

An *ab-initio* multiphonon approach to low- and high- energy nuclear spectroscopy

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Low-lying Nuclear spectroscopy approach: The Nuclear Shell Model

Infinite Space, A nucleons

$$H\Psi_\alpha = E_\alpha \Psi_\alpha$$



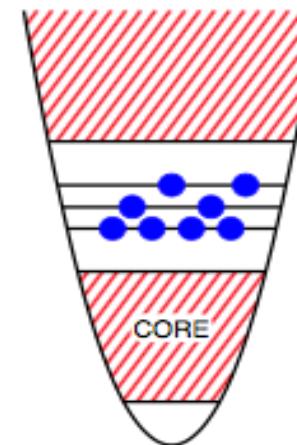
Model Space, v nucleons

$$H_{\text{eff}}\phi_\alpha = (T + V_{\text{eff}})\phi_\alpha = E_\alpha \phi_\alpha$$

The excitations of the core are absorbed by
the effective hamiltonian

$$V_{\text{eff}} = V + V \frac{Q}{E - H_0} V_{\text{eff}}$$

Problem :size of the Hamiltonian
matrix too large as v increases



SM describes only few low-
lying states.
It cannot describe most of
the collective states.

^{116}Sn , $N \sim 10^7$

^{130}Xe , $N \sim 10^9$

^{128}Xe , $N \sim 10^{10}$

Theoretical approaches to collective motions

Semiclassical

$$\alpha \quad \pi$$
$$[H, O_\lambda^\dagger] = \hbar\omega O_\lambda^\dagger$$

mapping

Harmonic appriximation

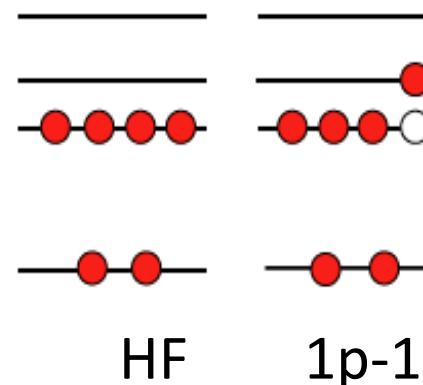
$$O_\lambda |0\rangle = 0$$

$$[H, O_\lambda^\dagger] |0\rangle = \omega_\lambda O_\lambda^\dagger |0\rangle = (E_\lambda - E_0) O_\lambda^\dagger |0\rangle$$

TDA

$$|\lambda\rangle = O_\lambda^\dagger |HF\rangle$$

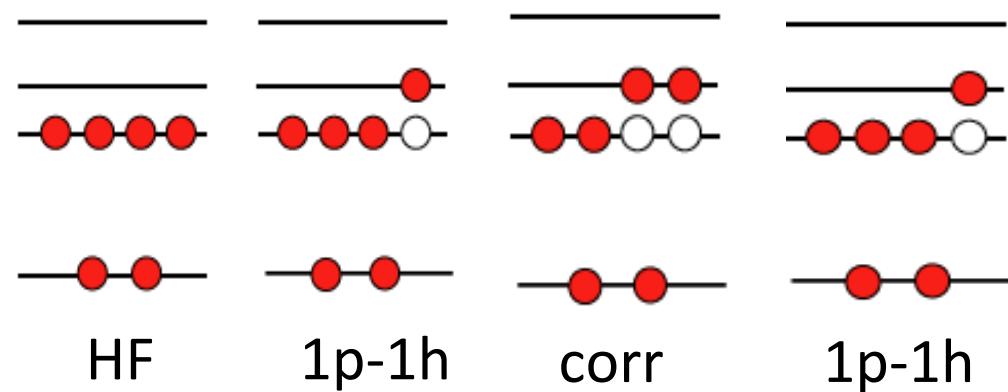
$$O_\lambda^\dagger = \sum_{ph} c_{ph}(\lambda) a_p^\dagger a_h$$



RPA

$$|\lambda\rangle = O_\lambda^\dagger |0_{corr}\rangle$$

$$O_\lambda^\dagger = \sum_{ph} [X_{ph}(\lambda) a_p^\dagger a_h - Y_{ph}(\lambda) a_h^\dagger a_p]$$



Beyond mean field: *Adopted Methods*

Non relativistic

- **qp-Phonon**

P. F. Bortignon et al. (Milano group)

- **2nd RPA**

R. Roth et al.

D. Gambacurta et al.

Phenomenological

- **QPM** (*Soloviev School (Dubna)*)

Relativistic: **RTBA**

E. Litvinova, P. Ring, D. Vretenar.....

Microscopic unified description: The Equation of Motion Phonon Method (EMPM)

Goal: Generate a multiphonon basis starting from TDA phonons

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Eigenvalue problem

$$H | \Psi_v \rangle = E_v | \Psi_v \rangle$$

$$| \Psi_v \rangle \in \mathcal{H} = \sum_n \bigoplus \mathcal{H}_n \quad \mathcal{H}_n \in | n; \beta \rangle \equiv n\text{-phonon basis states}$$

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An obvious, but unmanageable, multiphonon basis

$$| \lambda_1, \dots, \lambda_i, \dots \lambda_n \rangle = O_{\lambda_1}^\dagger \dots O_{\lambda_i}^\dagger \dots O_{\lambda_n}^\dagger | 0 \rangle$$

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A viable route

$$| \beta \rangle = \sum_{\lambda \alpha} C_{\lambda \alpha}^\beta O_{\lambda}^\dagger | \alpha \rangle$$

Construction of $|\beta\rangle$: EoM

Assuming $|\alpha\rangle$ known, we solve the **Eq. of Motion**

$$\left\langle \beta \left[H, O_\lambda^\dagger \right] \alpha \right\rangle = (E_\beta - E_\alpha) \left\langle \beta | O_\lambda^\dagger | \alpha \right\rangle$$



χ

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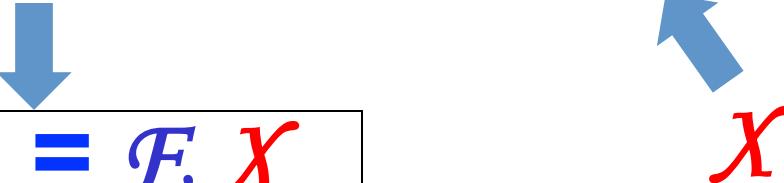


$$\mathcal{A} \chi = E \chi$$

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Where

$$\mathcal{A}_{(\lambda\alpha)(\lambda'\gamma)} = (E_\lambda + E_\alpha) \delta_{\alpha\gamma} \delta_{\lambda\lambda'} + \mathcal{V}_{(\lambda\alpha)(\lambda'\gamma)}$$

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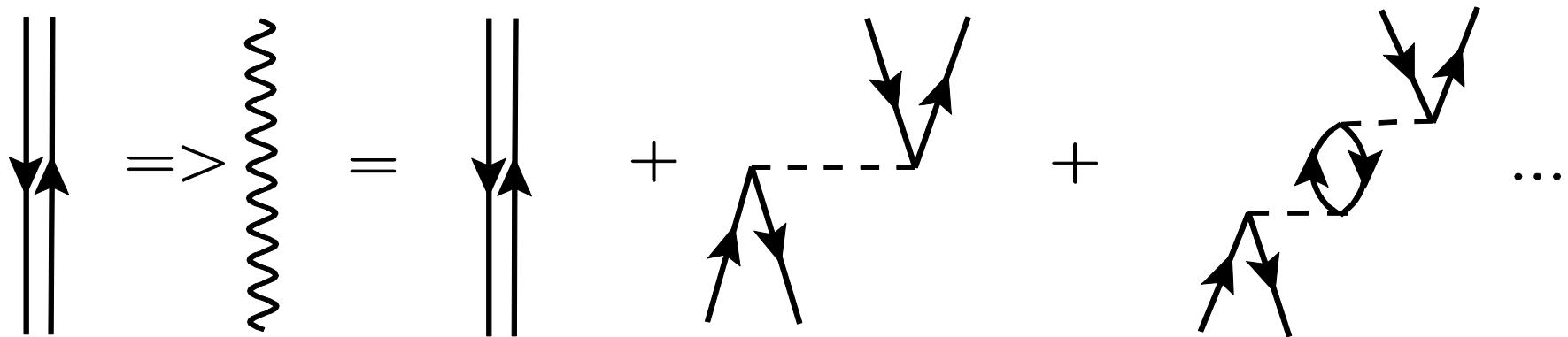
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TDA matrix

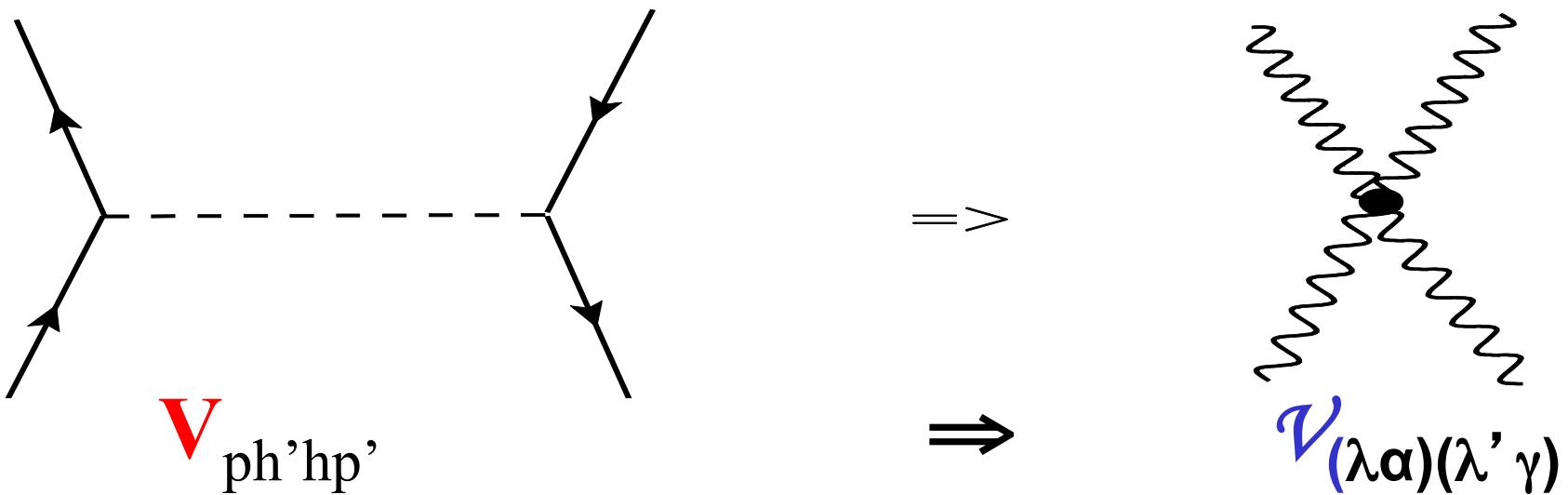
$$A_{(ph)(p'h')} = (\epsilon_p - \epsilon_h) \delta_{pp'} \delta_{hh'} + V_{ph'h'}$$

From p-h to TDA



$$\epsilon_p - \epsilon_h \Rightarrow E_\beta \\ |p\ h\rangle \Rightarrow |\beta\rangle = \sum_{ph} c_{ph}^\beta a_p^\dagger a_h |>$$

From TDA to EMPM



Construction of $|\beta\rangle$: EoM

Problem

$$\mathcal{A} \mathcal{X} = \mathcal{E} \mathcal{X}$$

is **not** a true Eigenvalue Eq.!

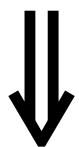
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is **not** a true Eigenvalue Eq.!

$\{\mathbf{O}_\lambda^\dagger |\alpha\rangle\}$ form a **non-orthogonal redundant** basis



$\mathcal{X} = \langle \beta | \mathbf{O}_\lambda^\dagger | \alpha \rangle$ is **not** a **true** expansion coefficient

Recipe for solving the problem

1° step

$$\mathcal{A} X = E X$$

$$X = \langle \beta | O_{\lambda}^{\dagger} | \alpha \rangle = \langle \beta | \lambda \alpha \rangle$$

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Expressing X in terms of C

$$X = D C$$

Where D is the metric matrix

$$D \equiv \{ \langle \lambda' \alpha' | \lambda \alpha \rangle \}$$

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$$[\mathcal{H} - E D] C = 0$$

where

$$\mathcal{H} = \mathcal{A} D$$

Recipe for solving the problem

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$$[\mathcal{H} - E D] C = 0$$

where

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But D is singular !

2° conclusive step: Choleski

$$[\mathcal{H} - \mathcal{E}\mathcal{D}] \mathbf{C} = \mathbf{0}$$

$$\mathcal{H} = \mathcal{A}\mathcal{D}$$

$$\mathcal{D} \equiv \{ < \lambda' a' \mid \lambda a > \}$$

2° conclusive step: Choleski

$$[\mathcal{H} - \mathcal{E}\mathcal{D}] \mathbf{C} = \mathbf{0}$$

$$\mathcal{H} = \mathcal{A}\mathcal{D}$$

Cholesky

$$\mathcal{D} \rightarrow \mathbf{D}$$

$$\mathcal{D} \equiv \{ < \lambda' \mathbf{a}' \mid \lambda \mathbf{a} > \}$$

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$$\mathcal{D} \rightarrow \mathbf{D}$$

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$$[\mathbf{H}^n - \mathbf{E}^n] \mathbf{C}^n = \mathbf{0}$$

$$\mathbf{H} = \mathbf{D}^{-1} \mathcal{A}\mathcal{D}$$

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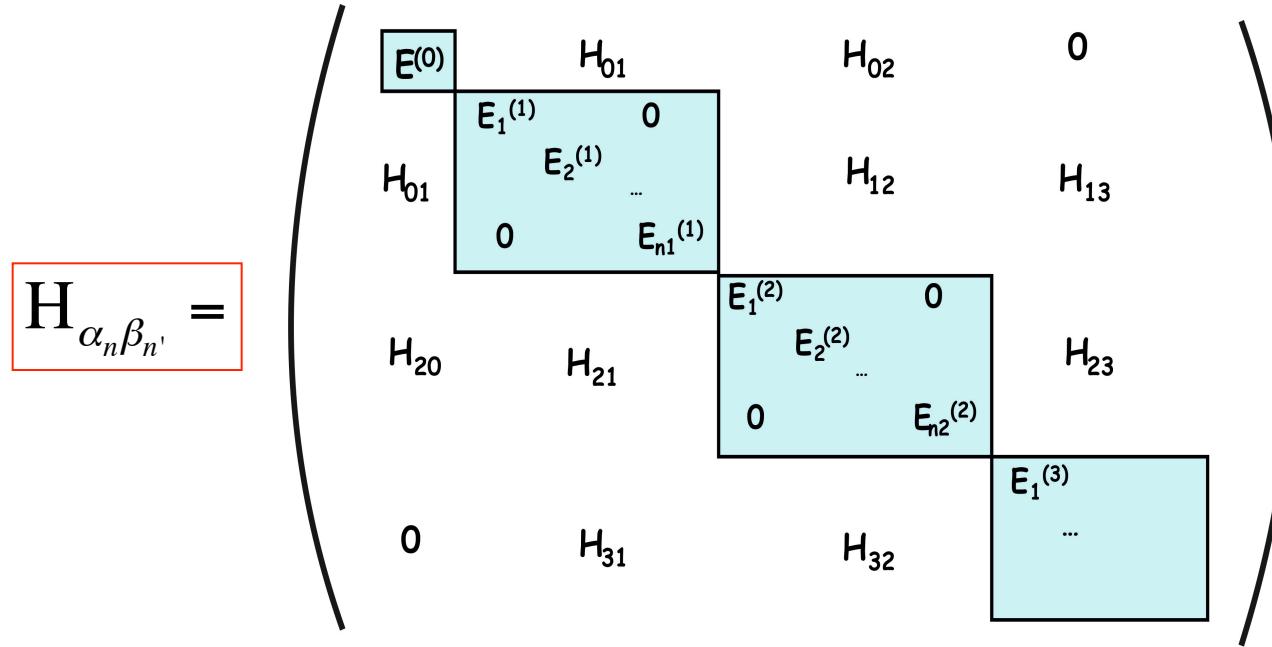


$$|\beta\rangle = \sum_{\lambda\alpha} C_{\lambda\alpha}^\beta \left| (\lambda \times \alpha)^\beta \right\rangle = \sum_{\lambda\alpha} C_{\lambda\alpha}^\beta \left(O_\lambda^\dagger |\alpha\rangle \right)^\beta$$

$$\langle \alpha_n | \beta_n, \rangle = \delta_{nn}, \quad \delta_{\alpha\beta}$$

Eigenvalue problem in the Multiphonon basis $\{|\alpha_n\rangle\}$

$$\sum_{n' \beta_{n'}} H_{\alpha_n \beta_{n'}} C_{\beta_{n'}}^\nu = \left[(E_{\alpha_n} - \varepsilon_\nu) \delta_{nn'} \delta_{\alpha_n \beta_{n'}} + V_{\alpha_n \beta_{n'}} \right] C_{\beta_{n'}}^\nu = 0$$



Final solution

$$|\Psi_\nu\rangle = \sum_{\alpha_n} C_{\alpha_n}^\nu |\alpha_n\rangle$$

$$|\alpha_n\rangle = \sum_{\lambda \alpha_{n-1}} C_{\lambda \alpha_{n-1}}^{\alpha_n} |(\lambda \times \alpha_{n-1})^{\alpha_n}\rangle$$

Implementation for even nuclei

Hamiltonian

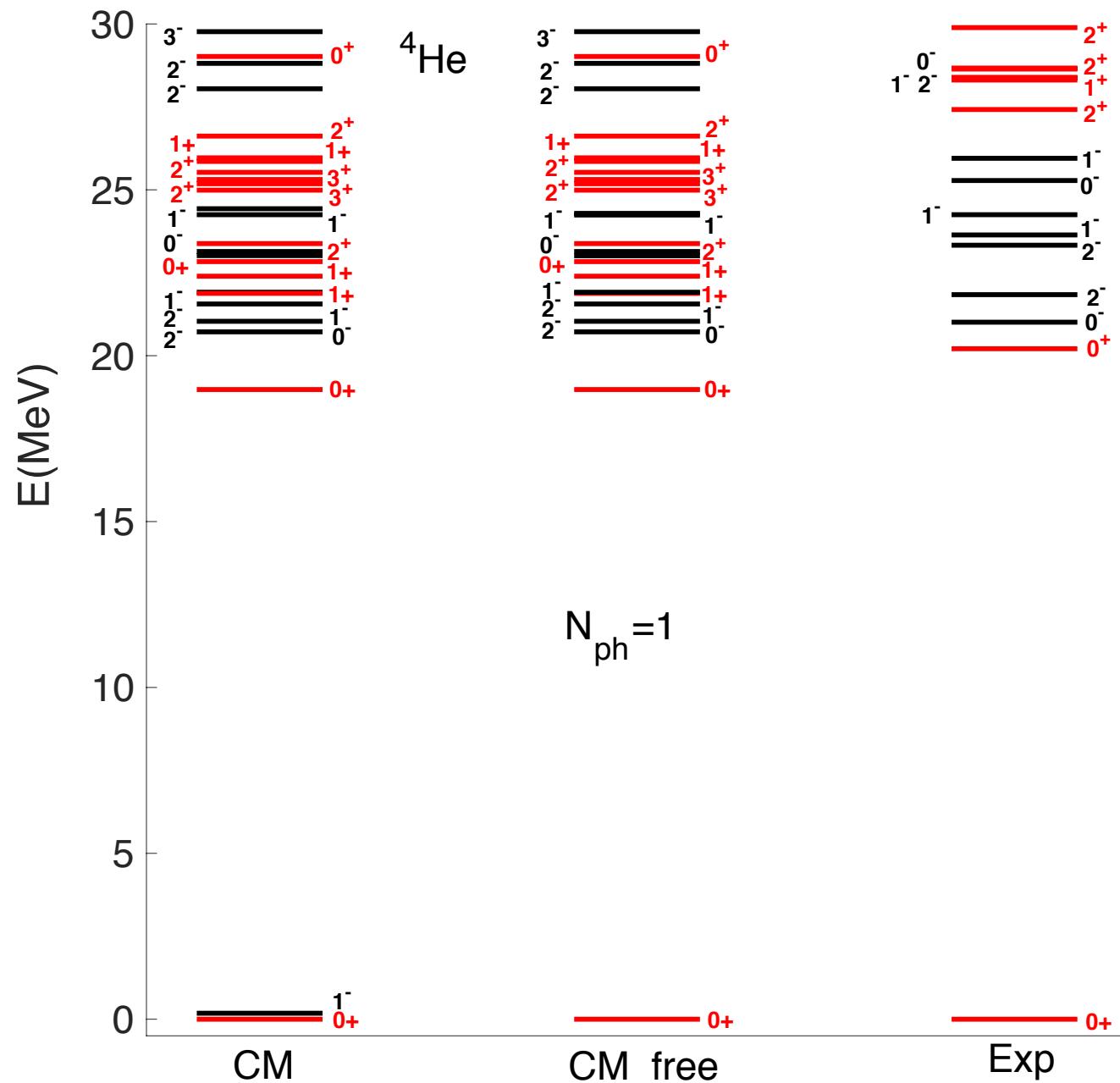
$$H = T_{\text{int}} + V$$

$$V = \text{NNLO}_{\text{sat}}$$

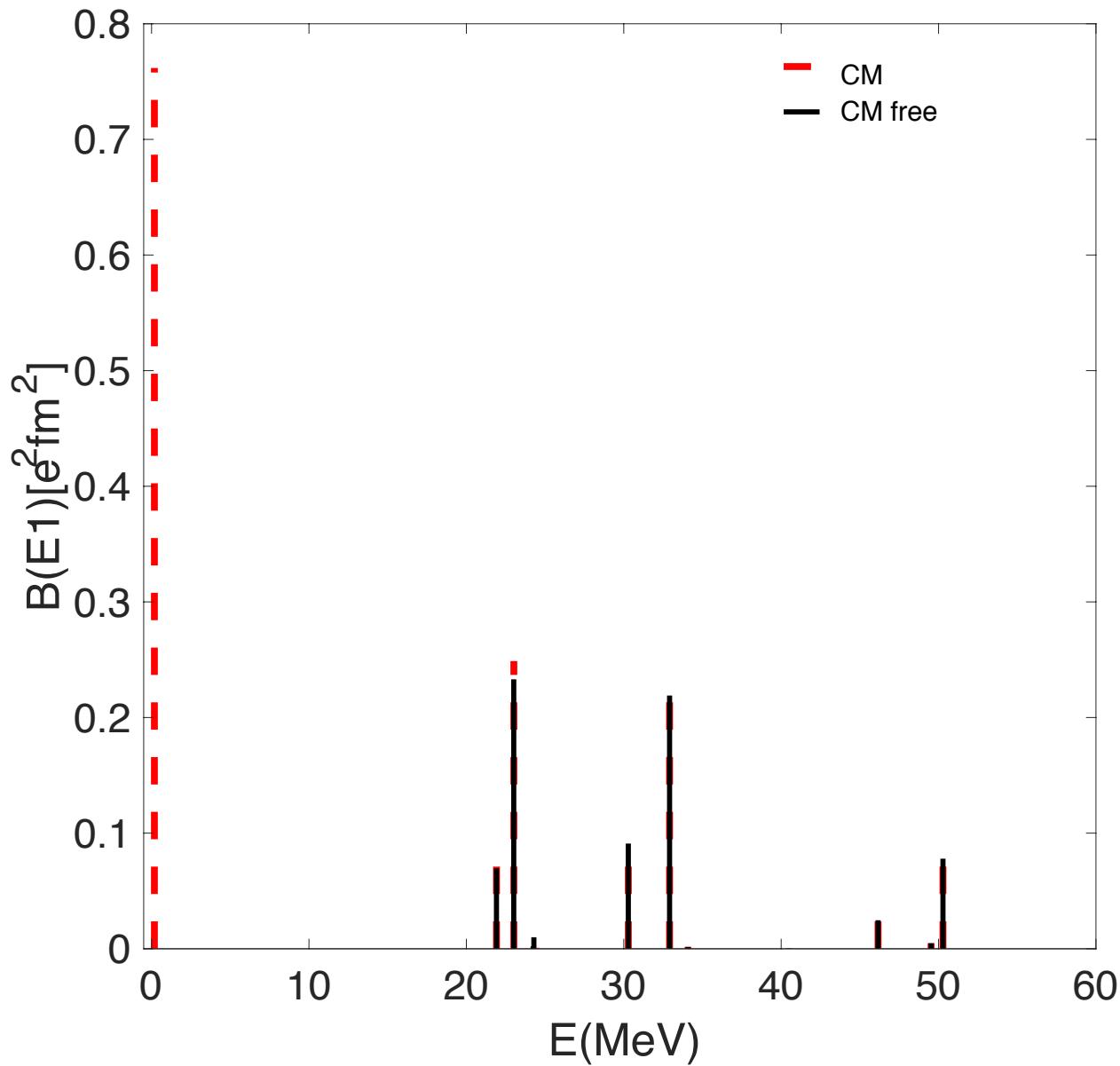
A. Ekstrom *et al.*, PRL 110, 192502 (2015)

- Perform **HF**
- Construct **TDA** phonons (**free** of **CM** spurious admixtures)
- Generate the multiphonon core basis $\{|\alpha_n\rangle\}$ (**free** of **CM** spurious admixtures)
- Full diagonalization

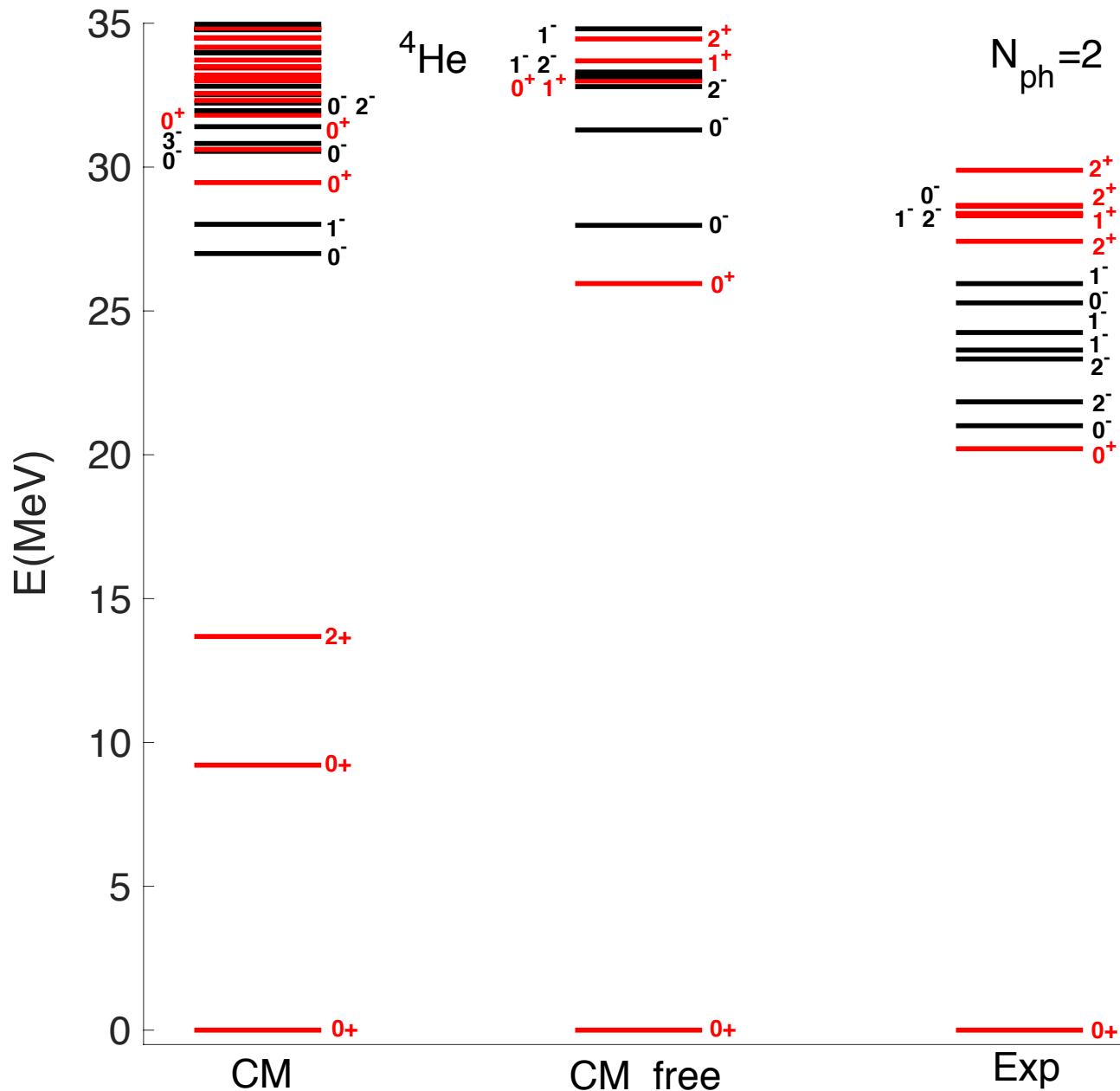
CM problem: Application to ${}^4\text{He}$



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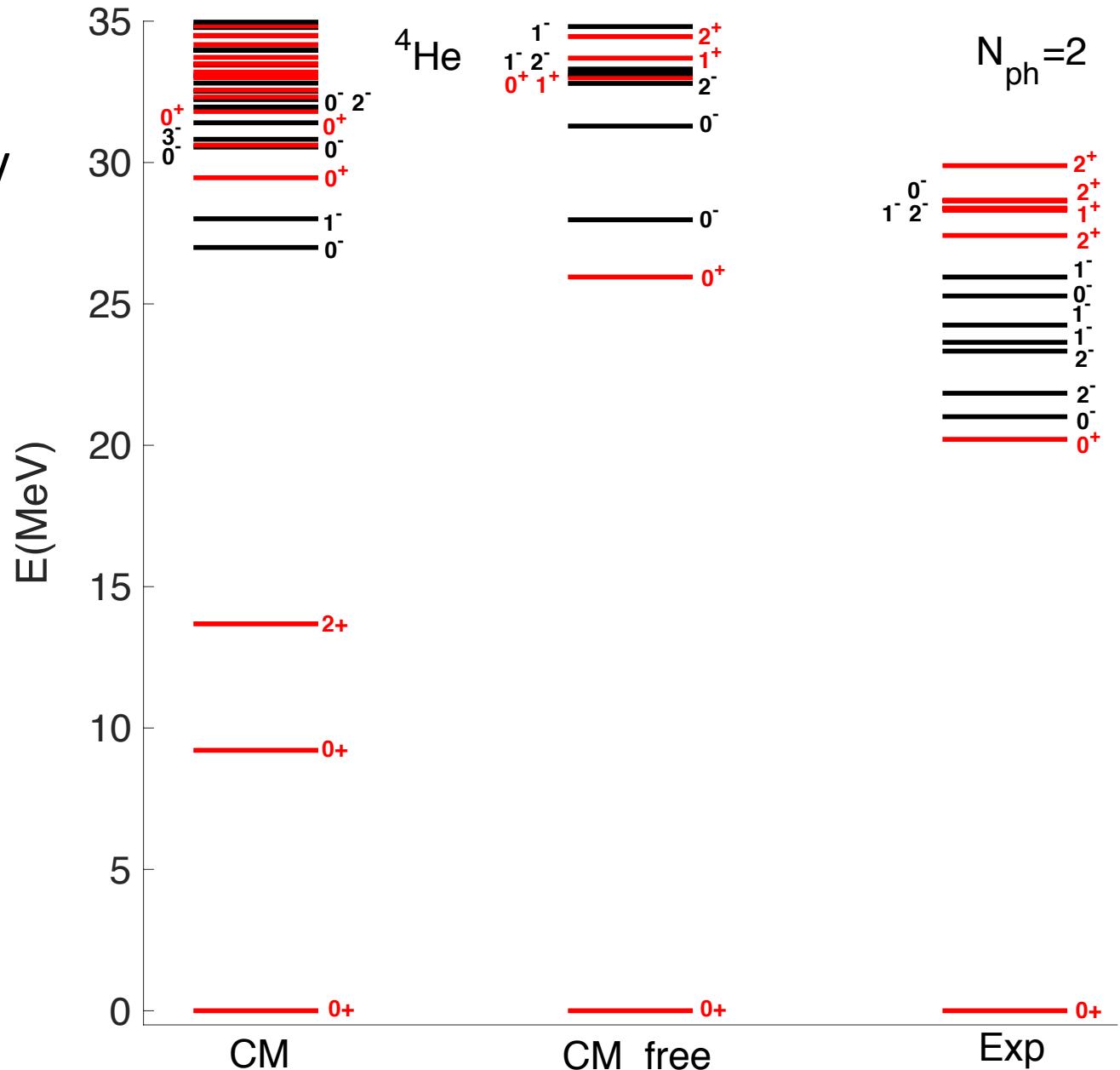
CM problem: Application to ${}^4\text{He}$

$$E_{HF} = -14.3712 \text{ MeV}$$

$$E_{corr} = -12.4943 \text{ MeV}$$

$$B_{th}/A = 6.7164 \text{ MeV}$$

$$B_{exp}/A = 7.073 \text{ MeV}$$



Numerical implementation

Hamiltonian

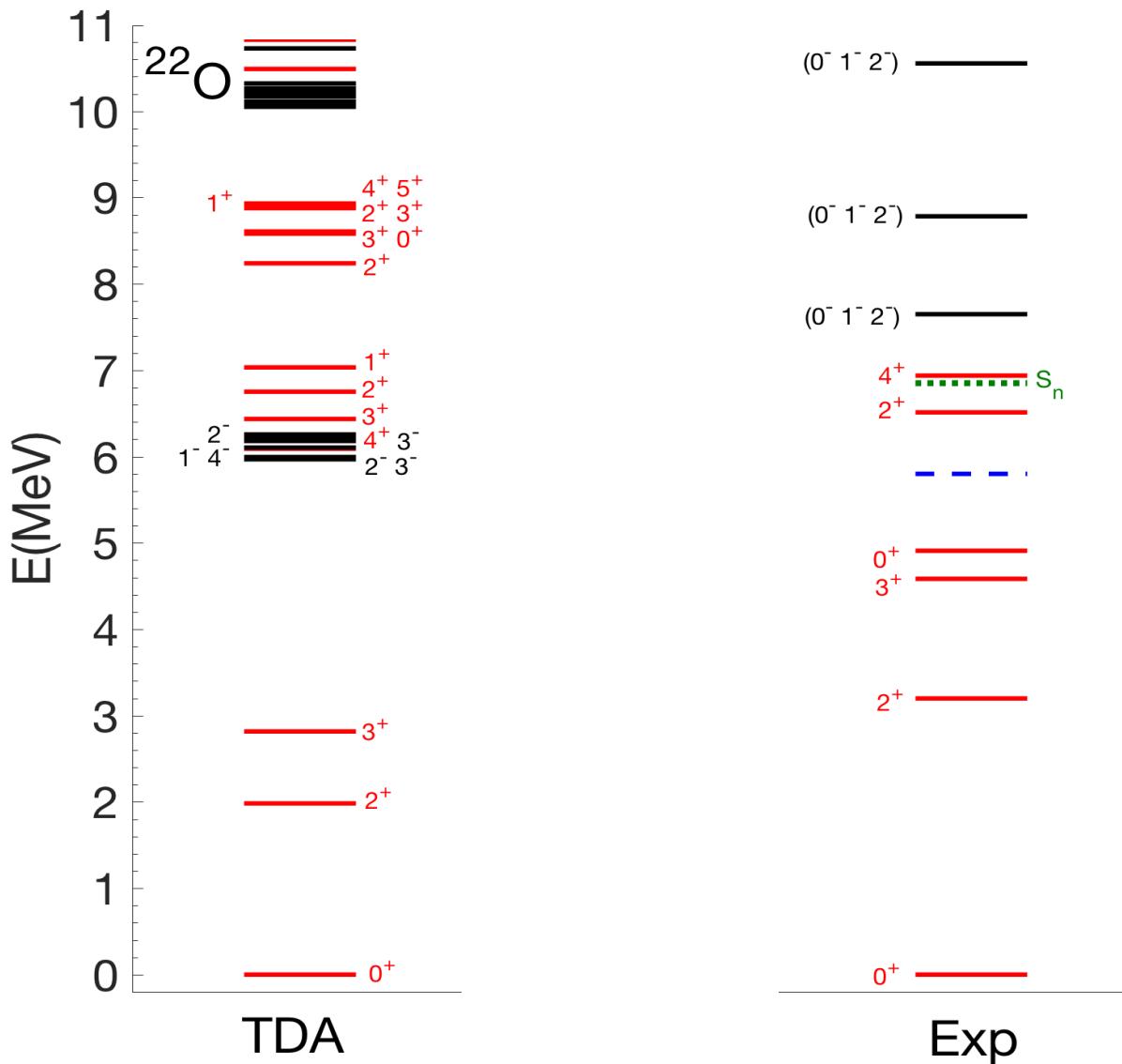
$$H = T_{\text{int}} + V_2$$

$$V_2 = \text{NNLO}_{\text{opt}}$$

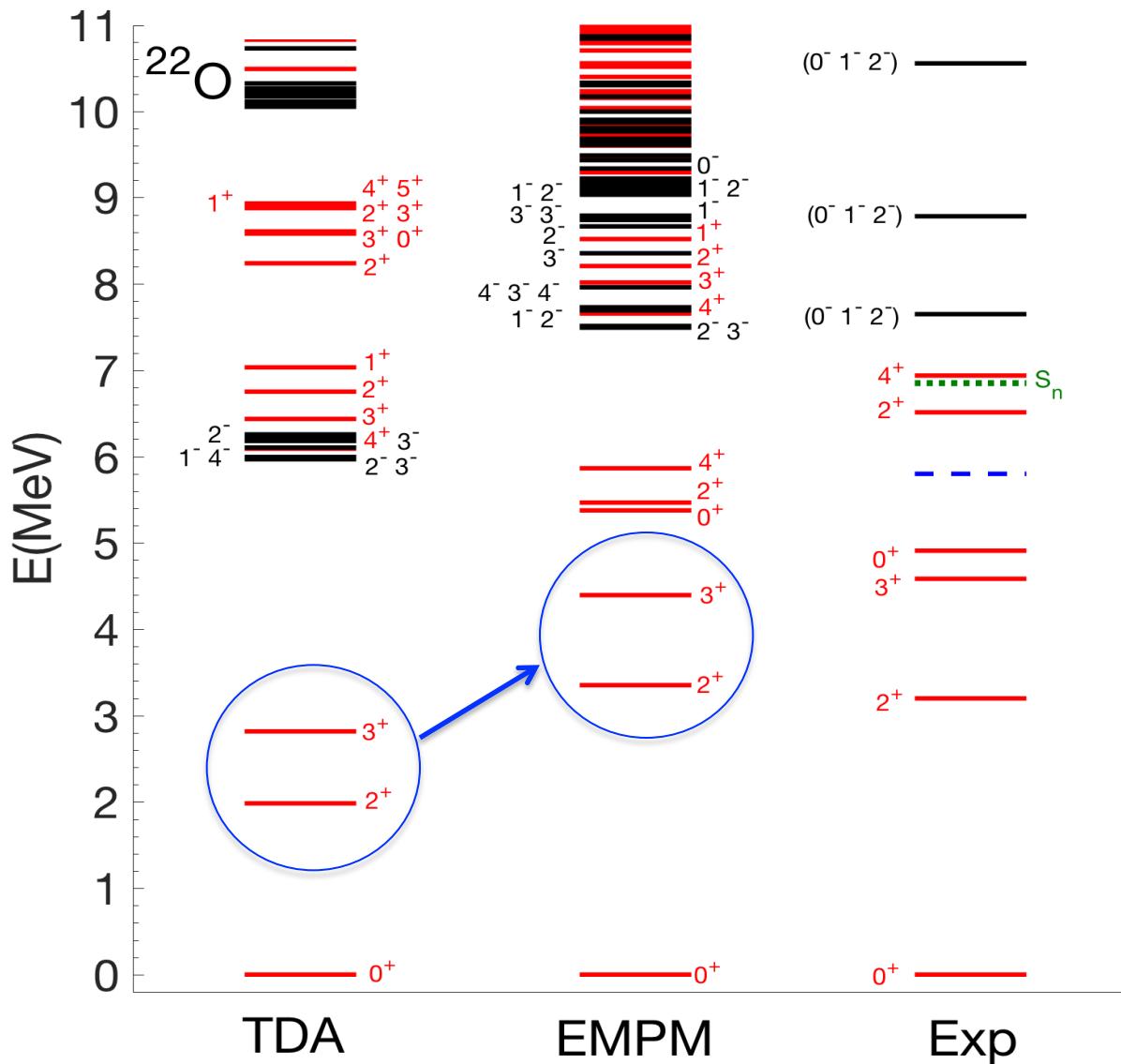
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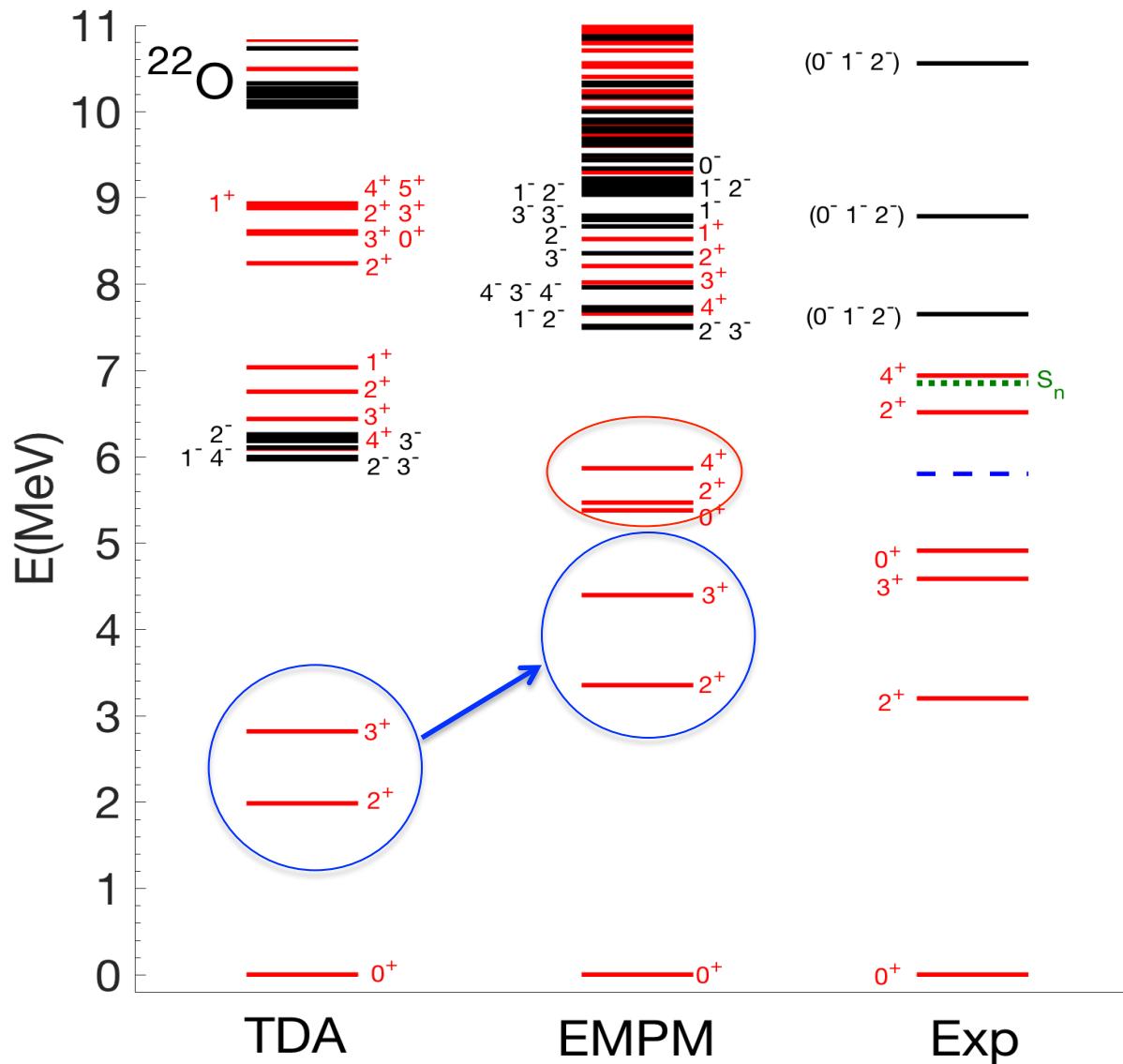
Application to the neutron rich ^{22}O



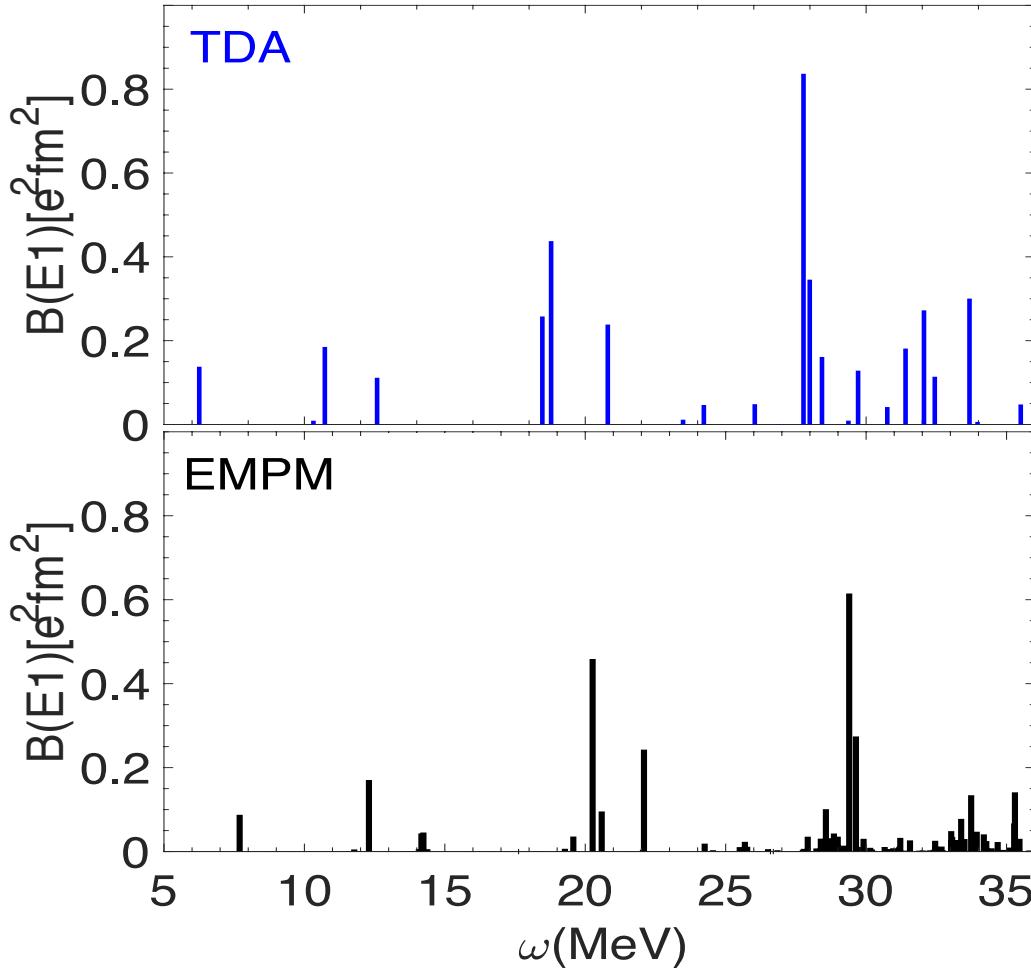
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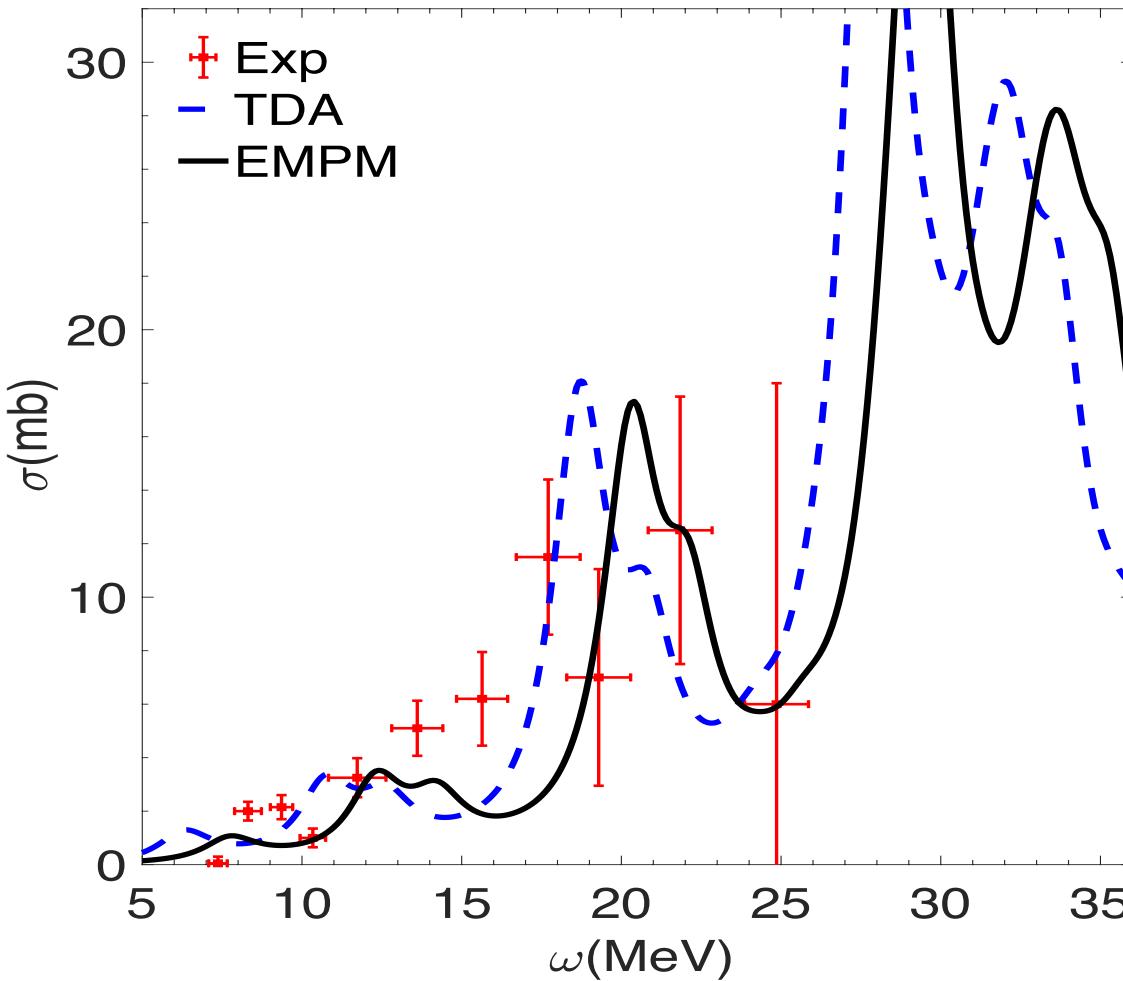


Dipole response in ^{22}O



$$B_\nu[E1](0_1^+ \rightarrow 1_\nu^-) = \left| \langle \Psi_{1^-}^\nu | M(E1) | \Psi_{0^+} \rangle \right|^2$$

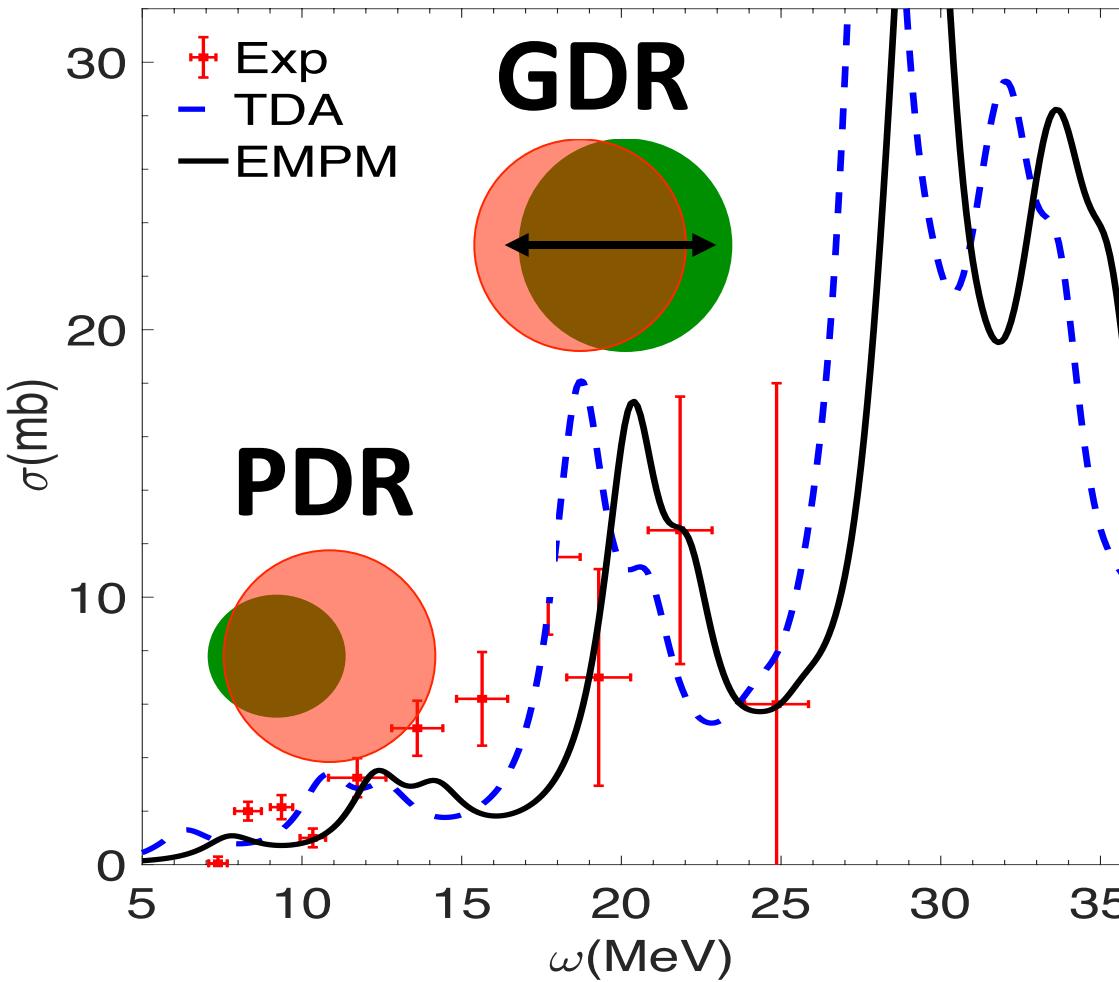
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$$S(E\lambda, \omega) = \sum_{\nu} B_{\nu}(E\lambda) \delta(\omega - \omega_{\nu}) \approx \sum_{\nu} B_{\nu}(E\lambda) \rho_{\Delta}(\omega - \omega_{\nu})$$

$$\sigma = \int_0^{\infty} \sigma(\omega) d\omega = \frac{16\pi^3}{9\hbar c} \int_0^{\infty} \omega S(E1, \omega) d\omega$$

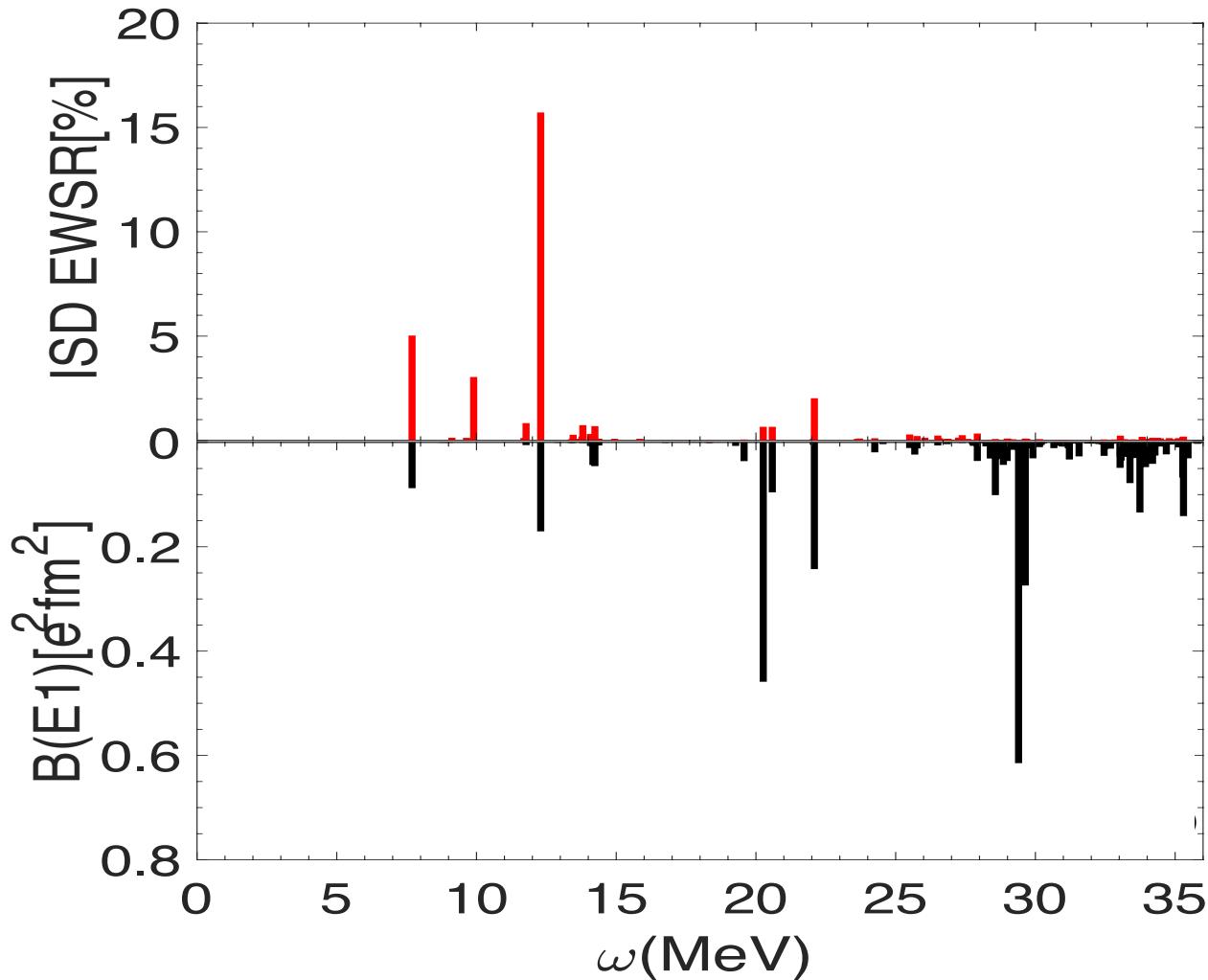
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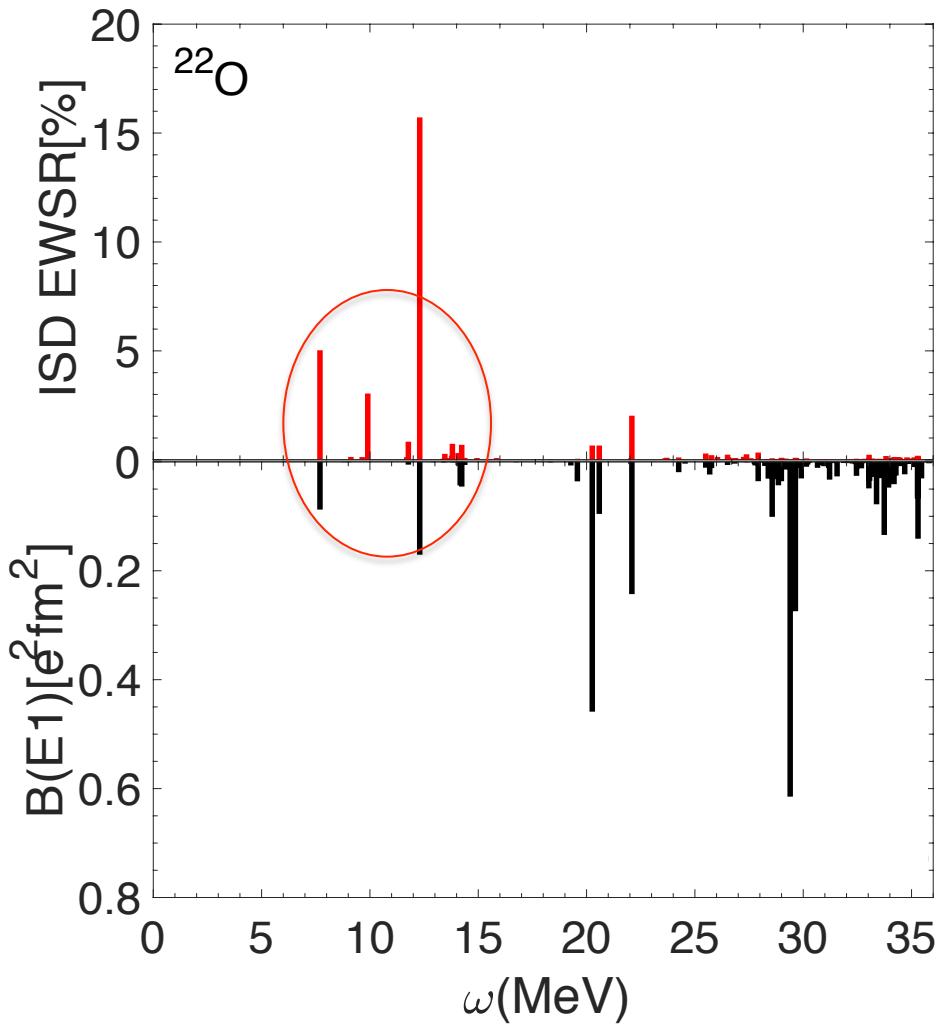
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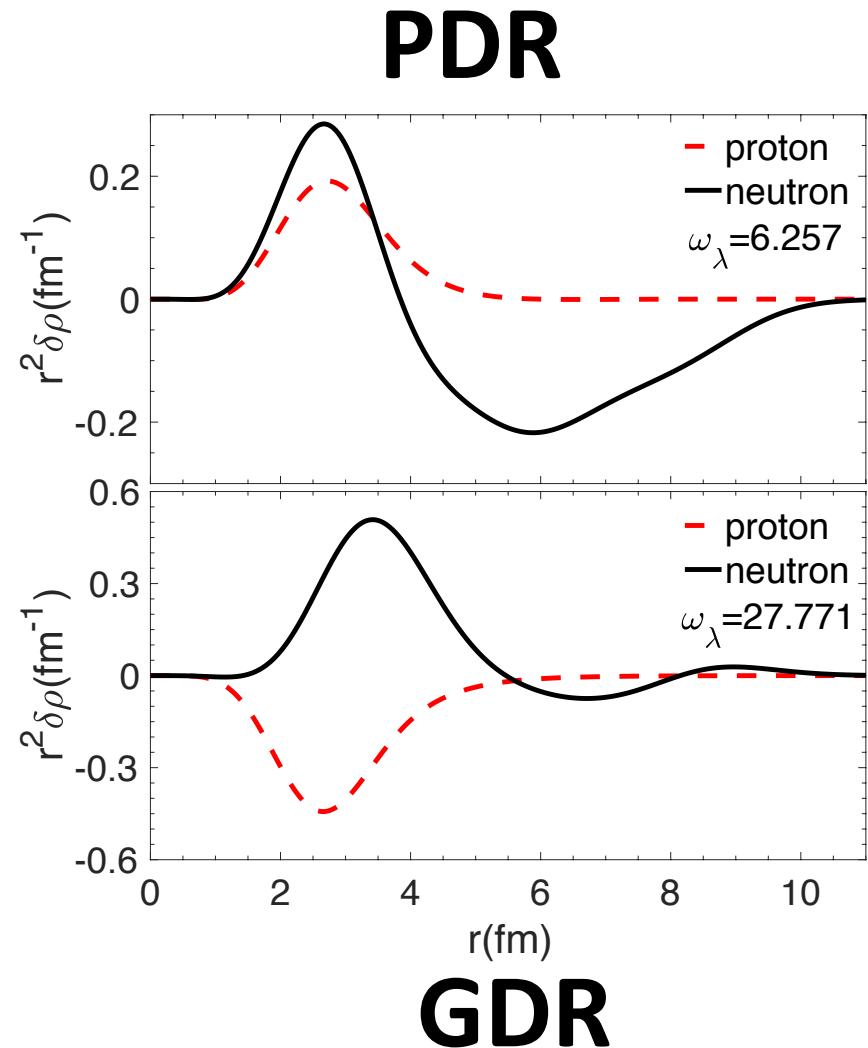
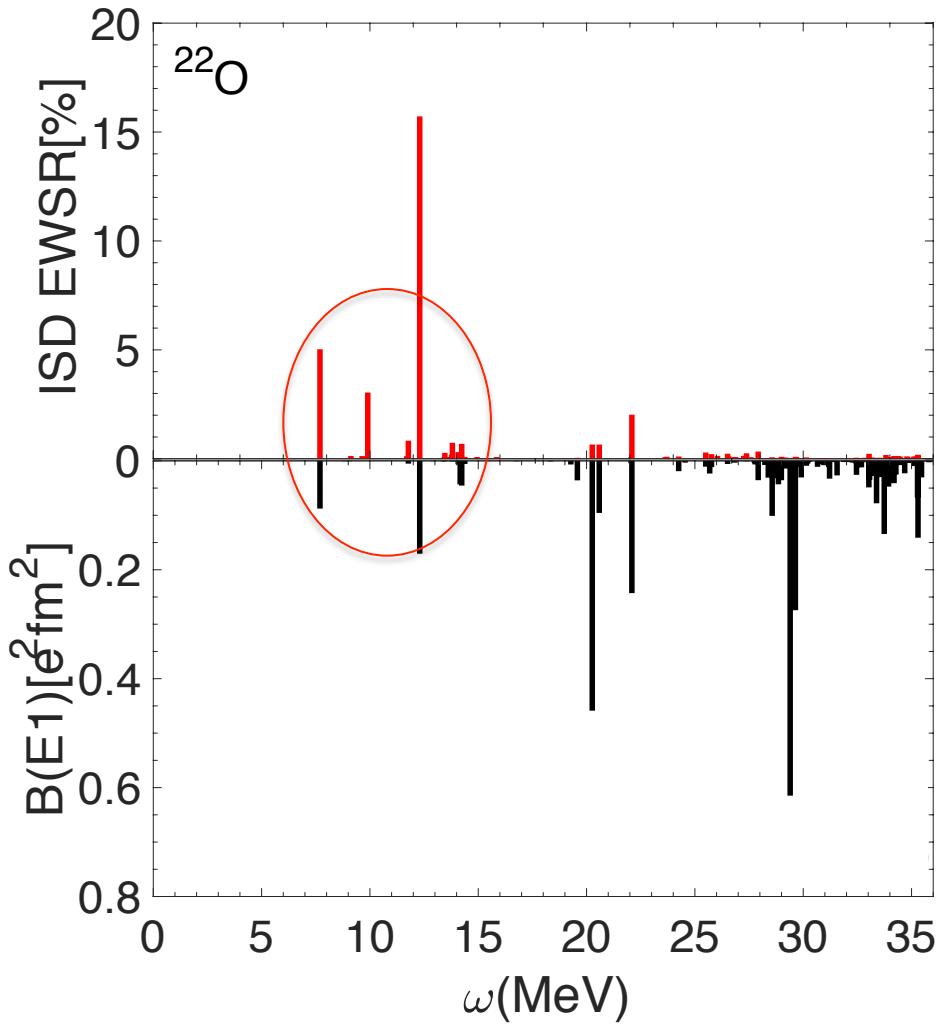
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EMPM for Odd nuclei: p (h) multiphonon scheme

$$|\lambda \alpha_{n-1}\rangle = O^\dagger_\lambda |\alpha_{n-1}\rangle \quad \longrightarrow \quad |\text{p } \alpha_n\rangle = a^\dagger_p |\alpha_n\rangle$$

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We solve the **Eq. of Motion**

$$\langle v_n | [H, a^\dagger_p] | \alpha_n \rangle = (E_{v_n} - E_{\alpha_n}) \langle v_n | a^\dagger_p | \alpha_n \rangle$$



χ

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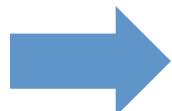
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Expressing x in terms of c

$$x = D c$$

x

$$\mathcal{A} x = E x$$



$$[\mathcal{H} - E D] C = 0$$

$$D \equiv \{ \langle p \alpha | p' \alpha' \rangle \}$$

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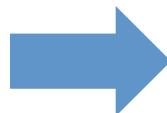
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Expressing x in terms of c

$$X = DC$$

$$X$$

$$\mathcal{A} X = E X$$



$$[\mathcal{H} - E\mathcal{D}] C = 0$$

$$\mathcal{D} \equiv \{ \langle p \alpha | p' \alpha' \rangle \}$$

Cholesky

$$\mathcal{D} \rightarrow D$$

EMPM for Odd nuclei: p (h) multiphonon scheme

$$[\mathbf{H} - \mathbf{E}] \mathbf{C} = \mathbf{0} \quad \mathbf{H} = \mathbf{D}^{-1} \mathcal{A} \mathcal{D}$$

EMPM for Odd nuclei: p (h) multiphonon scheme

$$[H - E] C = 0$$

$$H = D^{-1} \mathcal{A} \mathcal{D}$$



$$|\nu_n\rangle = \sum_{p\alpha} C_{p\alpha}^{\nu} a_p^\dagger |\alpha\rangle$$

$$\langle \nu'_{n'} | \nu_n \rangle = \delta_{\nu\nu'} \delta_{nn'}$$

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$$|\Psi_\Omega\rangle = \sum_{n\nu_n} C_{\nu_n}^\Omega |\nu_n\rangle$$

EMPM for Odd nuclei: p (h) multiphonon scheme

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$$\langle \nu'_{n'} | \nu_n \rangle = \delta_{\nu\nu'} \delta_{n'n}$$

$$|\Psi_\Omega\rangle = \sum_{n\nu_n} C_{\nu_n}^\Omega |\nu_n\rangle$$

- No approximations except for truncation !
- Pauli principle fully accounted for
- No redundant states!
- $|\nu_n\rangle$ form an orthonormal basis

Implementation for odd nuclei

Hamiltonian

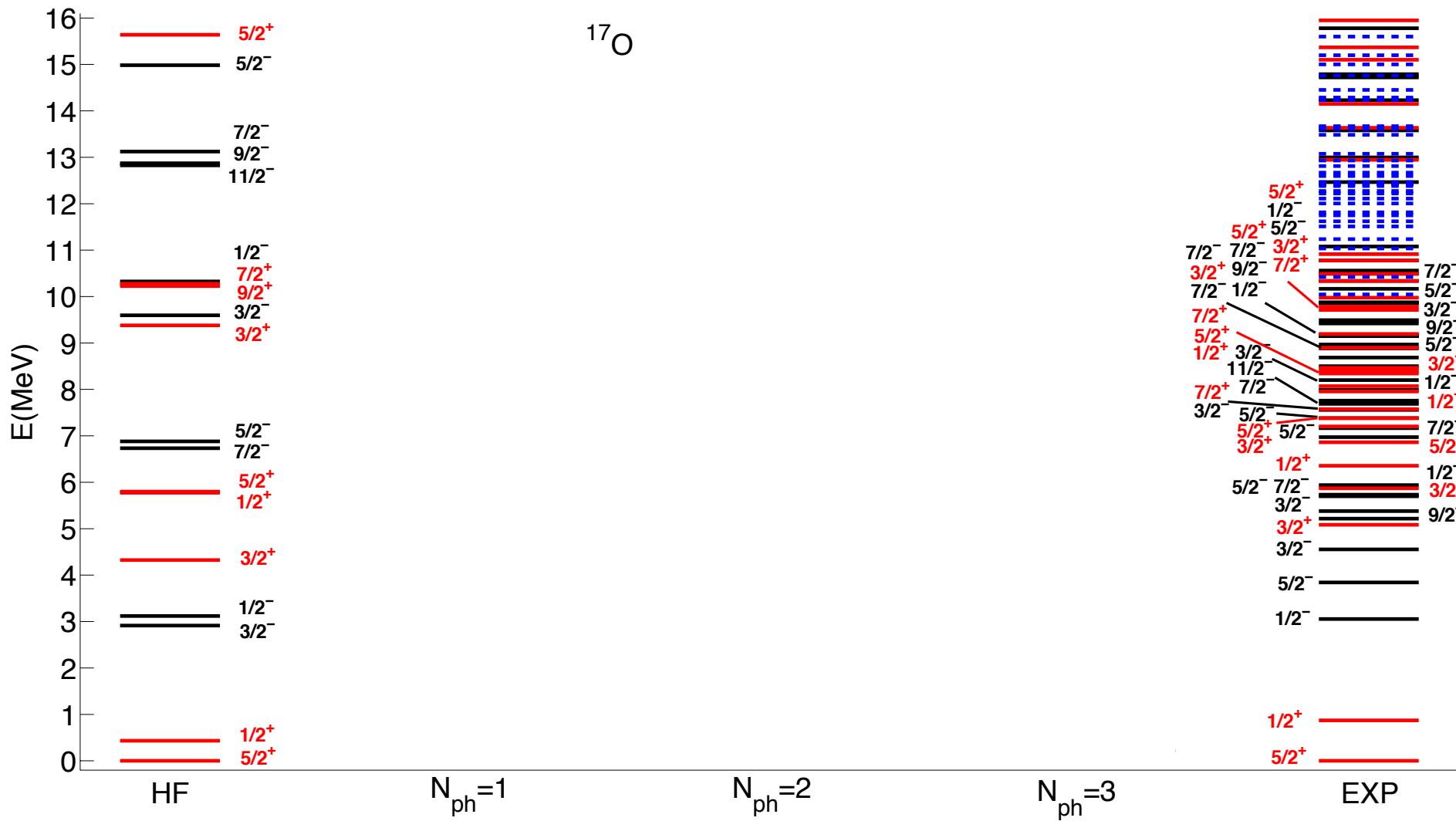
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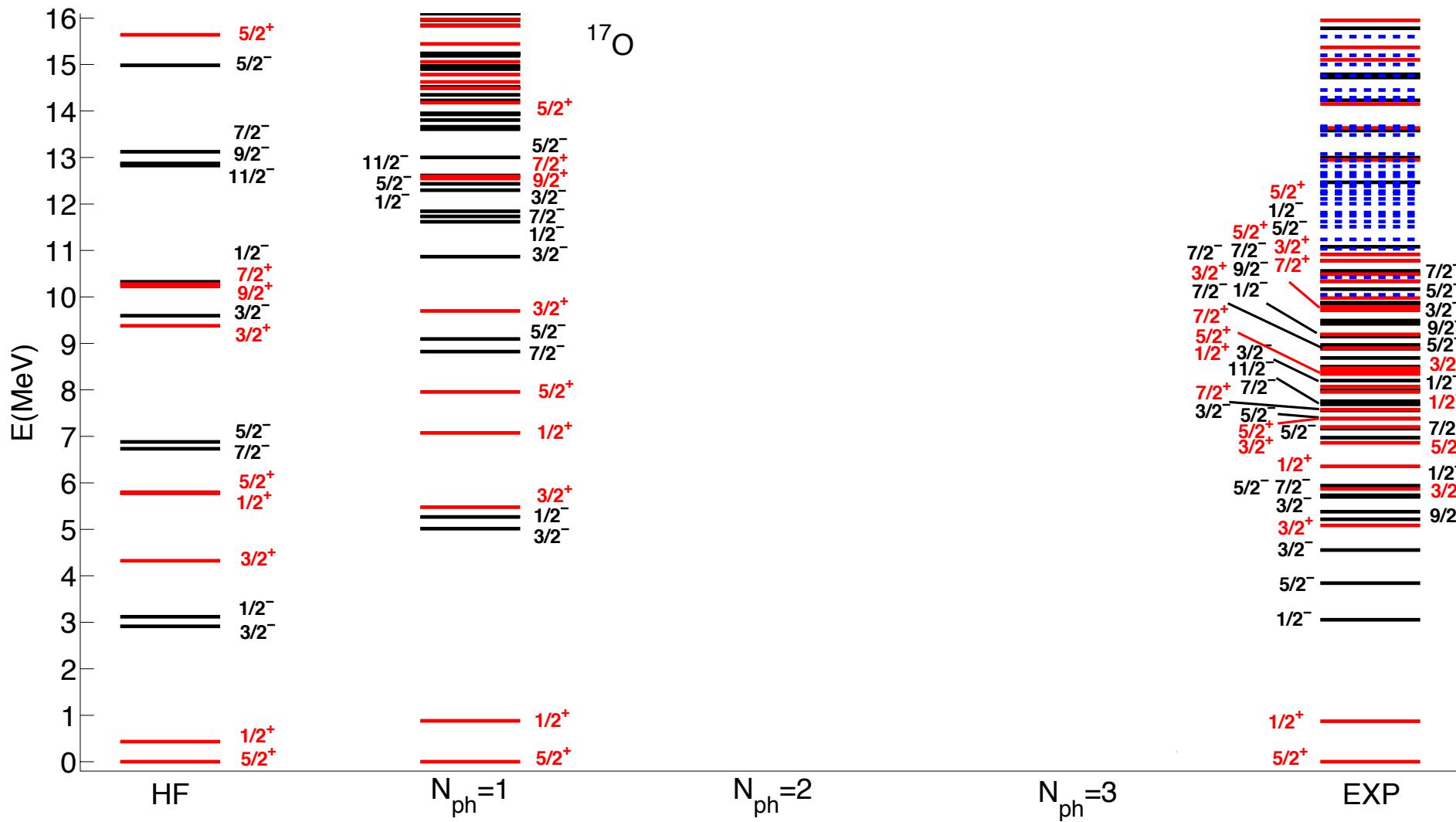
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- Generate the particle-phonon basis $\{|\nu_n\rangle\}$
- Full diagonalization

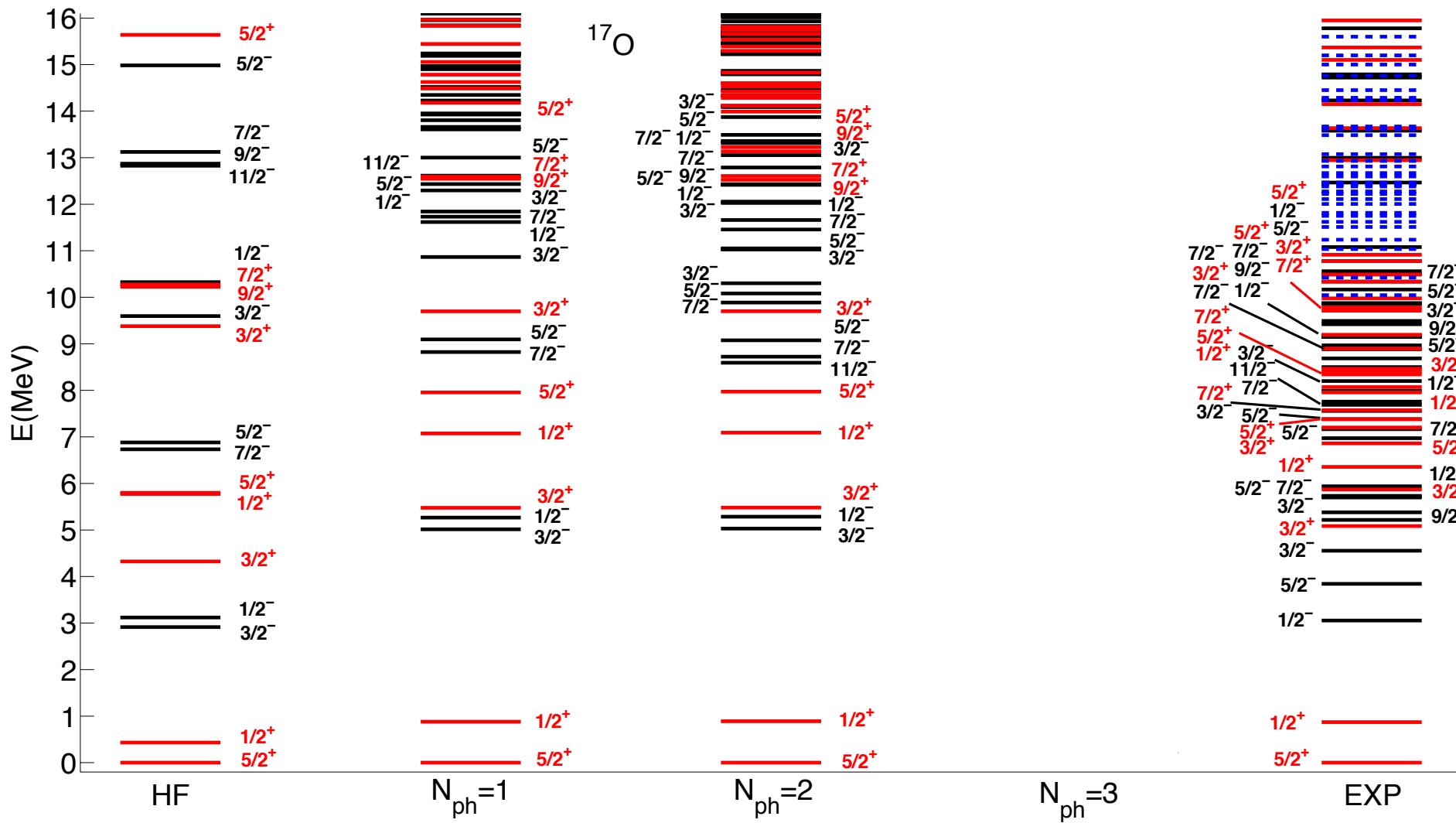
Application: ^{17}O spectra



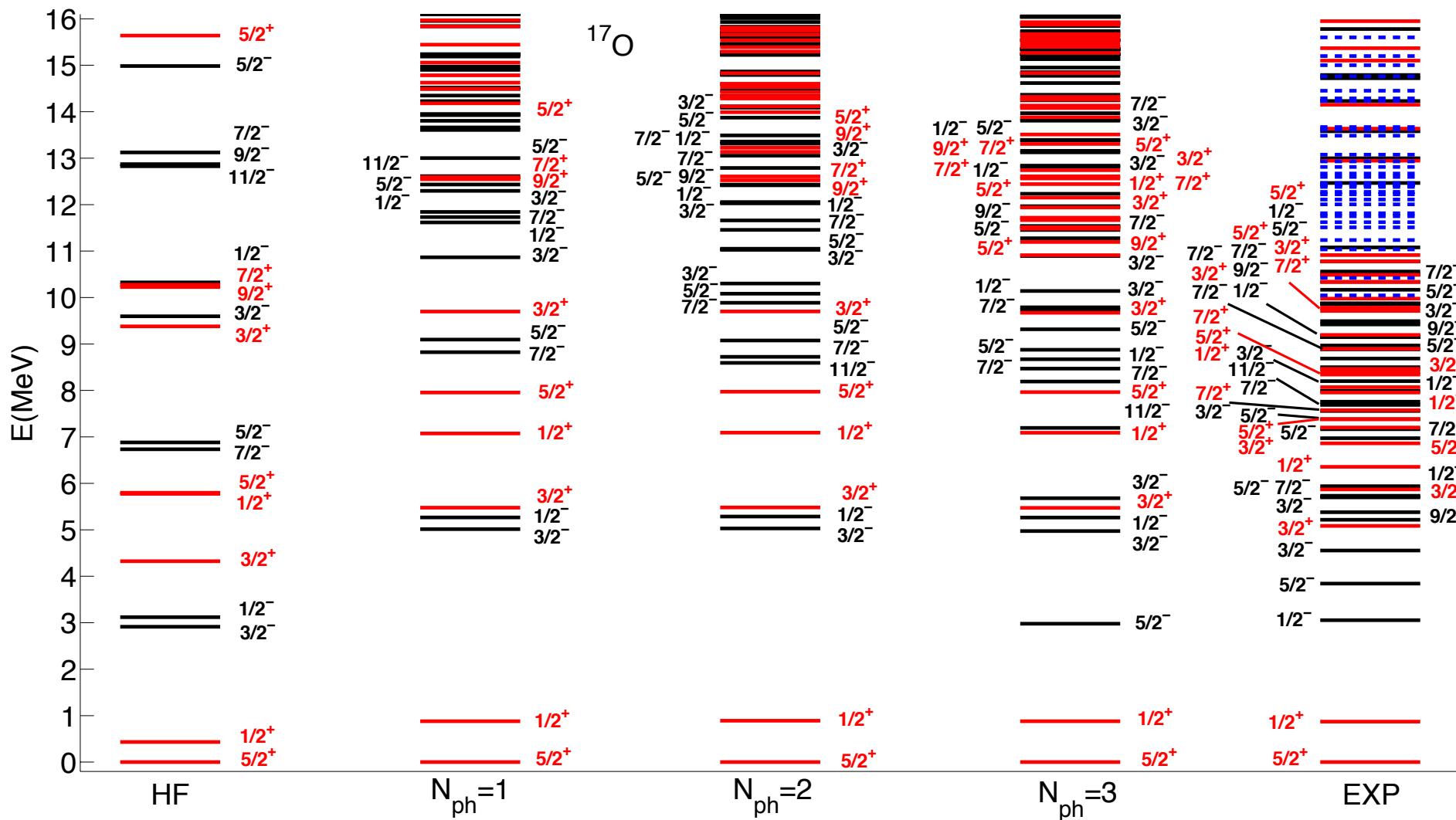
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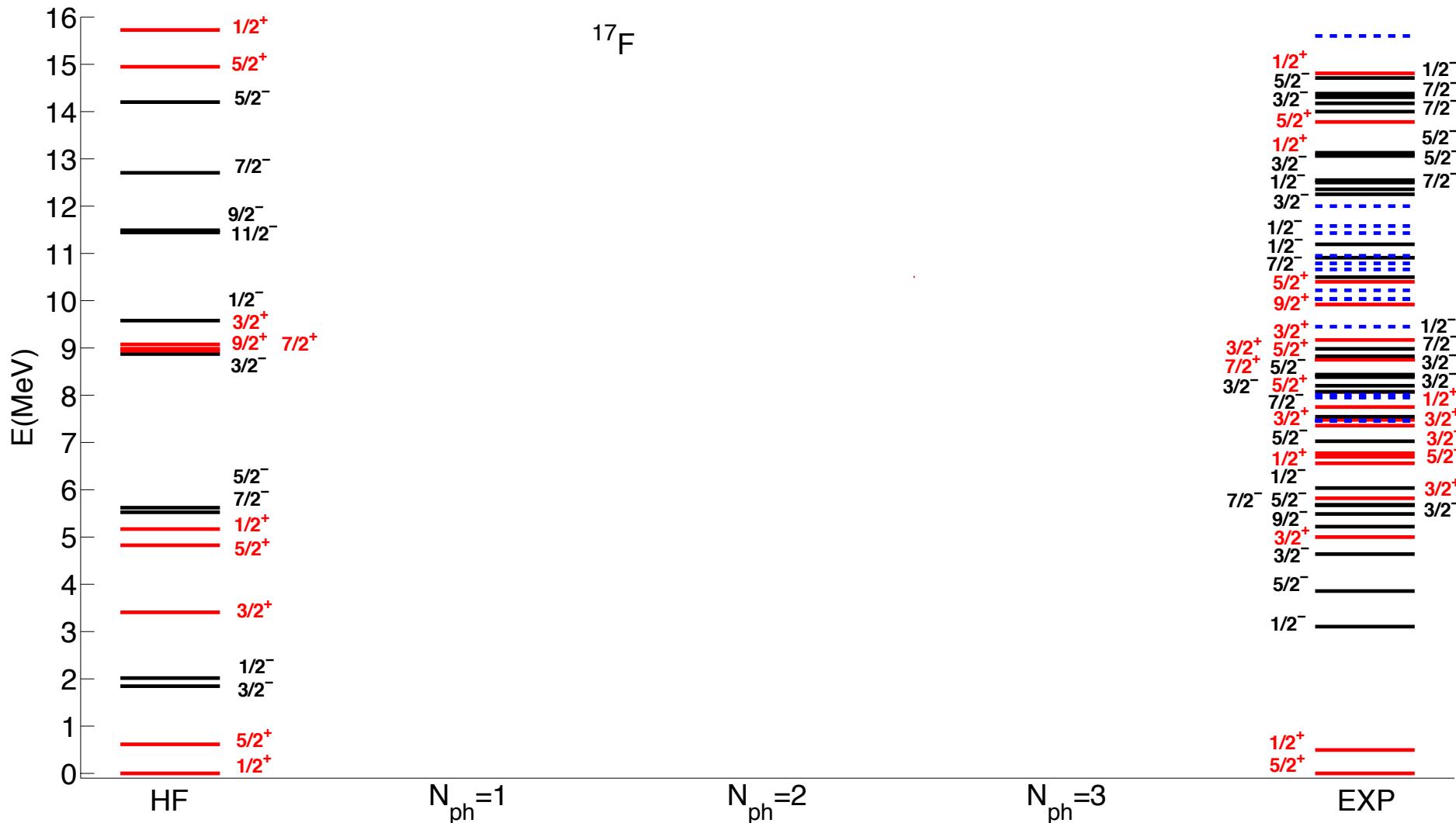
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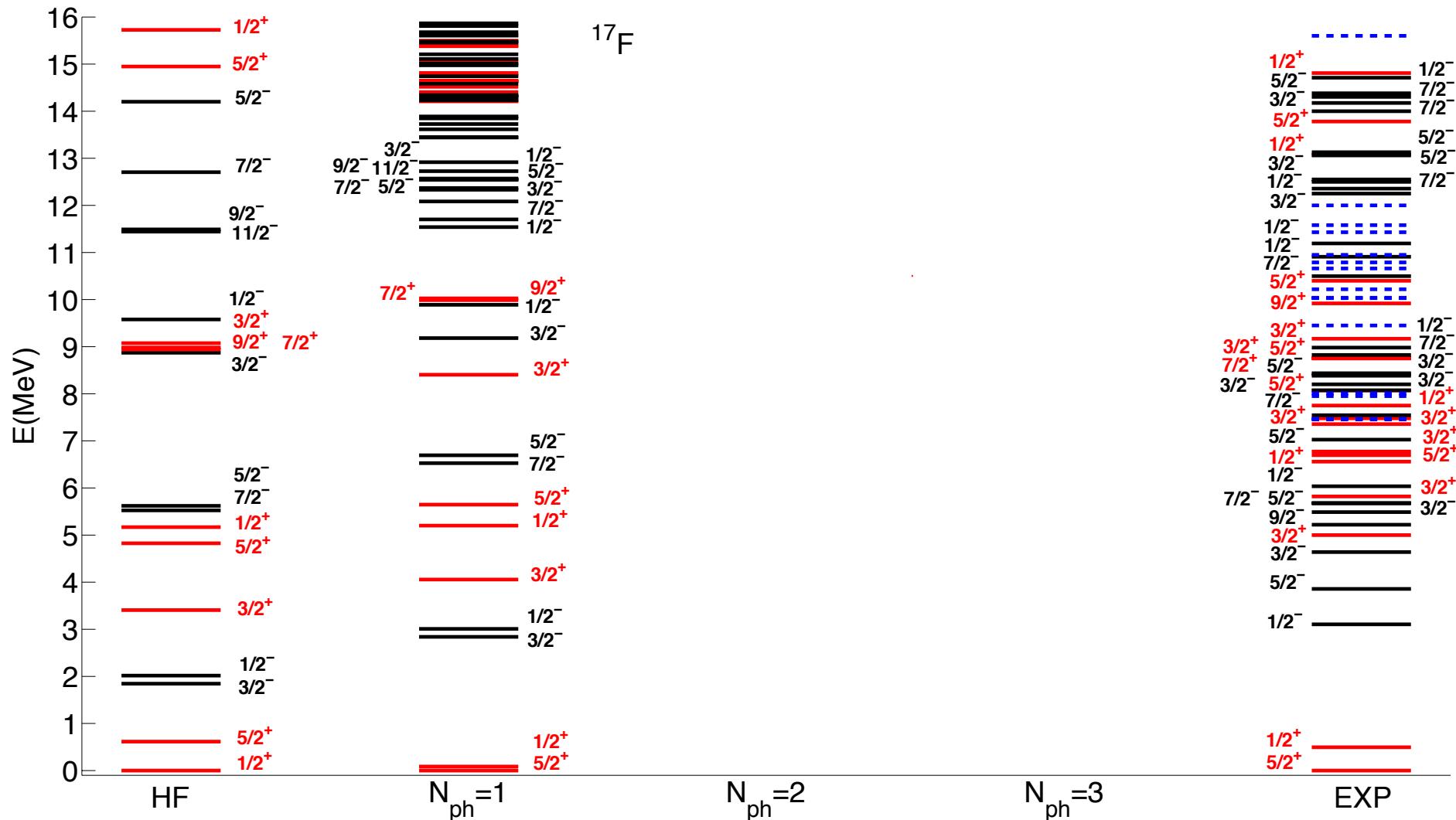
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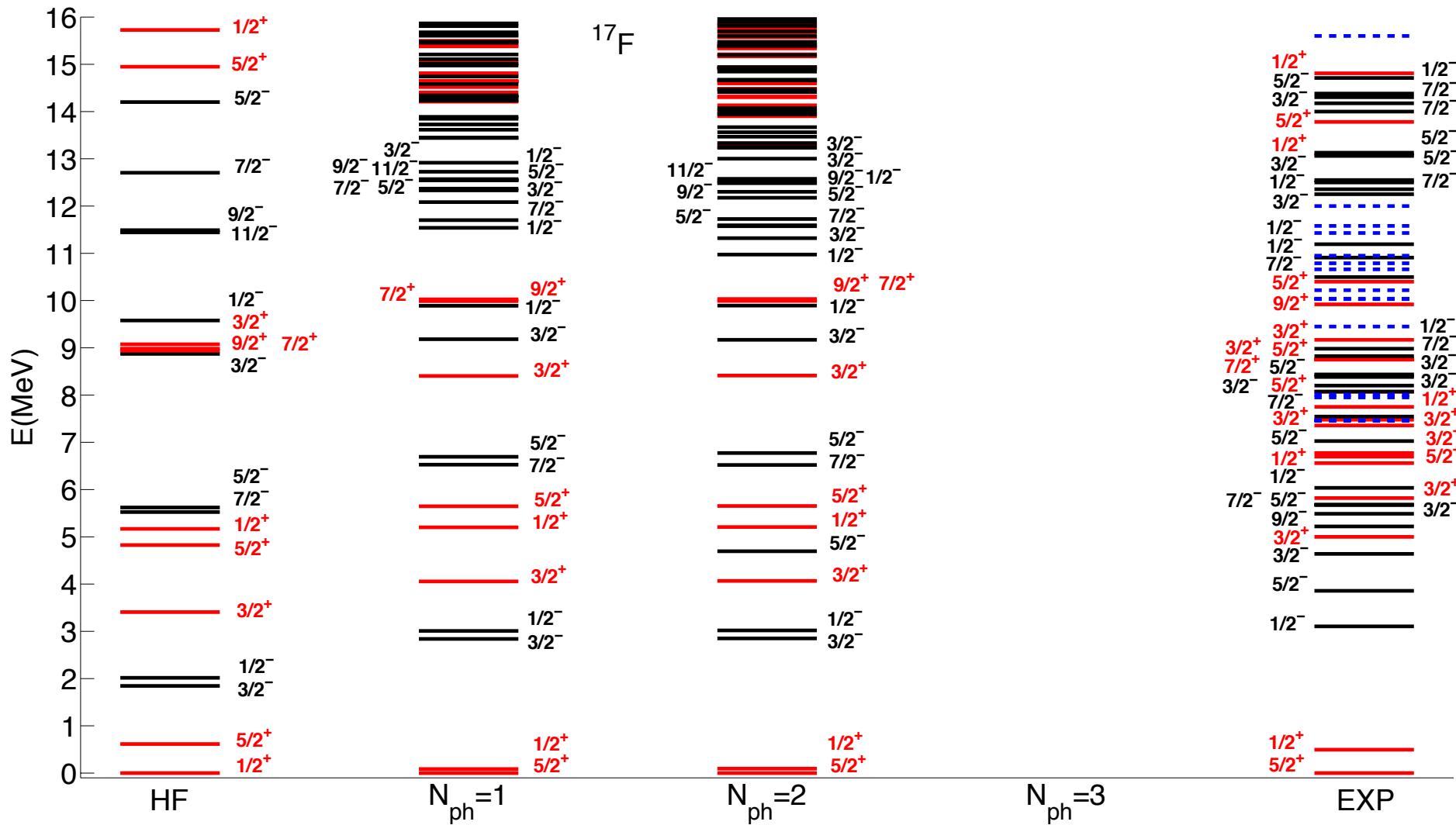
Application: ^{17}F spectra



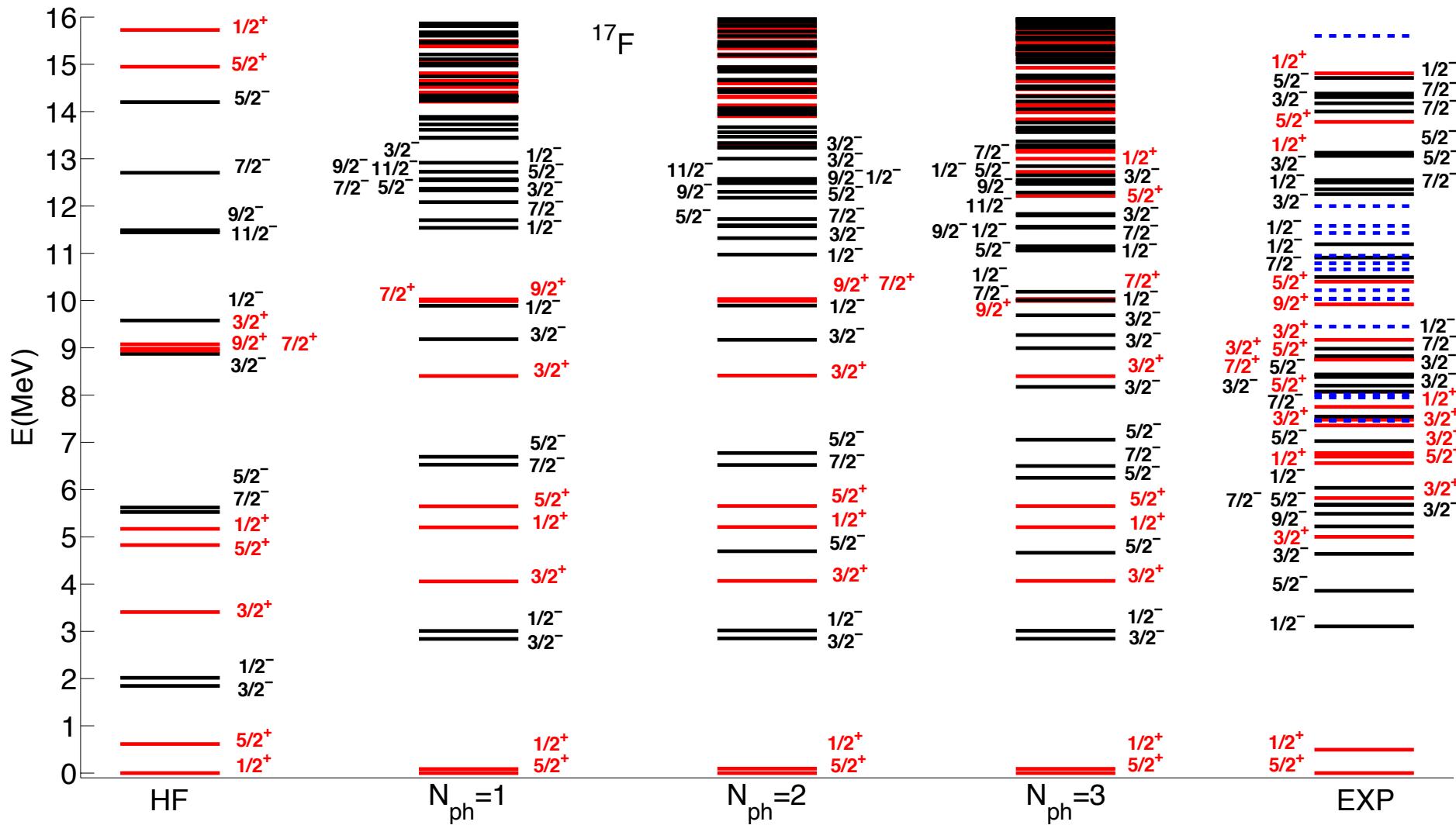
Application: ^{17}F spectra



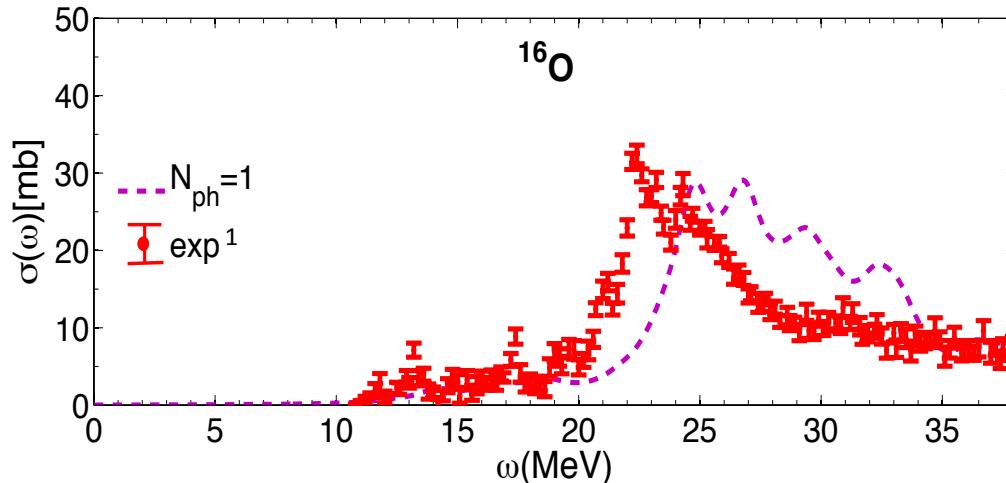
Application: ^{17}F spectra



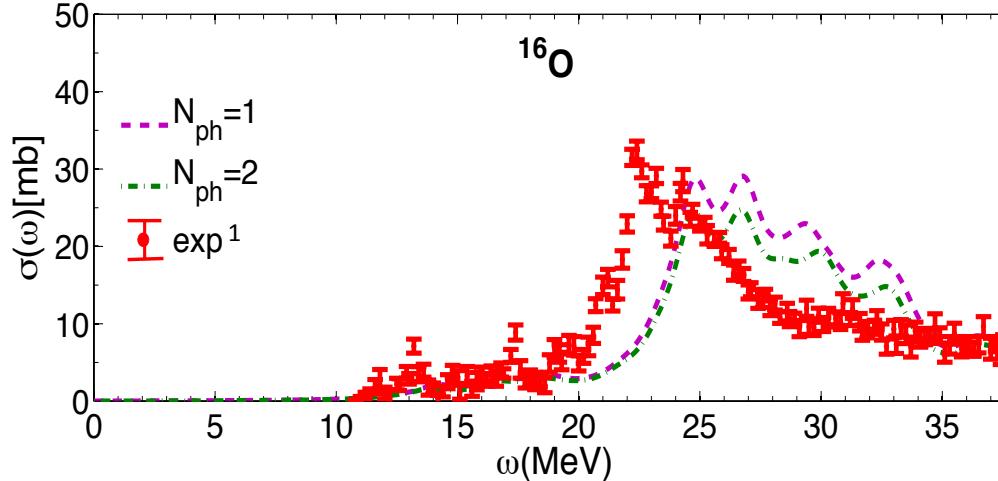
Application: ^{17}F spectra



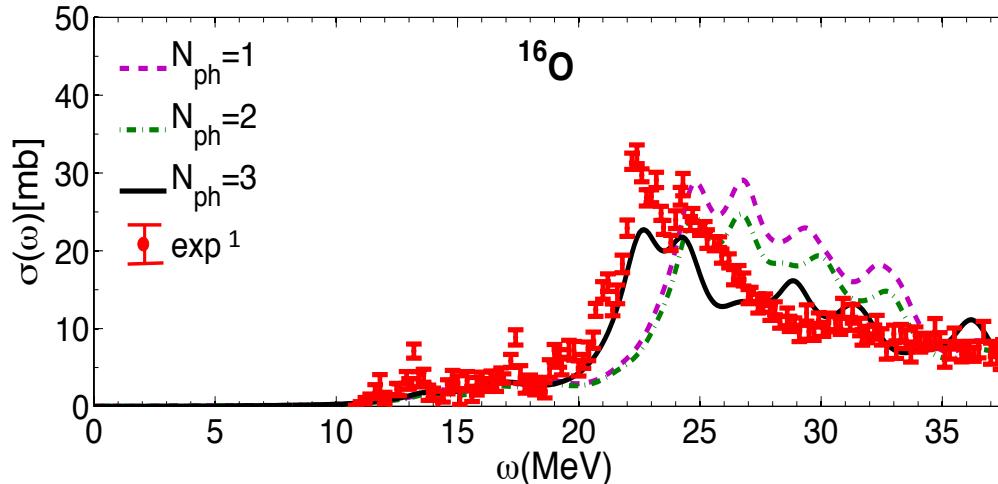
Cross section



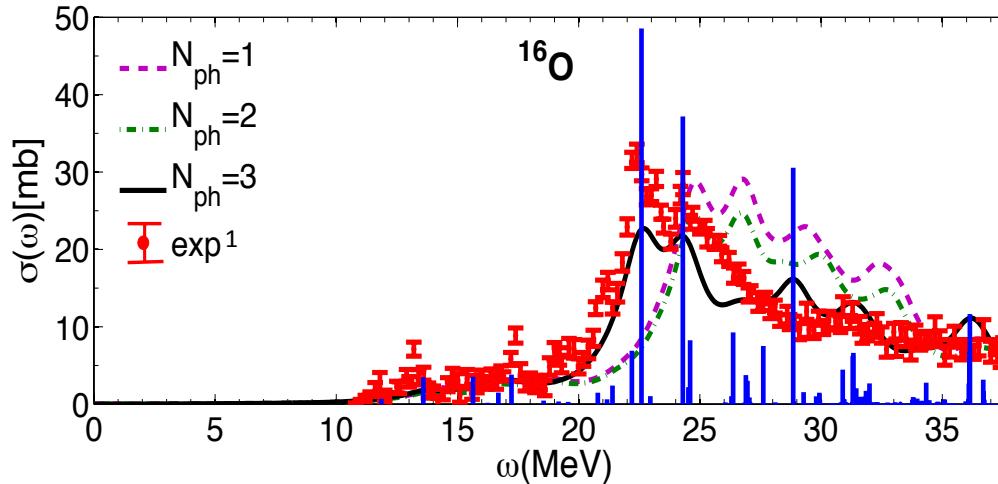
Cross section



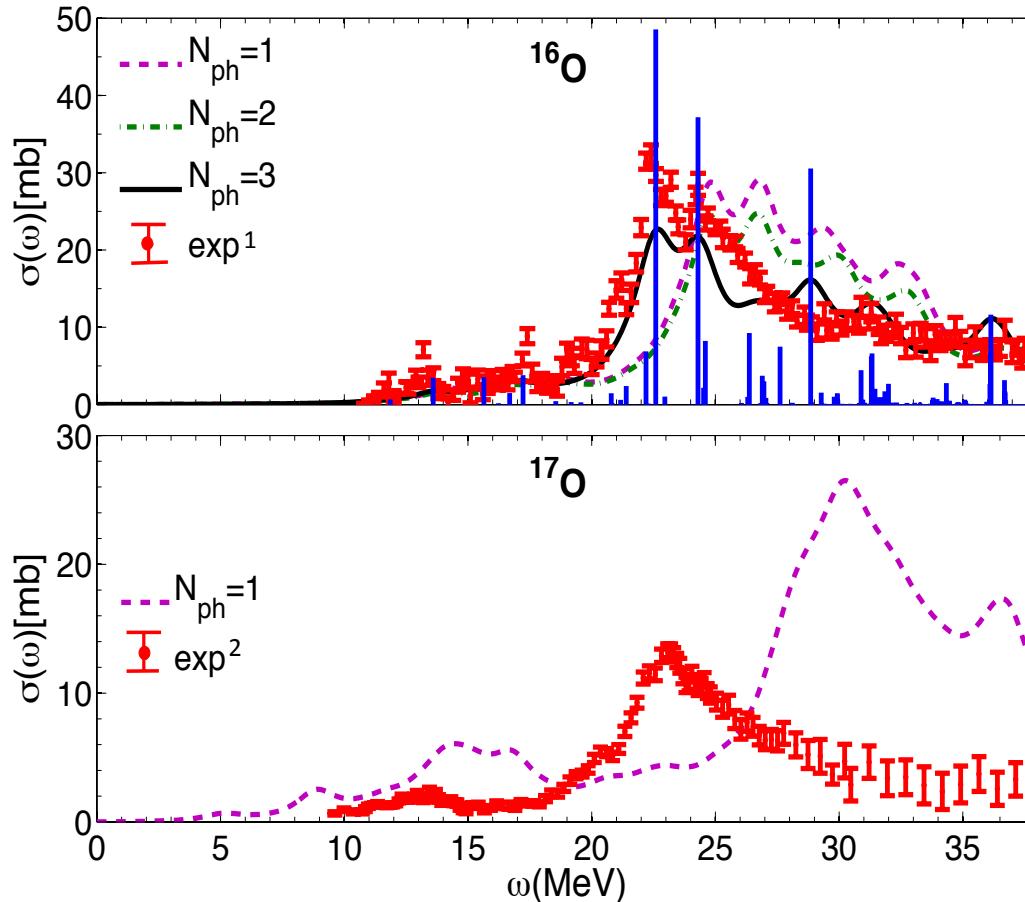
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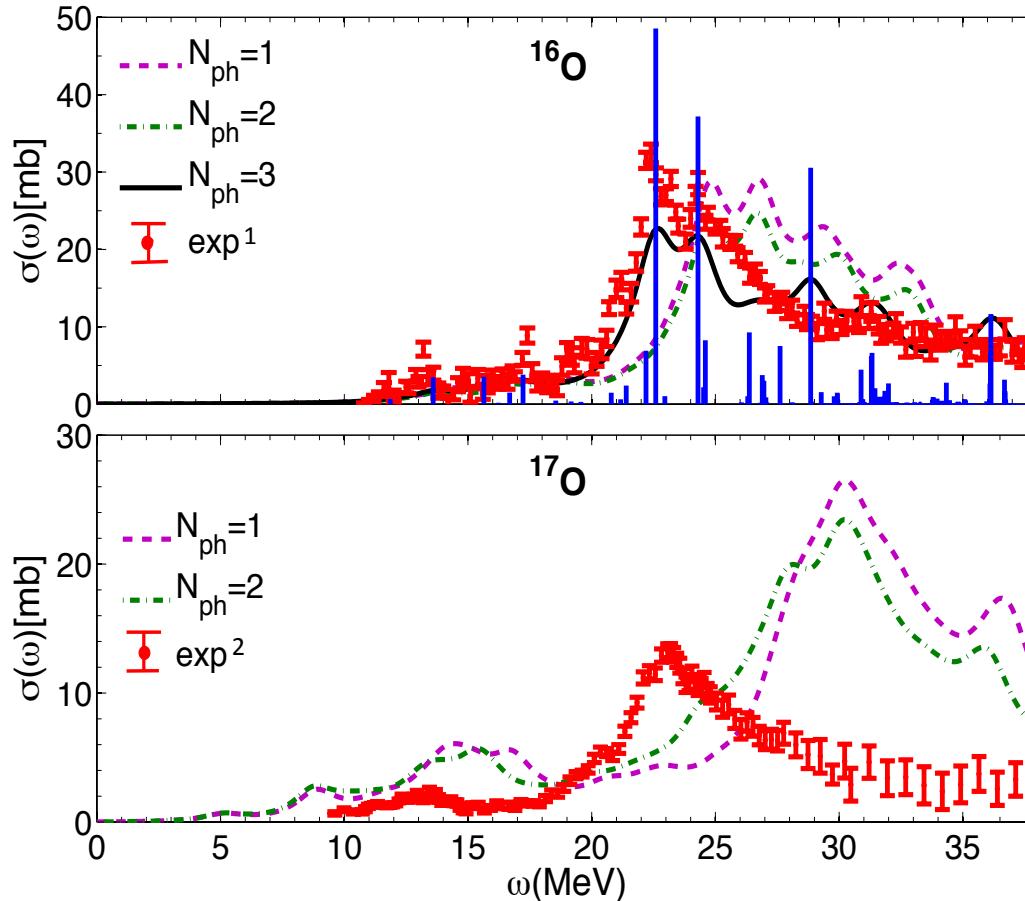
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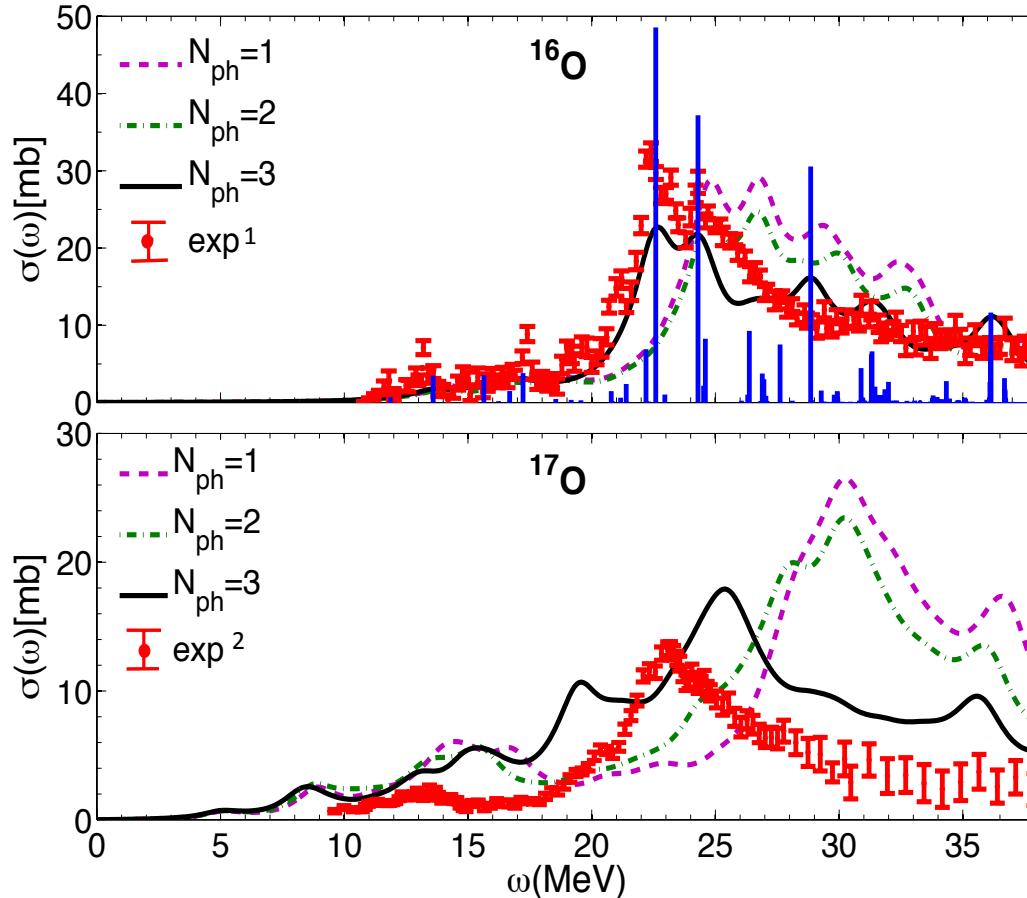
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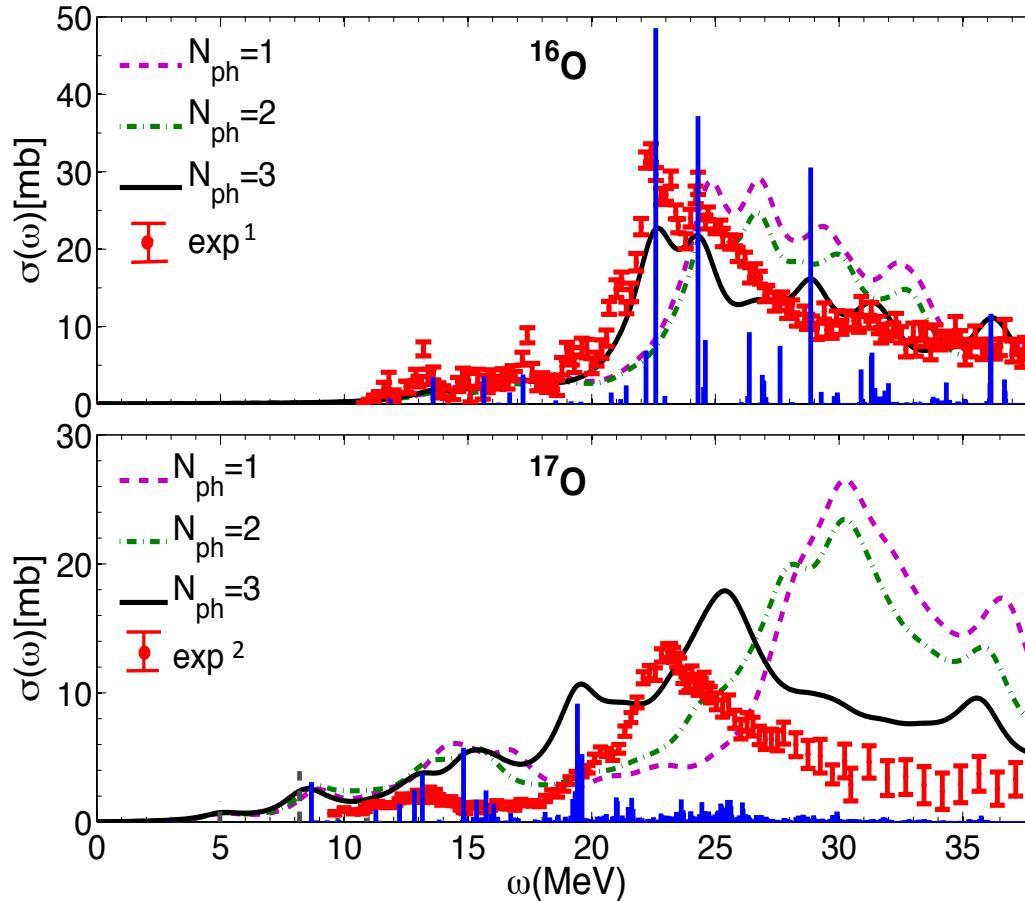
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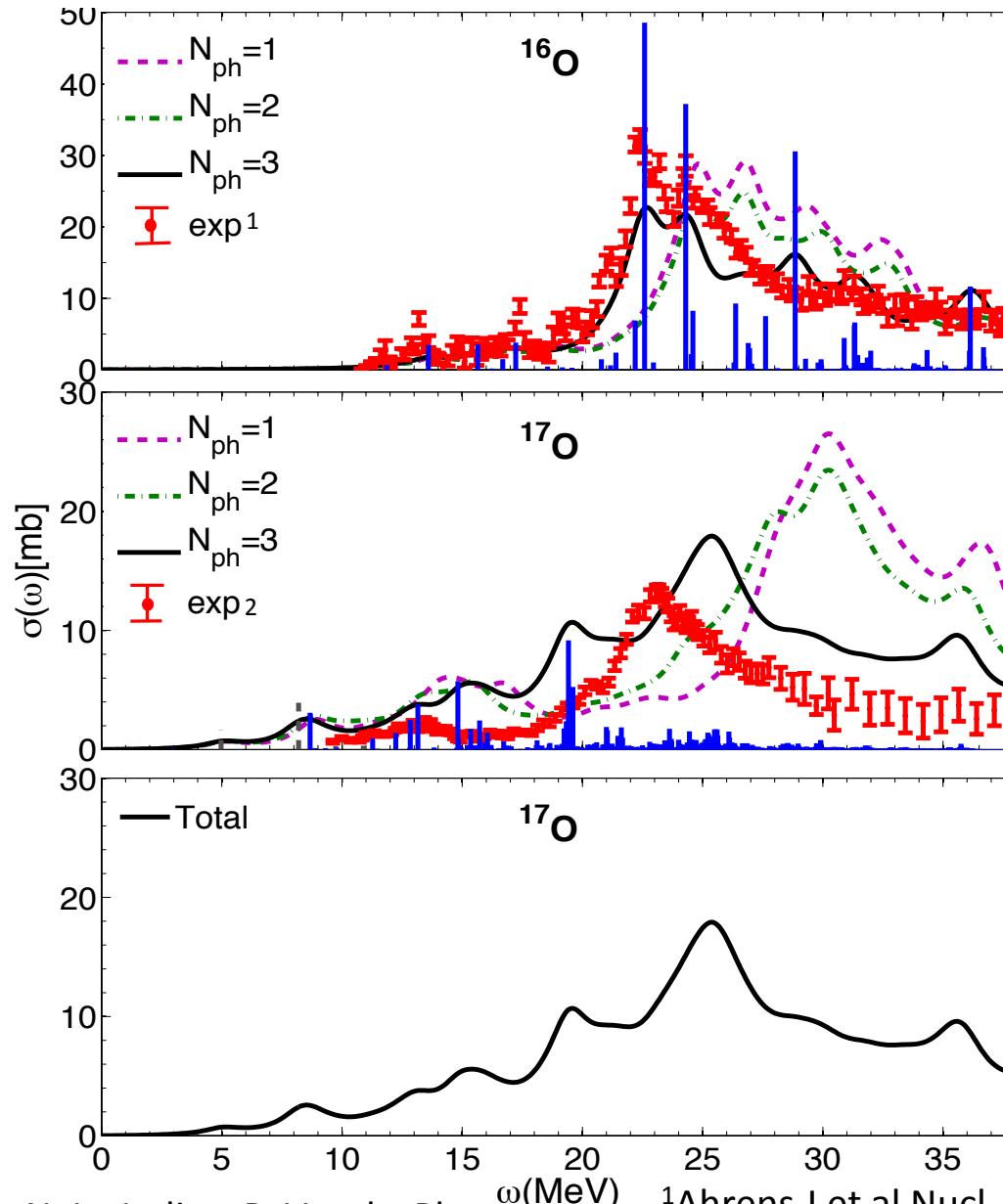
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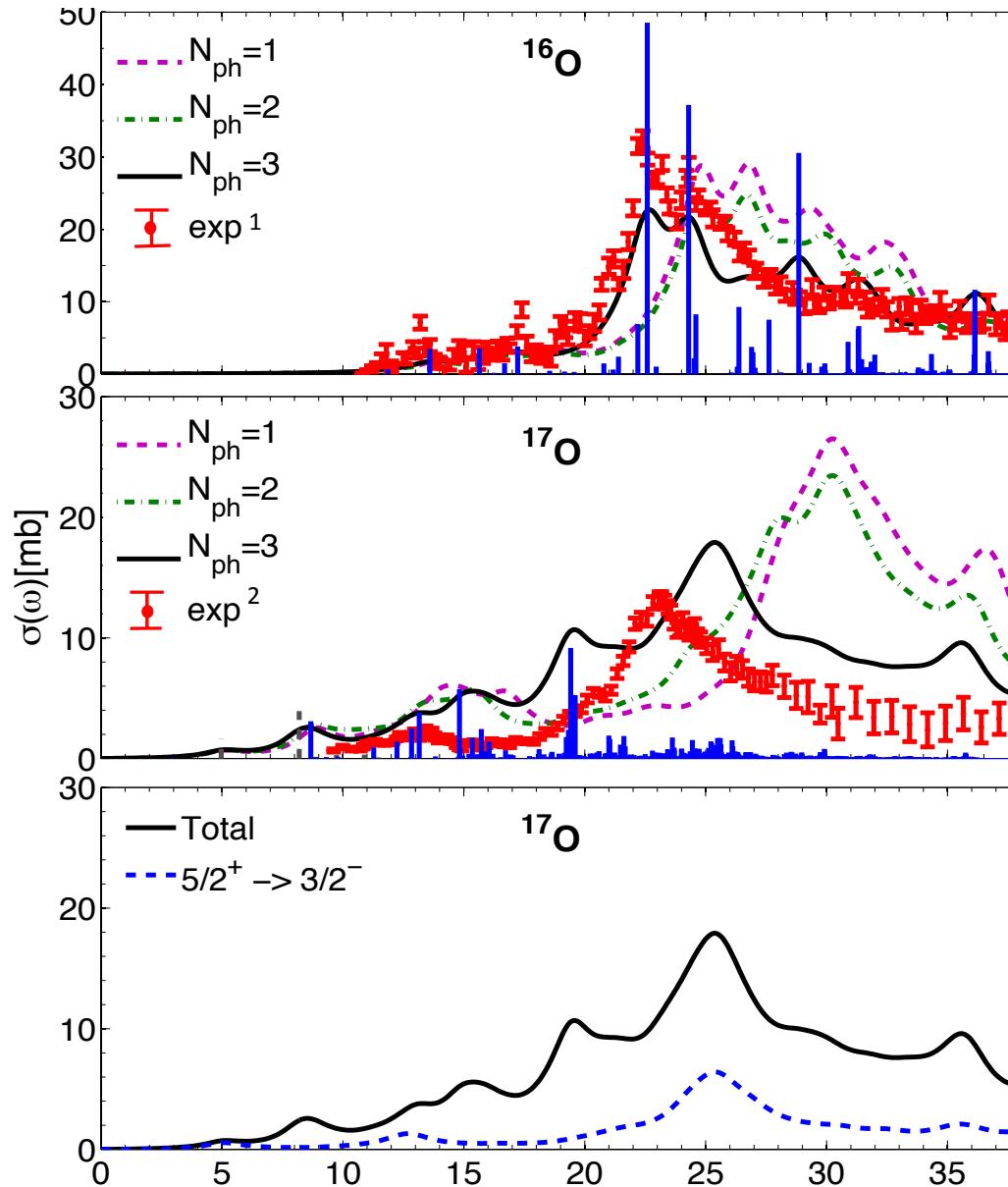
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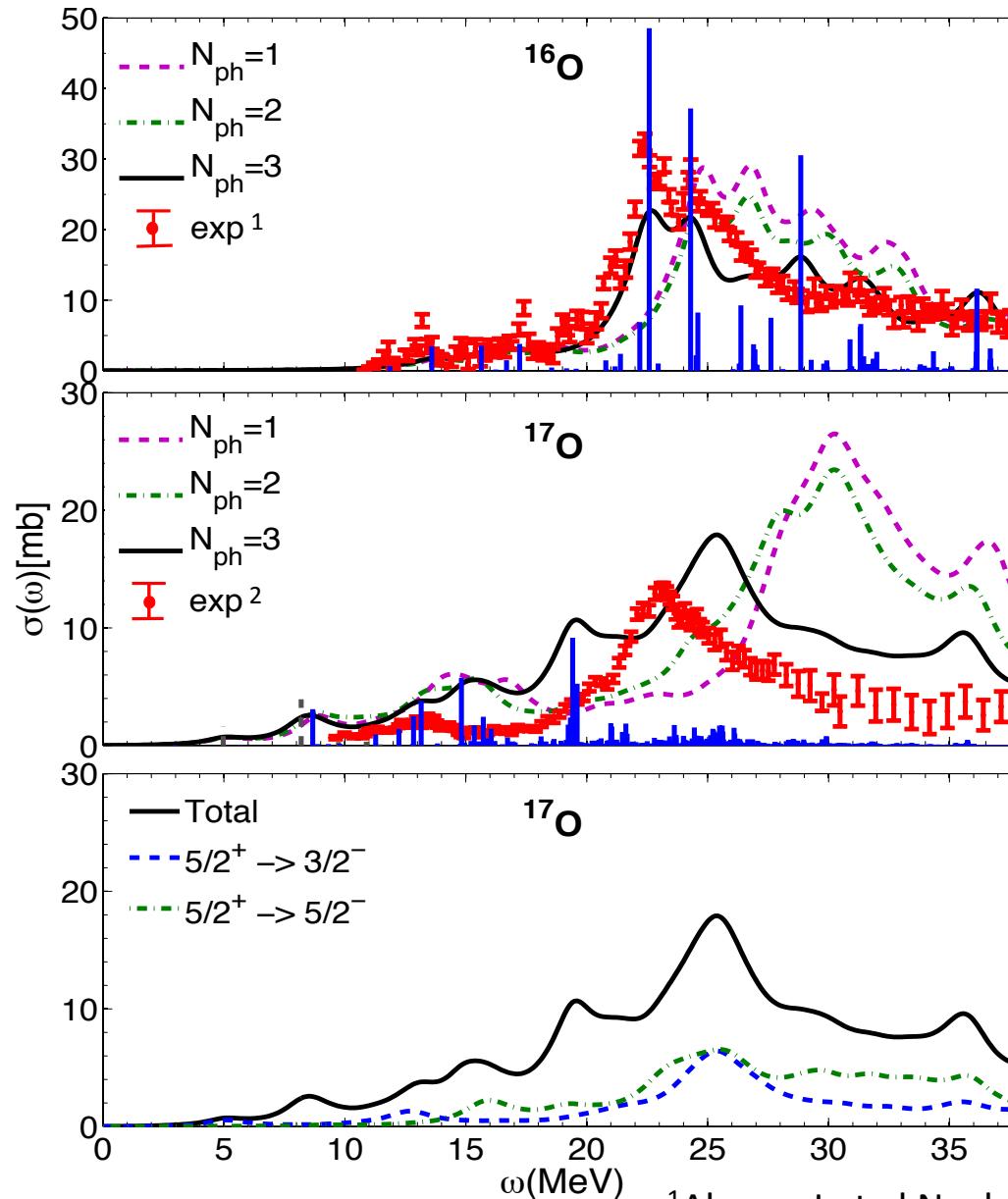
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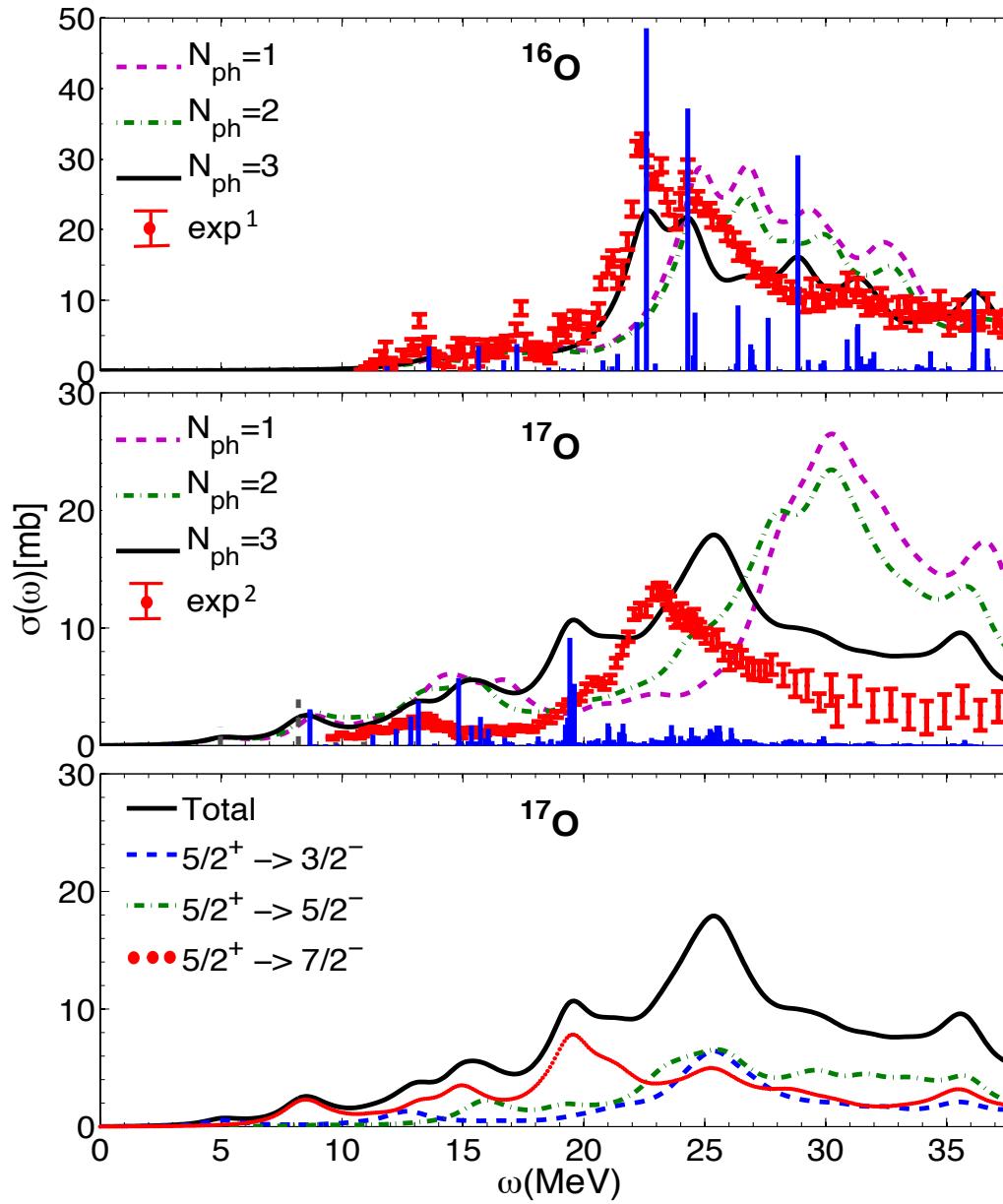
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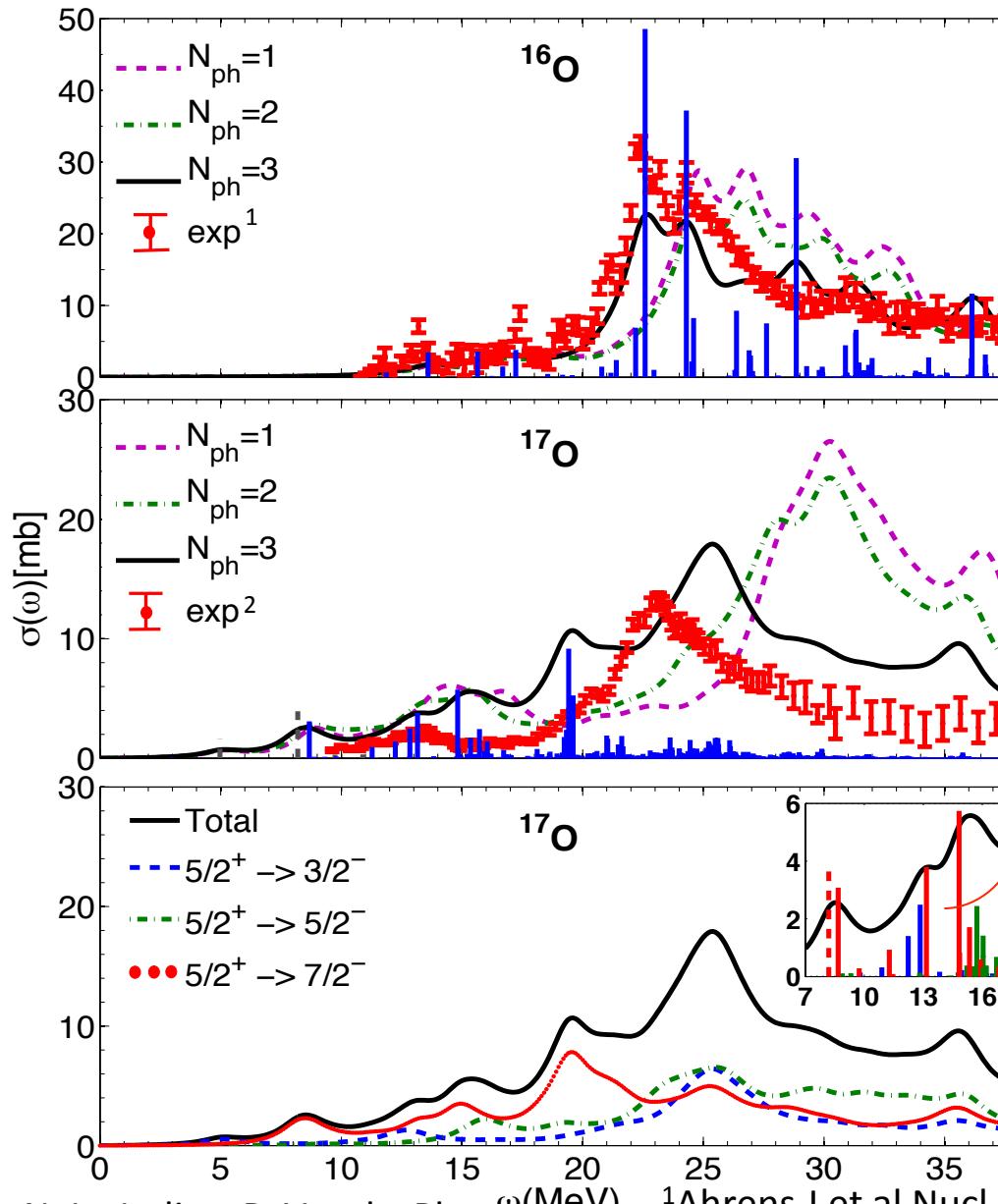
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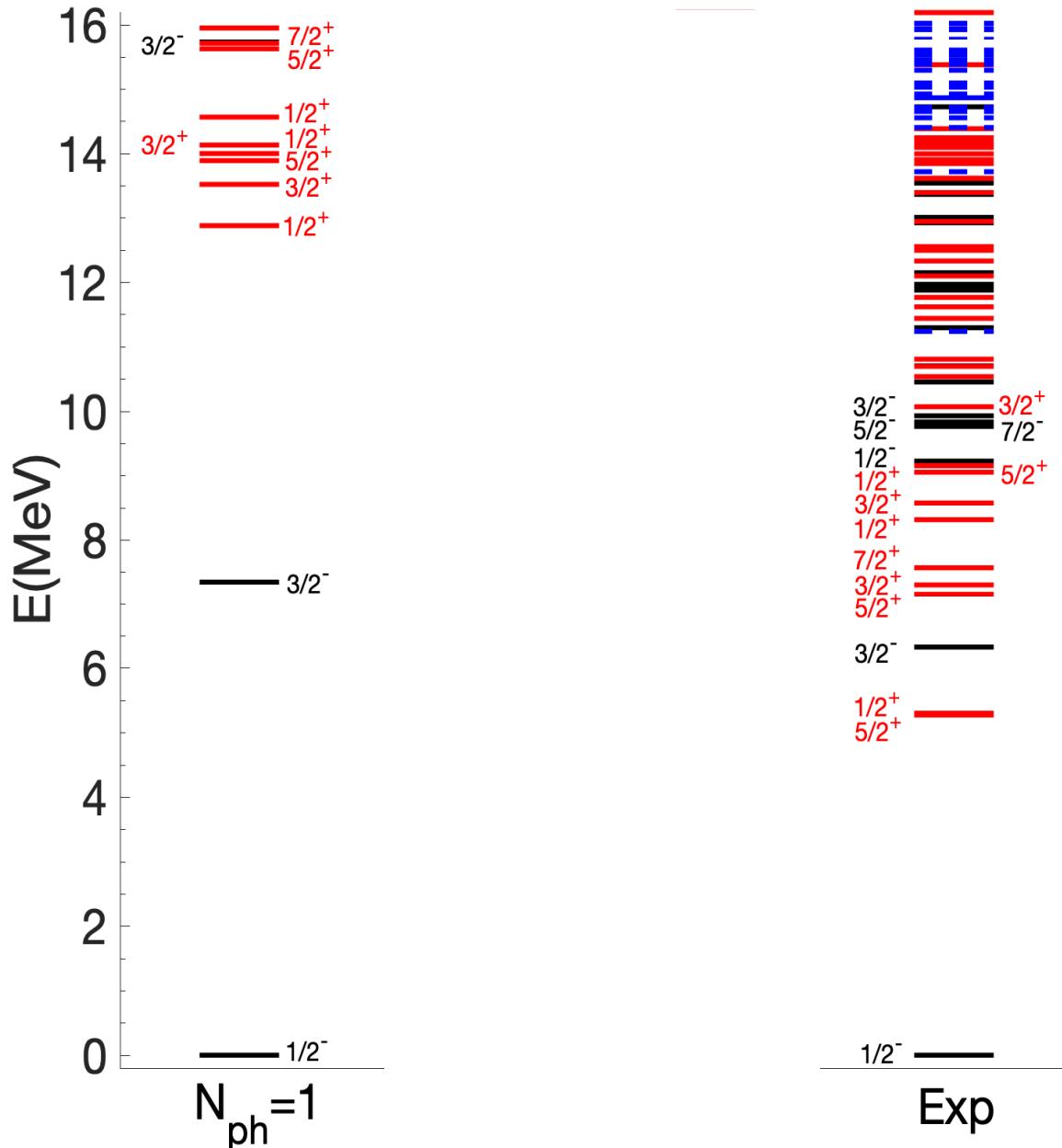
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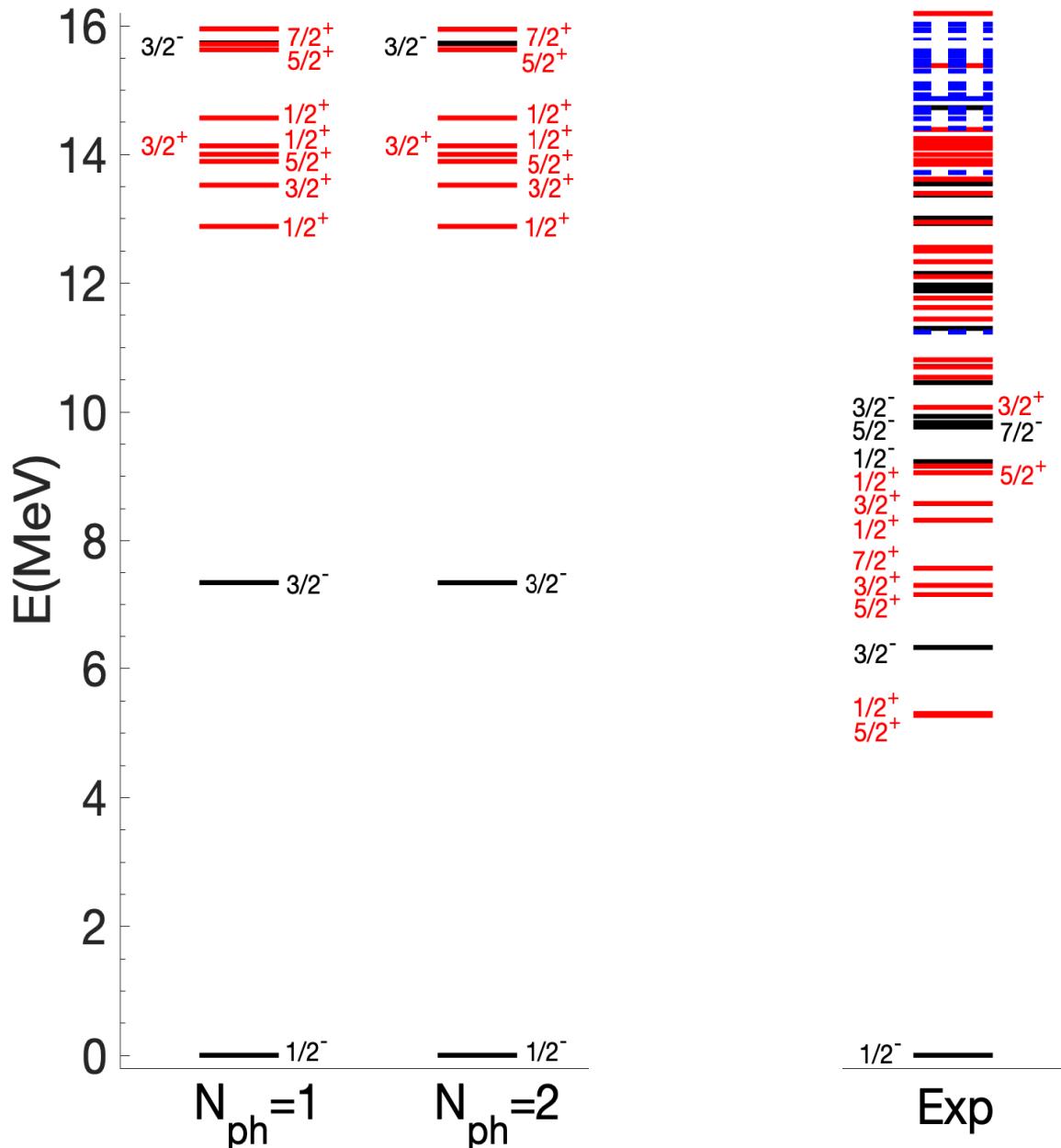
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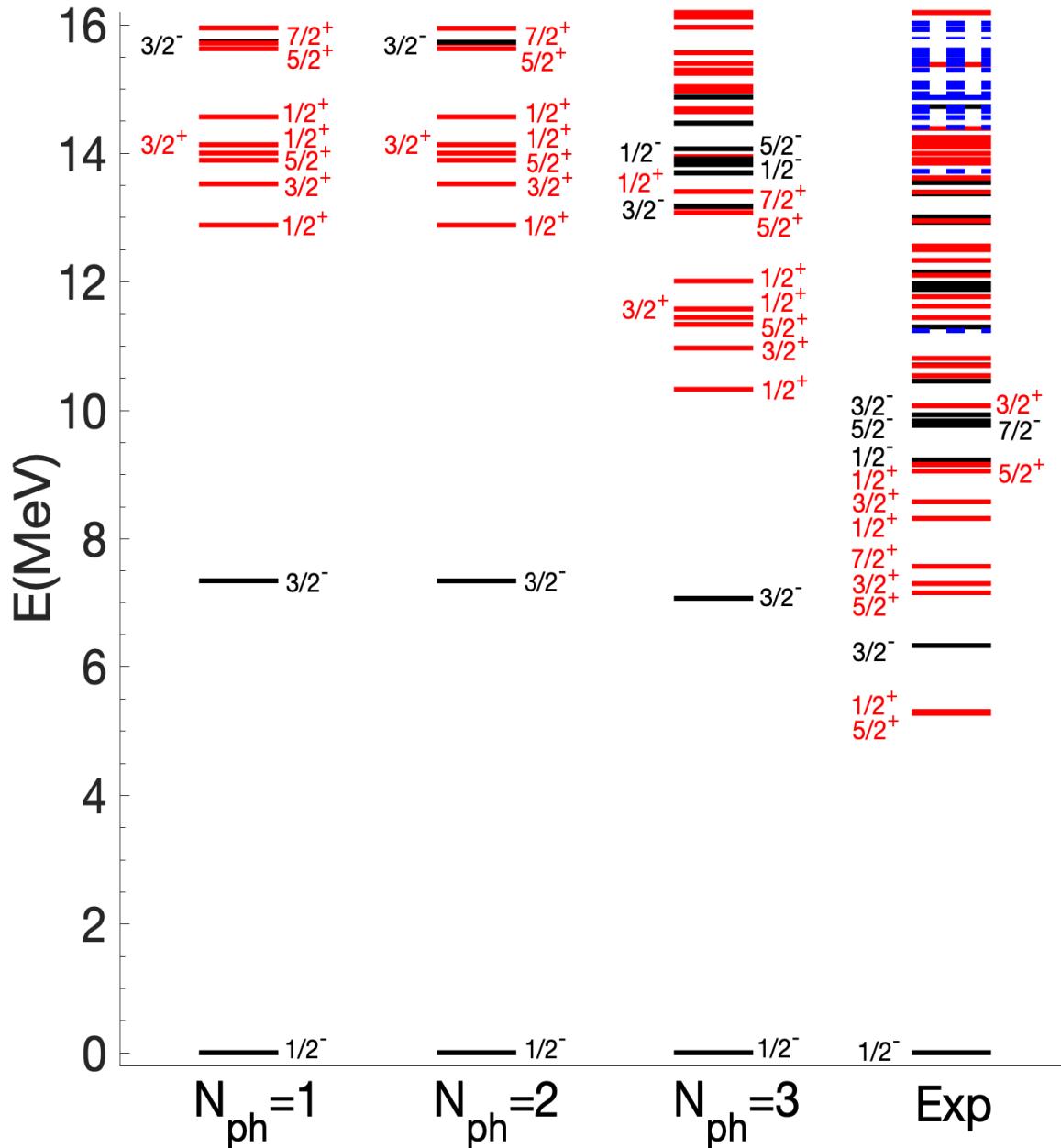
Application: ^{15}N spectra



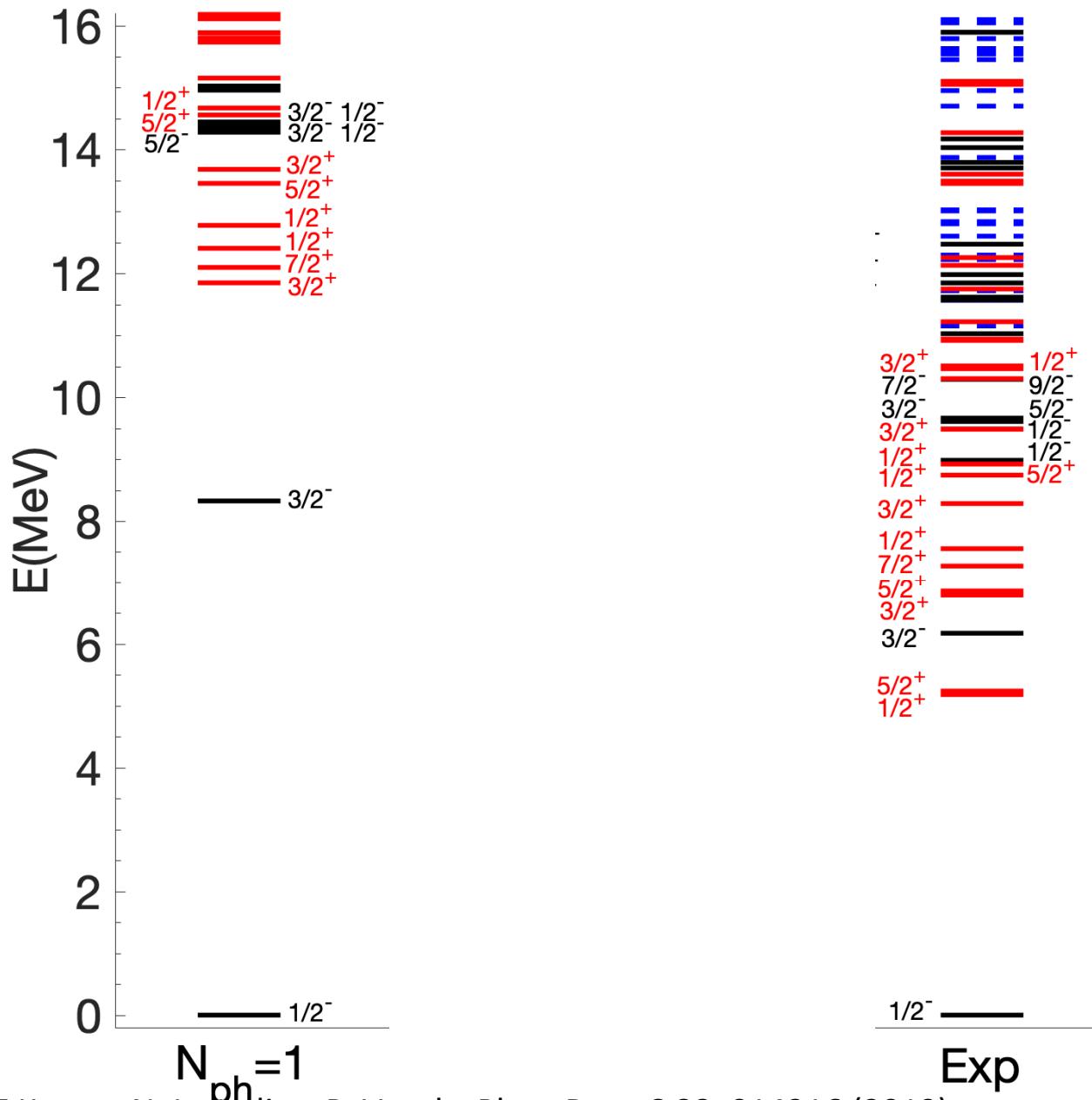
Application: ^{15}N spectra



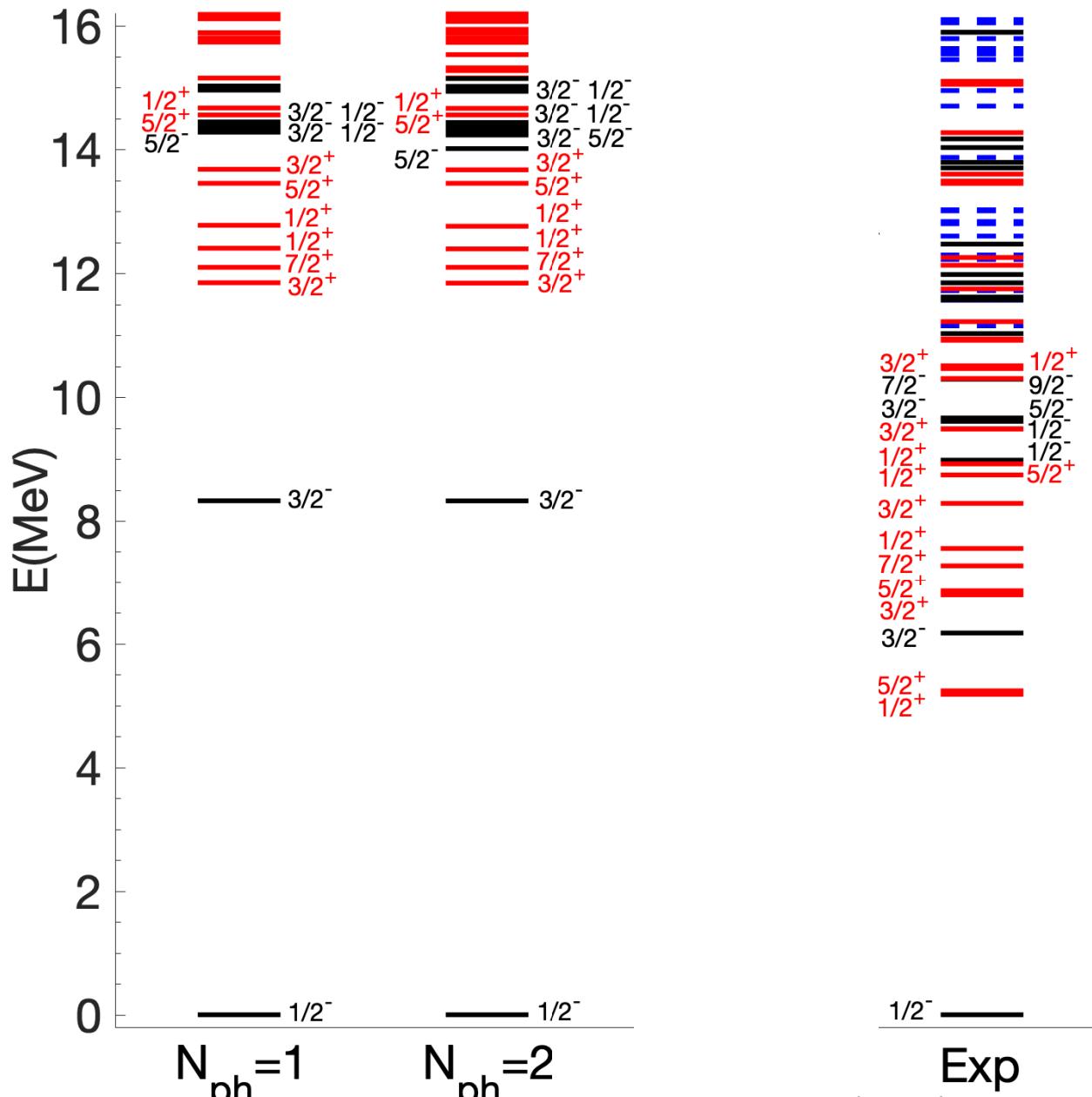
Application: ^{15}N spectra



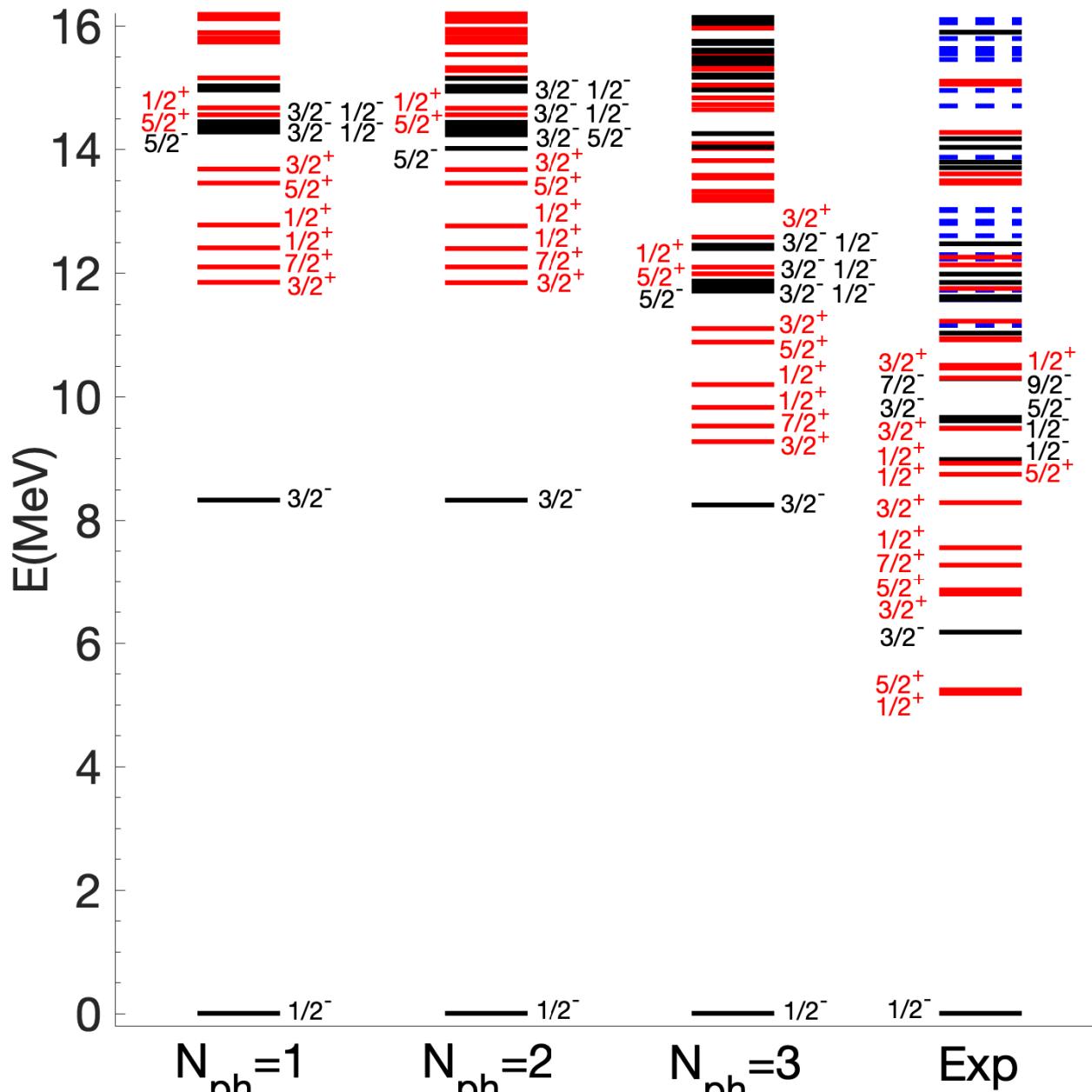
Application: ^{15}O spectra



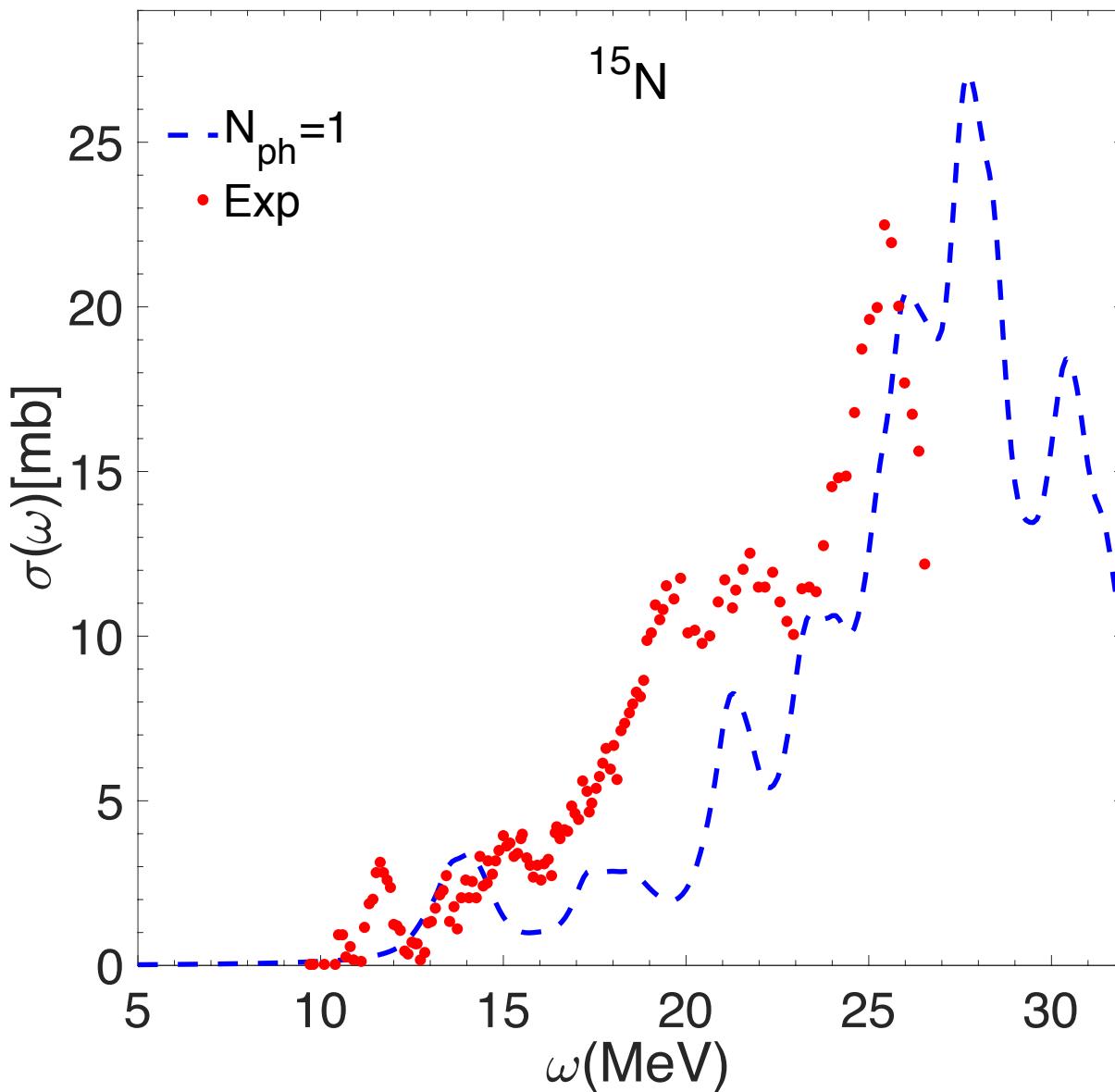
Application: ^{15}O spectra



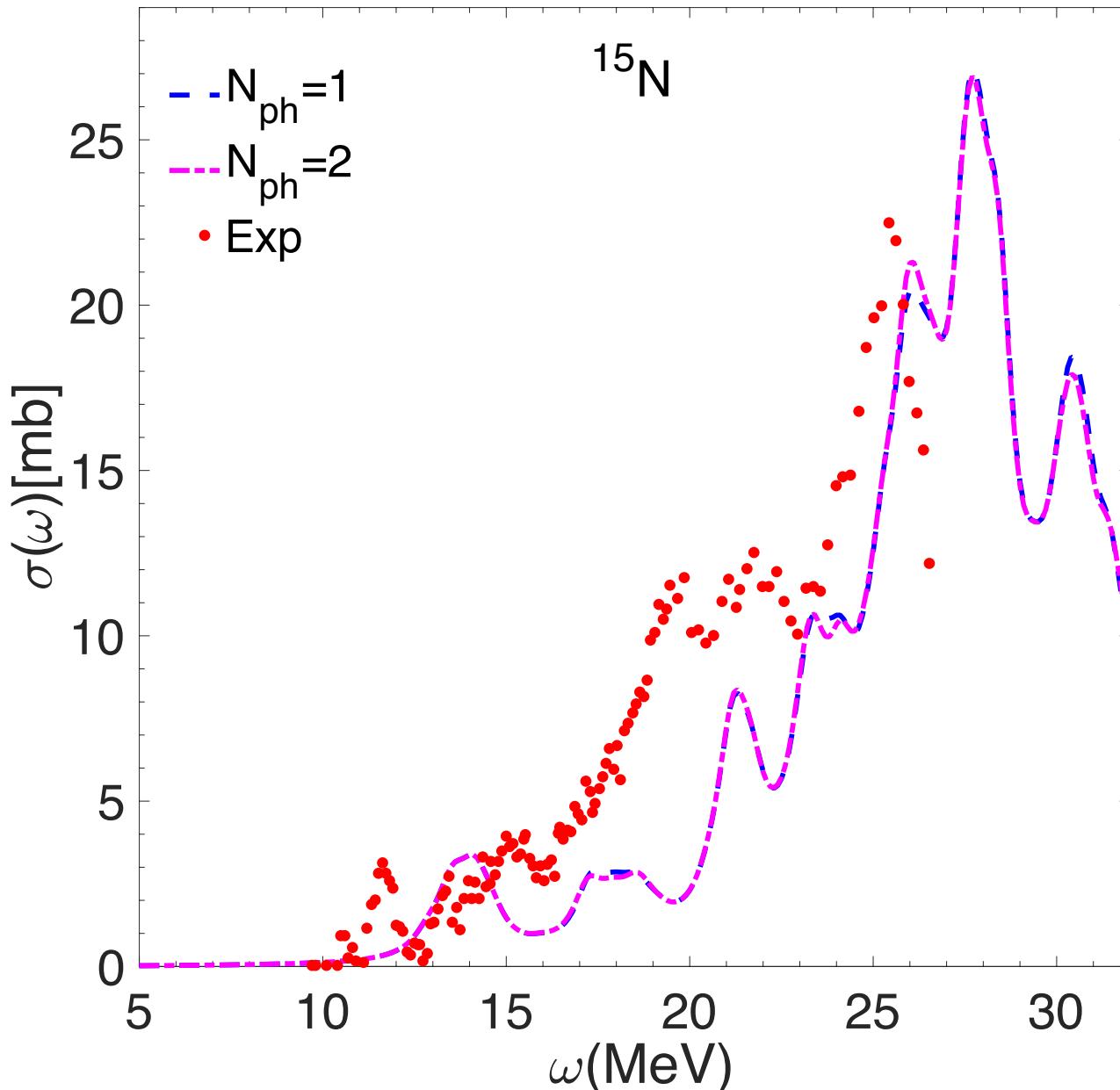
Application: ^{15}O spectra



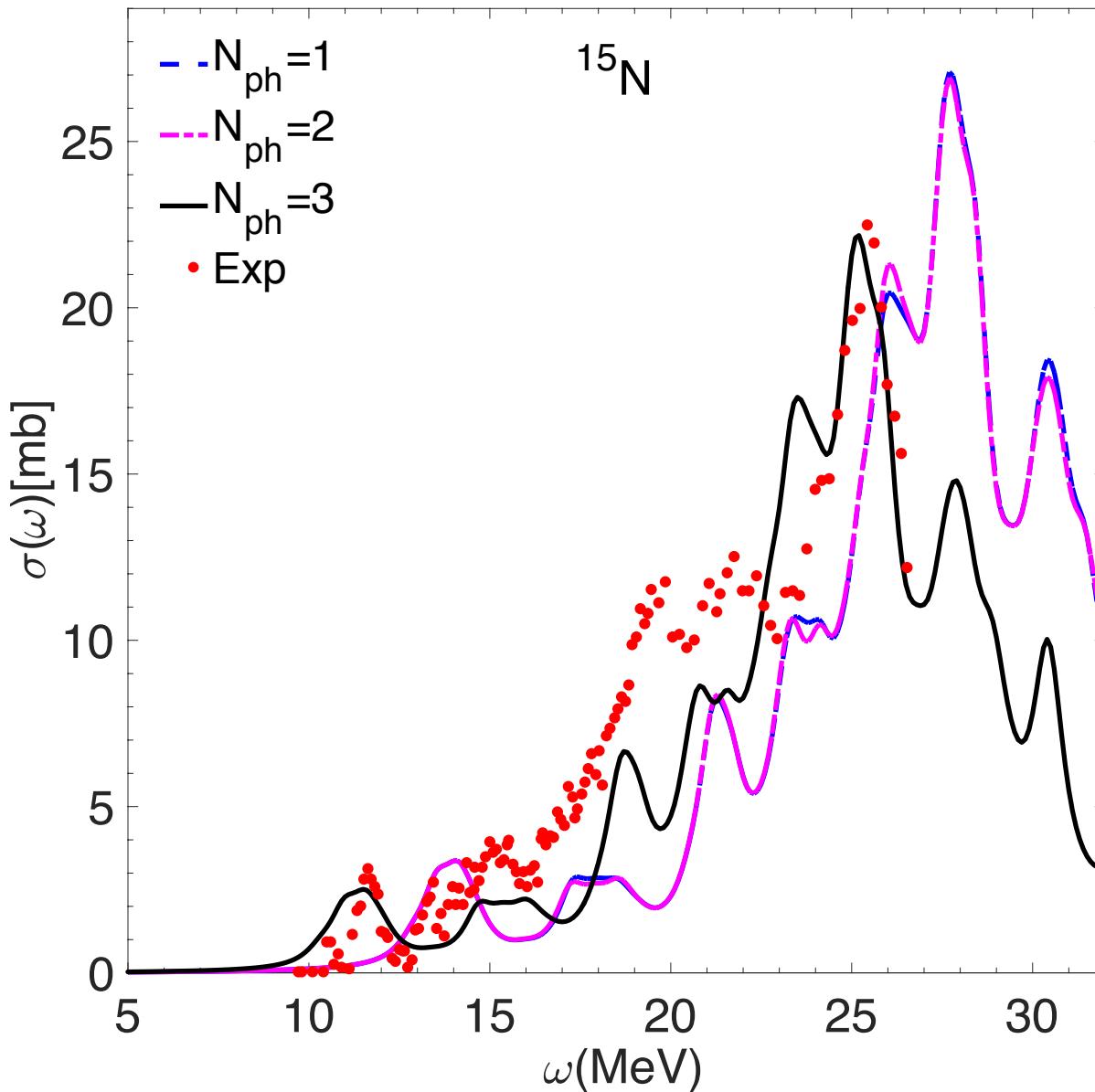
E1 response in ^{15}N



E1 response in ^{15}N



E1 response in ^{15}N



Ground state moments and transitions in A=15

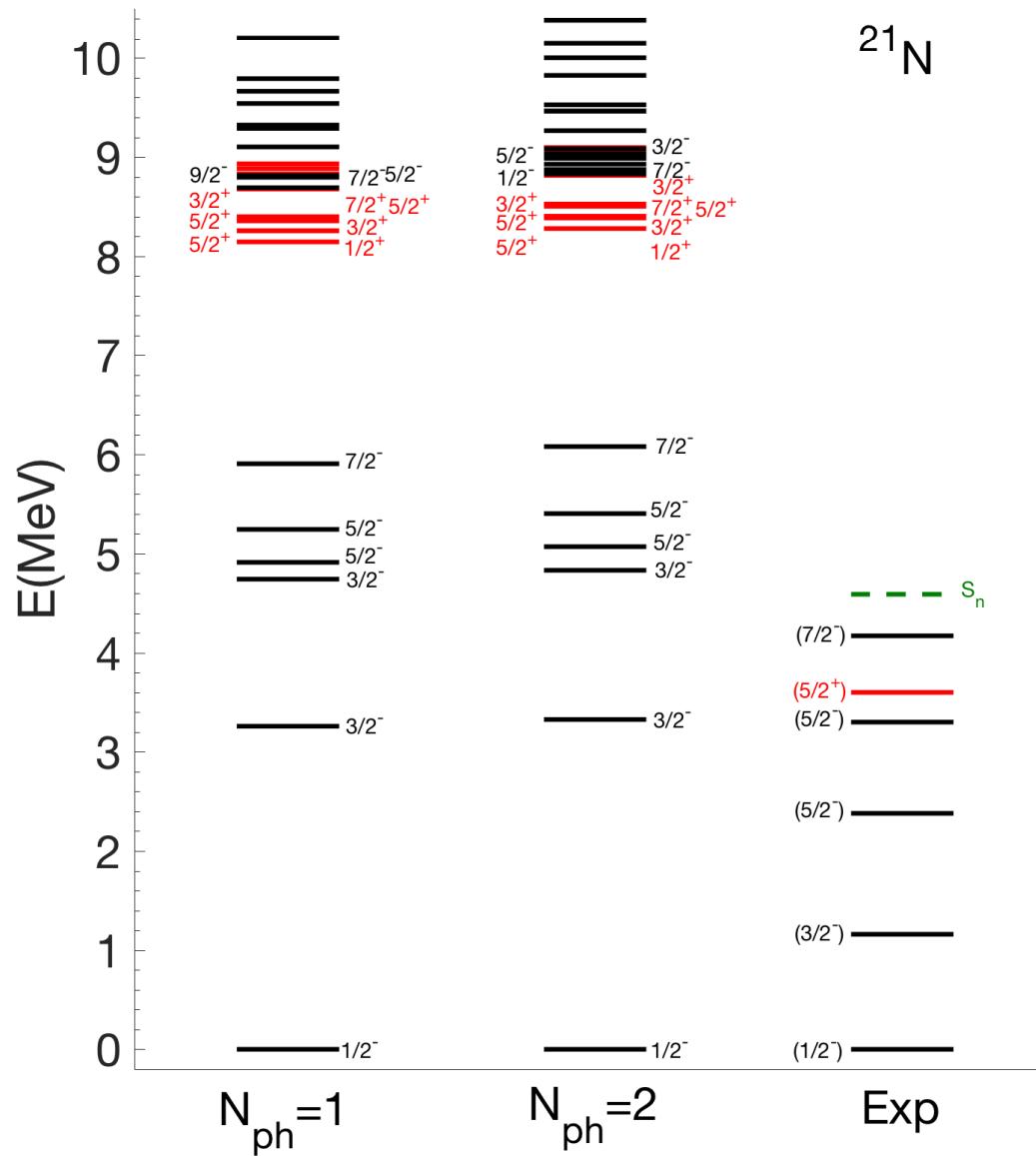
^{15}O

ν_f	Theory	Exp
μ	+0.5986	$\pm 0.7189(8)$
$\log ft$	3.650	3.637
$B(E1; 1/2_1^+ \rightarrow 1/2_1^-)$	0.06	$(1.4 \pm 0.2) \times 10^{-3}$
$B(E2; 3/2_1^- \rightarrow 1/2_1^-)$	0.03	>0.28
$B(E3; 5/2_1^+ \rightarrow 1/2_1^-)$	3.95	4 ± 2
$B(M1; 3/2_1^- \rightarrow 1/2_1^-)$	0.56	$>5.3 \times 10^{-2}$

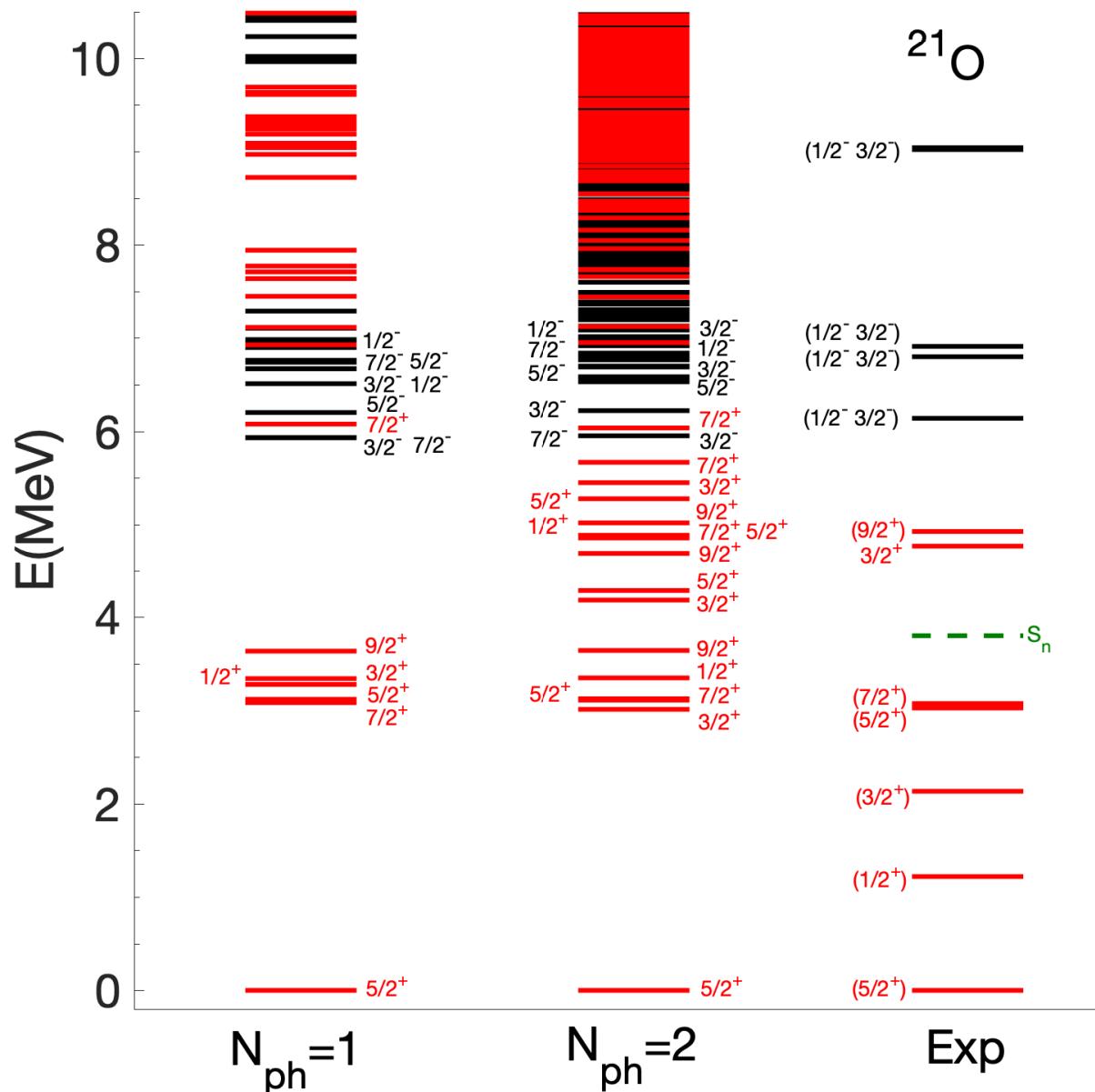
^{15}N

ν_f	Theory	Exp
μ	-0.2499	-0.2832
$B(E1; 1/2_1^+ \rightarrow 1/2_1^-)$	0.03	$(4.3 \pm 1.1) \times 10^{-4}$
$B(E2; 3/2_1^- \rightarrow 1/2_1^-)$	1.20	2.91 ± 0.24
$B(E3; 5/2_1^+ \rightarrow 1/2_1^-)$	3.11	7 ± 2
$B(M1; 3/2_1^- \rightarrow 1/2_1^-)$	0.687	0.578 ± 0.015

Application to neutron rich nuclei: ^{21}N spectra



Application to neutron rich nuclei: ^{21}O spectra



Ground state β decay of $^{21}\text{N} \rightarrow ^{21}\text{O}$

v_f	ω_f (MeV)	$\log(ft)$	$B(GT)$
$3/2^-$	5.95	7.14	0.00044
$3/2^-$	6.53	7.62	0.00015
$1/2^-$	6.69	8.29	0.00003
$1/2^-$	7.30	7.55	0.00016
$3/2^-$	10.02	5.60	0.01513
$3/2^-$	10.43	6.30	0.00306
$3/2^-$	13.05	5.70	0.00004
$3/2^+$	14.54	5.32	0.01221

v_f	ω_f (MeV) _{EXP}	$\log(ft)$ _{EXP}	$B(GT)$ _{EXP}
$(1/2^-, 3/2^-)$	6.14	5.44 ± 0.06	0.0224 ± 0.0032
$(1/2^-, 3/2^-)$	6.80	5.19 ± 0.06	0.0399 ± 0.0056
$(1/2^-, 3/2^-)$	6.91	5.44 ± 0.07	0.0224 ± 0.0035
$(1/2^-, 3/2^-)$	9.02	4.78 ± 0.06	0.1015 ± 0.0145
$(1/2^-, 3/2^-)$	9.04	4.62 ± 0.06	0.1462 ± 0.0206

Concluding remarks

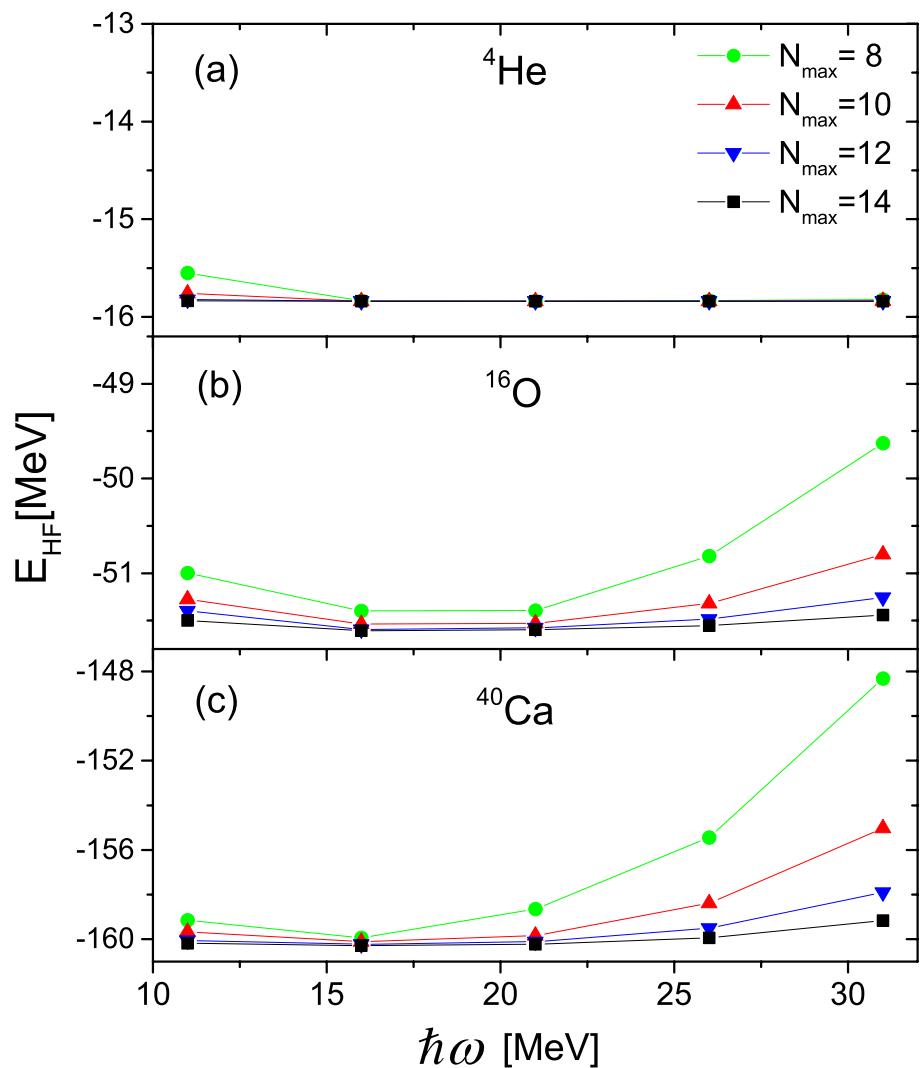
- We have seen that the multiphonon play an important role.
- Especially the three-phonon which through their coupling to one-phonon states enhance the density of states at low-energy in agreement with experiments.
- The $E1$ cross sections of ^{17}O and ^{15}N and ^{22}O are in fair agreement with the experimental data.

Serious discrepancies:

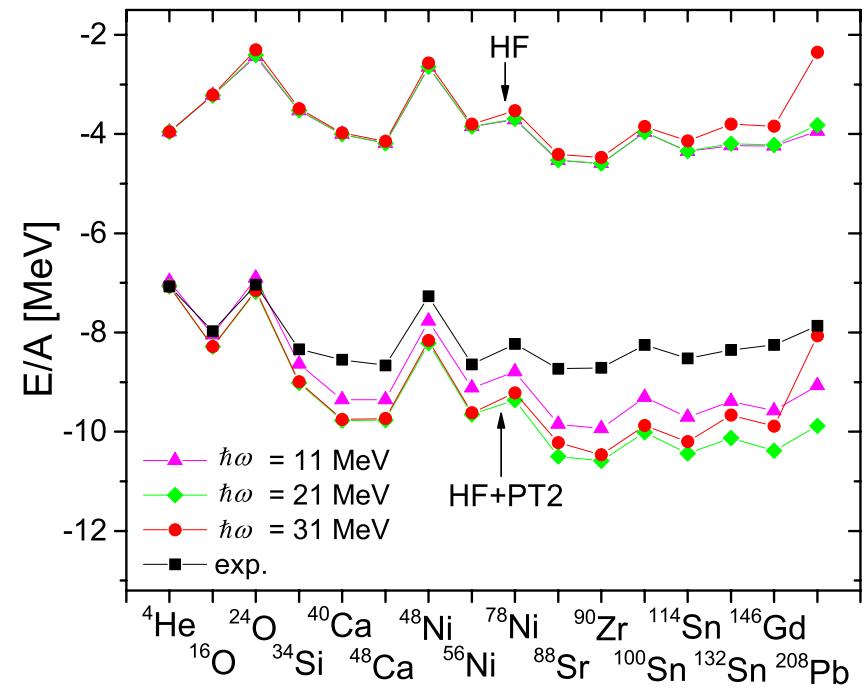
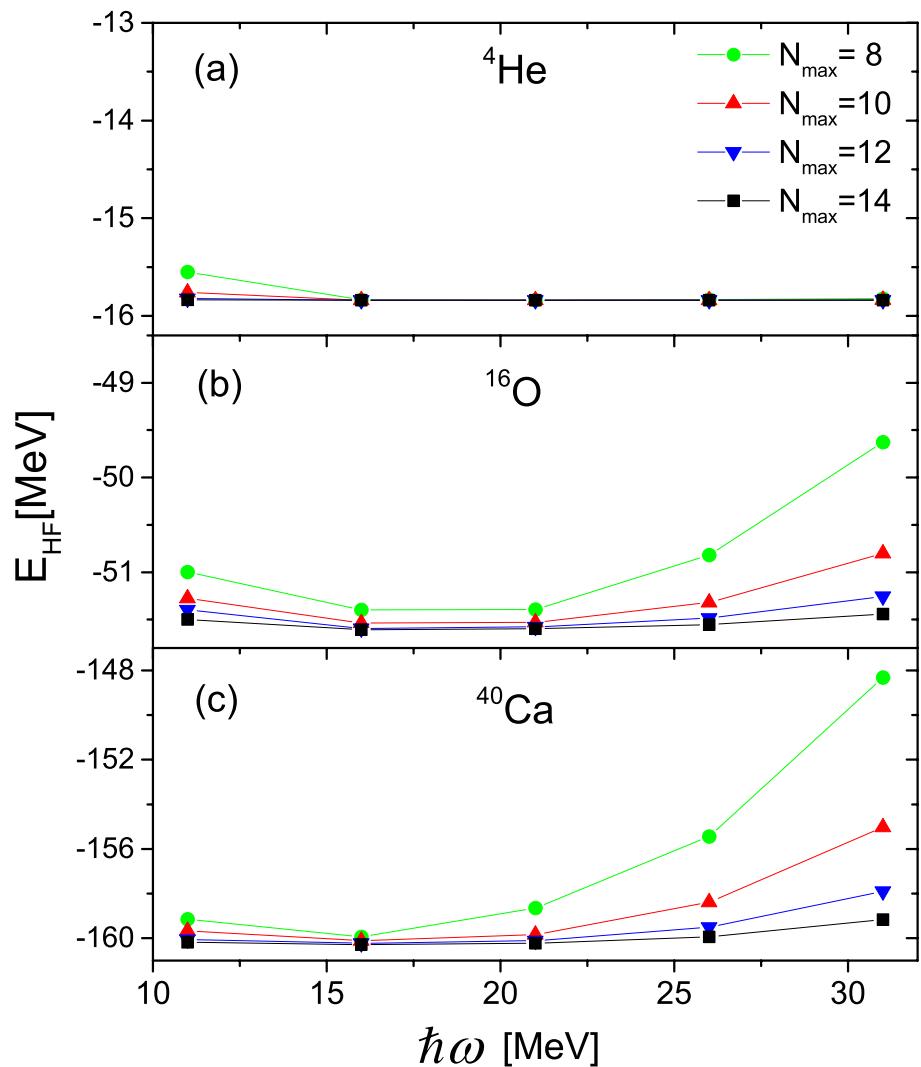
- The positive parity states of $A=17$ remain at high energy.
 - The low-lying states of $A=15$ are too high in energy.
 - **Recipies:**
- 1) Use of a better potential: Preliminary results using NNLO_{sat} potential, which includes explicitly the three-body force, are promising.
 - 2) Need of four phonon states

Thank you

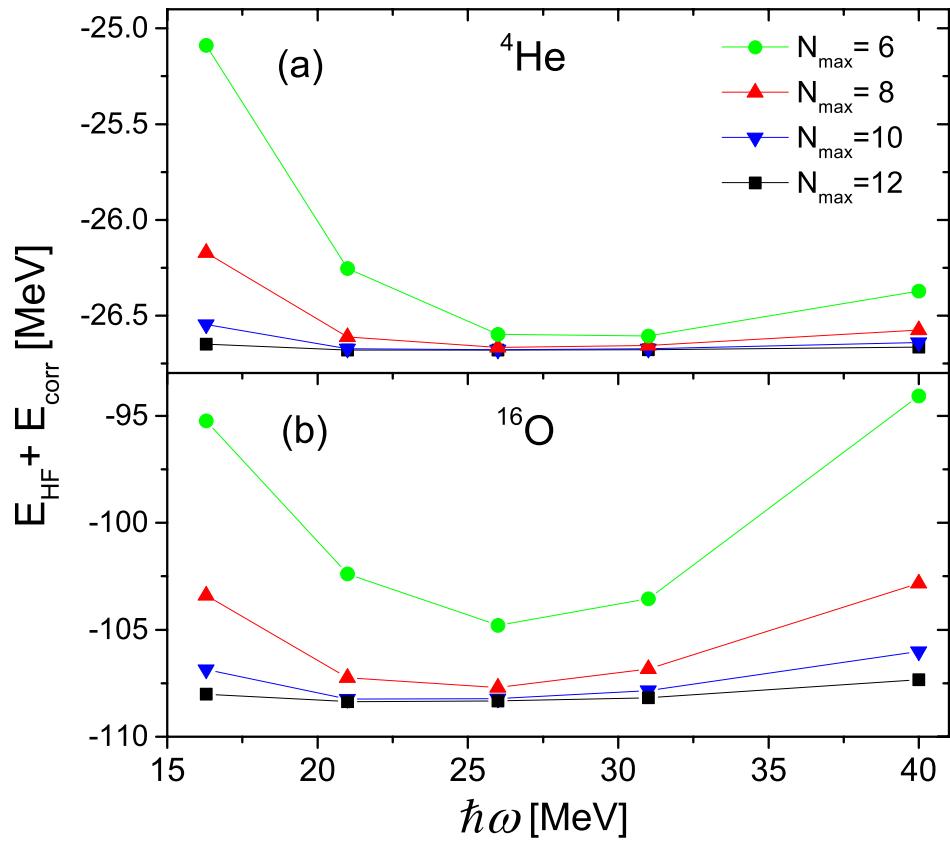
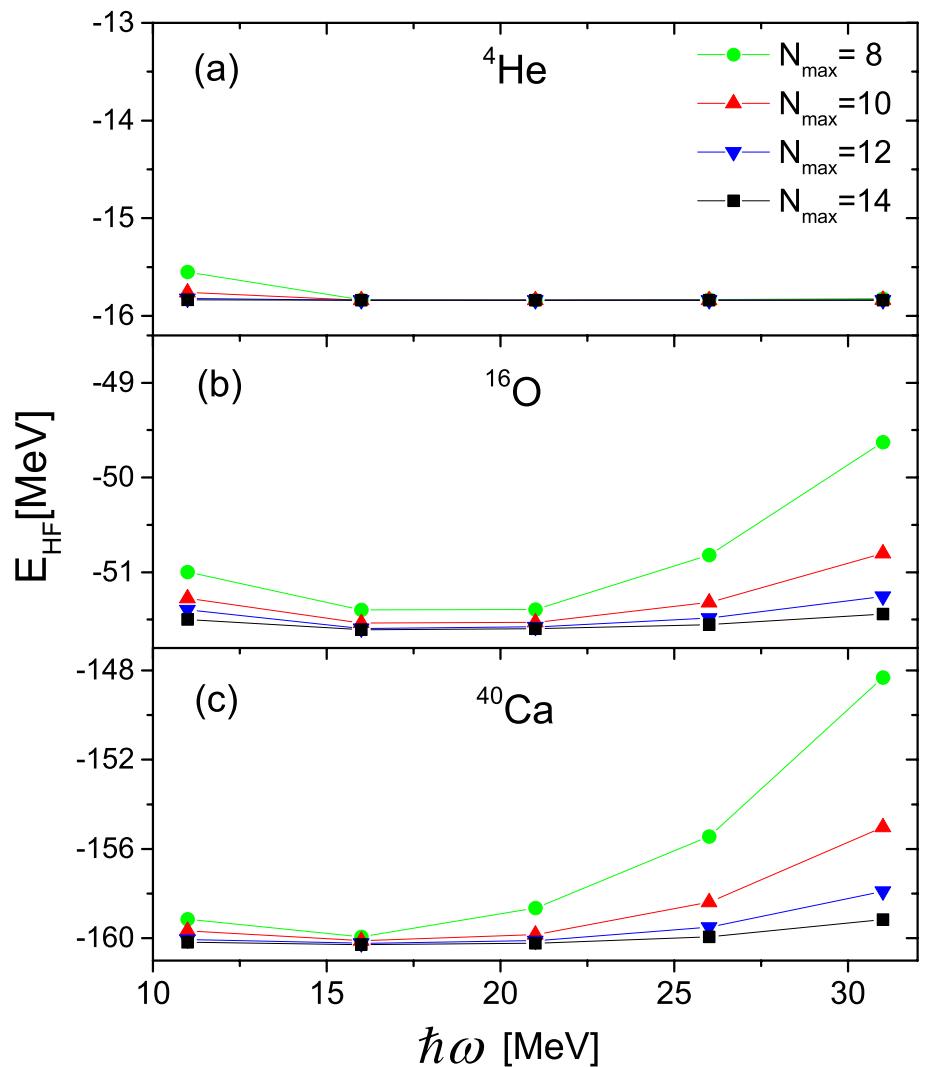
Ground state correlations



Ground state correlations



Ground state correlations



Ground state correlations

