

# Constraining light dark matter scenarios using long-lived particles

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NCBJ

23/06/2019



From 1807.00971 and 1907.xxxx

# Outline

Exploring the properties of sub-GeV dark matter through an example

Signatures of light dark sector at the intensity frontier

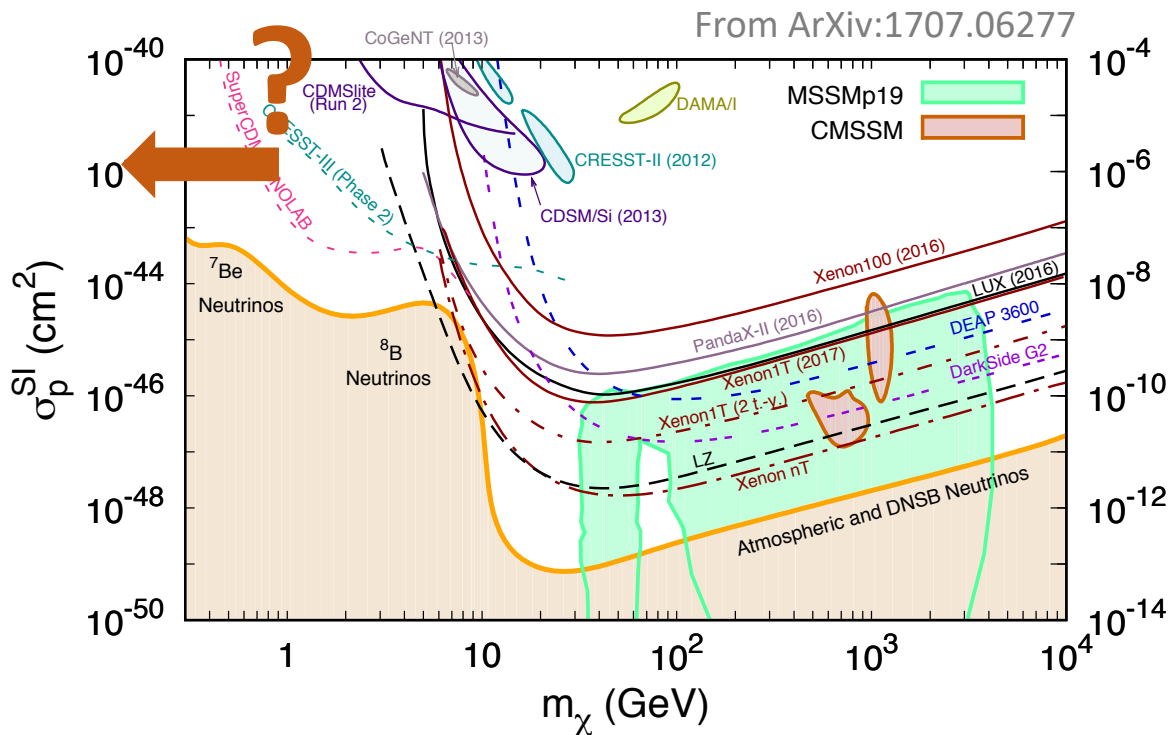
A broader setup: light dark sector effective theory and the intensity frontier

# Introduction

Thermal dark matter in the dark

# Light dark sectors and dark matter

- WIMP Dark matter is a mature area of research
- what about lighter (sub-GeV) DM ?



- Intense experimental effort in the intensity frontier
- $\approx 10$  relevant experiment in next 2 years

Information | References (339) | Citations (182) | Files | Plots

**US Cosmic Visions: New Ideas in Dark Matter 2017: Community Report**

Marco Battaglieri (SAC co-chair) (INFN, Genoa), Alberto Belloni (Coordinator) (Maryland U.), Aaron Chou (WG2 Convener) (Fermilab), Priscilla Cushman (Coordinator) (Minnesota U.), Bertrand Echenard (WG3 Convener) (Caltech), Rouven Essig (WG1 Convener) (SUNY, Stony Brook), Juan Estrada (WG1 Convener) (Fermilab), Jonathan L. Feng (WG4 Convener) (UC, Irvine), Brenna Flaugher (Coordinator), Patrick J. Fox (WG4 Convener) (Fermilab) *et al.* [Show all 251 authors](#)

Jul 14, 2017 - 113 pages

Information | References (355) | Citations (19) | Files | Plots

**Physics Beyond Colliders at CERN: Beyond the Standard Model Working Group Report**

J. Beacham (Ohio State U., Columbus (main)), C. Burrage (U. Nottingham), D. Curtin (Toronto U.), A. De Roeck (CERN), J. Evans (Cincinnati U.), J.L. Feng (UC, Irvine), C. Gatto (INFN, Naples & NIU, DeKalb), S. Gninenko (Moscow, INR), A. Hartin (U. Coll. London), I. Irastorza (U. Zaragoza, LFNAE) *et al.* [Show all 33 authors](#)

Jan 20, 2019 - 150 pages

- Not all DM-motivated, yet most can be used to constrained scenarios of sub-GeV, thermal DM

# What is required for such a light DM?

- Many different approaches, most include additional, light **dark sector**
  - Literature often focuses on finding new mechanisms → with great success

Impeded DM  
1609.02147,...

Co-scattering DM  
1705.08450, 1705.09292,...

iDM  
hep-ph/0101138, ...

Selfish DM  
1504.00361,...

Co-decaying DM  
1607.03110, ...

Secluded DM  
0711.4866, ...

Cannibal DM  
1602.04219, ...

Semi-annihilating DM  
1003.5912, ...

Forbidden DM  
Griest-Seckel, 1505.07107, ...

Boosted DM  
1405.7370, 1503.02669...

<Your choice> DM  
1906.xxxx

→ Seemingly only limited by our ability to invent new names ...

→ ... and by usually **a non-minimal accompanying dark sector**

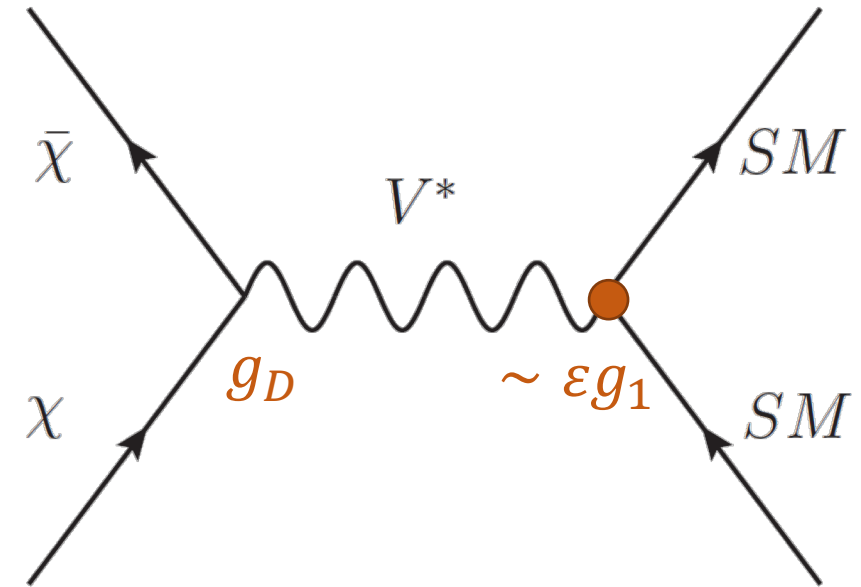
... and many other

# Exploring the properties of sub-GeV dark matter through an example

Model building, Relic density, Astrophysical bounds

# Building a toy-model for sub-GeV DM

- For this talk, let's try to build a simple, self-consistent dark sector model with sub-GeV dark matter → study detection prospects
- Suppose a vector mediator (dark photon), and (mostly in this talk) fermion dark matter
- Try to keep model building SM/WIMP-like then see the various possible regimes



$$\Omega h^2 \sim 0.1 \times \left( \frac{10^{-3}}{\epsilon} \right)^2 \left( \frac{0.1}{\alpha_D} \right) \left( \frac{25 \text{ MeV}}{M_\chi} \right)^2 \left( \frac{M_V}{75 \text{ MeV}} \right)^4$$

# Kinetic mixing and dark Higgs mechanism

- Coupling to SM obtained through “kinetic mixing” term

$$\mathcal{L}_{A'} = -\frac{1}{4} F'^{\mu\nu} F'_{\mu\nu} - \frac{1}{2} \frac{\epsilon}{\cos \theta_w} B_{\mu\nu} F'^{\mu\nu} + (D^\mu S)^* (D_\mu S) + \mu_S^2 |S|^2 - \frac{\lambda_S}{2} |S|^4$$

Kinetic mixing term

Dark Higgs potential, with SSB

The diagram illustrates a vertex where a dark photon (represented by a wavy line) couples to a Dirac fermion (represented by two fermion lines,  $\bar{\chi}$  and  $\chi$ ) and Standard Model particles (represented by two fermion lines, both labeled  $SM$ ). The vertex is labeled  $V^*$ . An orange arrow points from the text 'Kinetic mixing term' to the vertex, and a blue arrow points from the text 'Dark Higgs potential, with SSB' to the potential term in the equation.

- Anomaly cancellation -> Introduce a Dirac fermion dark matter  
 $\chi = (\chi_L, \bar{\chi}_R)$



# Fermion dark matter example

$$\mathcal{L}_{pDF}^{\text{DM}} = \bar{\chi} (i\not{D} - m_\chi) \chi + y_{SL} S \bar{\chi}^c P_L \chi + y_{SR} S \bar{\chi}^c P_R \chi + \text{h.c.}$$

- Yukawa couplings to the dark Higgs  $S$

→ Avoid Dirac DM (CMB exclusion)

→ After  $U(1)_D$  symmetry breaking, the dark matter acquires a Majorana mass

$$M_\chi = \begin{pmatrix} \sqrt{2} v_S y_{SL} & m_\chi \\ m_\chi & \sqrt{2} v_S y_{SR} \end{pmatrix}$$

$$M_V = g_{\alpha_D} q_S v_S \quad V$$

$$M_S = \sqrt{2} \lambda_S v_S \quad S$$

$$M_{\chi_2} - M_{\chi_1} = \sqrt{2} v_S (y_{SR} + y_{SL}) \quad \begin{matrix} \chi_2 \\ \chi_1 \end{matrix}$$

- After diagonalization → two Majorana fermions

# The case for small Higgs/Dark Higgs mixing

- Both the SM and dark Higgs potential must be minimised **simultaneously**

$$\mathcal{L} \supset \frac{\lambda_{SH}}{2} |S|^2 |H|^2 \quad \longrightarrow \quad \begin{aligned} v_S^2 &= \frac{1}{\lambda_S} \left( \mu_S^2 - \frac{\lambda_{SH}}{2\lambda_H} \mu_H^2 \right) \\ v_H^2 &\simeq \frac{\mu_H^2}{\lambda_H} \end{aligned}$$

- In order to avoid tuning the bare mass  $\mu_S$  against the SM Higgs contribution, this implies

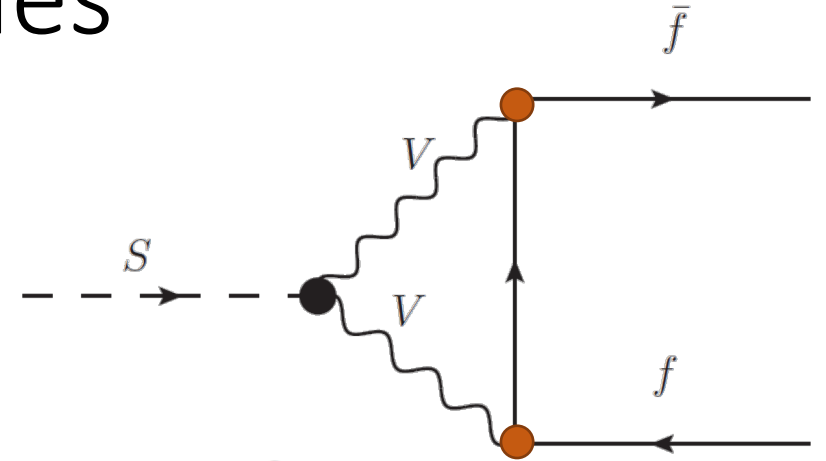
$$\lambda_{SH} \sim \frac{M_S^2}{v_H^2} \sim 10^{-8} - 10^{-6}$$

$$\tau_{S,H\text{mix}} \propto 1 \cdot 10^6 \text{ s} \times \left( \frac{100 \text{ MeV}}{M_S} \right) \left( \frac{100 \text{ MeV}}{M_V} \right)^2 \left( \frac{10^{-6}}{\lambda_{SH}} \right)^2 \left( \frac{q_S^2 \alpha_D}{\alpha_{\text{em}}} \right)$$

# A key aspect – Long-lived particles

1. Decay to SM particles is the only option

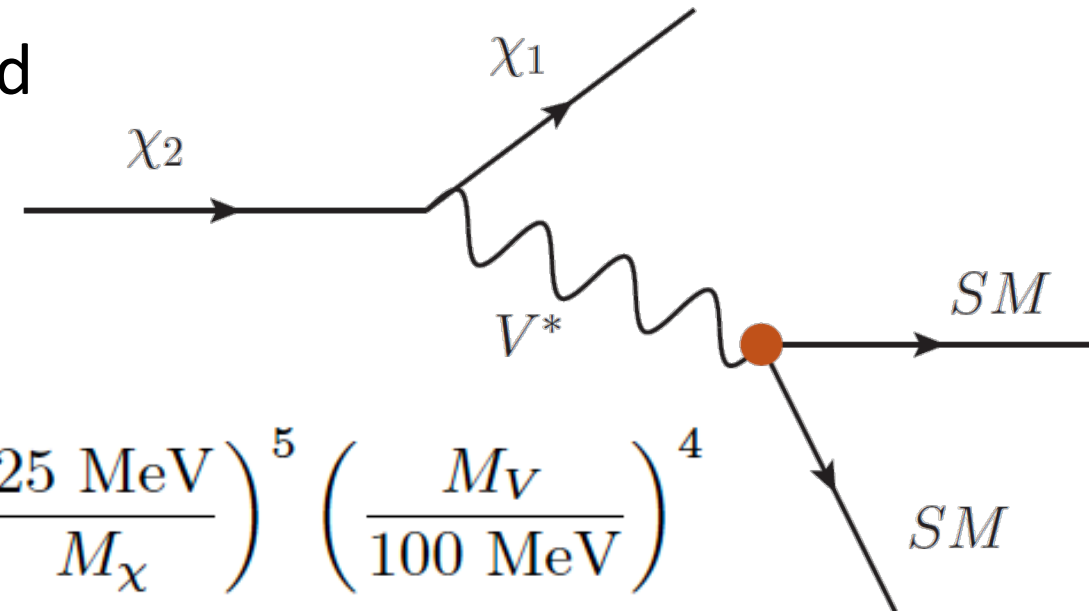
-> For instance **dark Higgs boson, dark photon**



$$\tau_S \propto 1 \text{ s} \times \left(\frac{\alpha'}{\alpha}\right) \times \left(\frac{10^{-3}}{\varepsilon}\right)^4 \left(\frac{100 \text{ MeV}}{M_S}\right) \left(\frac{M_V}{100 \text{ MeV}}\right)^2$$

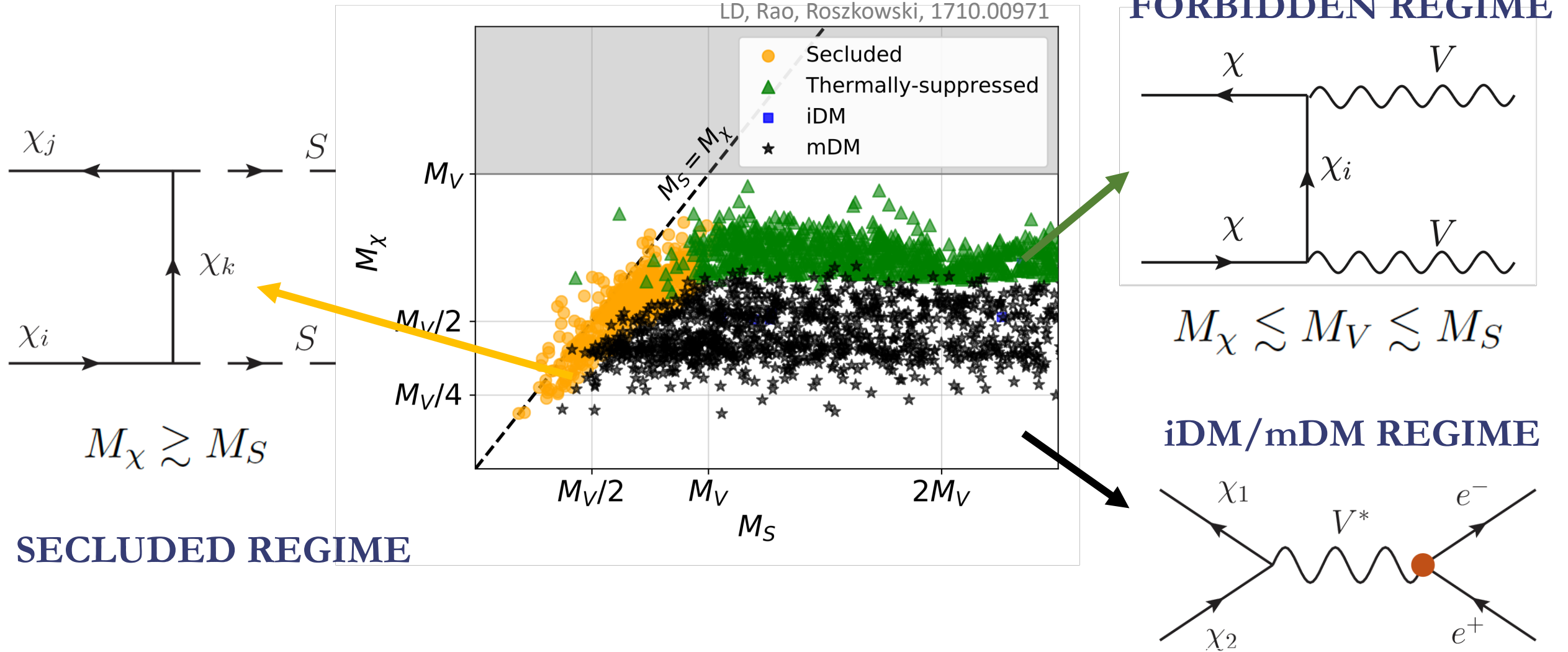
2. Internal dark sector decays must proceed through off-shell mediator

- For instance **dark Higgs boson**
- **heavy dark sector state**



$$c\tau_{\chi_2} \propto 100 \text{ m} \times \left(\frac{0.1}{\alpha_D}\right) \left(\frac{10^{-3}}{\varepsilon}\right)^2 \left(\frac{0.2M_\chi}{\Delta_\chi}\right)^5 \left(\frac{25 \text{ MeV}}{M_\chi}\right)^5 \left(\frac{M_V}{100 \text{ MeV}}\right)^4$$

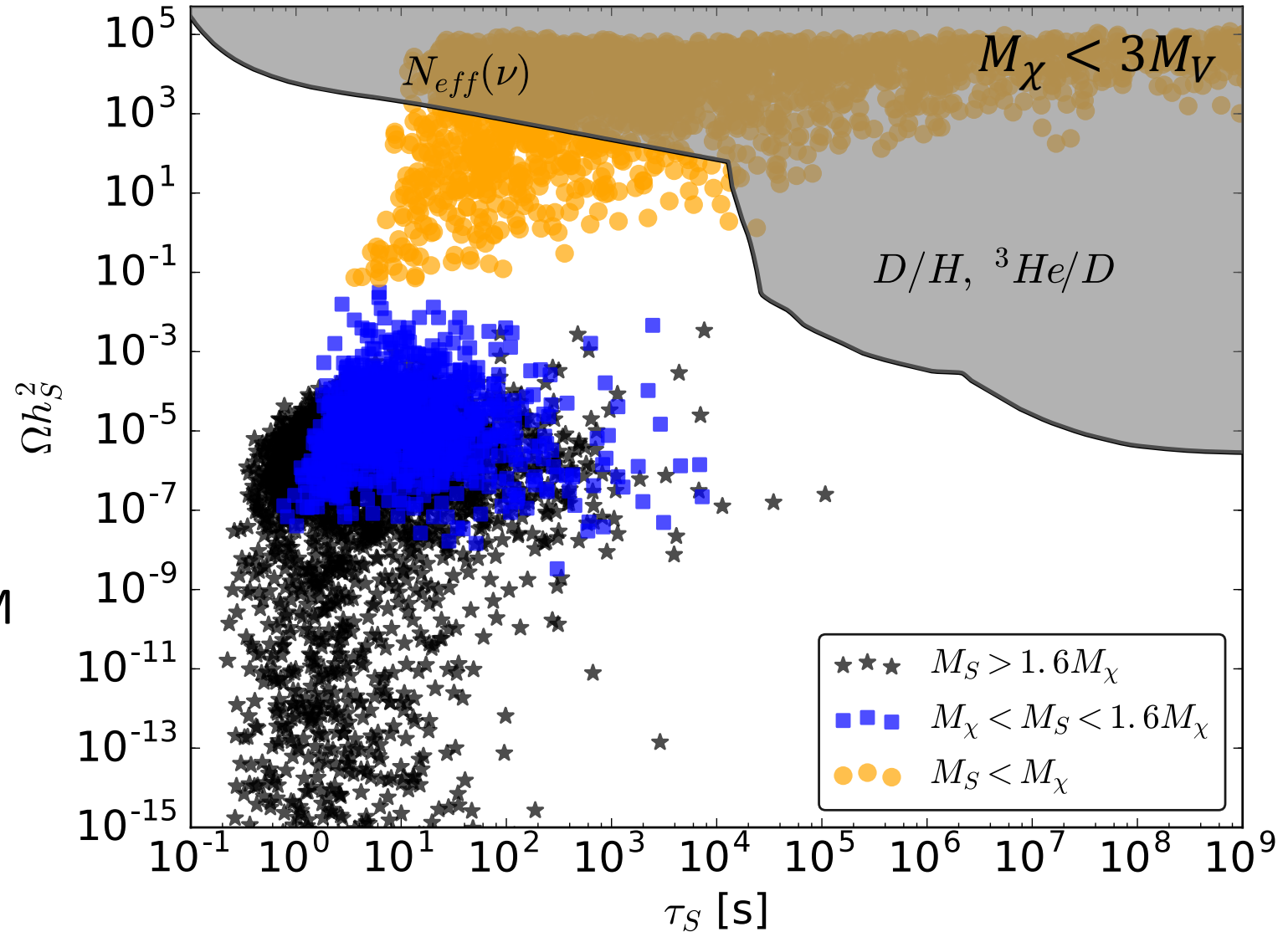
# Typical regimes with correct relic density



# BBN constraints

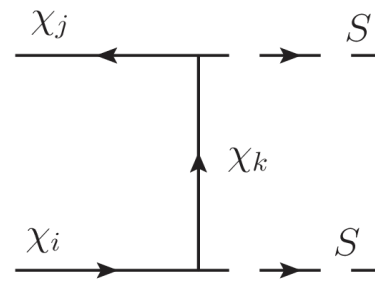
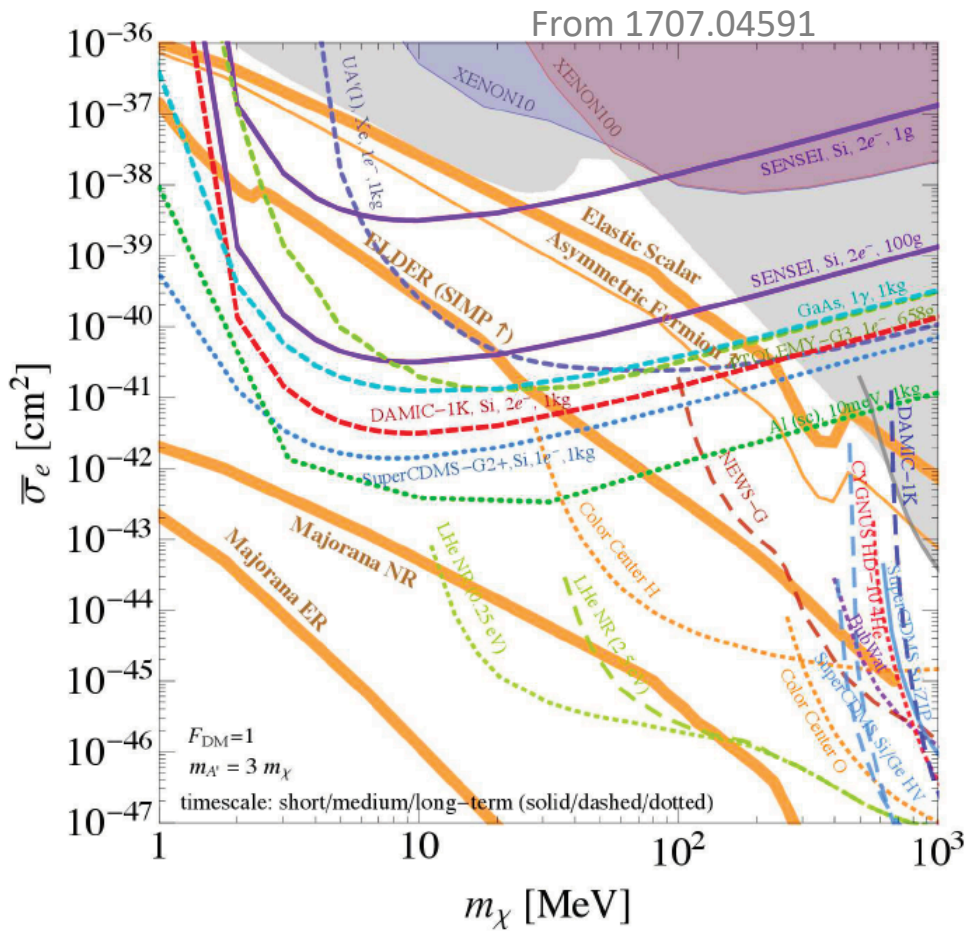
- Potentially large  $S$  metastable density after freeze-out
- Two relevant bounds
  - $N_{eff}$  neutrino from late time energy injection ( $\tau > 0.1s$ )
  - Light element abundances, e.g  $D/H$  ( $\tau > 10^4s$ )
  - CMB bounds on decaying DM
- Light dark Higgs only decay electromagnetically (electron-positron pairs)  $\rightarrow$  (relatively) weak bounds

LD, Rao, Roszkowski, 1710.00971

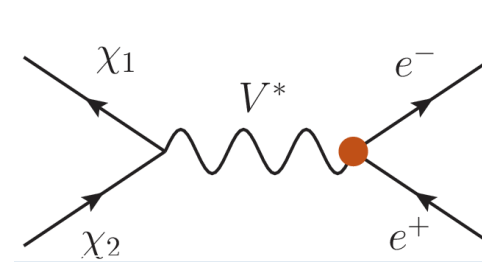


# Direct and indirect detection bounds

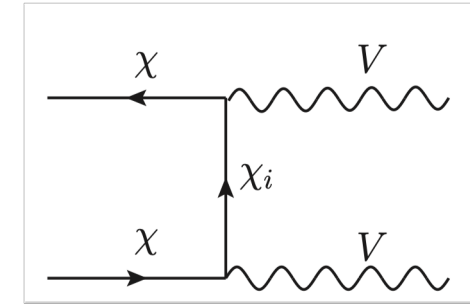
- Late-time annihilation
  - Standard CMB bound -> **s-wave annihilation channel forbidden**
  - either p-wave channel or irrelevant at CMB (e.g. co-annihilation)



P-wave



Co-annihilation



Thermal suppression

- Direct detection, **many** upcoming experiments
  - Standard searches currently only **relevant for scalar DM** case (since either inelastic DM or Majorana DM)

# Signatures of light dark sector at the intensity frontier

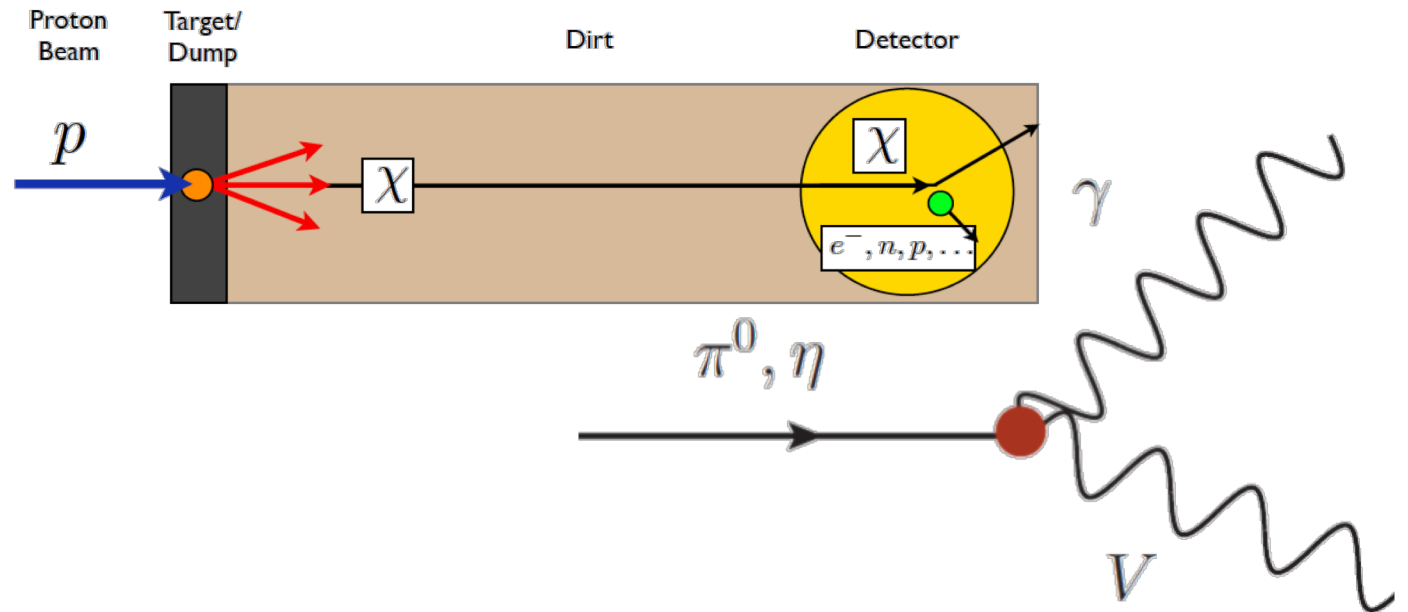
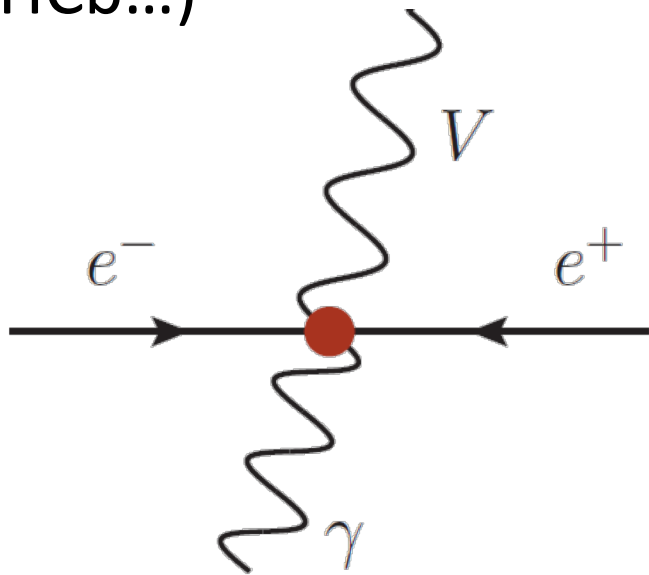
Production and detection at accelerator-based experiments

# Dark Sector searches - production

- Light dark sector particles may be accessible at the *intensity frontier*
  - « Low »-energy installations but with high intensity (e.g. LSND:  $10^{22} \pi^0$ )

Precision experiments at collider (e.g BaBar, BELLE, LHCb...)

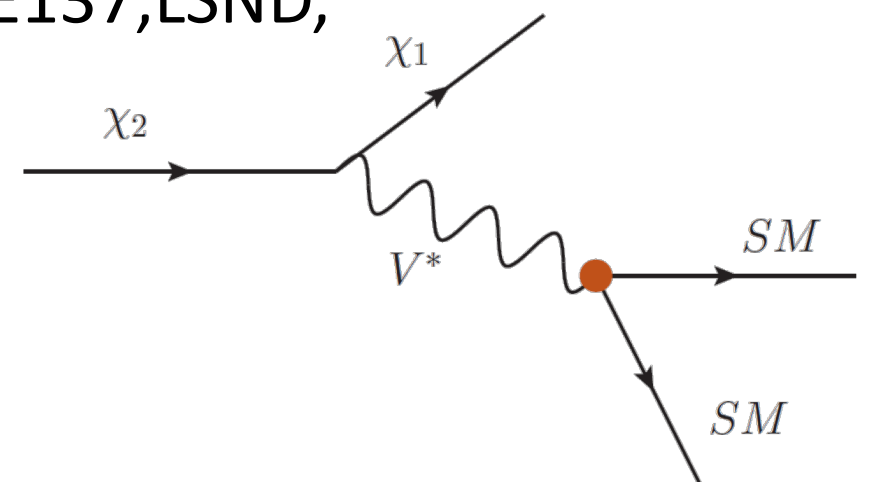
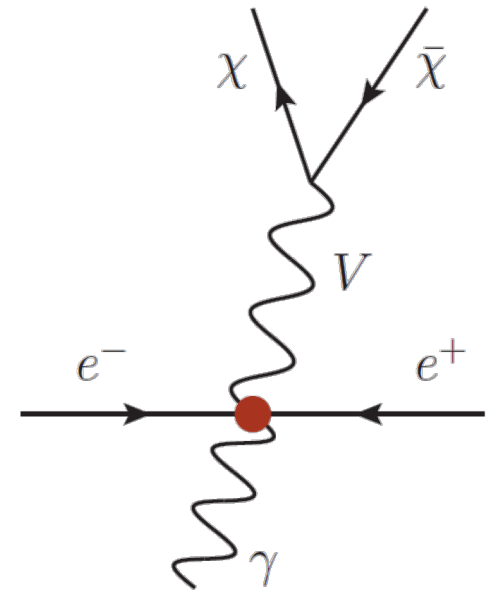
Beam-dump/fixed-target types of experiments (LSND, E137...)





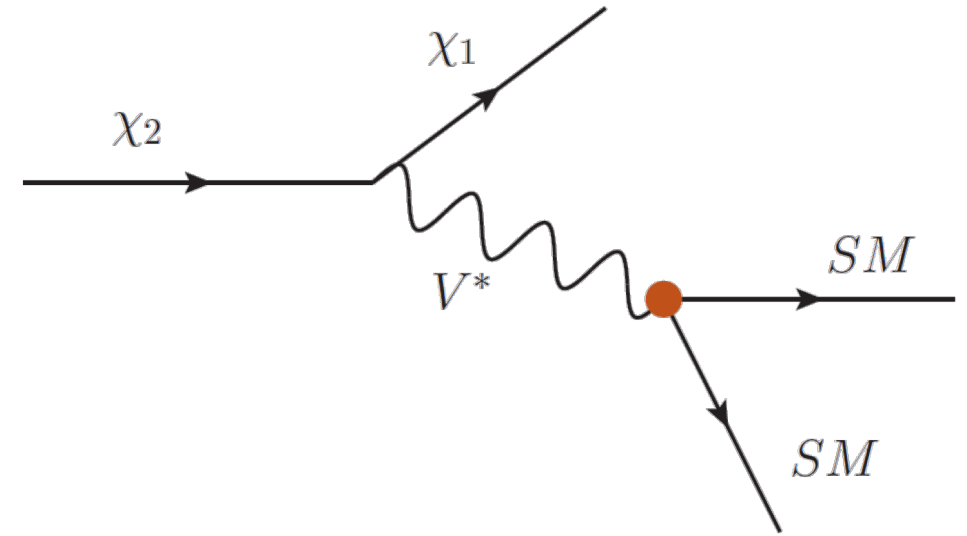
# Dark Sector searches -detection

- **Missing energy/ Invisible decay:** Mono-photon searches missing energy signature @ BaBar, NA64.
  - Mostly model independent,  $\varepsilon < 10^{-3}$
- **Dark sector beam production and detection**
  - **Scattering:** Searching for DM via scattering in fixed target/long-baseline neutrinos experiments (E137, LSND, miniBooNE ...)
  - **Dark sector visible decay:** Detection of an electron/positron pair in fixed target/beam dump experiments



# Heavy dark sector decays

- Dark sector decays through off-shell mediator  $\rightarrow$  typical decay length of order meter
  - $\rightarrow$  Decay into pair of electrons (no background)
  - $\rightarrow$  In optimum region, large portion of the heavy dark states decay in the detector

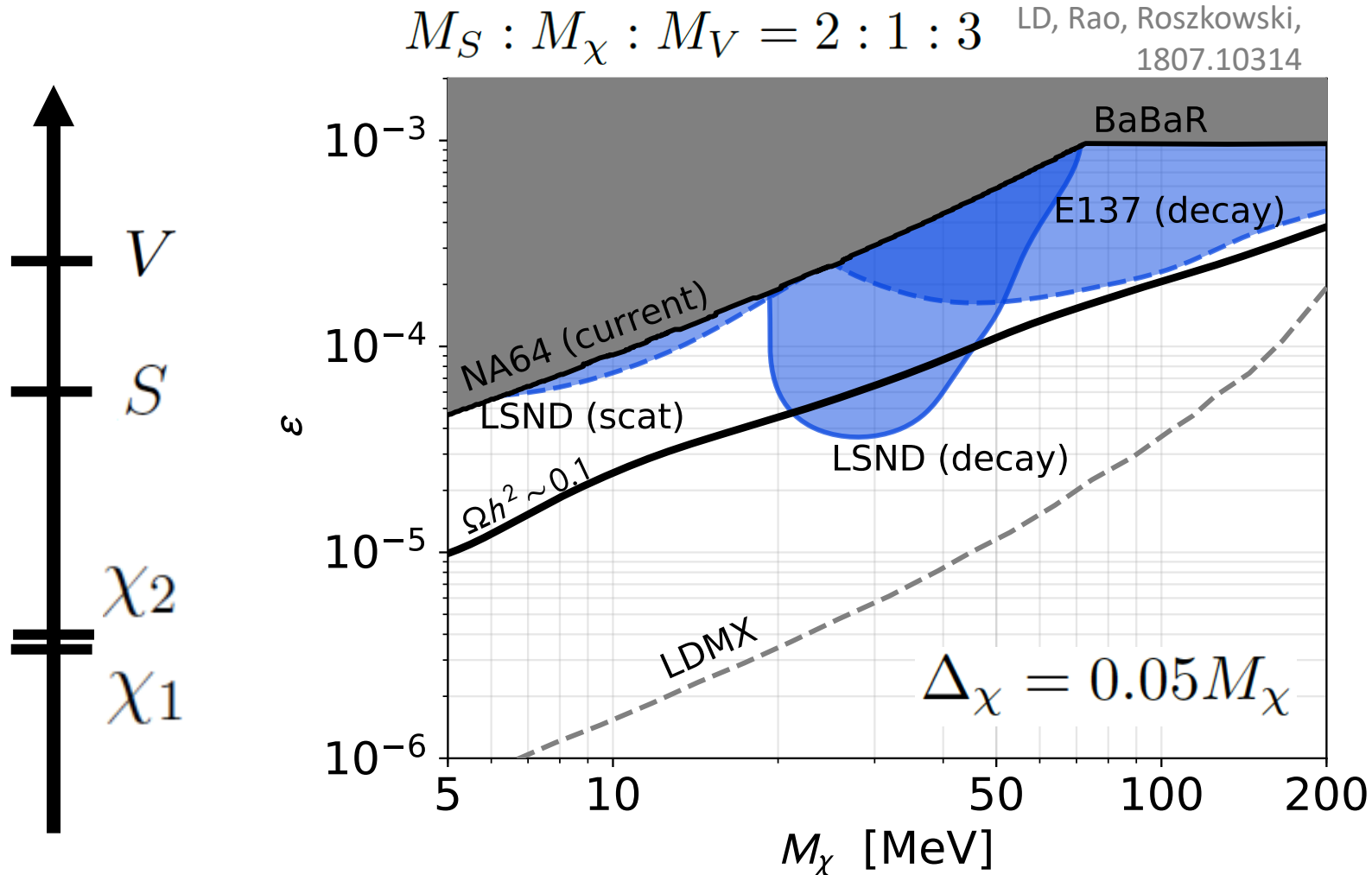


$$c\tau_{\chi_2} \propto 100 \text{ m} \times \left(\frac{0.1}{\alpha_D}\right) \left(\frac{10^{-3}}{\varepsilon}\right)^2 \left(\frac{0.2}{\Delta_\chi}\right)^5 \left(\frac{25 \text{ MeV}}{M_S}\right)^5 \left(\frac{M_V}{100 \text{ MeV}}\right)^4$$

- In the optimum case  $\text{NoE}_{\text{HDS}} \propto \varepsilon^2$
- Dark Higgs boson can belong to this category with  $S \rightarrow V e^+ e^-$  is opened

# Inelastic DM regime

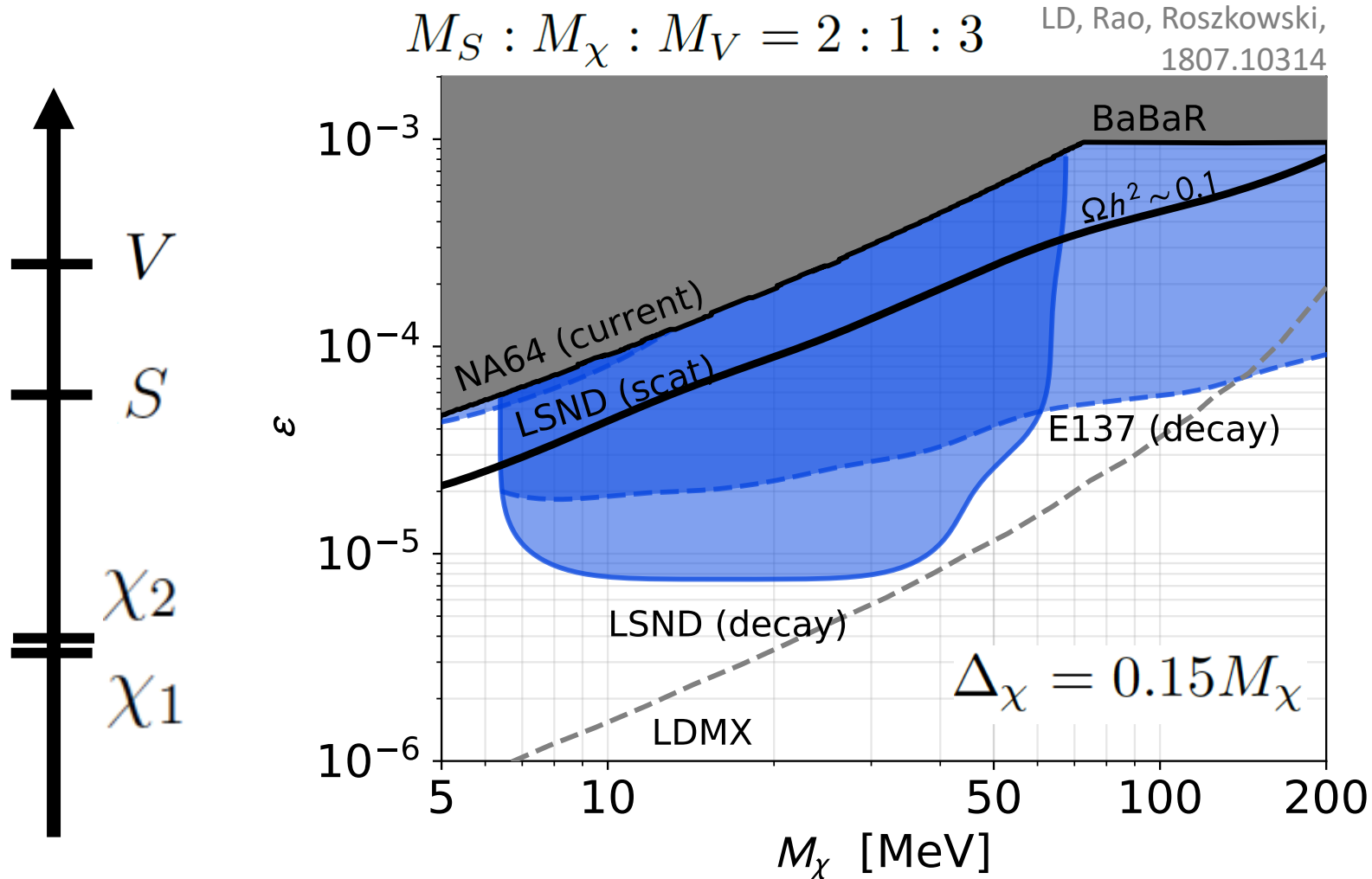
- Relic density fixed by s-channel, co-annihilation process:  $\chi_1\chi_2 \rightarrow e^+e^-$



- Main signatures:
  - Missing energy searches
  - $\chi_2 \rightarrow \chi_1 e^+ e^-$  decay
  - $\chi_1$  scattering
- When consider dark sector decays, decades-old experiment are still strongly ahead of current mono-photon searches!

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- Relic density fixed by s-channel, co-annihilation process:  $\chi_1\chi_2 \rightarrow e^+e^-$

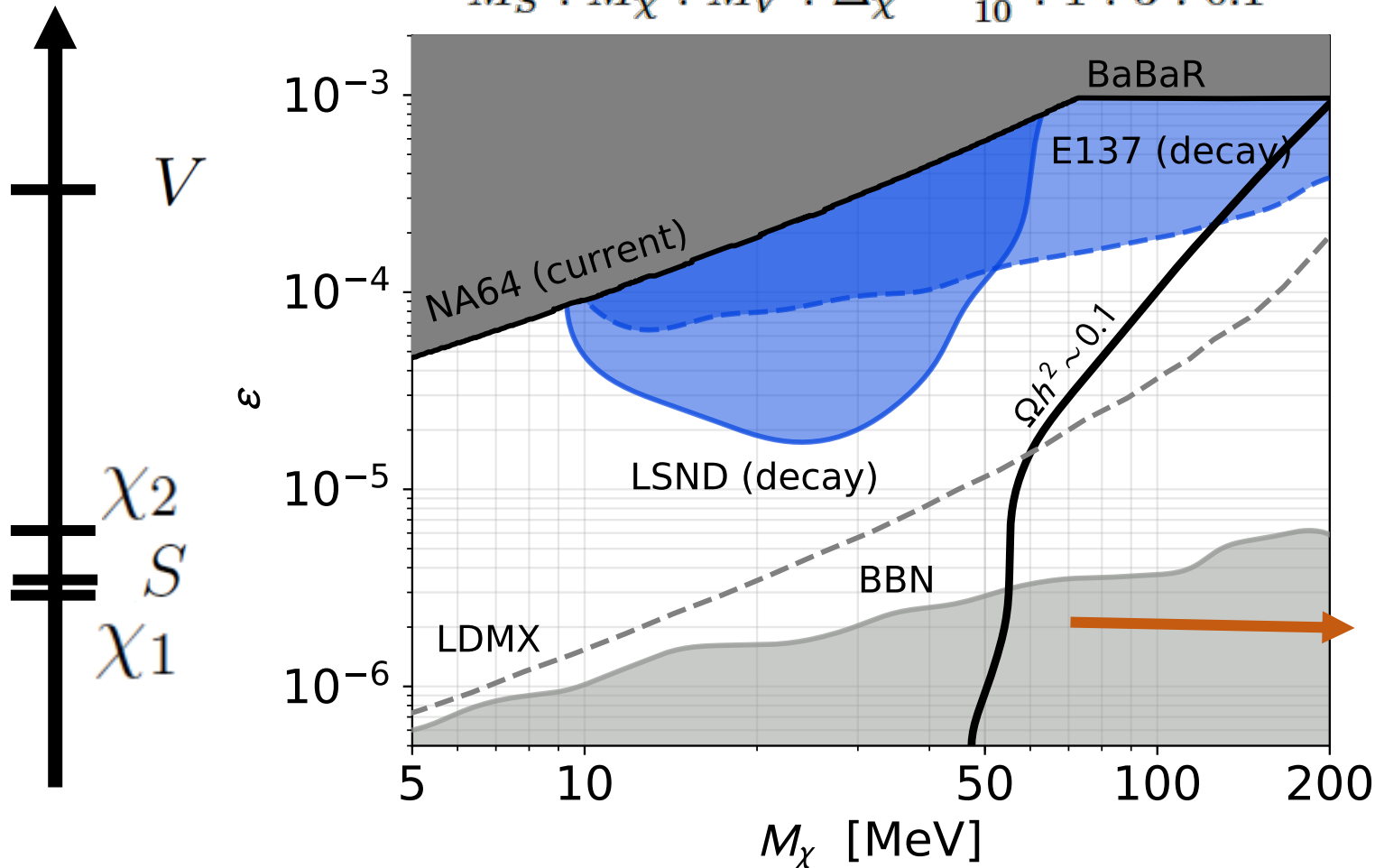


- Main signatures:
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  - $\chi_1$  scattering
- When consider dark sector decays, decades-old experiment are still strongly ahead of current mono-photon searches!

# Secluded DM regime

- Relic density mainly fixed by t-channel process,  $\chi\chi \rightarrow SS$

$$M_S : M_\chi : M_V : \Delta_\chi = \frac{9}{10} : 1 : 3 : 0.1$$



- Main signatures:
  - Missing energy searches
  - $\chi_2 \rightarrow \chi_1 e^+ e^-$  decay
  - $\chi_1$  scattering
- Dark Higgs boson is both very long-lived and abundantly produced  $\rightarrow$  BBN bound

Light dark sector effective theory and the  
intensity frontier:  
the fermion portal

# The off-shell effective approach

- Supposed a large mass splitting between dark matter and mediator scale  
 → Effective theory description relevant (same for leptons)

Vector operator

$$\sum_q \frac{g_q}{\Lambda^2} (\bar{\chi}_1 \gamma_\mu \chi_2) (\bar{q} \gamma^\mu q)$$

Axial vector operator

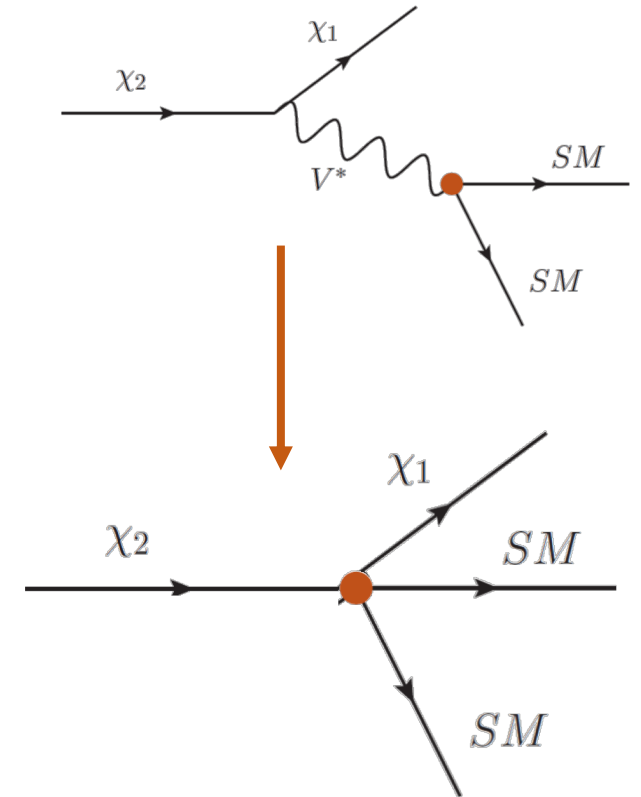
$$\sum_q \frac{\tilde{g}_q}{\Lambda^2} (\bar{\chi}_1 \gamma_\mu \gamma^5 \chi_2) (\bar{q} \gamma^\mu \gamma^5 q)$$

- Based on the above diagram, basic equivalence

$$\frac{\Lambda}{\sqrt{g}} \sim \frac{M_V}{\sqrt{\epsilon g_{De}}}$$



Could probe scale 2 to 3 orders of magnitudes larger than  $\chi$



# The off-shell effective approach (2)

- For off-shell mediator, long-lifetime comes from mass suppression
  - Typically need to go to the saturation limit  $M_1 \ll M_2$

$$c\tau^{\text{sat}} \sim 2 \text{ m} \times \left( \frac{\Lambda/\sqrt{g}}{1 \text{ TeV}} \right)^4 \left( \frac{1 \text{ GeV}}{M_{\chi_2}} \right)^5$$

- One can also search for “grey sector” with mixed coupling/heavy scale suppression

$$c\tau^{\text{PD}} \sim 375 \text{ m} \times \left( \frac{100 \text{ GeV}}{\Lambda} \right)^4 \left( \frac{1 \text{ GeV}}{M_{\chi_1}} \right)^5 \left( \frac{0.25}{\Delta_\chi} \right)^5 \left( \frac{0.01}{g} \right)^2$$

- Production is completely modified vs. the on-shell dark photon mediator
  - Strong suppression of lower energy processes  $\rightarrow$  typical suppression  $\propto \frac{M_m^4}{\Lambda^4}$
  - No CS from dark photon bremsstrahlung  $\rightarrow$  **Electron beam-dump hardly relevant**



# Meson decays in the fermion portal

- Depending on the nature of the operators, different production channels

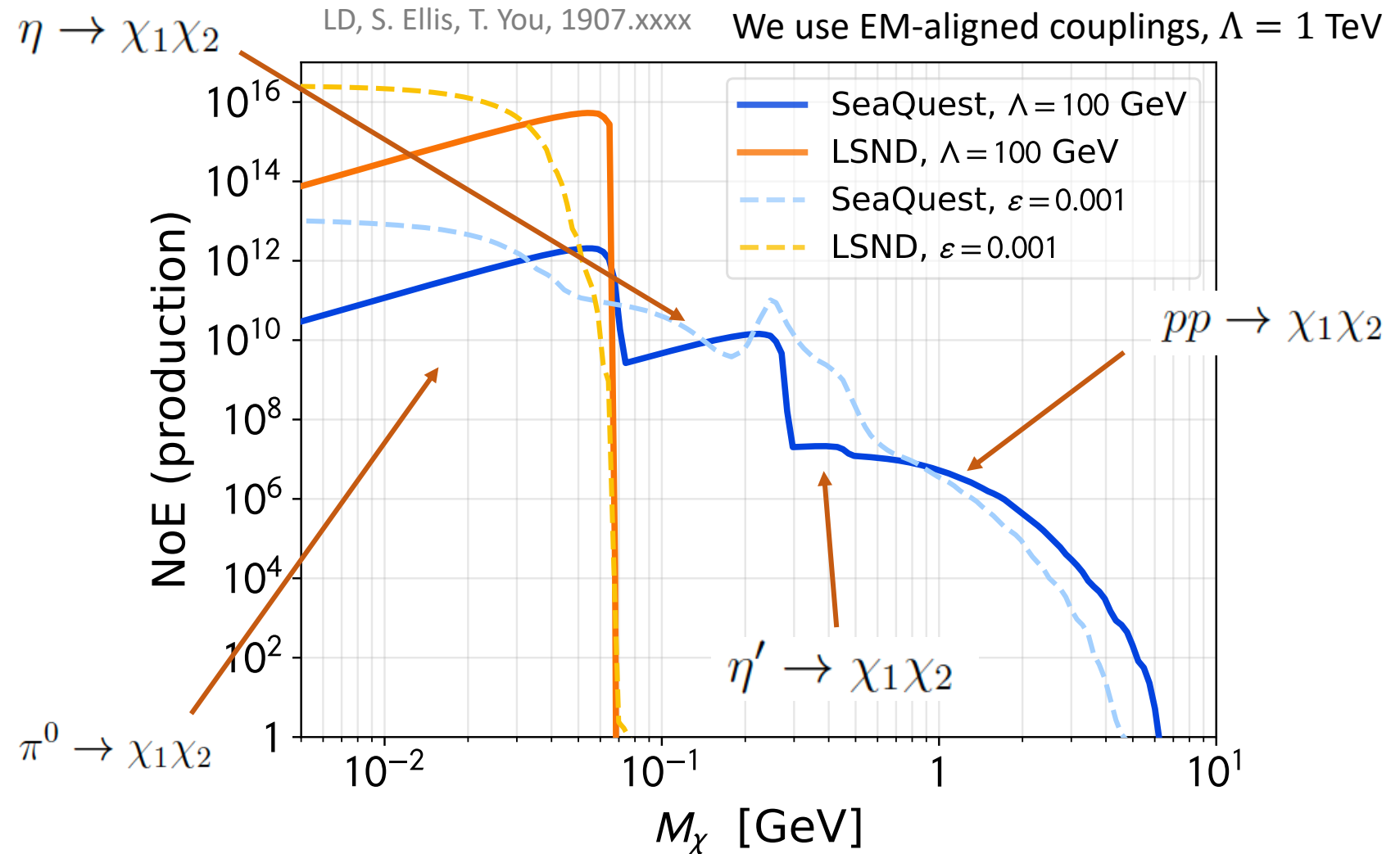
Vector operator

$$\pi^0, \eta, \eta' \rightarrow \gamma \chi \chi$$

$$\rho, \omega \rightarrow \chi \chi$$

Axial vector operator

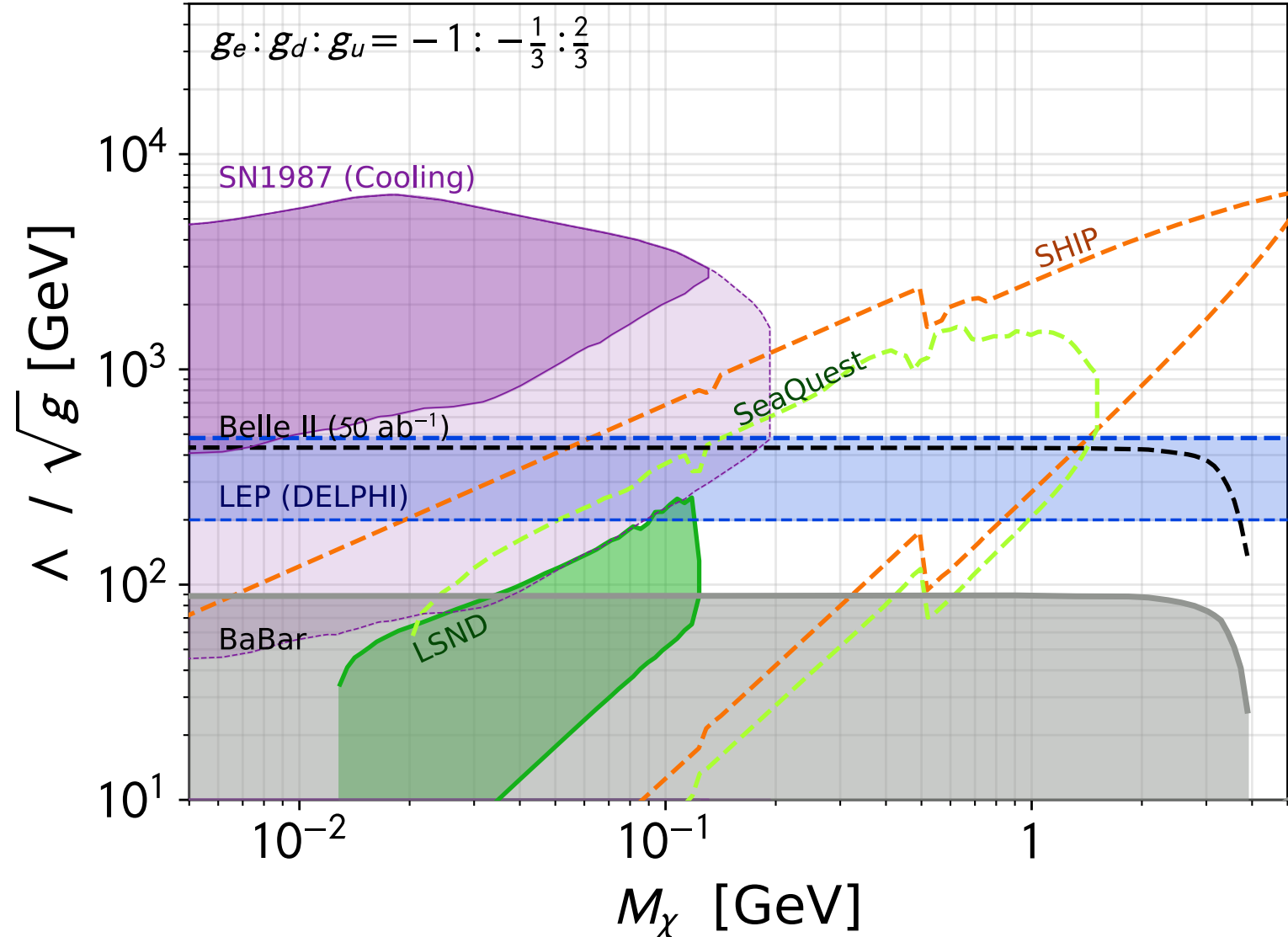
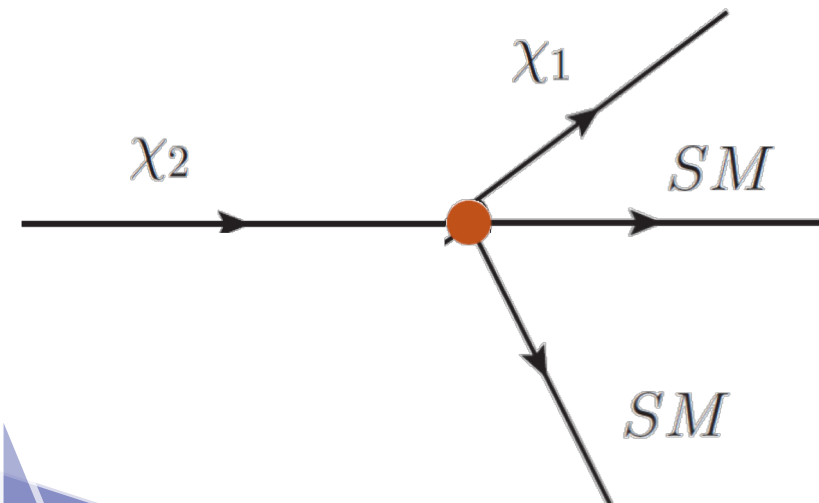
$$\pi^0, \eta, \eta' \rightarrow \chi \chi$$



# Limits in the vector case

LD, S. Ellis, T. You, 1907.xxxx

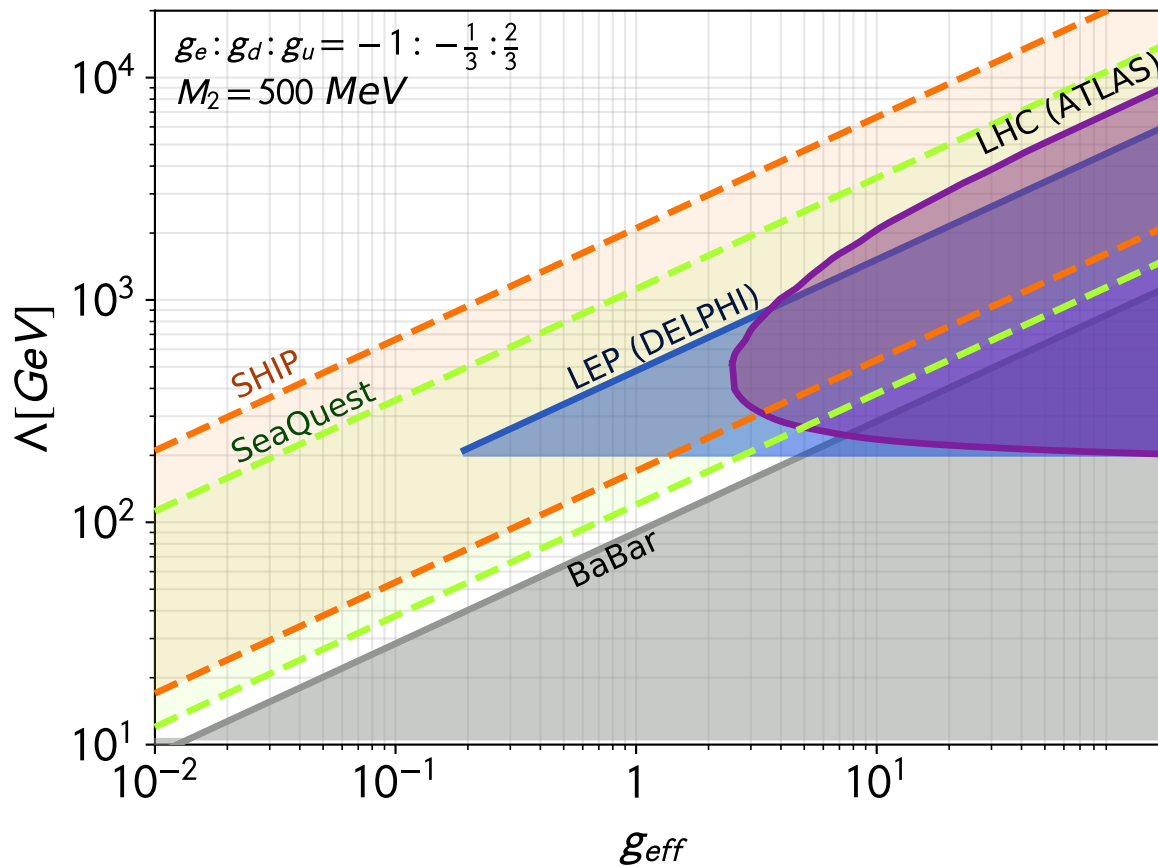
- We consider the saturation limit  $M_2 \gg M_1$
- Beam-dump experiment probing significantly above their energy



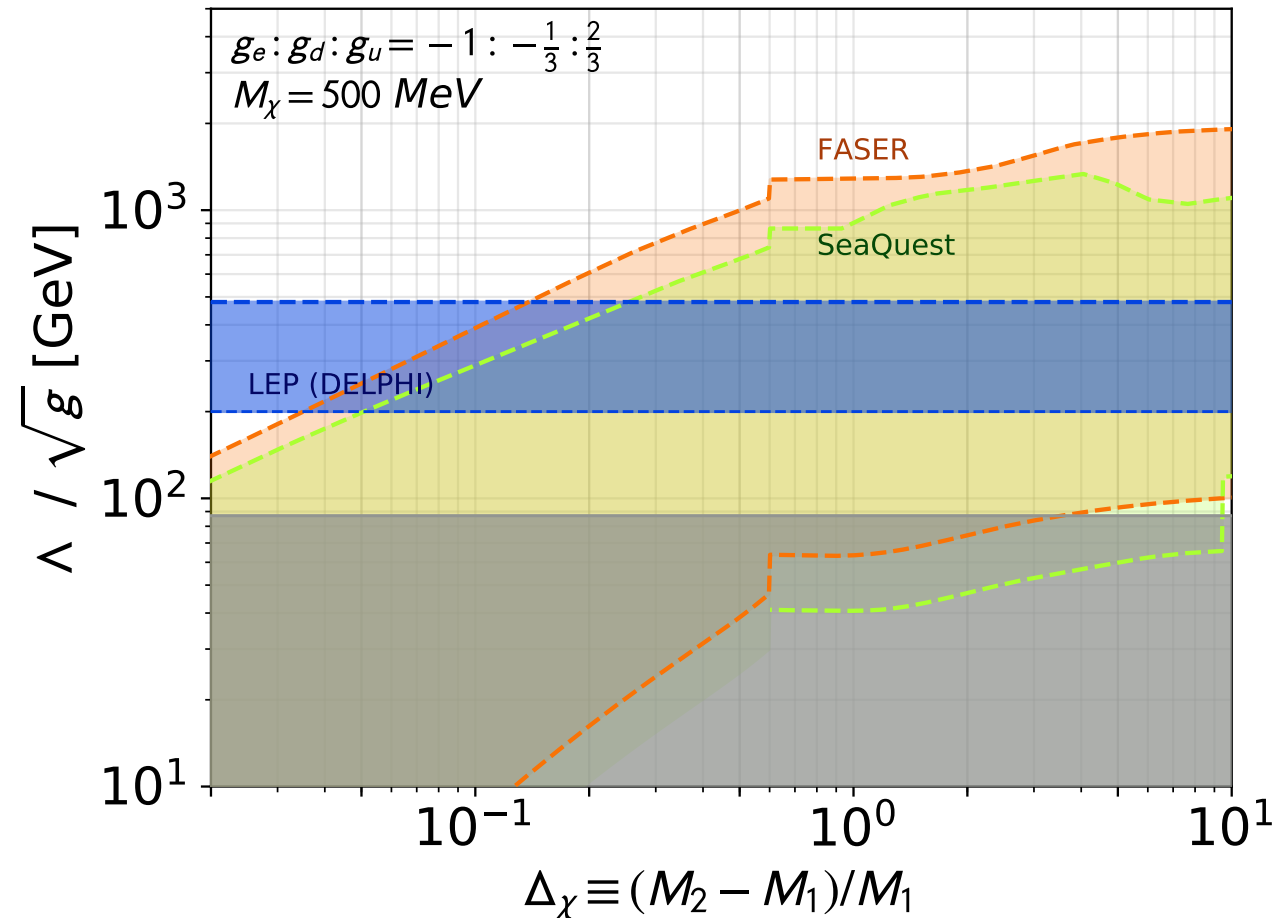
# Varying the couplings and splitting

- Reach of decay signatures depends strongly on  $\Delta_\chi$ , reach of collider mono-X on  $g_{eff}$

LD, S. Ellis, T. You, 1907.xxxx



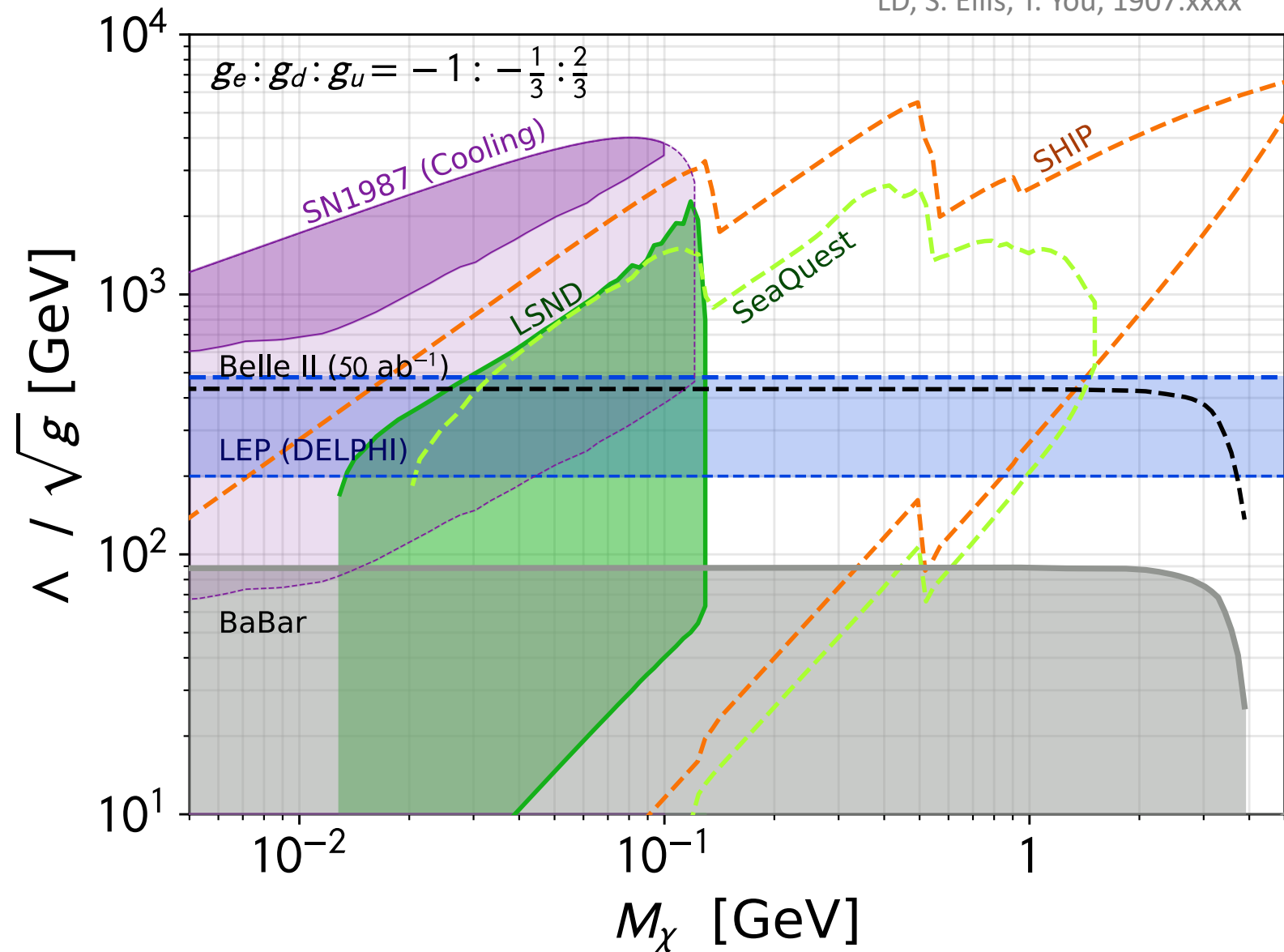
LD, S. Ellis, T. You, 1907.xxxx



# Limits in the axial-vector case

LD, S. Ellis, T. You, 1907.xxxx

- Mesons production strongly enhanced
- Better low-mass limits
- SN1987 based on  $\pi^0$  decay
  - > Very uncertain lower limit (also splitting-dependent)
- Notice the low limit for SHIP/SeaQuest



# Conclusion

# Looking forward ...

- Many upcoming relevant experiments:
  - **Neutrino** experiments -> the near detectors can search for dark sector particles
  - **Dark sector-oriented** -> looking for decays/ missing energy
  - **Flavour/ Rare mesons** decay -> Missing energy searches, invisible meson decay, etc...

Belle II, **BDX@Jlab**

**PADME**

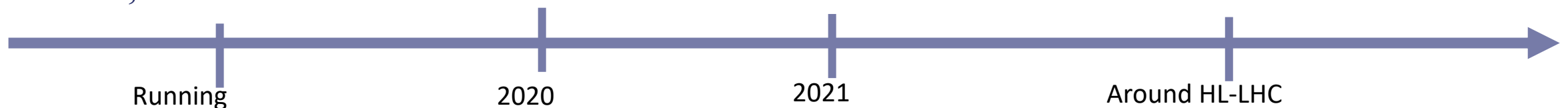
SBN, SeaQuest,

NA62, KOTO

LHCb Run 3,  
**FASER**, NA62+

**MAGIX**,  
**BDX @MESA**

**MATHUSLA**, SHIP,  
**CODEX-b**, KLEVER  
LBNF/DUNE



(Many missing, not all of them are funded yet...)

# Conclusion

- Light thermal dark matter candidates typically carry with them a dark sector **with long-lived particles**
- The dark sector allows more freedom in realizing the correct relic density.
- Lead to rich phenomenology in **intensity frontier experiments**
  - Accelerator searches, neutrino/dark sector/flavor experiments means bright prospects for dark sector searches
  - The accompanying dark sector offers both fascinating detection prospects and very strong astrophysical bounds (BBN ...)
- Dedicated analysis in such experiments have typically a strong discovery potential

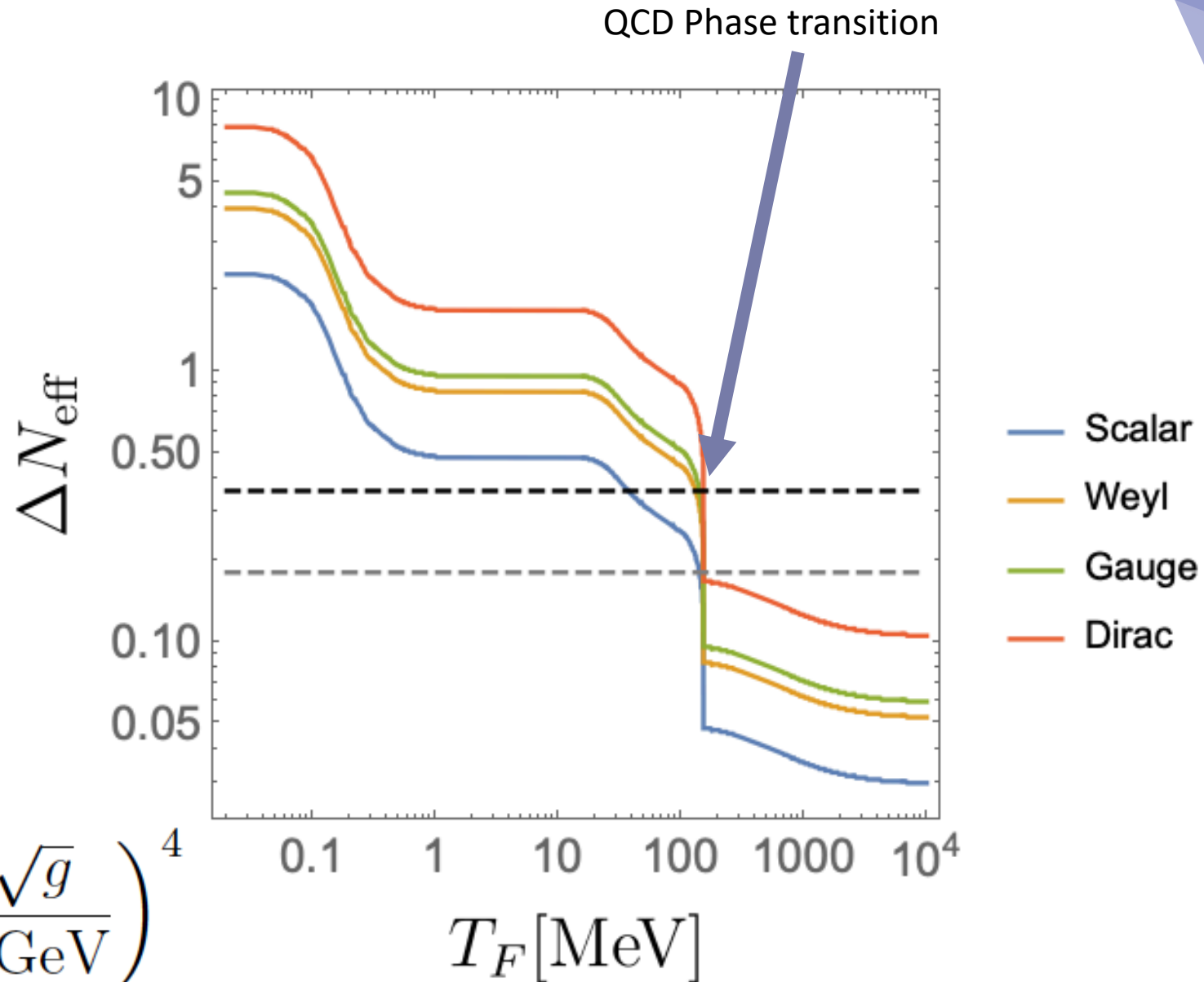
Backup slides



# Astrophysical limits

- For  $M_2 \gg M_1$ , the lightest dark sector can be relativistic relic
- More generally, one can still obtain dark matter candidate for iDM setup, e.g. for charm annihilation

$$\Omega h^2 \sim 0.3 \times \left( \frac{2 \text{ GeV}}{M_\chi} \right)^2 \left( \frac{\Lambda / \sqrt{g}}{500 \text{ GeV}} \right)^4$$

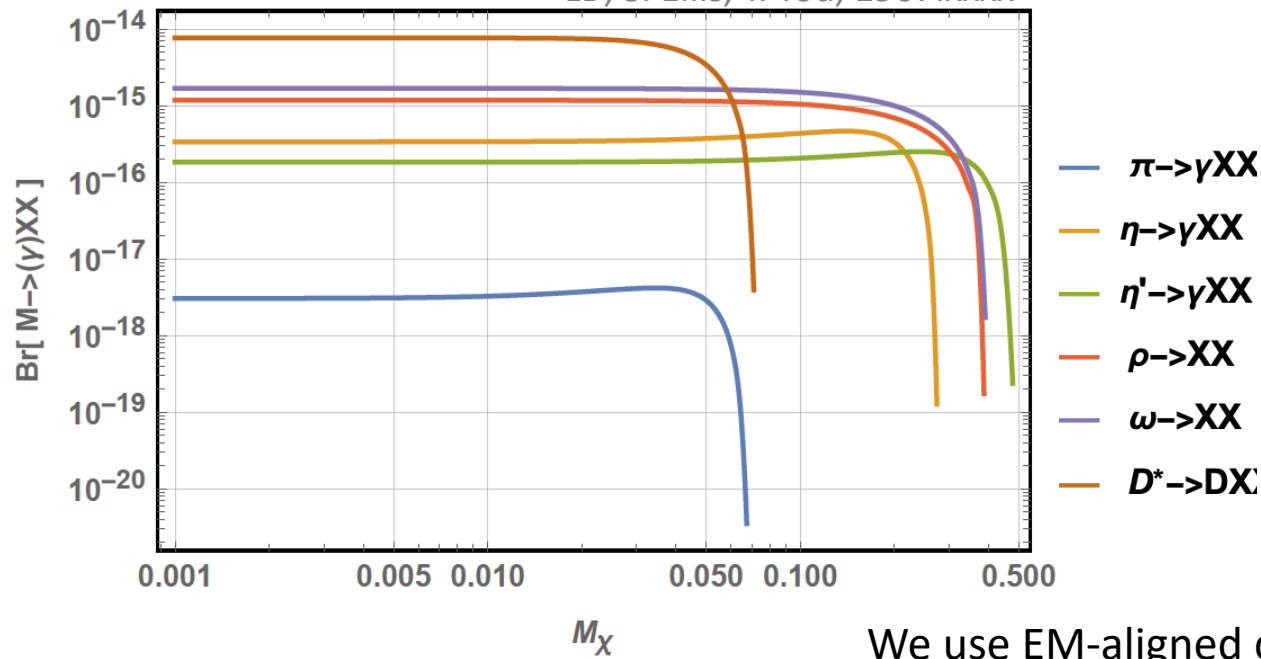


# Meson decays in the fermion portal

- Depending on the nature of the operators, different production channels

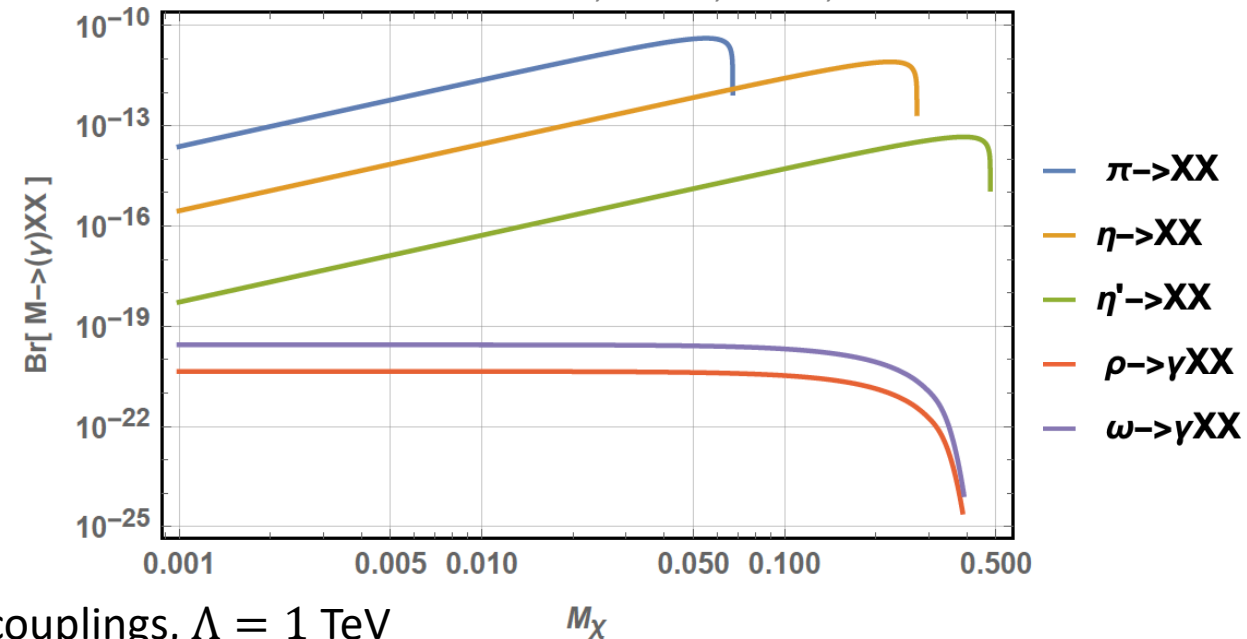
Vector operator  $\longrightarrow$   $\pi^0, \eta, \eta' \rightarrow \gamma\chi\chi$   
 $\rho, \omega \rightarrow \chi\chi$

LD, S. Ellis, T. You, 1907.xxxx



Axial vector operator  $\longrightarrow$   $\pi^0, \eta, \eta' \rightarrow \chi\chi$

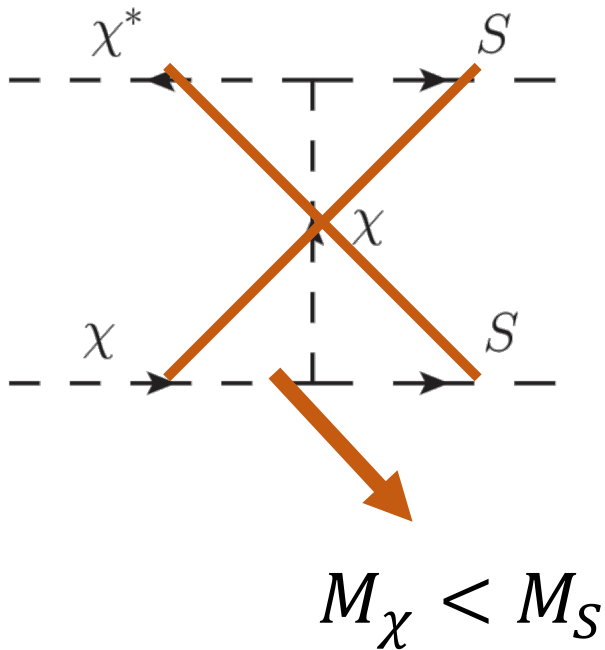
LD, S. Ellis, T. You, 1907.xxxx



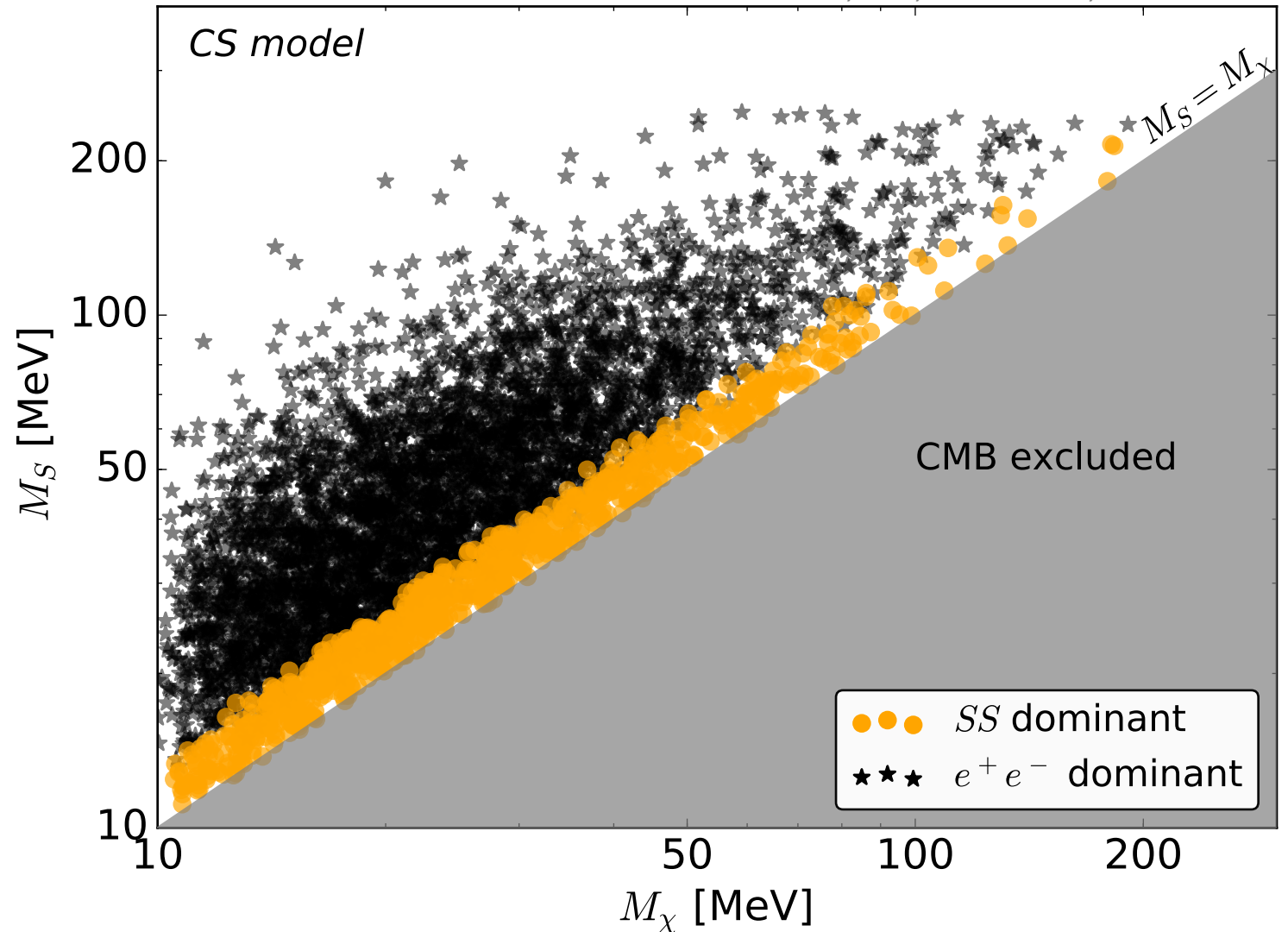
- Invisible meson decay limits relevant mostly for AV case, for  $\Lambda \sim 100 \text{ GeV}$

# CMB bounds -- the scalar case

- Planck measurements,  
→ forbids dark matter with s-wave annihilation and mass below the tens of GeV (1604.02457)

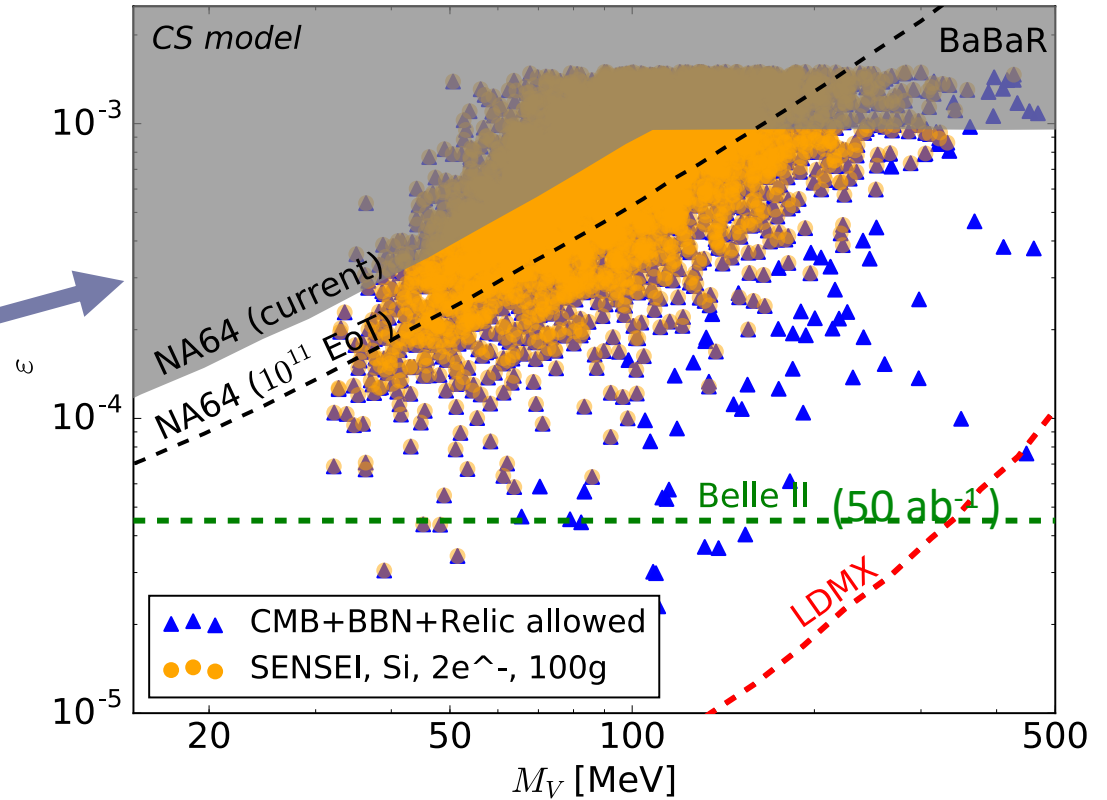
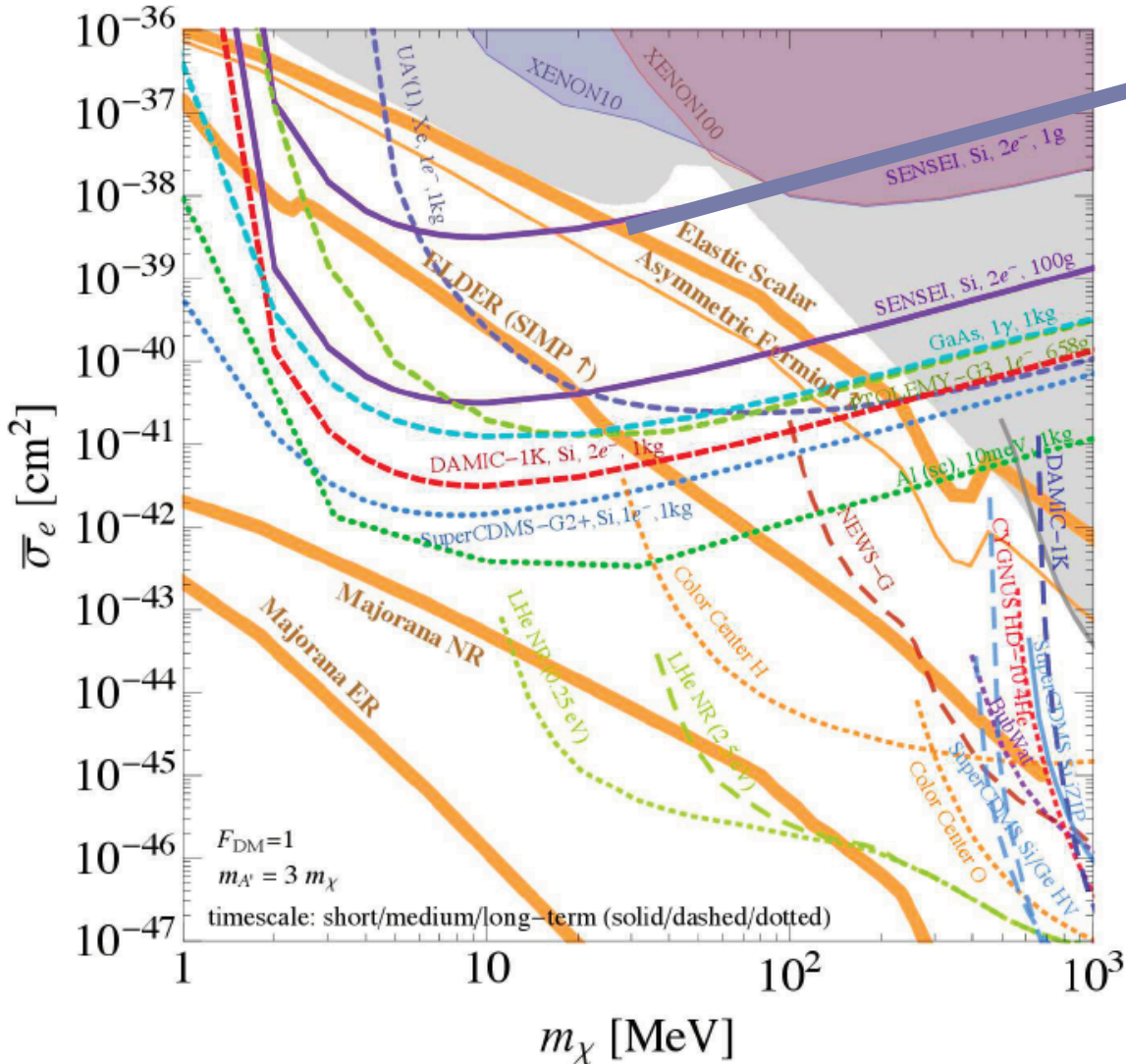


LD, Rao, Roszkowski, 1710.00971



# Direct Detection

From 1707.04591



- Interesting new development: boosted secondary DM flux (1810.10543, 1811.00520)
  - Collision with cosmic rays created boosted secondary flux
  - Allows better nucleon/electron recoil even for sub-GeV DM
- Very interesting prospects for iDM case, only partially explored though

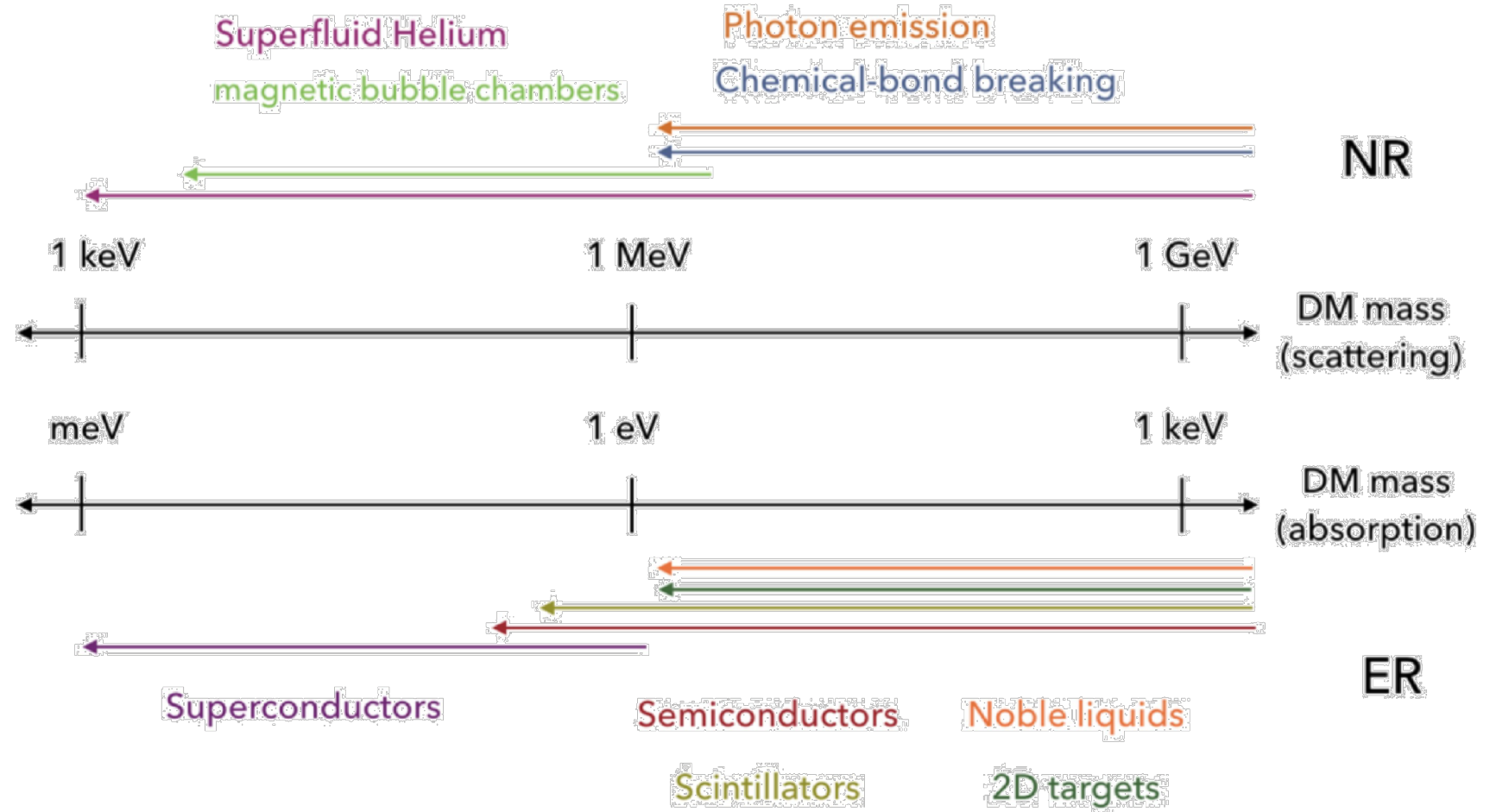
# Ranges for the bulk scan

Parameter	Range	Prior
$M_S$	5 MeV - 1 GeV	Log
$g_V$	0.01 - 2.5	Log
$M_V$	10 MeV - 500 MeV	Log
$\varepsilon$	$0.5 \times 10^{-6}$ - 0.001	Log
$M_\chi$	-250 MeV - 150 MeV	Linear
$\Delta_\chi$	$0.01 M_\chi $ - $10 M_\chi $	Log
$y_{DM}$	-2 - 2	Linear

Table 2. Input parameters for the scans presented in Figure 2.

# Detection strategies, sub-GeV DM

- Need for low-threshold experiments
  - DM-electron scattering
  - DM- low-Z elastic nucleus interactions
  - Bremsstrahlung in inelastic DM-nucleus scattering
- DAMIC, SENSEI, UA', NICE, SuperCDMS, NEWS



From  
arXiv:1707.04591

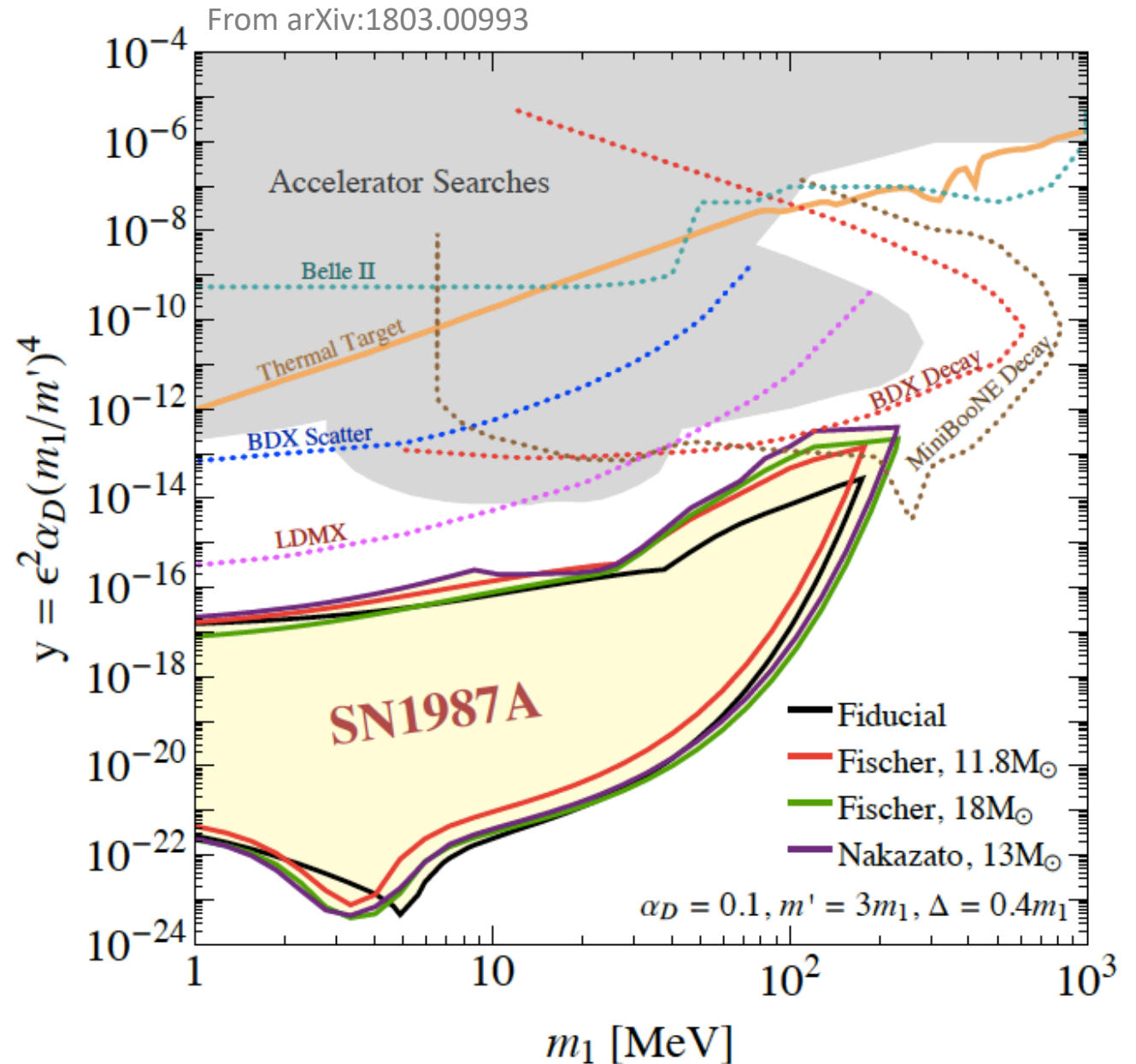
# SN 1987A bounds

- Typical bounds arise when DM do not scatter enough and escape the SN core and escape the SN core

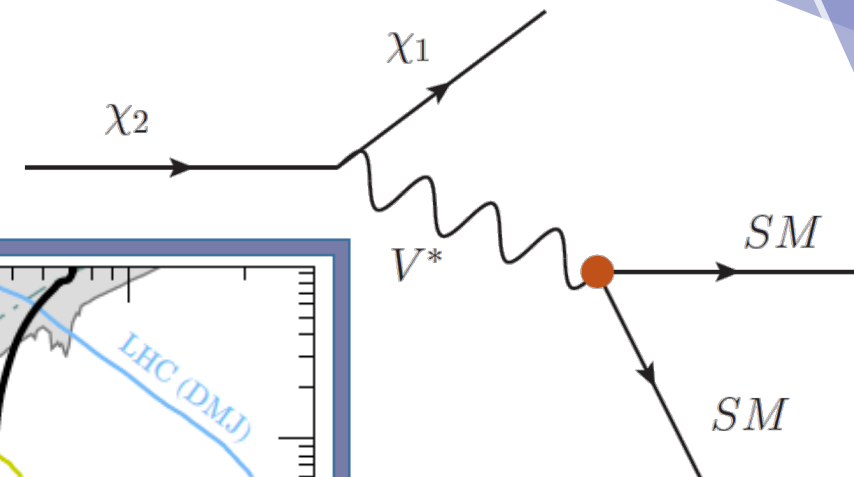
$$\alpha_D \epsilon^2 < O(\text{few}) \times 10^{-14}$$

- Not relevant for pseudo-Dirac case/Majorana case at the thermal target
- Dark Higgs bounds may be relevant at  $m_S < M_{\chi_1}$  or  $m_S < M_{\chi_2} - M_{\chi_1}$

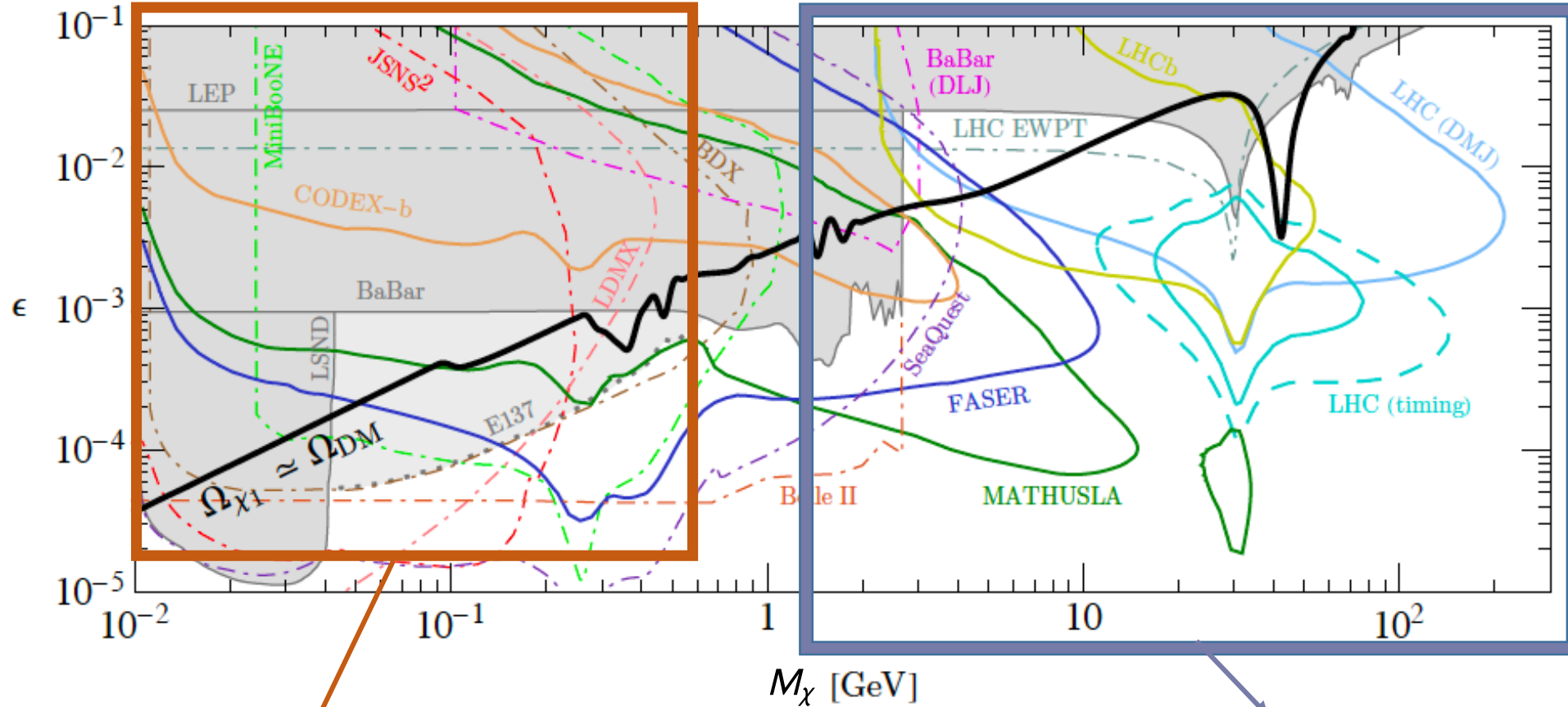
→ But scattering with DM halo inside the SN should be enough to trap it



# Heavy dark sector decays (2)



Fermionic iDM,  $m_{A'} = 3m_1$ ,  $\Delta=0.1$ ,  $\alpha_D=0.1$



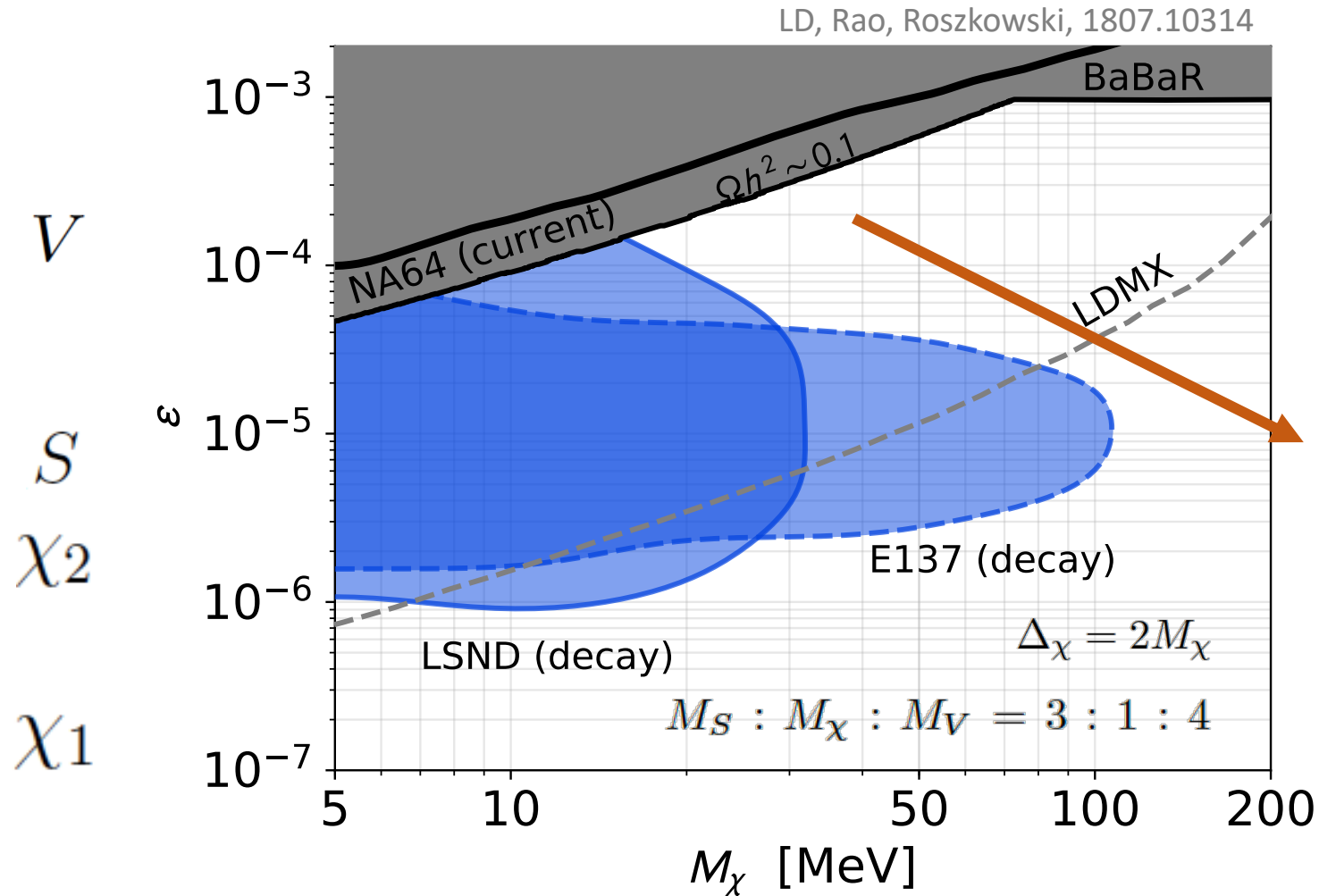
- **High-intensity frontier:** basically electrons physics

- “Grey sector”, LHC-related: Displaced muon jets, LLP...



# Heavy dark sector decay in mDM

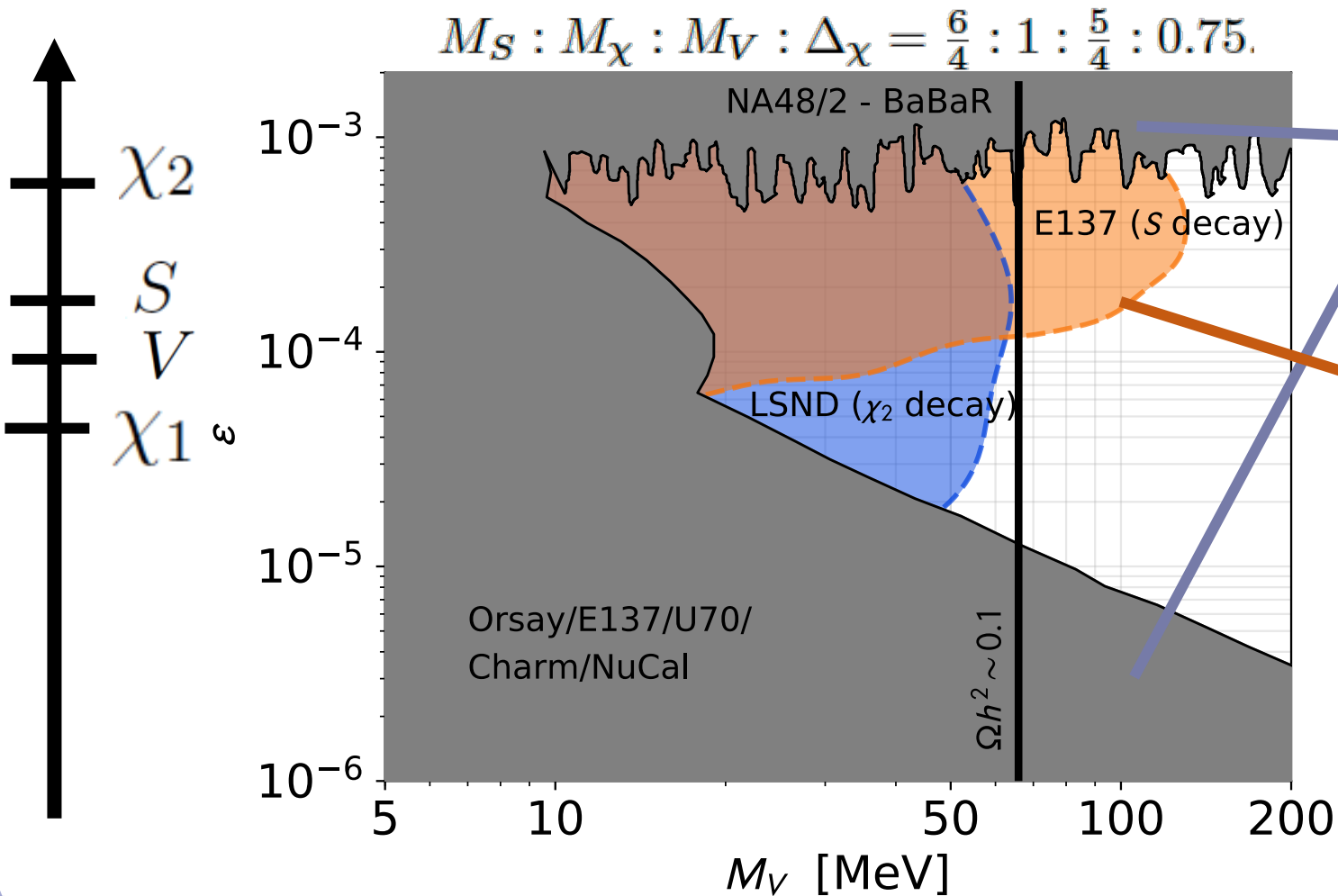
- Relic density fixed by s-channel, p-wave process:  $\chi\chi \rightarrow e^+e^-$



- Main signatures:
  - Missing energy searches
  - $\chi_2 \rightarrow \chi_1 e^+ e^-$  decay
- For large splitting, lifetime of order meter,  $c\tau_\chi < 1 m$ 
  - long baseline experiments like E137 miss critical part of the parameter space
  - Opportunity for others, MAGIX, SeaQuest, LHCb ...

# Forbidden DM regime

- Relic density fixed by thermally suppressed t-channel process,  $\chi\chi \rightarrow VV$



- $M_V \sim M_\chi \rightarrow$  Dark photon decays visibly
- $S \rightarrow V e^+ e^-$  kinematically open, dark Higgs lifetime shorter
  - $c \tau_S \sim 1 - 100 m$
- Thermal target independent of kinetic mixing