

Magnetic Susceptibility of neutron stars from QMC

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In collaboration with:

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Outline

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- (i) Overview.
- (ii) Introduction.
- (iii) Previous results on the magnetic susceptibility.
- (iv) Magnetic susceptibility.
- (v) Conclusions.

Overview

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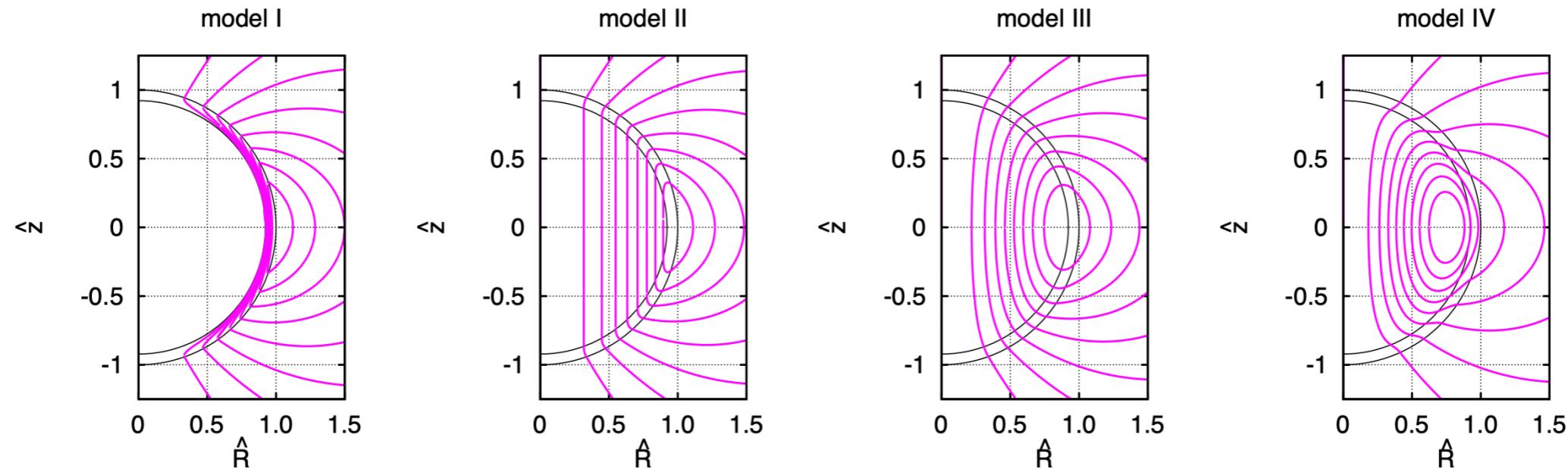
- ♦ Neutron star physics and presence of strong magnetic fields (magnetars) => spin polarization could play a role.
- ♦ We focus on densities typical of the outer core of the neutron stars.
- ♦ Use *ab initio* calculations to study ground state properties and compute magnetic susceptibility.
- ♦ Implications in the study of supernovae and protoneutrons stars.

Image by Jurik Peter | Shutterstock

Why are we doing this?

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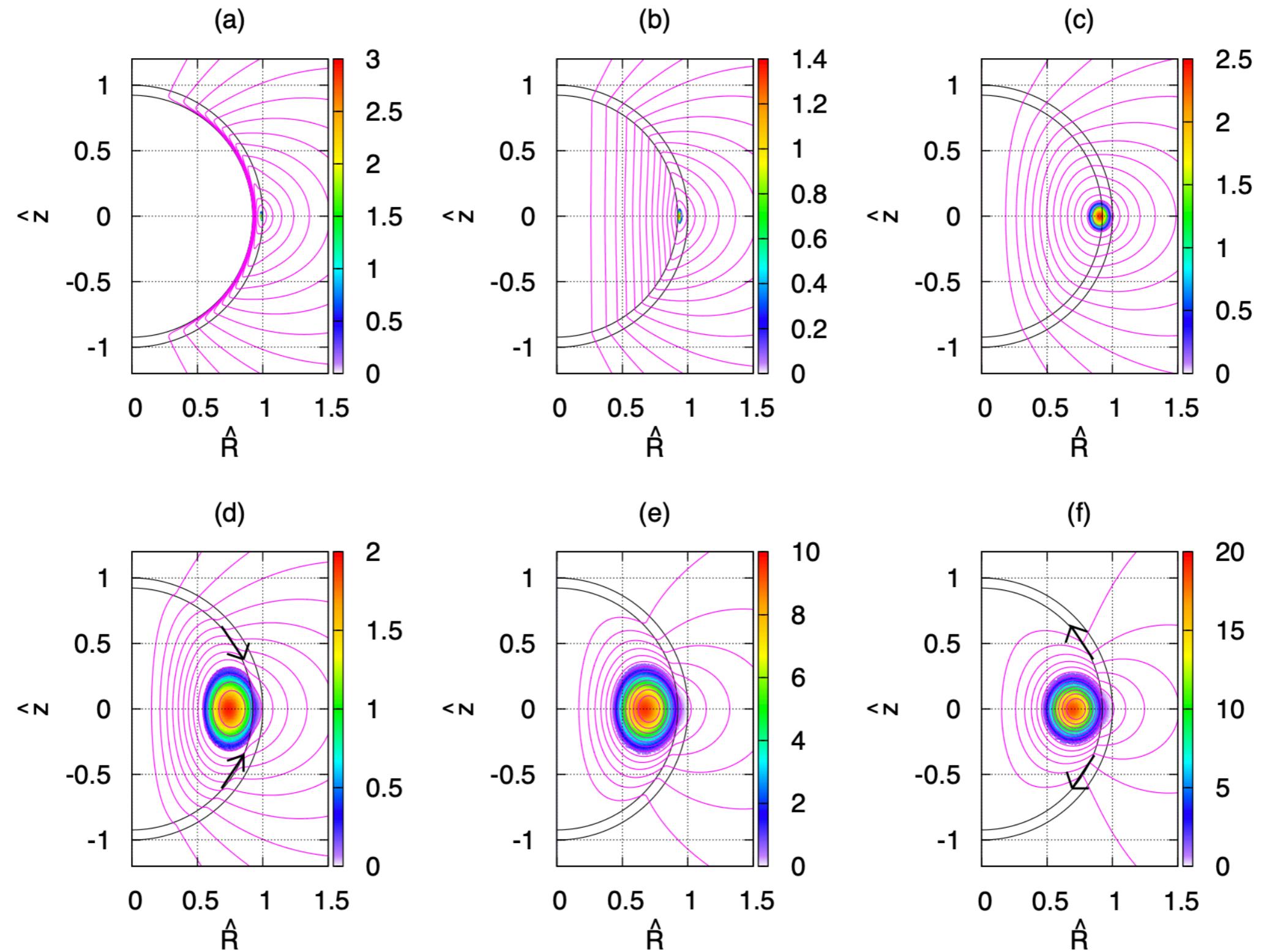
- Magnetars $\sim 10\%$ of young NS, with high magnetic field observed on the surface ($\sim 10^{14}/10^{15}$ G) \Rightarrow huge inferred internal magnetic field.



K. Fujisawa *et al.*, Mon. Not. 445, 2777–2793 (2014)

- Calculate magnetic susceptibility to include in simulation of NS mergers and supernovae explosions.

Why are we doing this?



Nuclear Hamiltonian

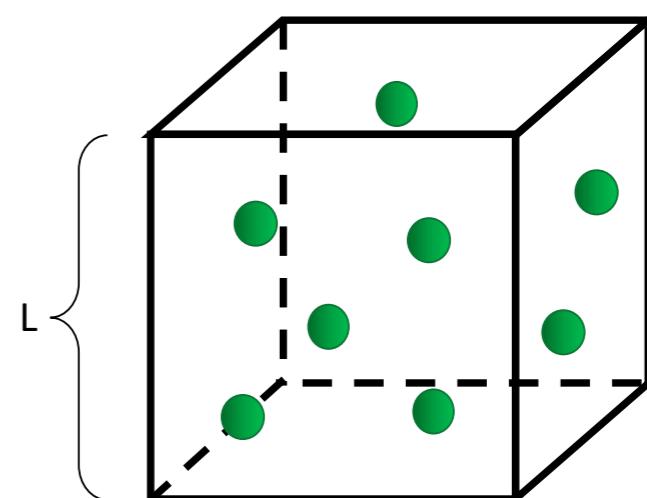
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$$H = T + \sum_{i < j} v_{ij} + \sum_{i < j < k} v_{ijk} + \dots$$

non-relativistic system with effective nucleon-nucleon (NN) force and three nucleon interaction (NNN).

Infinite matter \Leftrightarrow N particles in a box with PBC

Approximation: only neutrons.



Magnetic Susceptibility

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$$H = H_0 - \sum_i \vec{\sigma}_i \cdot \vec{b}$$

where we applied an external magnetic field $\vec{b} = \mu \vec{B}$ to the system.

$$\chi = -\rho \mu^2 \left. \frac{\partial^2 E_0(b)}{\partial b^2} \right|_{b=0}$$

Following approximation on S. Fantoni *et al.*, PRL 87, 181101 (2001).

Magnetic Susceptibility

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Using Pauli expansion in spin polarization: $\xi = - \left. \frac{\partial E_0(b)}{\partial b} \right|_{b=0}$

$$E(\xi) = E(0) - b\xi + \frac{1}{2}\xi^2 E''(0)$$

Minimizing the energy with respect to the spin polarization we obtain:

$$\chi = \mu^2 \rho \frac{1}{E''(0)}$$

Magnetic Susceptibility

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How to compute $E''(0)$?

Using chain rule (J_z is the spin asymmetry, i.e. $N_\uparrow - N_\downarrow$):

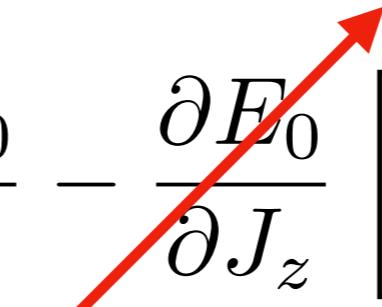
$$E''(0) = \left[\frac{\partial \xi}{\partial J_z} \right]^{-2} \left\{ \frac{\partial^2 E_0}{\partial J_z^2} - \frac{\partial E_0}{\partial J_z} \left[\frac{\partial \xi}{\partial J_z} \right]^{-1} \frac{\partial^2 \xi}{\partial J_z^2} \right\}$$

Magnetic Susceptibility

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Magnetic Susceptibility

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$$E''(0) = \left[\frac{\partial \xi}{\partial J_z} \right]^{-2} \frac{\partial^2 E_0}{\partial J_z^2}$$

Magnetic Susceptibility

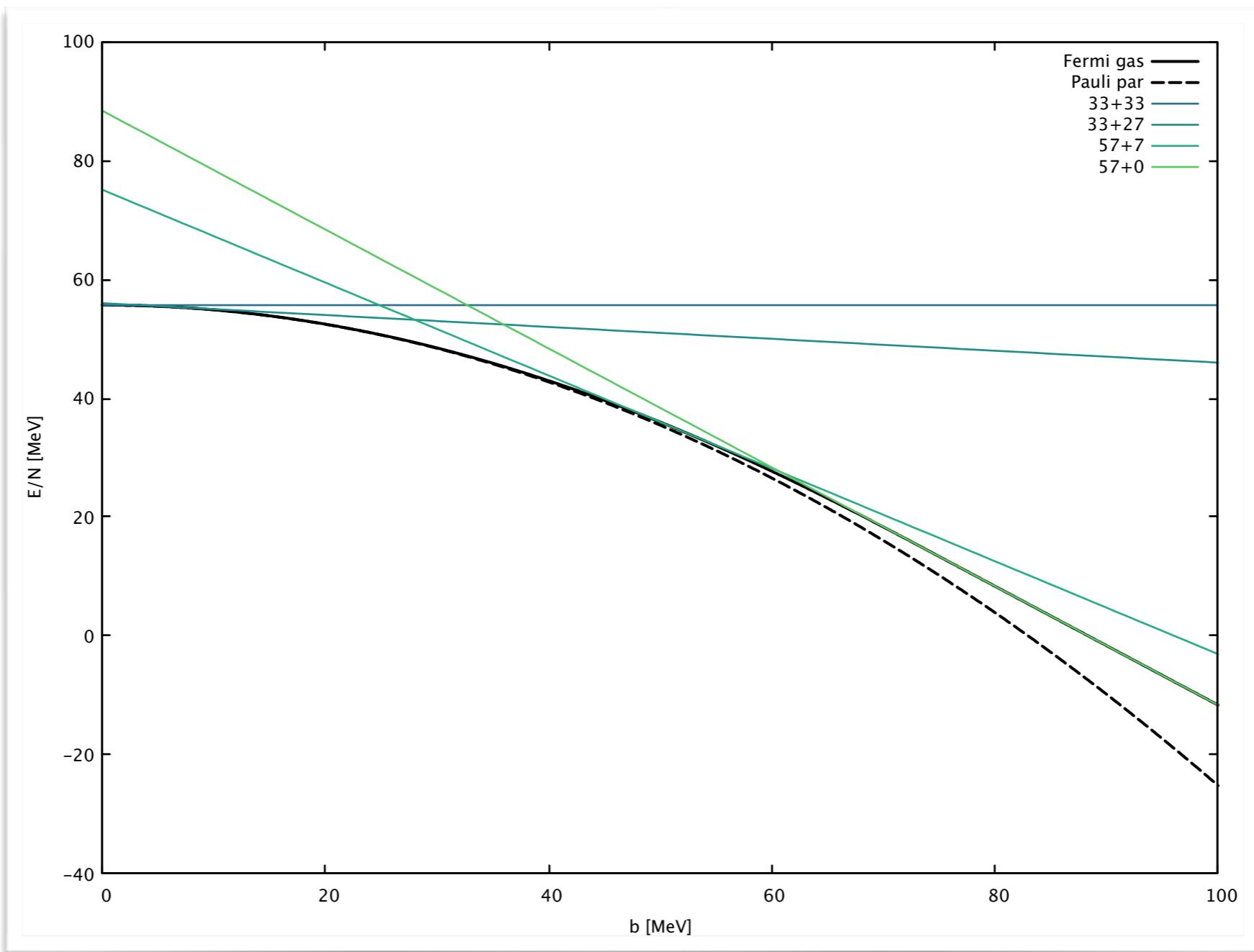
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Approximate expressions (exact for infinite system with low b and J_z):

$$\frac{\partial \xi}{\partial J_z} \approx \frac{E_0(J_z = J_{z0}, b = 0) - E_0(J_z = J_{z0}, b = b_0)}{J_{z0}b_0}$$

$$\frac{\partial^2 E_0}{\partial J_z^2} \approx 2 \frac{E_0(J_z = J_{z0}, b = 0) - E_0(J_z = 0, b = 0)}{J_{z0}^2}$$

Compute energies with AFDMC.



Magnetic Susceptibility

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This approximation is reasonable under these assumptions:

- (i) for $b = 0$, $E_0(J_z, b)$ is quadratic in J_z
- (ii) for a fixed J_z , $E_0(J_z, b)$ is linear in b
- (iii) the polarization is linear in J_z

Magnetic Susceptibility

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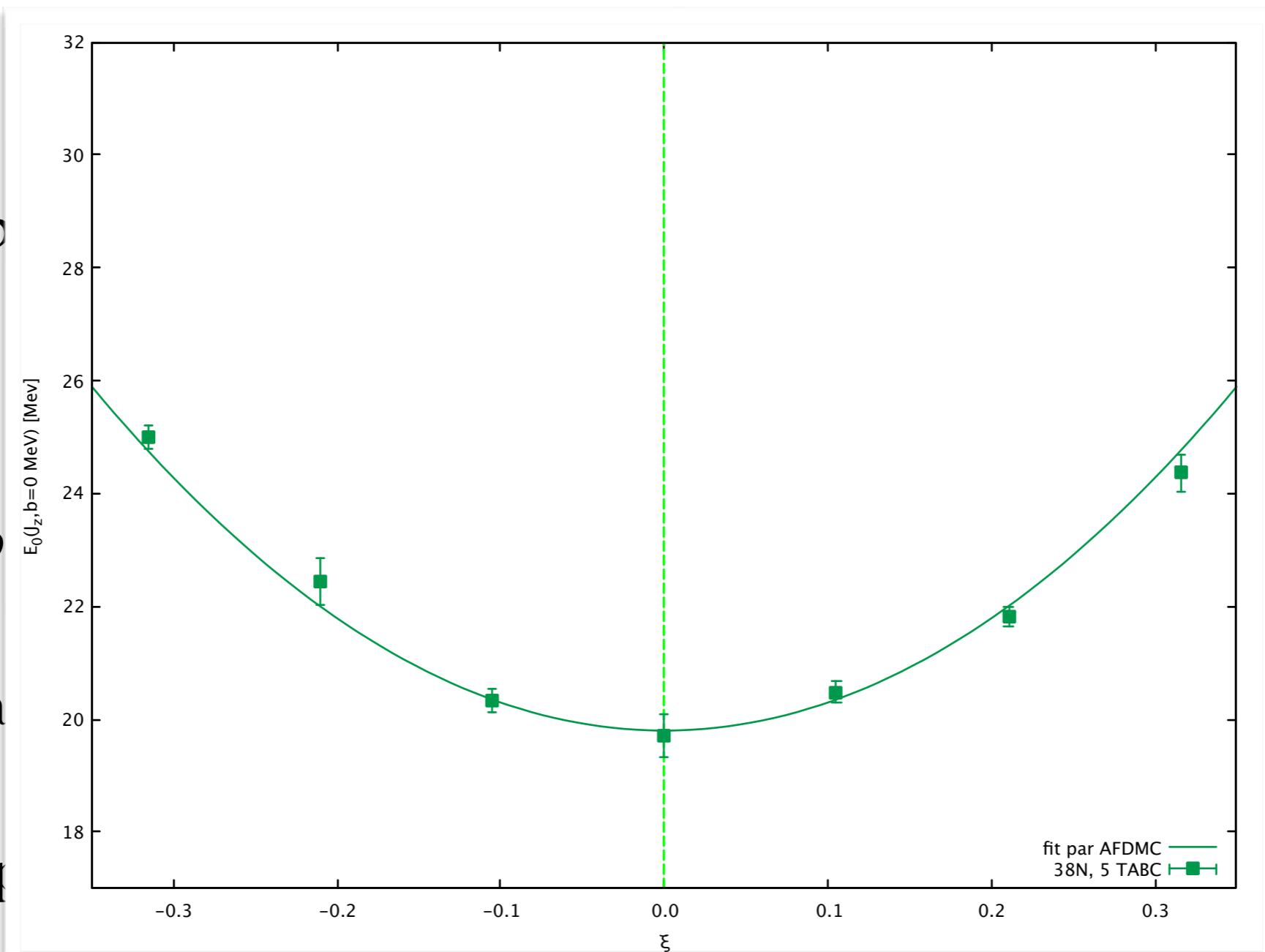
This app

ions:

(i) for b

(ii) for a

(iii) the p



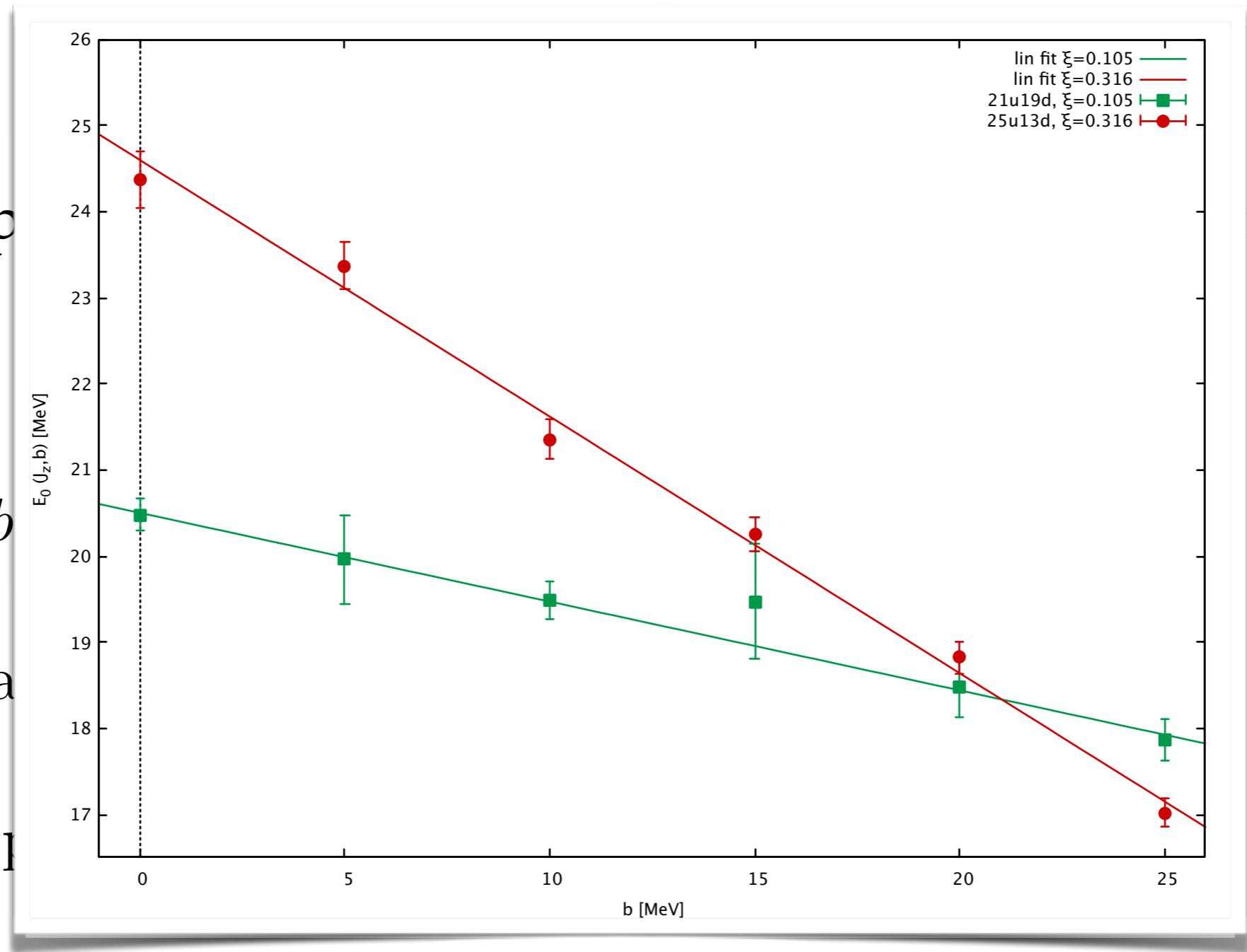
Magnetic Susceptibility

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This app

ions:

- (i) for b
- (ii) for a
- (iii) the 1



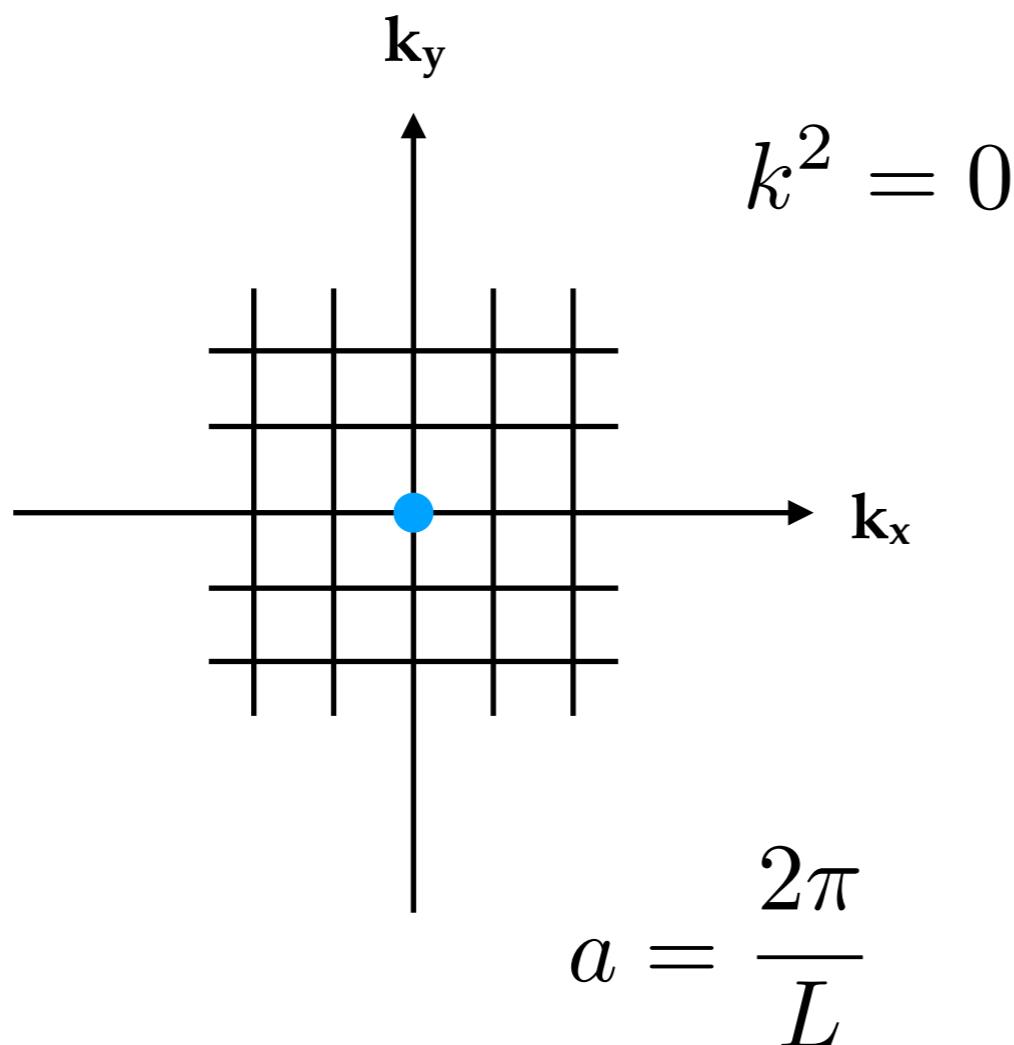
Beyond PBC: TABC

In collaboration with:

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◆ Diego Lonardoni @MSU/LANL

2D



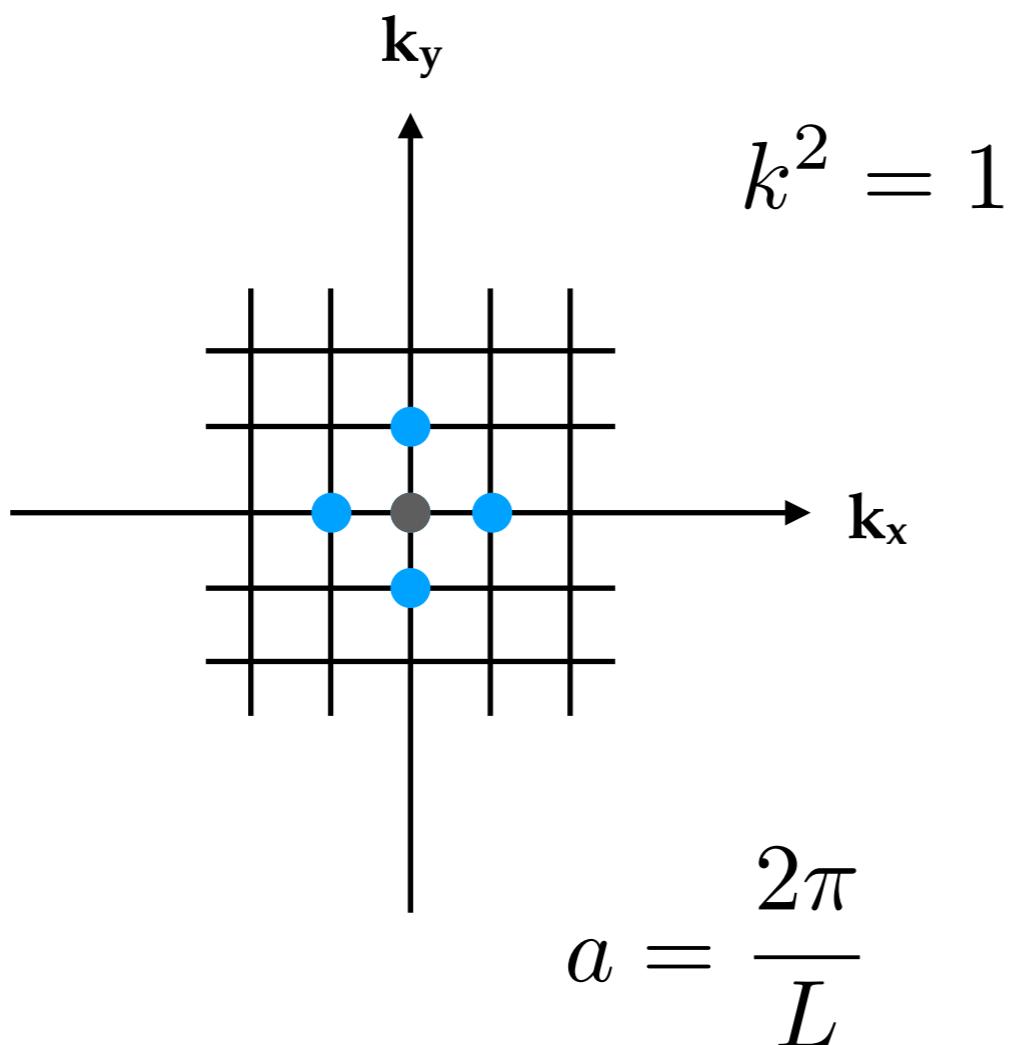
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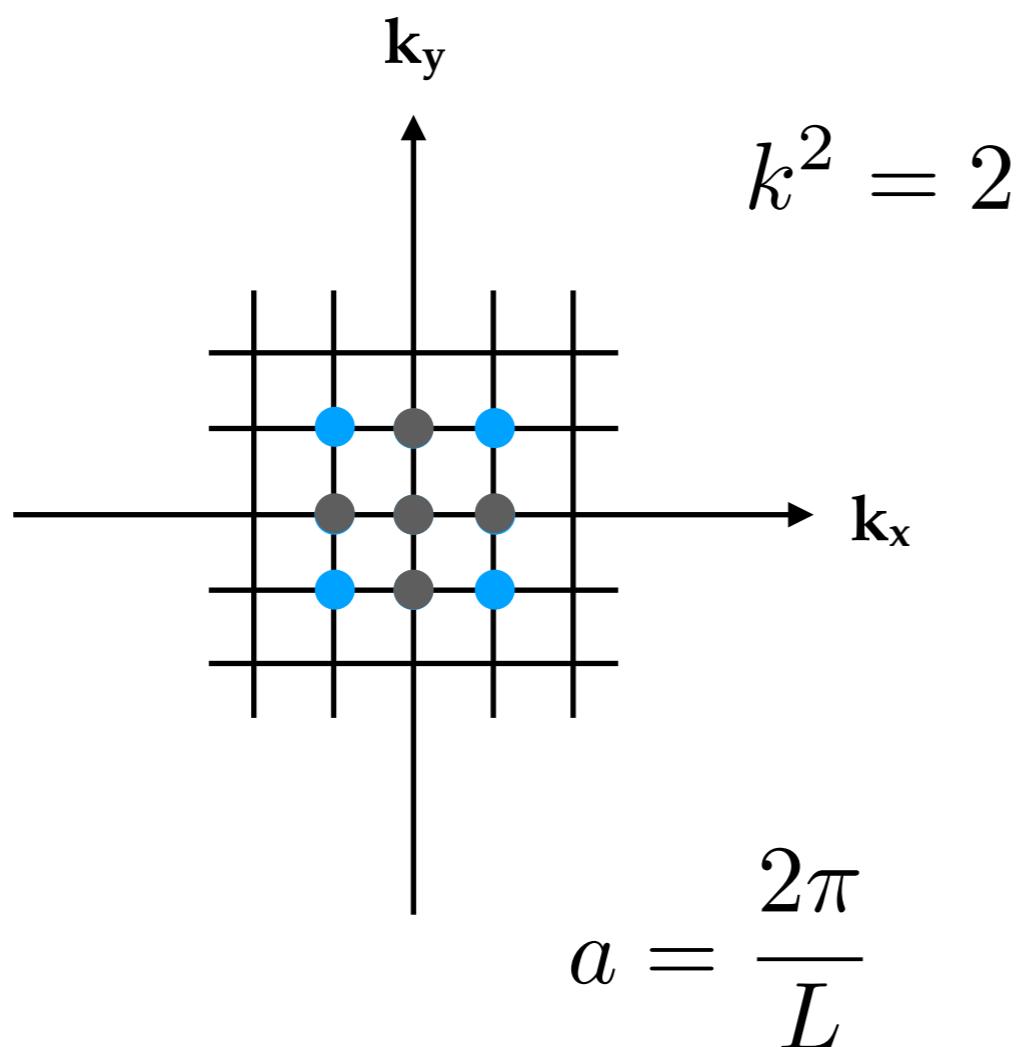
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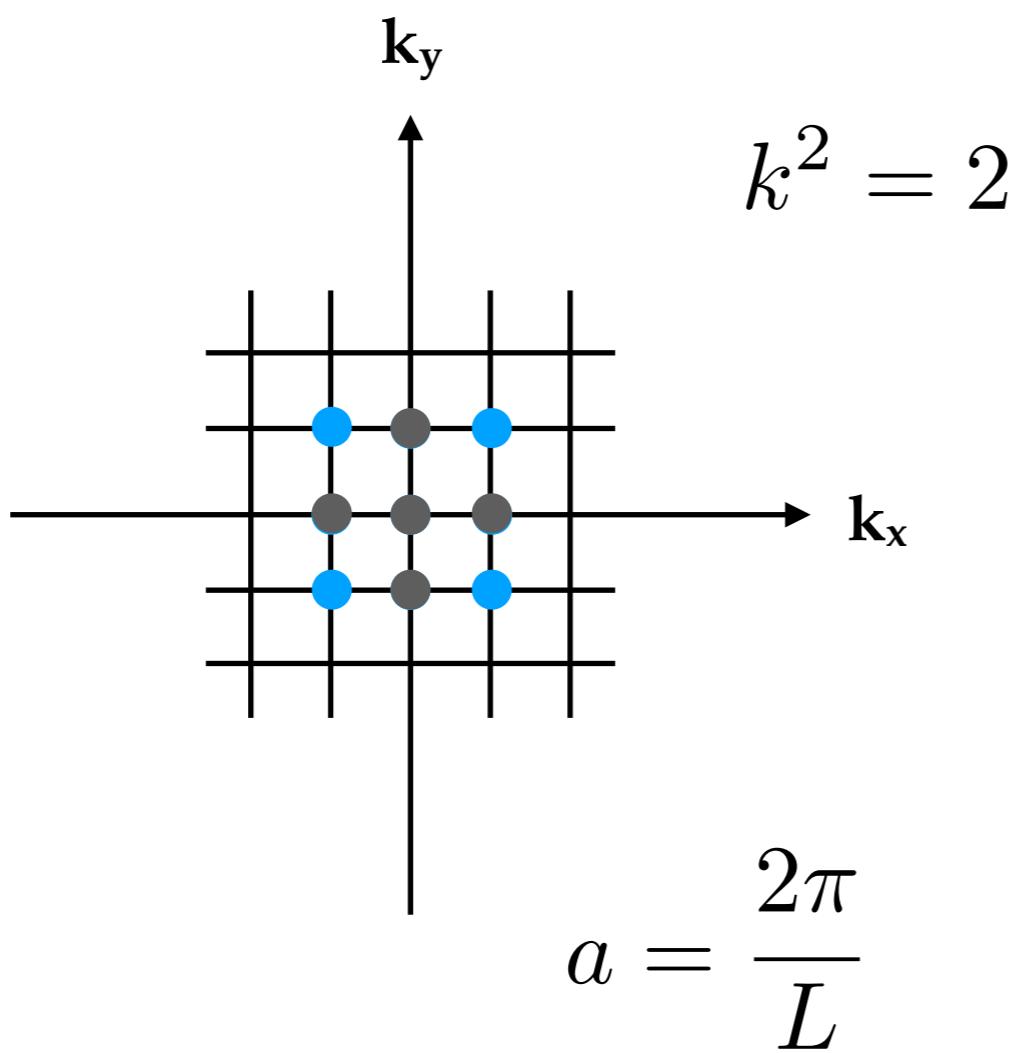
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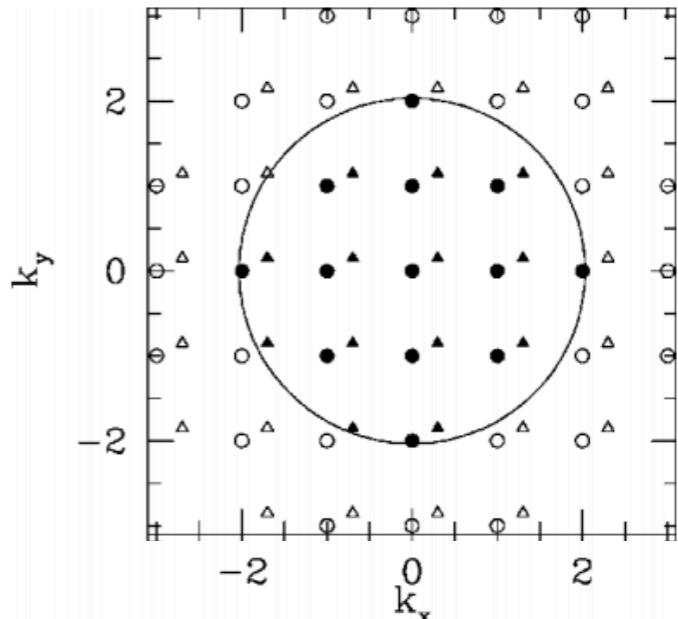
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2D



$$\theta_{TW} = \eta$$

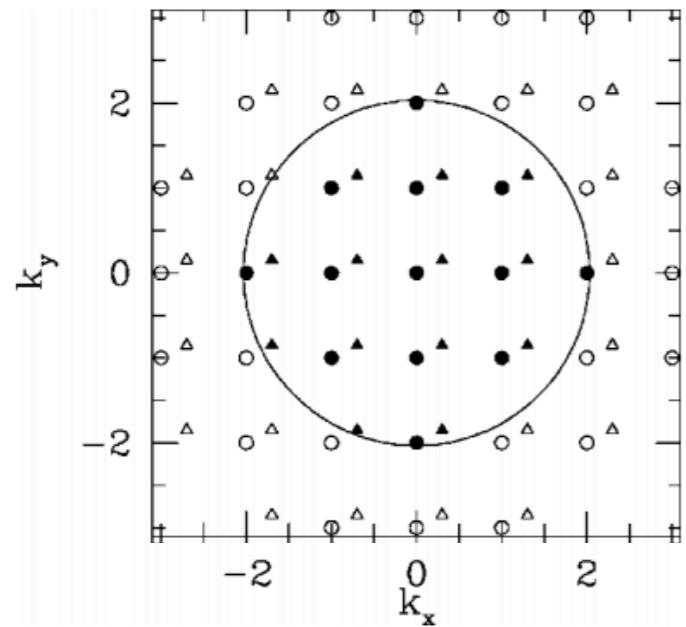
$$\eta \in (-0.5, 0.5)$$



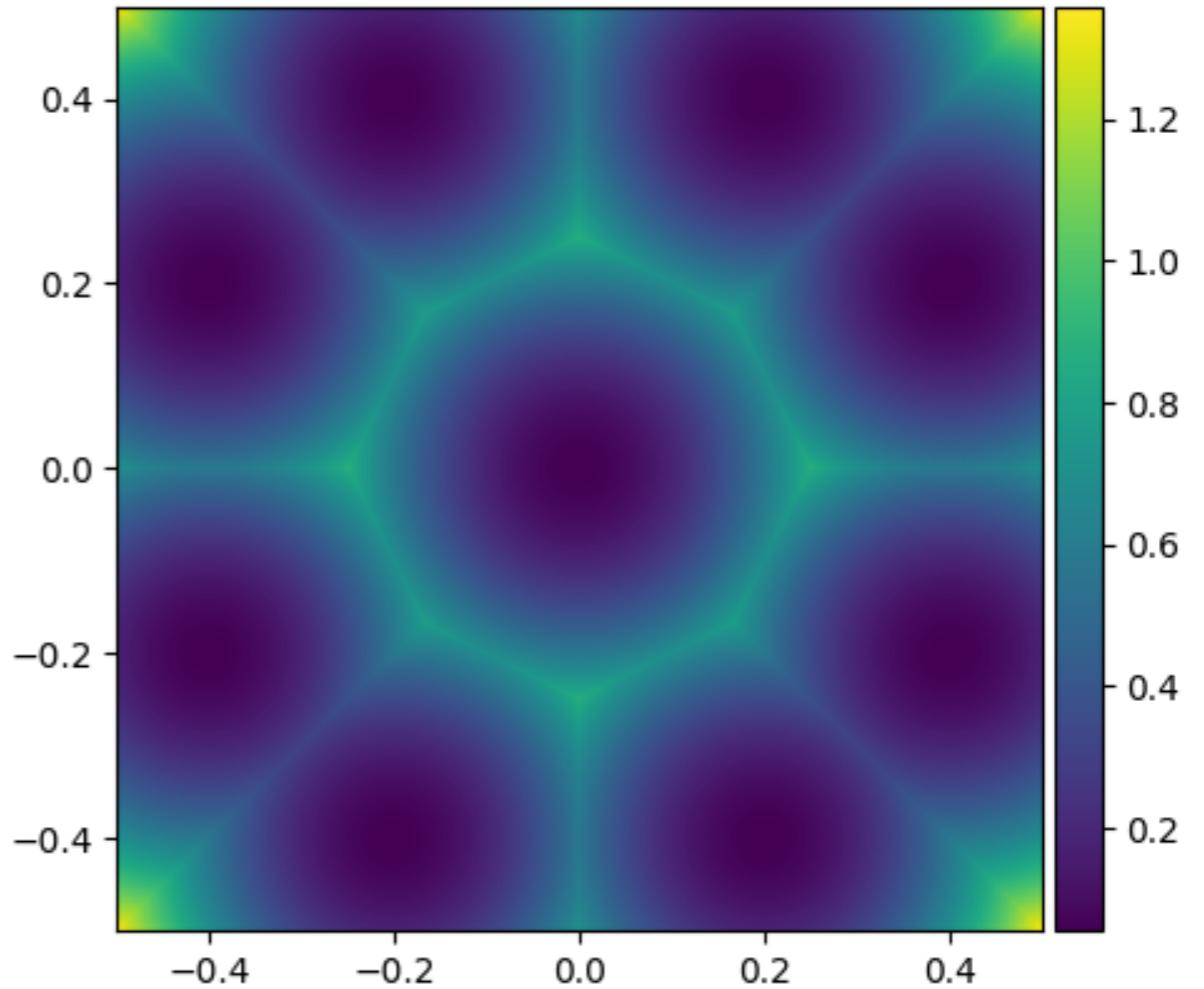
C. Lin *et al.*, Phys. Rev. E **64**, 016702 (2001)

Beyond PBC: TABC

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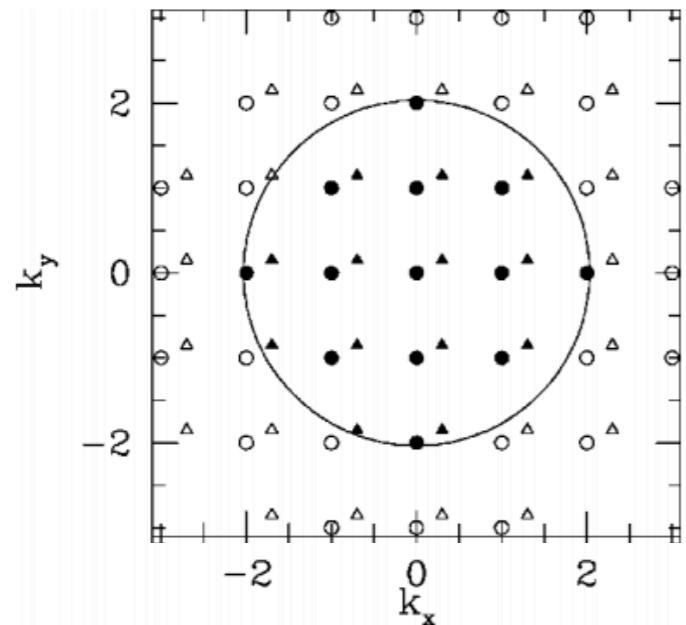
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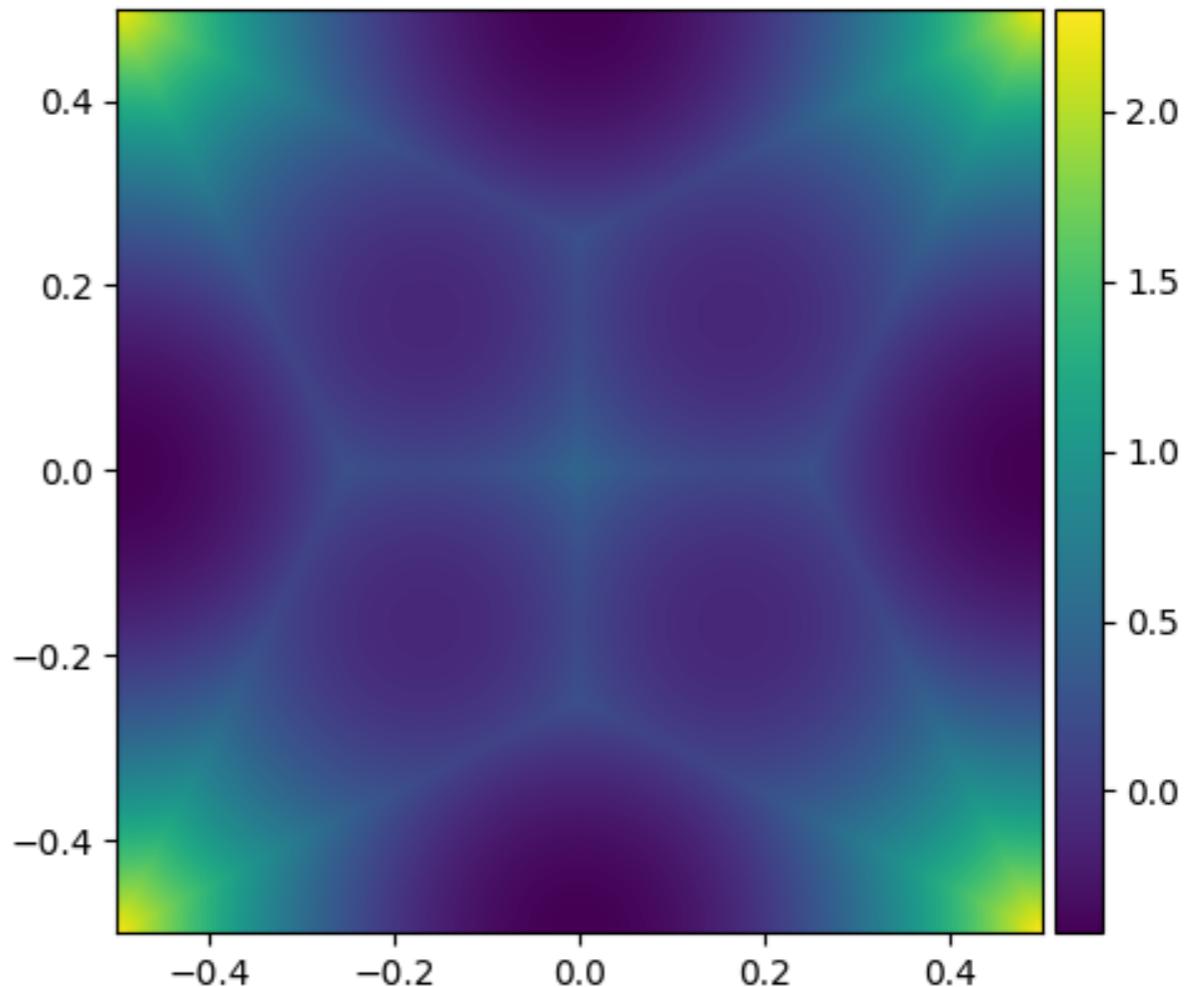
10 particles

Beyond PBC: TABC

14



C. Lin *et al.*, Phys. Rev. E **64**, 016702 (2001)

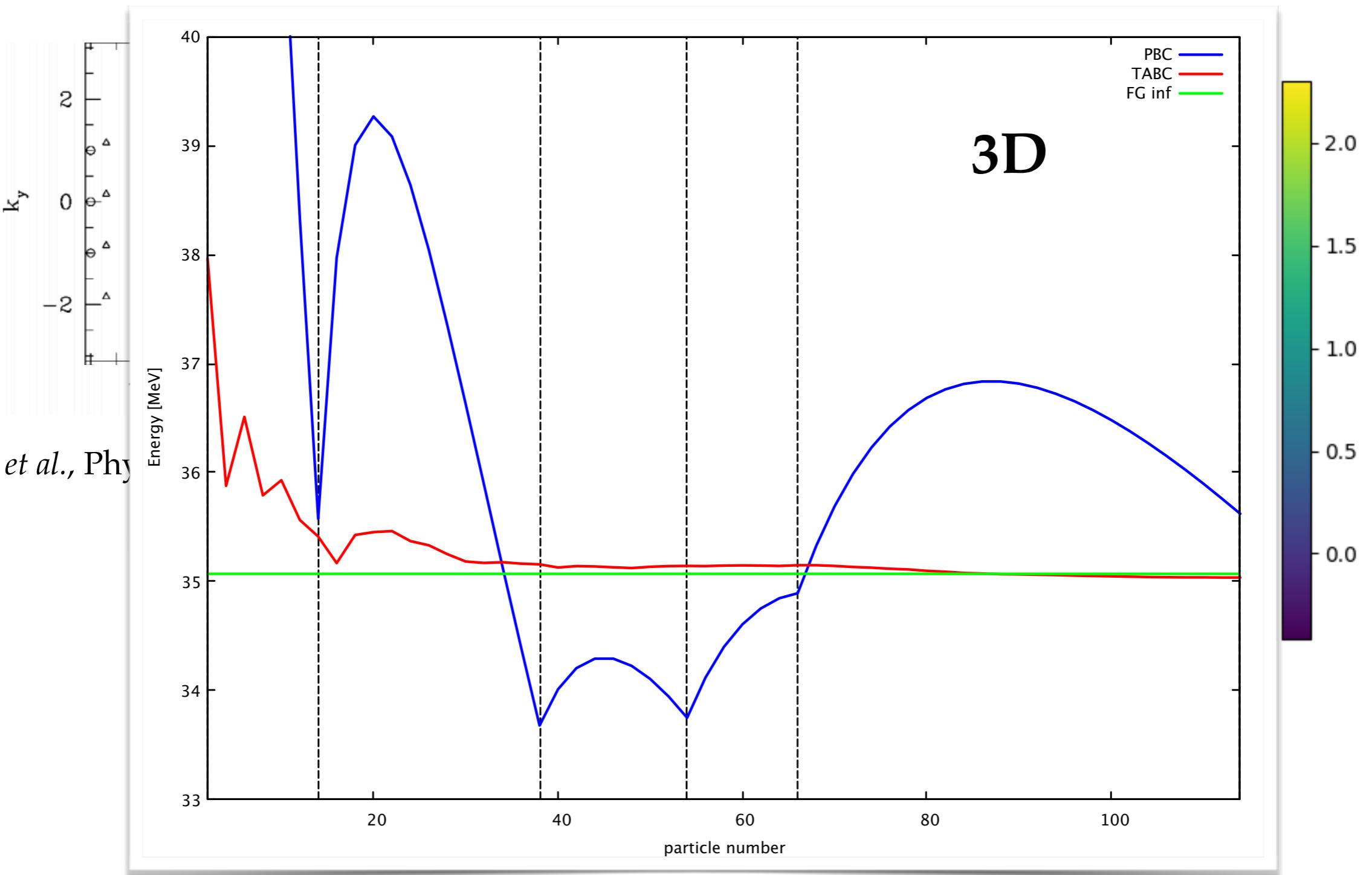


16 particles

Beyond PBC: TABC

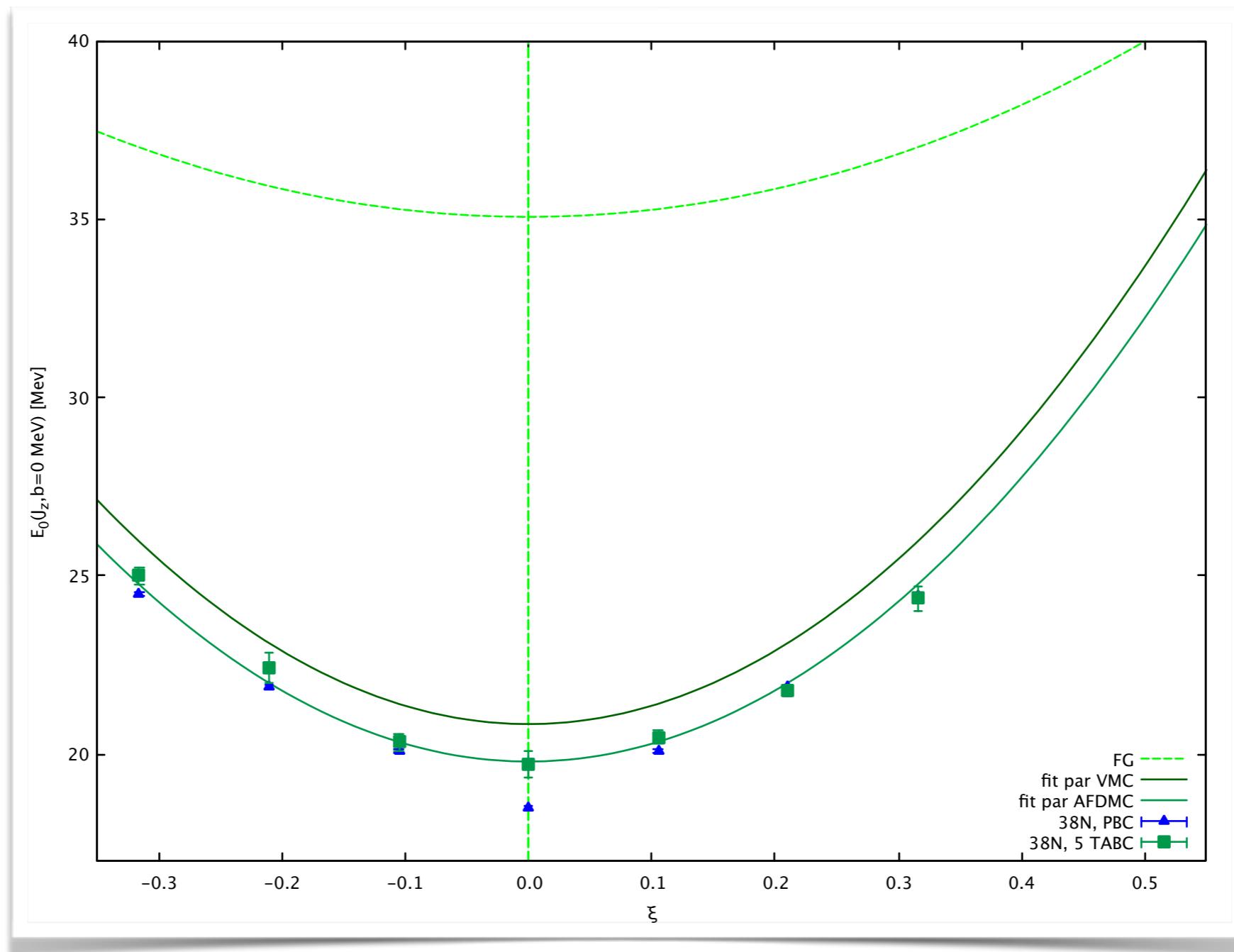
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C. Lin *et al.*, Phy



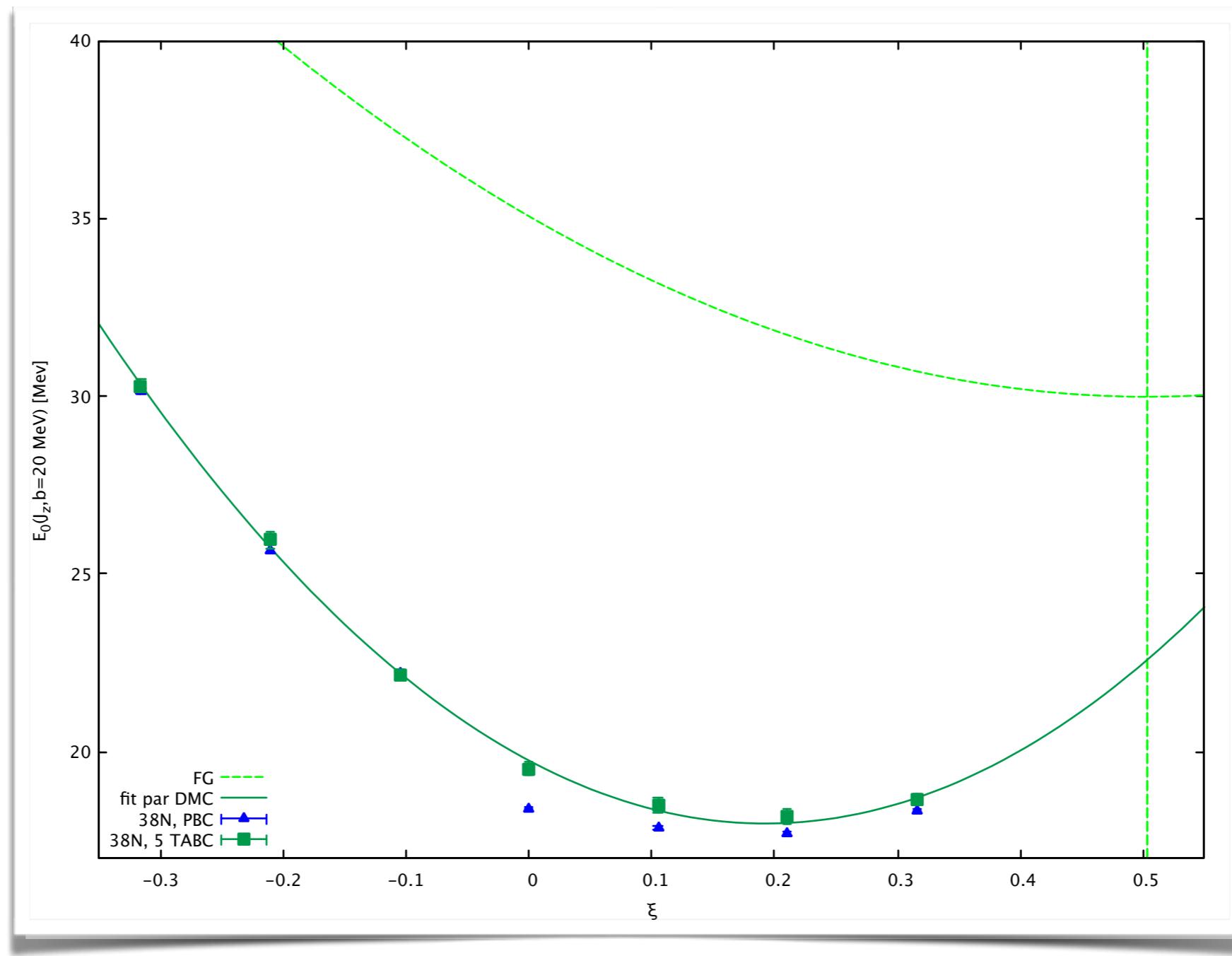
What can we do with TABC?

15



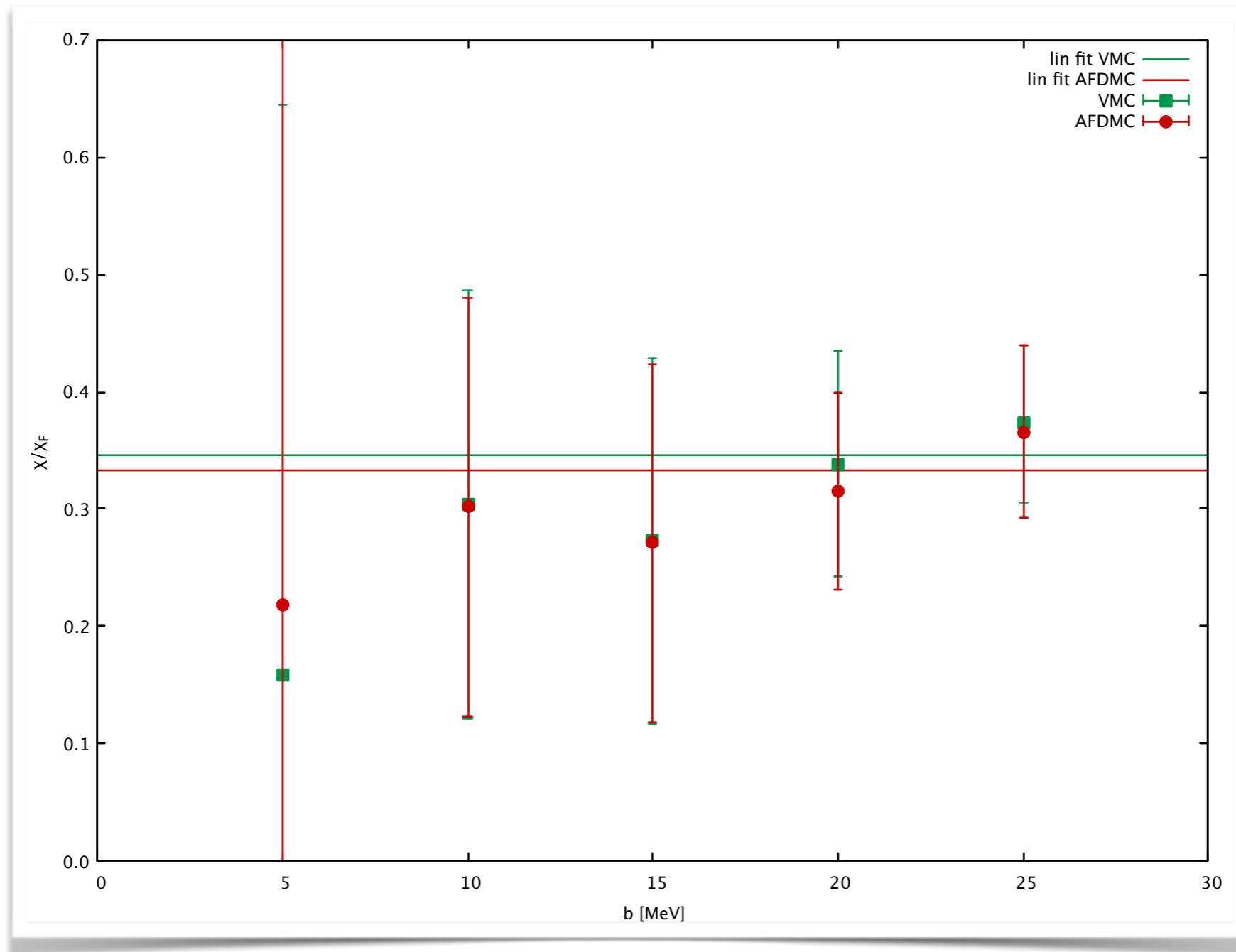
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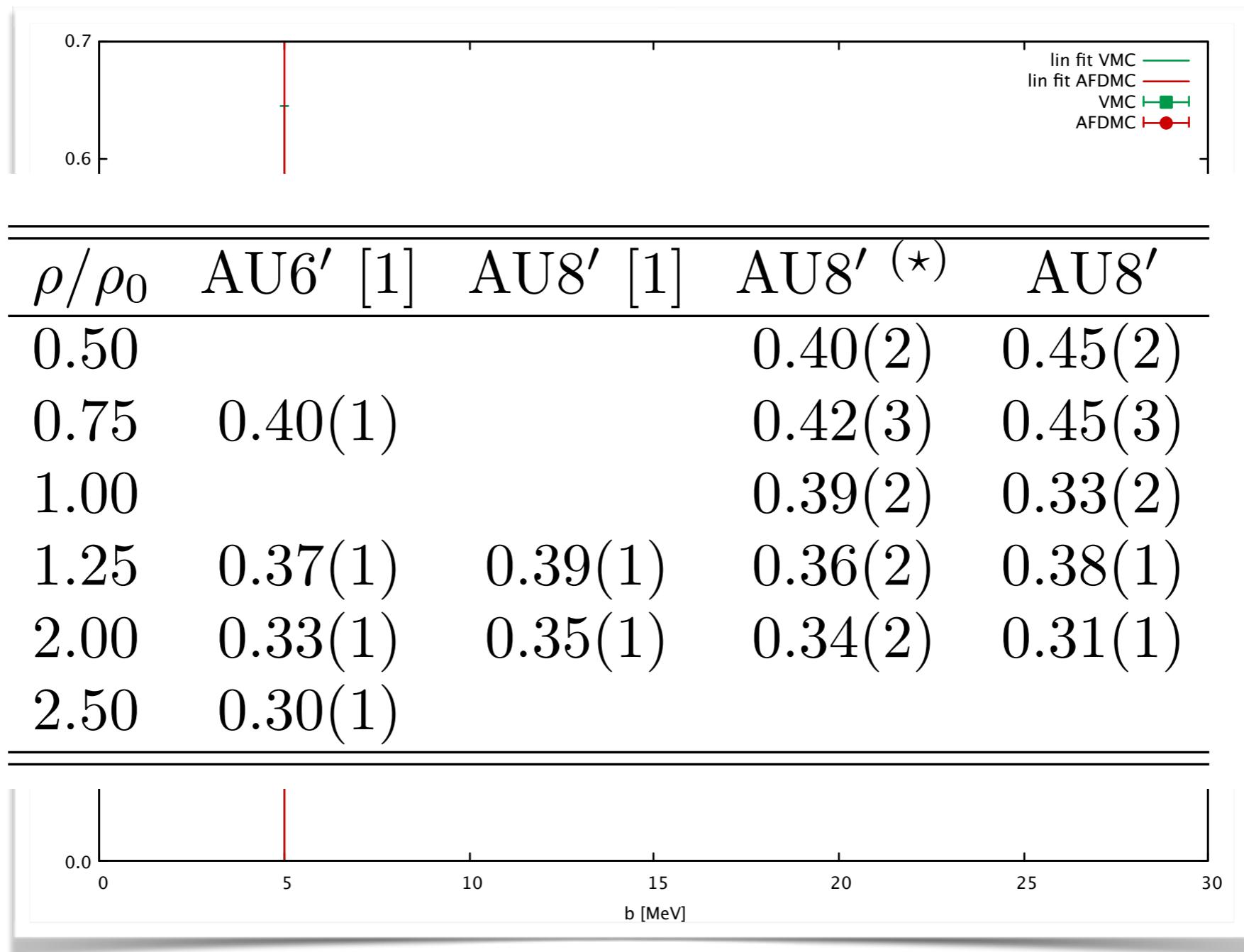
Results using approximation

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Results using approximation

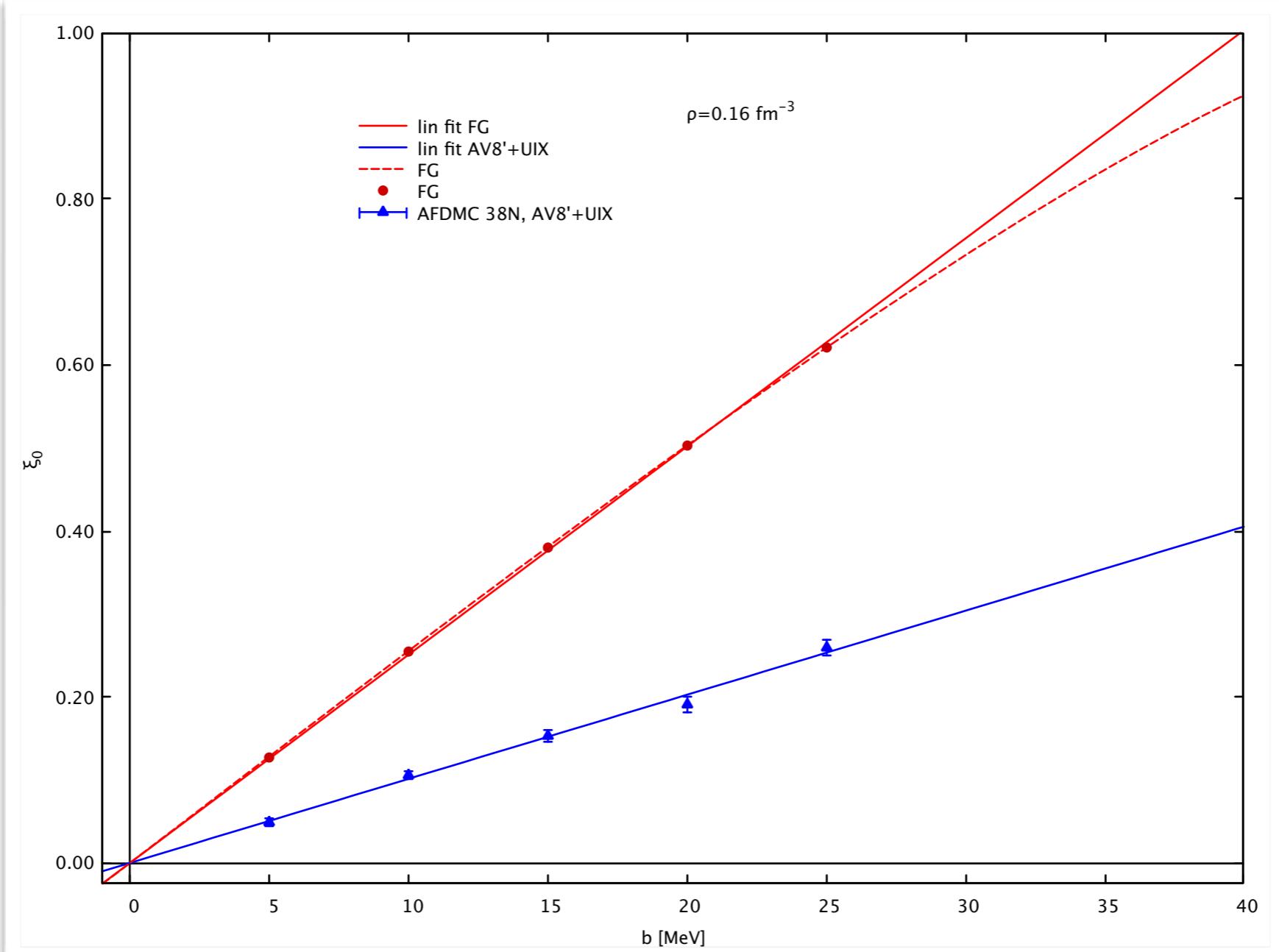
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We can do better

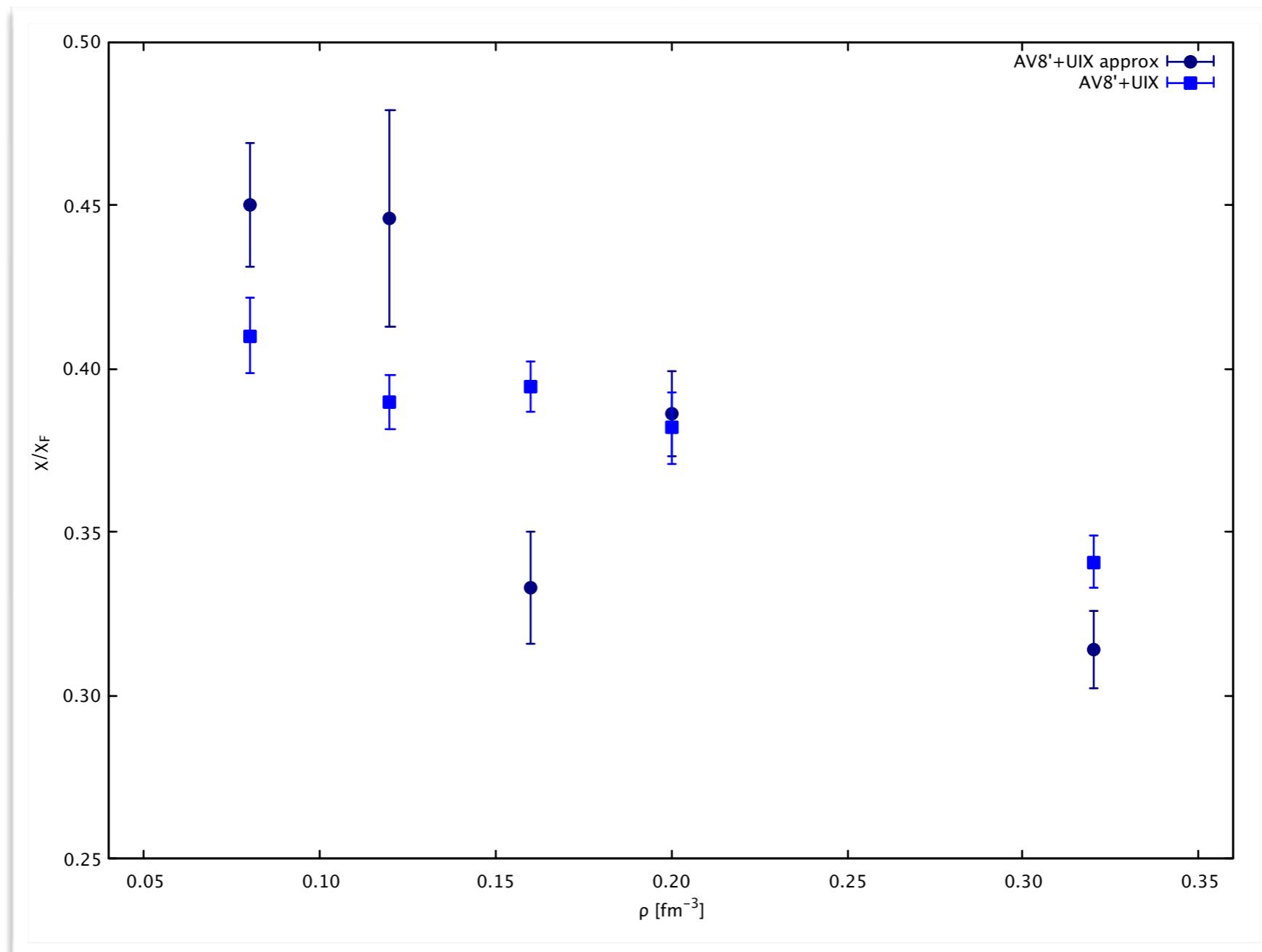
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$$\chi = -\rho\mu^2 \frac{\partial^2 E_0(b)}{\partial b^2} \Big|_{b=0} \quad \xi = -\left. \frac{\partial E_0(b)}{\partial b} \right|_{b=0} \rightarrow \chi = \rho\mu^2 \frac{\partial \xi}{\partial b}$$



Magnetic susceptibility ratio

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Nuclear Potentials

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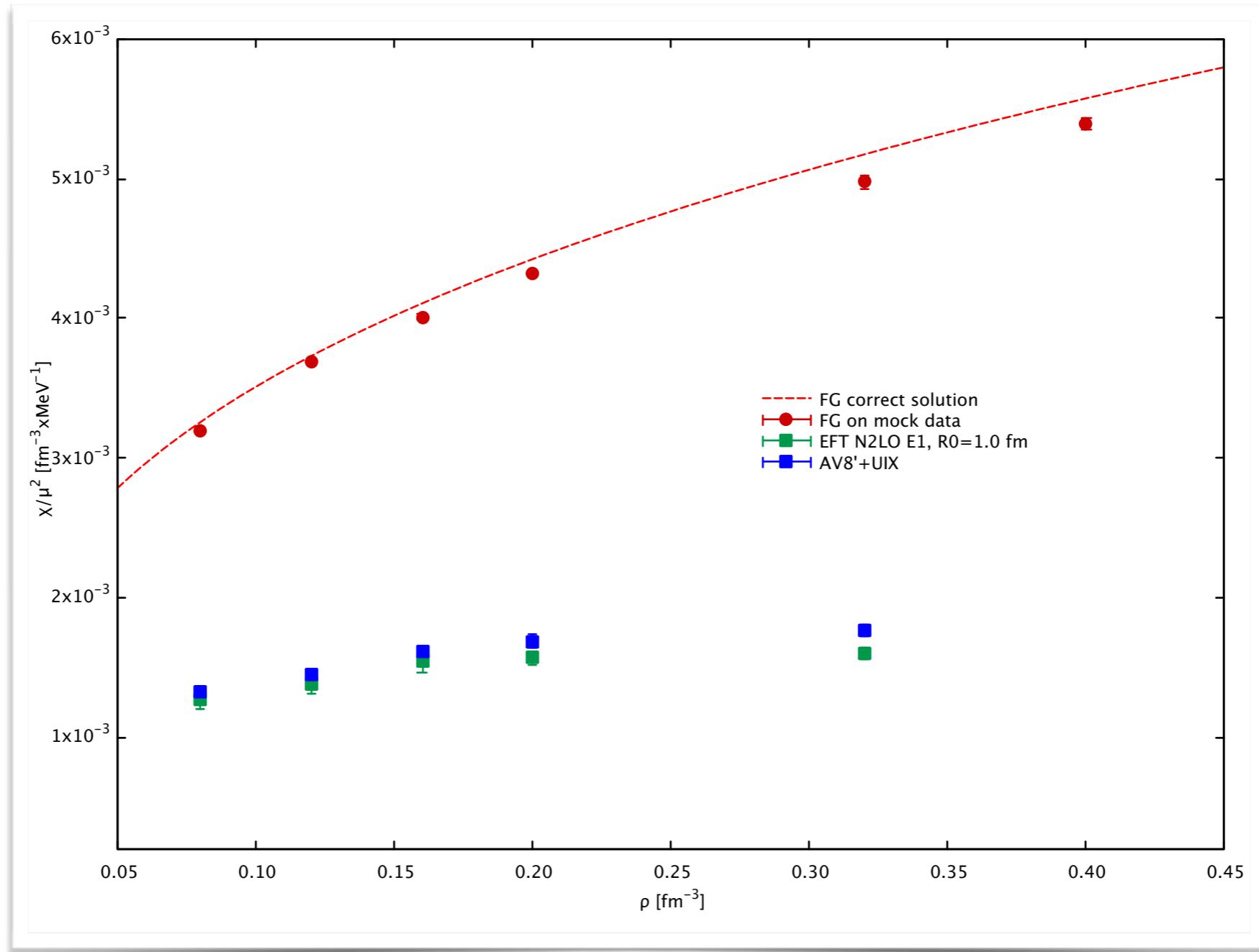
Two kinds of potential:

- Phenomenological AV8'+UIX.
- Local Chiral Effective Field Theory at N2LO
[Lynn et al. PRL 116, 062501 (2016)].

Ground State energies computed by means of Auxiliary Field Diffusion Monte Carlo (AFDMC).

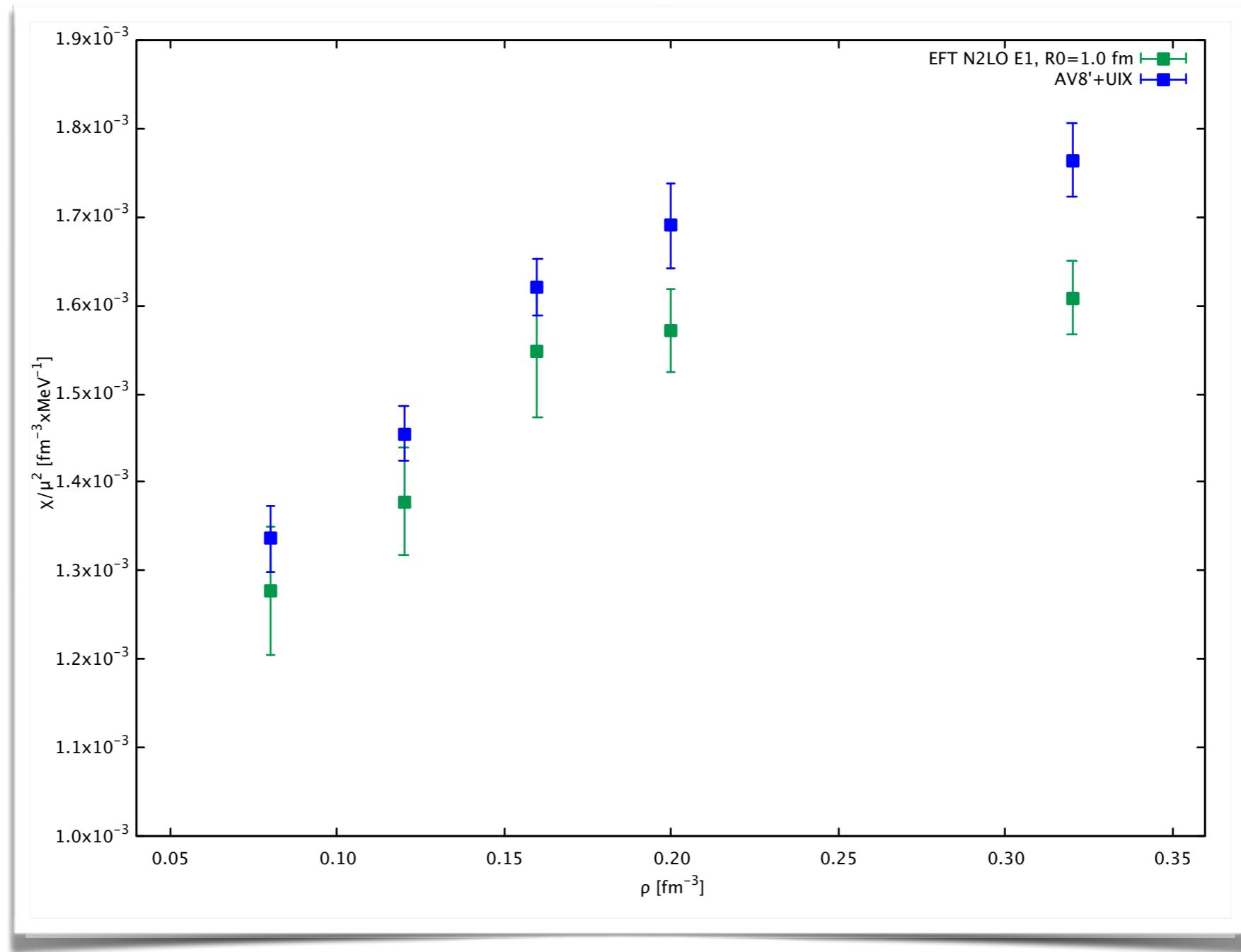
Magnetic susceptibility

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Magnetic susceptibility

20



Simple model

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Assuming an energy-density function of the kind:

$$E(\rho, \xi) = E_0(\rho) + \xi^2 (E_1(\rho) - E_0(\rho))$$

And recalling that the magnetic susceptibility can be computed as:

$$\chi = \mu^2 \rho \frac{1}{E''(0)}$$

Simple model

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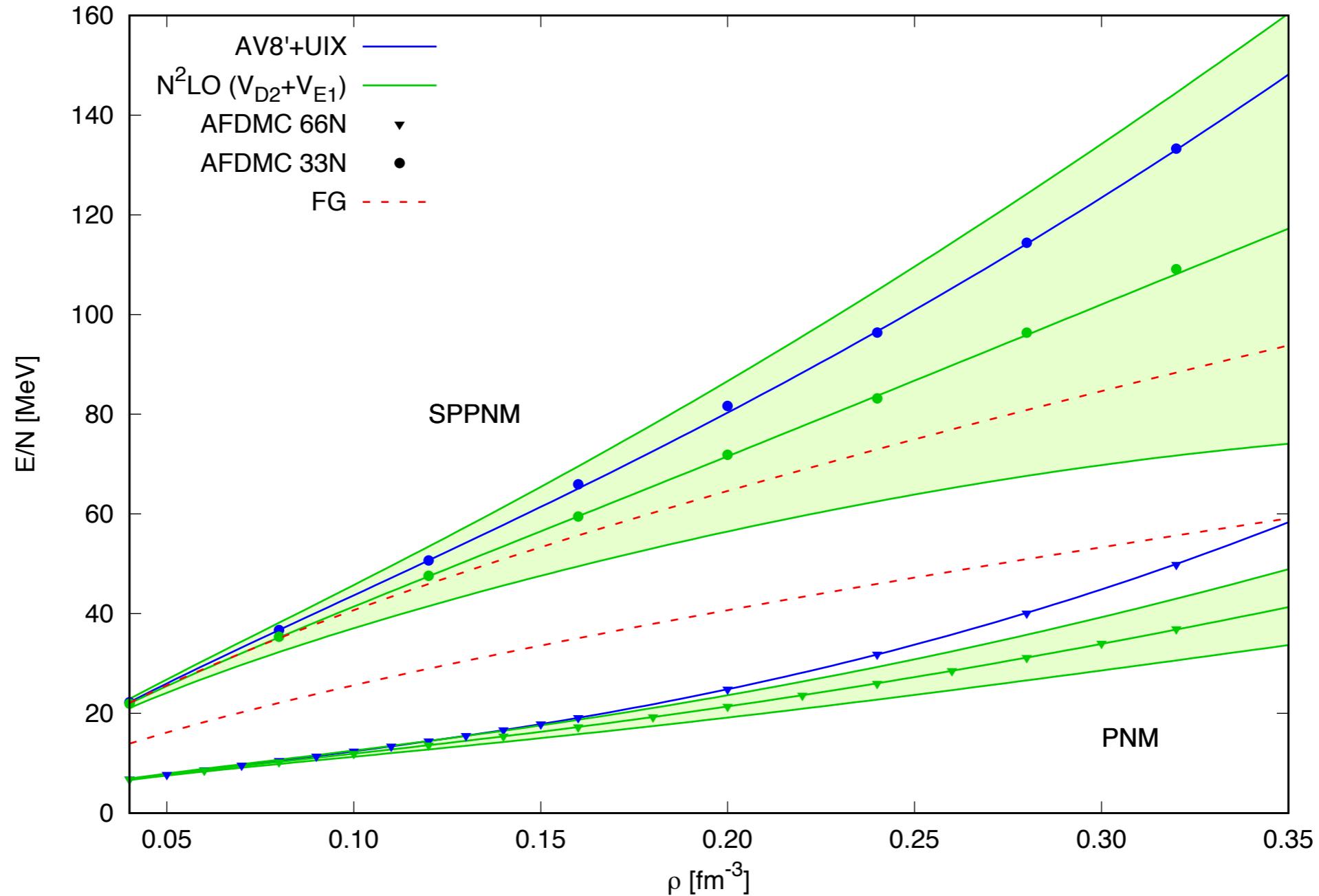
$$E(\rho, \xi) = E_0(\rho) + \xi^2 \underbrace{(E_1(\rho) - E_0(\rho))}_{SSE(\rho)}$$

And recalling that the magnetic susceptibility can be computed as:

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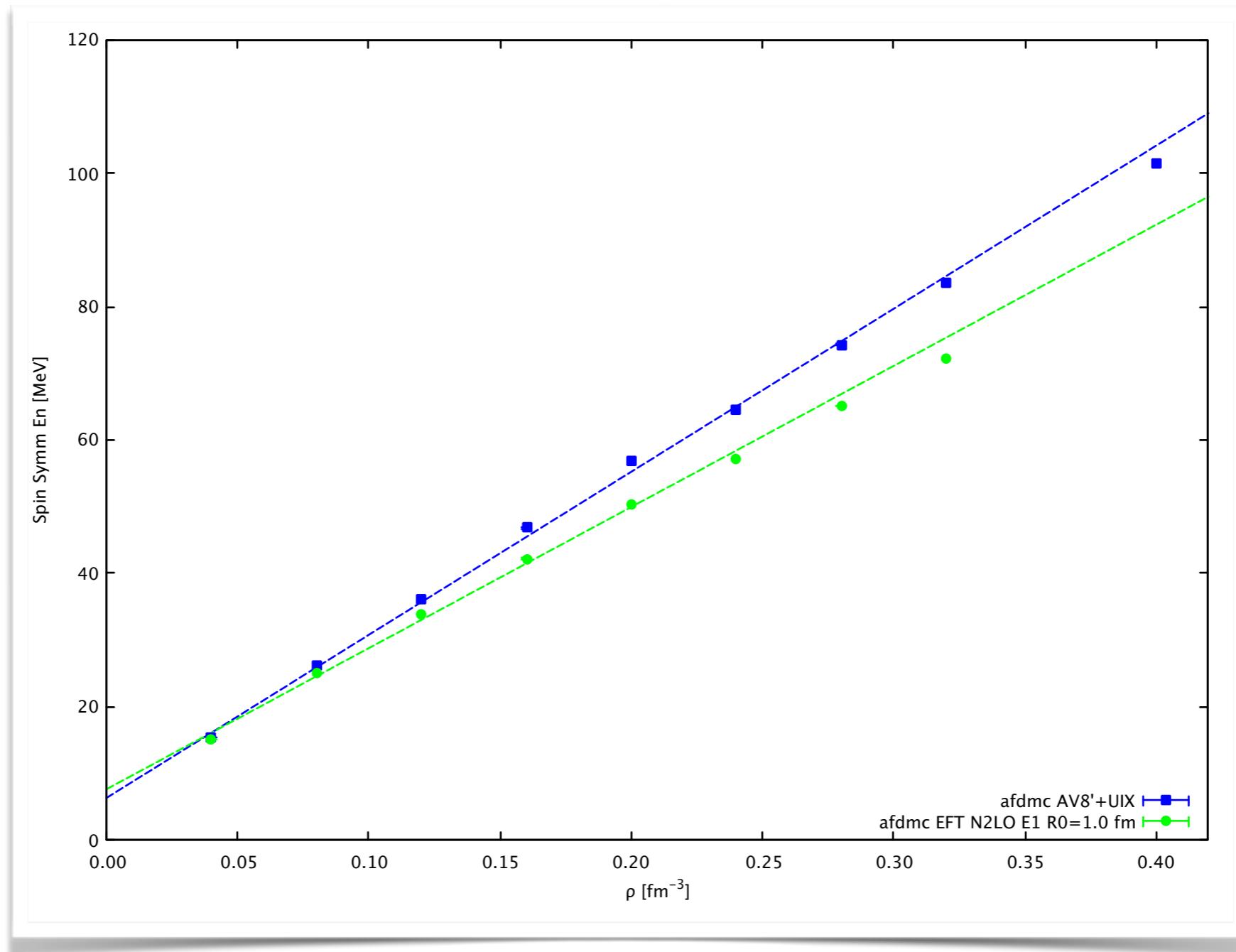
Simple model

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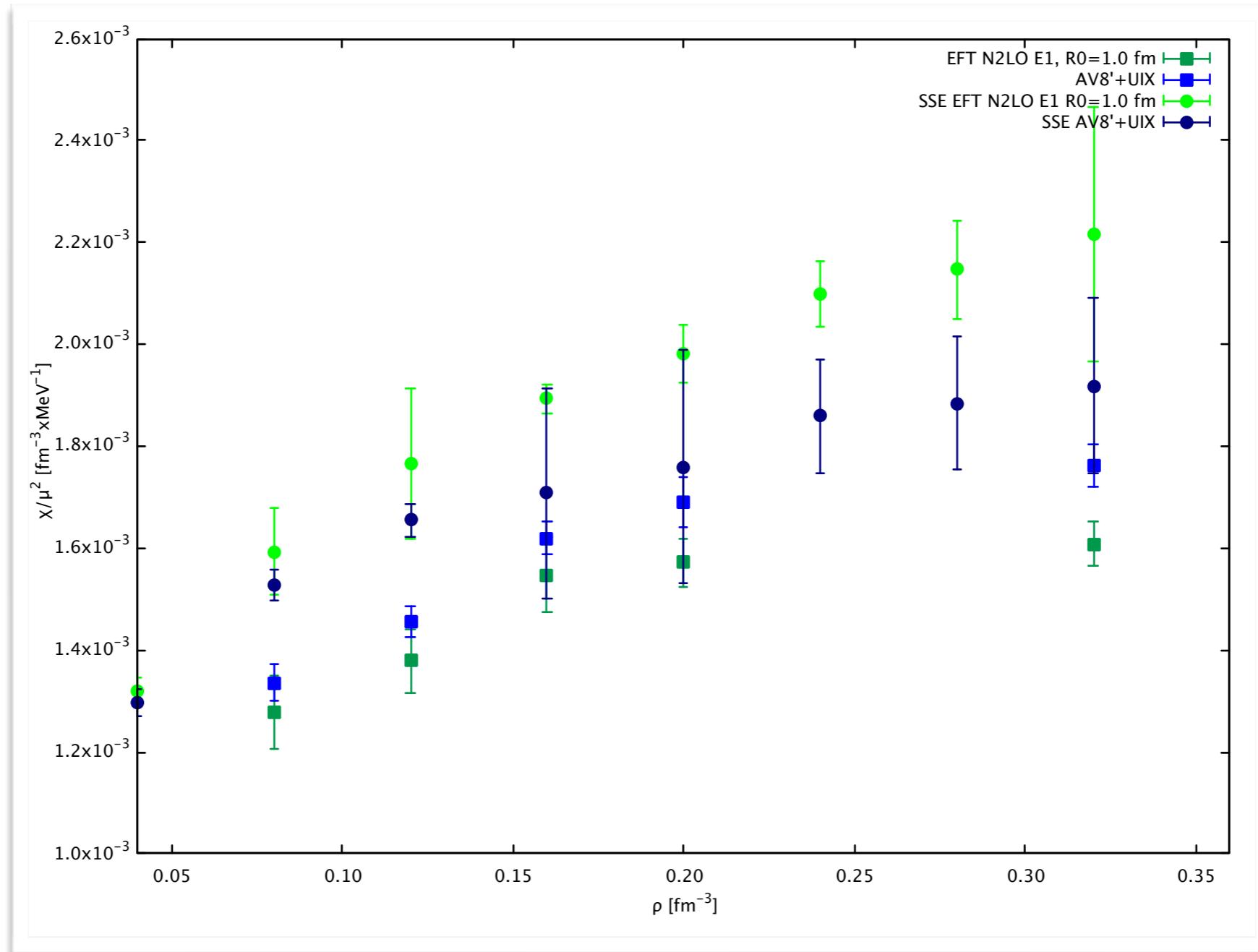


Simple model

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Simple model



Summary

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- ◆ Using TABC to study partially polarized systems reducing finite-size effects and arbitrary particle numbers.
- ◆ Magnetic susceptibility computed ‘correctly’ with *ab initio* methods.
- ◆ Phenomenological and Chiral interactions give similar results, some differences at high densities.