

CP violation in hadronic τ decays

Oscar Vives

U. València and IFIC



XI SuperB Workshop

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- Another possibility is to use a triple product in the angular distribution which does not need a strong phase

CP violation

In the SM, τ decays through a virtual W into $f + \nu_\tau$.

$$\mathcal{H}_{\text{eff}}^{\text{SM}} = \frac{G_F}{\sqrt{2}} \sin \theta_c \bar{\nu}_\tau \gamma_\mu (1 - \gamma_5) \tau \quad \bar{s} \gamma^\mu (1 - \gamma_5) u + \text{h.c.},$$

the hadronic matrix element $J^\mu \equiv \langle f | \bar{s} \gamma^\mu (1 - \gamma_5) u | 0 \rangle$ has several form factors for different final states.

We need a different contribution to the effective Hamiltonian with different weak phases.

Moreover a different strong phase requires two different hadronic intermediate states, different hadronic operators.

Phenomenological analysis of CP violation in τ decays

1. G. Bonvicini *et al.* [CLEO Collaboration], “Search for CP Violation in $\tau^- \rightarrow K\pi\nu_\tau$ Decays,” *Phys. Rev. Lett.* **88**, 111803 (2002) [arXiv:hep-ex/0111095].
2. A. Datta, K. Kiers, D. London, P. J. O’Donnell and A. Szykman, “CP violation in hadronic tau decays,” *Phys. Rev. D* **75** (2007) 074007 [Erratum-ibid. *D* **76** (2007) 079902] [arXiv:hep-ph/0610162].
3. K. Kiers, K. Little, A. Datta, D. London, M. Nagashima and A. Szykman, “CP violation in $\tau^- \rightarrow K\pi\pi\nu_\tau$,” *Phys. Rev. D* **78** (2008) 113008 [arXiv:0808.1707 [hep-ph]].

Literature

4. D. Kimura, K. Y. Lee, T. Morozumi and K. Nakagawa, “CP violation of $\tau^- \rightarrow K\pi(\eta, \eta')\nu$ decays,” arXiv:0808.0674 [hep-ph].
5. D. Kimura, K. Nakagawa, T. Morozumi and K. Y. Lee, “Direct CP Violation In Hadronic Tau Decays,” Nucl. Phys. Proc. Suppl. **189** (2009) 84.
6. D. Delepine, “CP violation in semi-leptonic tau decays,” AIP Conf. Proc. **917** (2007) 90 [arXiv:hep-ph/0702107].

K pi pi nu

Assume an additional charged Higgs contribution.

$$\mathcal{H}_{\text{eff}}^{\text{NP}} = \frac{G_F}{\sqrt{2}} \sin \theta_c [\eta_S \bar{\nu}_\tau (1 + \gamma_5) \tau \bar{S} U + \eta_P \bar{\nu}_\tau (1 + \gamma_5) \tau \bar{S} \gamma_5 U]$$

NP effects incorporated in the SM scalar form factor:

$$B_4 \rightarrow \tilde{B}_4 = B_4 + \frac{f_H}{m_\tau} \eta_P,$$

1. Rate asymmetry

$$A_{CP}^{(0)} = \frac{\Gamma - \bar{\Gamma}}{\Gamma + \bar{\Gamma}} \propto \frac{B_4 f_H \eta_P \sin(\delta_4 - \delta_H) \sin \phi_H}{2\Gamma}$$

but B_4 is small...rate asymmetry expected to be small...

K pi pi nu

2. Modified rate asymmetry and triple product asymmetry

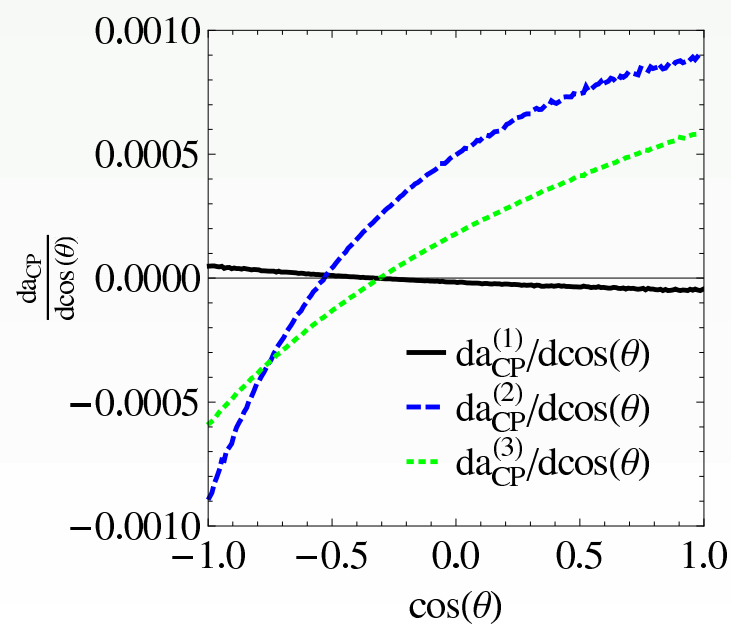
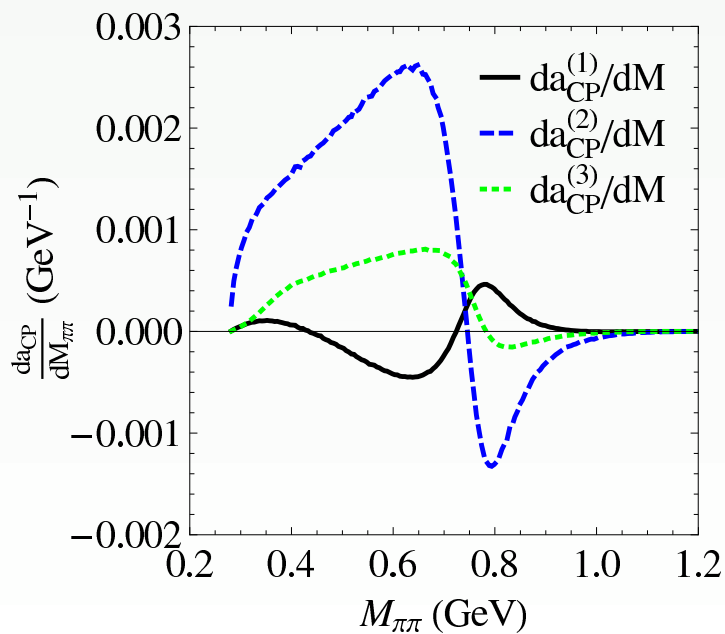
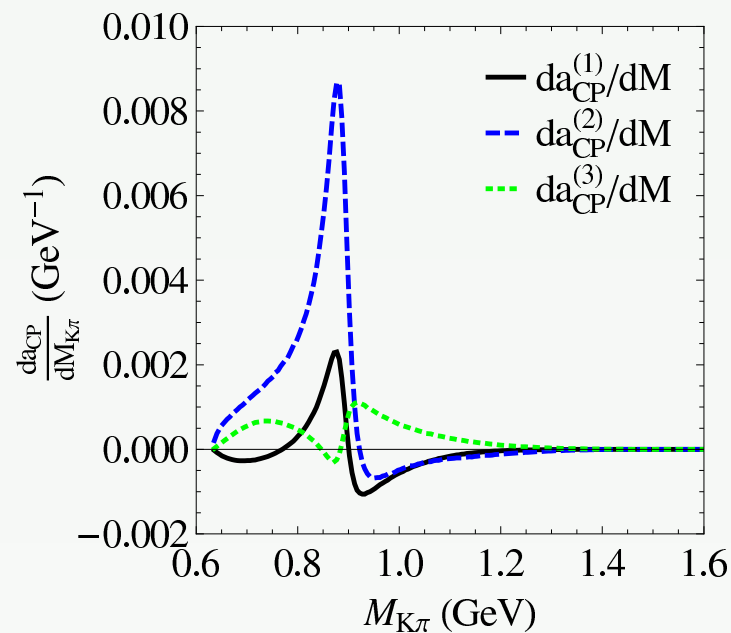
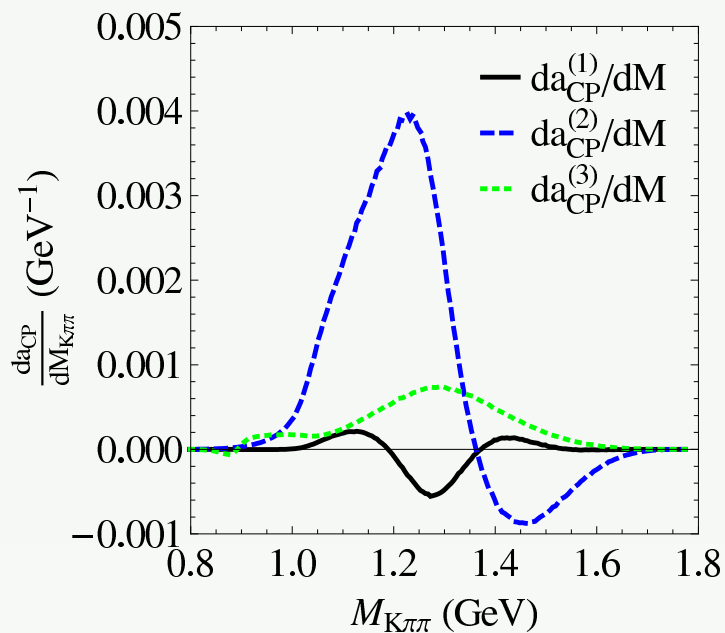
$$A_{CP}^{(i)} = \frac{1}{\Gamma + \bar{\Gamma}} \int \left(\frac{d\Gamma_i}{dQ^2 ds_1 ds_2} - \frac{d\bar{\Gamma}_i}{dQ^2 ds_1 ds_2} \right) dQ^2 ds_1 ds_2 .$$

with

$$\frac{d\Gamma_1}{dQ^2 ds_1 ds_2} = A(Q^2) \left[\langle \bar{K}_3 \rangle \text{Re}(B_1 B_3^*) - \langle \bar{K}_2 \rangle \text{Re}(B_2 B_4^*) \right] ,$$

$$\frac{d\Gamma_2}{dQ^2 ds_1 ds_2} = A(Q^2) \left[\langle \bar{K}_3 \rangle \text{Re}(B_2 B_3^*) + \langle \bar{K}_2 \rangle \text{Re}(B_1 B_4^*) \right] ,$$

$$\frac{d\Gamma_3}{dQ^2 ds_1 ds_2} = A(Q^2) \left[\langle \bar{K}_3 \rangle \text{Im}(B_1 B_2^*) + \langle \bar{K}_2 \rangle \text{Im}(B_3 B_4^*) \right] ,$$



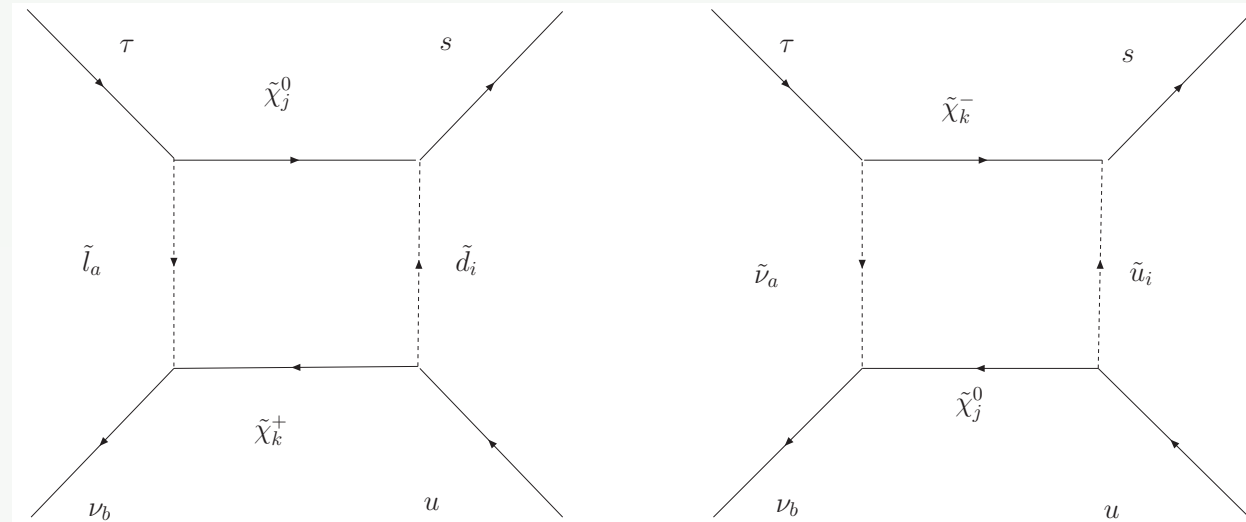
K pi pi nu

where the $A_{CP}^{(i)}$ are obtained from $a_{CP}^{(i)}$ multiplying it by $f_H \text{Im}(\eta_P)$.

Taking $\eta_P \simeq 1$ and phases $O(1)$, we get $f_H \text{Im}(\eta_P) \simeq 10$.

\Rightarrow CP asymmetries visible with
“large” non-standard Higgs couplings...

Tree vs. Loop



Example

Estimate :

$$C_{scal}^{SUSY} / C^{SM} \sim \frac{g^4 y_\tau (\delta_{LR}^d)_{32} V_{13}^{CKM}}{16\pi^2 G_F M_{SUSY}^2 V_{12}^{CKM}}$$

$$\sim 10^{-5} (\delta_{LR}^d)_{32} \tan \beta \frac{M_W^2}{M_{SUSY}^2} \leq 10^{-6}$$