

The ${}^6\text{Li}$ ground state within the Hyperpsherial Harmonics basis

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- The Hyperspherical Harmonic(HH) method
- The ${}^6\text{Li}$ ground state
 - Convergence properties
 - Electromagnetic structure
 - $\alpha + d$ structure
- Conclusions and future prospective

$$H = \sum_i \frac{p_i^2}{2M} + \sum_{i < j} V(i, j) + \sum_{i < j < k} W(i, j, k) + \dots$$

Search for accurate solution of $H\Psi = E\Psi$

- Variational approach
- Expansion of Ψ on the basis of **Hyperspherical Harmonics** functions
- [A. Kievsky, S. Rosati, M. Viviani, L.E. Marcucci, and L. Girlanda J. Phys. G, **35**, 063101 (2008)]
- Applied for $A = 3, 4$ and now also for $A = 6$

The HH method general characteristics

- Jacobi vectors \Rightarrow CoM decoupling
- HH functions

$$L^2(\Omega)\mathcal{Y}_{[K]}(\Omega) = K(K + D - 2)\mathcal{Y}_{[K]}(\Omega)$$

- Transformation Coefficients
 - $\mathcal{Y}_{[K]}(\Omega') = \sum_{[K']}^{K=K'} a_{[K],[K']} \mathcal{Y}_{[K']}(\Omega)$
 - antisymmetrization
 - overcomplete basis \Rightarrow orthogonalization

$$\langle \Phi_\alpha | V(1, 2) | \Phi_\beta \rangle = \sum_{[\alpha'], [\beta']} a_{[\alpha], [\alpha']} a_{[\beta], [\beta']} V_{[\alpha'], [\beta']}(1, 2)$$

All correlations generated by the potential are constructed through the TC

- the coupling of quantum number are the same for **any** potential
- **No need to save the matrix elements**
only the combination of the transformation coefficients
- ~ 3 hours for constructing and diagonalize the Hamiltonian

- We will use SRG evolved N^3LO_{500} NN interaction [1-2]
 - The Coulomb interaction is included as “bare” (not SRG evolved)
 - SRG evolution parameter $\Lambda = 1.2, 1.5, 1.8 \text{ fm}^{-1}$
- Explorative study with $NNLO_{sat}$ [3]
- No 3-body forces (for now)
- We compute the mean values of “bare” operators

[1] S.K. Bogner, R.J. Furnstahl, and R.J. Perry, *PRC* **75**, 061001(R) (2007)

[2] D.R. Entem and R. Machleidt, *PRC* **68**, 041001(R) (2003)

[3] A. Ekström, *et al.*, *PRC* **91**,051301 (2015)

Selection of the states

- Include all the states up to a $K_{max} \Rightarrow$ FAIL!!
- NOT all the states gives the same contribution \Rightarrow division by class [1]
 - Centrifugal barrier $\Rightarrow \ell_1 + \dots + \ell_5 \leq 2$
 - two-body correlations are more important
- The ${}^6\text{Li}$ ground state is a $J^\pi = 1^+$ state

wave	class	int.	K_{max}
S	C1	two-body	14
	C3	many-body	10
D	C2	two-body	12
	C4	many-body	10
P	C5	all	8

[1] M. Viviani, *et al.*, PRC 71,024006 (2005)

- For each channel we define

$$\Delta E_{ch}(K) = E_{ch}(K) - E_{ch}(K + 2)$$

- Fit the quantity ΔE_{ch} with [1]

$$\Delta E_{ch}(K) = A_{ch} e^{-b_{ch}K} (1 - e^{-2b_{ch}})$$

- Missing energy for channel

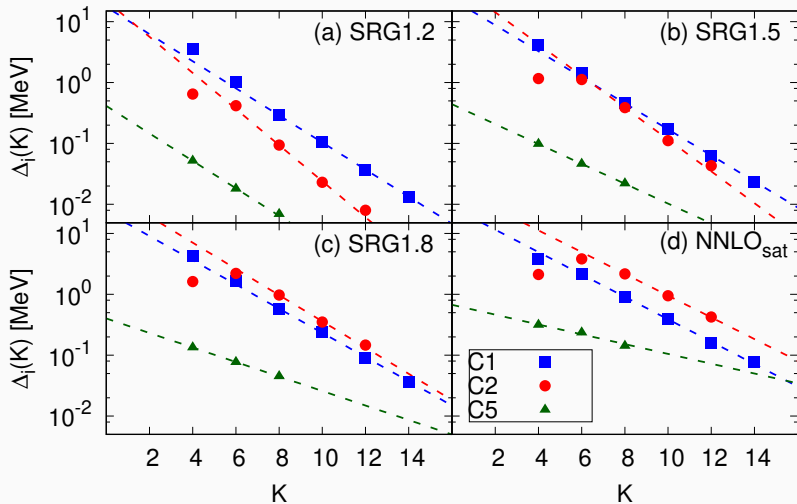
$$\Delta E_{ch}(\infty) = \sum_{K=\bar{K}}^{\infty} \Delta E_{ch}(K) = A_{ch} e^{-b_{ch}\bar{K}}$$

- \bar{K} maximum K used for the channel
- Extrapolated energy

$$E(\infty) = E(\bar{K}) - \sum_{ch} \Delta E_{ch}(\infty)$$

[1] S.K. Bogner *et al.*, NPA **801**, 21 (2008)

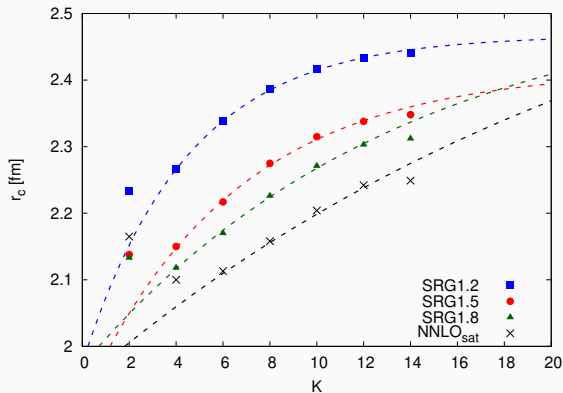
Convergence of the classes



Λ [fm $^{-1}$]	$E(\bar{K})$ [MeV]	$E(\infty)$ [MeV]	Ref. [1]	Exp.
SRG1.2	-31.735	-31.767(7)	-31.85(5)	-31.99
SRG1.5	-32.699	-32.789(15)	-33.00(5)	-31.99
SRG1.8	-32.093	-32.305(25)	-32.8(1)	-31.99
NNLO $_{sat}$	-29.77	-30.71(15)	–	-31.99

- The errors come from the fit
- Results of Ref. [1] extrapolated from $N_{max} = 10$

[1] E.D Jungerson, P. Navrátil and R.J. Furnstahl, PRC **83**, 034301 (2011)



	SRG1.2	SRG1.5	SRG1.8	NNLO _{sat}	Ref. [1]	Exp.
r_c [fm]	2.47(1)	2.42(3)	2.54(11)	2.7(3)	2.40(6)	2.540(28)

[1] CDB2k-SRG1.5 C. Forssén, E. Caurier, P. Navrátil, PRC **71**, 021303 (2009)

$${}^6\text{Li} \simeq \alpha + d \Rightarrow \mu_z({}^6\text{Li}) \simeq \mu_z(d)$$

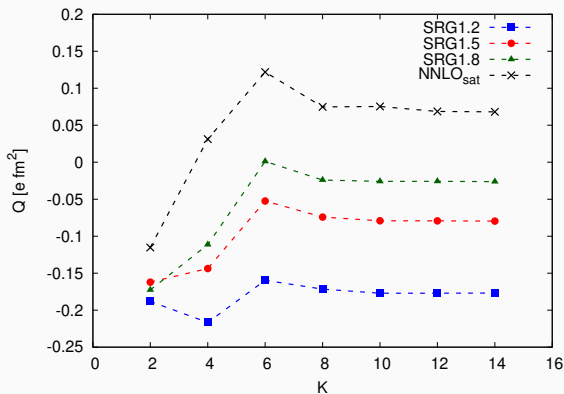
Experiment tell us $\mu_z({}^6\text{Li}) < \mu_z(d)$

	$\mu_z(d)$	$\mu_z({}^6\text{Li})$
SRG1.2	0.872	0.865
SRG1.5	0.868	0.860
SRG1.8	0.865	0.856
NNLO _{sat}	0.860	0.850
Exp.	0.857	0.822

- Negative contribution only from the $L = 2$ $S = 1$ component
⇒ NOT SUFFICIENT
- We need two body currents contribution!! [1]

[1] R. Schiavilla, *et al.*, PRC **99**, 034005 (2019)

Electric quadrupole moment



	SRG1.2	SRG1.5	SRG1.8	NNLO _{sat}	Ref. [1]	Exp.
Q [e fm ²]	-0.177	-0.080	-0.026	+0.068	-0.066(40)	-0.0806(8)

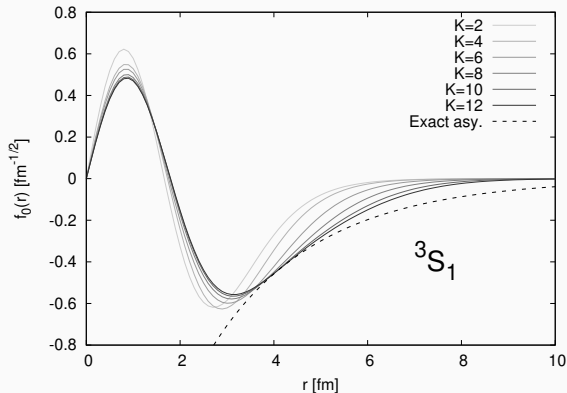
- Large cancellations between different K

Contribution of the various waves

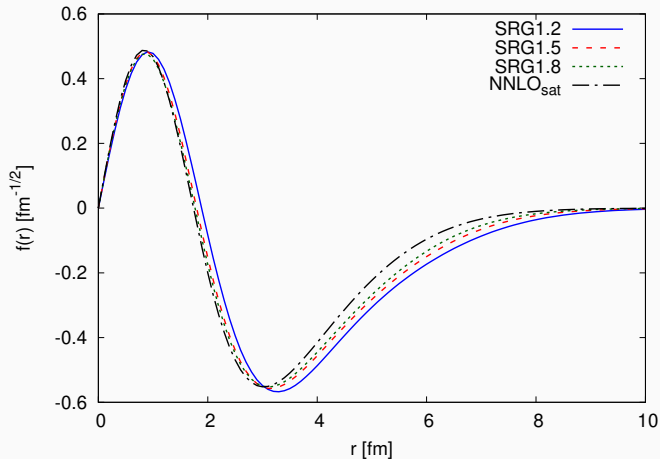
	$S - D$	$D - D$	$P - P$	$P - D$
SRG1.2	-0.173	-0.022	0.009	0.009
SRG1.5	-0.080	-0.021	0.012	0.010
SRG1.8	-0.028	-0.020	0.012	0.010
NNLO _{sat}	0.058	-0.016	0.015	0.011

- Direct connection with the strength of the tensor term in the potential
- Two-body currents contribution could be necessary!!

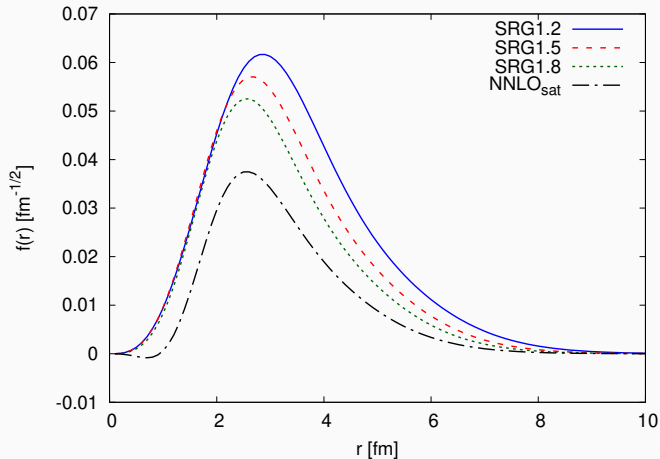
$$\frac{f_L(r)}{r} = \langle [(\Psi_\alpha \otimes \Psi_d)_S Y_L(\hat{r})]_J | \Psi_{6\text{Li}} \rangle$$

 $N^3\text{LO500-SRG1.5}$

$$\frac{f_L(r)}{r} = \langle [(\Psi_\alpha \otimes \Psi_d)_S Y_L(\hat{r})]_J | \Psi_{6\text{Li}} \rangle$$



$$\frac{f_L(r)}{r} = \langle [(\Psi_\alpha \otimes \Psi_d)_S Y_L(\hat{r})]_J | \Psi_{6\text{Li}} \rangle$$



ANC and spectroscopic factor (preliminary)

$$C_L = \frac{f_L(r)}{W_{-\eta, L+1/2}(2kr)} \quad S_L = \int_0^\infty dr |f_L(r)|^2$$

	B_c [MeV]	C_0 [$\text{fm}^{-1/2}$]	C_2 [$\text{fm}^{-1/2}$]	$S_0 + S_2$
SRG1.2	2.985(7)	-3.96(2)	0.101(1)	0.914
SRG1.5	2.385(15)	-3.05(3)	0.056(1)	0.871
SRG1.8	1.684(7)	-2.34(6)	0.030(1)	0.841
Exp.	1.4743	-2.91(9)	0.077(18)	0.85(4)

- The ANCs depend on the binding energy B_c and only
- The spectroscopic factor represents the percentage of $\alpha + d$ clusterization of ${}^6\text{Li}$

$$B_c = B_{{}^6\text{Li}} - B_\alpha - B_d, \quad k = \sqrt{2\mu B_c/\hbar^2} \quad \text{and} \quad \eta = 2.88\mu/k$$

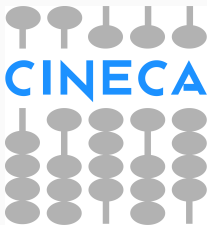
Conclusions and future prospective

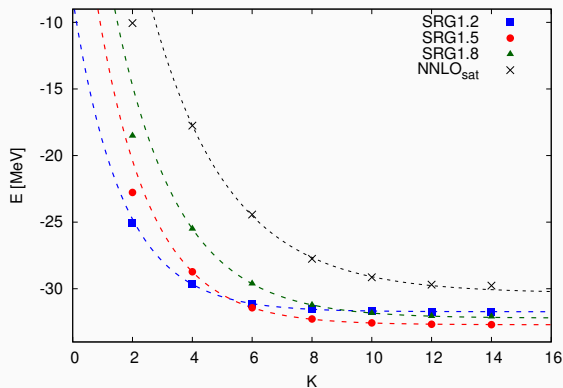
- Extension of the HH method to $A > 4$ is possible
 - Time consumption for coupling the TC $\sim 30.000h$
 \Rightarrow technically possible to increase the basis
 - ~ 3 human hours for building and solving the ${}^6\text{Li}$ Hamiltonian
 \Rightarrow easy to test various potentials
- Ground state of ${}^6\text{Li}$ within the HH approach
 - Nicely covered for SRG potentials
 - Magnetic dipole moment and electric quadrupole moment
 - strongly depend on tensor forces
 - need two-body currents!!
 - Towards scattering: the $\alpha + d$ clusterization
 - ANCs and spectroscopic factors

Pisa group

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Calculations supported by





- Exponential behavior [1]

$$E(K) = E(\infty) + Ae^{-bK}$$

[1] S.K. Bogner *et al.*, NPA **801**, 21 (2008)