

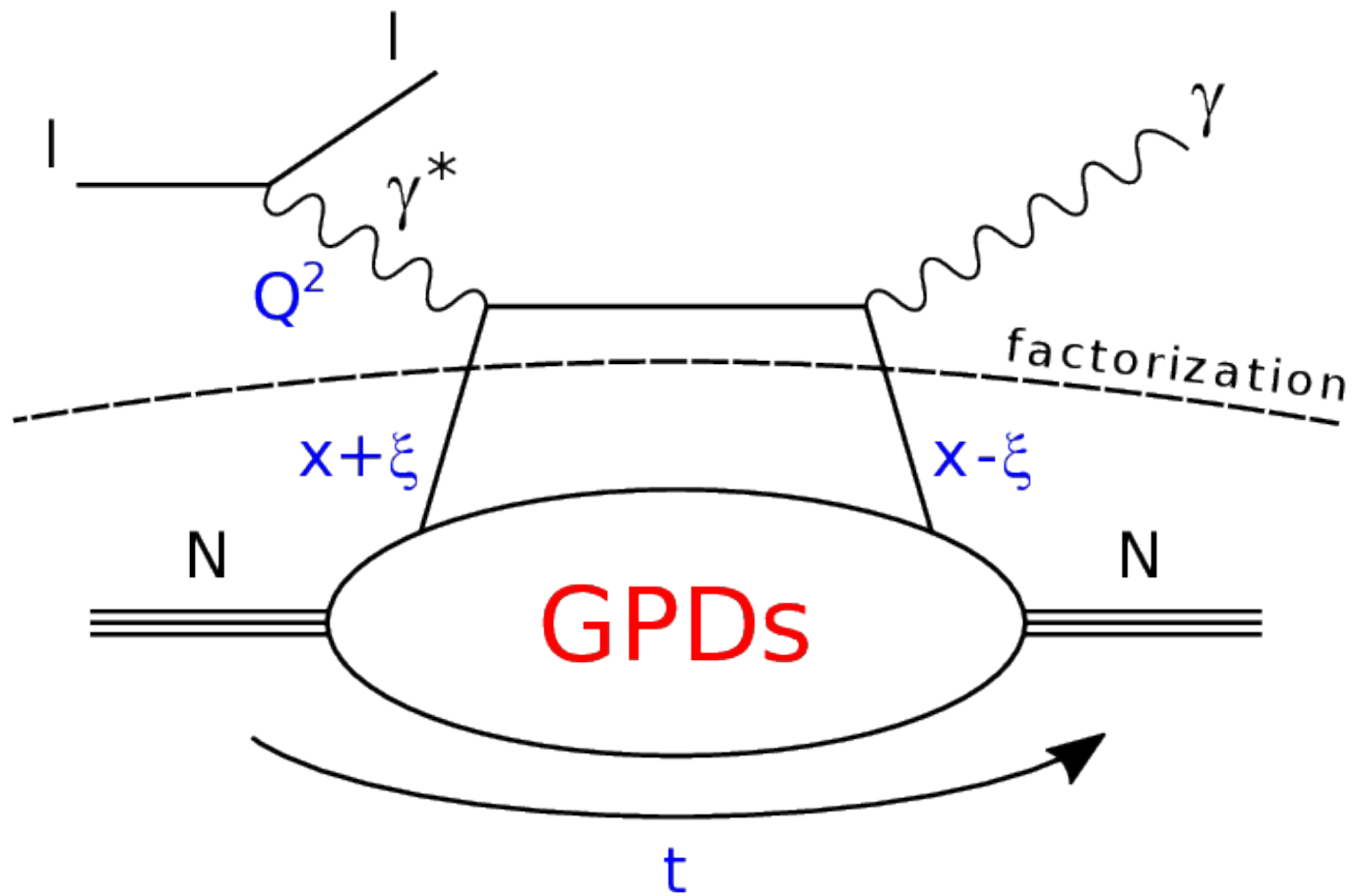
Addressing the problem of model dependency in GPD phenomenology



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Deeply Virtual Compton Scattering (DVCS)



factorisation for $|t|/Q^2 \ll 1$

Chiral-even GPDs:
(helicity of parton conserved)

$H^{q,g}(x,\xi,t)$	$E^{q,g}(x,\xi,t)$	for sum over parton helicities
$\tilde{H}^{q,g}(x,\xi,t)$	$\tilde{E}^{q,g}(x,\xi,t)$	for difference over parton helicities
nucleon helicity conserved	nucleon helicity changed	

Reduction to PDF:

$$H(x, \xi = 0, t = 0) \equiv q(x)$$

Polynomiality - non-trivial consequence of Lorentz invariance:

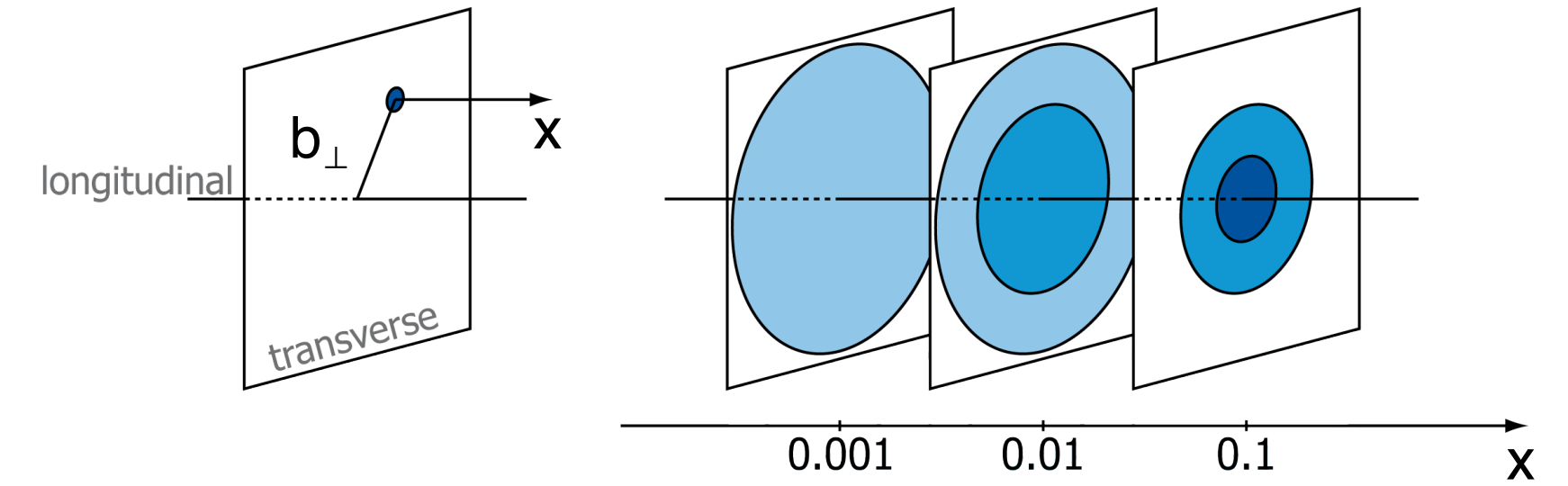
$$\mathcal{A}_n(\xi, t) = \int_{-1}^1 dx x^n H(x, \xi, t) = \sum_{\substack{j=0 \\ \text{even}}}^n \xi^j A_{n,j}(t) + \text{mod}(n, 2) \xi^{n+1} A_{n,n+1}(t)$$

Positivity bounds - positivity of norm in Hilbert space, e.g.:

$$|H(x, \xi, t)| \leq \sqrt{q\left(\frac{x+\xi}{1+\xi}\right) q\left(\frac{x-\xi}{1-\xi}\right) \frac{1}{1-\xi^2}}$$

Nucleon tomography:

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta}{4\pi^2} e^{-i\mathbf{b}_\perp \cdot \Delta} H^q(x, 0, t = -\Delta^2)$$



Energy momentum tensor in terms of form factors (OAM and mechanical forces):

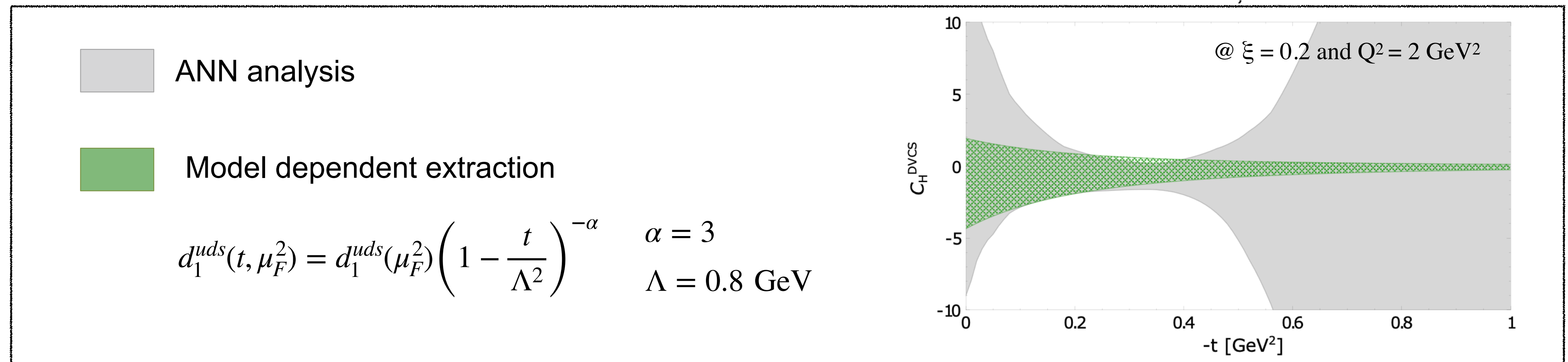
$$T^{\mu\nu} = \begin{bmatrix} \text{Energy density } T^{00} & \text{Momentum density } T^{01} & T^{02} & T^{03} \\ T^{10} & T^{11} & T^{12} & T^{13} \\ T^{20} & T^{21} & T^{22} & T^{23} \\ T^{30} & T^{31} & T^{32} & T^{33} \end{bmatrix}$$

Shear stress
Normal stress

Energy flux Momentum flux

$$\langle p', s' | \hat{T}^{\mu\nu} | p, s \rangle = \bar{u}(p', s') \left[\frac{P^\mu P^\nu}{M} A(t) + \frac{\Delta^\mu \Delta^\nu - \eta^{\mu\nu} \Delta^2}{M} C(t) + M \eta^{\mu\nu} \bar{C}(t) + \right. \\ \left. \frac{P^\mu i \sigma^{\nu\lambda} \Delta_\lambda}{4M} [A(t) + B(t) + D(t)] + \frac{P^\nu i \sigma^{\mu\lambda} \Delta_\lambda}{4M} [A(t) + B(t) - D(t)] \right] u(p, s)$$

- Despite a substantial progress in both measurement and description of exclusive processes, and in lattice-QCD the problem of the model dependency of GPDs is still poorly addressed.
- Exceptions:
 - probing nucleon tomography at low-xB (see: [N. d'Hose's talk](#))
 - extraction of D-term (see: [Nature 570 \(2019\) 7759, E1](#), [EPJC 81 \(2021\) 4, 300 and below](#))



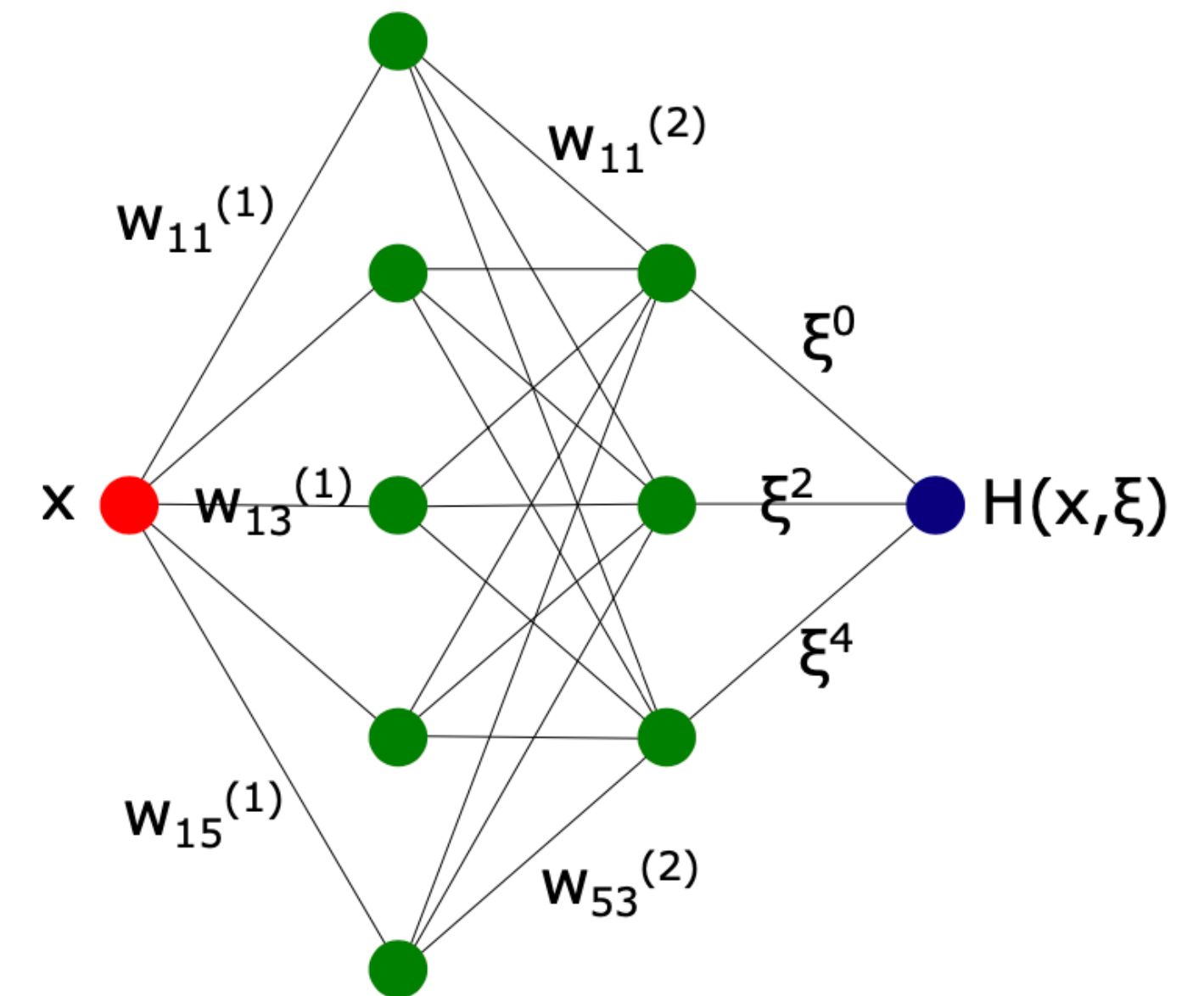
- No GPD models that could be considered non-parametric → no tools to study model dependency of the extraction of GPDs, nucleon tomography and orbital angular momentum (see: [EPJC 82 \(2022\) 3, 252 and this talk](#))

- Polynomiality:

$$\mathcal{A}_n(\xi) = \int_{-1}^1 dx x^n H(x, \xi) = \sum_{\substack{j=0 \\ \text{even}}}^n \xi^j A_{n,j} + \text{mod}(n, 2) \xi^{n+1} A_{n,n+1}$$

suggests that true degrees of freedom of GPDs are $A_{n,j}$ coefficients

- This leads us to the moment problem
→ reconstruction of GPDs from their moments
- We address this problem with ANNs
- Drawback of this method:
one can not keep PDF singularity for only $x=0$ and $\xi=0$
- See EPJC 82 (2022) 3, 252 and backup slides for more details



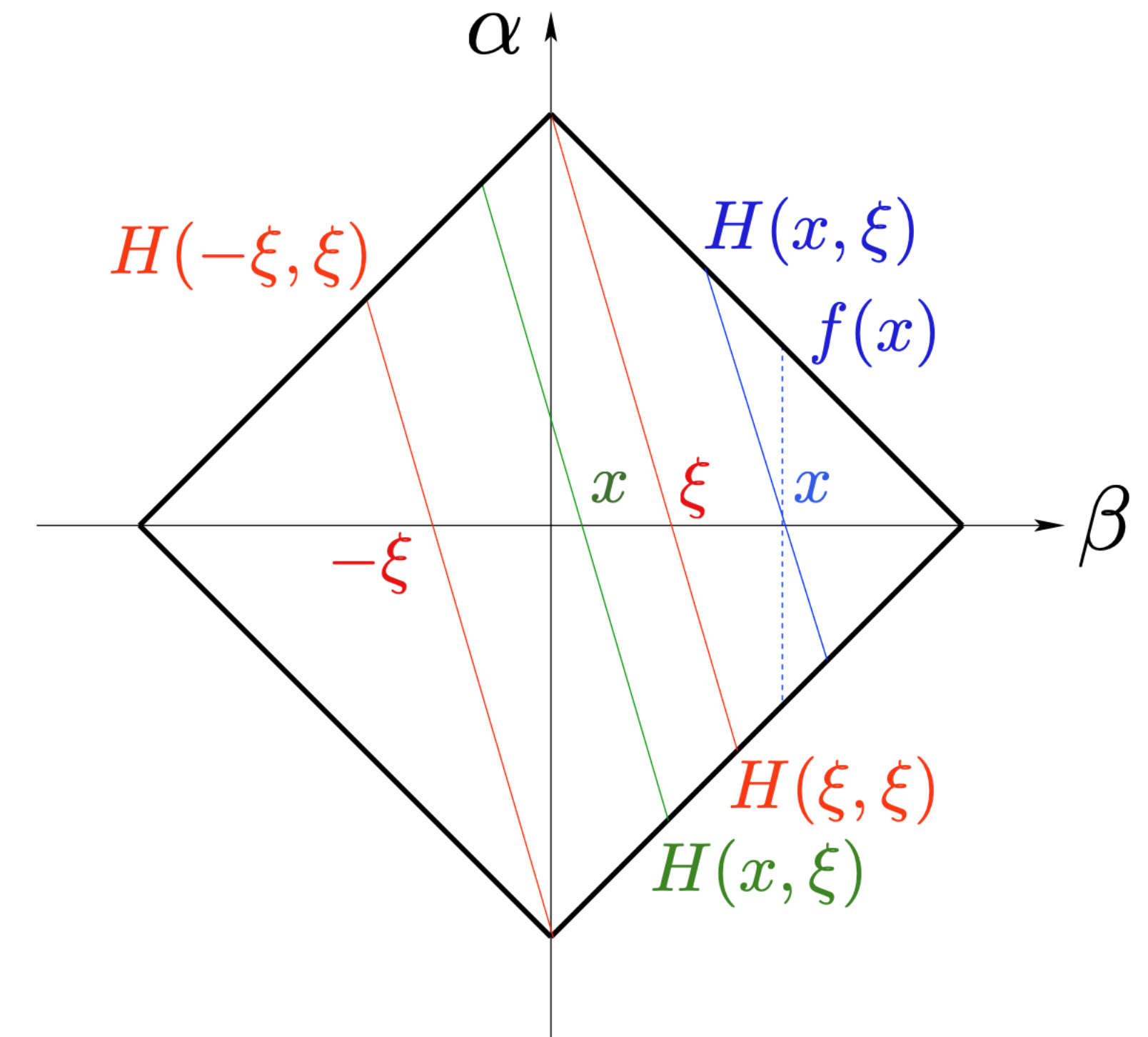
Double distribution:

$$H(x, \xi, t) = \int d\Omega F(\beta, \alpha, t)$$

where:

$$d\Omega = d\beta d\alpha \delta(x - \beta - \alpha\xi)$$

$$|\alpha| + |\beta| \leq 1$$



from PRD83, 076006, 2011

Double distribution:

$$(1 - x^2)F_C(\beta, \alpha) + (x^2 - \xi^2)F_S(\beta, \alpha) + \xi F_D(\beta, \alpha)$$

Classical term:

$$F_C(\beta, \alpha) = f(\beta)h_C(\beta, \alpha)\frac{1}{1 - \beta^2}$$

$$f(\beta) = \text{sgn}(\beta)q(|\beta|)$$

$$h_C(\beta, \alpha) = \frac{\text{ANN}_C(|\beta|, \alpha)}{\int_{-1+|\beta|}^{1-|\beta|} d\alpha \text{ANN}_C(|\beta|, \alpha)}$$

Shadow term:

$$F_S(\beta, \alpha) = f(\beta)h_S(\beta, \alpha)$$

$$f(\beta) = \text{sgn}(\beta)q(|\beta|)$$

$$h_S(\beta, \alpha)/N_S = \frac{\text{ANN}_S(|\beta|, \alpha)}{\int_{-1+|\beta|}^{1-|\beta|} d\alpha \text{ANN}_S(|\beta|, \alpha)} - \frac{\text{ANN}_{S'}(|\beta|, \alpha)}{\int_{-1+|\beta|}^{1-|\beta|} d\alpha \text{ANN}_{S'}(|\beta|, \alpha)}.$$

$$\text{ANN}_{S'}(|\beta|, \alpha) \equiv \text{ANN}_C(|\beta|, \alpha)$$

D-term:

$$F_D(\beta, \alpha) = \delta(\beta)D(\alpha)$$

$$D(\alpha) = (1 - \alpha^2) \sum_{\substack{i=1 \\ \text{odd}}} d_i C_i^{3/2}(\alpha)$$

Shadow term is closely related to the so-called **shadow GPDs**

Shadow GPDs have considerable size and:

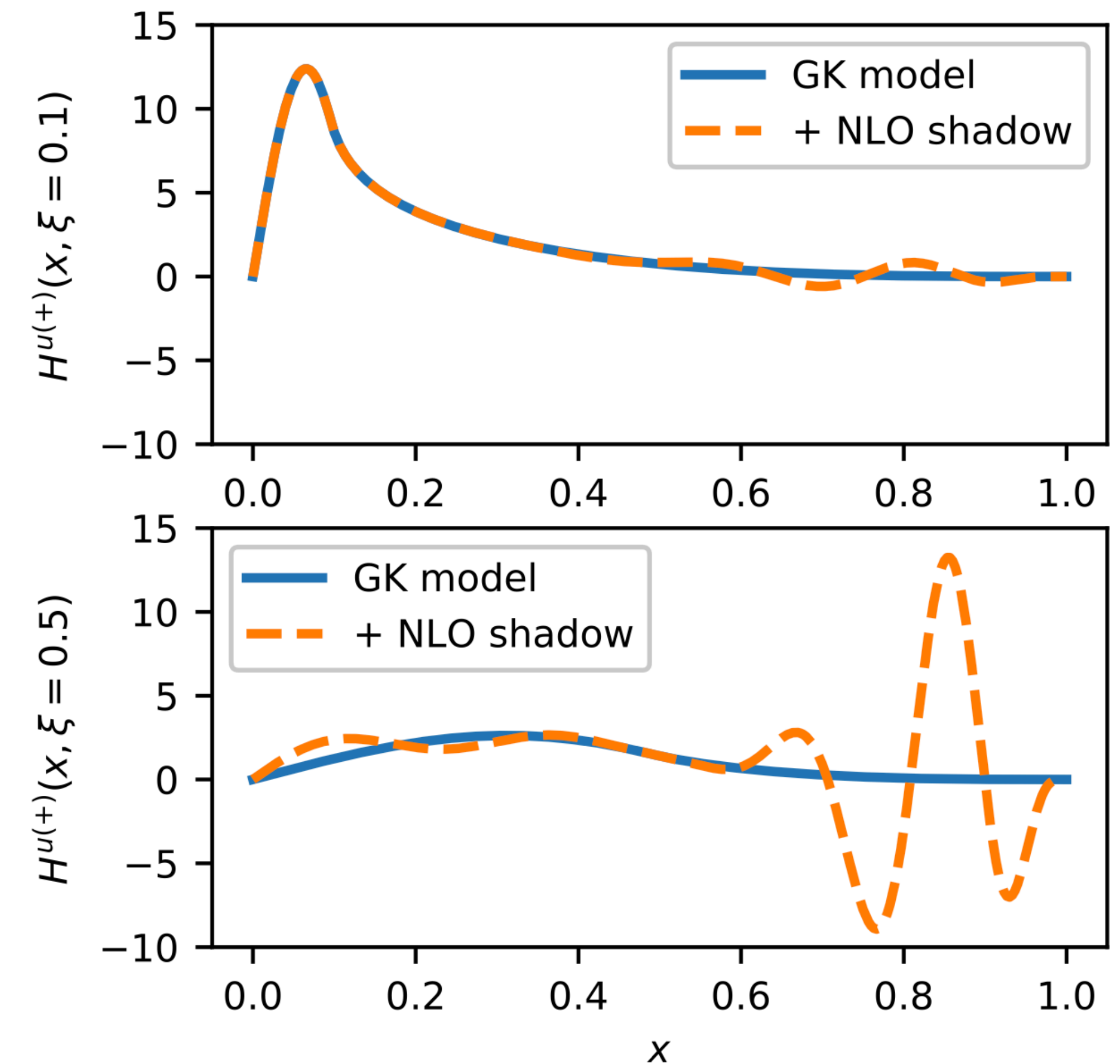
- at the initial scale do not contribute to both PDFs and CFFs
- at some other scale they contribute negligibly

making the deconvolution of CFFs ill-posed

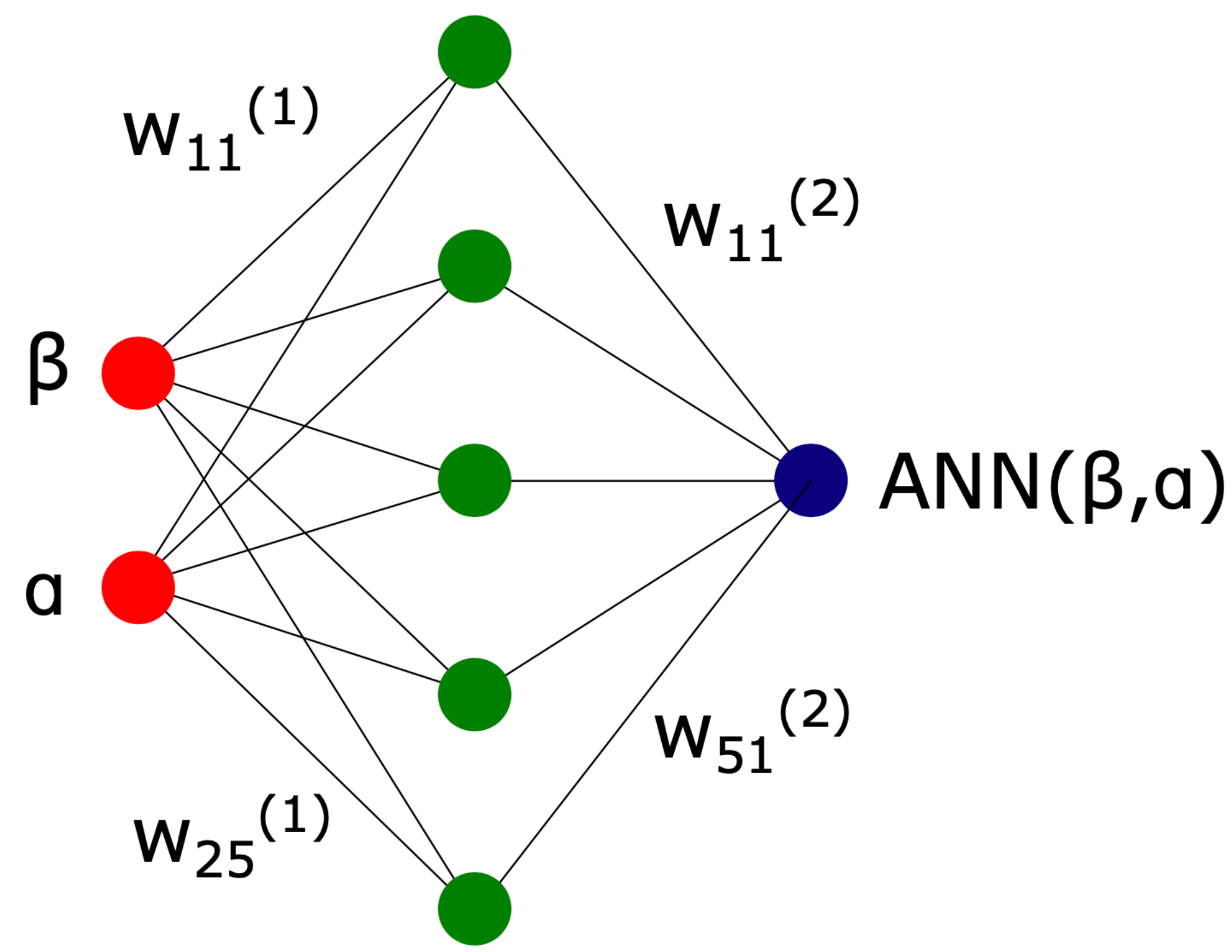
We found such GPDs for both LO and NLO

For more details see:

→ [V. Bertone et al., *PRD* 103 \(2021\) 11, 114019](#)



Our ANNs:

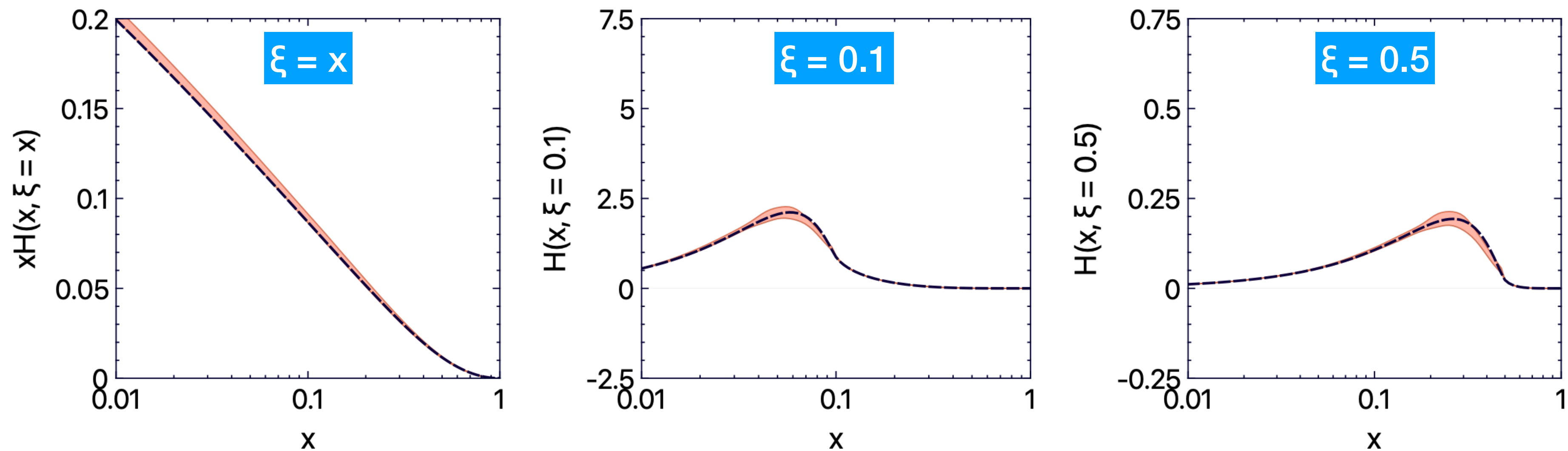


Requirements:

- symmetric w.r.t. α
- symmetric w.r.t. β
- vanishes at $|\alpha| + |\beta| = 1$

Activation function:

$$\left(\varphi_i \left(w_i^\beta |\beta| + w_i^\alpha \alpha / (1 - |\beta|) + b_i \right) - \varphi_i \left(w_i^\beta |\beta| + w_i^\alpha + b_i \right) \right) + (w^\alpha \rightarrow -w^\alpha)$$




Conditions:

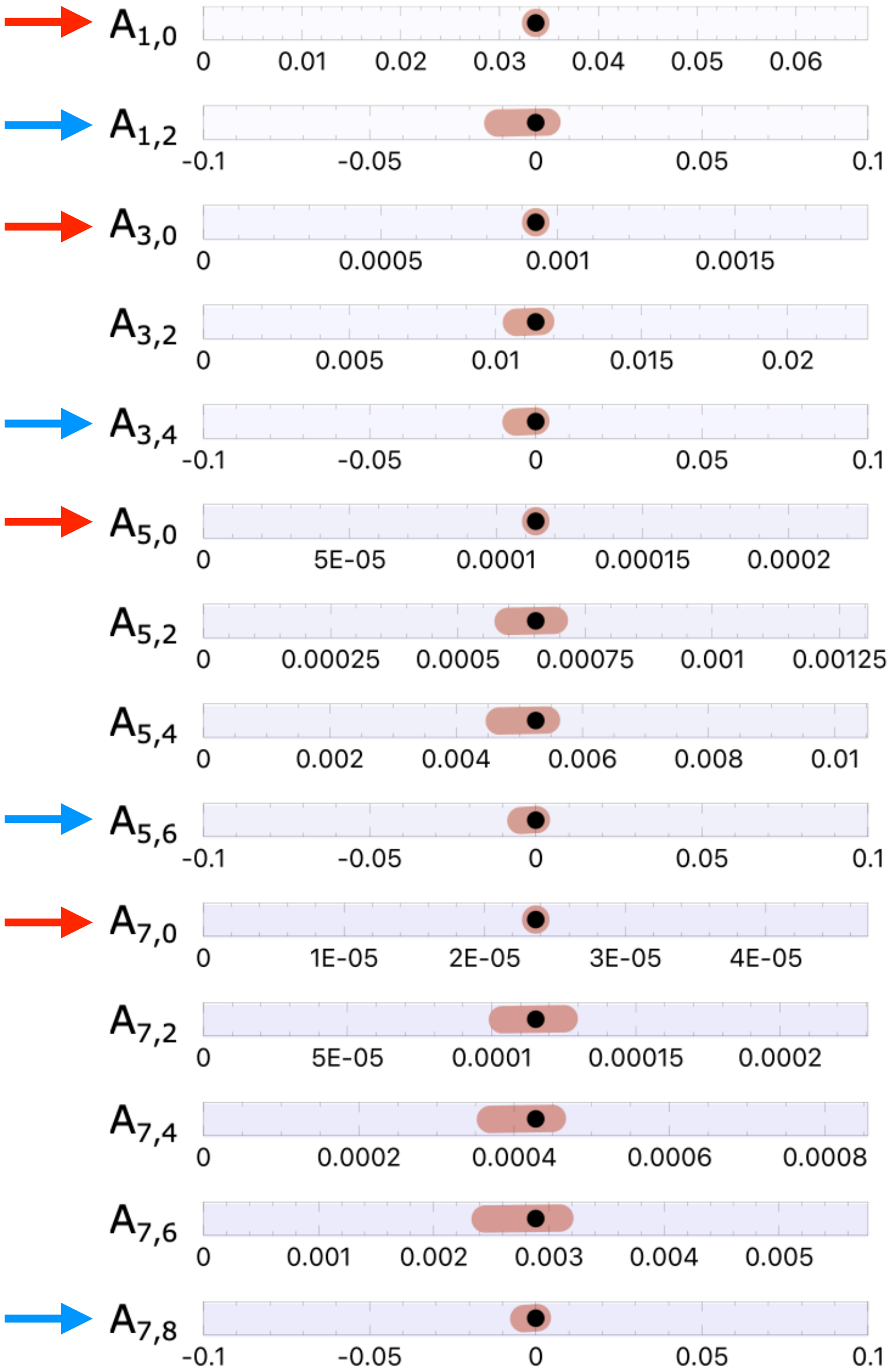
- Input: 400 $x \neq \xi$ points generated with GK model
- Positivity not forced

Technical detail of the analysis:

- Minimisation with genetic algorithm
- Replication for estimation of model uncertainties
- “Local” detection of outliers
- Dropout algorithm for regularisation

--- GK

 ANN model
68% CL
 $F_C + F_S + F_D$



● GK

ANN model
68% CL
 $F_C + F_S + F_D$

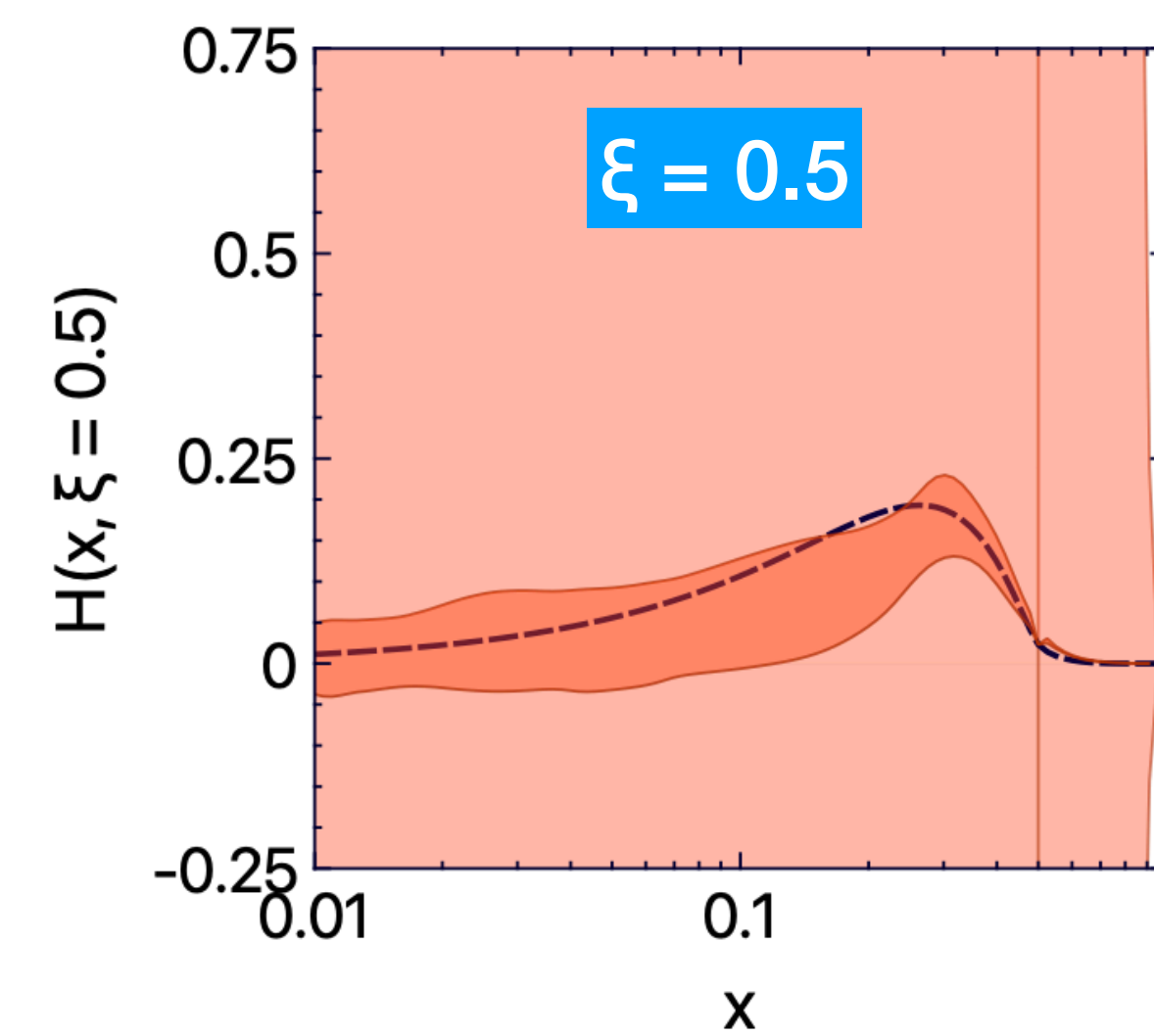
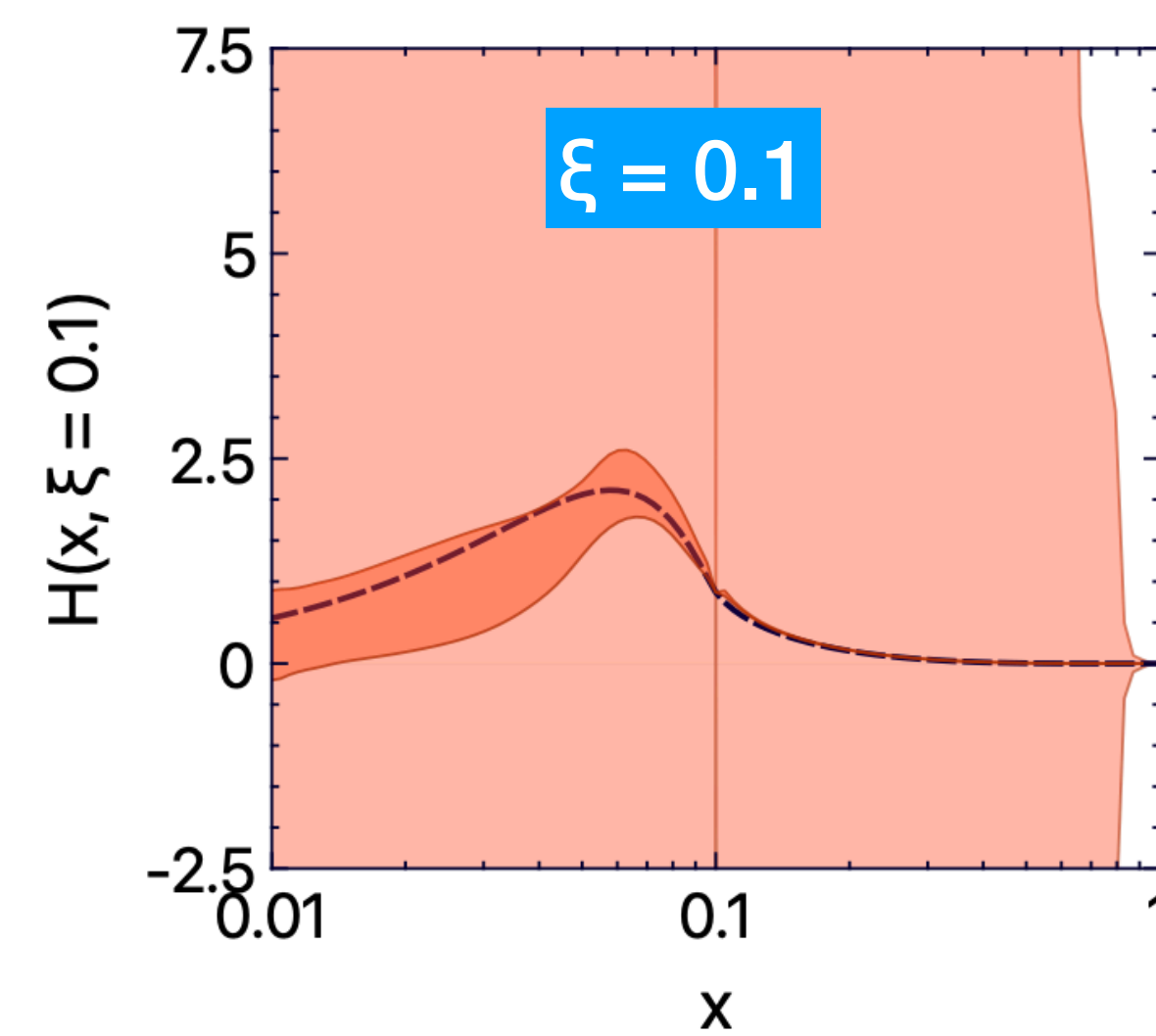
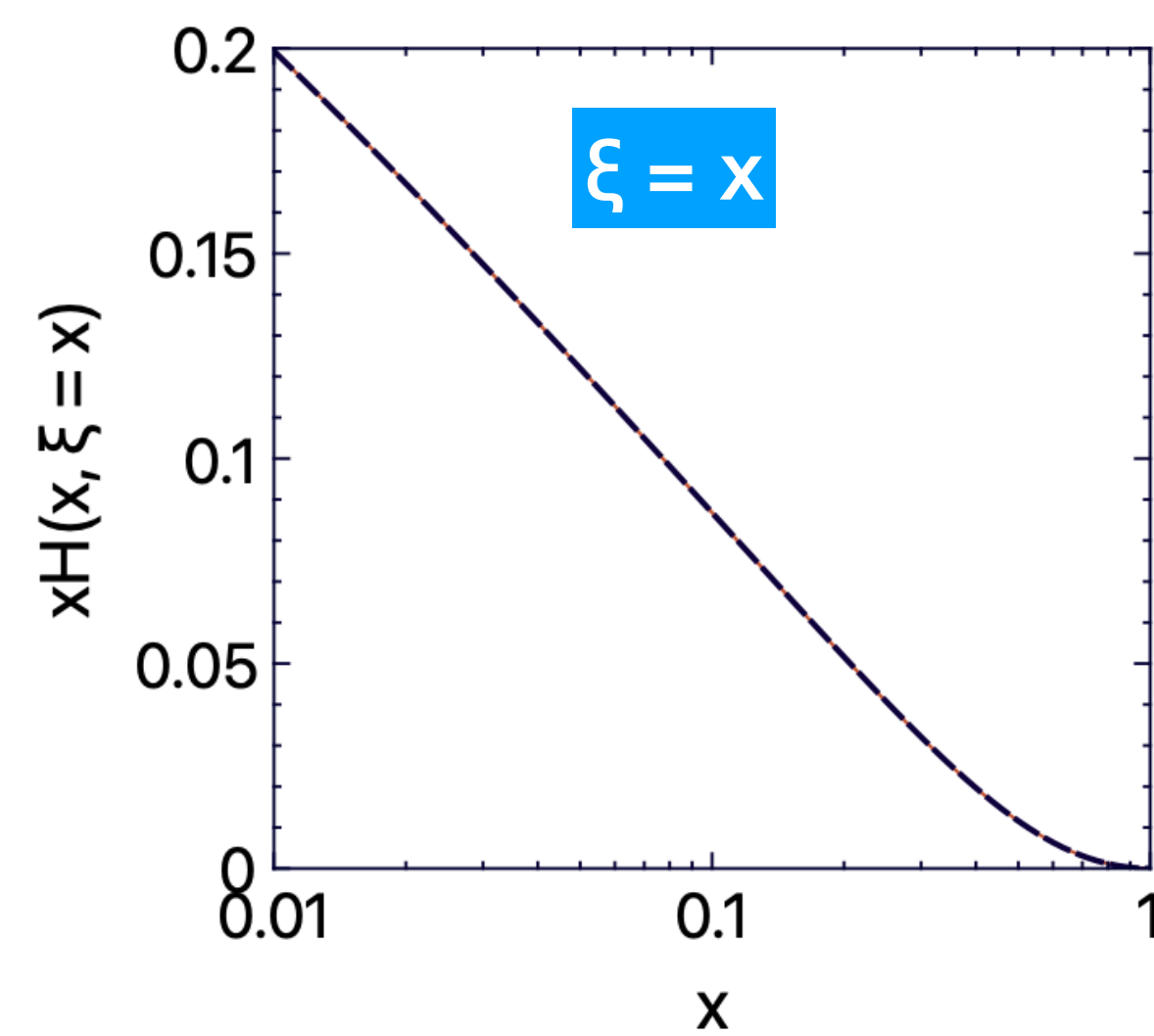
Mellin mom. coefficients:

→ related to PDF

→ related to D-term

Conditions:


- Input: 400 $x \neq \xi$ points generated with GK model
- Positivity not forced




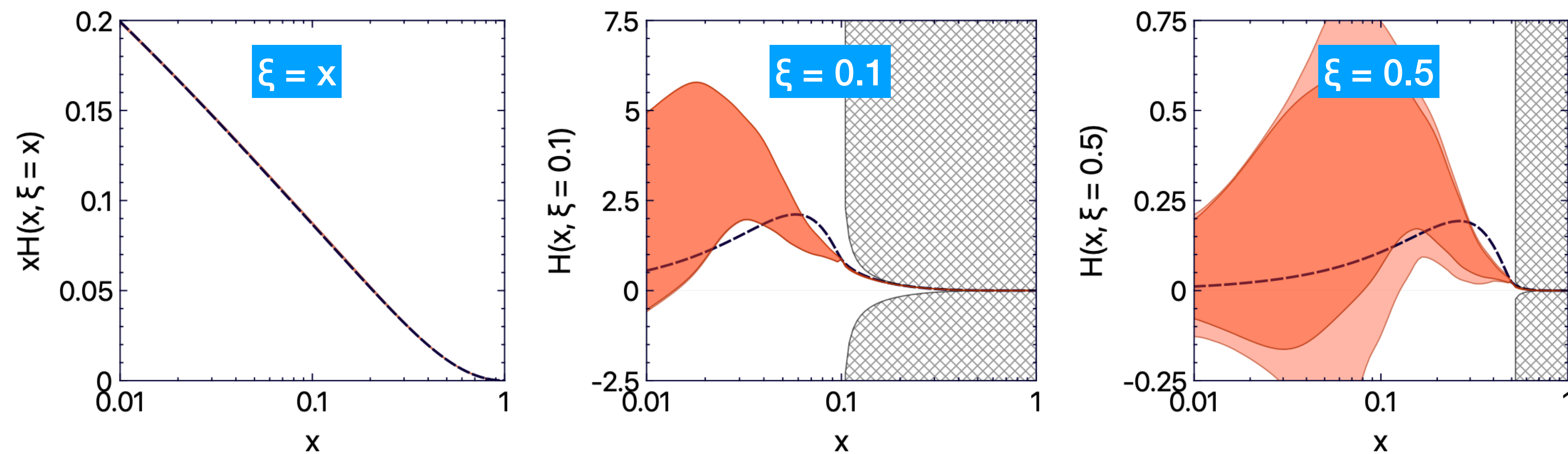
Conditions:

- Input: 200 $x = \xi$ points generated with GK model
- Positivity not forced

--- GK

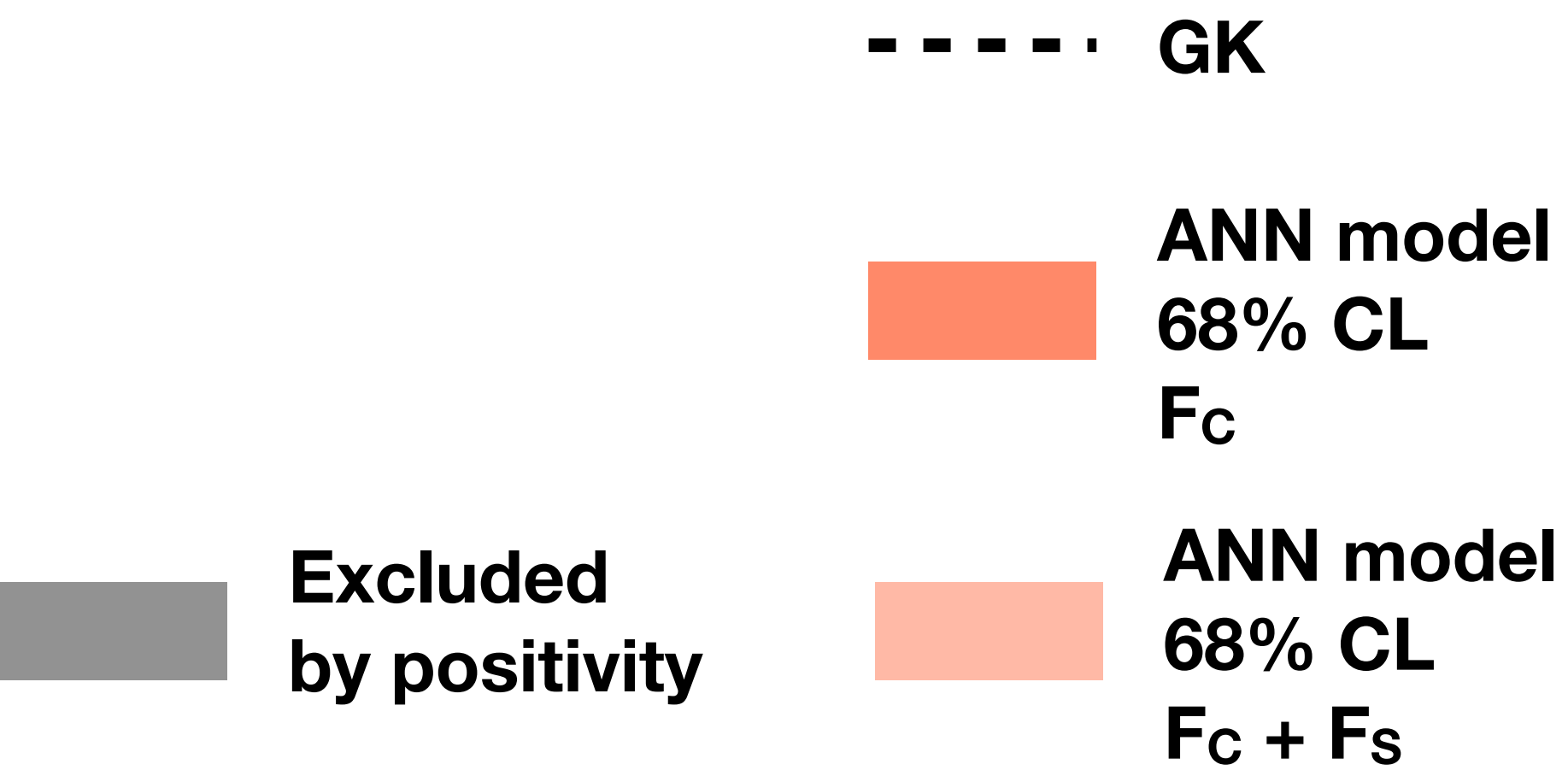
 ANN model
68% CL
 F_c

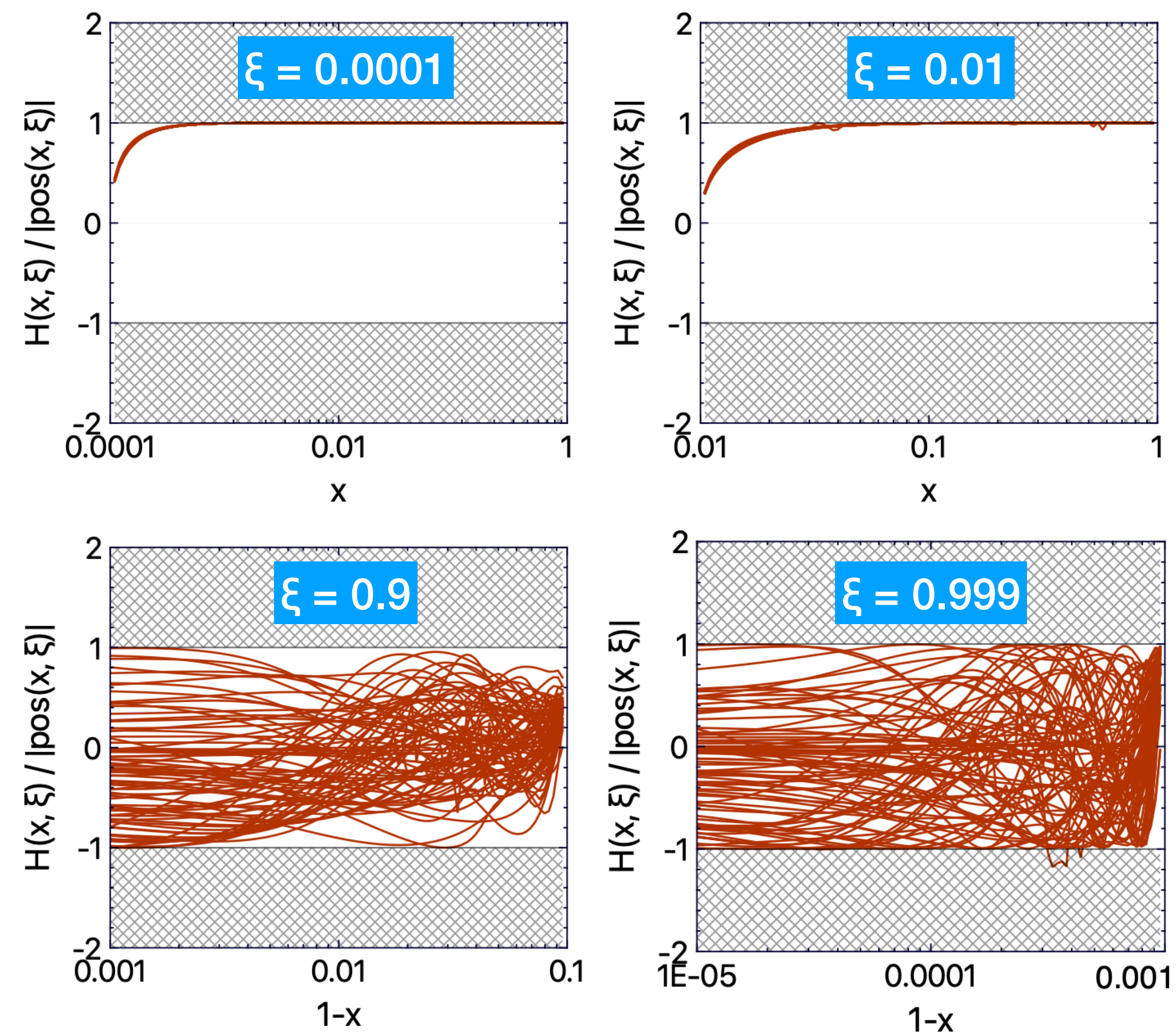
 ANN model
68% CL
 $F_c + F_s$



Conditions:

- Input: 200 $x = \xi$ points generated with GK model
- Positivity **forced**





Conditions:

- Input: 200 $x = \xi$ points generated with GK model
- Positivity **forced**



- For the first time, we propose modelling GPDs based on ANNs
→ new, nontrivial and timely analysis
- Our modelling fulfils all theory-driven constraints (including positivity)
→ subject not touched enough in the current literature
- Can easily accommodate lattice-QCD results
→ important to include additional sources of GPD information
- **These is the new tool to address the long-standing problem of model dependency of GPDs**

- Novel MC generator called EpIC released
 - E.C. Aschenauer et al., hep-ph/2205.01762
 - <https://pawelsznajder.github.io/epic>
- EpIC is based on PARTONS (note: v3 version of PARTONS is now available!)
 - B. Berthou, EPJC 78 (2018) 6, 478
 - <https://partons.cea.fr>
- EpIC is characterised by:
 - flexible architecture that utilises a modular programming paradigm
 - a variety of modelling options, including radiative corrections
 - multichannel capability (initial version includes DVCS, TCS and DVMP)
- **This is the new tool to be use in the precision era commenced by the new generation of experiments**



Backup

Polynomiality:

$$\mathcal{A}_n(\xi) = \int_{-1}^1 dx x^n H(x, \xi) = \sum_{\substack{j=0 \\ \text{even}}}^n \xi^j A_{n,j} + \text{mod}(n, 2) \xi^{n+1} A_{n,n+1}$$

Let us express GPD by:

$$H^N(x, \xi) = \sum_{\substack{j=0 \\ \text{even}}}^N f_j(x) \xi^j$$

only even j as there is no odd power of ξ in polynomiality expansion

Support:

$$f_j(-1) = f_j(1) = 0$$

we want GPDs to vanish at |x| = 1

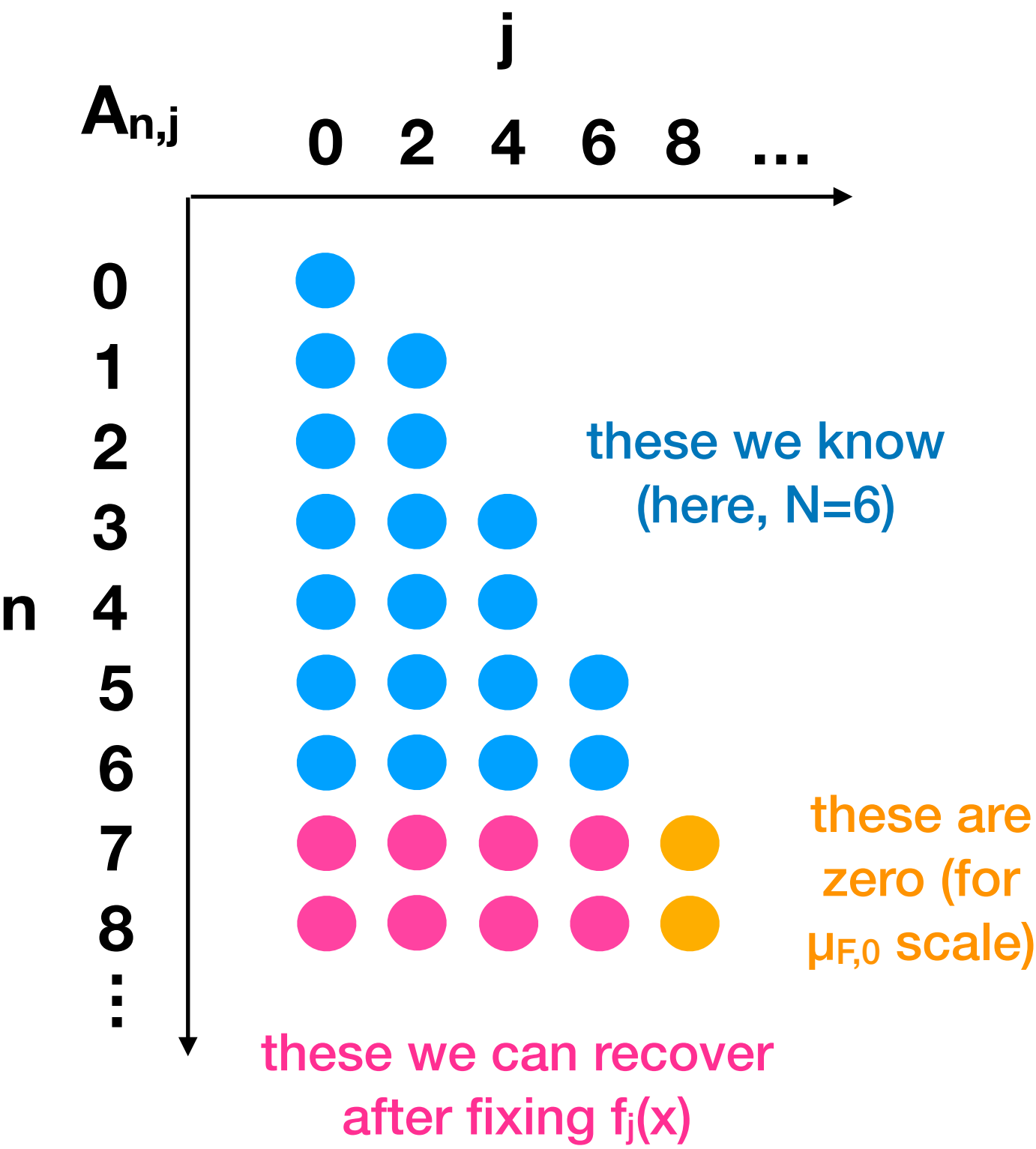
Mellin coefficients:

$$A_{n,j} = \int_{-1}^1 dx x^n f_j(x)$$

choice of f_j(x) functional form is arbitrary

where e.g.:

$$A_{0,2} = \int_{-1}^1 dx f_0(x) = 0$$



Backup

Polynomial basis: This basis leads to Dual Parameterisation → M. Polyakov, A. Shuvaev, hep-ph/0207153

Any attempt of describing GPDs by orthogonal polynomials will lead to this basis

$$f_j(x) = \sum_{i=0}^{N+2} w_{i,j} x^i$$

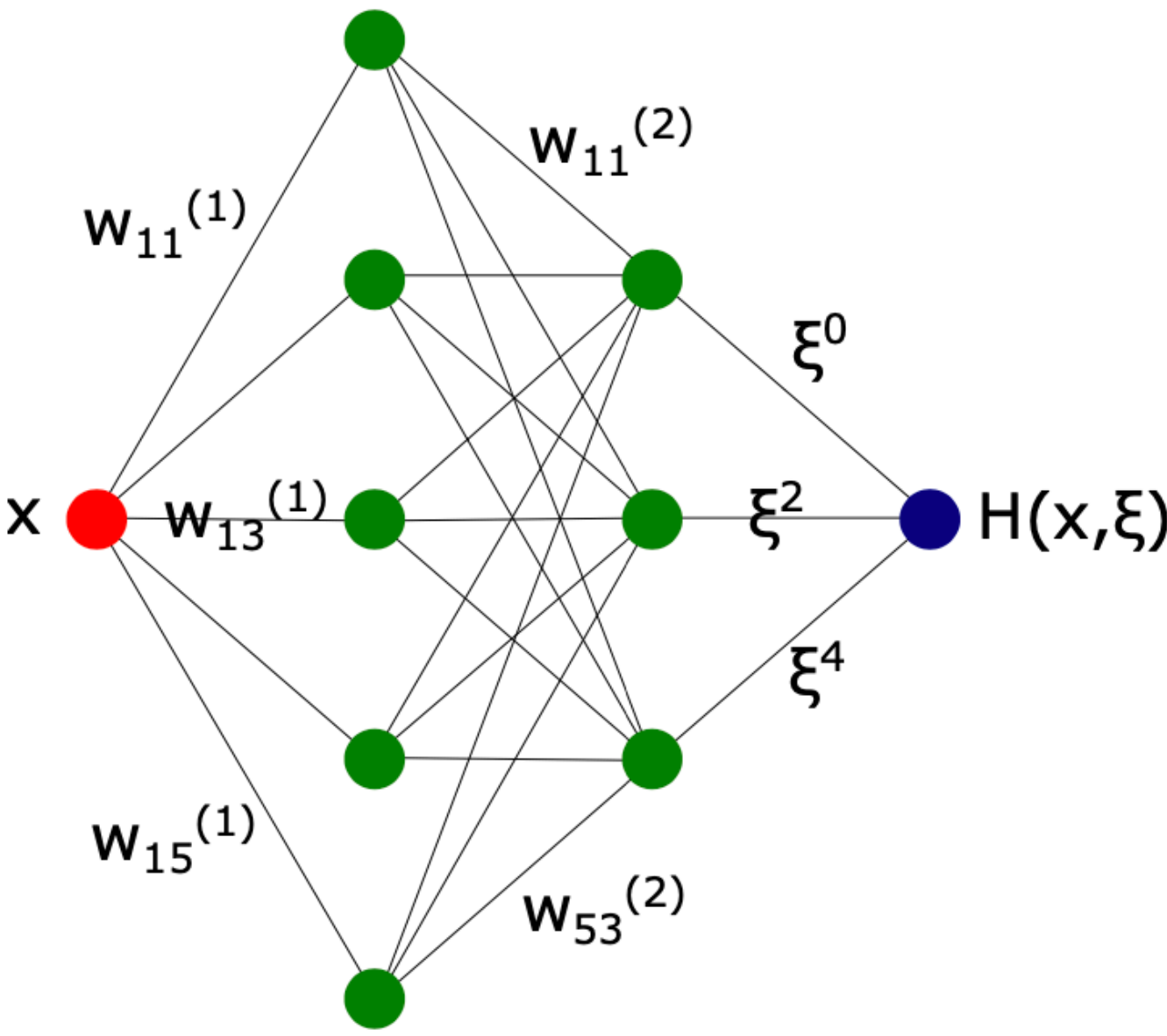
GPD will be expressed by sum of monomials $x^i \xi^j$

ANN basis: New!

We can describe GPD by a single ANN

$$f_j(x) = \text{ANN}_j(x)$$

GPD will be expressed by sum of ANNs multiplied by ξ^j



Backup

$\xi = 0$

$\xi = 0.5$

Test model
(see e.g.: [hep-ph/2110.06052](https://arxiv.org/abs/hep-ph/2110.06052)):

$$H_\pi(x, \xi) =$$

$$\Theta(x - |\xi|) \frac{30(1 - x)^2(x^2 - \xi^2)}{(1 - \xi^2)^2} +$$

$$\Theta(|\xi| - |x|) \frac{15(1 - x)(\xi^2 - x^2)(x + 2x\xi + \xi^2)}{2\xi^3(1 + \xi)^2}$$

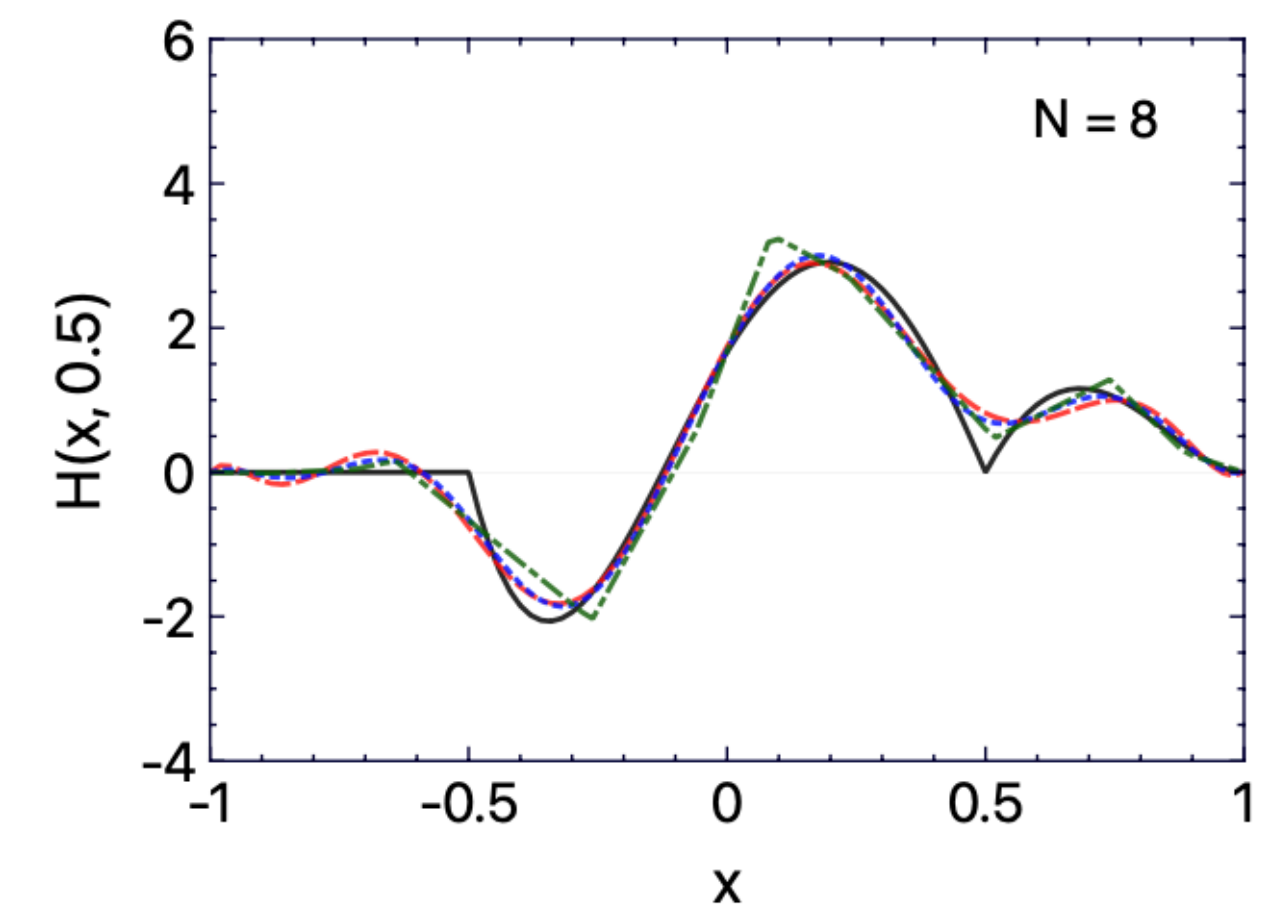
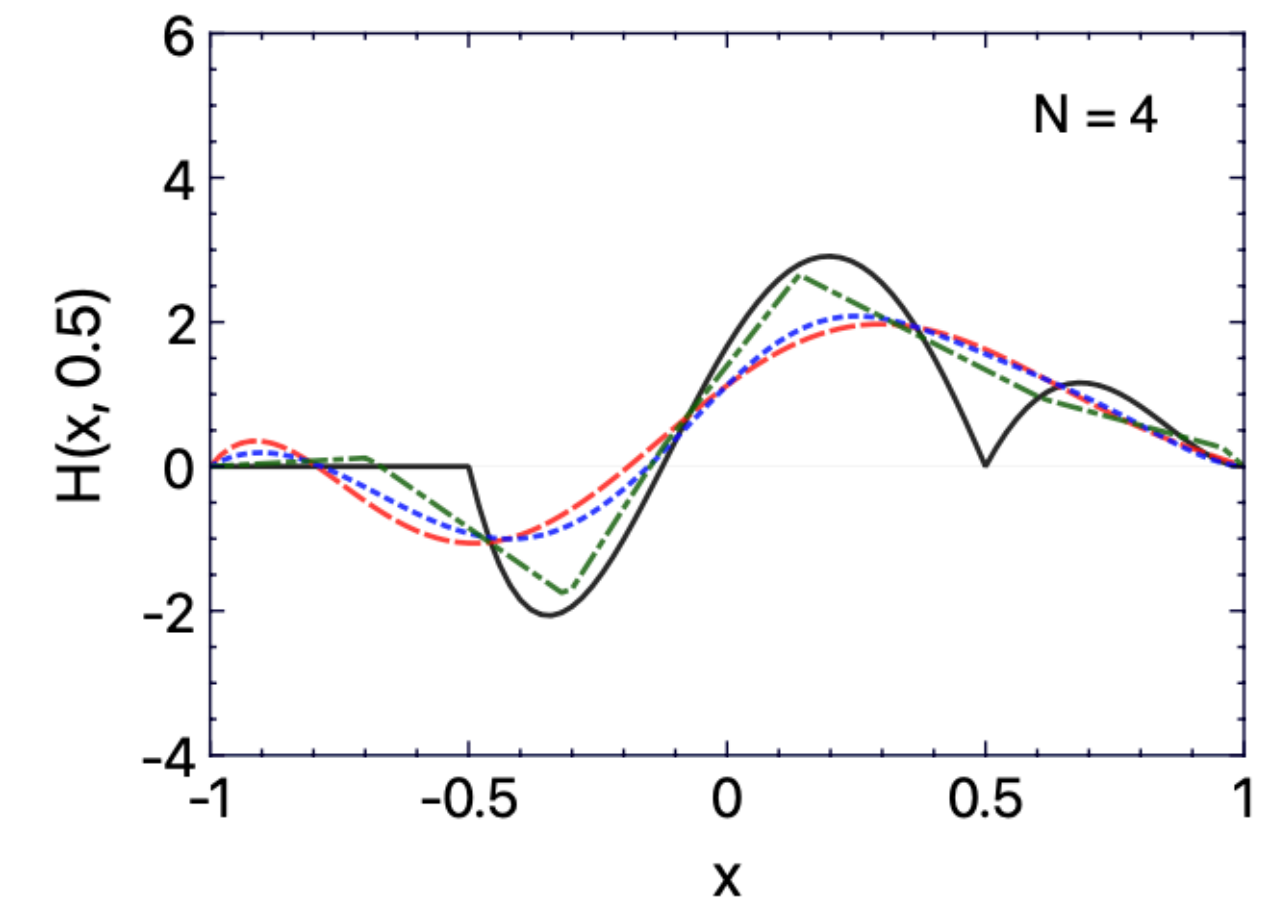
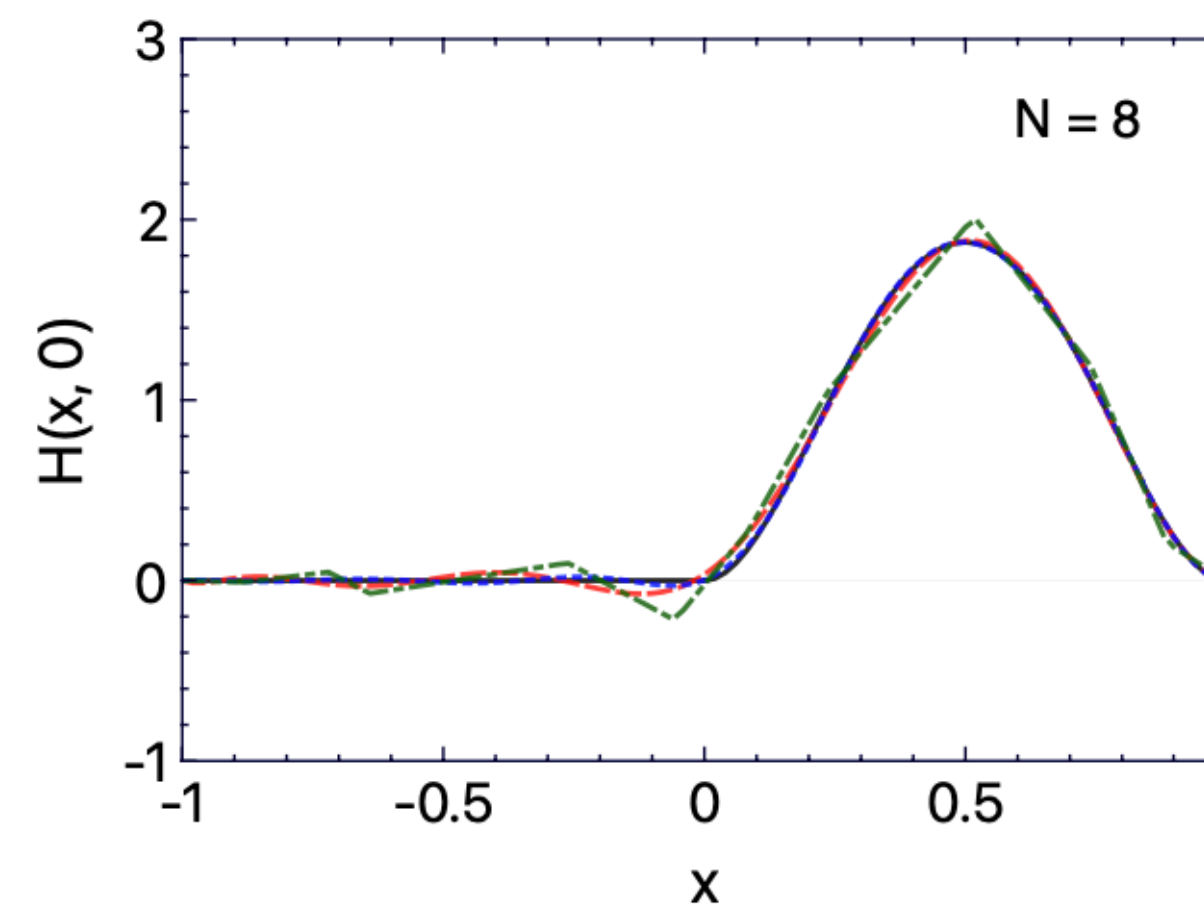
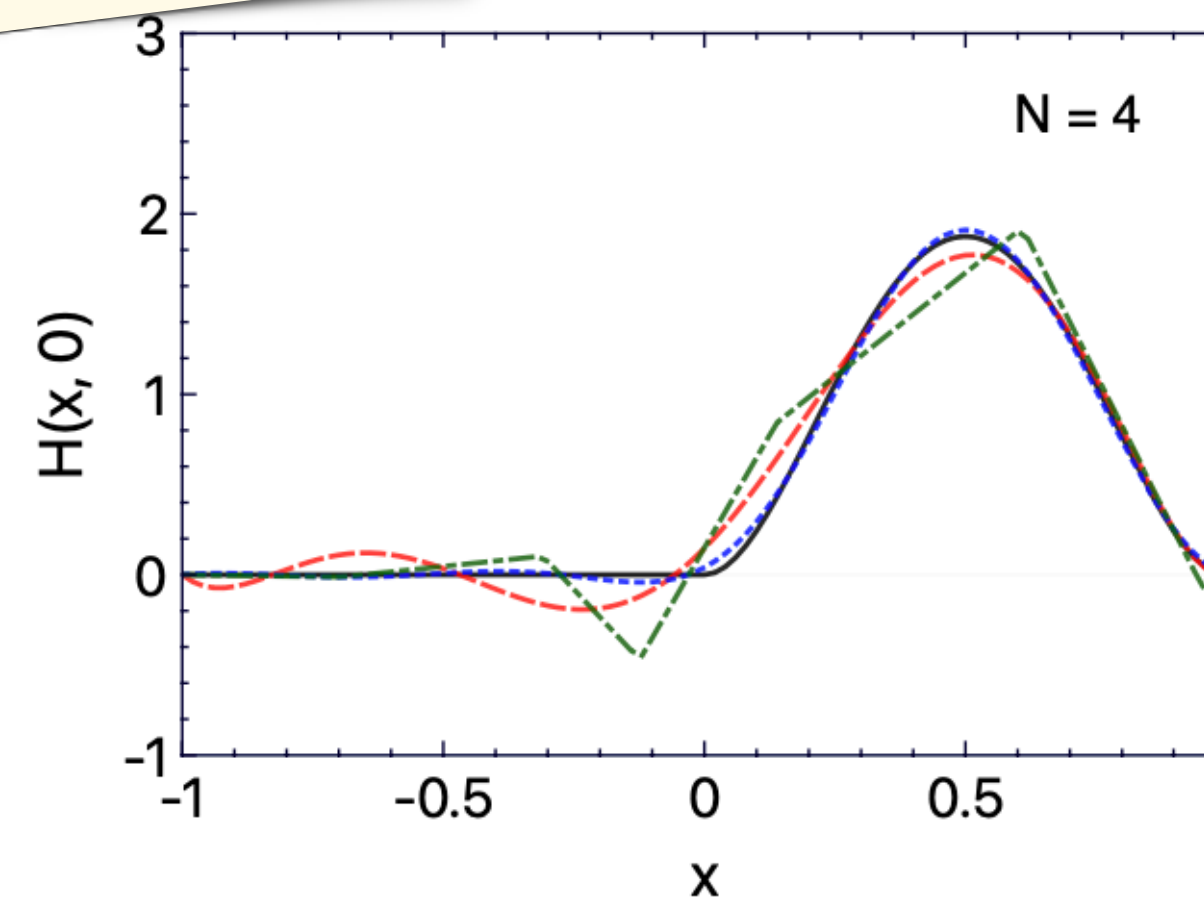
Polynomial basis

ANN basis - sigmoid

$$\varphi_k^{(2)}(\cdot) = \frac{1}{1 + \exp(-(\cdot))}$$

ANN basis - ReLU

$$\varphi_k^{(2)}(\cdot) = (\cdot) \Theta(\cdot)$$



Note:

- positivity not enforced here
- few extensions of this modelling possible, see the next slide

Basic:

$$H(x, \xi) = \sum_{\substack{j=0 \\ \text{even}}}^N f_j(x) \xi^j$$

With explicit PDF:

$$H(x, \xi) = q(x) + \sum_{\substack{j=2 \\ \text{even}}}^N f_j(x) \xi^j$$

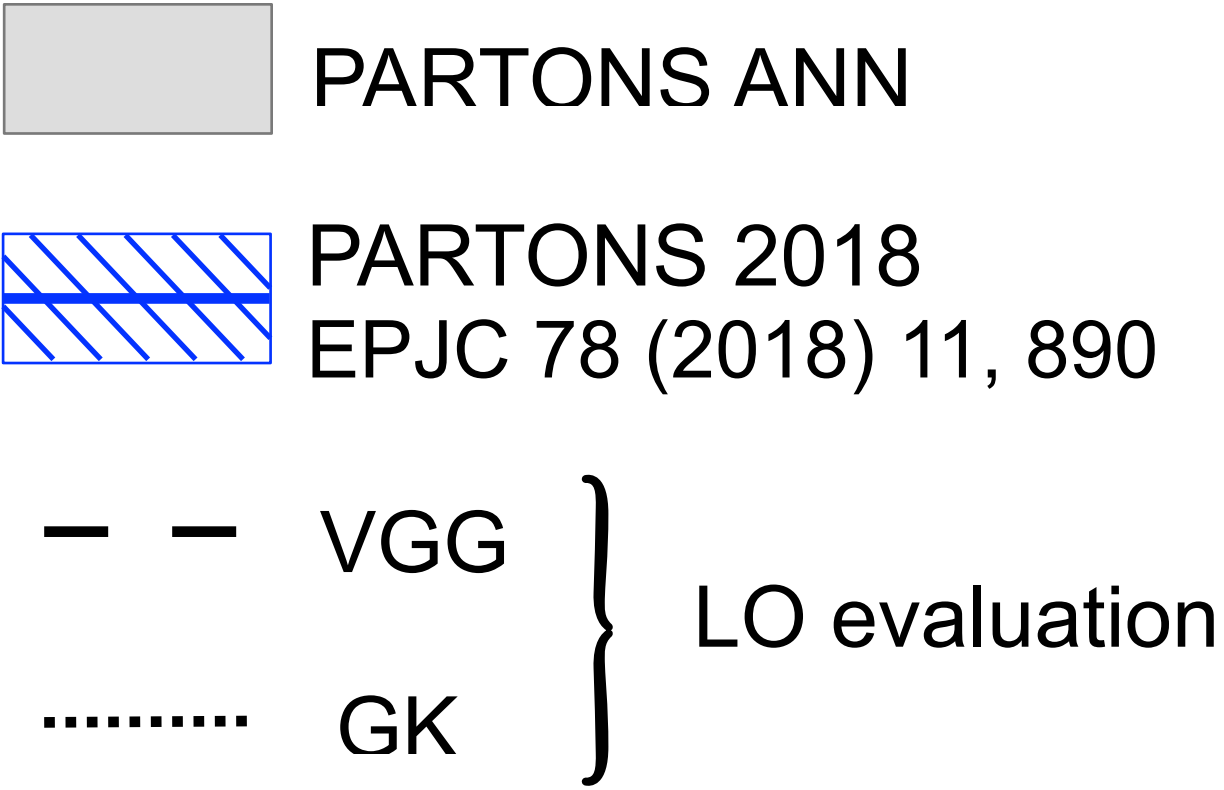
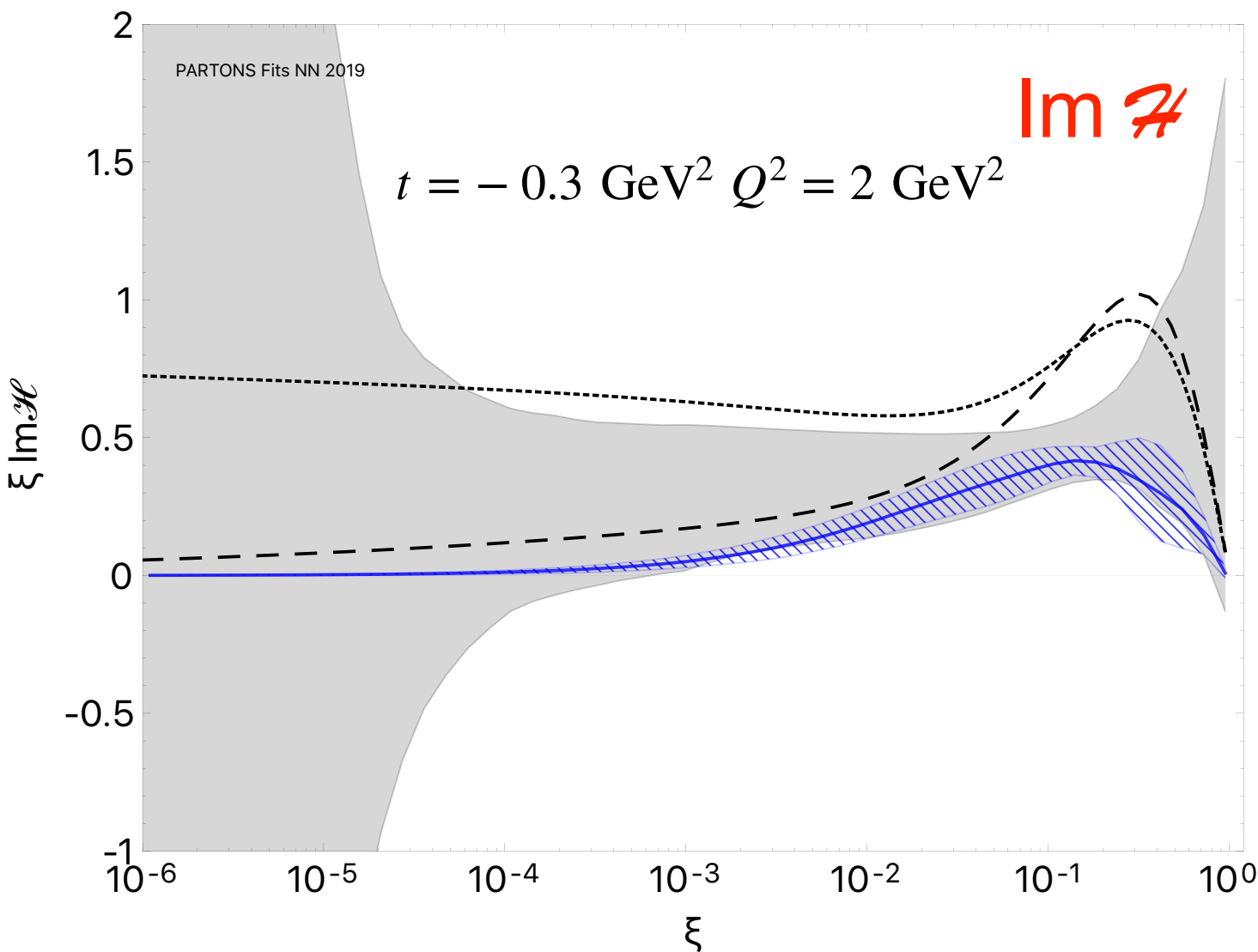
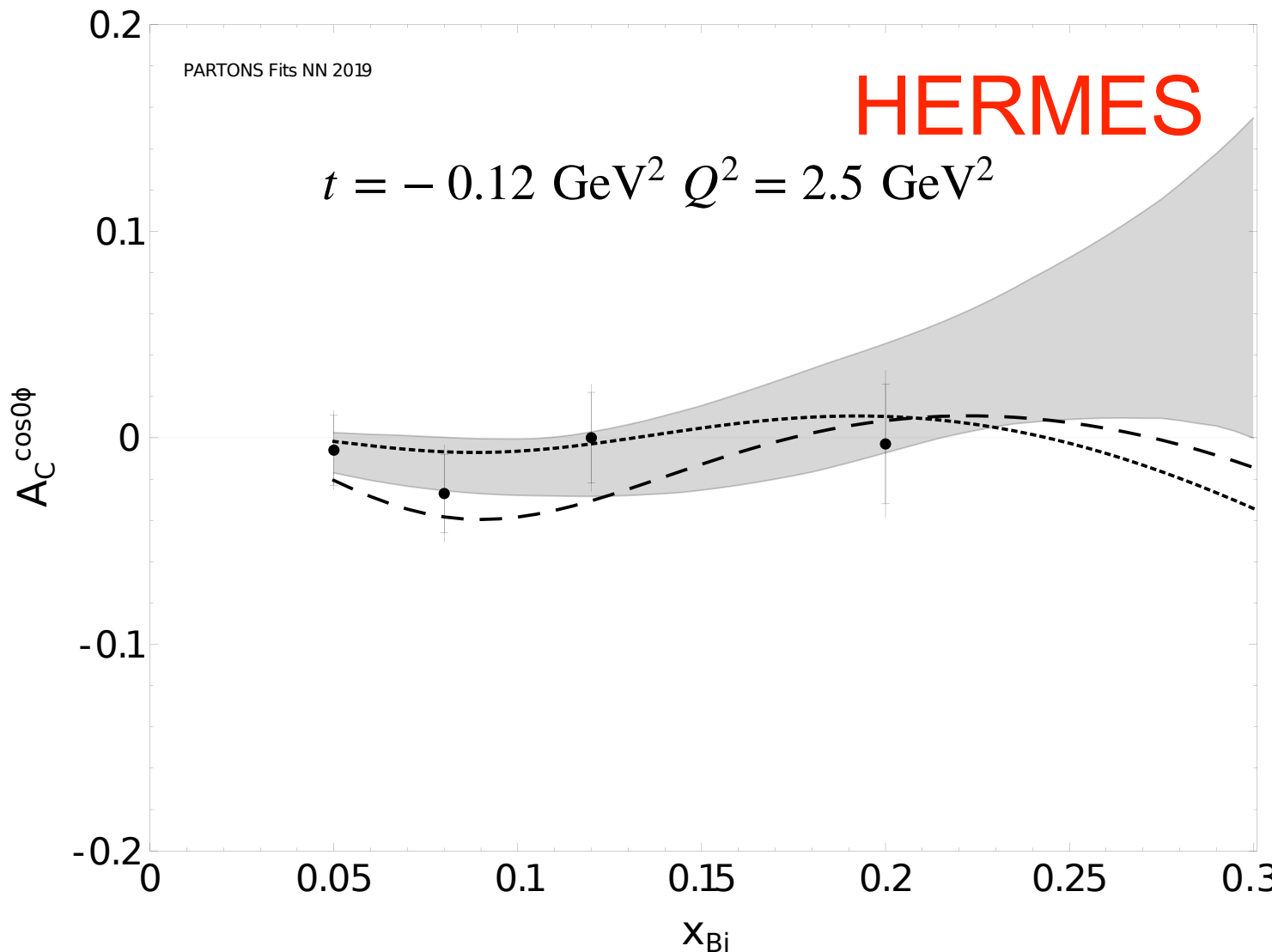
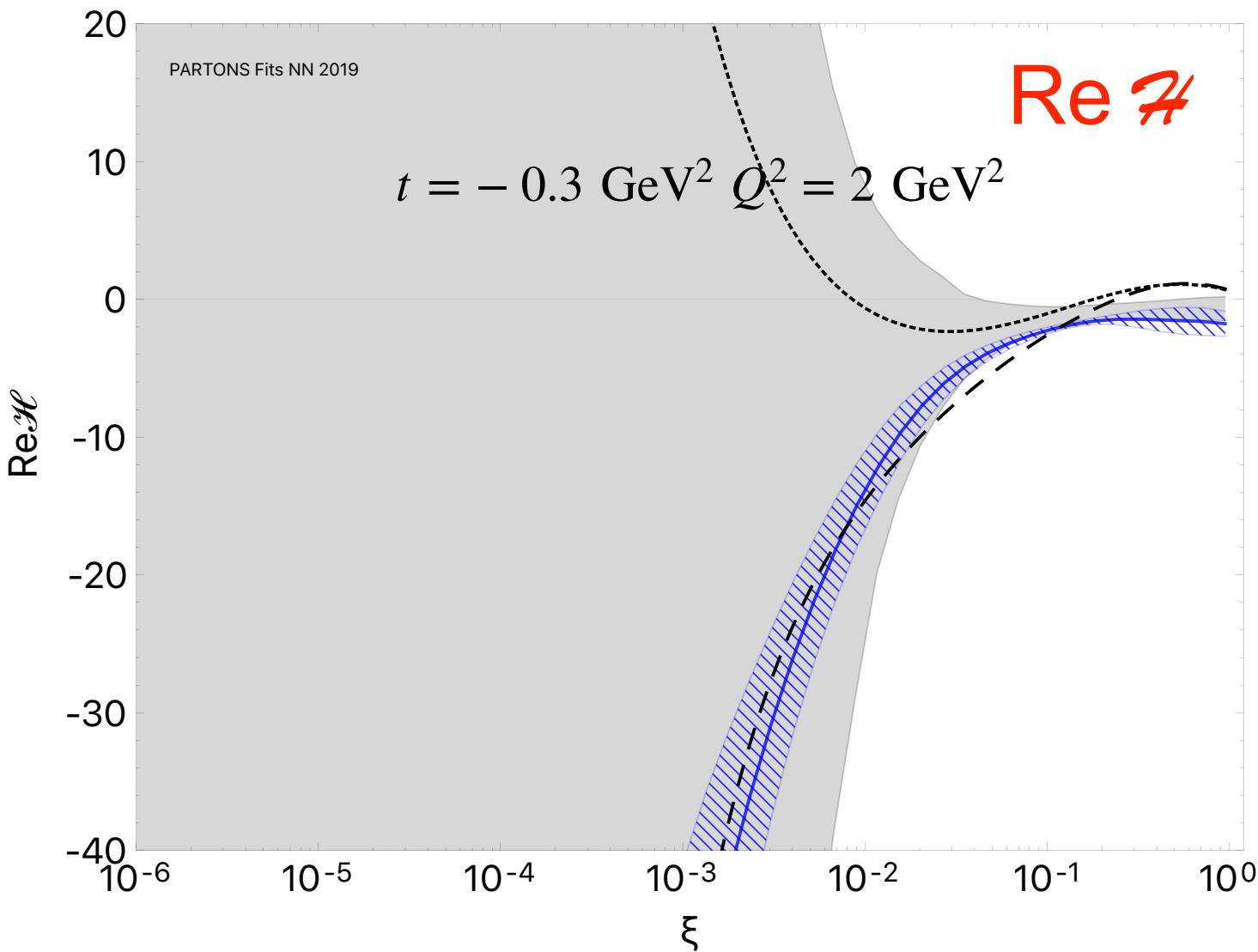
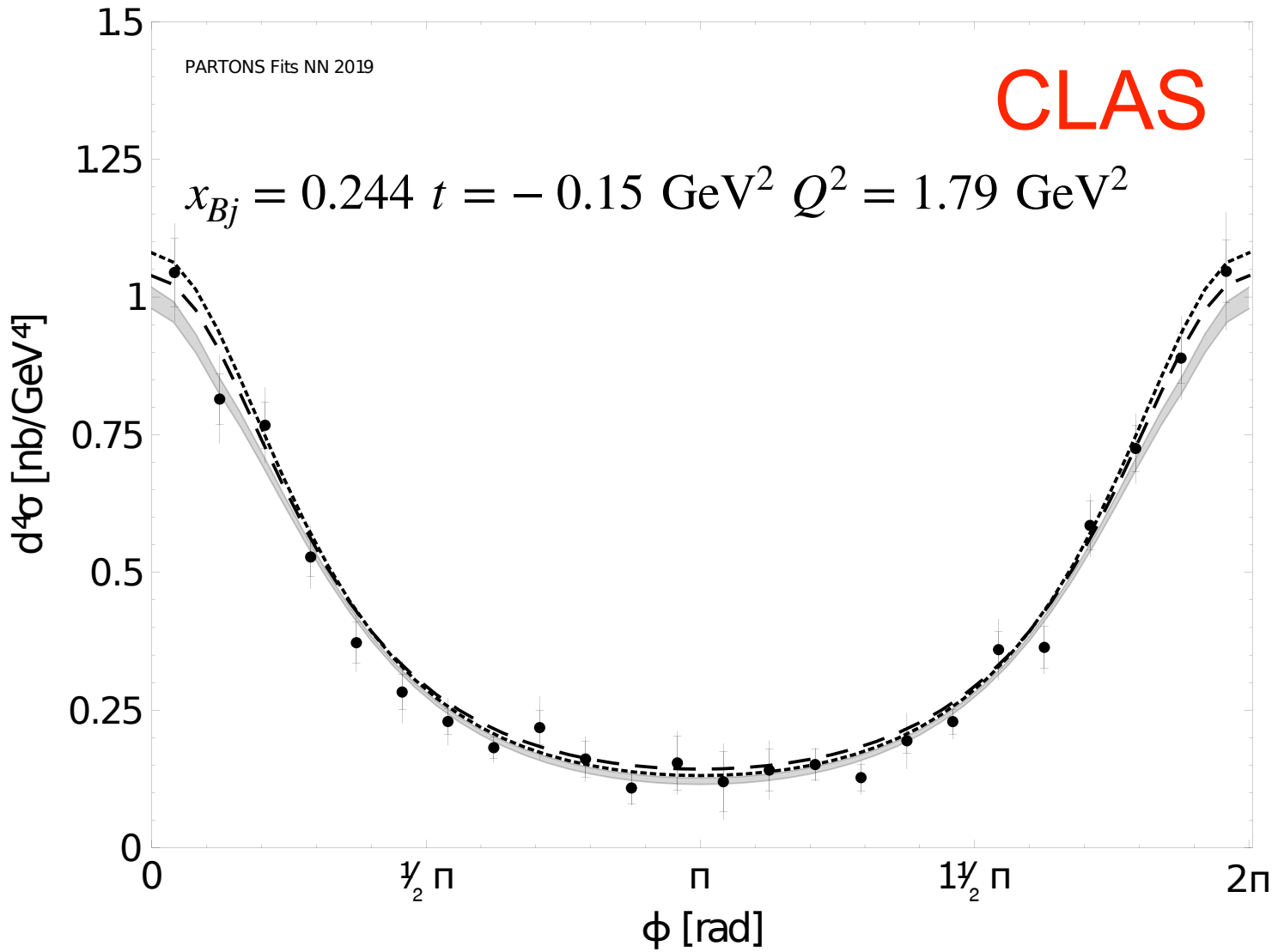
Vanishing at $x=x_i$:

$$H(x, \xi) = (x^2 - \xi^2) \sum_{\substack{j=0 \\ \text{even}}}^N f_j(x) \xi^j$$

With D-term:

$$H(x, \xi) = D_{\text{term}}(x/\xi) + \sum_{\substack{j=0 \\ \text{even}}}^N f_j(x) \xi^j$$

Backup



$\xi \approx x_{Bj}/(2 - x_{Bj})$

Dispersion relation:

$$\mathcal{C}_H(t, Q^2) = \text{Re } \mathcal{H}(\xi, t, Q^2) - \frac{1}{\pi} \int_0^1 d\xi' \text{Im } \mathcal{H}(\xi', t, Q^2) \left(\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'} \right)$$

Relation between subtraction constant and D-term ($z=x/\xi$):

$$\mathcal{C}_H(t, Q^2) \stackrel{LO}{=} 2 \sum_q e_q^2 \int_{-1}^1 dz \frac{D_{\text{term}}^q(z, t, \mu_F^2 \equiv Q^2)}{1 - z}$$

Decomposition into Gegenbauer polynomials:

$$D_{\text{term}}^q(z, t, \mu_F^2) = (1 - z^2) \sum_{\text{odd } n} d_n^q(t, \mu_F^2) C_n^{3/2}(z)$$

Finally:

$$\mathcal{C}_H(t, Q^2) \stackrel{LO}{=} 4 \sum_q e_q^2 \sum_{\text{odd } n} d_n^q(t, \mu_F^2 \equiv Q^2)$$

Connection to EMT FF:

$$d_1^q(t, \mu_F^2) = 5C_q(t, \mu_F^2)$$

Master formula:

$$\text{Re}\mathcal{H}(\xi, t, Q^2) - \frac{1}{\pi} \int_0^1 d\xi' \text{Im} \mathcal{H}(\xi, t, Q^2) \left(\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'} \right) \stackrel{LO}{=} 4 \sum_q e_q^2 \sum_{\text{odd } n} d_n^q(t, \mu_F^2 \equiv Q^2)$$

Extraction of subtraction constant from DVCS data requires:

- integral over ξ (alternatively: x_{Bj} or ν) between ε and 1
 - good knowledge of both Re and Im parts of CFF H
- $\varepsilon = 10^{-6}$

Model assumptions to extract EMT FF C from subtraction constant:

- truncation to d1
- sensitivity to gluon contribution via evolution

$$C_H(t, Q^2) = 4 \sum_q e_q^2 d_1^q(t, \mu_F^2 \equiv Q^2)$$

$$d_1^G(t, \mu_{F,0}^2) = 0 \quad \mu_{F,0}^2 = 0.1 \text{ GeV}^2$$

- symmetry of light quark contributions

$$d_1^u(t, \mu_F^2) = d_1^d(t, \mu_F^2) = d_1^s(t, \mu_F^2) \equiv d_1^{uds}(t, \mu_F^2)$$

- tripole Ansatz for t-dependence

$$d_1^{uds}(t, \mu_F^2) = d_1^{uds}(\mu_F^2) \left(1 - \frac{t}{\Lambda^2} \right)^{-\alpha} \quad \begin{array}{l} \alpha = 3 \\ \Lambda = 0.8 \text{ GeV} \end{array}$$

Backup

Subtraction constant:

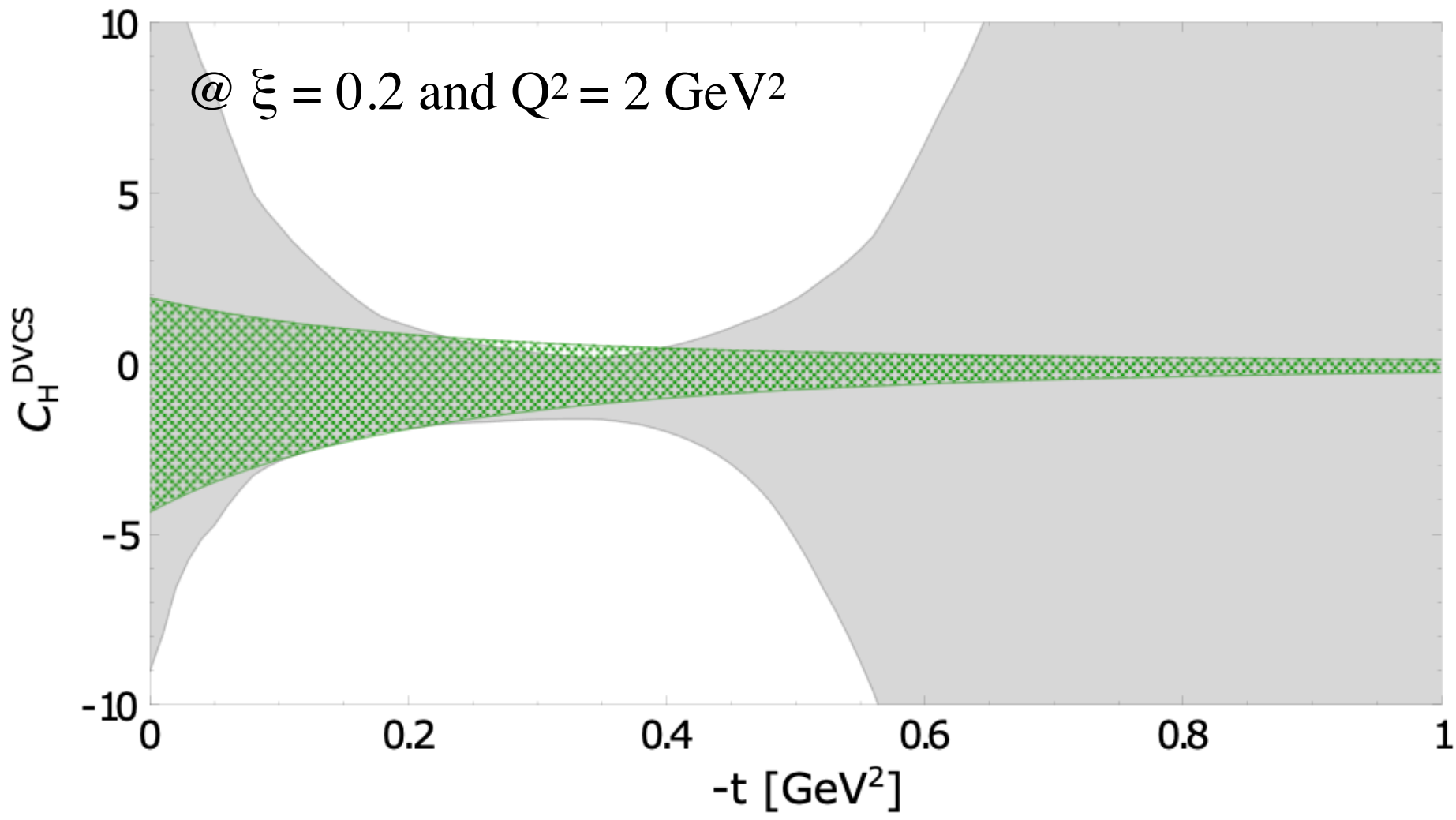


ANN analysis



Model dependent extraction

$$d_1^{uds}(t, \mu_F^2) = d_1^{uds}(\mu_F^2) \left(1 - \frac{t}{\Lambda^2}\right)^{-\alpha} \quad \begin{matrix} \alpha = 3 \\ \Lambda = 0.8 \text{ GeV} \end{matrix}$$



Results:

Parameter	Value
$d_1^{uds}(\mu_F^2 = 2 \text{ GeV}^2)$	-0.5 ± 1.2
$d_1^c(\mu_F^2 = 2 \text{ GeV}^2)$	-0.0020 ± 0.0053
$d_1^g(\mu_F^2 = 2 \text{ GeV}^2)$	-0.6 ± 1.6