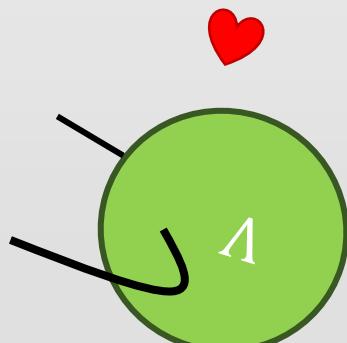
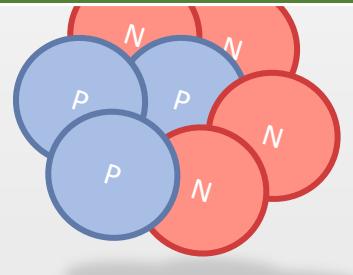
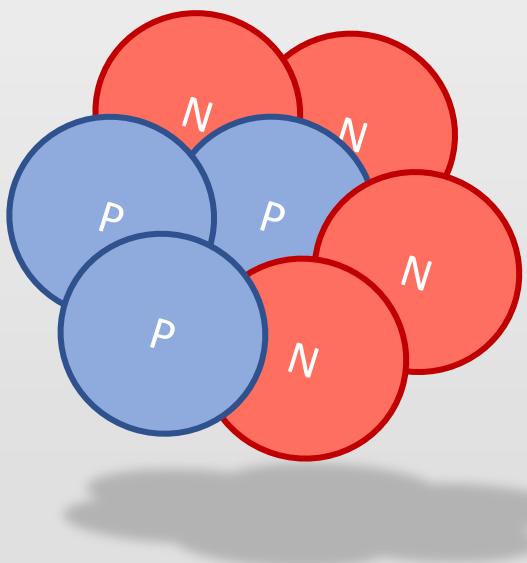


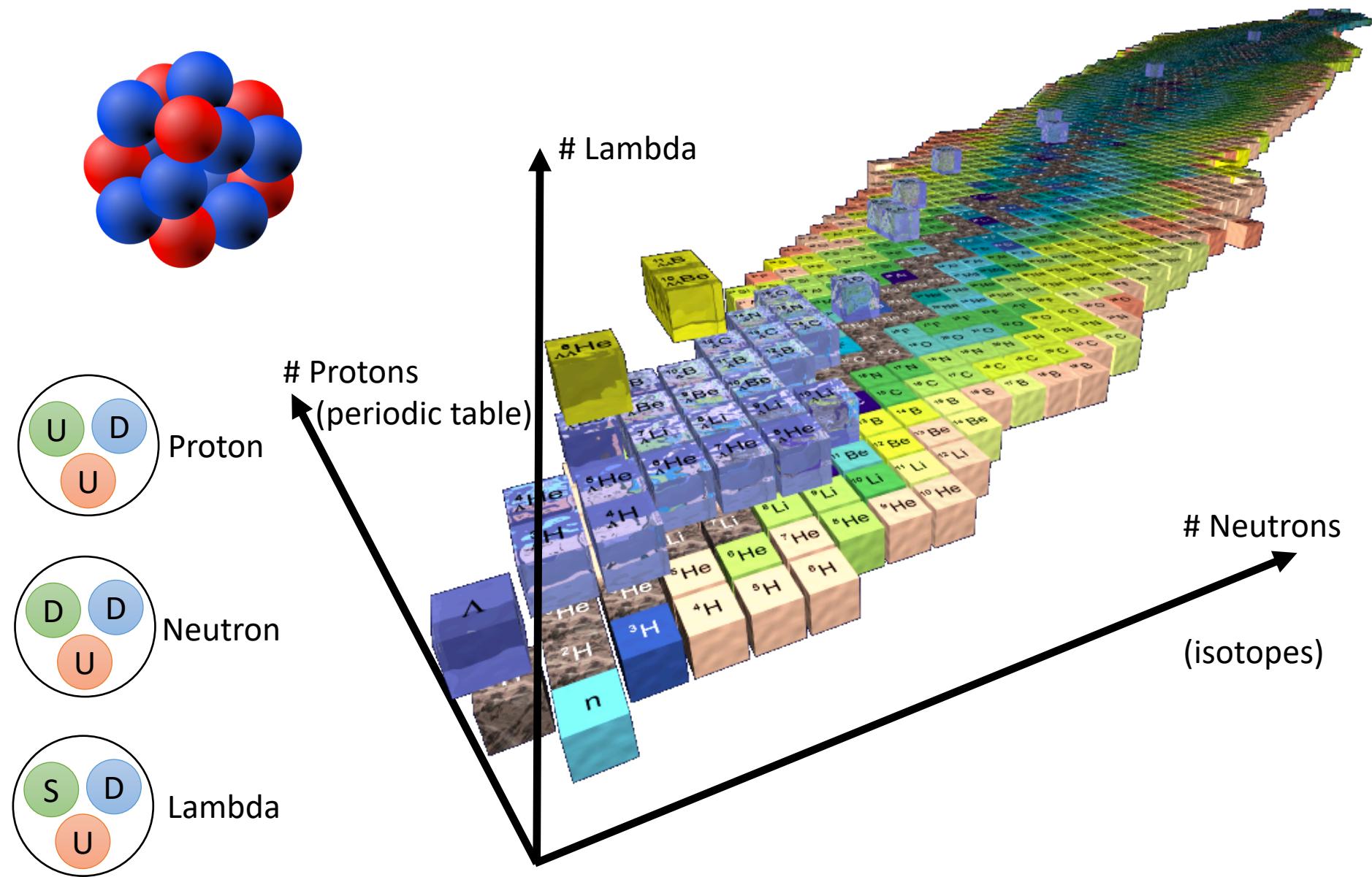
# Single- and double- $\Lambda$ hypernuclei in EFT( $\pi$ )



Lorenzo Contessi  
Martin Shäfer  
Nir Barnea  
Avraham Gal  
Jiří Mareš



THE HEBREW  
UNIVERSITY  
OF JERUSALEM



## Phaseshift shortage

$N - \Lambda$

$\Lambda - \Lambda$

Experimentally **not known**  
Theoretically **debated**

## Unknown FEW-BODY

$nn\Lambda$

$n\Lambda\Lambda$

${}^5_{\Lambda\Lambda}H$

$np\Lambda\Lambda$

$nn\Lambda\Lambda$

J-PARC P75 proposal

*Theoretical challenging*

---- ----

*Few input data to tune theories*



*Minimal number of free parameters*

## Known FEW-BODY

${}^4_{\Lambda}H(S=0)$        ${}^3_{\Lambda}H$

${}^4_{\Lambda}H(S=1)$        ${}^5_{\Lambda}He$

Experimentally known  
**Theoretically hard to be described all together**

## Known Double- $\Lambda$

${}^6_{\Lambda\Lambda}He$        ${}^{10}_{\Lambda\Lambda}Be$

${}^{11}_{\Lambda\Lambda}Be$

# New experiments are planned

STAR collaboration

J-parc

HALQCD

BES III

J-Lab

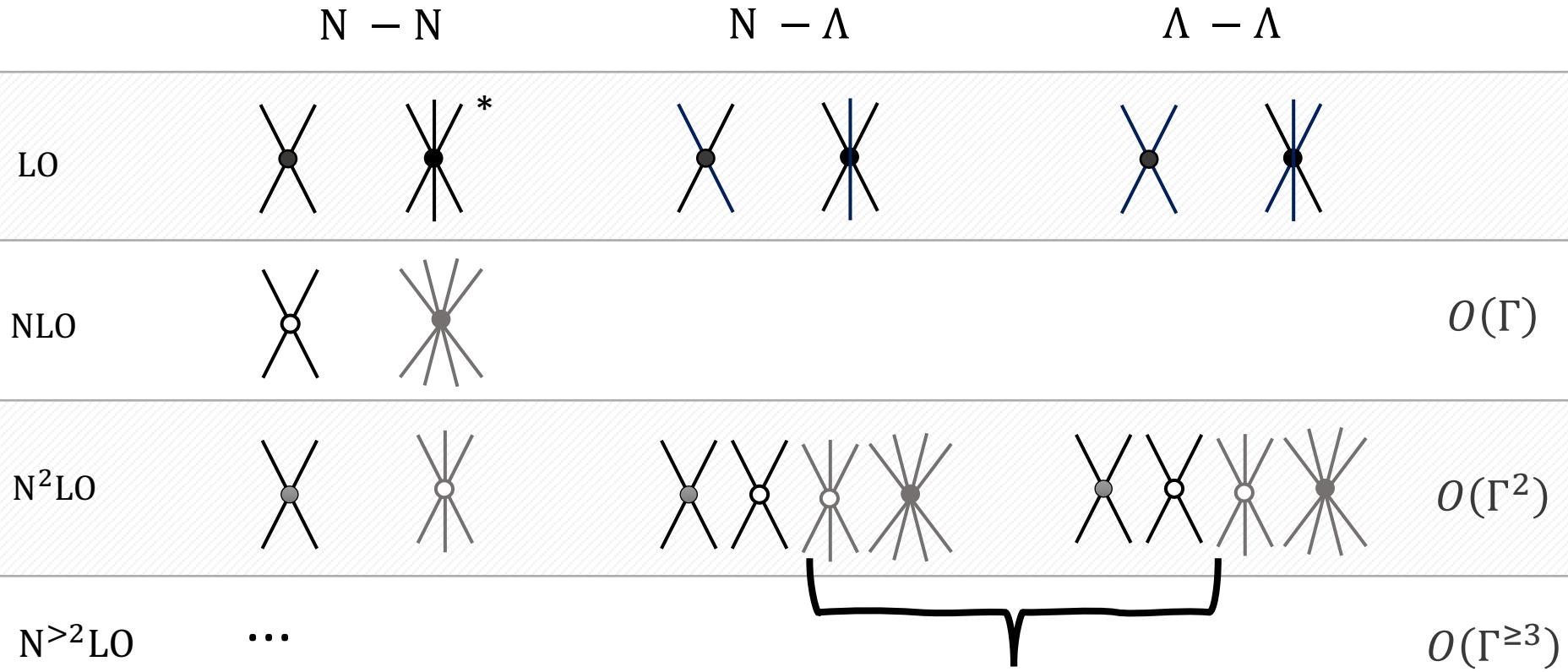
PANDA

LHC

# Abundant open queries

- **Description** of few-body hypernuclei
- **Double  $\Lambda$**  hypernuclei description
- **Life time** of  ${}^3_{\Lambda}\text{H}$  and  ${}^3_{\Lambda}\text{n}$
- **Charge symmetry breaking** ( ${}^4_{\Lambda}\text{H} - {}^4_{\Lambda}\text{He}$ )
- **$\Lambda^*(1405)$  matter**
- **Neutron star equation of state**
- ...

# Pionless powercounting



$$\Gamma_{NN} = \frac{Q}{m_\pi} = 0.5 \sim 0.8$$

$$\Gamma_{N\Lambda} = \frac{Q}{2 m_\pi^{**}} \sim 0.2$$

B. Bazak, Four-Body Scale in Universal Few-Boson Systems, PRL 122.143001 (2019)

G.P. Lepage, How to renormalize the Schrodinger equation (1997)

van Kolck, U. Nucl.Phys. A645 (1999) 273-302

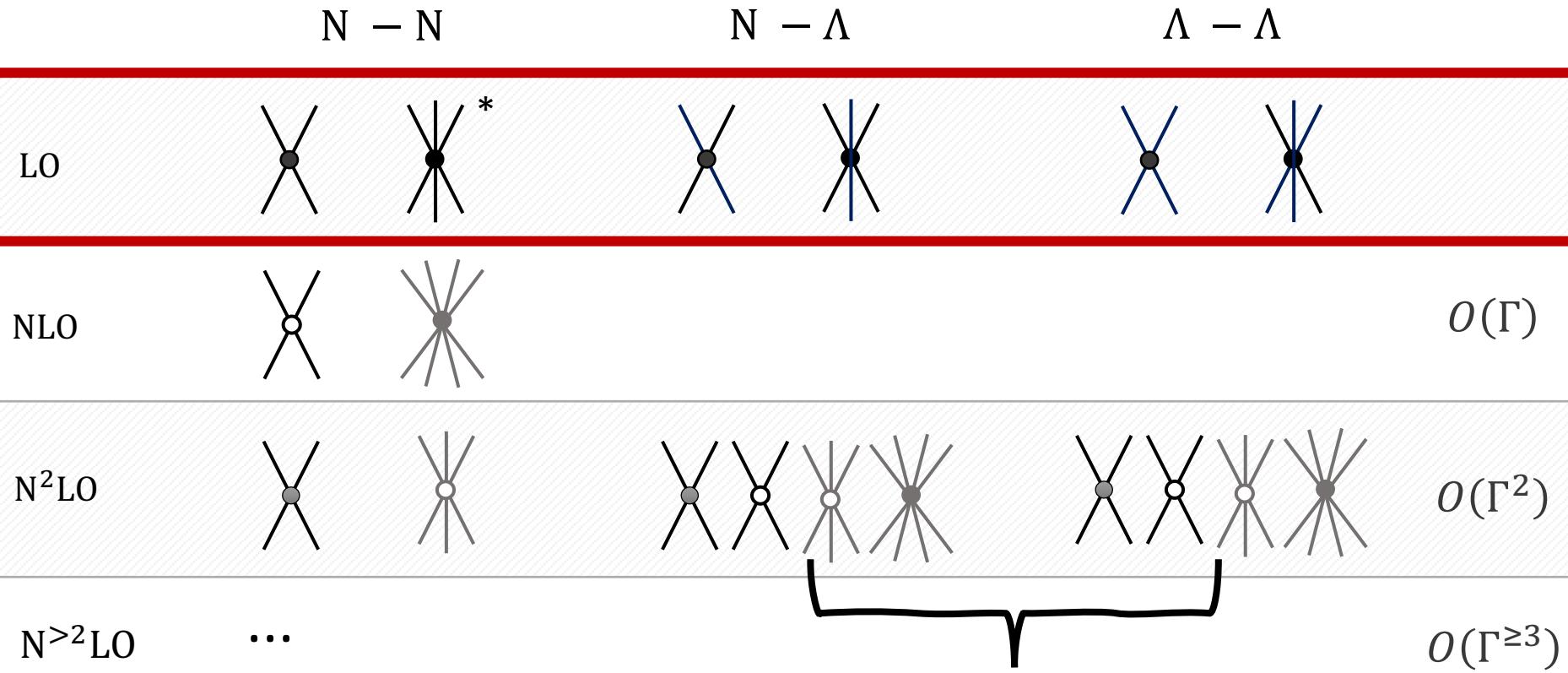
Chen, Jiunn-Wei et al. Nucl.Phys. A653 (1999)

S. König, H. W. Grießhammer, H. W. Hammer, and U. van Kolck J. Phys. G43, 055106 (2016)

\* Three-body force is necessary  
to avoid Thomas collapse

\*\* OPE not allowed

# Pionless powercounting



B. Bazak, Four-Body Scale in Universal Few-Boson Systems, PRL 122.143001 (2019)

G.P. Lepage, How to renormalize the Schrodinger equation (1997)

van Kolck, U. Nucl.Phys. A645 (1999) 273-302

Chen, Jiunn-Wei et al. Nucl.Phys. A653 (1999)

S. König, H. W. Grießhammer, H. W. Hammer, and U. van Kolck J. Phys. G43, 055106 (2016)

Cortona 2019 - Lorenzo Contessi

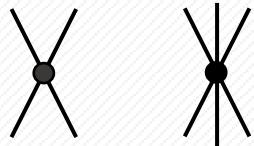
\* Three body force is necessary  
to avoid Thomas collapse

\*\* OPE not allowed

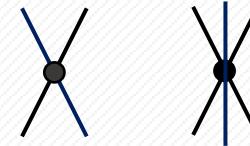
# Pionless powercounting

N - N      N -  $\Lambda$        $\Lambda$  -  $\Lambda$

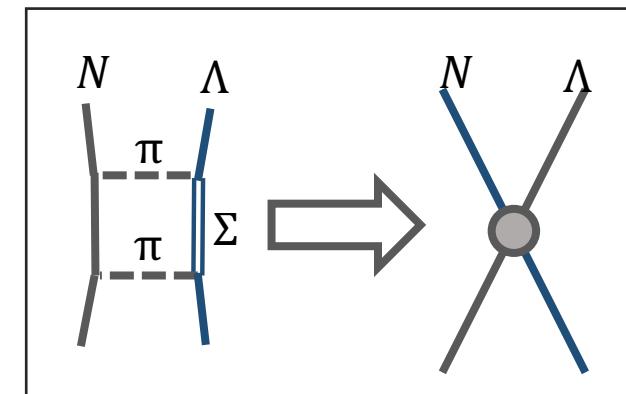
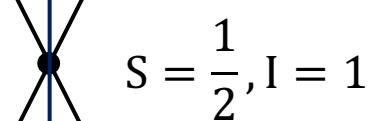
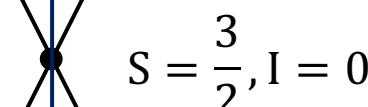
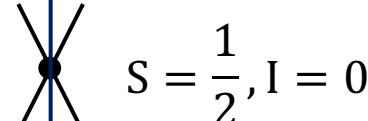
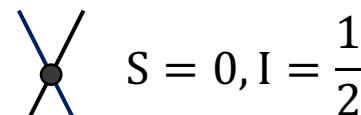
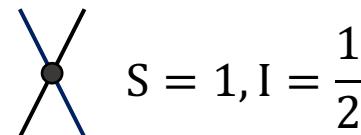
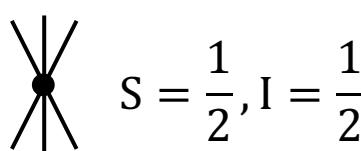
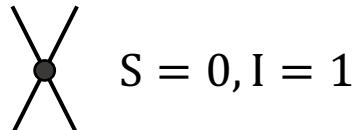
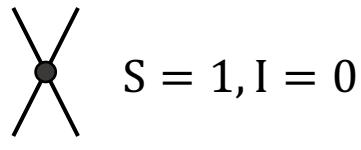
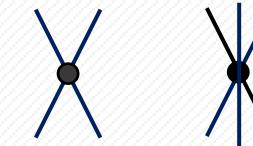
LO



N -  $\Lambda$



$\Lambda$  -  $\Lambda$



LO



Pionless at Leading Order



## Interaction:

$$V = \sum_{ij} C_p \delta^p(\vec{r}_i, \vec{r}_j) + \sum_{ijk} D_p \delta^p(\vec{r}_i, \vec{r}_j, \vec{r}_k)$$

LO

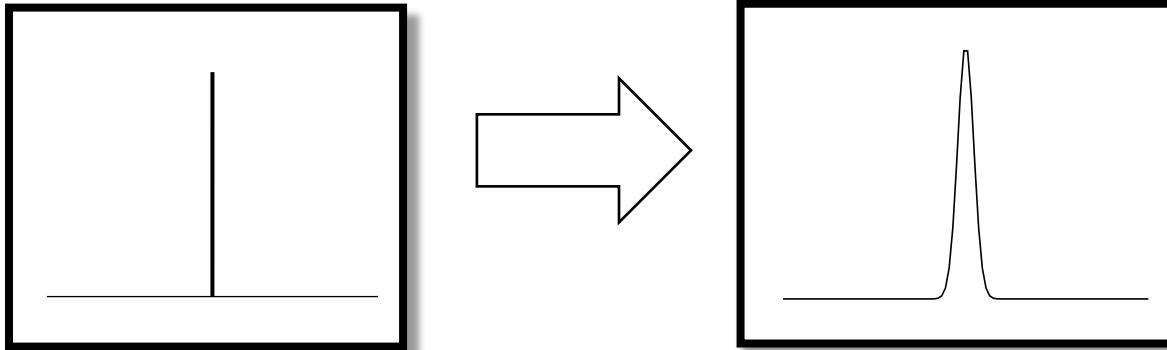


Pionless at Leading Order

**Interaction:**

$$V = \sum_{ij} C_p \delta^p(\vec{r}_i, \vec{r}_j) + \sum_{ijk} D_p \delta^p(\vec{r}_i, \vec{r}_j, \vec{r}_k)$$

$p$  represent projections  
In the 2- and 3-body S-wave channels.

**Regularization / Renormalization required***Observables are cut-off dependent:*

$$O_\lambda = O_\infty + \frac{\alpha}{\lambda} + \frac{\beta}{\lambda^2} + \frac{\gamma}{\lambda^3} + \dots$$

- $\lambda \rightarrow \infty *:$
- Regularization/model independent
  - Observables are  $\lambda$  dependent

\*  $\lambda \gg M$

LO

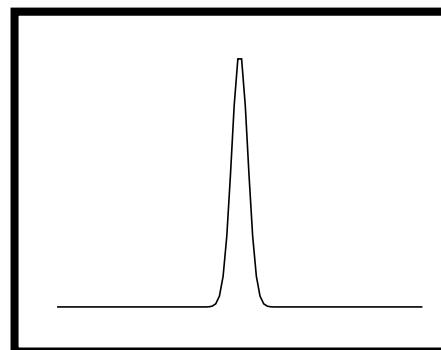
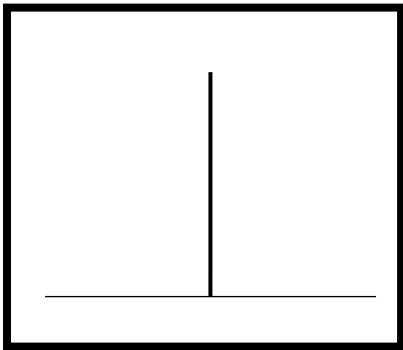


Pionless at Leading Order

**Interaction:**

$$V = \sum_{ij} C_p^\lambda P^p e^{-\frac{\lambda^2 r_{ij}^2}{4}} + \sum_{ijk} D_p^\lambda P^p \sum_{cyc} \left[ e^{-\frac{\lambda^2 (r_{ij}^2 + r_{ik}^2)}{4}} \right]$$

Projections  
in channels

**Regularization / Renormalization required**

$$\lambda \rightarrow \infty *:$$

- Regularization/model independent
- Observables are  $\lambda$  dependent

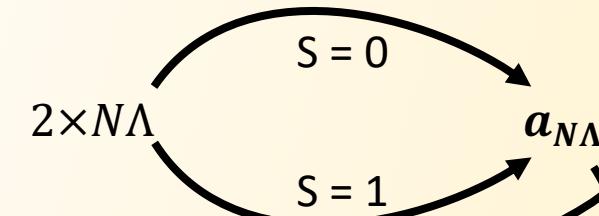
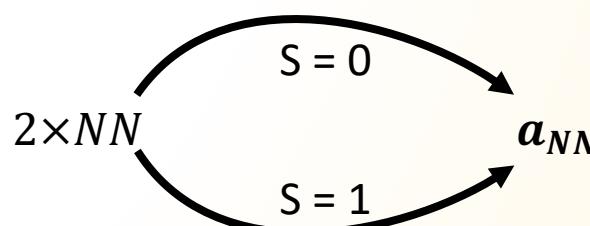
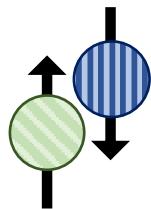
*Observables are cut-off dependent:*

$$O_\lambda = O_\infty + \frac{\alpha}{\lambda} + \frac{\beta}{\lambda^2} + \frac{\gamma}{\lambda^3} + \dots$$

$$* \lambda \gg M$$

# Fitting input

## Two body



$\Lambda\Lambda \longrightarrow a_{\Lambda\Lambda}$

No precise  
Experimental data

$NNN$

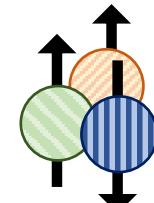
$^3\text{H}$

$^3\text{H}$

$^3\cancel{\text{H}}(S = 2/3)$

$^3\cancel{\Lambda}\text{He}$

## Three body



$3 \times N\Lambda\Lambda$

$^4\Lambda\text{H}(S = 0, I = 1/2)$

$^4\Lambda\text{H}(S = 1, I = 1/2)$

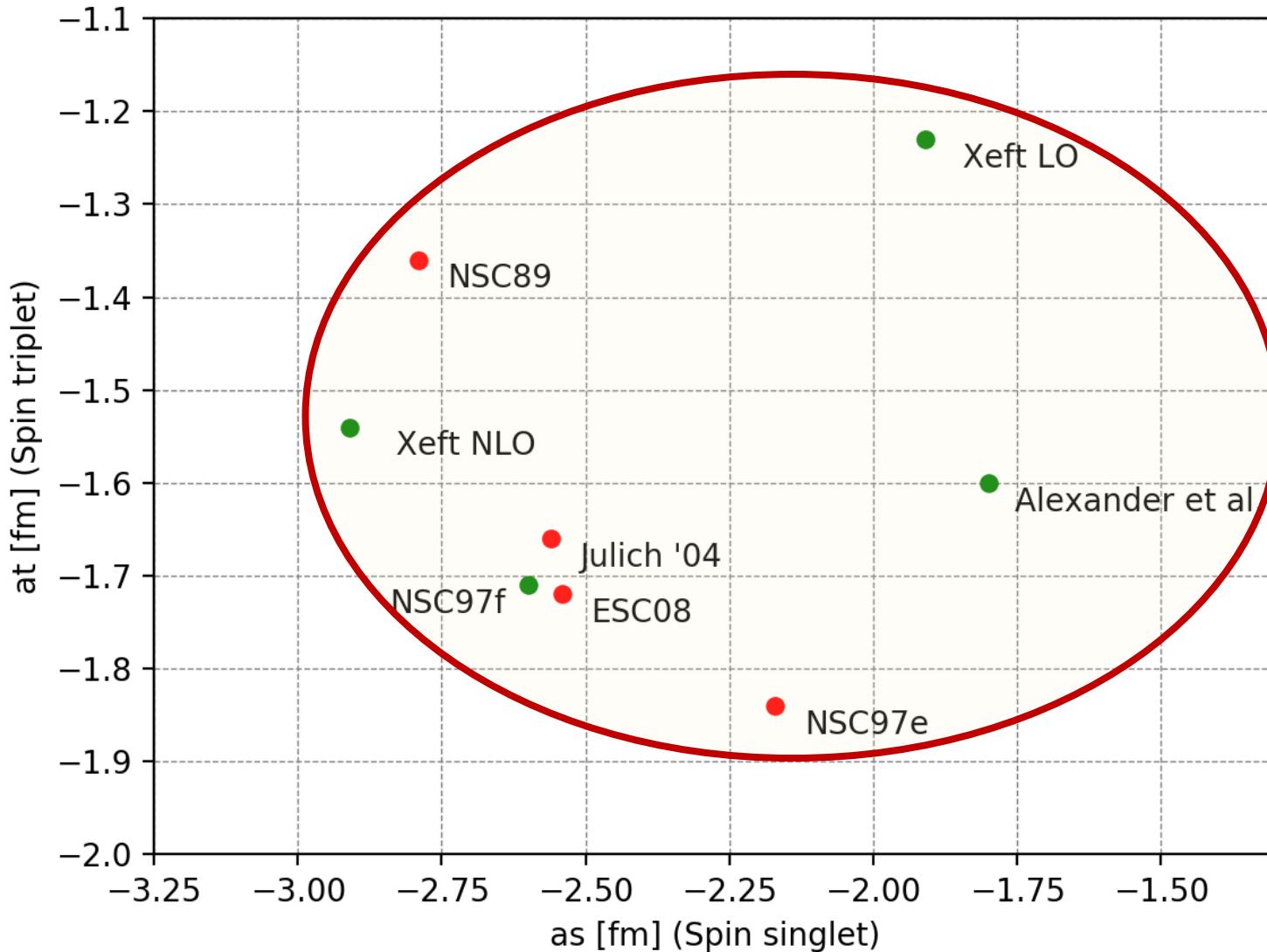
$N\Lambda\Lambda$

$^6\cancel{\Lambda}\Lambda\text{He}$

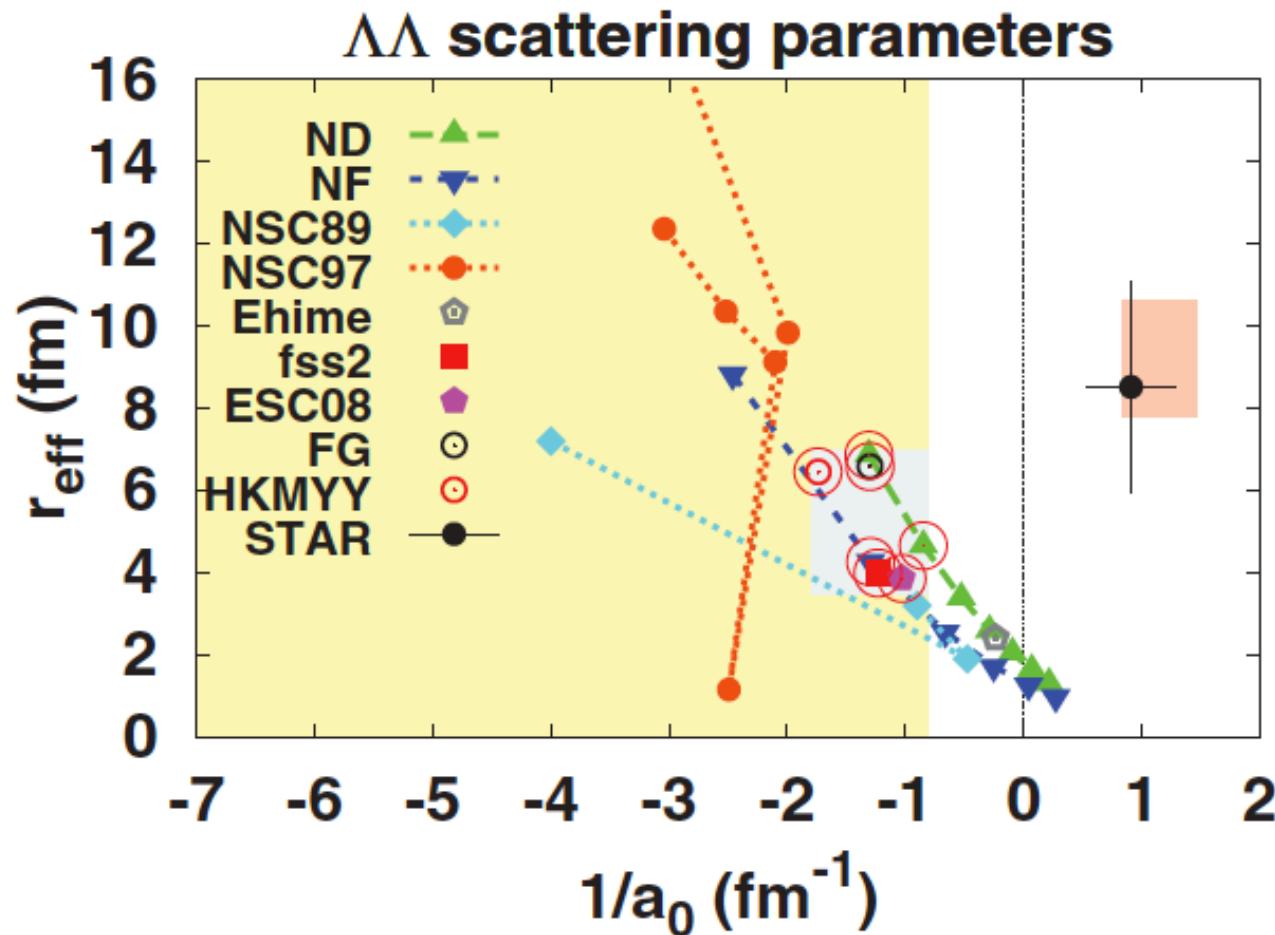
$^6\Lambda\Lambda\text{He}$

# $N\Lambda$ scattering length

A. Gal et al. - Strangeness in nuclear physics - Rev.Mod.Phys. 88 (2016) no.3, 035004



# $\Lambda\Lambda$ Scattering data

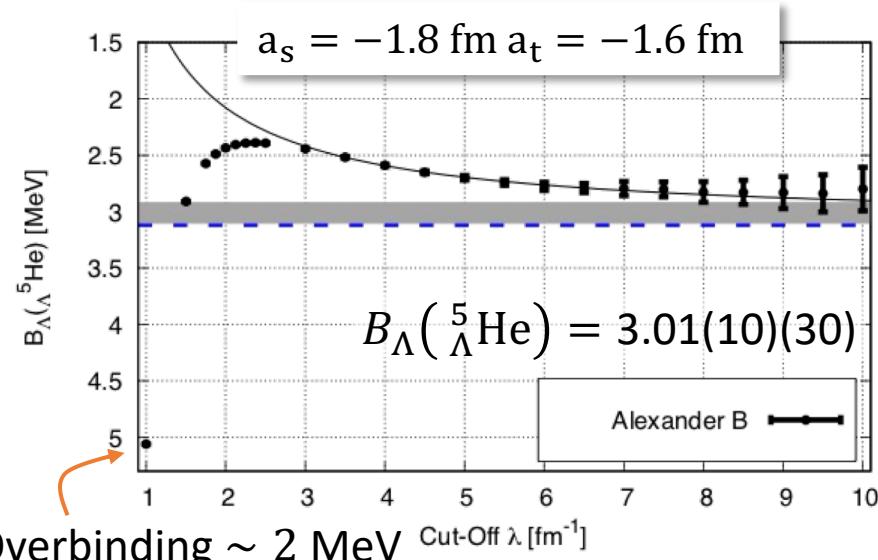


# Results

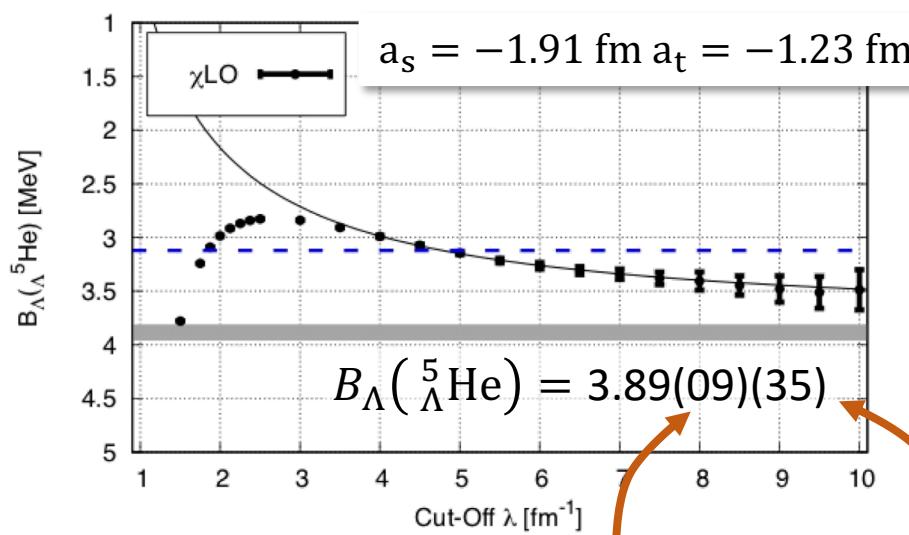
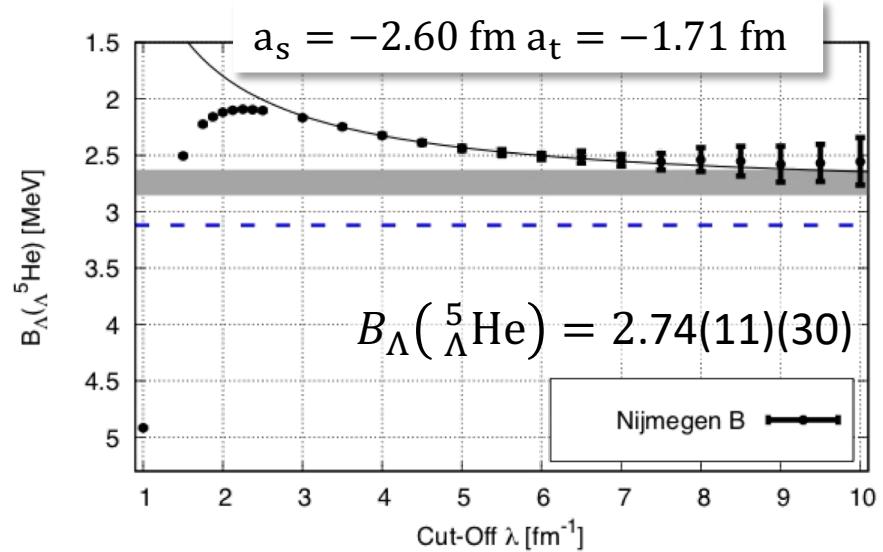
- ❖  $^5_{\Lambda}\text{He}$
- ❖  $^5_{\Lambda\Lambda}\text{H}$
- ❖  $np\Lambda\Lambda$

# ${}^5_{\Lambda}\text{He}$ : $\Lambda$ separation energy

15

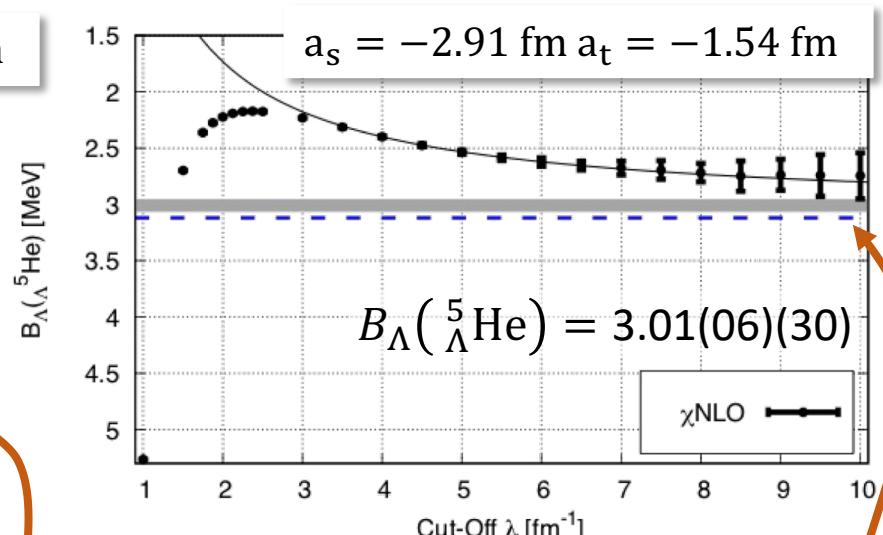


Overbinding  $\sim 2$  MeV



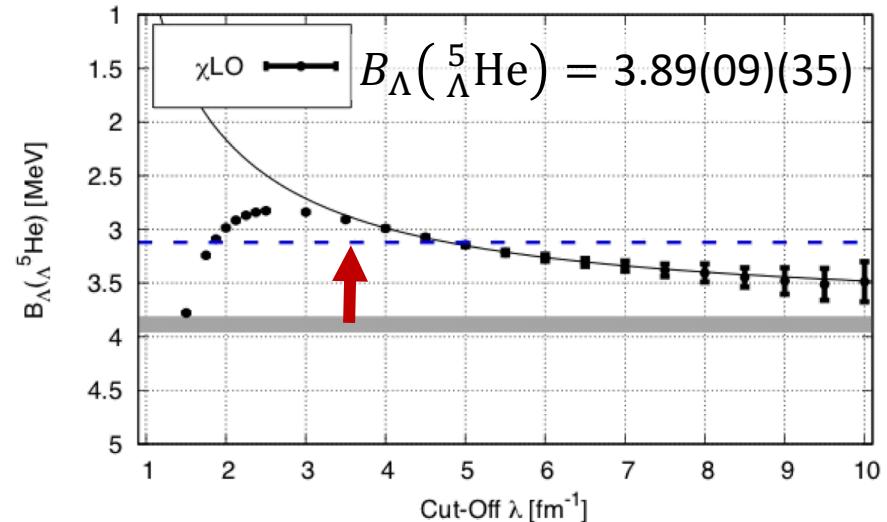
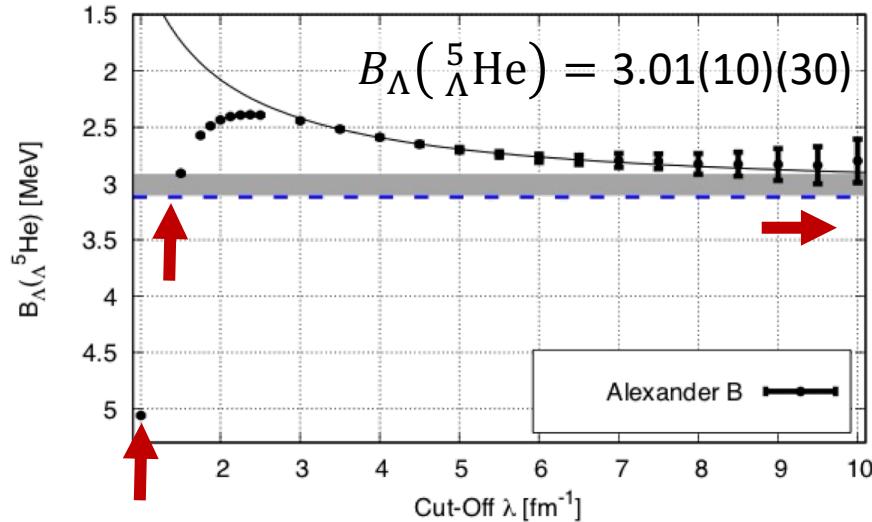
Numerical uncertainty

Theory uncertainty



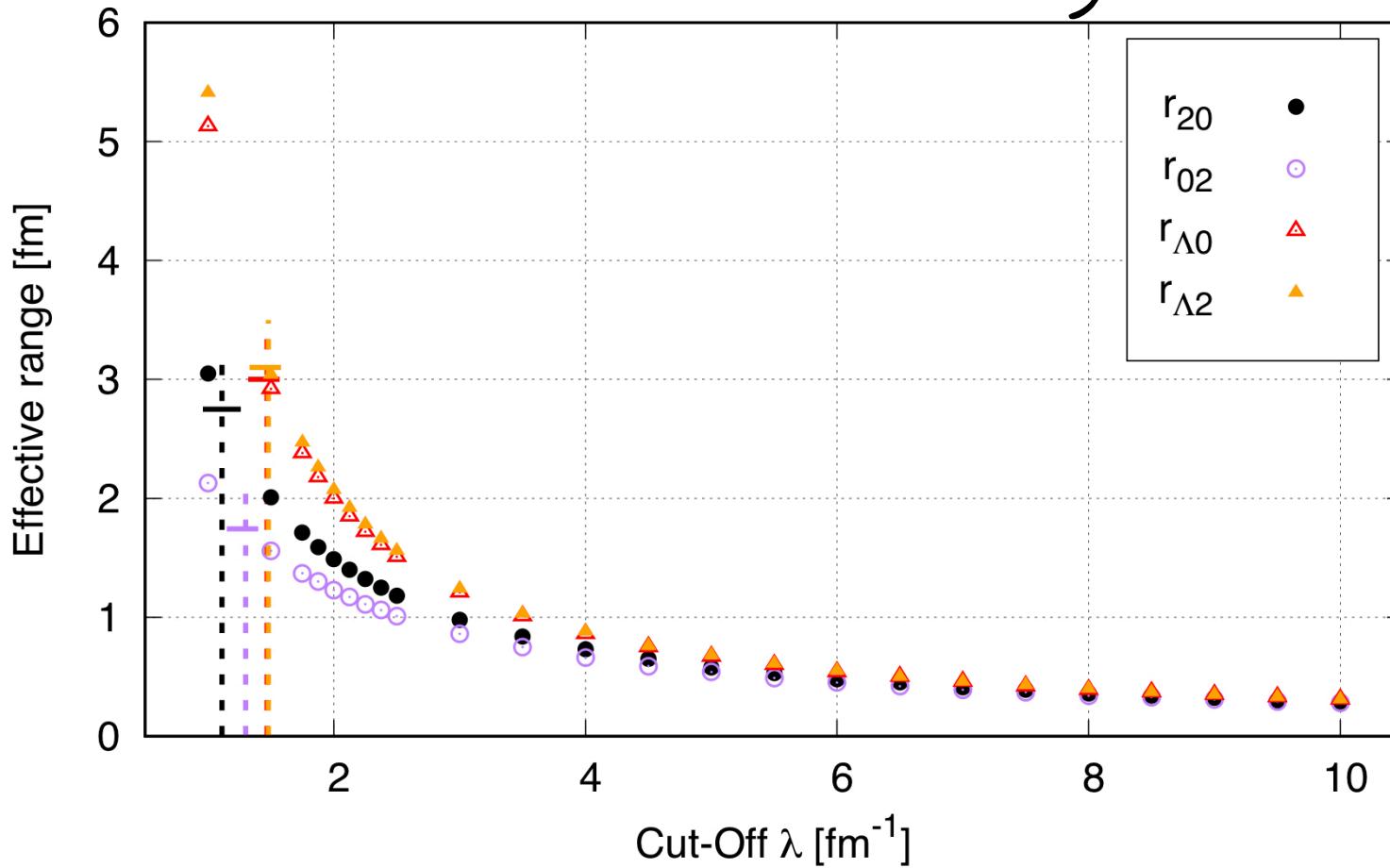
Experimental: 3.12(2) MeV

# Relevant cut-offs



- $\lambda \gg M \sim 2 m_{\pi}$  Is the **standard choice** for a renormalizable EFT.
- $2 < \lambda < 4 \text{ fm}^{-1}$  is between  $M \sim 2 m_{\pi}$  and the closest **not-included vector meson**.
- $\lambda \sim 1 \text{ fm}^{-1} \rightarrow$  **Many similarities** with models that **overbind**  ${}_{\Lambda}^5\text{He}$ .
- $\lambda \sim 1.5 \text{ fm}^{-1}$  **describes**  $r_0$  and effectively takes into account sub-leading orders.

# Effective range



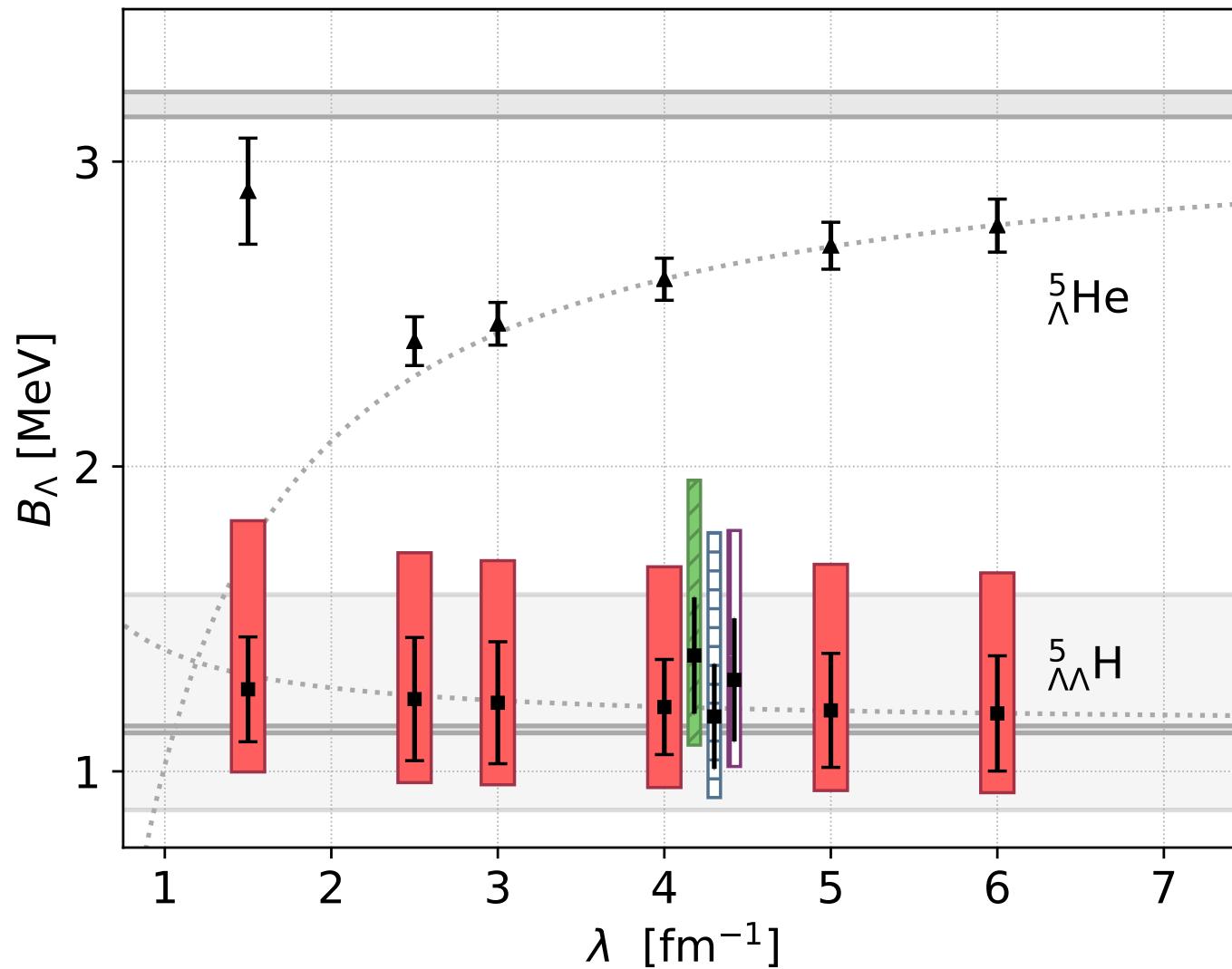
Dashed lines represent cut-off that fit the experimental values

$$\lambda(r_0^{(20)}) = 1.11 \text{ fm}^{-1},$$

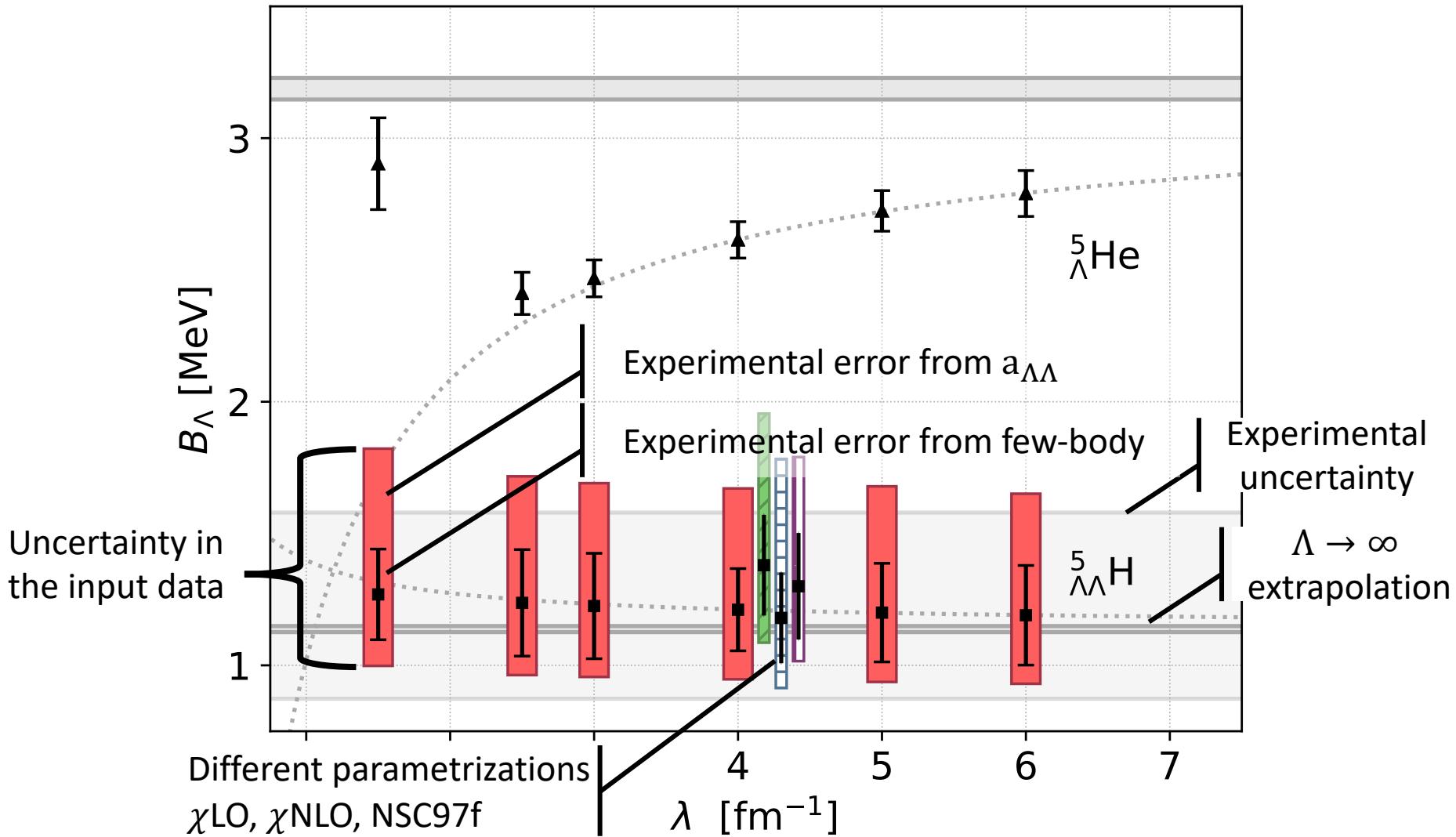
$$\lambda(r_0^{(02)}) = 1.30 \text{ fm}^{-1},$$

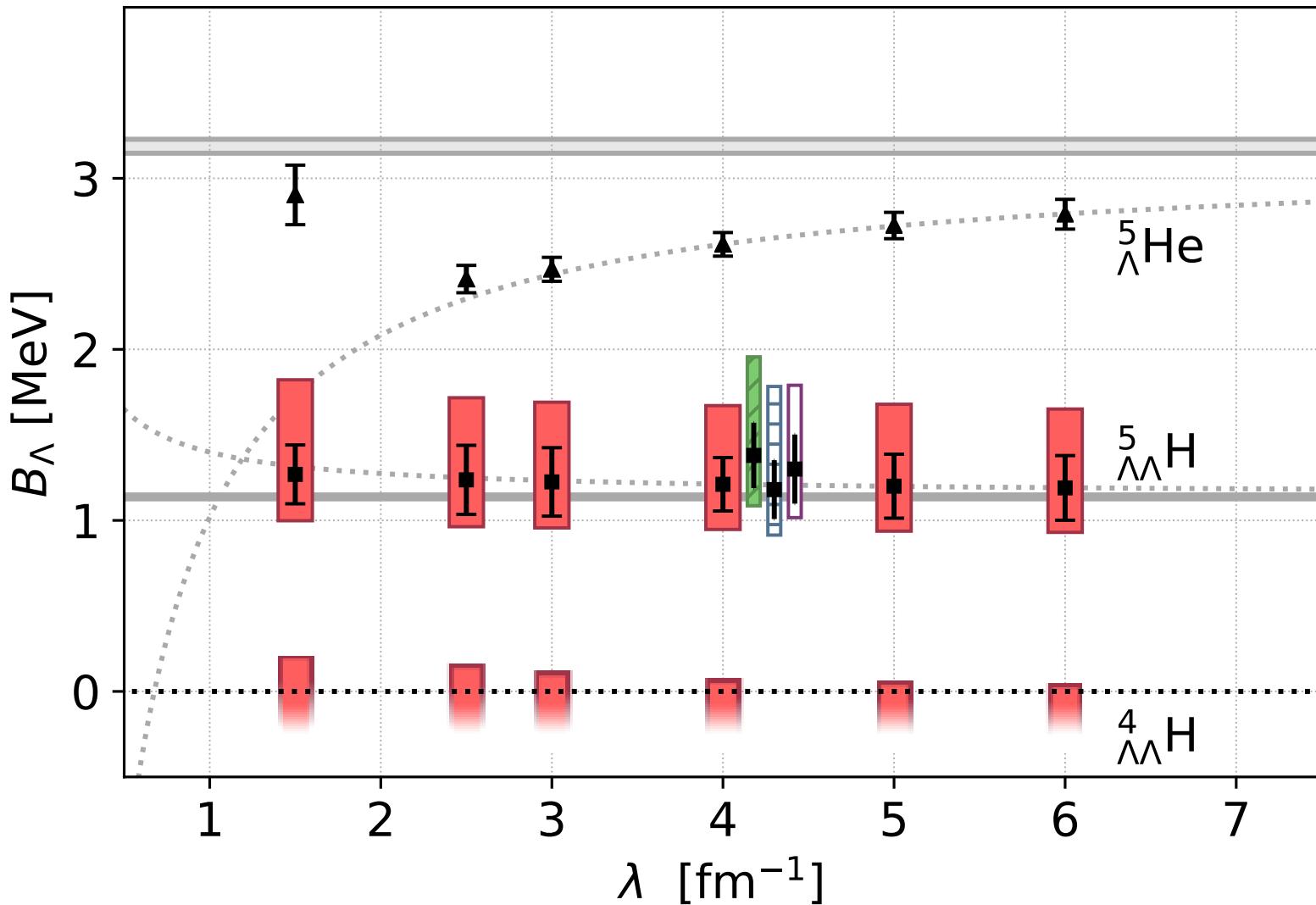
$$\lambda(r_0^{(\Lambda 0)}) = 1.47 \text{ fm}^{-1},$$

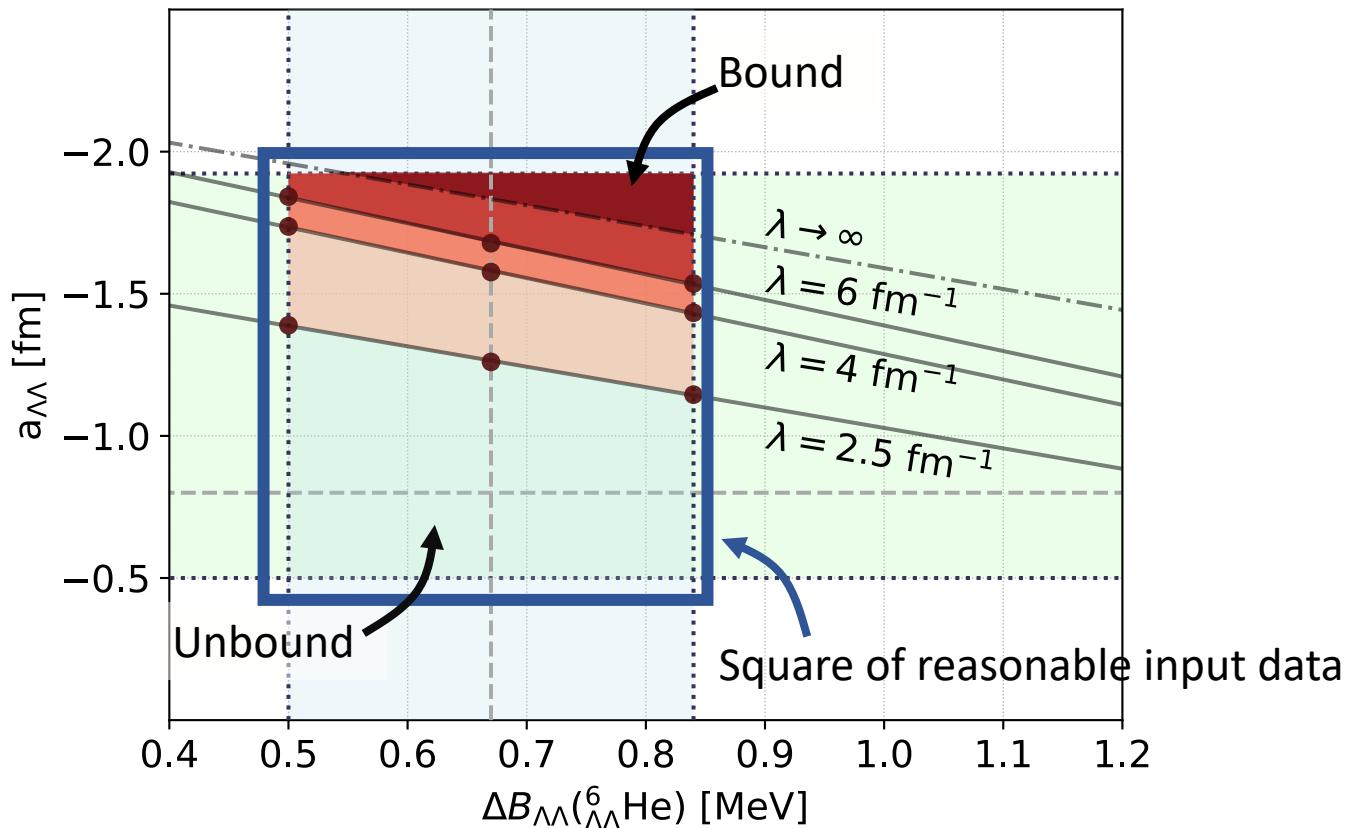
$$\lambda(r_0^{(\Lambda 2)}) = 1.48 \text{ fm}^{-1}.$$

$^5_{\Lambda\Lambda}\text{H}$ 

$$B_\Lambda(\Lambda\Lambda^5\text{H}) = 1.14(1)^{+(44)}_{-(26)} \text{ MeV}$$



$^4_{\Lambda\Lambda}\text{H}$ 

$^4_{\Lambda\Lambda}\text{H}$ 


$$\Delta B_{\Lambda\Lambda}(^6_{\Lambda\Lambda}\text{He}) = B_{\Lambda\Lambda}(^6_{\Lambda\Lambda}\text{He}) - 2 B_{\Lambda}(^5_{\Lambda}\text{He})$$

$^4_{\Lambda\Lambda}\text{H}$  is bound/unbound depending  
to the theory input

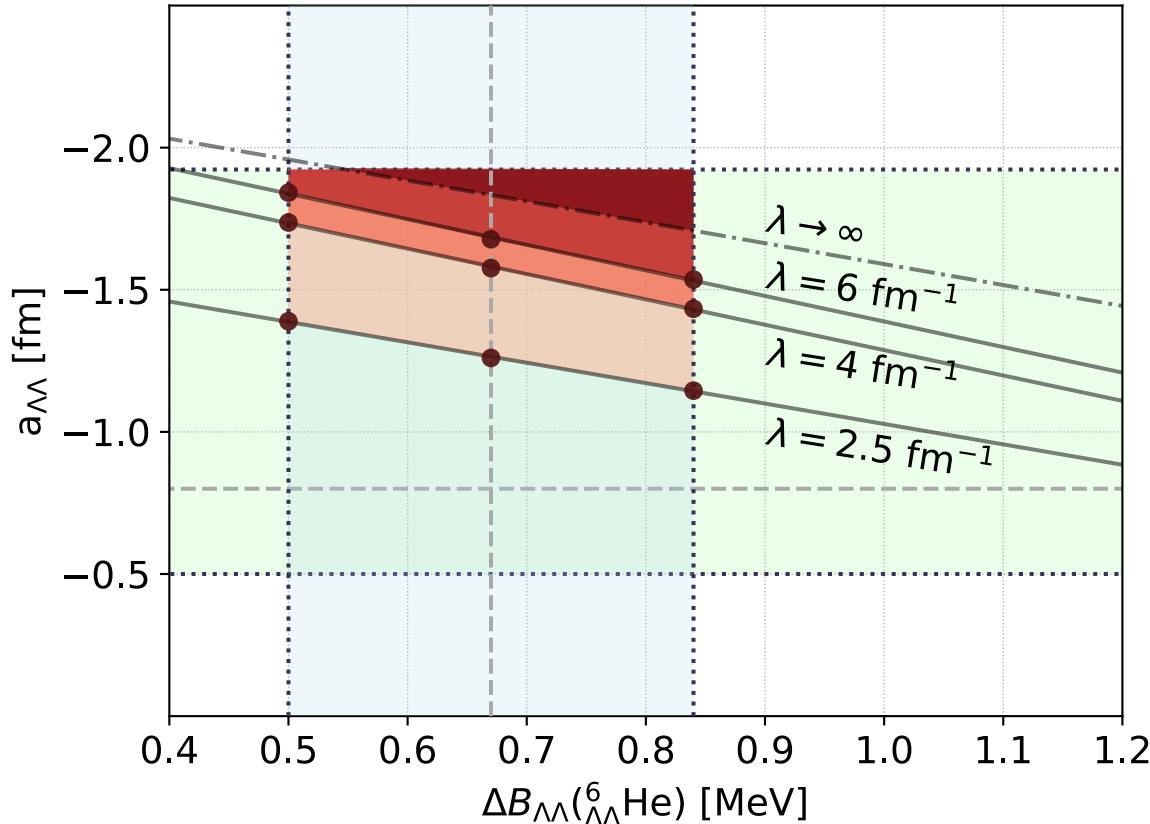
$^4_{\Lambda\Lambda}H$ 


TABLE I:  $\Lambda$  separation energies  $B_\Lambda(^A_{\Lambda\Lambda}Z)$  for  $A=3-6$ , calculated using  $a_{\Lambda\Lambda}=-0.8$  fm, cutoff  $\lambda=4$  fm $^{-1}$  and the Alexander[B]  $\Lambda N$  interaction model [18]. In each row a  $\Lambda\Lambda N$  LEC was fitted to the underlined binding energy constraint.

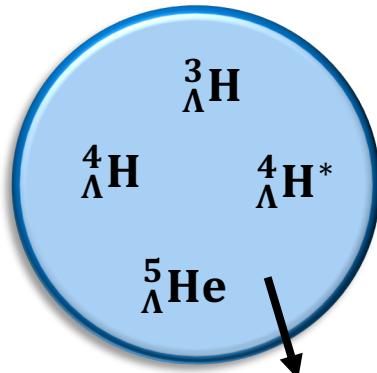
Constraint (MeV)	$^3_{\Lambda\Lambda}n$	$^4_{\Lambda\Lambda}n$	$^4_{\Lambda\Lambda}H$	$^5_{\Lambda\Lambda}H$	$^6_{\Lambda\Lambda}He$
$\Delta B_{\Lambda\Lambda}(^6_{\Lambda\Lambda}He)=\underline{0.67}$	—	—	—	1.21	3.28
$B_\Lambda(^4_{\Lambda\Lambda}H)=\underline{0.05}$	—	—	0.05	2.28	4.76
$B(^4_{\Lambda\Lambda}n)=\underline{0.10}$	—	0.10	0.86	4.89	7.89
$B(^3_{\Lambda\Lambda}n)=\underline{0.10}$	0.10	15.15	18.40	22.13	25.66

Cortona 2019 - Lorenzo Contessi

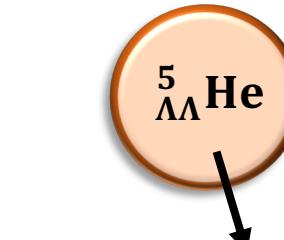
**Not bound**

$nn\Lambda$   
 $n\Lambda\Lambda$   
 $nn\Lambda\Lambda$

# In a nutshell

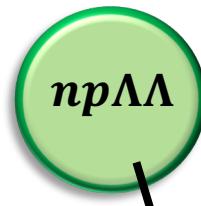


It is possible to describe them all together.  
 ( No overbinding problem! )

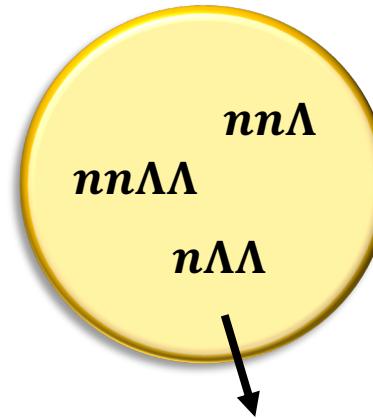


Solidly bound!

$$B_\Lambda(^5_{\Lambda\Lambda} He) = 1.14(1)^{+(44)}_{-(26)}$$



Bound?  
 Need more precise  
 $\Lambda\Lambda$  scattering data!



Unbound

# Conclusions

## General

- $\pi$ -EFT can be applied **successfully** to  $\Lambda$  hypernuclei:  
(no catastrophic failure, truncation error of  $\sim 10\%$  at LO).
- 7 new input data that **can** be fix on **experimental** data!
- **Overcomes overbinding** problem (comprehensive description of  $A \leq 5$   $\Lambda$ -hyperons)

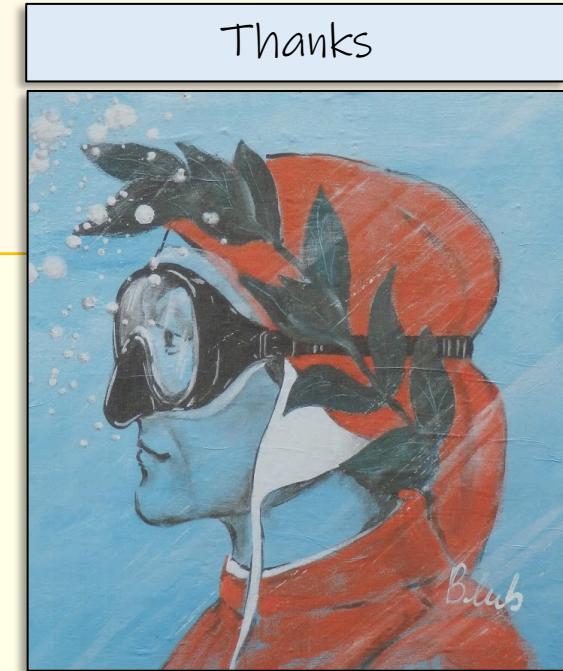
## Predictions

- No boundstate in  $nn\Lambda$ ,  $np\Lambda$  ( $S = \frac{3}{2}$ ),  $n\Lambda\Lambda$  or  $nn\Lambda\Lambda$
- $np\Lambda\Lambda$  might be bound for large  $a_{\Lambda\Lambda} < -1.5$  fm
- ${}^5_{\Lambda\Lambda}\text{He}$  bound ( $B({}^5_{\Lambda\Lambda}\text{He}) = 1.14(1)^{+(44)}_{-(26)} \text{ MeV}$ )

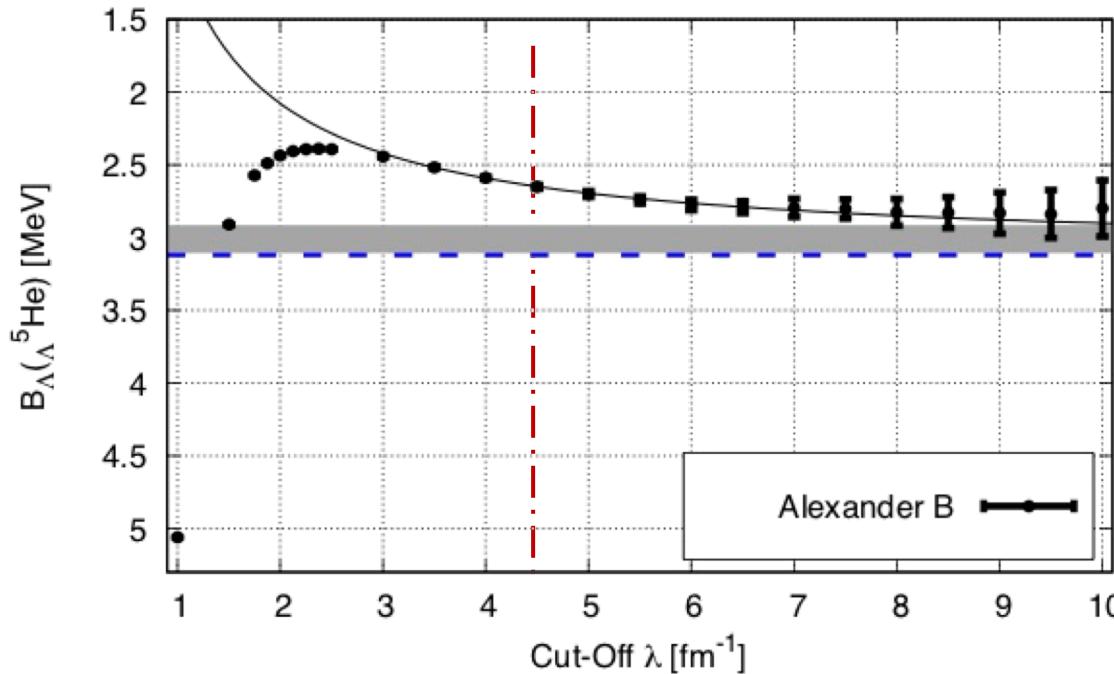
## Prospective

- **Extend** this approach to  $A > 6$  systems.
- Nuclear **NLO**.
- Include **subleading contributions**  
(explicit  $\Xi$  mixing, effective range, .. ).

Thanks

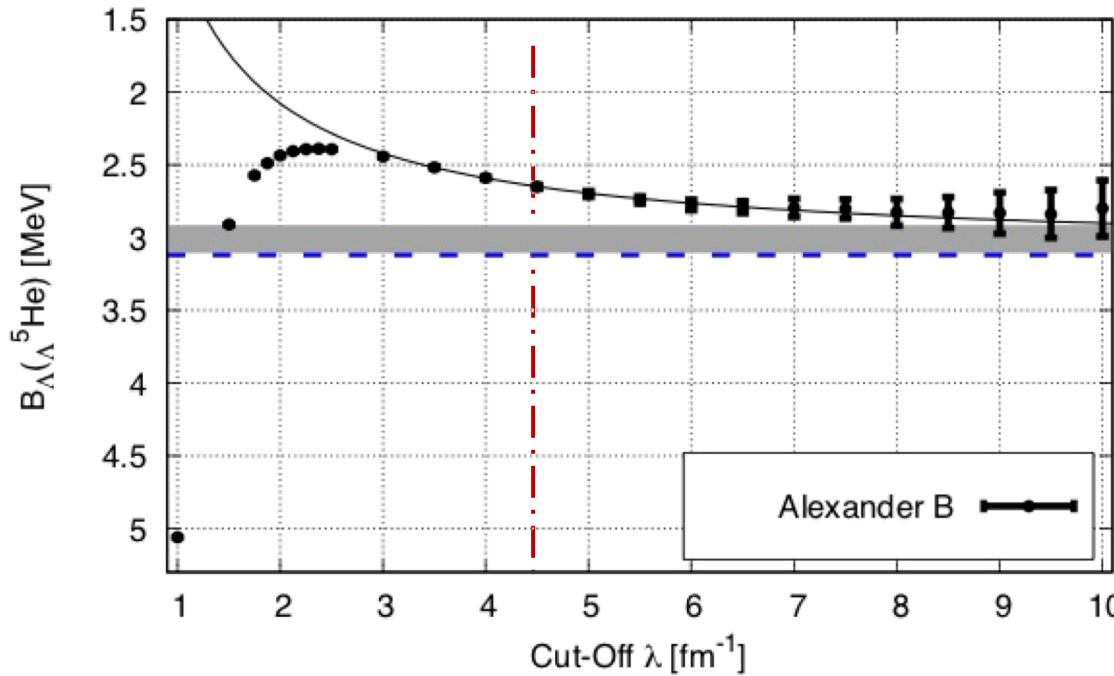


# Error sources



<b>Numeric error</b>	(From computational method)	$\sim 0.01 \text{ MeV}$
<b>Extrapolation uncertainty</b>	$(O_\lambda = O_\infty + \frac{\alpha}{\lambda} + \frac{\beta}{\lambda^2} + \frac{\gamma}{\lambda^3} + \dots)$	From $\lambda = 4 \text{ fm}^{-1} \sim 0.1 \text{ MeV}$
<b>Fitting uncertainty</b>	(Random error of the fit procedure)	$\sim 0.01 \text{ MeV}$
<b>Experimental uncertainty</b>	(Error propagated from experiments)	$\sim 6 \%$
<b>Truncation error</b>	$(\sim \left(\frac{Q}{M}\right)^2)$	$\sim 9 - 50 \%$ (Depends on the scales)

# Error sources



<b>Numeric error</b>	(From computational method)	$\sim 0.01 \text{ MeV}$
<b>Extrapolation uncertainty</b>	$(O_\lambda = O_\infty + \frac{\alpha}{\lambda} + \frac{\beta}{\lambda^2} + \frac{\gamma}{\lambda^3} + \dots)$	From $\lambda = 4 \text{ fm}^{-1} \sim 0.1 \text{ MeV}$
<b>Fitting uncertainty</b>	(Random error of the fit procedure)	$\sim 0.01 \text{ MeV}$
<b>Experimental uncertainty</b>	(Error propagated from experiments)	$\sim 6 \%$
<b>Truncation error</b>	$(\sim \left(\frac{Q}{M}\right)^2)$	$\sim 9\% \text{ (Depends on the scales)}$

# $^5_{\Lambda}\text{He}$ : Overbinding problem

	$B_{\Lambda}(^3_{\Lambda}H)$	$B_{\Lambda}(^4_{\Lambda}H_{g.s.})$	$B_{\Lambda}(^4_{\Lambda}H_{exc.})$	$B_{\Lambda}(^5_{\Lambda}He)$
Exp.	0.13(5) [4]	2.16(8) [5]	1.09(2) [6]	3.12(2) [4]
DHT [7]	0.10	2.24	0.36	$\geq 5.16$
AFDMCa	-	1.97(11) [8]	-	5.1(1) [9]
AFDMCb'	0.23(9) [13]	1.95(9) [13]	-	2.60(6) [13]
$\chi$ EFTa	0.11 [10]	2.31 (3) [11]	0.95(15) [11]	5.82(2) [12]
$\chi$ EFTb	-	2.13 (3) [11]	1.39(15) [11]	4.43(2) [12]

All the energies are in MeV.

- [7] R.H. Dalitz, R.C. Herndon, and Y.C. Tang, Nucl. Phys. B 47, 109 (1972).
- [8] D. Lonardoni, F. Pederiva, and S. Gandolfi, Phys. Rev. C 89, 014314 (2014).
- [9] D. Lonardoni, S. Gandolfi, and F. Pederiva, Phys. Rev. C 87, 041303(R) (2013).
- [10] R. Wirth et al., Phys. Rev. Lett. 113, 192502 (2014).
- [11] D. Gazda and A. Gal, Phys. Rev. Lett. 116, 122501 (2016); D. Gazda and A. Gal, Nucl. Phys. A 954, 161 (2016).
- [12] R. Wirth and R. Roth, Phys. Lett. B 779, 336 (2018). We thank Roland Wirth for providing us with these values.
- [13] D. Lonardoni arXiv:1711.07521v2 & Private communication.
- [15] H. Nemura, Y. Akaishi, and Y. Suzuki, Phys. Rev. Lett. 89, 142504 (2002); see also Y. Akaishi, T. Harada.

# $N\bar{\Lambda}$ scattering data

**Alexander et al. :**

$$a_s = -1.8 \text{ fm}$$

$$a_t = -1.6 \text{ fm}$$

**Sechi-Zorn et al. :**

$$0 > a_s > -9 \text{ fm}$$

$$-0.8 > a_t > -3.2 \text{ fm}$$

G. Alexander, U. Karshon, A. Shapira, et al. Phys. Rev. 173, 1452 (1968)

Sechi-Zorn, B., B. Kehoe, J. Twitty, and R. A. Burnstein, 1968, Phys. Rev. 175, 1735.

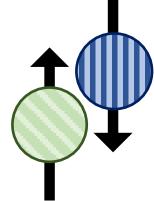
TABLE VII  $\Lambda N$  scattering lengths and effective ranges (in fm) for several  $YN$  interaction models. For the EFT models, these refer to  $\Lambda p$  and to cutoff parameter of 600 MeV.

Model	Reference	$a^s$	$r_0^s$	$a^t$	$r_0^t$
NSC89	Maessen, Rijken, and de Swart (1989)	-2.79	2.89	-1.36	3.18
NSC97e	Rijken, Stoks, and Yamamoto (1999)	-2.17	3.22	-1.84	3.17
NSC97f	Rijken, Stoks, and Yamamoto (1999)	-2.60	3.05	-1.71	3.33
ESC08c	Nagels, Rijken, and Yamamoto (2015b)	-2.54	3.15	-1.72	3.52
Jülich '04	Haidenbauer and Meißner (2005)	-2.56	2.75	-1.66	2.93
EFT (LO)	Polinder, Haidenbauer, and Meißner (2006)	-1.91	1.40	-1.23	2.20
EFT (NLO)	Haidenbauer <i>et al.</i> (2013)	-2.91	2.78	-1.54	2.72

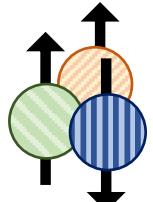
A. Gal et al. - Strangeness in nuclear physics - Rev.Mod.Phys. 88 (2016) no.3, 035004

LEC	State	Fitting	
$C_{02}$	$S = 1, I = 0$	$^2\text{H}$	✓ Boundstate
$C_{20}$	$S = 0, I = 1$	$\text{N} - \text{N}$	✓
$C_{01}$	$S = 1, I = \frac{1}{2}$	$\Lambda - \text{N}$	✗ Scattering
$C_{21}$	$S = 0, I = \frac{1}{2}$	$\Lambda - \text{N}$	✗
$C_{00}$	$S = 0, I = 0$	$\Lambda - \Lambda$	✗

**Two body**

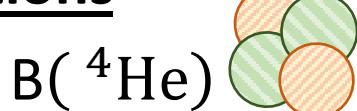
**Three body**



Boundstates

LEC	State	Fitting
$D_{11}$	$S = \frac{1}{2}, I = \frac{1}{2}$	$^3\text{H}$ ✓
$D_{01}$	$S = \frac{1}{2}, I = 0$	$^3\Lambda\text{H}$ ✓
$D_{03}$	$S = \frac{1}{2}, I = 1$	$^4\Lambda\text{H}_{S=0, I=\frac{1}{2}}$ ✓
$D_{21}$	$S = \frac{3}{2}, I = 0$	$^4\Lambda\text{H}_{S=1, I=\frac{1}{2}}$ ✓
$D_{11}^{\Lambda\Lambda N}$	$S = \frac{1}{2}, I = \frac{1}{2}$	$^6\Lambda\Lambda\text{He}$ ✓

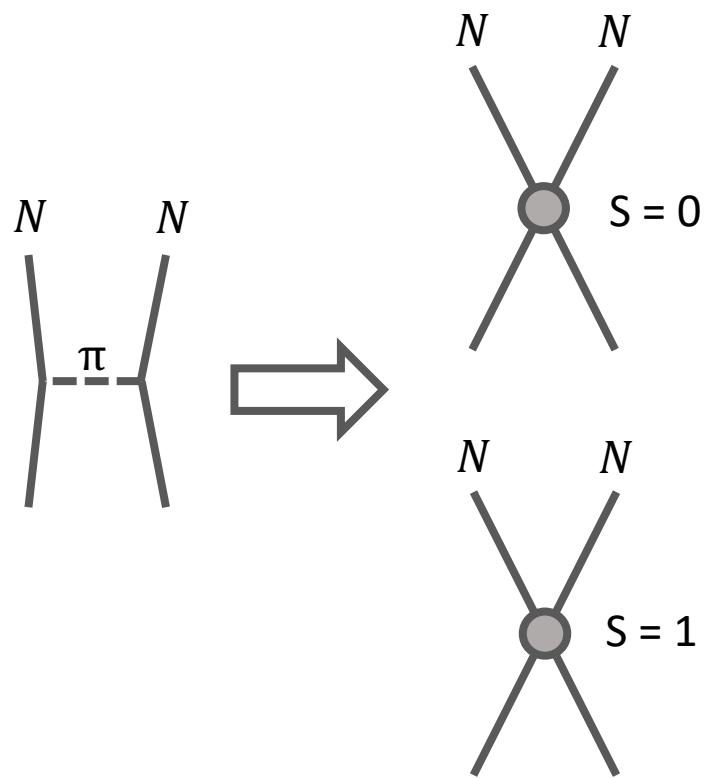
## Predictions



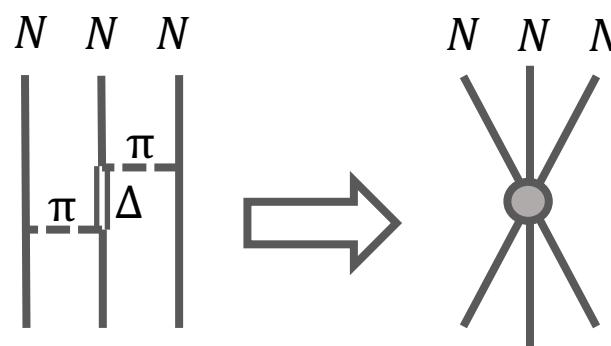
# $\pi$ -EFT ( $N$ )

$M$  = Theory break-scale  
 $Q$  = Typical exchanged momentum  
 $B$  = Typical binding per particle<sup>30</sup>

## 2-Body



## 3-Body



$$B \sim 7 \text{ MeV}$$

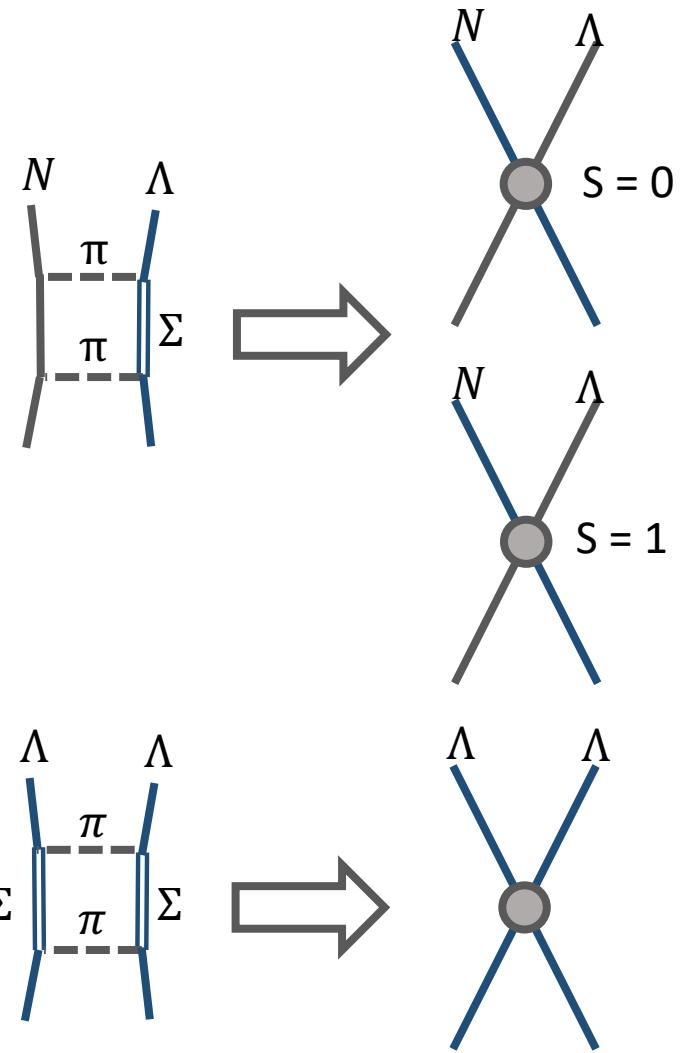
$$M \sim m_\pi$$

$$\delta LO = \left( \frac{Q}{M} \right) \sim 50\%$$

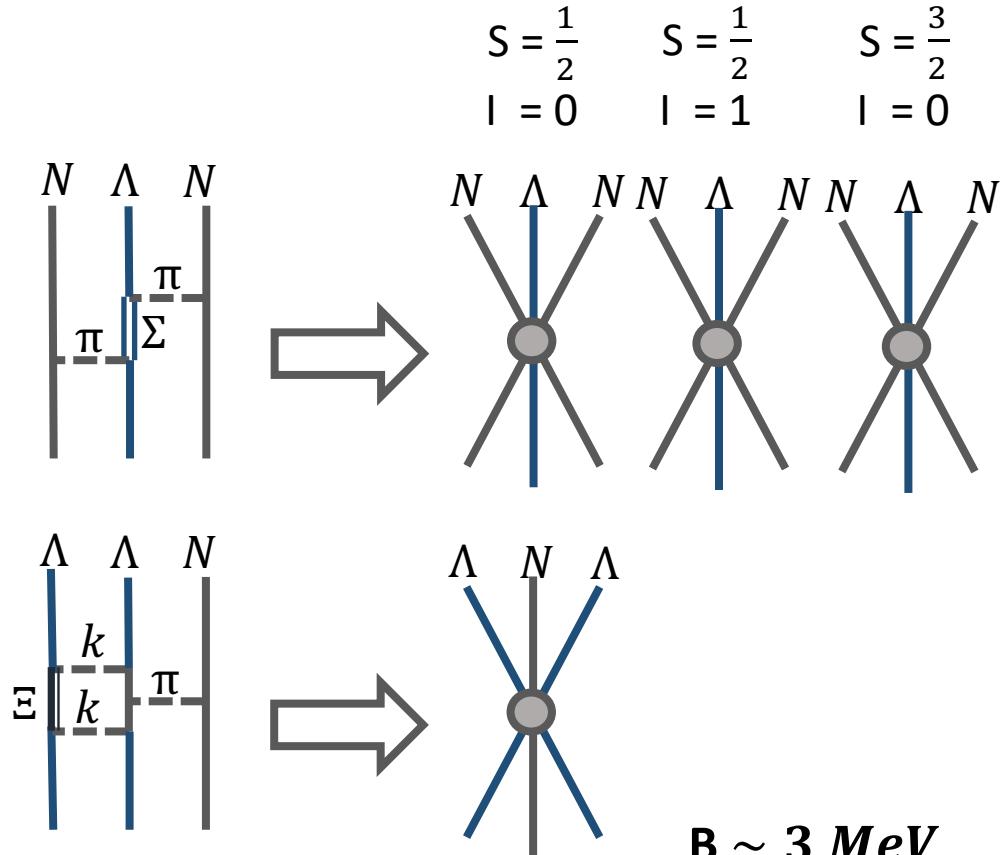
# $\pi$ -EFT ( $\Lambda$ )

$M$  = Theory break-scale  
 $Q$  = Typical exchanged momentum  
 $B$  = Typical binding per particle

## 2-Body



## 3-Body



$$B \sim 3 \text{ MeV}$$

$$M \sim 2 m_\pi$$

$$\delta LO = \left( \frac{Q}{M} \right)^2 \sim 9\%$$

$$V_{2b}^\lambda = \sum_{ij} e^{-\left(\frac{r_{ij}\lambda}{2}\right)^2} [C_{10}^\lambda P_{[S=1,I=0]}^{NN} + C_{01}^\lambda P_{[S=0,I=1]}^{NN}]$$

### Three body force

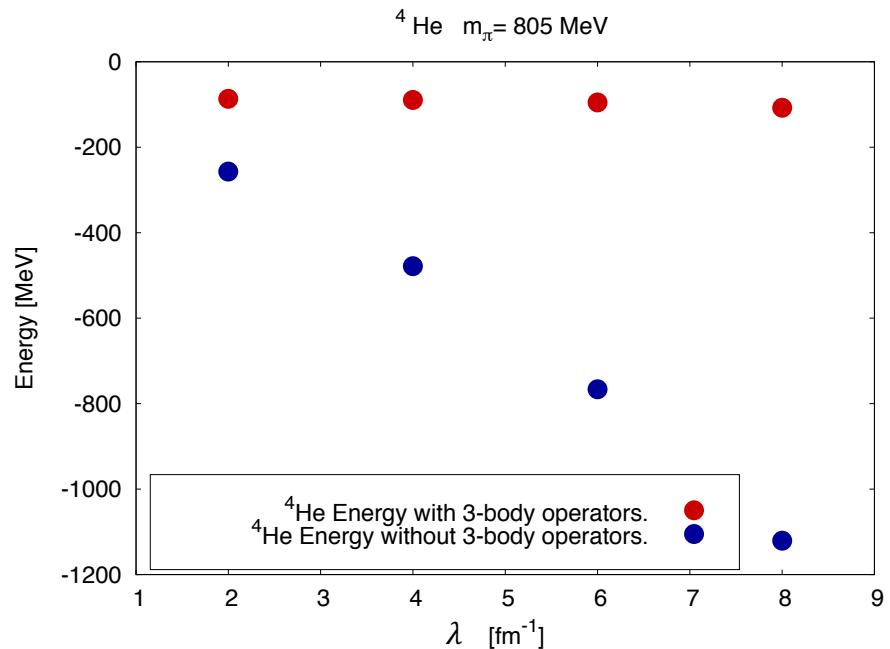
The energy of the  $A \geq 3$  depends on  $\lambda$

When  $\lambda \rightarrow \infty \rightarrow E_{3b} \rightarrow -\infty$

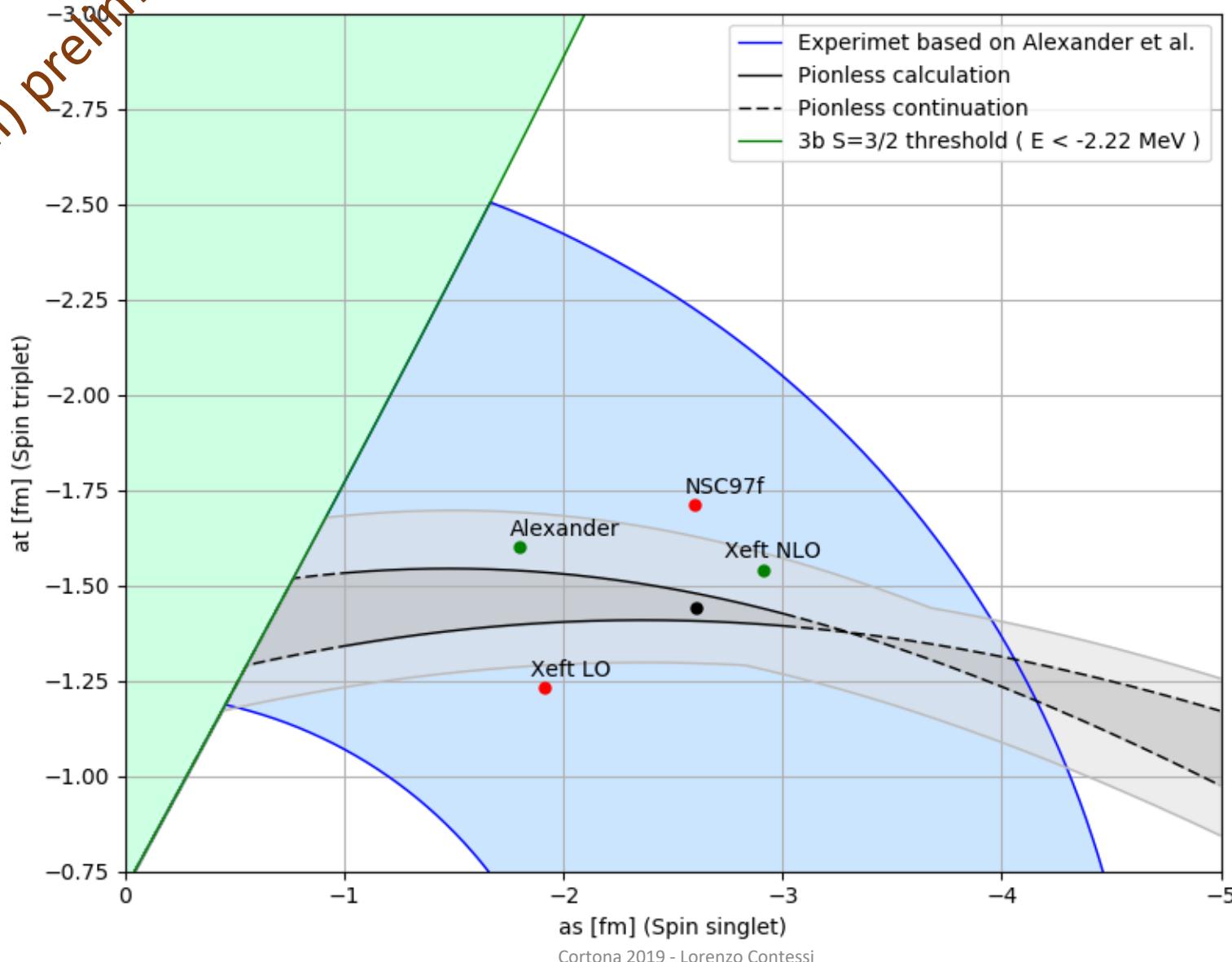
**NO PREDICTIVE POWER!**

**Three-body** force stabilize  
 $A \geq 3$  systems

The three body force require  
Regularization/Renormalization as well.



(Still) preliminary



(Still) preliminary

