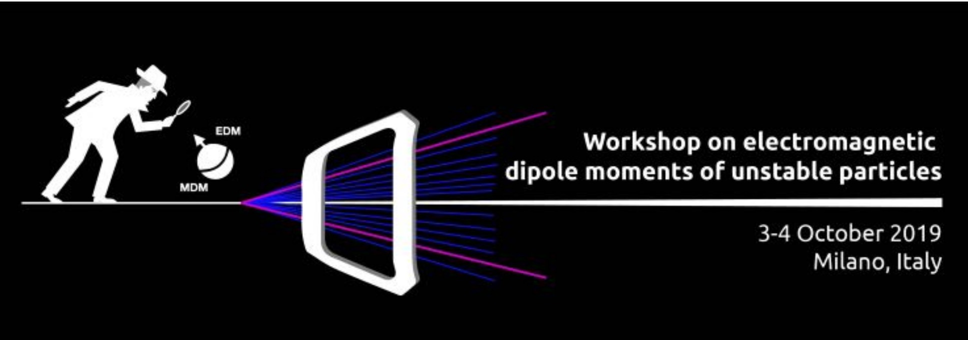
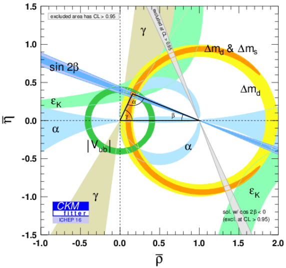


Motivation

Quark-flavour and CP violation in the SM:

- CKM describes flavour and CP violation
- Extremely constraining, one phase
- Especially, K and B physics agree
- Only tensions so far $(R_{K,K^*}, P_5', B \to D^{(*)}\tau\nu, g_\mu - 2, \dots)$

➥ Works well!



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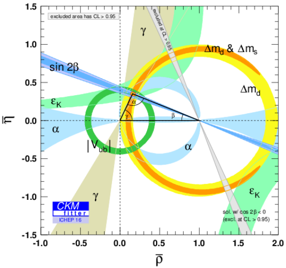
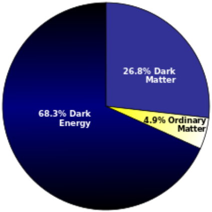


We expect new physics (ideally at the (few-)TeV scale):

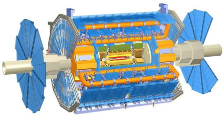
- Baryon asymmetry of the universe
- Hierarchy problem
- Dark matter and energy
- ...

➥ So where is it?



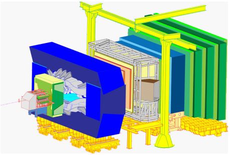
The Quest for New Physics

Three of the main strategies (missing are e.g. ν, DM, astro,. . .):



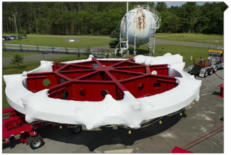
Direct search:

- Tevatron, LHC
- Maximal energy fixed



Indirect search, flavour violating:

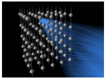
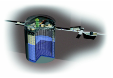
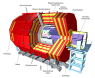
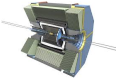
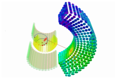
- LHCb, Belle II, BES III, NA62, MEG, . . .
- Maximal reach flexible



Indirect search, flavour diagonal:

- EDM experiments, g-2, . . .
- Maximal reach flexible, complementary to flavour-violating searches

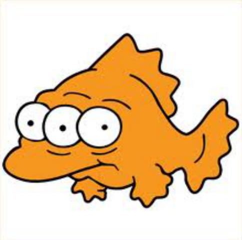
A new era in particle physics!

The curious case of the One-Higgs-Doublet Model

EDMs are finite in the SM...

...but flavour-sector of the SM is special (→):



- Unique connection between Flavour- and CP-violation
- FCNCs highly suppressed, $\sim \Delta m^2/M_W^2$
  - ➥ $\Delta m^2/M_W^2 \sim 10^{-25}$ for ν in the loop!
- FConservingNCs with CPV as well:
  - ➥ $d_e^{SM} \lesssim 10^{-38} e\,\mathrm{cm}$ [Khriplovich/Pospelov '91]

EDMs are quasi-nulltests of the SM!

NP models typically do not exhibit such strong cancellations

- ➥ Background-free precision-laboratories for NP (assuming dynamical solution for strong CP)
- ➥ EDMs $\sim CPV/\Lambda^2$ (interference with SM, *e.g.* LFV $\sim 1/\Lambda^4$)

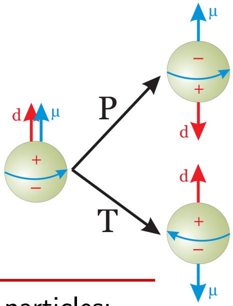
Here: focus as much as possible on model-independent statements

Back to basics: EDMs

Classically: $\mathbf{d} = \int d^3 r \rho(\mathbf{r})\mathbf{r}$, $U = \mathbf{d}\cdot\mathbf{E}$

QM: non-degenerate ground state implies $\mathbf{d} \sim \mathbf{j}$

- $\mathbf{d} \neq \mathbf{0}$ implies T- and P-violation!
- CP-violation for conserved CPT
- Search for linear shift $U = d\,\mathbf{j}\cdot\mathbf{E}$



Non-relativistic neutral system of point-like particles:
Potential EDMs of constituents are shielded! [Schiff'63]

- Sensitivity stems from violations of the assumptions
  - Paramagnetic systems: relativistic enhancement
  - Diamagnetic systems: finite-size effects

Shielding can be reversed, *e.g.* $d_A^{\mathrm{para}} \sim \mathcal{O}(100) \times d_e$!
[Sandars'65,'66]

EDMs and New Physics: Generalities

Sakharov's conditions ('67):
NP models necessarily involve new sources of CPV!

- This does not *imply* sizable EDMs
- However, typically (too) large EDMs in NP models

➥ Generic one-loop contributions excluded
($\rightarrow$ SUSY CP-problem)

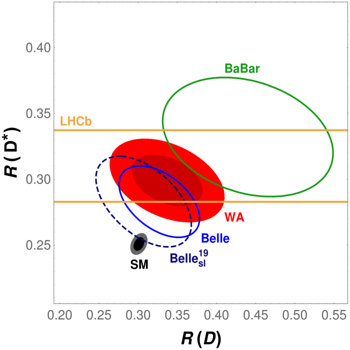
➥ EDMs test combination of flavour- and CPV-structure

EDMs important on two levels:

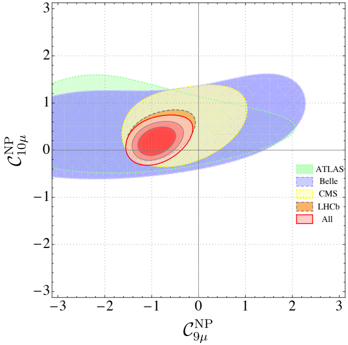
- "Smoking-gun-level": Visible EDMs proof for NP
- Quantitative level:
  Setting limits/determining parameters
  ➥ Theory uncertainties are important!

Flavour anomalies and EDMs

$b \to c\tau\nu$ [Murgui+'19]



$b \to s\ell^+\ell^-$ [Algueró+'19]



- Presently $\sim 3\sigma$ and $\sim 5\sigma$ from SM predictions
- No indication of CPV
  - ➥ Why is this relevant for EDMs?

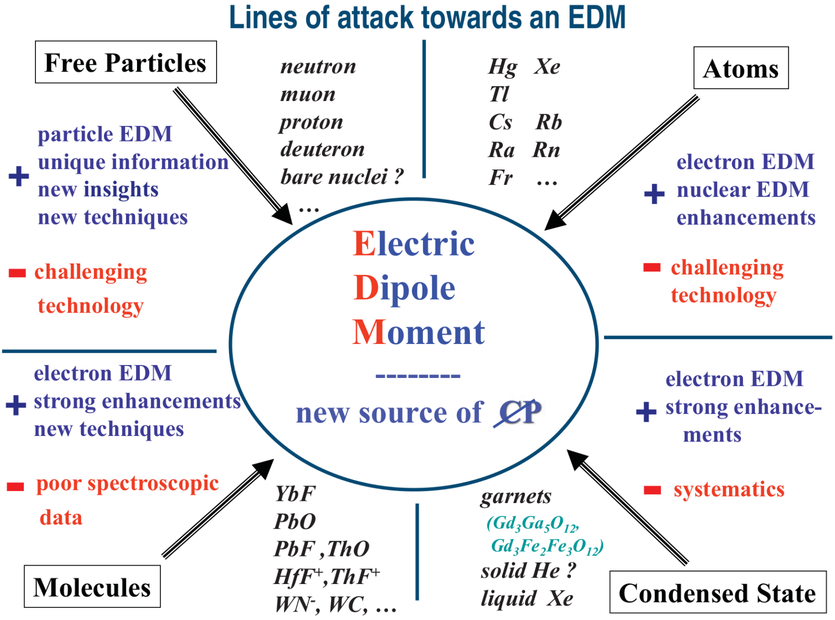
➥ Both imply lepton-flavour-non-universality (LFNU)!

➥ Often implicitly assumed in NP scenarios (at least in the past)

➥ Decouples e, μ, τ EDMs, no scaling with masses

- ➥ Increased importance of explicit μ, τ-EDM measurements!

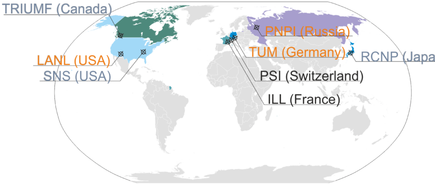
Experimental approaches [K. Jungmann'13 in Annalen der Physik]



Experimental status [see talk by P. Schmidt-Wellenburg]

Neutron EDM:

- $|d_n| \leq 3.6 \times 10^{-26} e\,\mathrm{cm}\,(95\%\mathrm{CL})$
  [Pendlebury+'15,Baker+'06]
- Worldwide effort aiming at $(10 \rightarrow 0.1) \times 10^{-27} e\,\mathrm{cm}$
- UCN sources critical problem



[P.Schmidt-Wellenburg'16]

Paramagnetic systems:

- Atomic: $|d_{\mathrm{Tl}}| \leq 9.6 \times 10^{-25} e\,\mathrm{cm}\,(95\%\,\mathrm{CL})$ [Regan+'02]
- Molecular: $|\omega_{\mathrm{ThO}}| \leq 1.1\,\mathrm{mrad/s}\,(95\%\,\mathrm{CL})$ [ACME'18]
- Ionic: HfF^+, $|\omega_{\mathrm{HfF}}| \leq 7.9\,\mathrm{mrad/s}\,(90\%\,\mathrm{CL})$ [Cairncross+'17]

Diamagnetic systems:

- $|d_{\mathrm{Hg}}| \leq 7.4 \times 10^{-30} e\,\mathrm{cm}\,(95\%\,\mathrm{CL})$ [Graner+'16]
- Ongoing: Xe, Hg, exploit octupole deformation, *e.g.* Ra, Rn,. . .

Solid state systems: $|d_e| \leq 6.1 \times 10^{-24-25} e\,\mathrm{cm}$ [Eckel+'12,Kim+'15]

Storage rings: $|d_\mu| \leq 1.9 \times 10^{-19} e\,\mathrm{cm}$ [Bennett+'08]

Collider: $|d_\tau| \leq 3.4 \times 10^{-17} e\,\mathrm{cm}$ [Belle'03]

Relating NP parameters and experiment

- Most stringent constraints from neutron, atoms and molecules
  ➥ Shielding typically applies

Atomic level

$\Downarrow$

Nuclear Level

$\Downarrow$

Hadronic level

$\Downarrow$

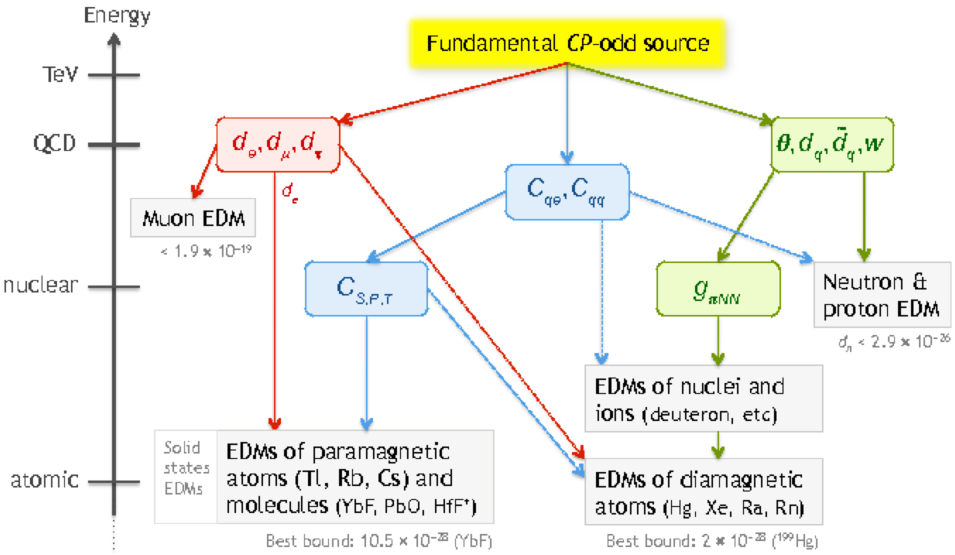
Effective Theory with (C)EDMs of fermions, $\mathcal{O}_W, \ldots$

$\Downarrow$

Parameters of your favourite NP model

- Each step potentially involves large uncertainties!
- 4/5 model-independent $\Rightarrow$ series of EFTs [e.g. deVries+'11]
- Limits usually displayed as allowed regions
  ➥ Conservative uncertainty estimates important

Schematic EFT framework [Pospelov/Ritz'05, Hoecker'12]



The EDM in heavy paramagnetic systems

Two main contributions, enhanced by Z^3: [Sandars'65, Flambaum'76]

➡ A single measurement does not restrict d_e directly

- C_S: CP-odd electron-nucleon interaction
- Atoms: typically polarized in external field
- Molecules: aligned in external field
  ➡ Exploit huge internal field

For molecules: energy shift $\Delta E = \hbar\omega$ with

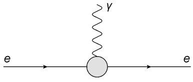
$$\omega_M[\mathrm{mrad/s}] = \alpha_M^{d_e} d_e + \alpha_M^{C_S} C_S\,.$$

Molecule	$\alpha_M^{d_e}/10^{-27}e\mathrm{cm}$	$\alpha_M^{C_S}/10^7$
HfF^+	34.9 ± 1.4	32.0 ± 1.3
ThO	120.6 ± 4.9	181.6 ± 7.3

[Results entering: Skripnikov'17,Fleig'17,Denis/Fleig'16,Skripnikov'16

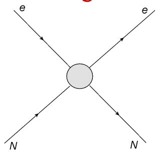
Averages: Fleig/MJ'18]

d_e



$\bar{e}F_{\mu\nu}\sigma^{\mu\nu}\gamma_5 e$

C_S

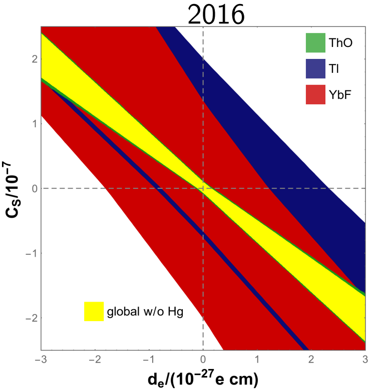


$(\bar{e}i\gamma_5 e)(\bar{N}N)$

Model-independent extraction of d_e and C_S

In principle: two unknowns, three measurements (Tl,YbF,ThO)

➥ Extract d_e, C_S model-independently [Dzuba et al.'11,MJ'13]



Problem: Aligned constraints

➥ weak limits

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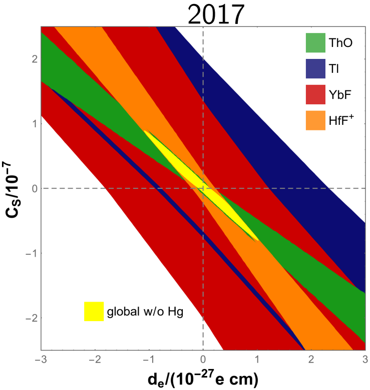

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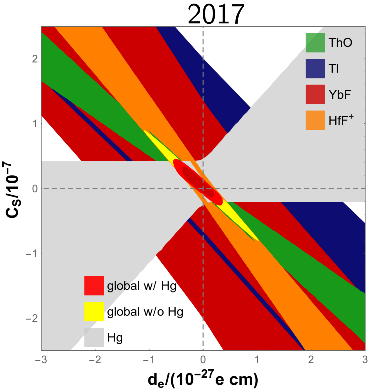

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Partial resolution: HfF^+ result

Mercury bound ∼ orthogonal!

Assumption: C_S, d_e saturate $d_{\rm Hg}$

➥ Conservative

[Fleig,MJ'18]

$d_e \leq 3.8 \times 10^{-28} e\ \mathrm{cm}$

$C_S \leq 2.7 \times 10^{-8}$

Yields model-independent limit on every paramagnetic system!

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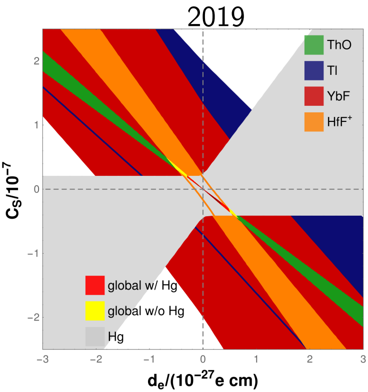

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Future measurements aim at precision beyond present constraints!

- Help to resolve the alignment problem
- Requires precision measurements of low-Z and high-Z elements

EDMs of diamagnetic systems and nucleons

Situation more complicated than for paramagnetic systems:

- Potential SM contribution: $\bar{\theta}$ ($\rightarrow$ strong CP puzzle)
- Contributions from $\bar{\theta}, d_q, \tilde{d}_q, w, C_{S,P,T}, C_{qq}$
  - ➥ Interpretation usually model-dependent
    (for model-independent prospects: [Chupp/Ramsey-Musolf'14])

Complementary measurements, different sources possible/likely

- $|d_{Hg}| \leq 7.4 \times 10^{-30} e\,\mathrm{cm}$ [Graner et al. '16] , very constraining
  Problem: QCD and nuclear theory uncertainties (x00%!)
  - ➥ No conservative constraint on CEDMs left! [MJ/Pich'13]
- $|d_n| \leq 3.6 \times 10^{-26} e$ cm [Pendlebury'15]
  Theory in better shape, still $\mathcal{O}(100\%)$ uncertainties
  [Pospelov/Ritz'01,Hisano et al'12,Demir et al'03,'04,de Vries et al'11]

Progress in theory necessary to fully exploit these measurements
Unique: orders-of-magnitude improvement w/o new measurement!

The role of Mercury in determining the electron EDM

Mercury is a diamagnetic system, many contributions

➥ Why is it shown in the paramagnetic global fit? [MJ'13]

- Shielding of C_S and d_e effective (even vanishing at LO)
  - ➥ Schiff moment contribution expected to be dominant
  - ➥ d_e, C_S only a fraction of the total EDM
- ➥ Assuming d_e, C_S to saturate the exp. limit is conservative

New calculation of the $\mathbf{C_S}$ coefficient [Fleig/MJ'18]

LO contribution vanishes

➥ Triple perturbative expansion necessary:

1. External electric field (here: included in basis set)
2. Hyperfine splitting
3. d_e/C_S

$$\alpha_{C_S} = -2.8(6) \times 10^{-22}\, e\ \text{cm}$$

α_{d_e} w.i.p., so far old calculation [Martensson-Pendrill/Oster'85] + conservative error estimate

The importance of multiple measurements

Only pattern of CPV observables allows for model-differentiation!
➥ There is no single "best" measurement!

Paramagnetic systems:

- 1 significant measurement NP
- 2 determine ideally d_e *and* C_S
- More for consistency (unless MQM is relevant)

Diamagnetic systems, nucleons/baryons, light nuclei:

- 1 significant measurement: $\bar{\theta}$ possible explanation
- 2 should tell $\bar{\theta}$ from other sources
- Many more to identify model-independently CPV strucuture

➥ We need as many measurement as possible!
➥ Ideally very different systems
➥ Try to find P-,T-odd measurements besides EDMs

EDMs in NP Models

EDM constraints forbid generic CPV contributions up to two loops
➥ huge scales or highly specific structure!

- hardly testable elsewhere
- simple power-counting insufficient (UV sensitivity)

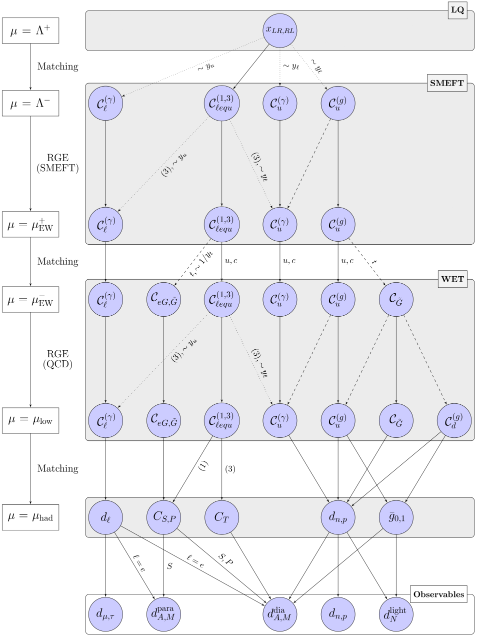
➥ Model-independent analyses difficult

EDMs unique, both blessing and curse

- some model-independent relations exist, *e.g.* to β decay [Khriplovich'91, see also e.g. Dekens/Vos'15]
- strong (model-dependent) constaints of related observables

➥ Consider models or subsets of model-independent framework

EDMs in sLQ models [Dekens/de Vries/MJ/Vos'18]



Cascade of EFTs:

Example: R_2 LQ

Tree-level: semileptonic operators

1-loop (matching + running):
Dipole operators are generated

Below $\mu_{\rm EW}$: gluonic operators added

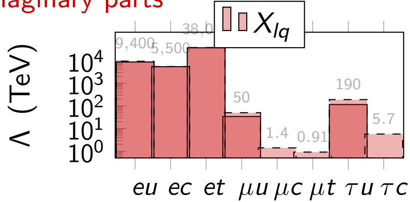
$\mu_{\rm low} \sim 1$ GeV: $\rightarrow$ hadronic operators

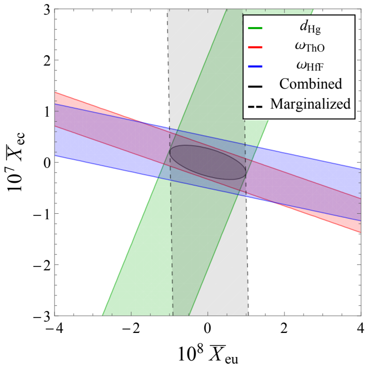
- enter EDM calculations
  ($\rightarrow$ atomic + nuclear MEs)
- MEs have large uncertainties

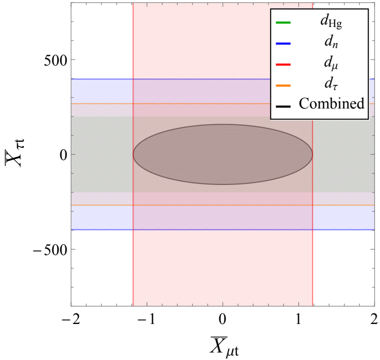
Phenomenological consequences

Most observables constrain (mainly) real parts
- EDMs constrain complementarily imaginary parts

Flavour-dependence of constraints
- Vastly different magnitudes
- Most relevant observables differ
- Complementarity of measurements!



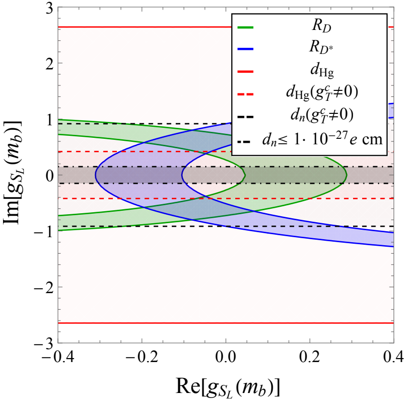




Relation to $R(D) - R(D^*)$ flavour anomaly

R_2 LQ part of NP model for flavour anomalies: [Bečirević+'18]

- Generates $C_{S_L} \sim 4C_T$ ($@\mu_{\rm LQ}$)
- Explanation of $R(D^{(*)})$ possible, but requires imaginary part
- The same coupling combination yields $(\bar{c}\sigma^{\mu\nu}\gamma_5 c)(\bar{\tau}\sigma_{\mu\nu}\tau)$
  - ➥ Generates charm ($+\ \tau$) EDMs + Weinberg operator
  - ➥ Bounds from neutron + Hg EDMs



2 effects:

1. Weinberg operator: smaller effect (outer line)
2. Charm EDM: depends on charm tensor-current neutron ME
   1 calculation [Alexandrou+'17]
   ➥ compatible with 0

Future EDM experiments or lattice can improve this

Conclusions

- EDMs unique tests of NP models
- Model-independent constraints on NP parameters difficult
  - ➥ Need (at least) as many experiments as (eff.) parameters
- Quantitative results require close look at theory uncertainties
  - ➥ Use conservative limits, allowing for cancellations
  - ➥ For *e.g.* $d_n, d_{\rm Hg}$ bottleneck! Chance for nuclear theory
- Robust, model-independent limit on electron EDM (C_S not model-independently negligible):

$$|d_e| \leq 3.9 \times 10^{-28} e\,\mathrm{cm} \quad (95\%\,\mathrm{CL})$$

- Flavour anomalies killed LFU paradigm
  - ➥ Increased importance of μ, τ EDM
- EDMs in scalar LQ models
  - ➥ Demonstrate this point
  - ➥ Every measurement important for at least one coupling!
- Plethora of new results to come
  - ➥ Might turn limits into determinations!

Backup slides

- EDM EFT framework
- 2HDM Framework
- Limits on $|d_e|$ and $|C_S|$
- Expected limits from paramagnetic systems

Framework

Effective Lagrangian at a hadronic scale:

$$\mathcal{L} = - \sum_{f=u,d,e} \left[\frac{d_f^\gamma}{2} \mathcal{O}_f^\gamma + \frac{d_f^C}{2} \mathcal{O}_f^C \right] + C_W \mathcal{O}_W + \sum_{i,j=(q,l)} C_{ij} \mathcal{O}_{ij}^{4f} \,,$$

in the operator basis

$$\mathcal{O}_f^\gamma = ie\bar{\psi}_f F^{\mu\nu} \sigma_{\mu\nu} \gamma_5 \psi_f \,, \qquad \mathcal{O}_f^C = ig_s \bar{\psi}_f G^{\mu\nu} \sigma_{\mu\nu} \gamma_5 \psi_f \,,$$

$$\mathcal{O}_W = +\frac{1}{3} f^{abc} G^a_{\mu\nu} \tilde{G}^{\nu\beta,b} G_\beta{}^{\mu,c} \,, \qquad \mathcal{O}_{ij}^{4f} = (\bar{\psi}_i \psi_i)(\bar{\psi}_j i \gamma_5 \psi_j)$$

Options for matrix elements:

- Naive dimensional analysis[Georgi/Manohar '84] : only order-of-magnitude estimates
- Baryon χPT: not applicable for all the operators
- QCD sum rules: used here [Pospelov et al.] , uncertainties large

Framework for 2HDM contributions

In 2HDMs, CPV in new interactions can generate EDMs!

Parametrization for $H^\pm$ Yukawas, ς_i complex:

$$\mathcal{L}_Y^{H^\pm} = -\frac{\sqrt{2}}{v}\, H^+ \left\{ \bar{u} \left[V \varsigma_d M_d \mathcal{P}_R - \varsigma_u\, M_u^\dagger V \mathcal{P}_L \right] d + \bar{\nu} \varsigma_l M_l \mathcal{P}_R l \right\} + \text{ h.c.}$$

- General for coupling matrices ς_i (M_i choice of normalization)
- Numbers ς_i: Aligned 2HDM [Pich/Tuzon'09,MJ/Pich/Tuzon'10]
- Easily matched on your favourite model

For mass eigenstates $\varphi_i^0 = \{h, H, A\}$, $\mathcal{M}^2_{\rm diag} = \mathcal{R}\mathcal{M}^2\mathcal{R}^T$, we have

$$\begin{aligned} \mathcal{L}_Y^{\varphi_i^0} &= -\frac{1}{v} \sum_{\varphi, f} \varphi_i^0\, \bar{f}\, y_f^{\varphi_i^0} M_f \mathcal{P}_R f \; + \text{ h.c.}\,, \\ y_f^{\varphi_i^0} &= \mathcal{R}_{i1} + (\mathcal{R}_{i2} \pm i\, \mathcal{R}_{i3}) \left(\varsigma_{F(f)}^{(*)} \right)_{ff} \quad \text{for} \quad F(f) = d, l(u)\,. \end{aligned}$$

For neutrals: additional CPV contributions from the potential!

Why 2HDM?

Model-independent NP analysis: Too many parameters in general

EW symmetry breaking mechanism still not completely fixed:

- 1HDM minimal and elegant, but "unlikely" (SUSY,GUTs,...)
- 2HDM "next-to-minimal":
  - ρ-parameter "implies" doublets
  - low-energy limit of more complete NP models
    ➥ Model-independent element
  - simple structure, but interesting phenomenology
  - important effects in flavour observables
- Plethora of 2HDMs:
  ➥ differ in their suppression mechanism for FCNCs

Could explain tensions in the flavour sector (e.g. $B \to D^{(*)}\tau\nu$)

Not an attempt at a complete theory!

Framework for 2HDM contributions

The CPV interactions of the 2nd doublet can generate EDMs

General parametrization for $H^\pm$ Yukawas, ς_i complex matrices:

$$\mathcal{L}_Y^{H^\pm} = -\frac{\sqrt{2}}{v} H^+ \left\{ \bar{u} \left[V \varsigma_d M_d \mathcal{P}_R - \varsigma_u M_u^\dagger V \mathcal{P}_L \right] d + \bar{\nu} \varsigma_l M_l \mathcal{P}_R l \right\} + \text{ h.c.}$$

- Induce couplings like W-exchange, just with a charged Higgs ($M_{H^\pm} \gtrsim m_t$)
- Easily matched on your favourite model
  - ➥ M_i only choice of normalization
- $\varsigma_i \to$ numbers: Aligned 2HDM [Pich/Tuzon'09,MJ/Pich/Tuzon'10]
  - ➥ Comparisons with flavour data in this model

Neutral Higgs exchanges: couplings $y_i^0(\varsigma_i, V)$

➥ Additional CPV contributions from the potential

➥ Analysis depends on many unknown parameters

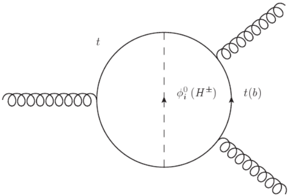
EDMs in 2HDMs

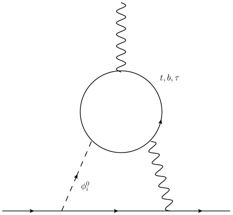
From necessary flavour suppression for a viable model:

- One-loop (C)EDMs: controlled (not tiny) [e.g. Buras et al. '10]
- 4-quark operators small (no $\tan^3\beta$-enhancement)

➥ Two-loop graphs dominant
[Weinberg '89, Dicus '90, Barr/Zee '90, Gunion/Wyler '90,. . .]

  - Weinberg diagram important for neutron EDM
  - Barr-Zee(-like) diagrams dominate other EDMs





Paramagnetic systems: tree-level can be relevant ($C_S \times Z^3$)
(light-quark mass $\times$ tree) vs. (top mass $\times$ two-loop)

Neutral Higgs contributions in general 2HDMs [MJ/Pich'13]

Contributions typically involve the following sum:
(f,f': fermions, F(f): family of the fermion)

$$\sum_i \mathrm{Re}\left(y_f^{\varphi_i^0}\right)\mathrm{Im}\left(y_{f'}^{\varphi_i^0}\right) = \pm\,\mathrm{Im}\left[(\varsigma^*_{F(f)})_{ff}(\varsigma_{F(f')})_{f'f'}\right]$$

- R.h.s. independent of the Higgs potential
- Vanishes for equal fermions (universality: equal family)
- Modified by mass-dependent weight factors...
  - ➥ but holds for degenerate masses and decoupling limit

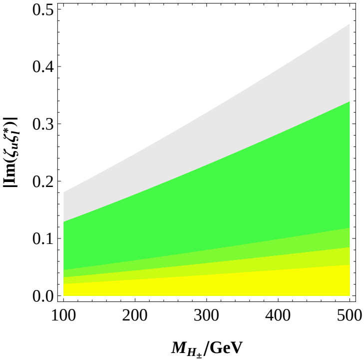
CPV in the potential tends to have smaller impact

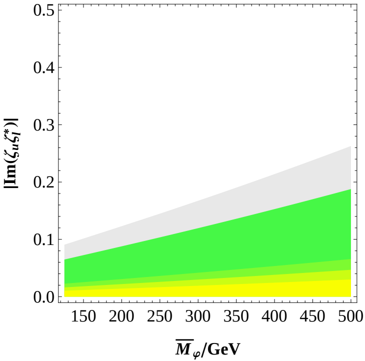
➥ Approximation for phenomenological analysis:

$$\sum_i f(M_{\varphi_i^0})\mathrm{Re}\left(y_f^{\varphi_i^0}\right)\mathrm{Im}\left(y_{f'}^{\varphi_i^0}\right) \to \pm\, f(\overline{M}_{\varphi})\mathrm{Im}\left[(\varsigma^*_{F(f)})_{ff}(\varsigma_{F(f')})_{f'f'}\right].$$

Bounds from the electron EDM

- Contributions via Barr-Zee diagrams [Bowser-Chao et al.'97]
- Sensitivity to $d_e \sim \mathrm{Im}(\varsigma^*_{u,33}\varsigma_{l,11})$
- Bounds $\mathrm{Im}(\varsigma^*_u \varsigma_l) \lesssim \mathcal{O}(0.05)$
  - Strong despite two-loop suppression and mass factors
- Implies $\mathrm{Im}(\varsigma_l \varsigma^*_u)/M^2_{H^\pm} \leq \times 10^{-5}\mathrm{GeV}^{-2}$ (universal ς_i's)
  - A factor 1000 stronger than (semi)leptonic constraints!

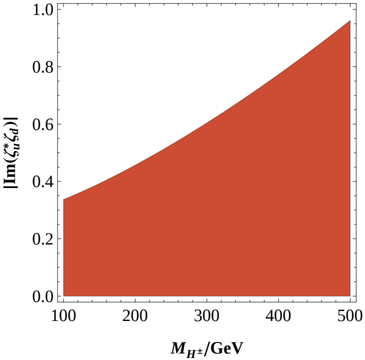


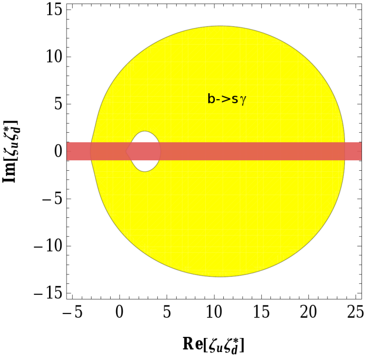


Bounds from the neutron EDM

- Size of Weinberg (charged) and Barr-Zee (neutral) similar
- So far no fine-tuning necessary
- Next-generation experiments will test critical parameter space
- Constraint from Hg potentially a few times stronger
- Comparison with $b \to s\gamma$: large impact! [MJ/Pich'14,MJ/Li/Pich'12]

  ➡ EDMs restrict CPV in other modes





Theory uncertainties and the EDM of Mercury

- Extremely precise atomic EDM limit:
  $|d_{Hg}| \leq 3.1 \times 10^{-29} e\,\mathrm{cm}$ [Griffith et al. '09]
- However: difficult diamagnetic system
  - Shielding efficient $\rightarrow$ sensitivity $\sim d_n, d_{Tl}$

$$d_{Hg} \overset{Atomic}{=} d_{Hg}(S, C^N_{S,P}) \overset{Nuclear}{=} d_{Hg}(\bar{g}_{\pi NN}, C^{p,n}_{S,P})$$
$$\overset{QCD}{=} d_{Hg}(d^C_f, C_{qq'}, C^q_{S,P})$$

  - Uncertainties:
    Atomic$\sim$ 20%, Nuclear$\sim x$00%, QCD sum rules$\sim$ 100 − 200%
  - ➡ No conservative constraint on CEDMs left! [MJ/Pich'13]

$$d_{\mathrm{Hg}} = \Big\{-(1.0 \pm 0.2)\Big((1.0 \pm 0.9)\,\bar{g}^{(0)}_{\pi NN} + 1.1\,(1.0 \pm 1.8)\,\bar{g}^{(1)}_{\pi NN}\Big) + (1.0 \pm 0.1) \times 10^{-5}\,[-4.7\,C_S + 0.49\,C_P]\Big\} \times 10^{-17}\; e\,\mathrm{cm}\,,$$

Progress in theory necessary to fully exploit precision measurements of diamagnetic EDMs

The EDM of the Neutron

Explicit expressions for the neutron EDM [MJ/Pich'13 (refs therein)]

$$
\begin{aligned}
d_n\left(d_q^\gamma, d_q^C\right)/e &= \left(1.0^{+0.5}_{-0.7}\right)\ \Big[1.4\left(d_d^\gamma(\mu_h) - 0.25\, d_u^\gamma(\mu_h)\right) \\
&\quad + 1.1\ \left(d_d^C(\mu_h) + 0.5\, d_u^C(\mu_h)\right)\Big]\ \frac{\langle\bar{q}q\rangle(\mu_h)}{(225\ \mathrm{MeV})^3}\,, \\
|d_n(C_W)/e| &= \left(1.0^{+1.0}_{-0.5}\right)\ 20\ \mathrm{MeV}\ C_W\,, \\
|d_n(C_{bd})/e| &= 2.6\ \left(1.0^{+1.0}_{-0.5}\right) \times 10^{-3}\ \mathrm{GeV}^2 \left(\frac{C_{bd}(\mu_b)}{m_b(\mu_b)} + 0.75\,\frac{C_{db}(\mu_b)}{m_b(\mu_b)}\right)\,.
\end{aligned}
$$

Chances and challenges for nuclear theory

Some more detail:

- Measurements with neutral atoms (now) or ions (future)
- Atomic theory relates d_A to P-,T-odd nuclear moments
  1. Schiff moment: typically dominant in diamagnetic systems
  2. MQM: relevant in paramagnetic systems
  3. EDM: typically shielded, but relevant for ions
- Nuclear theory relates nuclear moments to hadronic operators
  1. EDMs of neutron and proton $d_{n,p}$
  2. CP-violating pion-nucleon interactions $\bar{g}_{\pi NN}$
  3. Four-nucleon contact terms (C_{4N})
- QCD relates hadronic operators to quark-level operators

➥ Nuclear theory essential *e.g.* for world's best EDM limit (Hg)

Challenge: calculate $S, M, d_N(d_{n,p}, \bar{g}_{\pi NN}, C_{4N})$ for $A \sim 200$

Hg: sign of $\bar{g}^{(1)}_{\pi NN}$ unclear → no constraint

$S(d_{n,p})$: 1. just d_n 2. shell model → $S(d_{n,p})$ 3. can we do better?

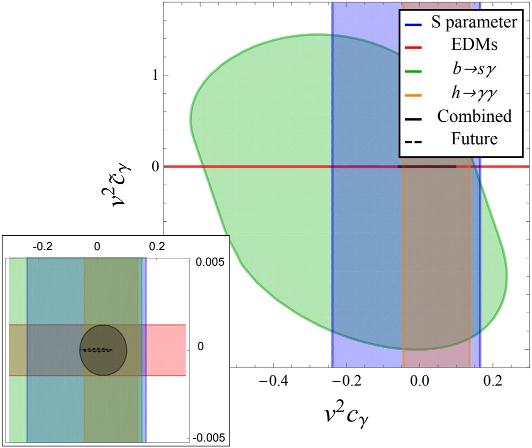
Unique chance: orders of magnitude *without a new experiment!*

Connecting high- and low-energy observables with EFTs

Example from [Cirigliano et al.'16] :

Consider chirality-flipping SMEFT operators with top and Higgs

➥ Affect EDMs, Higgs observables, flavour, ...



[Cirigliano+'16]

- CPV dominated by EDMs
- collider observables complementary

➥ significant progress expected for both

Turning the argument around

Other limits not relevant to global fit

➥ Use results to conservatively bound their EDMs

System	Indirect bound	Present/Expected limit
Cs	$[-3.1, 2.2]$	1400 [Murthy+'89] /1
Rb	$[-0.8, 0.5]$	10^8 [Ensberg+'67] /0.1
	unpublished:	(1200) [Huang-Hellinger'87]
Fr	$[-3.2, 4.2]$	—/1

Bounds on $|d_X|$ in $10^{-26}e$ cm

➥ **Several orders of magnitude below present limits!**

Experiments aiming at even better sensitivity:

➥ Important progress to be expected

➥ In case of a violation of the above limits:
Highly-tuneed cancellations or experimental problem