

Hadron production in electron-positron annihilation

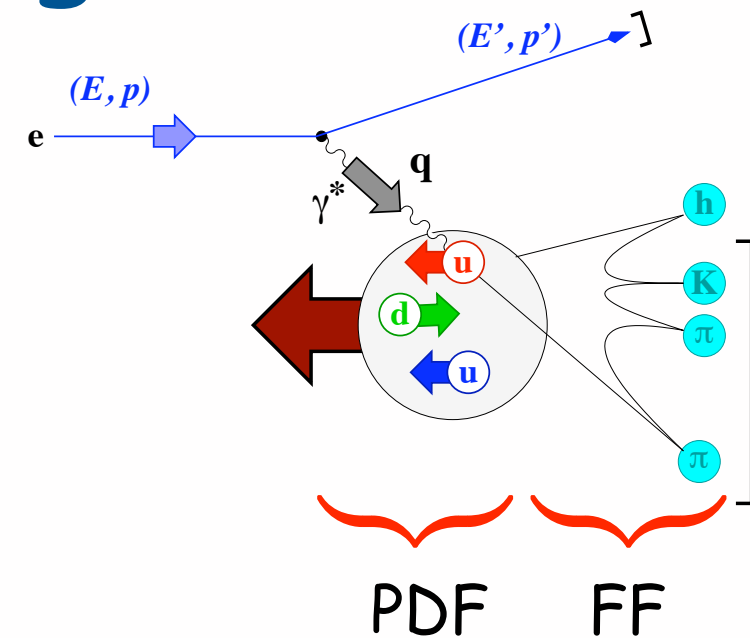
highlights from the past five years

Probing parton dist's through fragmentation

nucleon pol.

quark pol.

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

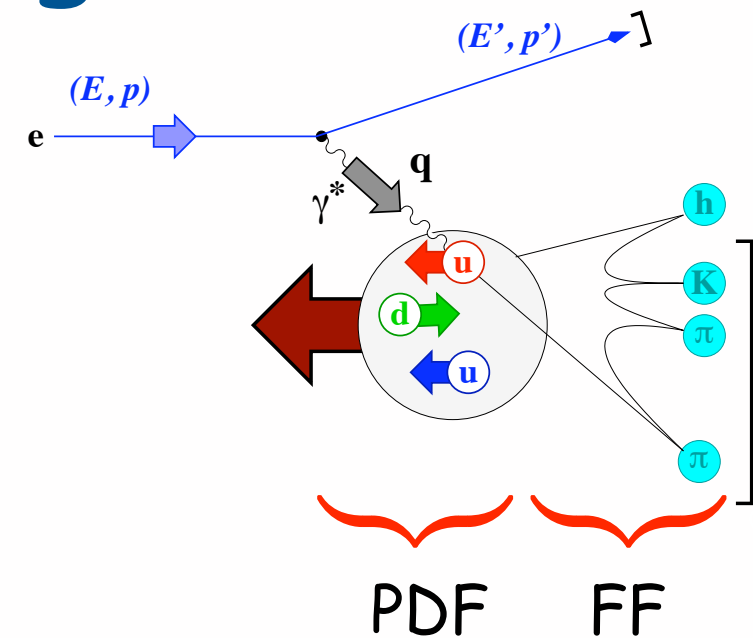


in SIDIS*) couple PDFs to:

*) semi-inclusive DIS with unpolarized final state

Probing parton dist's through fragmentation

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



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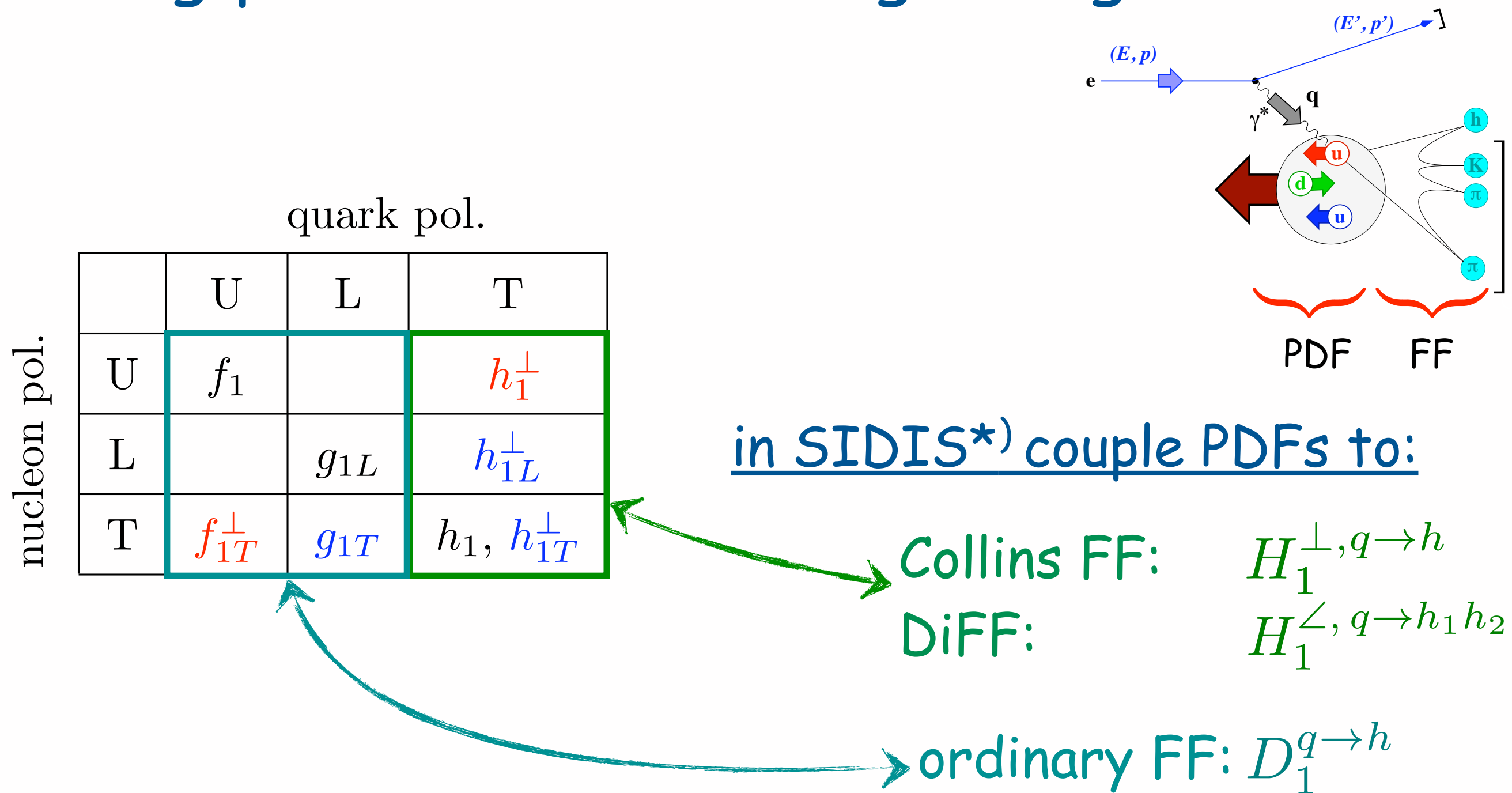
Collins FF:
DiFF:

$$H_1^\perp, q \rightarrow h$$

$$H_1^\angle, q \rightarrow h_1 h_2$$

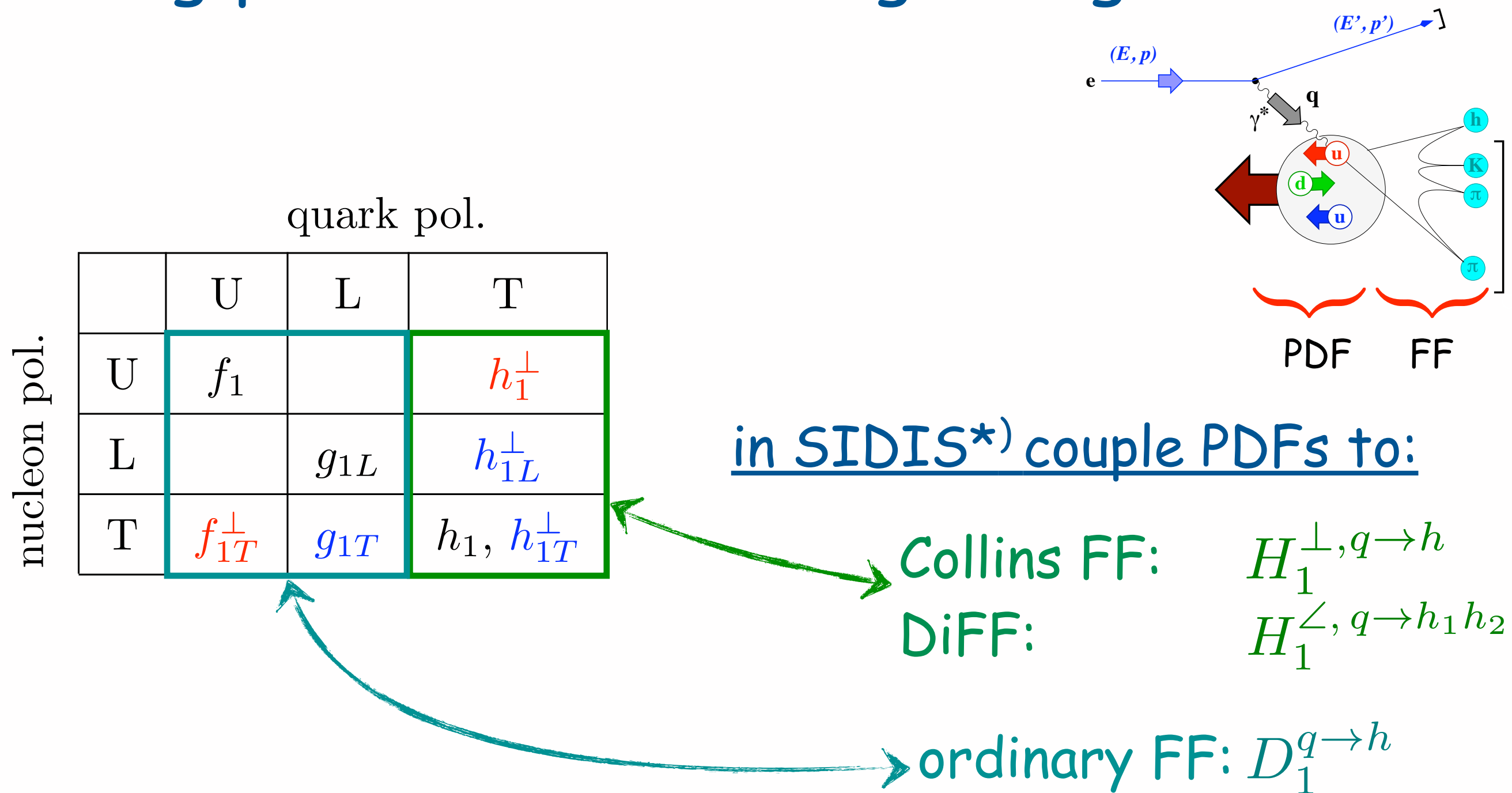
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Probing parton dist's through fragmentation

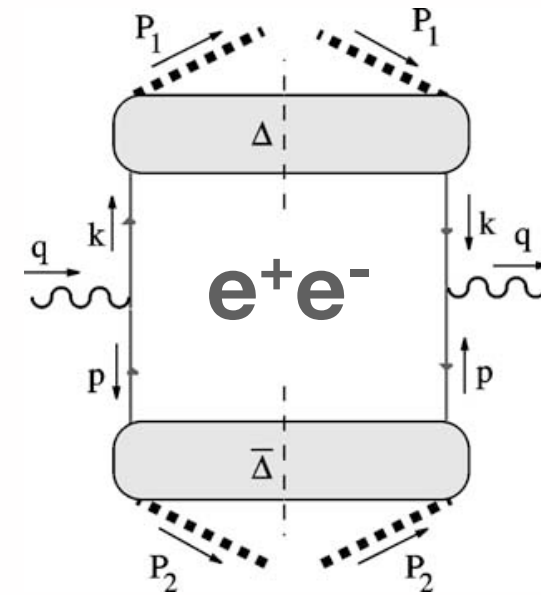


⇒ FFs act as quark flavor-tagger and polarimeter

*) semi-inclusive DIS with unpolarized final state

fragmentation in e^+e^- annihilation

- single-inclusive hadron production,
 $e^+e^- \rightarrow hX$
- D_1 fragmentation fctn.
- D_{1T}^\perp spontaneous transv. pol.

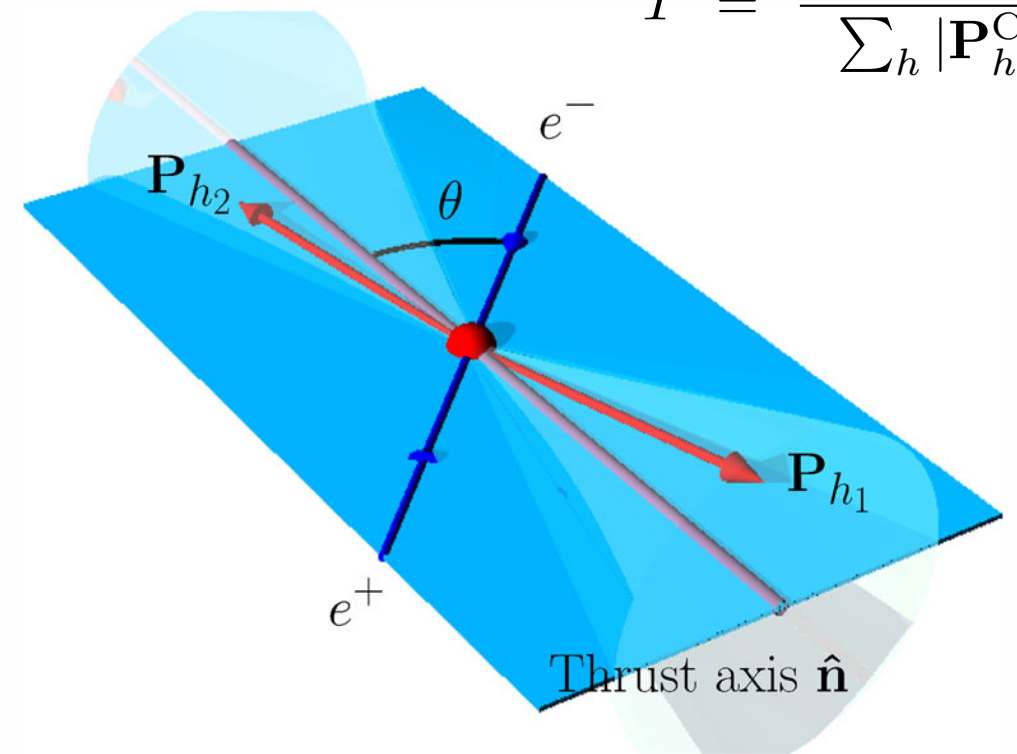


fragmentation in e^+e^- annihilation

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- D_1 fragmentation fctn.
- D_{1T^\perp} spontaneous transv. pol.
- inclusive “back-to-back” hadron pairs, $e^+e^- \rightarrow h_1h_2X$
- product of FFs
- flavor, transverse-momentum, and/or polarization tagging

Thrust (axis):

$$T \stackrel{\text{max}}{=} \frac{\sum_h |\mathbf{P}_h^{\text{CMS}} \cdot \hat{\mathbf{n}}|}{\sum_h |\mathbf{P}_h^{\text{CMS}}|}$$

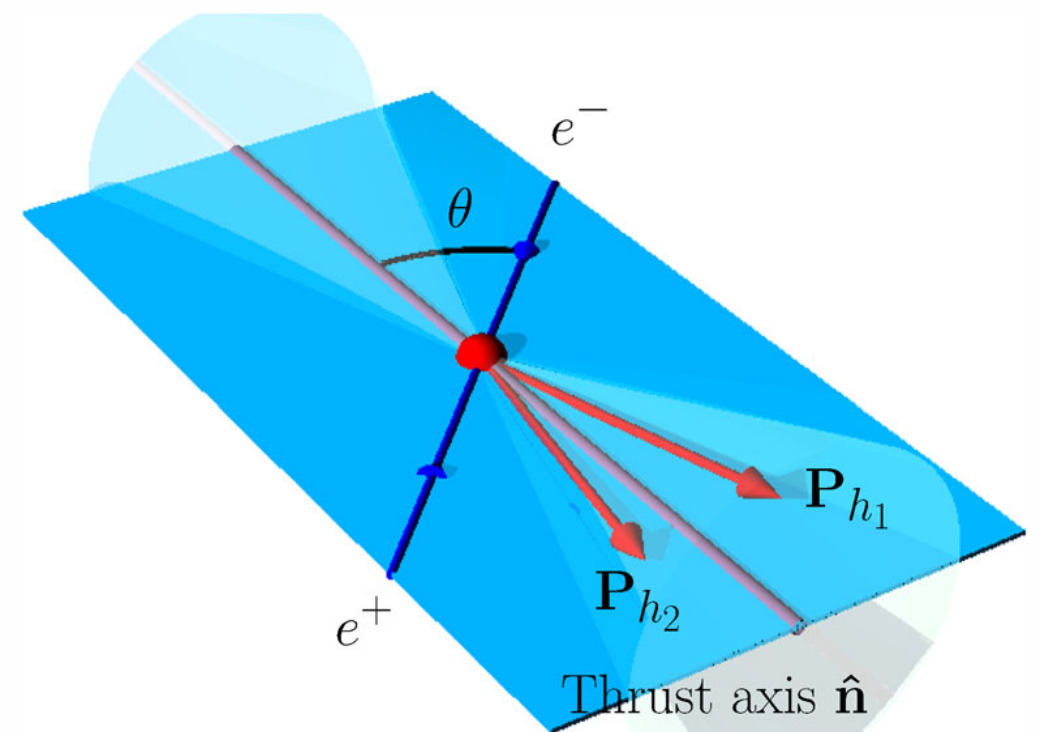
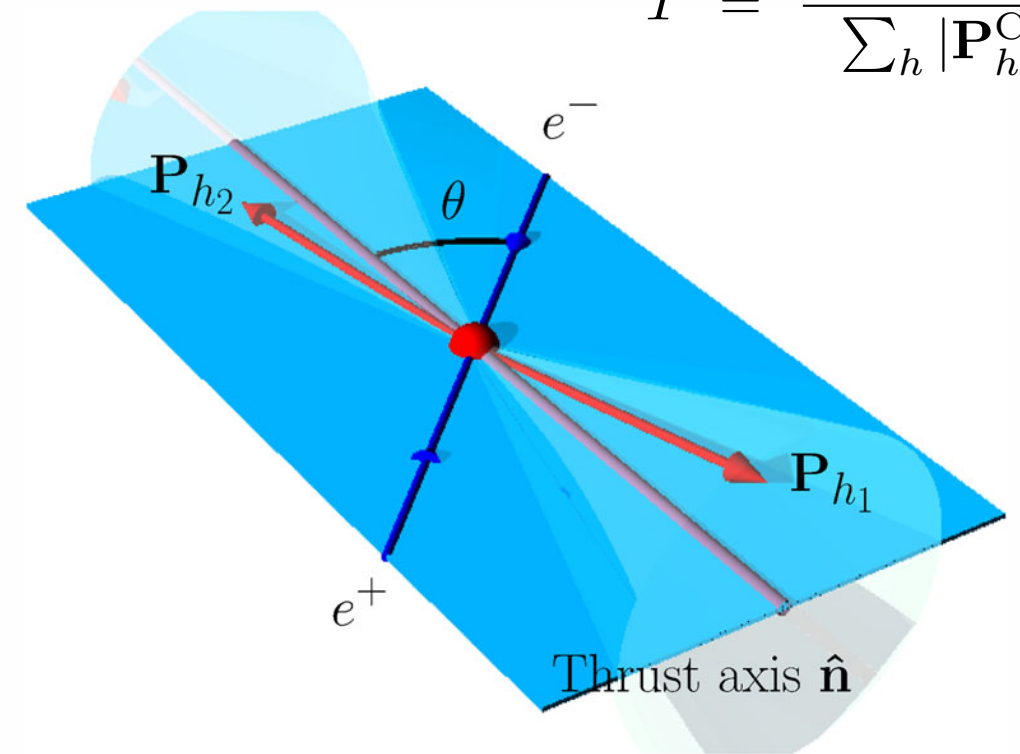


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- inclusive “back-to-back” hadron pairs, $e^+e^- \rightarrow h_1h_2X$
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- inclusive same-hemisphere hadron pairs, $e^+e^- \rightarrow h_1h_2X$
- dihadron fragmentation

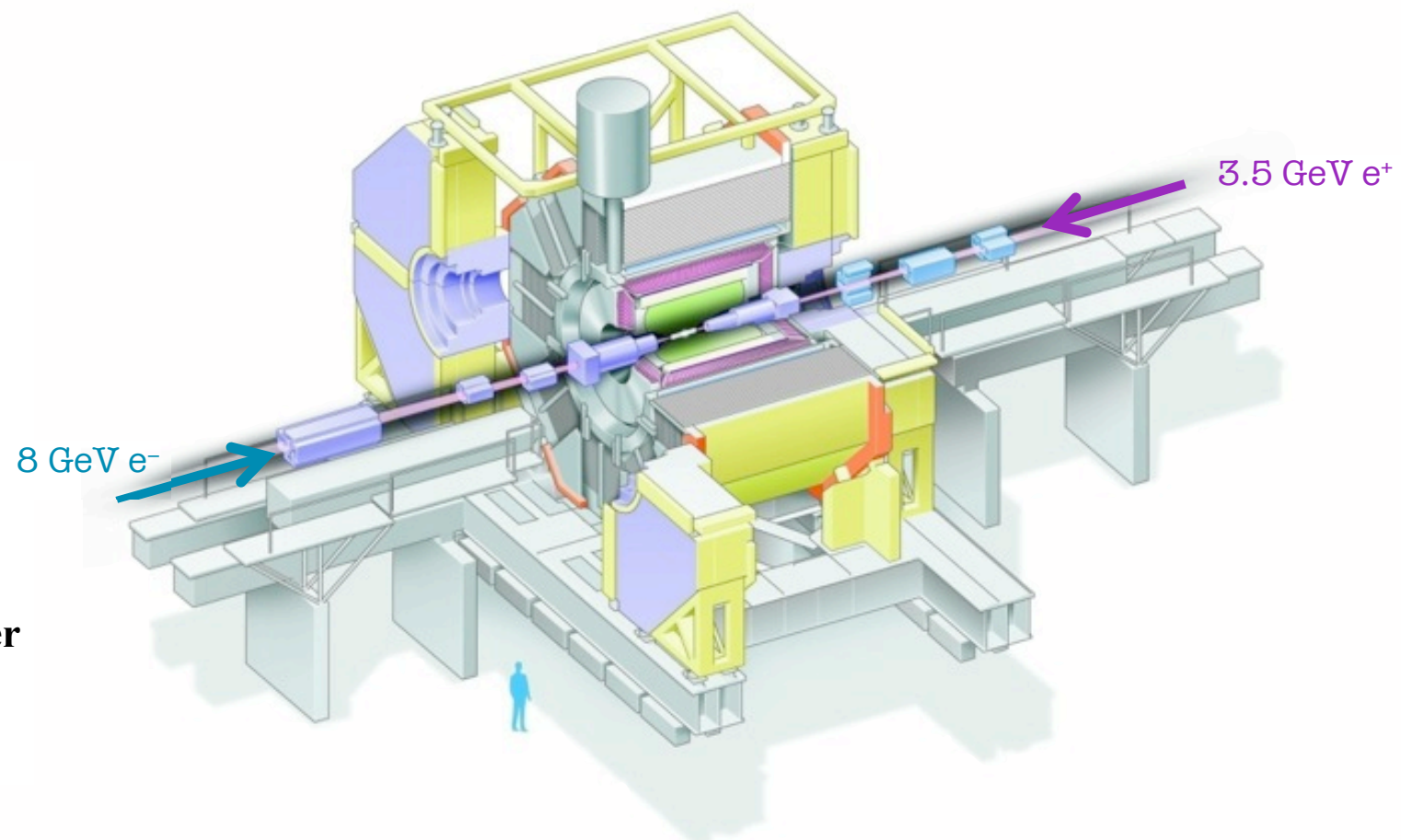
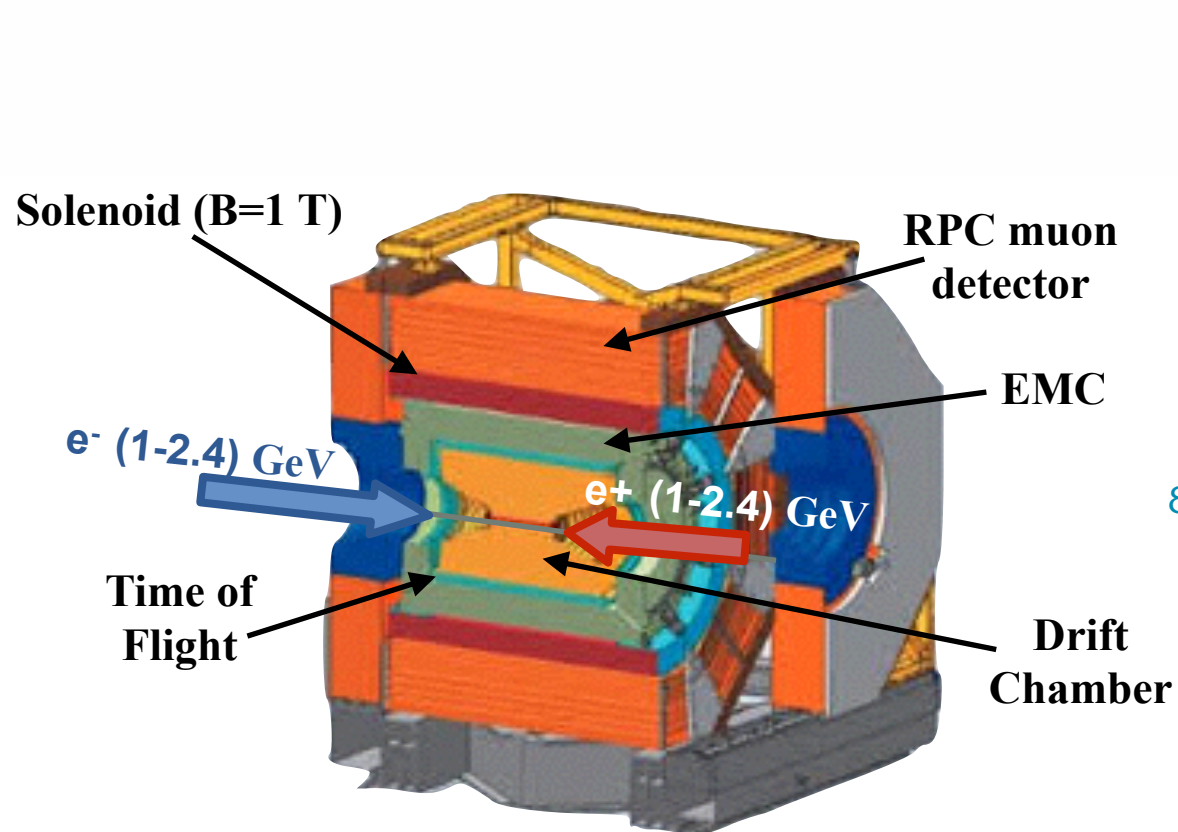
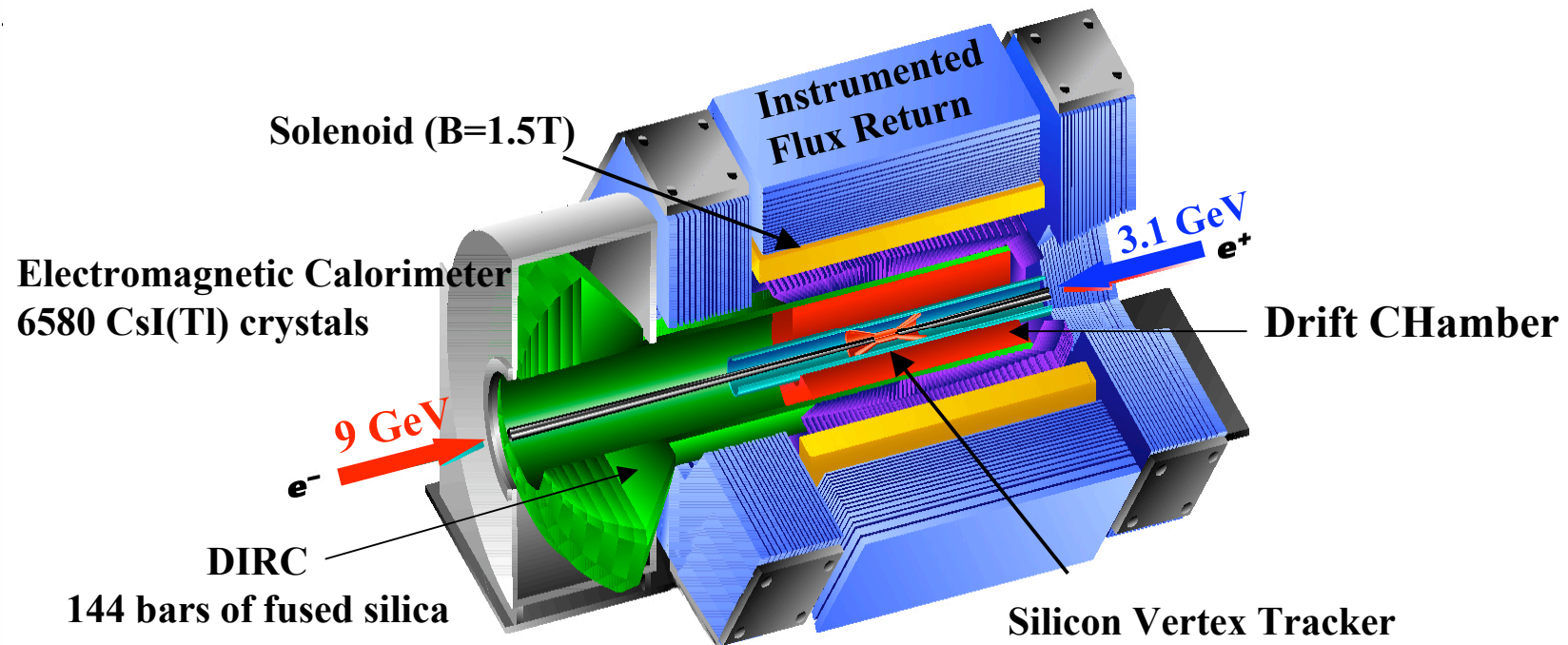
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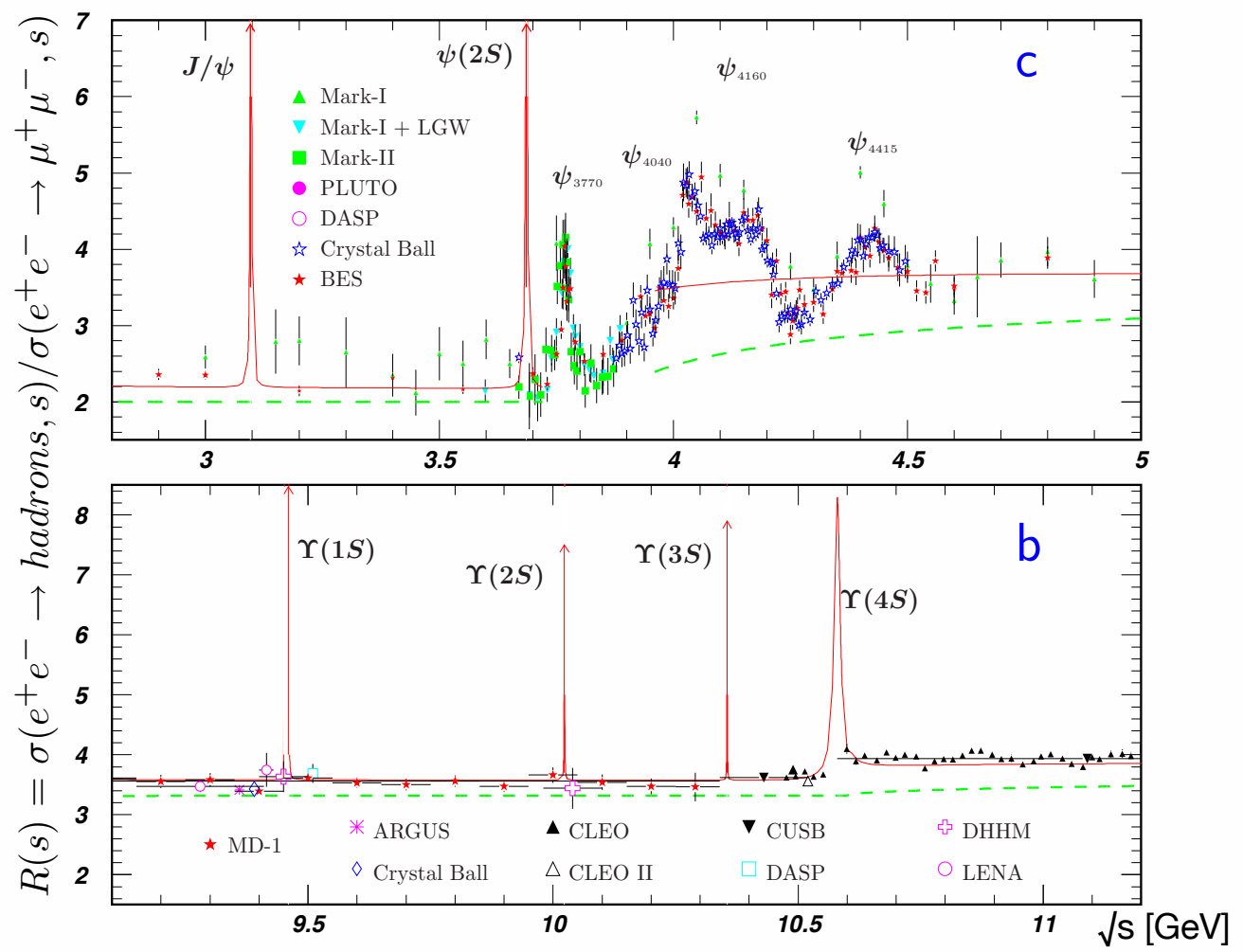
e^+e^- annihilation at BaBar, Belle, and BESIII

- BaBar/Belle: asymmetric beam-energy e^+e^- collider near/at $\Upsilon(4S)$ resonance (10.58 GeV)
- BESIII: symmetric collider with $E_e=1\ldots2.4$ GeV



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- BESIII: symmetric collider with $E_e=1...2.4$ GeV
- integrated luminosities:



	$\Upsilon(4S)$ on resonance	$\Upsilon(4S)$ off resonance	other
BaBar	424.2 fb ⁻¹	43.9 fb ⁻¹	
Belle	(140+571) fb ⁻¹	(15.6+73.8) fb ⁻¹	~180 fb ⁻¹ @ $\Upsilon(nS)$
BESIII			~62 pb ⁻¹ @3.65 GeV *)

from hadron yields to cross sections

- hadron yields undergo series of corrections
 - particle (mis)identification [e.g., not every identified pion was a pion]
 - smearing unfolding [e.g., measured and true momentum might differ]
 - non- $q\bar{q}$ processes [e.g., two-photon processes, $\Upsilon \rightarrow BB, \dots$]
 - “ 4π ” correction [selection criteria and limited geometric acceptance]
 - QED radiation [initial-state radiation (ISR)]
 - optional: weak-decay removal (e.g., “prompt fragmentation”)

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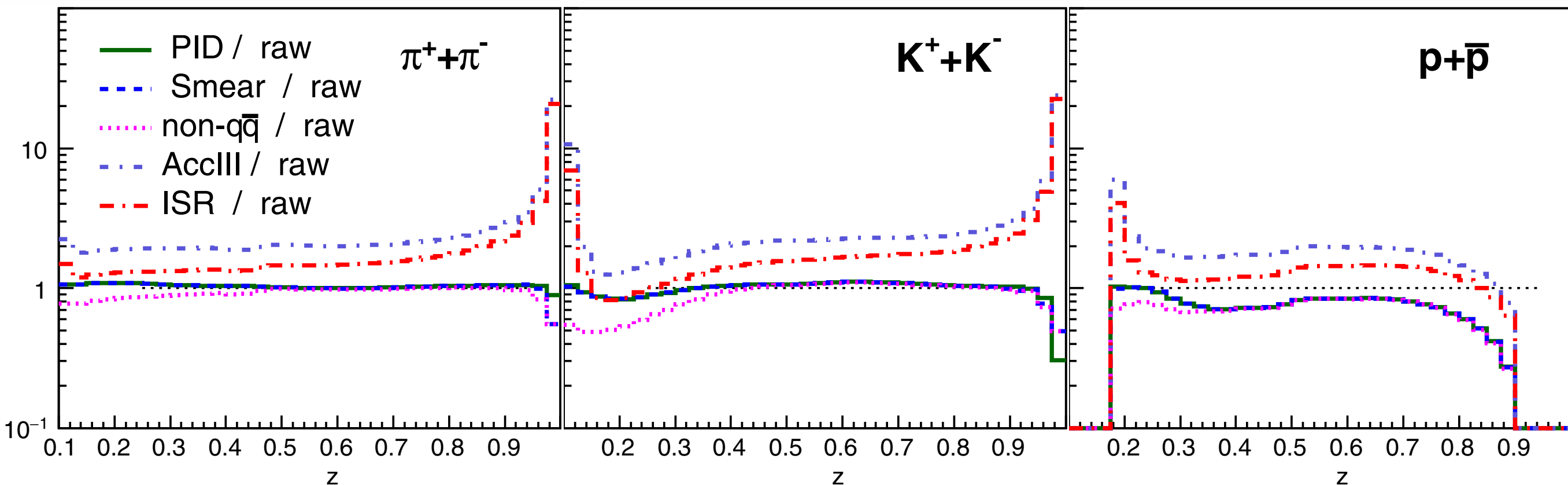
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- partially different approaches in different experiments/analyses

from hadron yields to cross sections

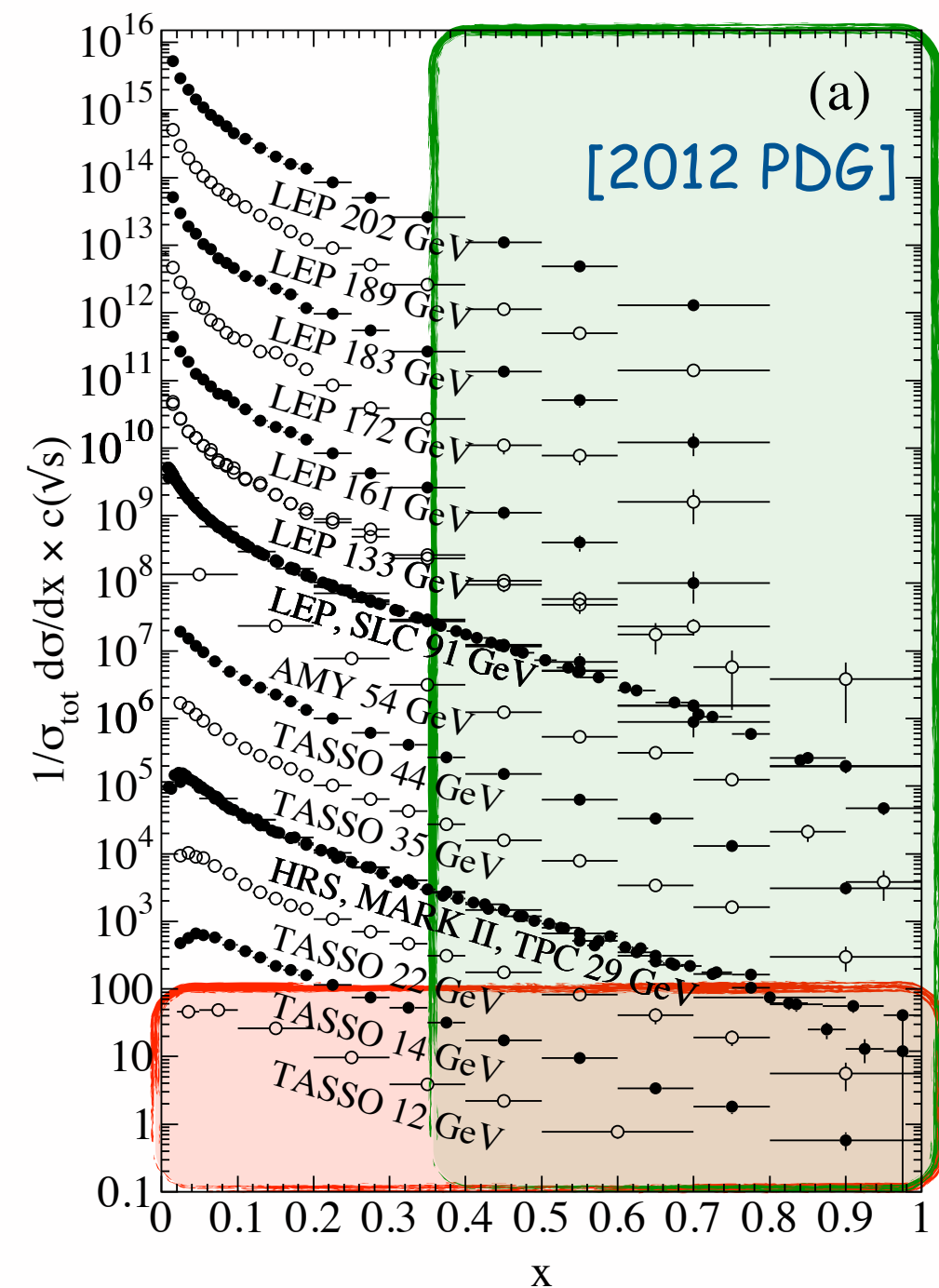
- **example:** single-hadron inclusive cross sections
- cumulative effect of correction steps



- largest effect for mesons from acceptance and ISR correction
- larger PID correction for protons than for mesons

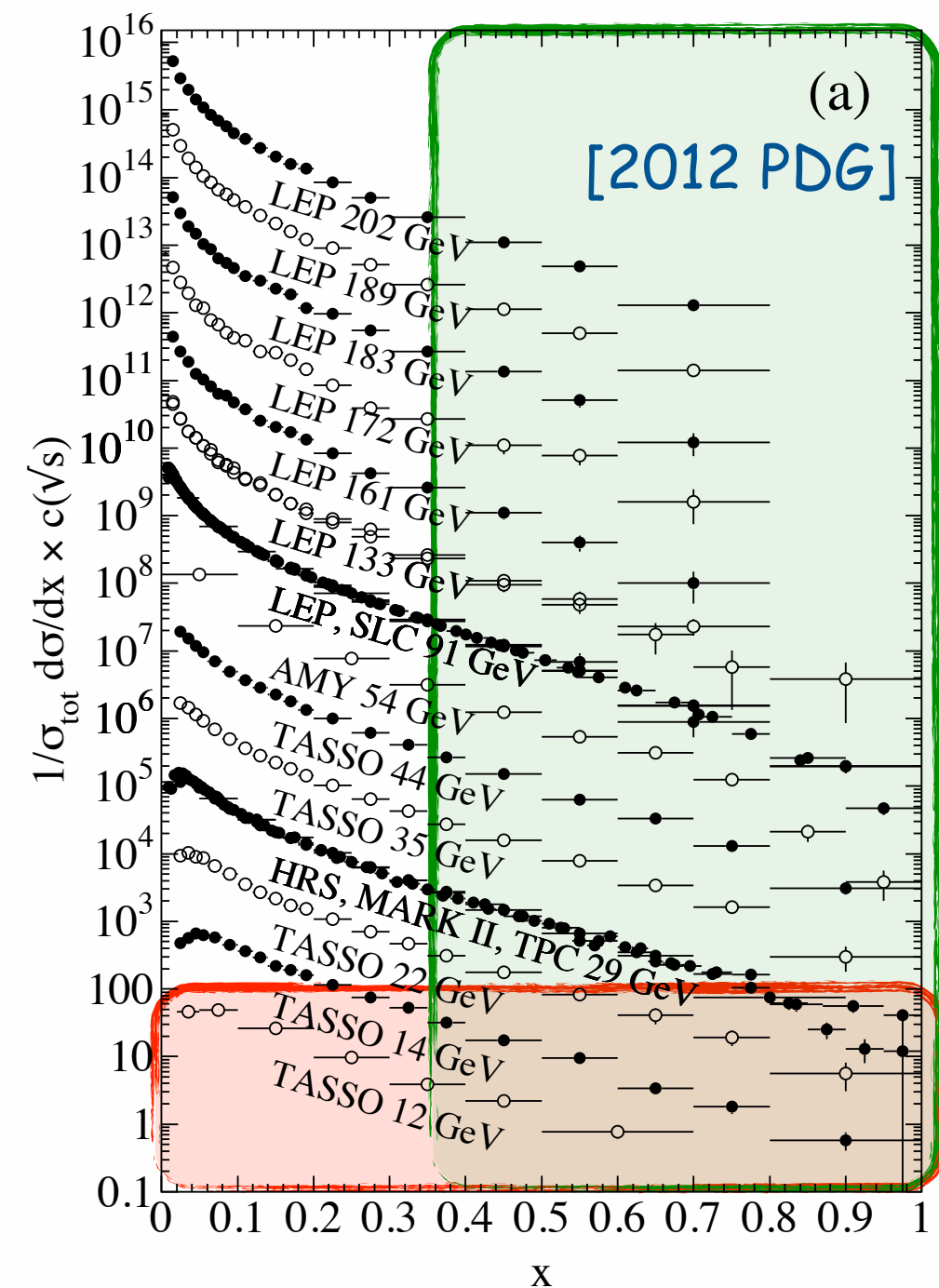
single-hadron production

- before 2013: lack of precision data at (moderately) high z and at low \sqrt{s}
- limits analysis of evolution and gluon fragmentation
- limited information in kinematic region often used in semi-inclusive DIS



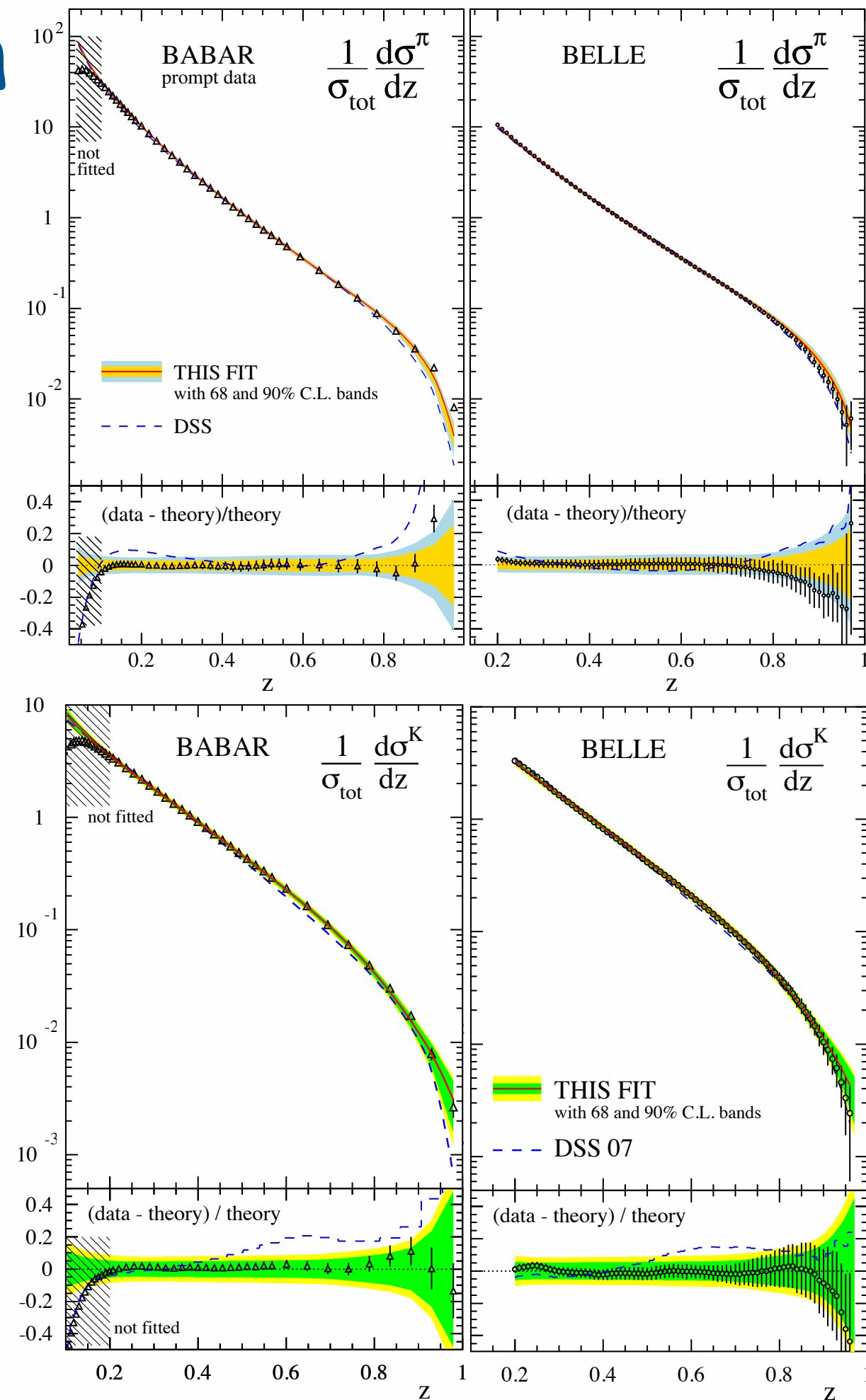
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- by now also results from BaBar and Belle:
 - BaBar Collaboration, Phys. Rev. D88 (2013) 032011: π^\pm , K^\pm , $p+p$
 - Belle Collaboration, Phys. Rev. Lett. 111 (2013) 062002: π^\pm , K^\pm
 - Belle Collaboration, Phys. Rev. D92 (2015) 092007: π^\pm , K^\pm , $p+p$



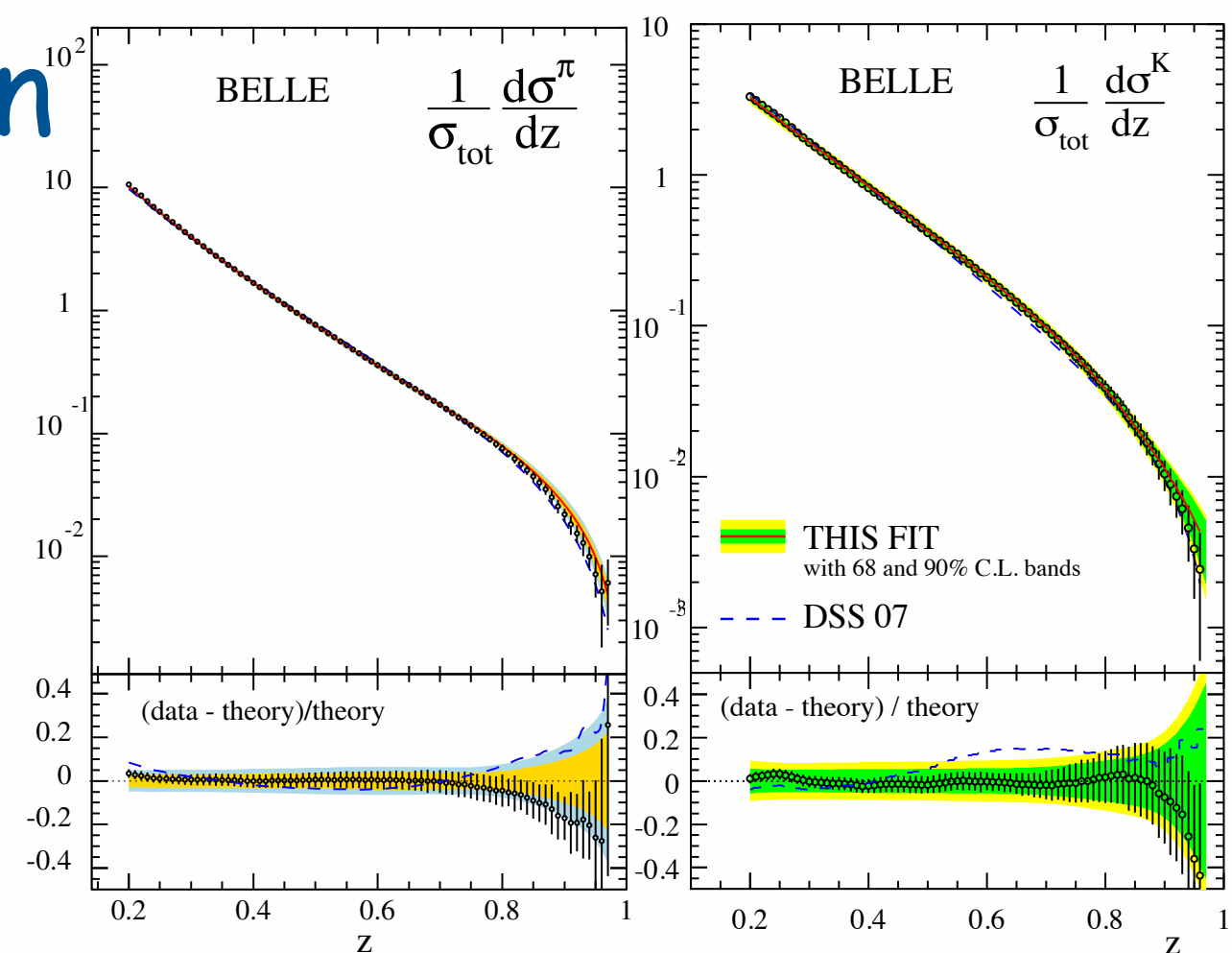
single-hadron production

- very precise data for charged pions and kaons
- Belle data available up to very large z ($z < 0.98$)
- included in recent DEHSS fits
- slight tension at low- z for BaBar and high- z for Belle



single-hadron production

- very precise data for charged pions and kaons
- Belle data available up to very large z ($z < 0.98$)
- included in recent DEHSS fits [e.g. PRD 91, 014035 (2015)]
- Belle radiative corrections “undone” in FF fits

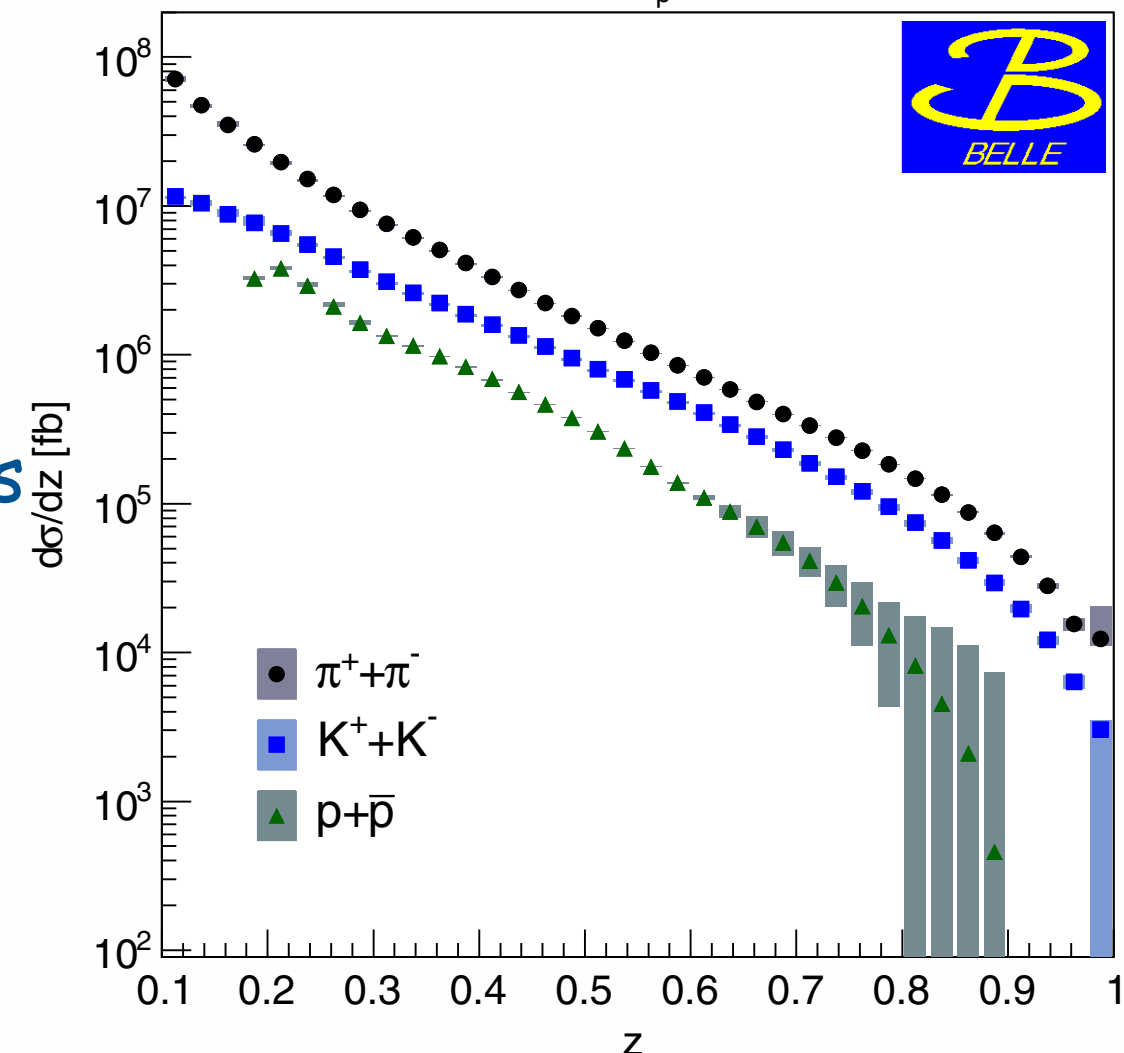
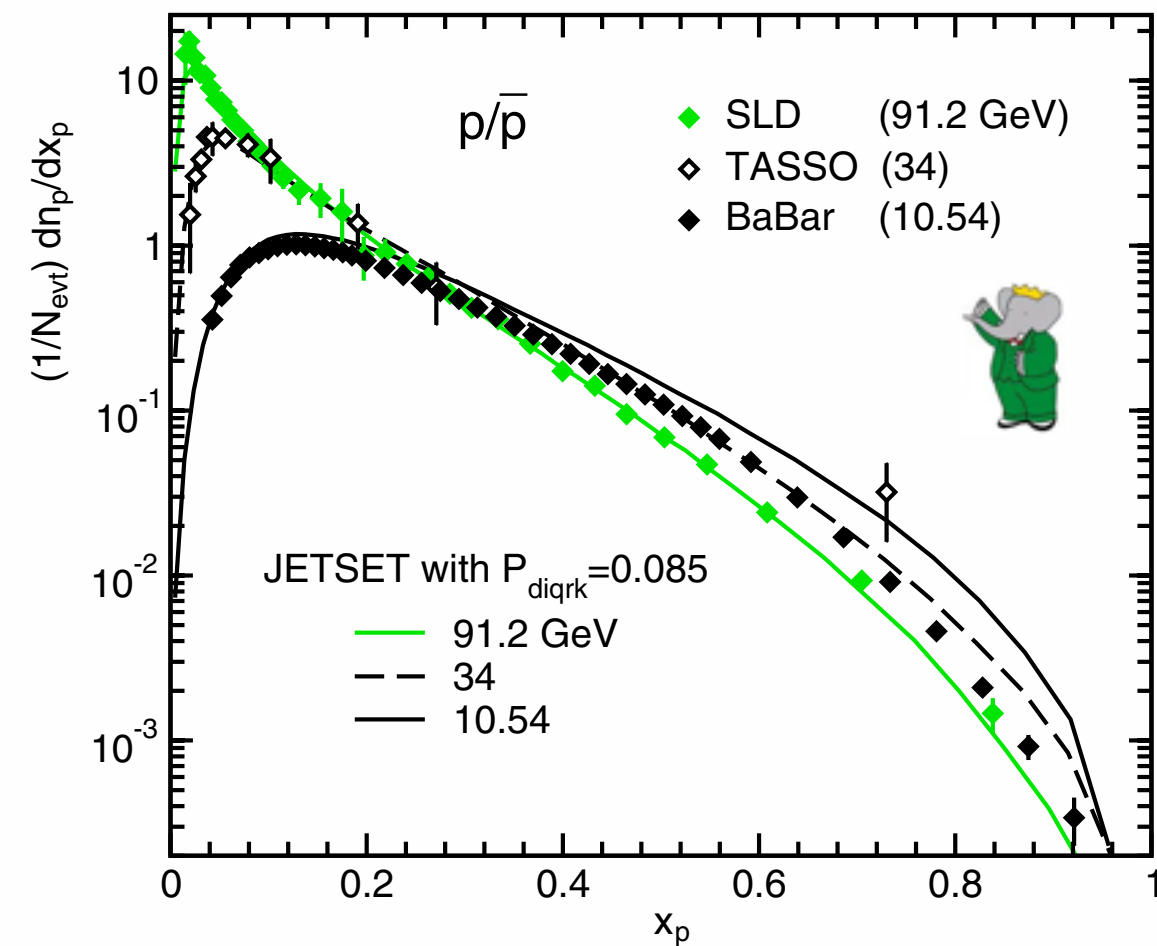


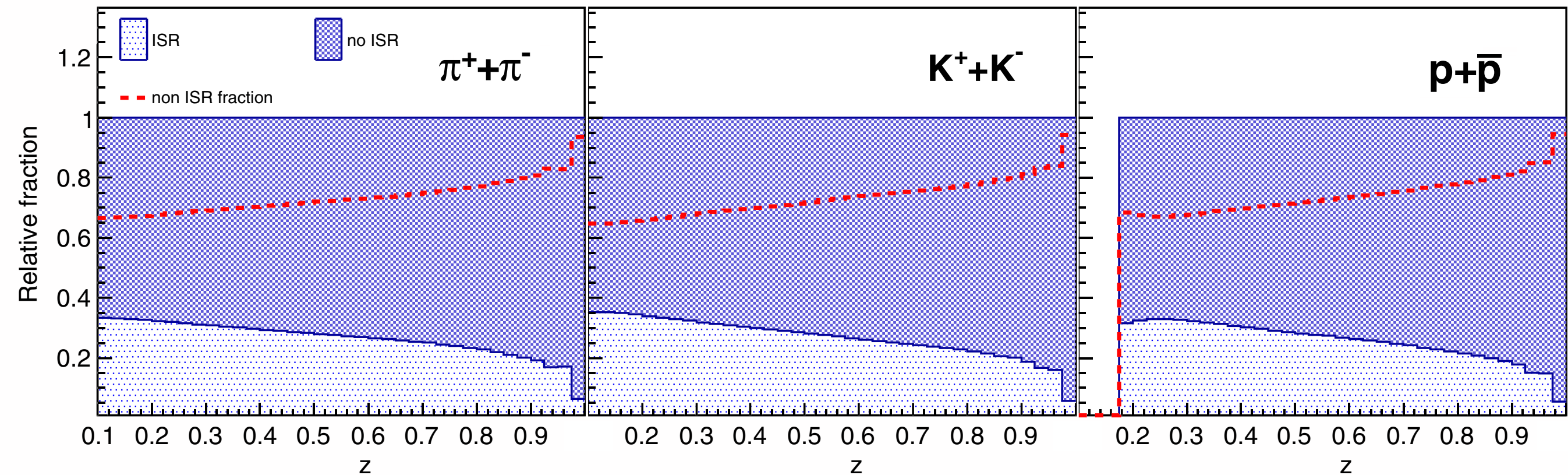
[EPJC 77 (2017) 516, NNFF1.0]

In the case of the BELLE experiment we multiply all data points by a factor $1/c$, with $c = 0.65$ for charged pions and kaons [69] and with c a function of z for protons/antiprotons [53]. This correction is required in order to treat the BELLE data consistently with all the other SIA measurements included in NNFF1.0. The reason is that a kinematic cut on radiative photon events was applied to the BELLE data sample in the original analysis instead of unfolding the radiative QED effects. Specifically, the energy scales

single-hadron production

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- included in recent DEHSS fits [e.g. PRD 91, 014035 (2015)]
- Belle radiative corrections “undone” in FF fits
- Belle data for protons and anti-protons
 - not (yet) included in DEHSS, but in NNFF 1.0 [EPJC 77 (2017) 516]
 - similar z dependence as pions
 - about $\sim 1/5$ of pion cross sections



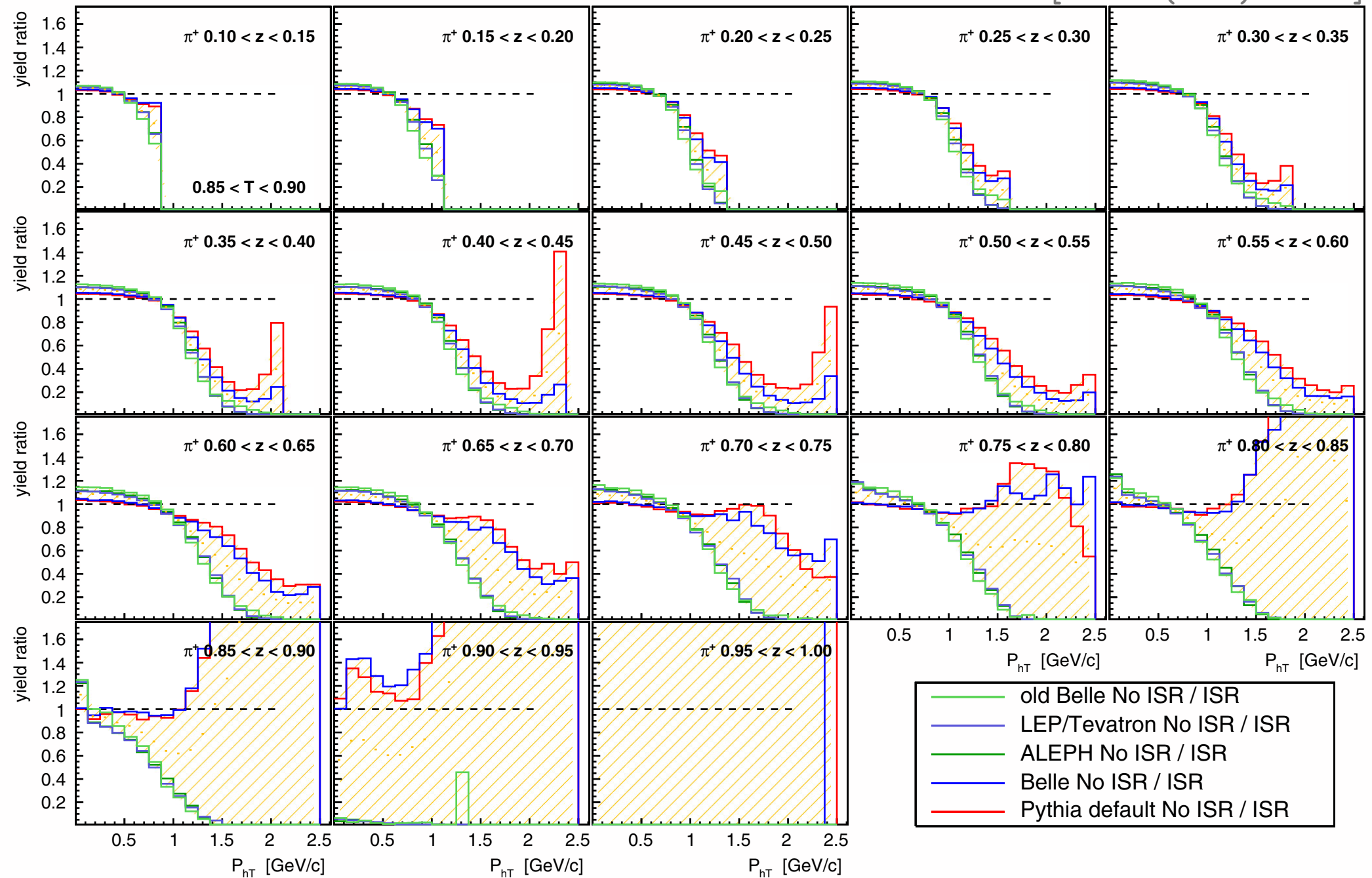


- relative fractions of hadrons as a function of z originating from ISR or non-ISR events (\equiv energy loss less than 0.5%)
- large non-ISR fraction at large z , as otherwise not kinematically reachable (remember $z = E_h / 0.5\sqrt{s}_{\text{nominal}}$)
- keep only fraction of events w/o initial-state photon emission
→ strictly speaking not single-inclusive annihilation

ISR corrections - PRD99 (2019) 112006



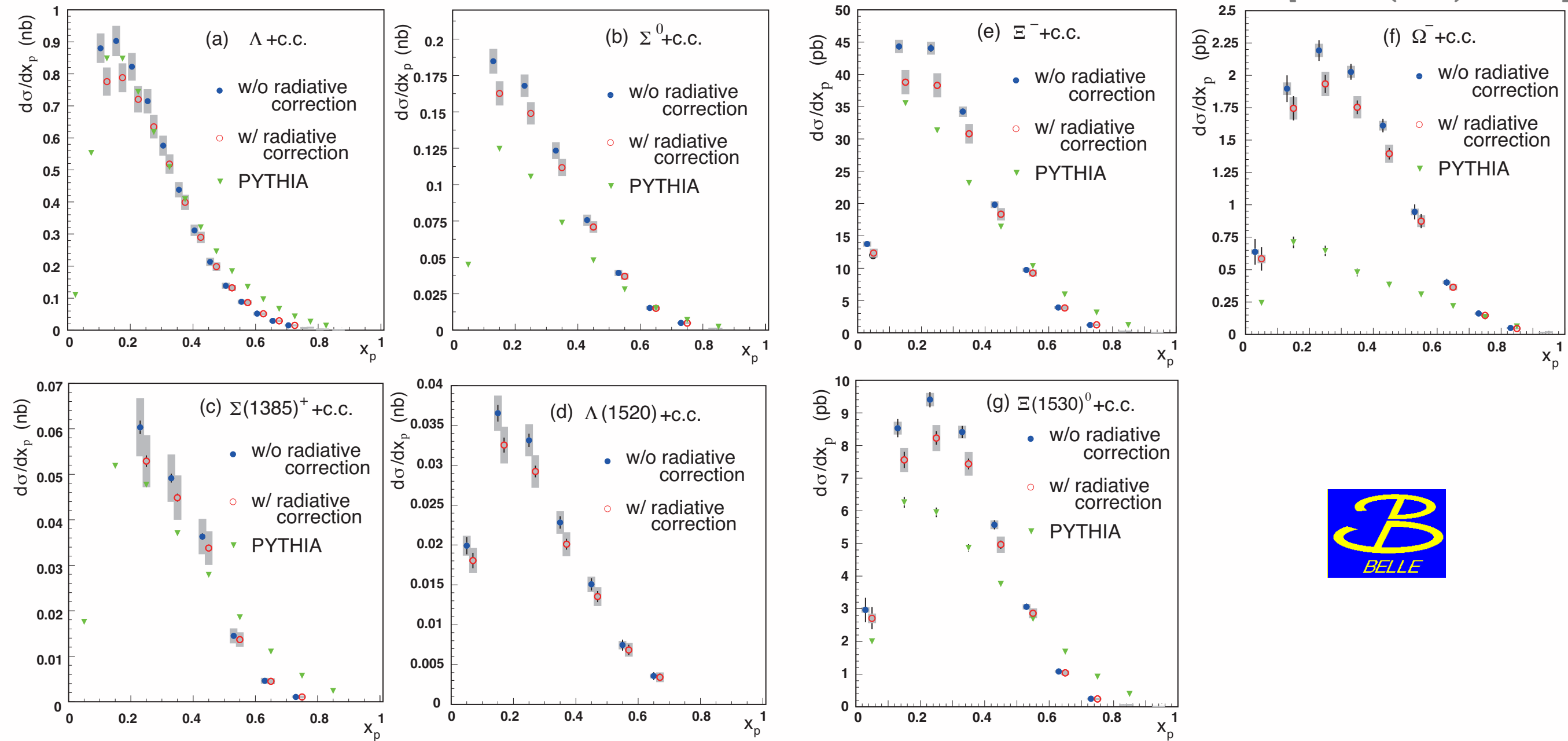
[PRD 99 (2019) 112006]



- non-ISR / ISR fractions based on PYTHIA switch MSTP(11)
[e.g., for transverse-momentum dependence of incl. hadron production]
- PYTHIA model dependence; absorbed in systematics by variation of tunes

inclusive hyperon production

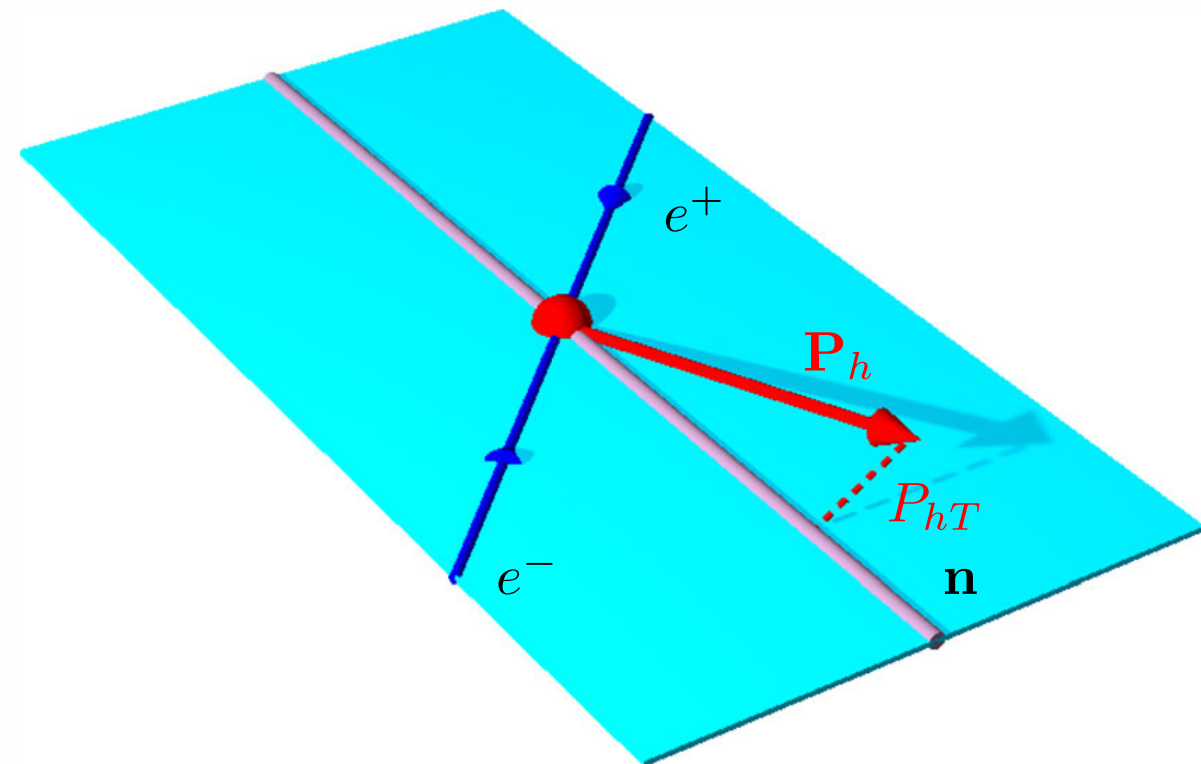
[PRD 97 (2018) 072005]



- Λ production reasonably well described by Pythia
- less satisfactory for heavier hyperons
- fails to describe Ω^- production

inclusive hadrons - transverse momentum

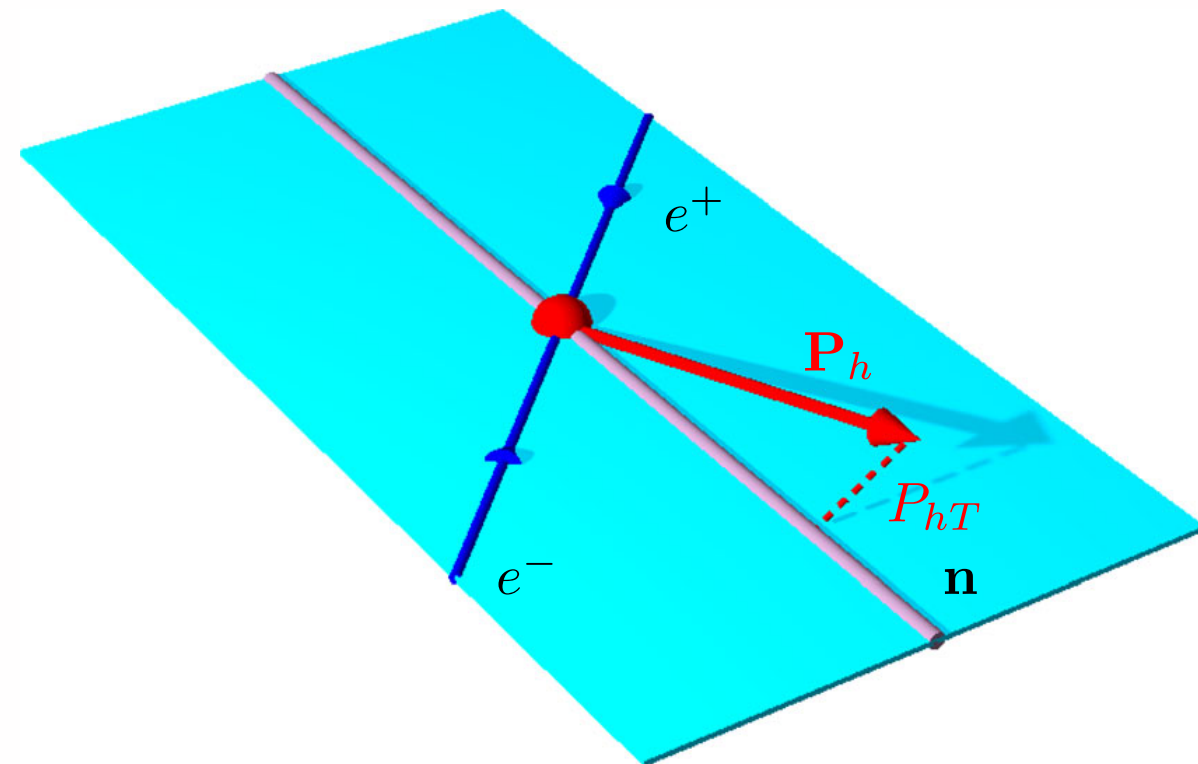
- quasi-inclusive hadron production gives access to transverse momentum in fragmentation
- transverse momentum measured with respect to thrust axis \mathbf{n}
- involves sum over all final-state particles in event
- event selection and hadron distributions dependent on thrust value T required
 - low thrust \rightarrow more spherical
 - high thrust \rightarrow highly collimated



$$T \stackrel{\text{max}}{=} \frac{\sum_h |\mathbf{P}_h^{\text{CMS}} \cdot \hat{\mathbf{n}}|}{\sum_h |\mathbf{P}_h^{\text{CMS}}|}$$

inclusive hadrons - transverse momentum

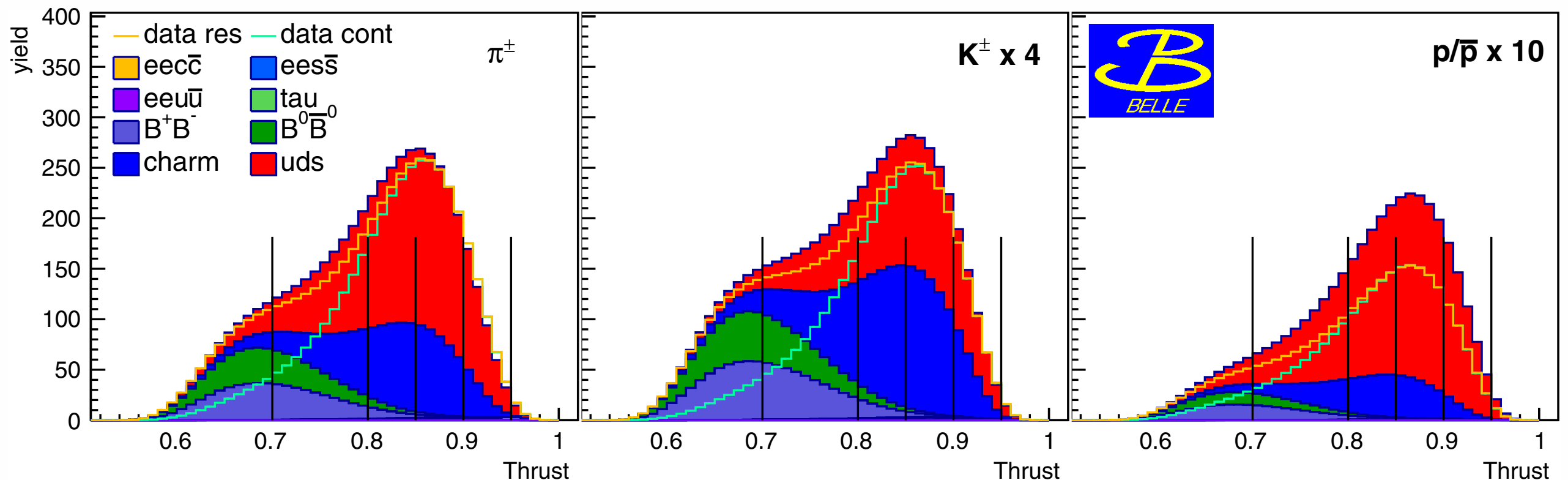
- quasi-inclusive hadron production gives access to transverse momentum in fragmentation
- transverse momentum measured with respect to thrust axis \mathbf{n}
- analysis performed differential in z and P_{hT} in various slices in T (18x20x6 bins)
- correction steps similar as for P_{hT} -integrated cross sections (using the "MSTP(11) radiative corrections")
- Gaussian fit to transverse-momentum distribution provided for all hadrons in (z,T) -bins



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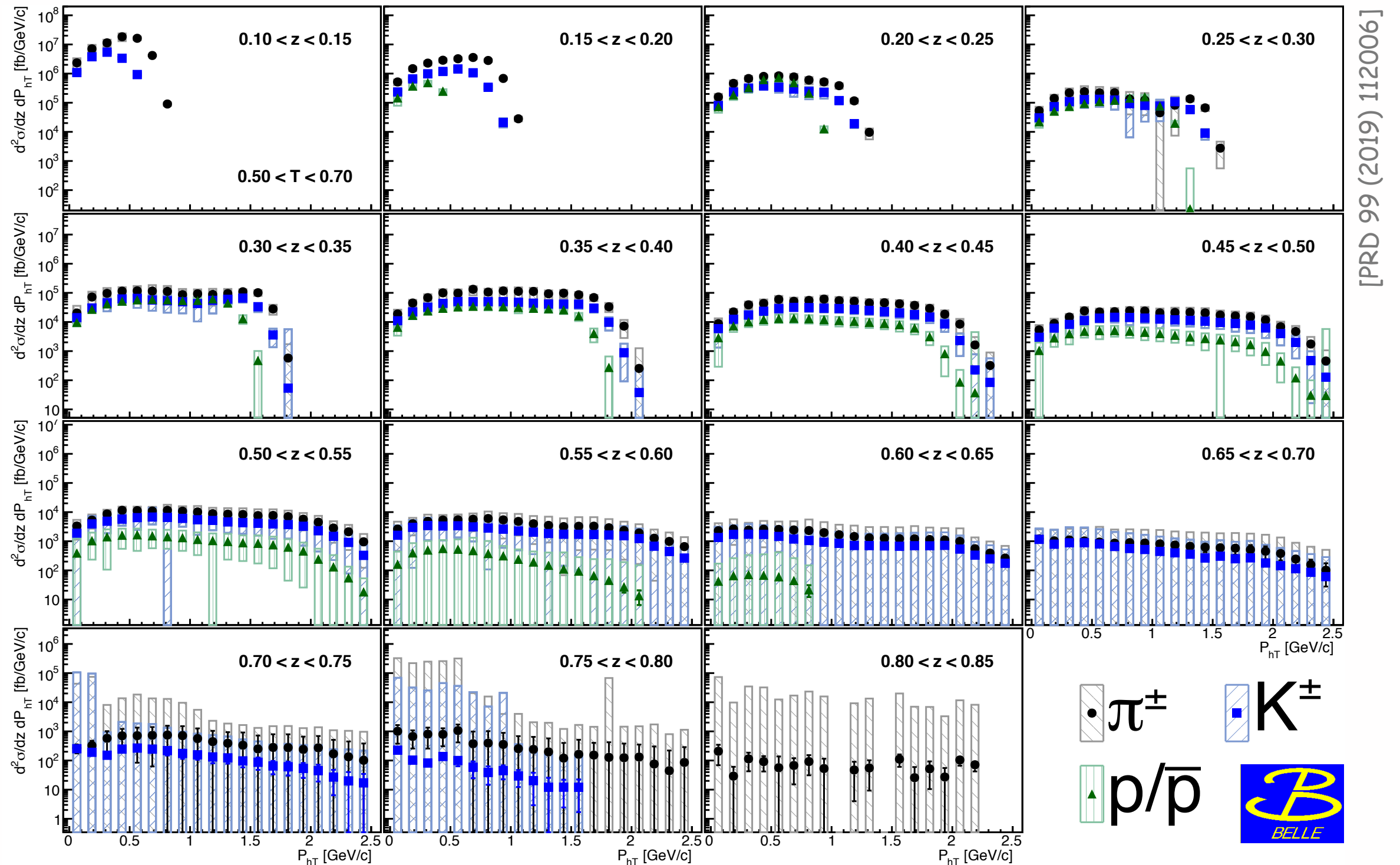
thrust distribution: process contributions

[Belle, PRD 99 (2019) 112006]



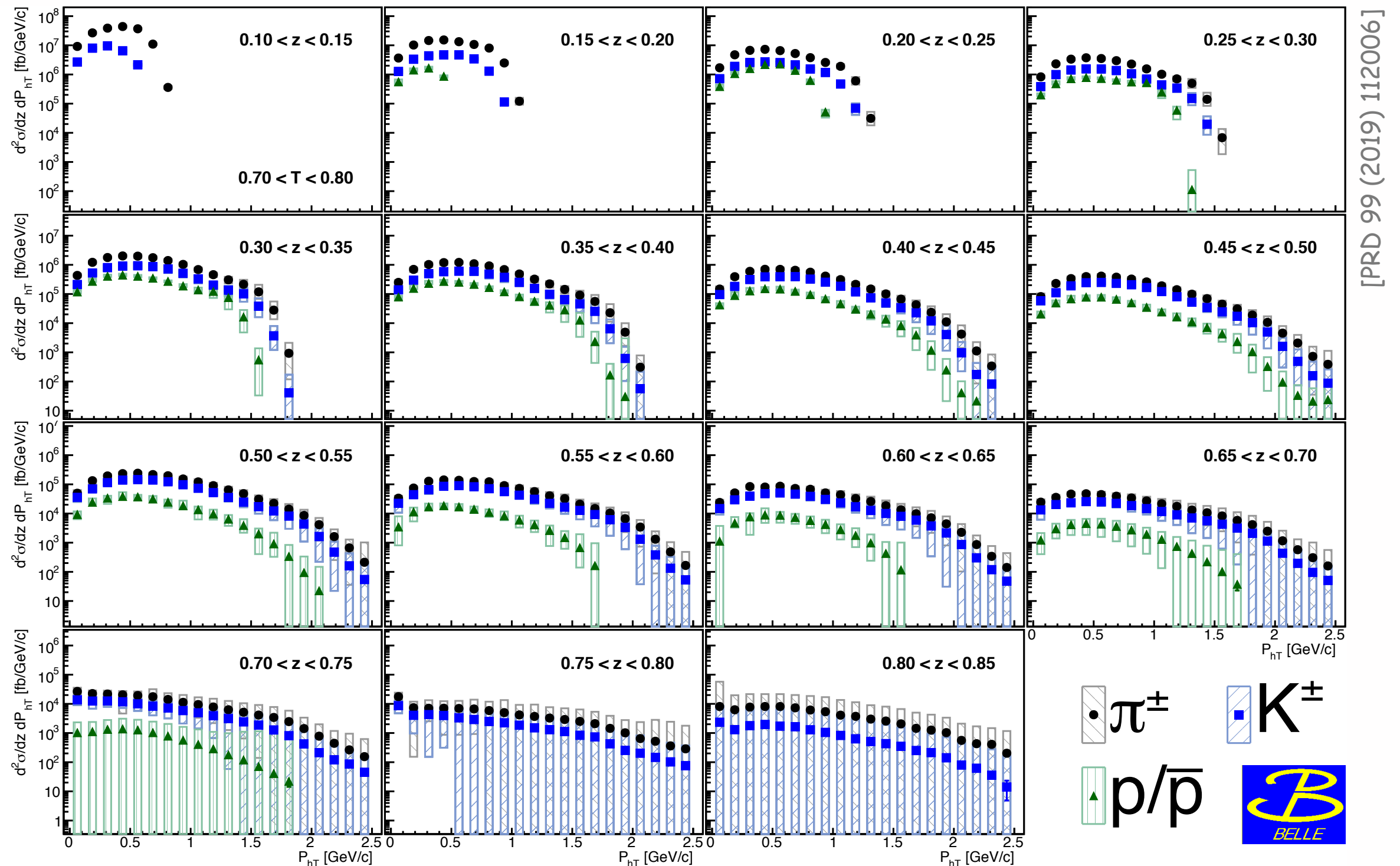
- large contribution from BB at lower thrust
- large thrust dominated by uds and charm fragmentation
(at very large T significant τ contribution for pions, not visible here)
- concentrate mainly on $0.85 < T < 0.9$ bin, though others available as well

transverse-momentum distributions



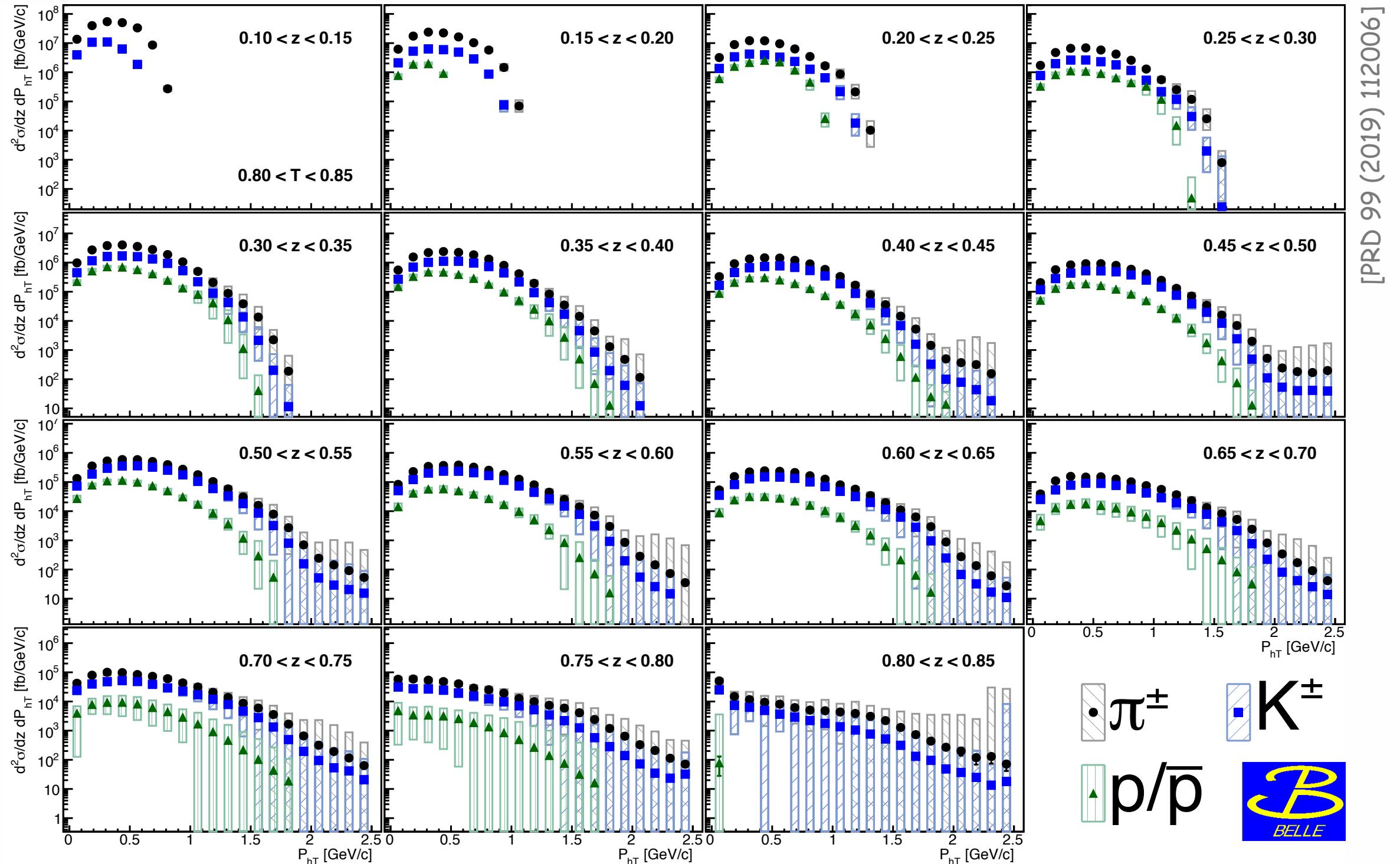
- lowest T bin -> rather spherical events
- transverse momenta almost uniformly distributed in medium-z bins
- faster drop for heavier hadrons

transverse-momentum distributions



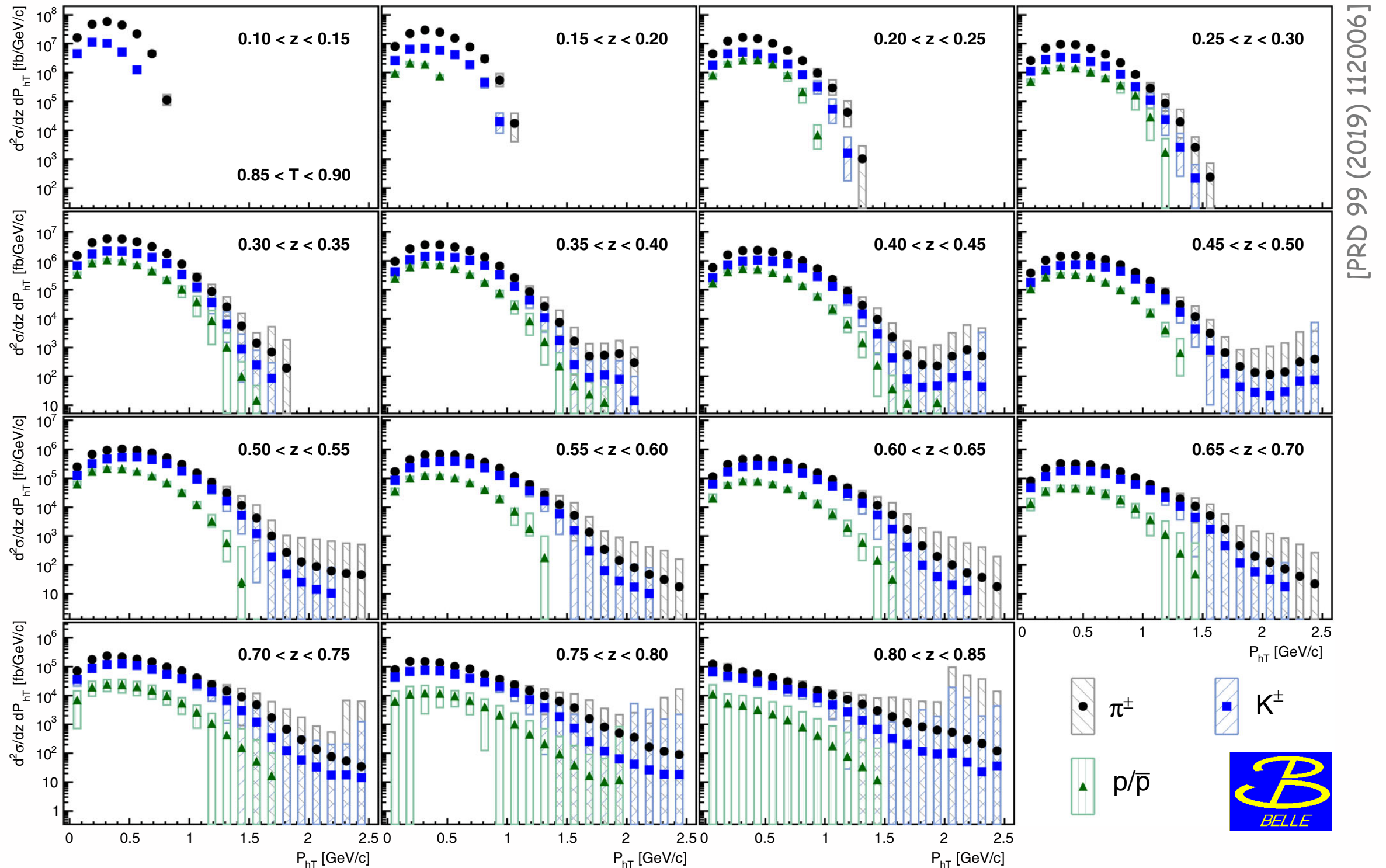
- 0.7 < T < 0.8 -> particles already more collimated
- transverse momenta already more Gaussian distributed
- large-z region with large uncertainties

transverse-momentum distributions



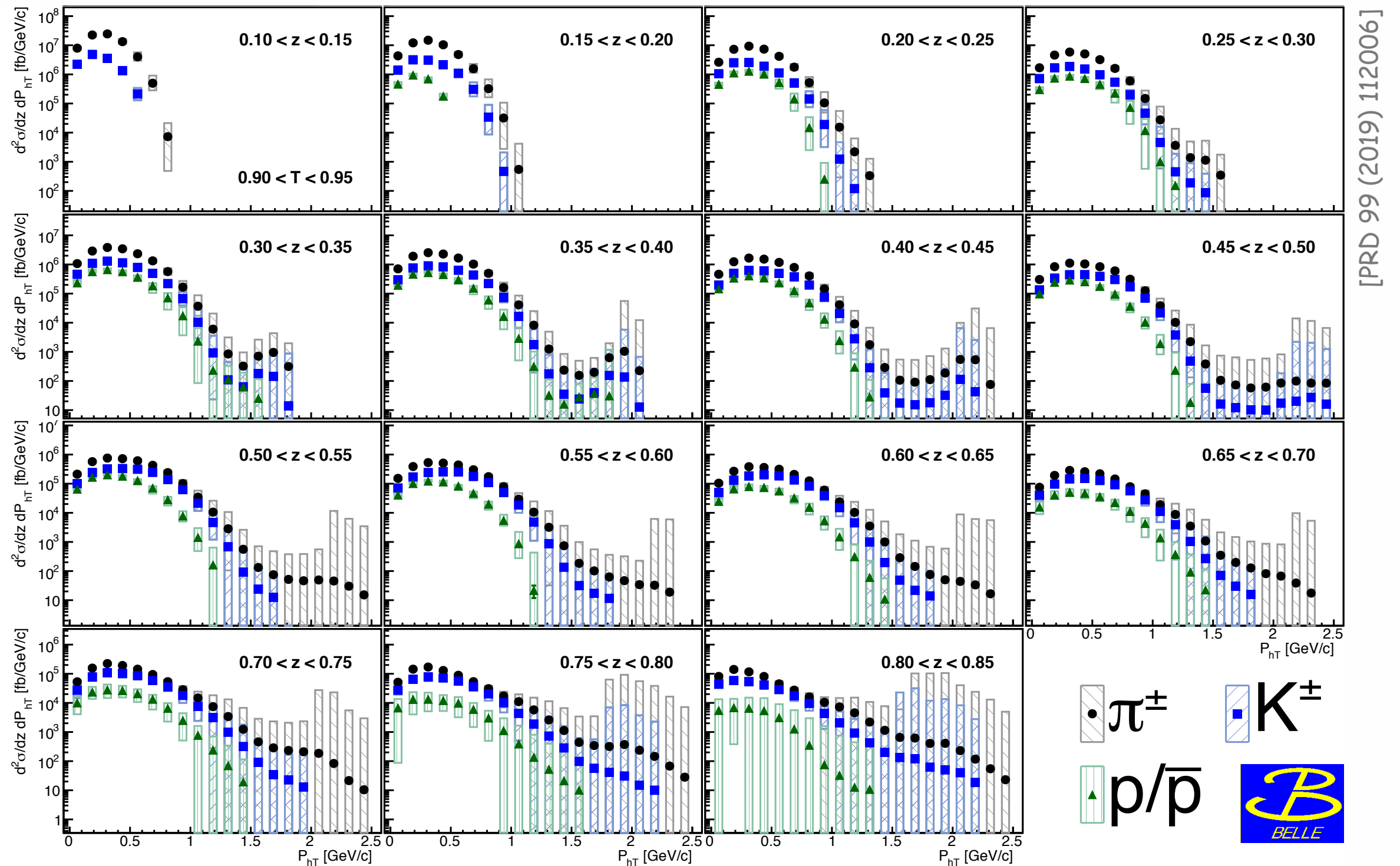
- 0.8 < T < 0.85
- transverse momenta mostly Gaussian distributed
- possible deviations for large- P_{hT} tails [but also larger uncertainties]

transverse-momentum distributions



- $0.85 < T < 0.9$
- transverse momenta mostly Gaussian distributed; widths narrowing
- possible deviations for large- P_{hT} tails [but also larger uncertainties]

transverse-momentum distributions

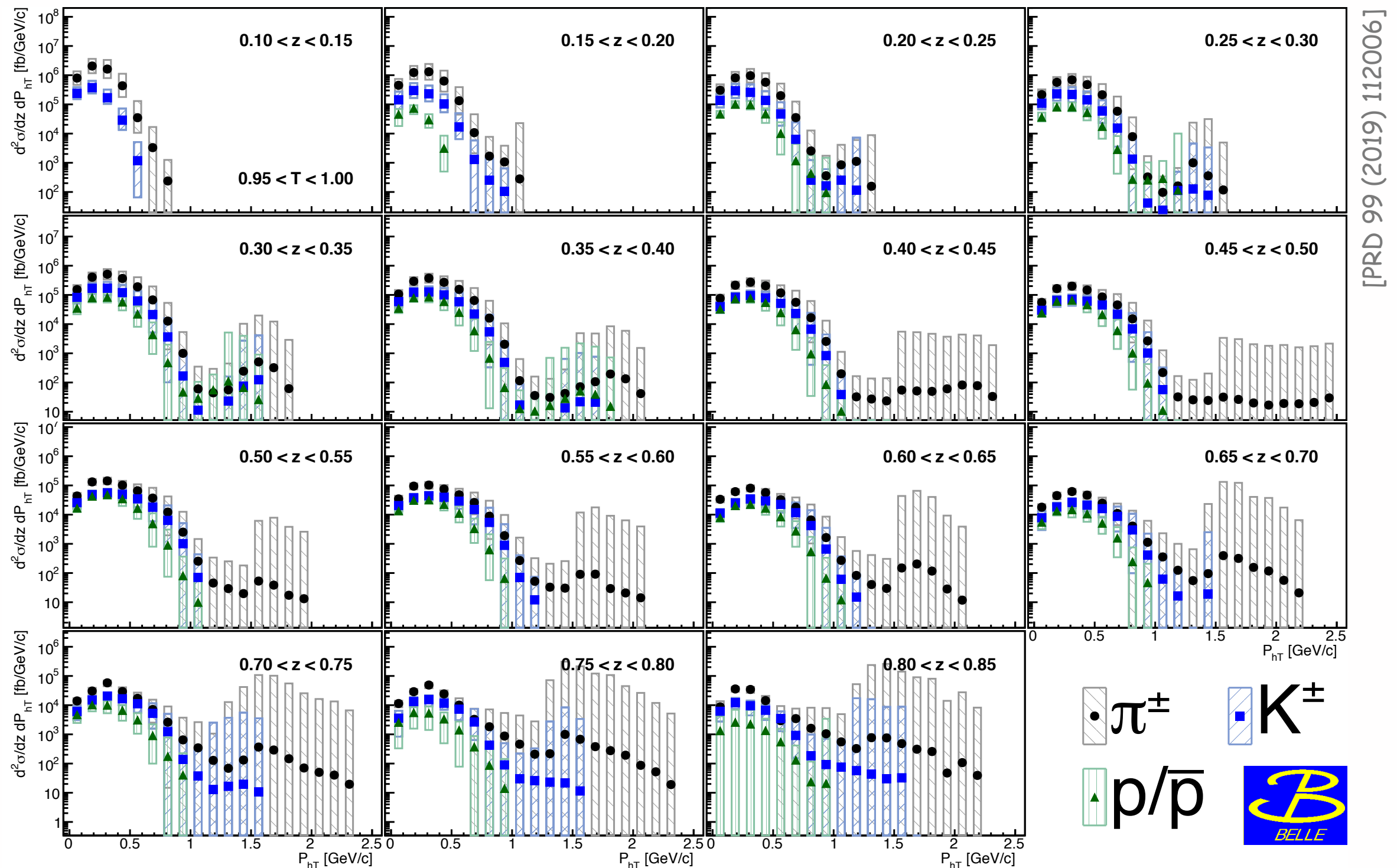


● $0.9 < T < 0.95$

● transverse momenta mostly Gaussian distributed; widths even narrower

● possible deviations for large- P_{hT} tails [but also larger uncertainties]

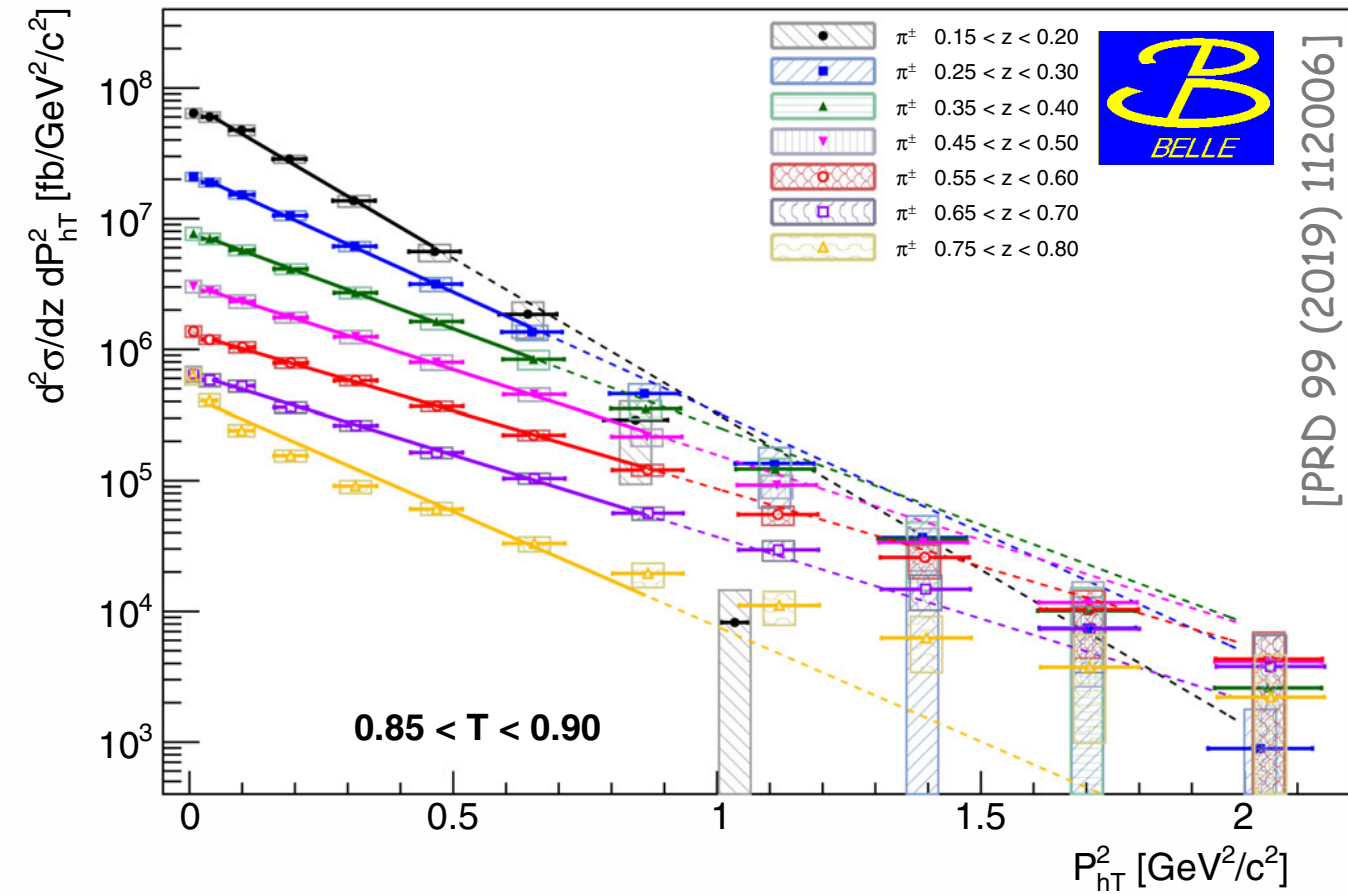
transverse-momentum distributions



- $0.95 < T < 1.0$
- transverse momenta mostly Gaussian distributed
- widths very narrow as particles very collimated

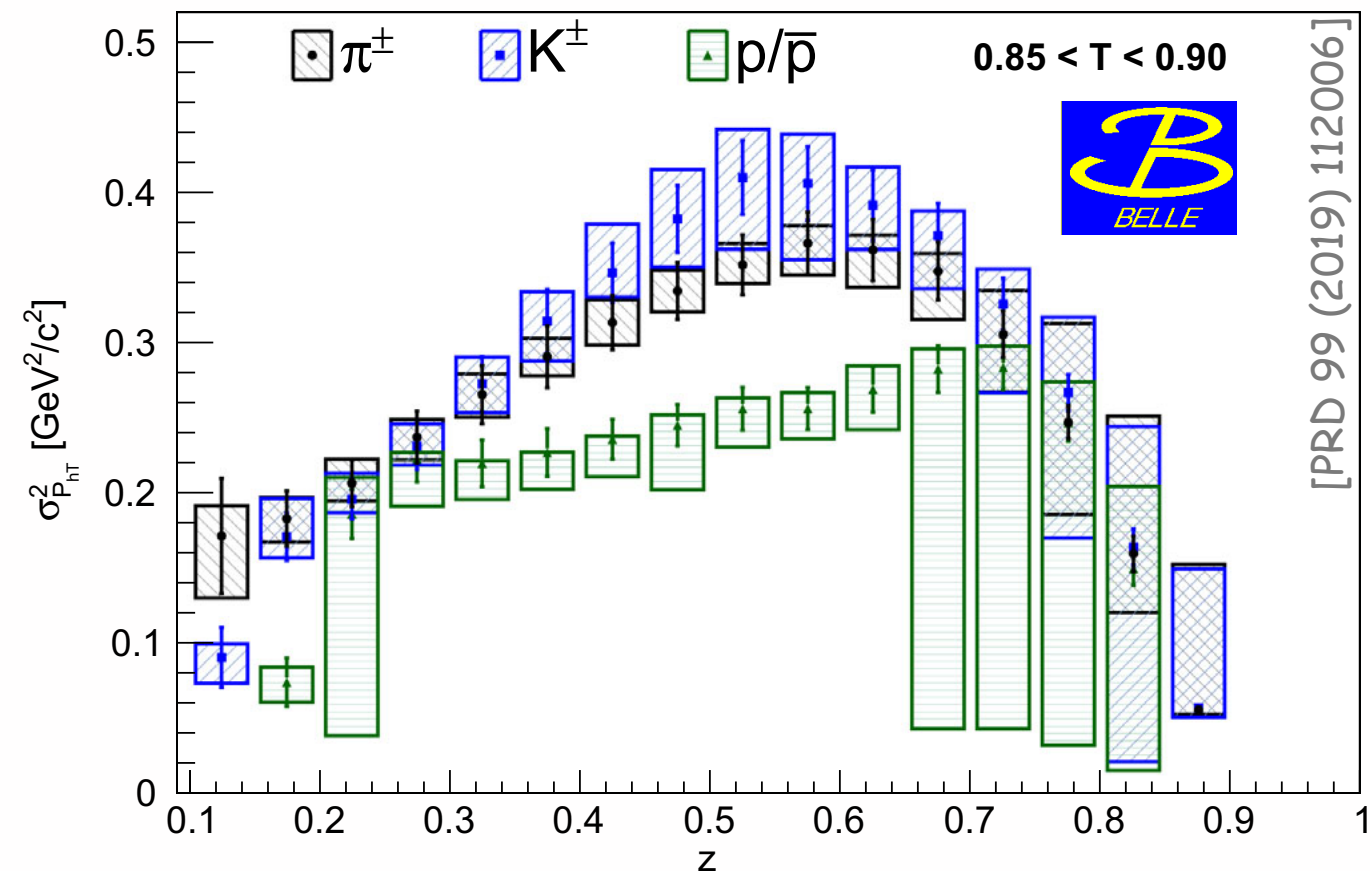
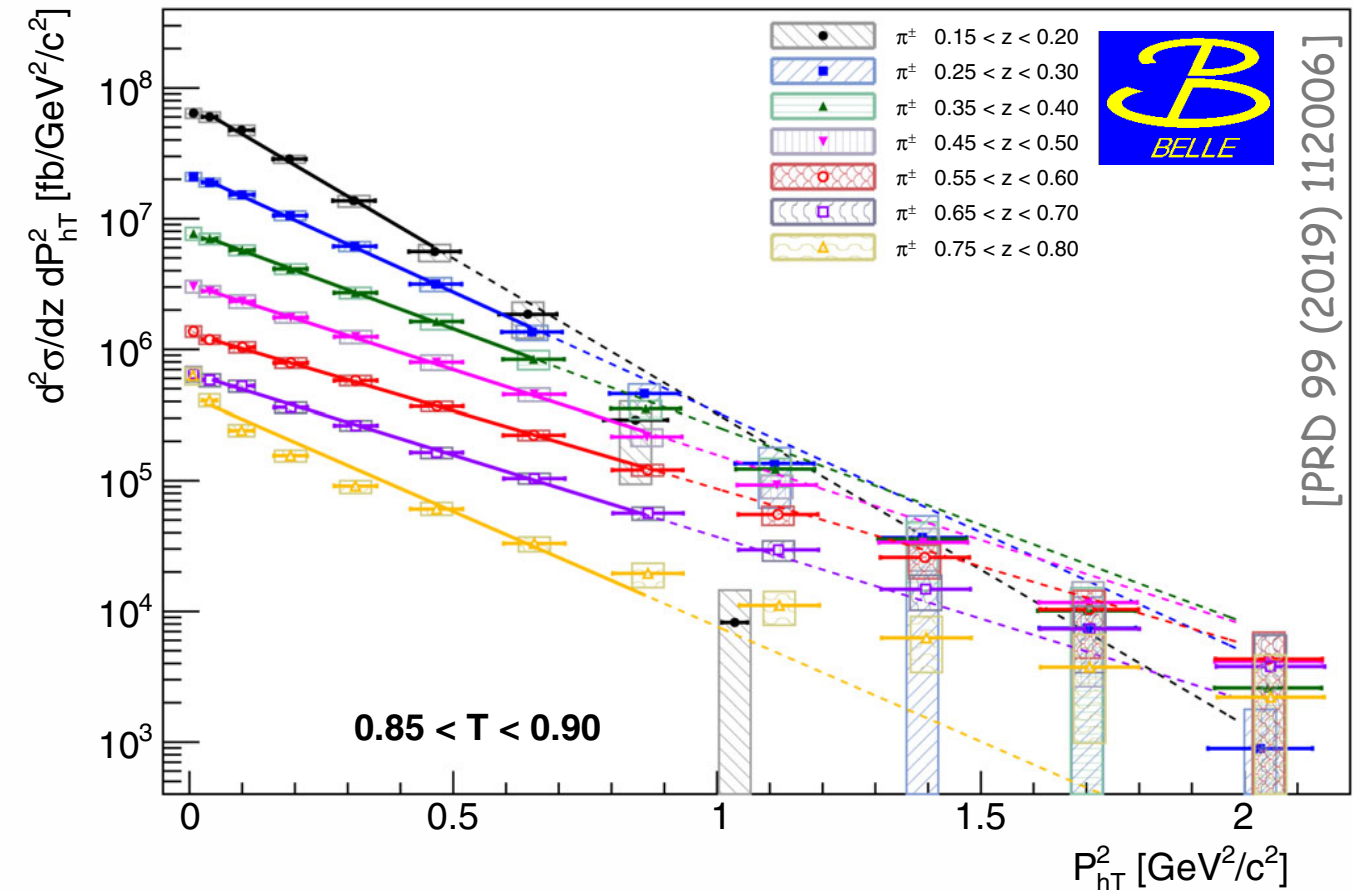
transverse-momentum: Gaussian widths

- $0.85 < T < 0.90$
- fit Gauss to low- P_{hT} data
- mostly well described with possible exception at high z
- deviation from Gauss at large P_{hT}
- clear increase of width with z for low values of z



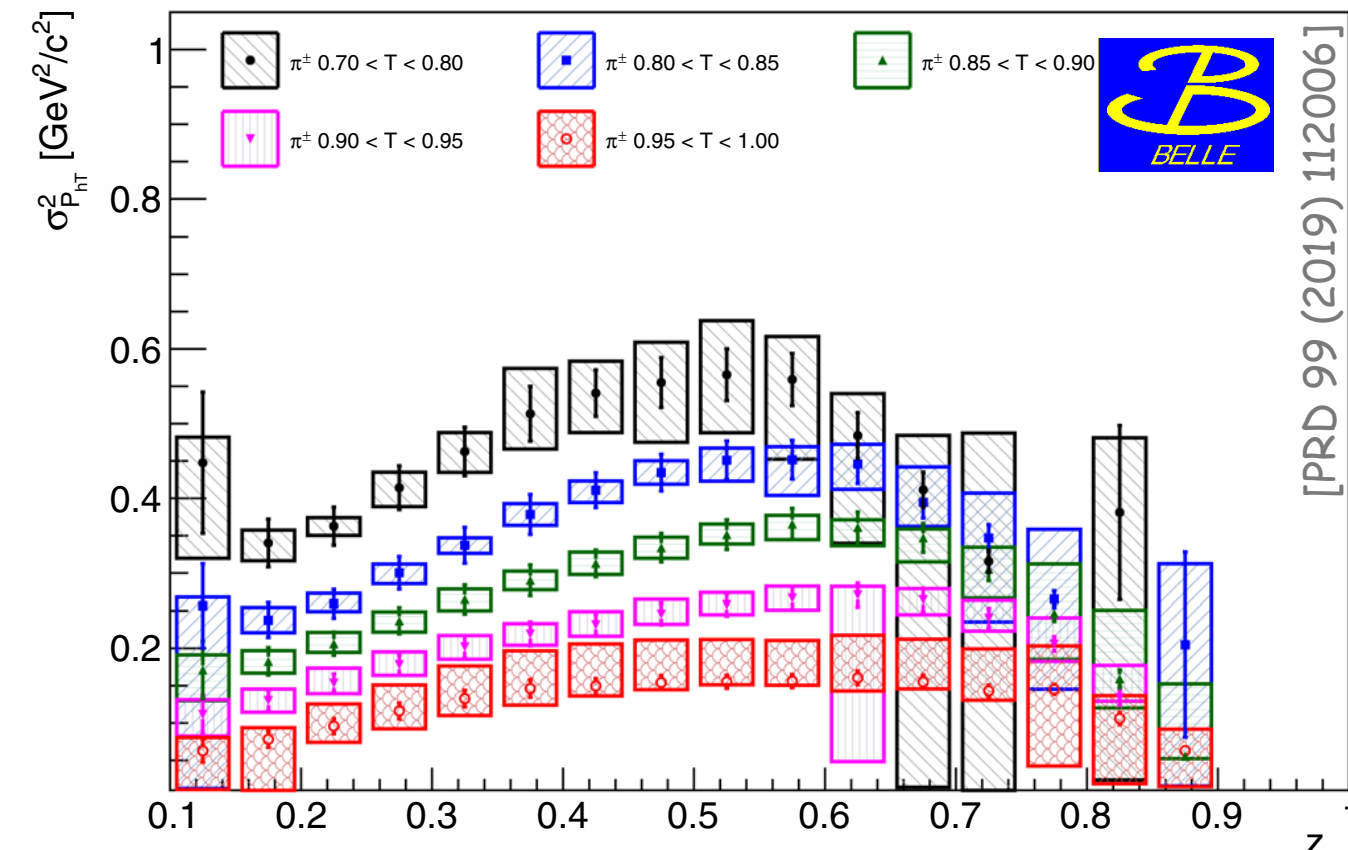
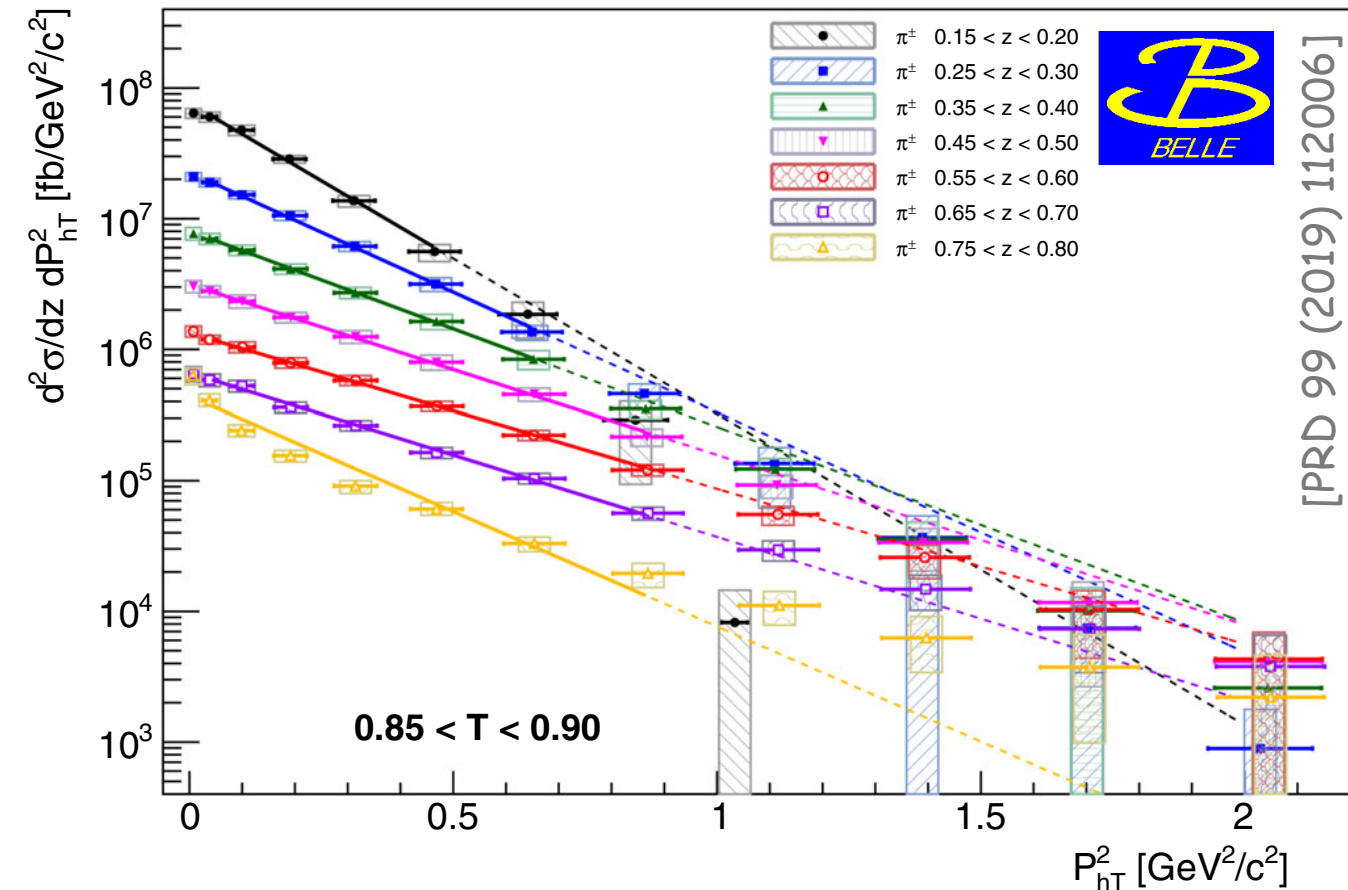
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- deviation from Gauss at large P_{hT}
- clear increase of width with z for low values of z
- Gaussian widths as function of z
 - general increase with z with turnover at larger values of z for mesons
 - protons with smaller width and more linear rise with z



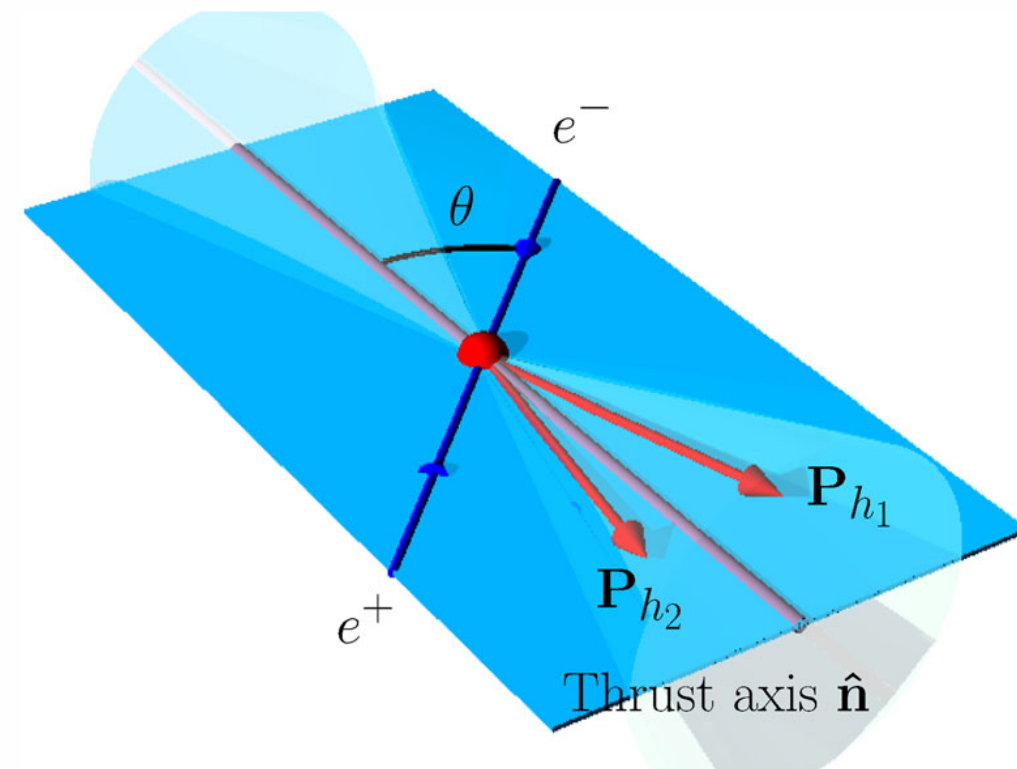
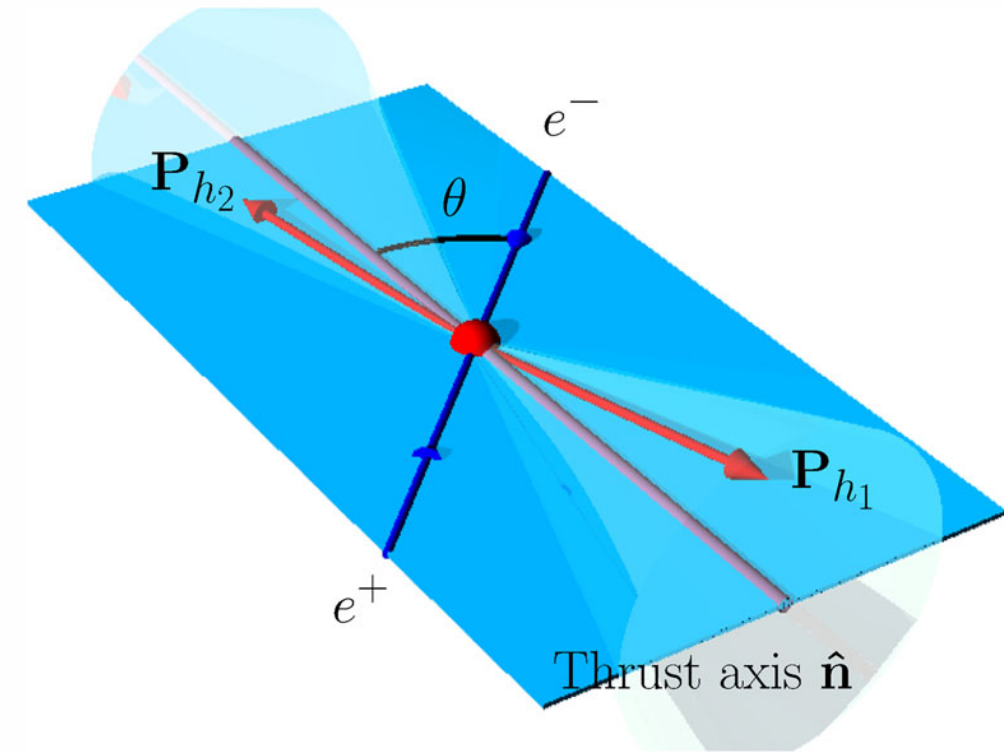
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- fit Gauss to low- P_{hT} data
- mostly well described with possible exception at high z
- deviation from Gauss at large P_{hT}
- clear increase of width with z for low values of z
- Gaussian widths depend on z and T
- general increase with z with turnover at larger values of z
- clear decrease of widths with increase of T
- particles more and more collimated



hadron-pair production

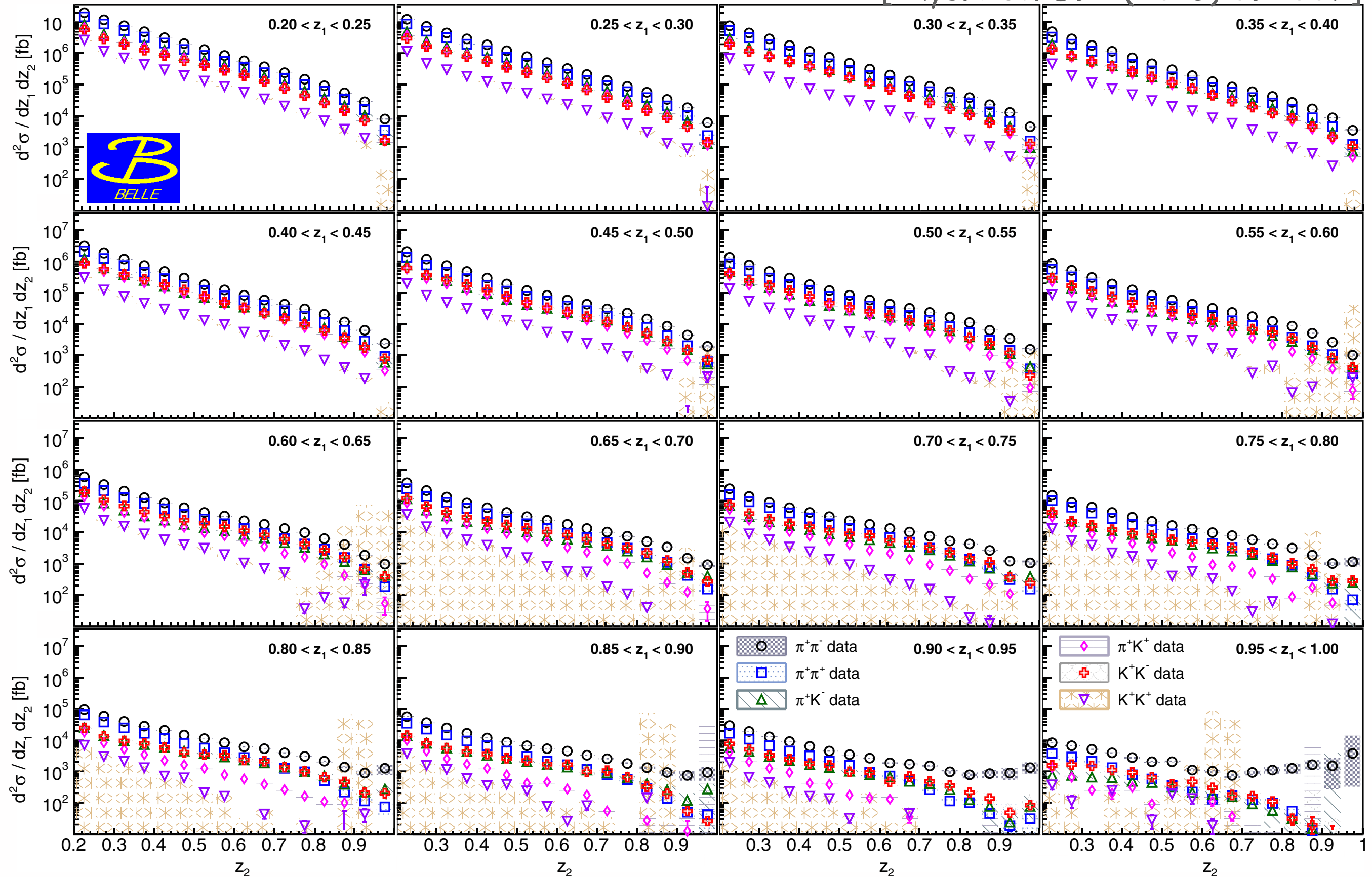
- single-hadron production has low discriminating power for parton flavor
- can use 2nd hadron in opposite hemisphere to “tag” flavor
- mainly sensitive to product of single-hadron FFs
- if hadrons in same hemisphere:
dihadron fragmentation
 - a la de Florian & Vanni [Phys. Lett. B 578 (2004) 139]
 - a la Collins, Heppelmann & Ladinsky [Nucl. Phys. B 420 (1994) 565]; Boer, Jacobs & Radici [Phys. Rev. D 67 (2003) 094003]
- raises question of defining hemispheres



no hemisphere selection

hadron-pair production

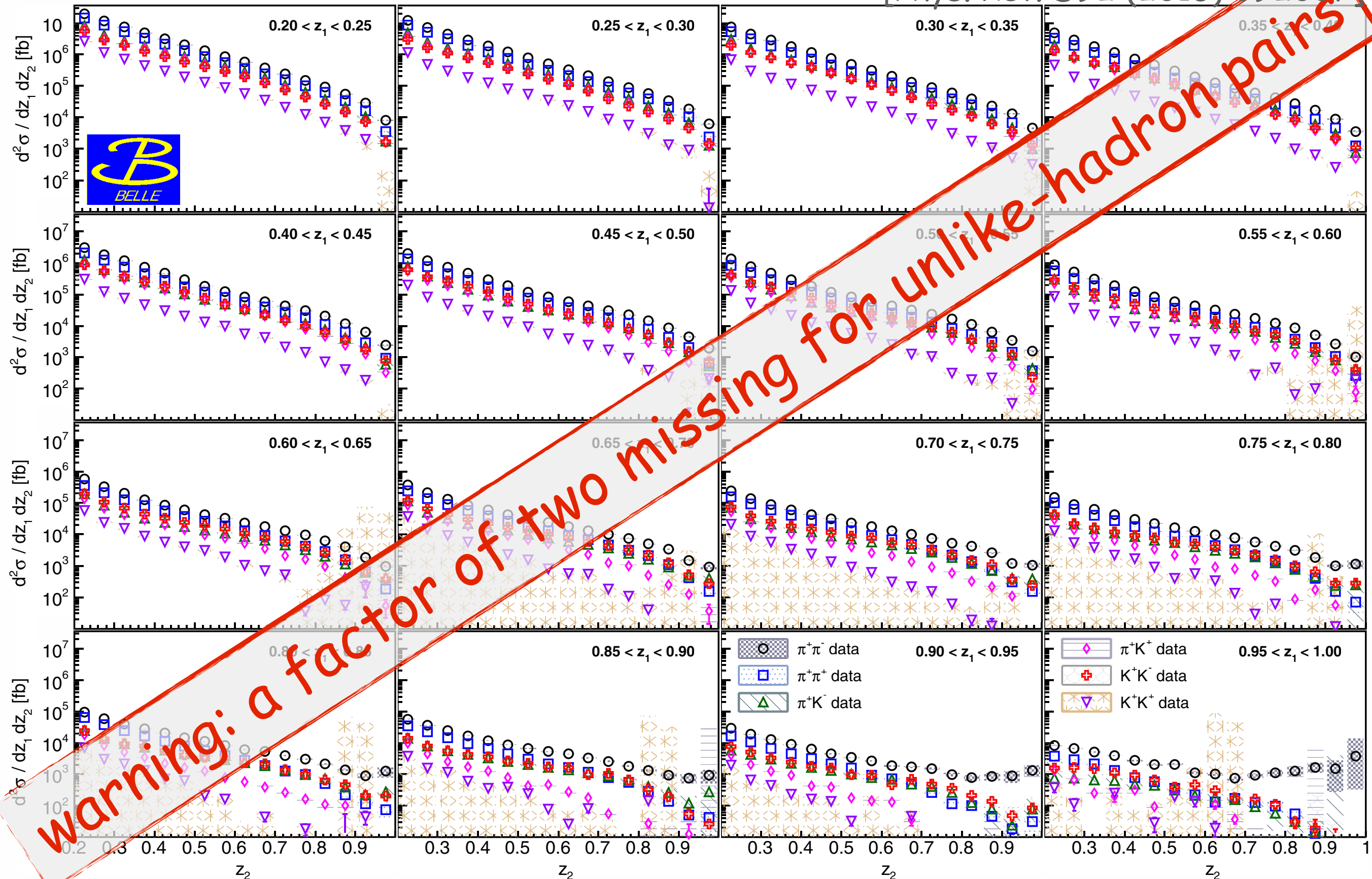
[Phys. Rev. D92 (2015) 092007]



hadron-pair production

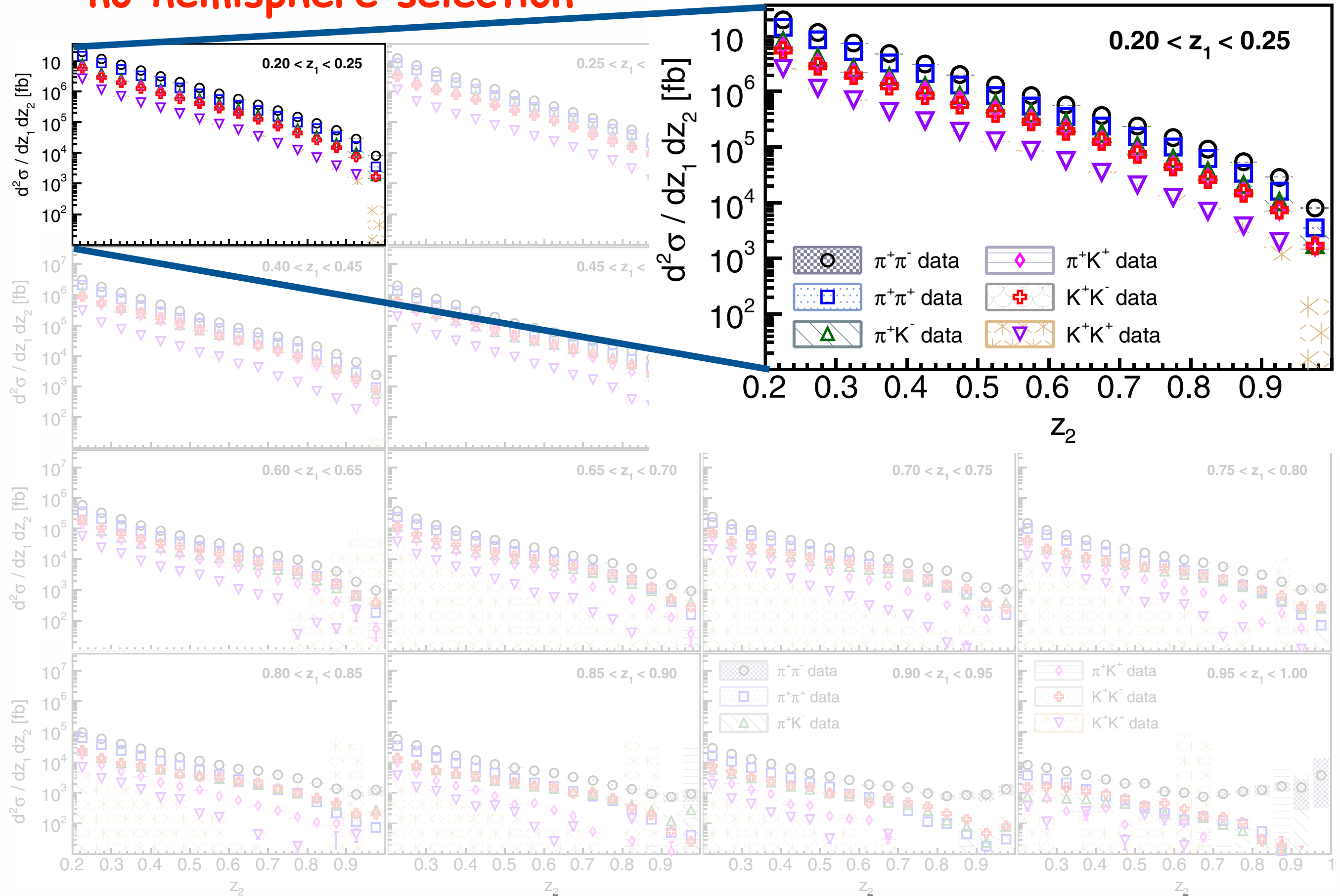
no hemisphere selection

[Phys. Rev. D92 (2015) 092007]



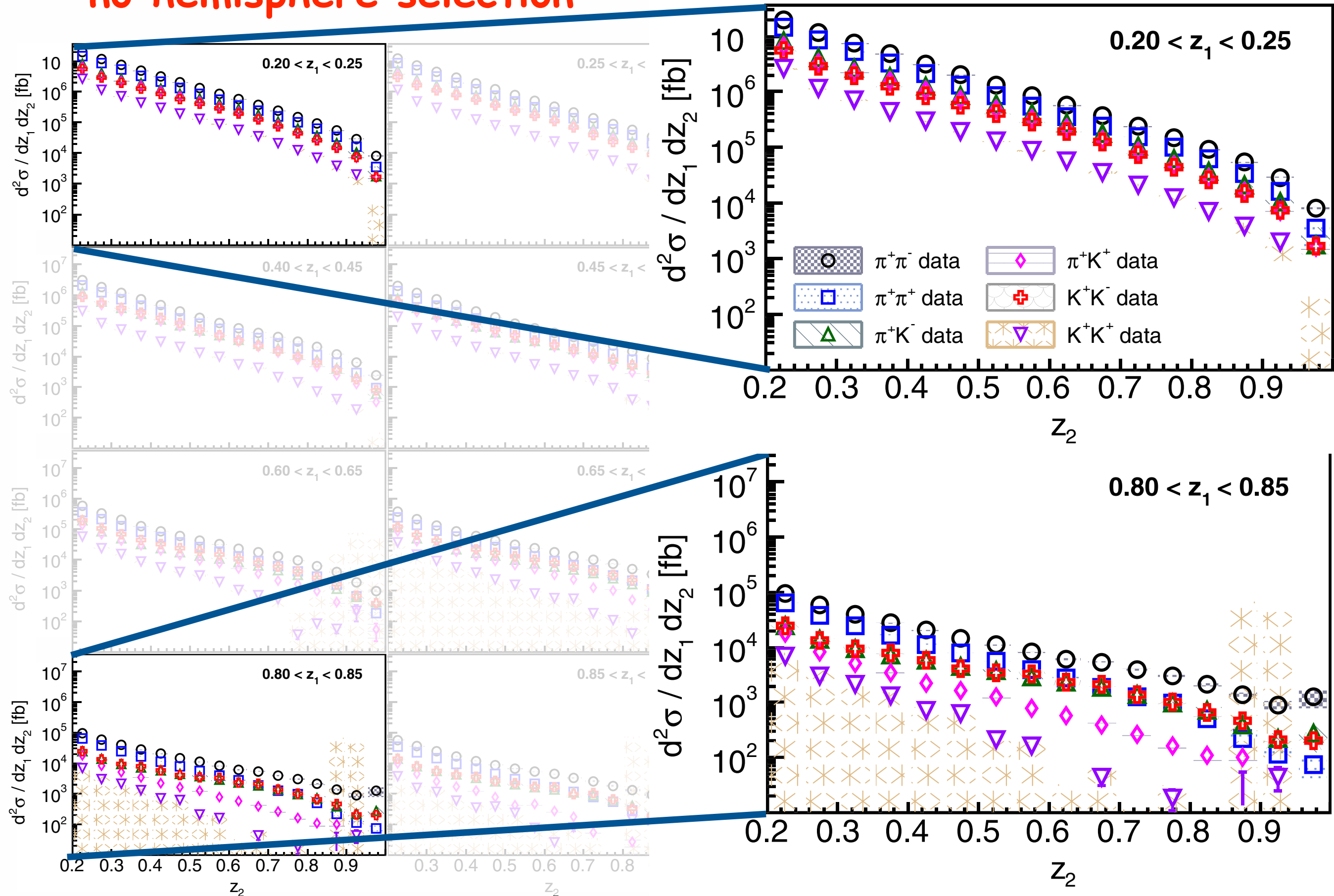
hadron-pair production

no hemisphere selection



hadron-pair production

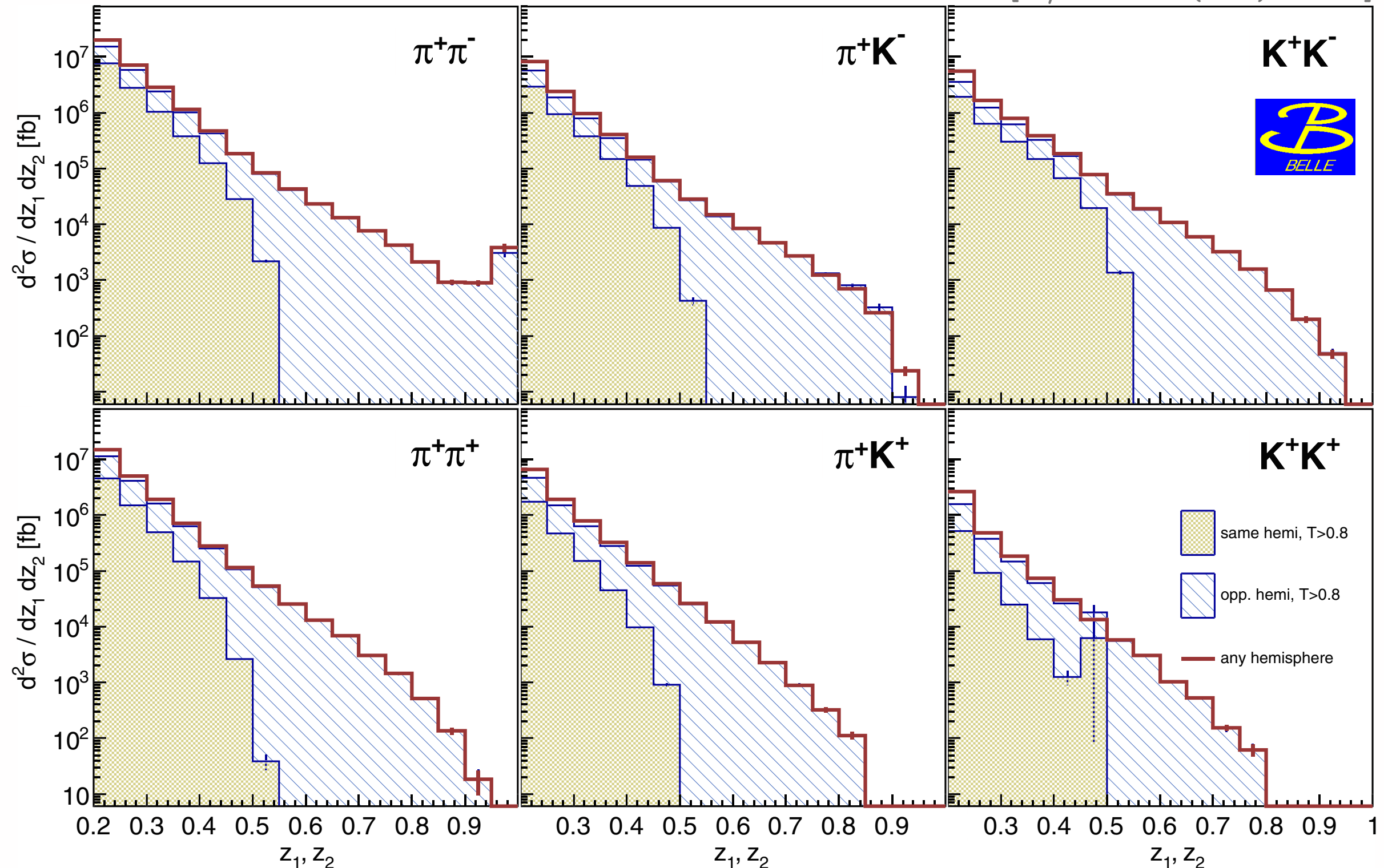
no hemisphere selection



hadron-pairs: topology comparison

- any hemisphere vs. opposite- & same-hemisphere pairs
- same-hemisphere pairs with kinematic limit at $z_1=z_2=0.5$

[Phys. Rev. D92 (2015) 092007]

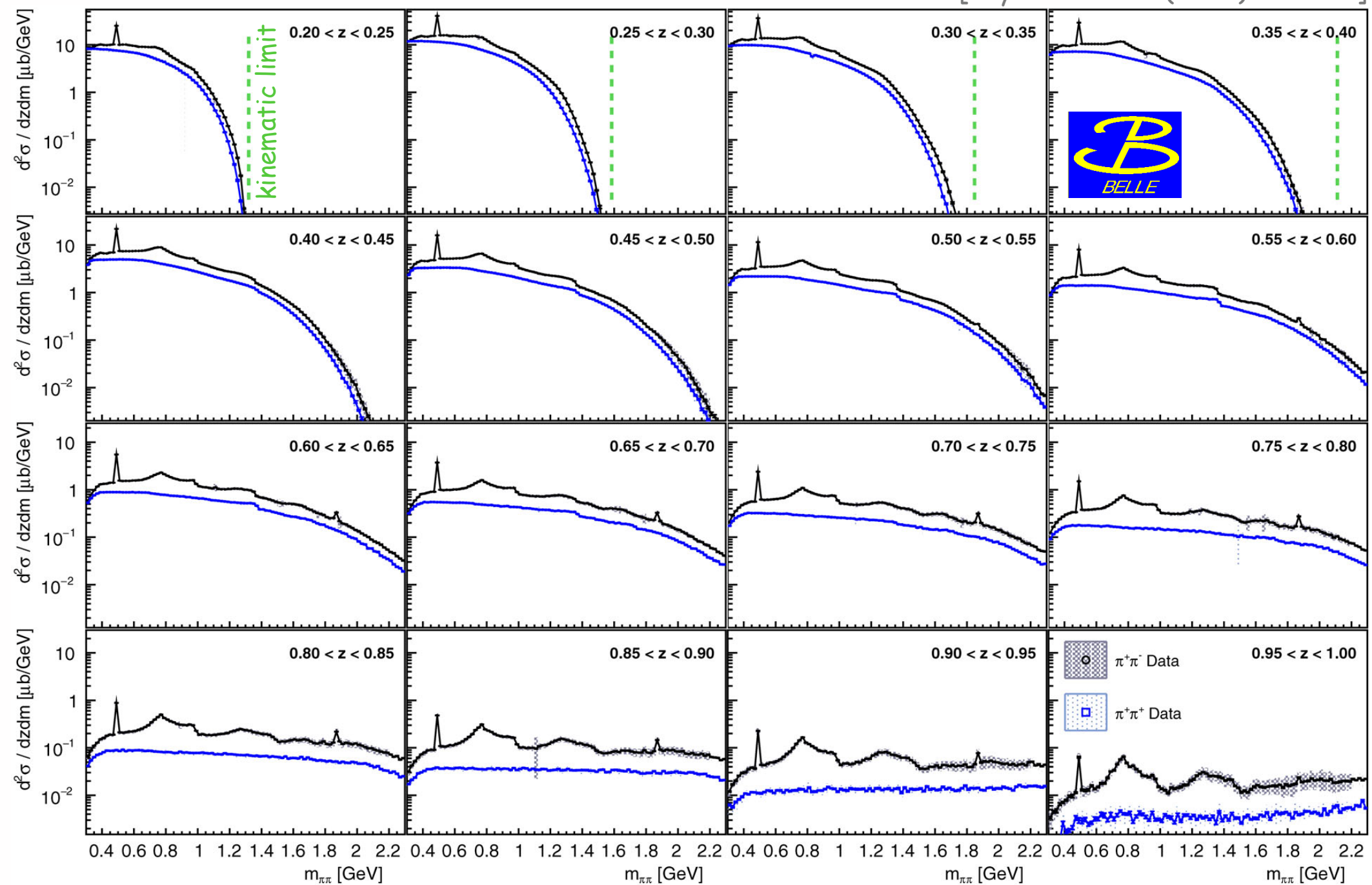
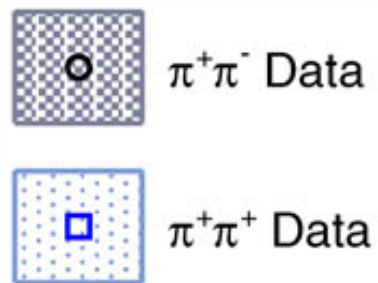


same-hemisphere data: $M_{h_1 h_2}$ dependence

[Phys. Rev. D96 (2017) 032005]

unlike-sign
hadron pairs

like-sign
hadron pairs

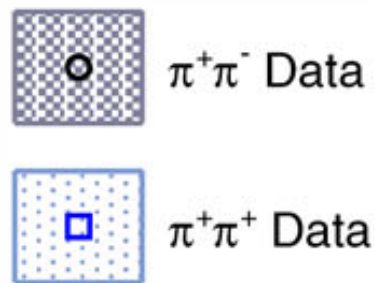


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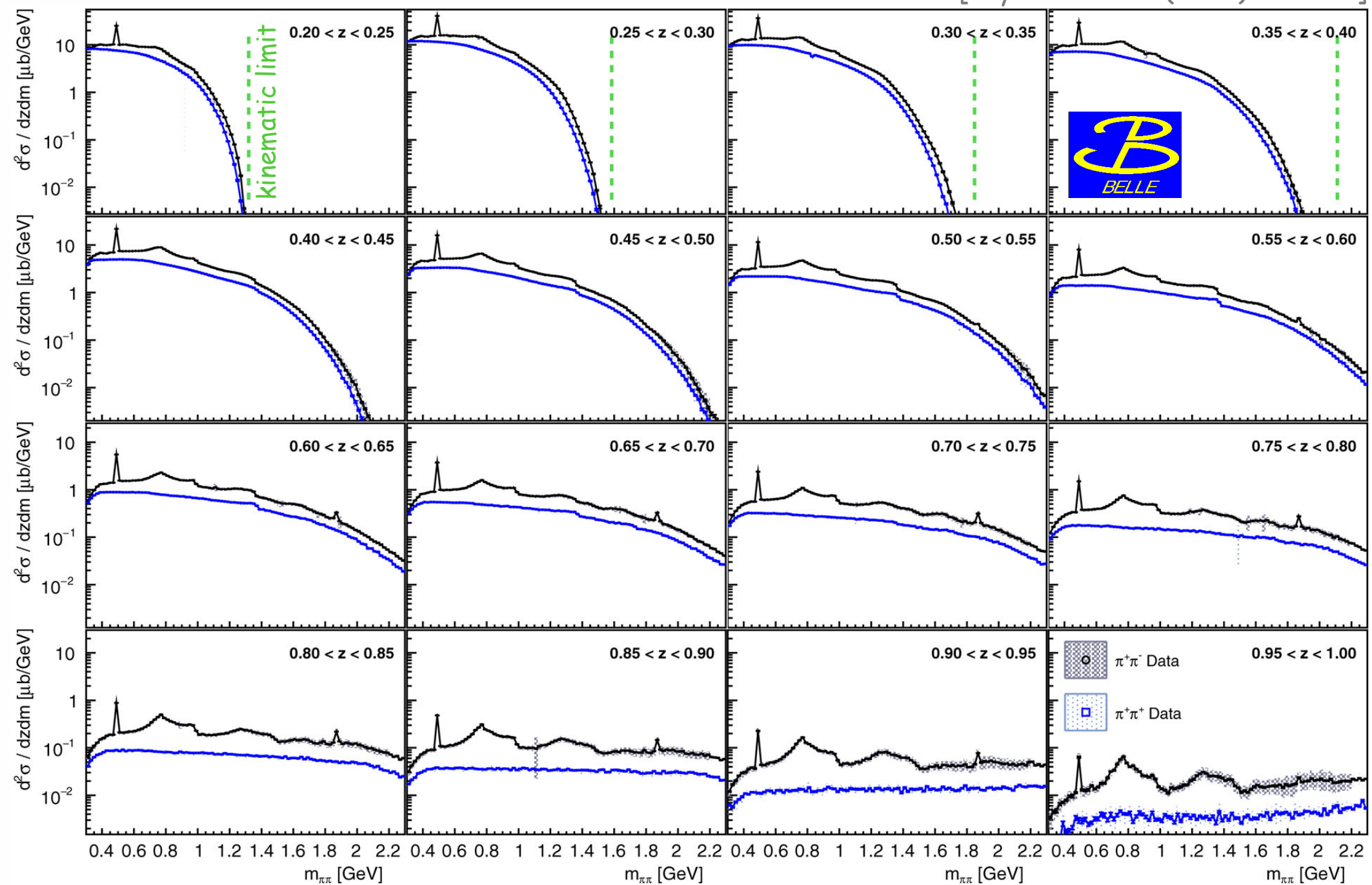
unlike-sign
hadron pairs

like-sign
hadron pairs



$T > 0.8$

$z_{1,2} > 0.1$

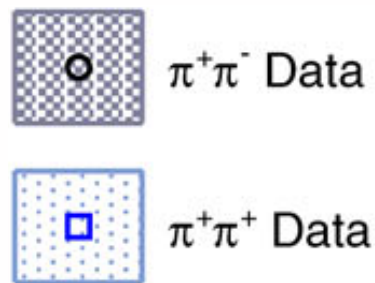


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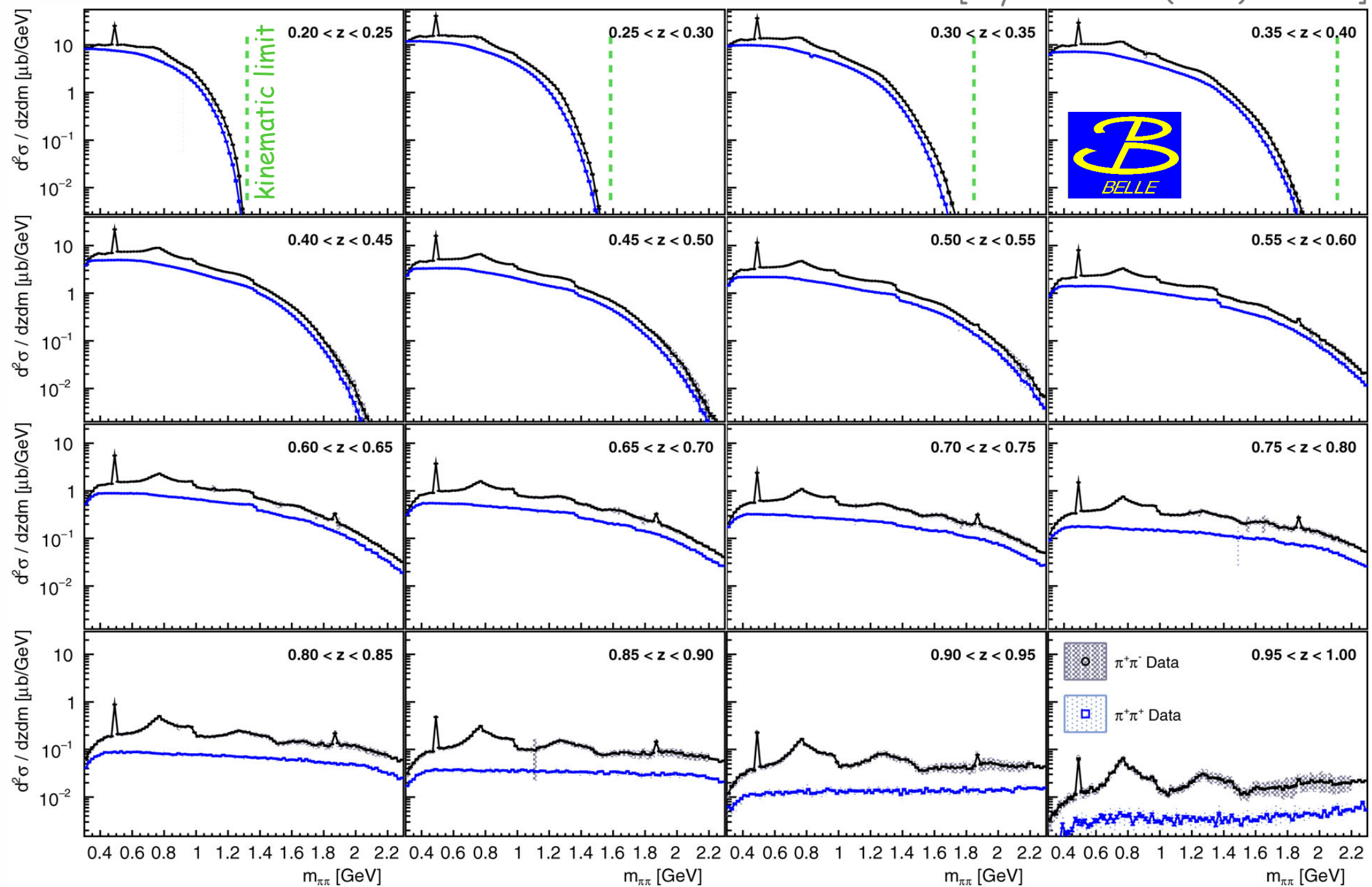
unlike-sign
hadron pairs

like-sign
hadron pairs



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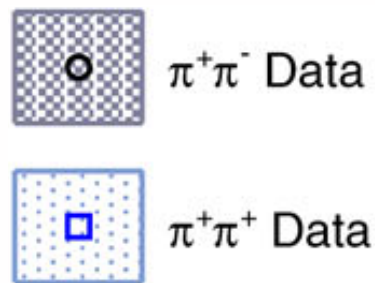
- unlike-sign pairs with clear decay and resonance structure: K_s , ρ^0 ...
- like-sign pairs with much smoother and smaller cross sections

same-hemisphere data: $M_{h_1 h_2}$ dependence

[Phys. Rev. D96 (2017) 032005]

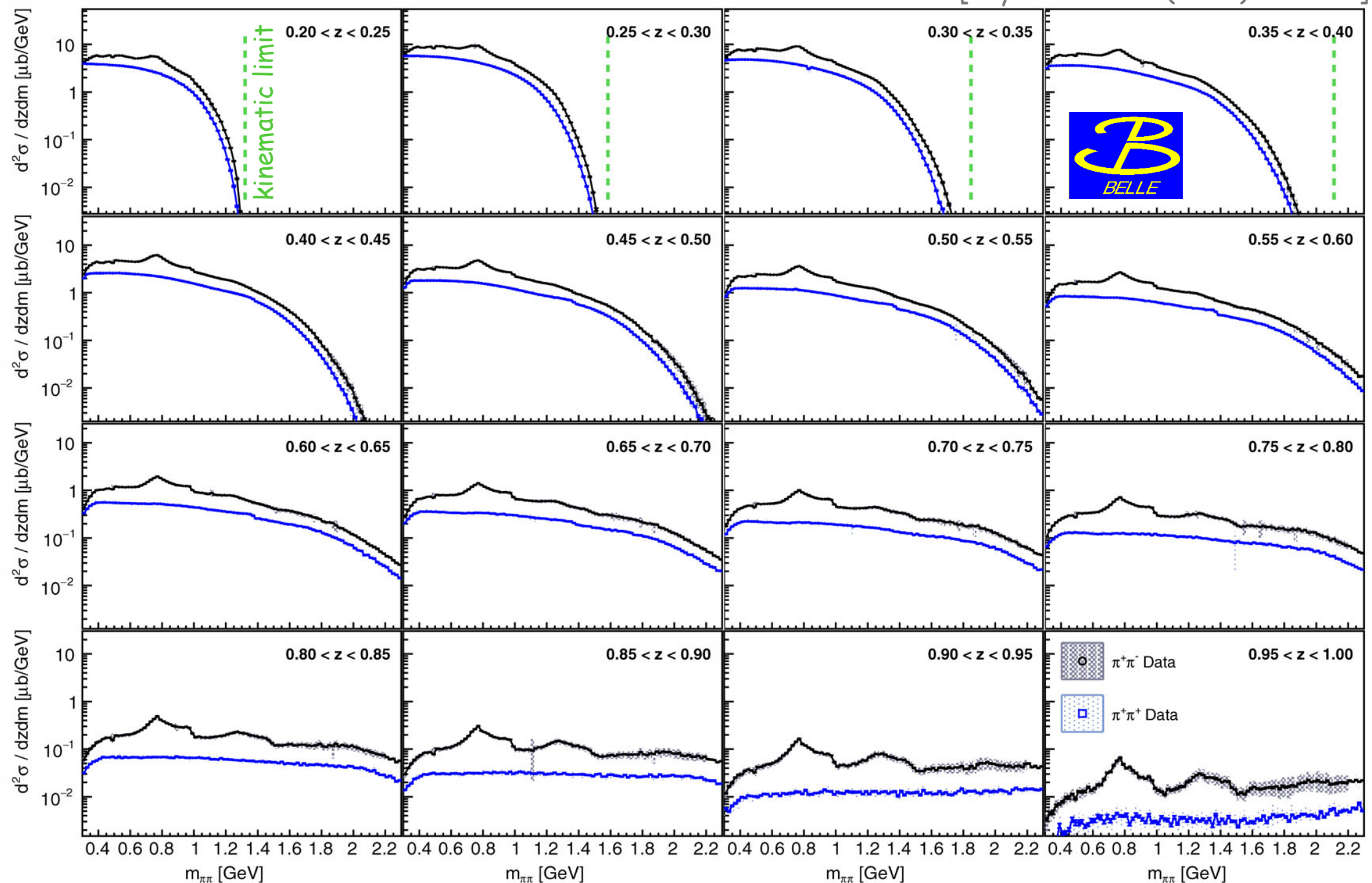
unlike-sign
hadron pairs

like-sign
hadron pairs



$T > 0.8$

$z_{1,2} > 0.1$

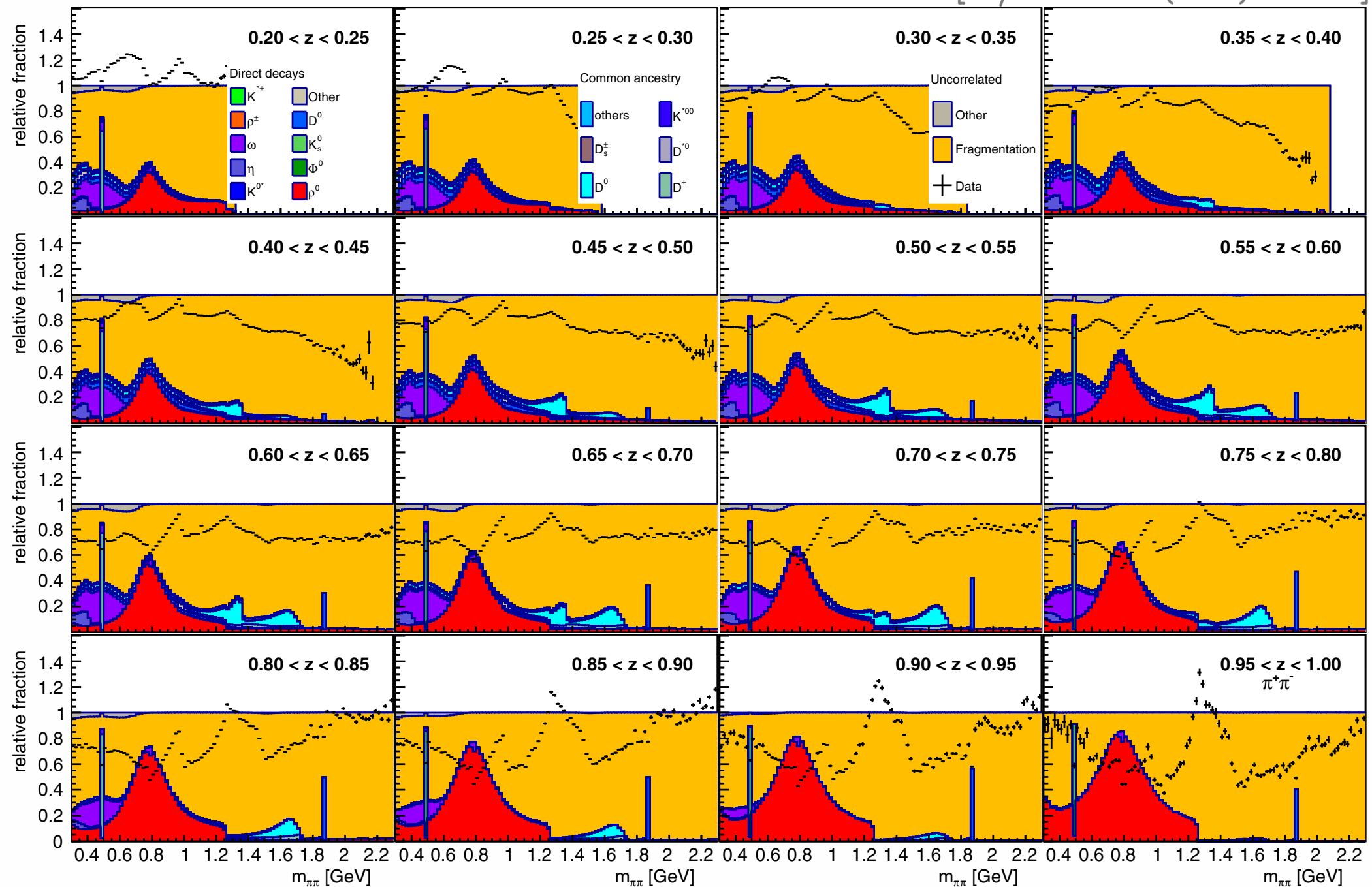


- cross sections after (MC-based) removal of weak-decay contributions
- relies on good description of those channels in PYTHIA

same-hemisphere data: $M_{h_1 h_2}$ dependence

unlike-sign
pion pairs

[Phys. Rev. D96 (2017) 032005]



$T > 0.8$
 $z_{1,2} > 0.1$

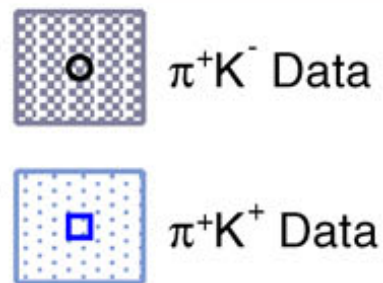
- decomposition based on PYTHIA simulation
- clear differences in invariant-mass dependence between MC and data

same-hemisphere data: $M_{h_1 h_2}$ dependence

[Phys. Rev. D96 (2017) 032005]

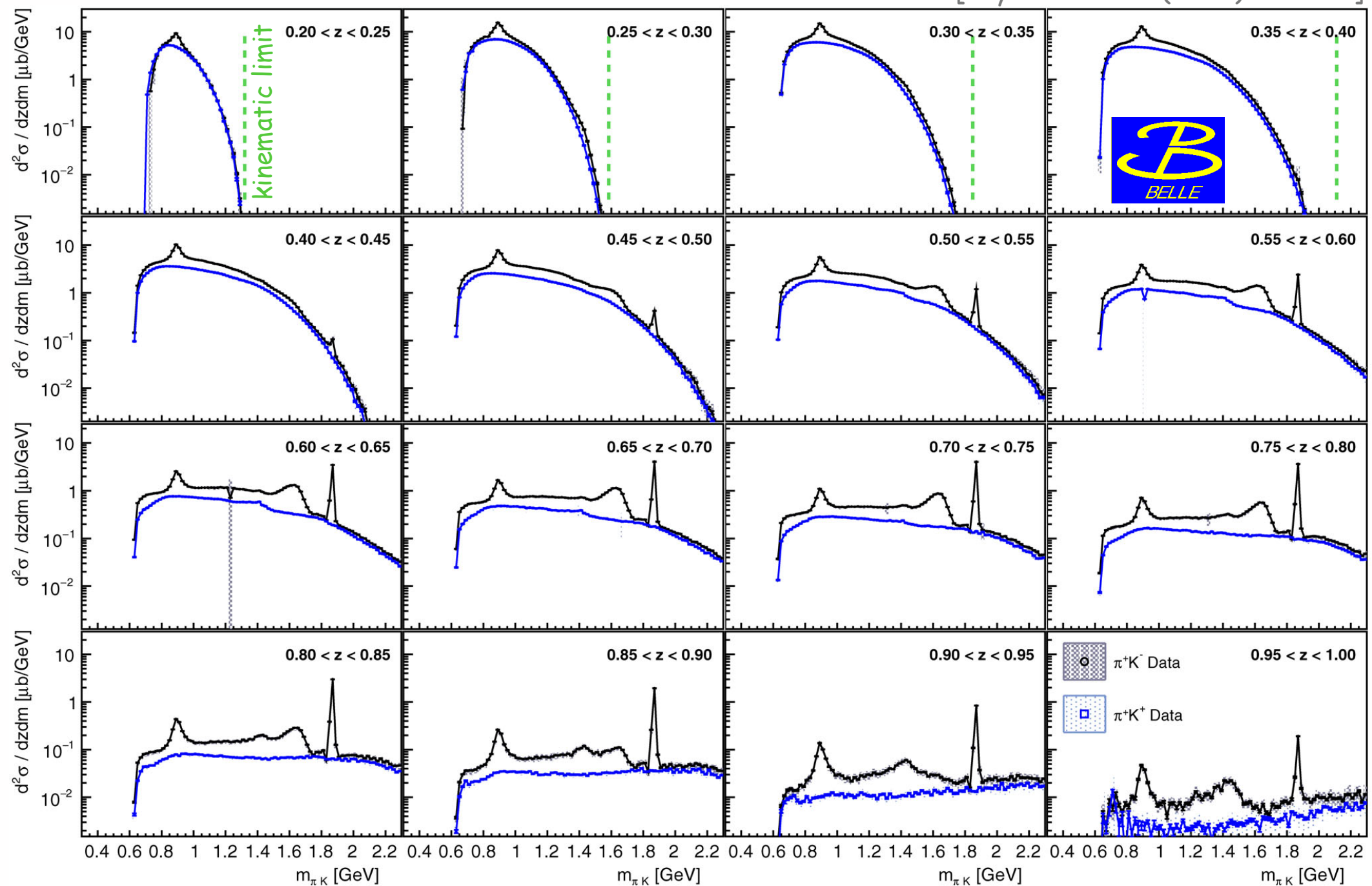
unlike-sign
hadron pairs

like-sign
hadron pairs



$T > 0.8$

$z_{1,2} > 0.1$



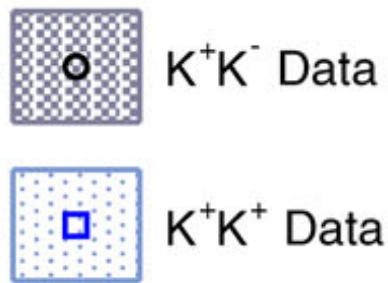
● unlike-sign πK pairs with clear K^* and increased D-decay contributions

same-hemisphere data: $M_{h_1 h_2}$ dependence

[Phys. Rev. D96 (2017) 032005]

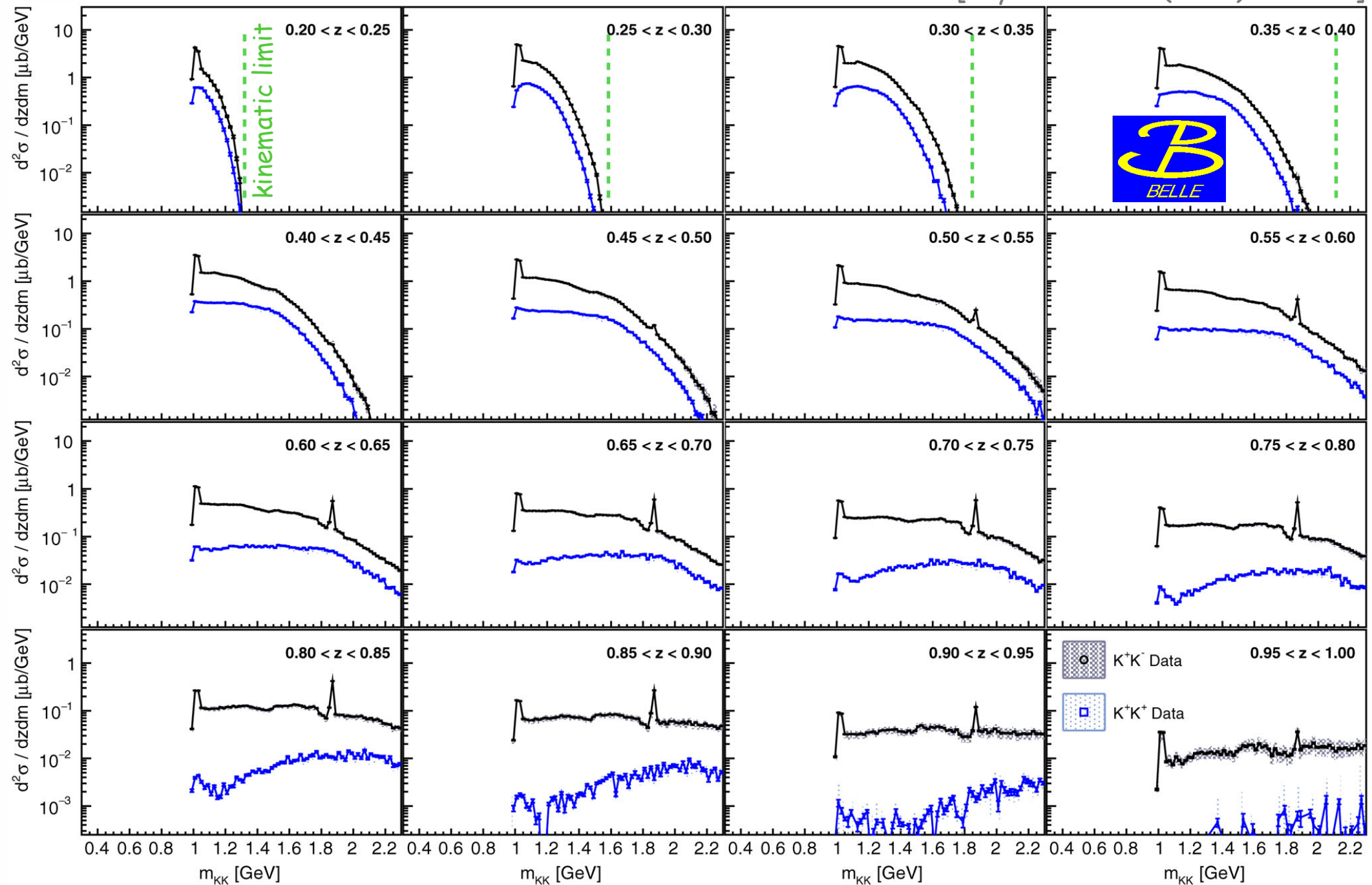
unlike-sign
hadron pairs

like-sign
hadron pairs



$T > 0.8$

$z_{1,2} > 0.1$



- unlike-sign kaon pairs with (again) a decay structure (e.g. ϕ and D)
- like-sign kaon pairs strongly suppressed at larger z

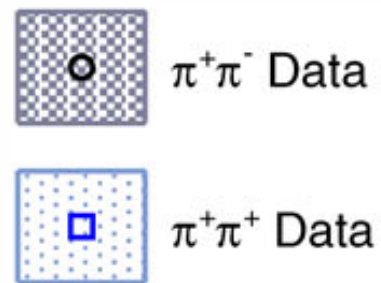
some more details

same-hemisphere data: $M_{h_1 h_2}$ dependence

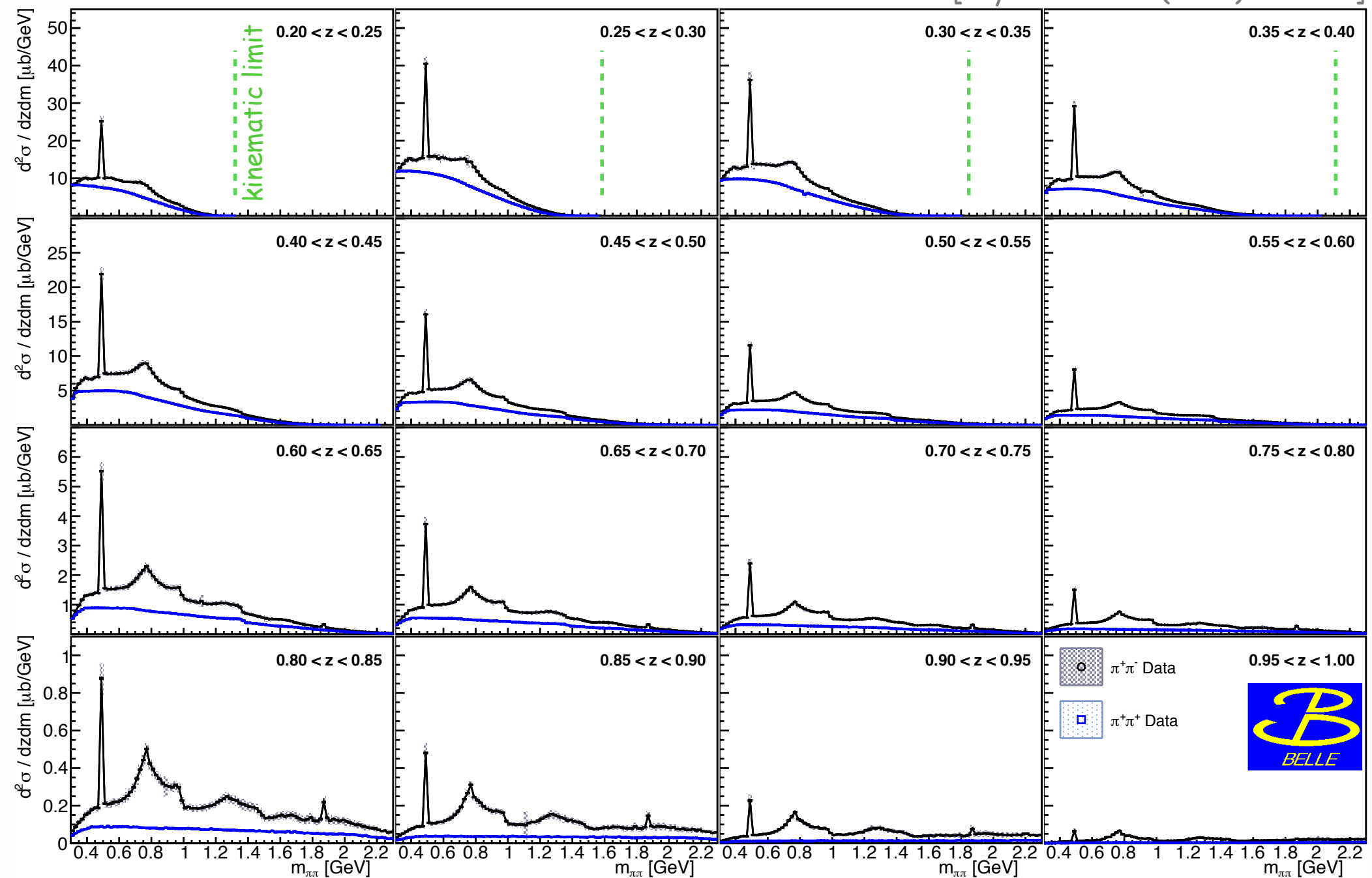
[Phys. Rev. D96 (2017) 032005]

unlike-sign
hadron pairs

like-sign
hadron pairs



$T > 0.8$
 $z_{1,2} > 0.1$





same-hemisphere data: M_{h1h2} dependence

[Phys. Rev. D96 (2017) 032005]

unlike-sign
hadron pairs

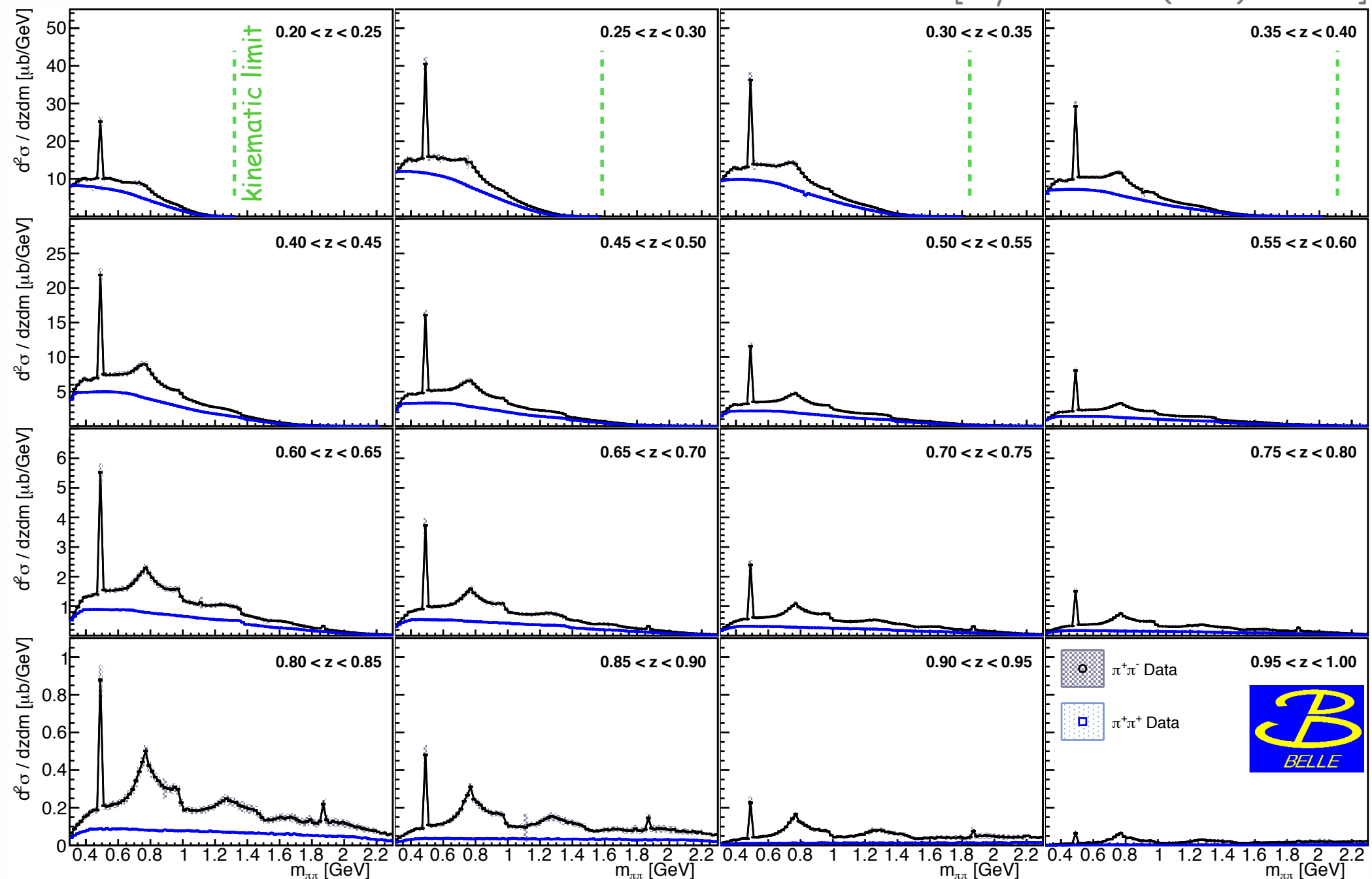
like-sign
hadron pairs

 $\pi^+\pi^-$ Data

 $\pi^+\pi^+$ Data

$T > 0.8$

$z_{1,2} > 0.1$



● **thrust** very useful experimentally to suppress BG and to define hemispheres

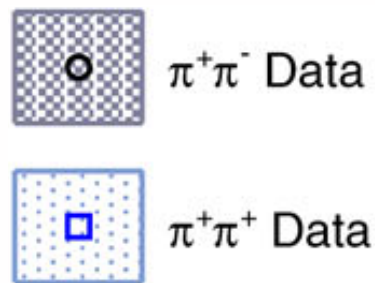
● potentially difficult to incorporate in phenomenology (unlike thrust axis?)

same-hemisphere data: $M_{h_1 h_2}$ dependence

[Phys. Rev. D96 (2017) 032005]

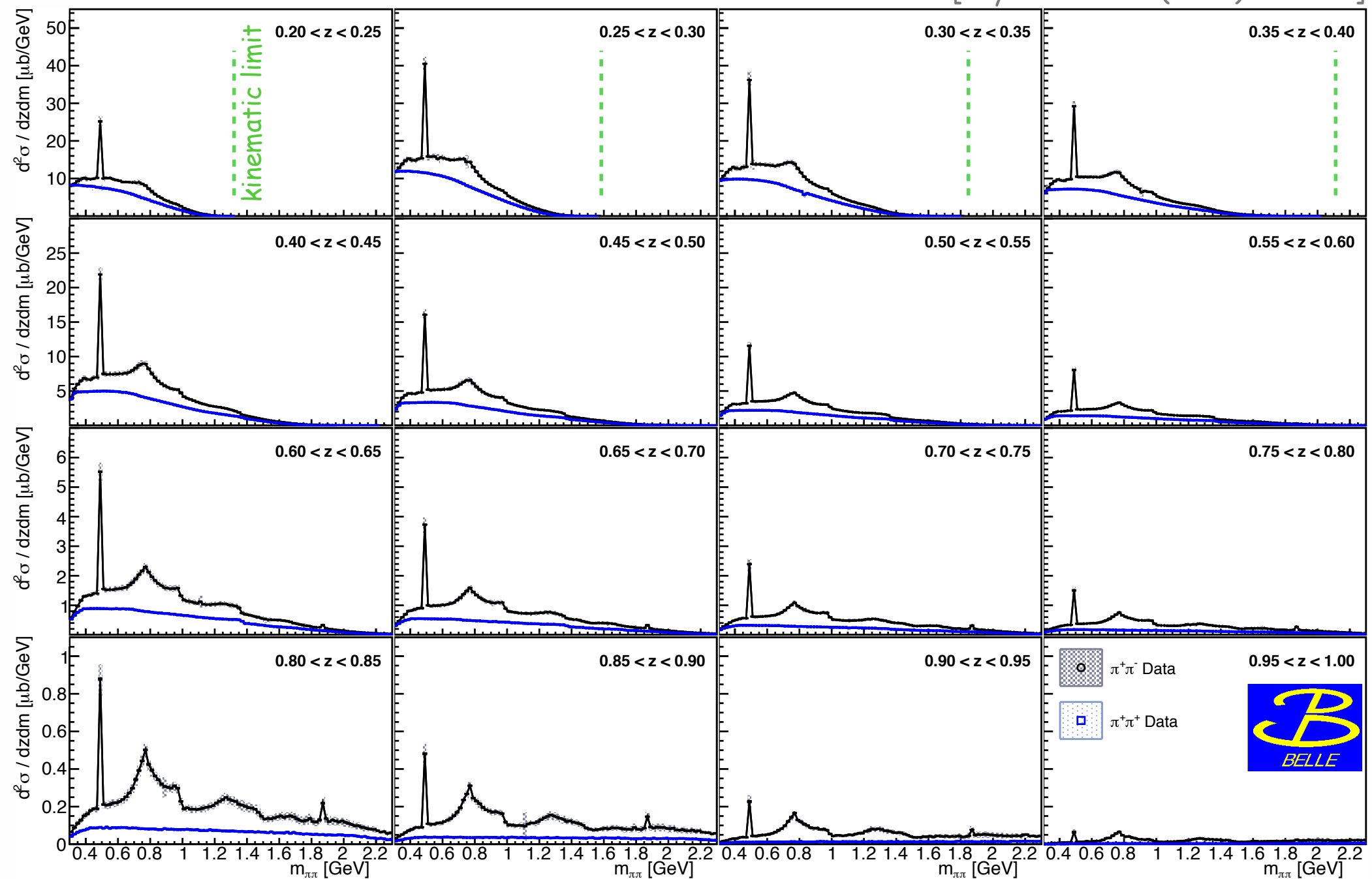
unlike-sign
hadron pairs

like-sign
hadron pairs



$T > 0.8$

$z_{1,2} > 0.1$

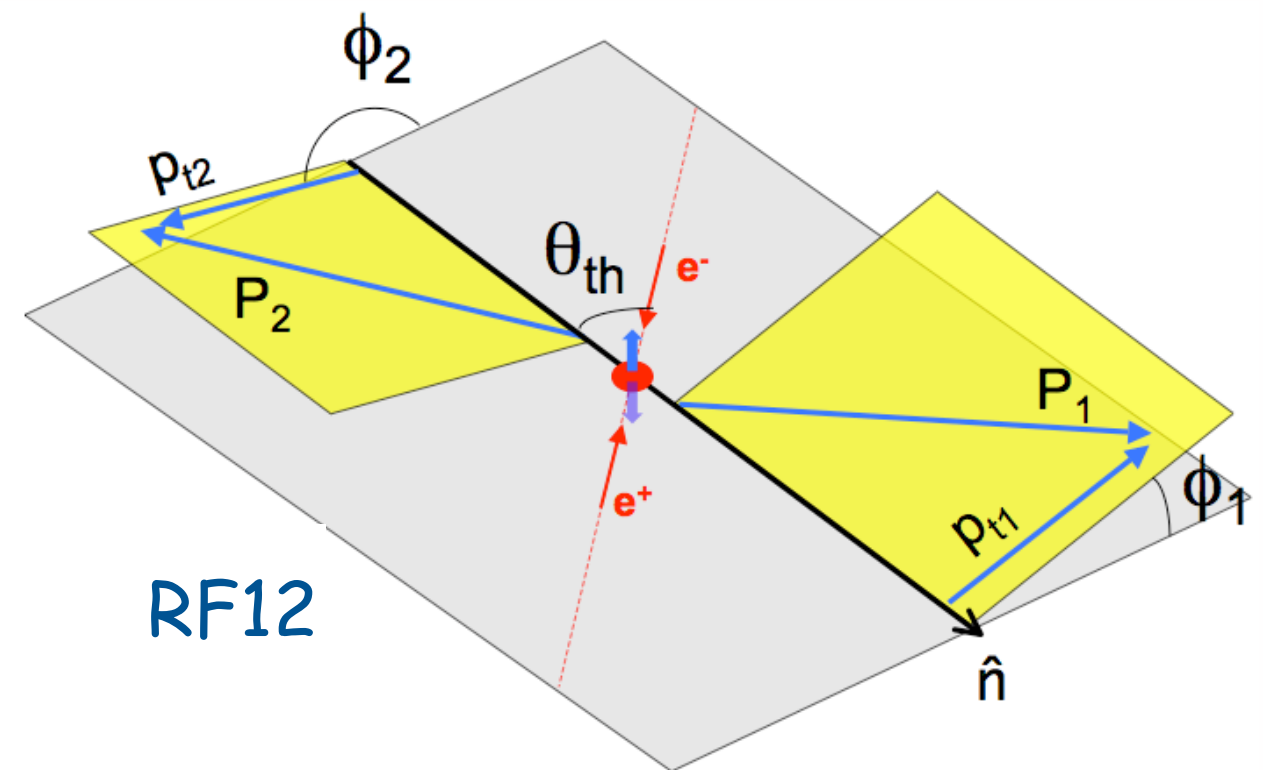
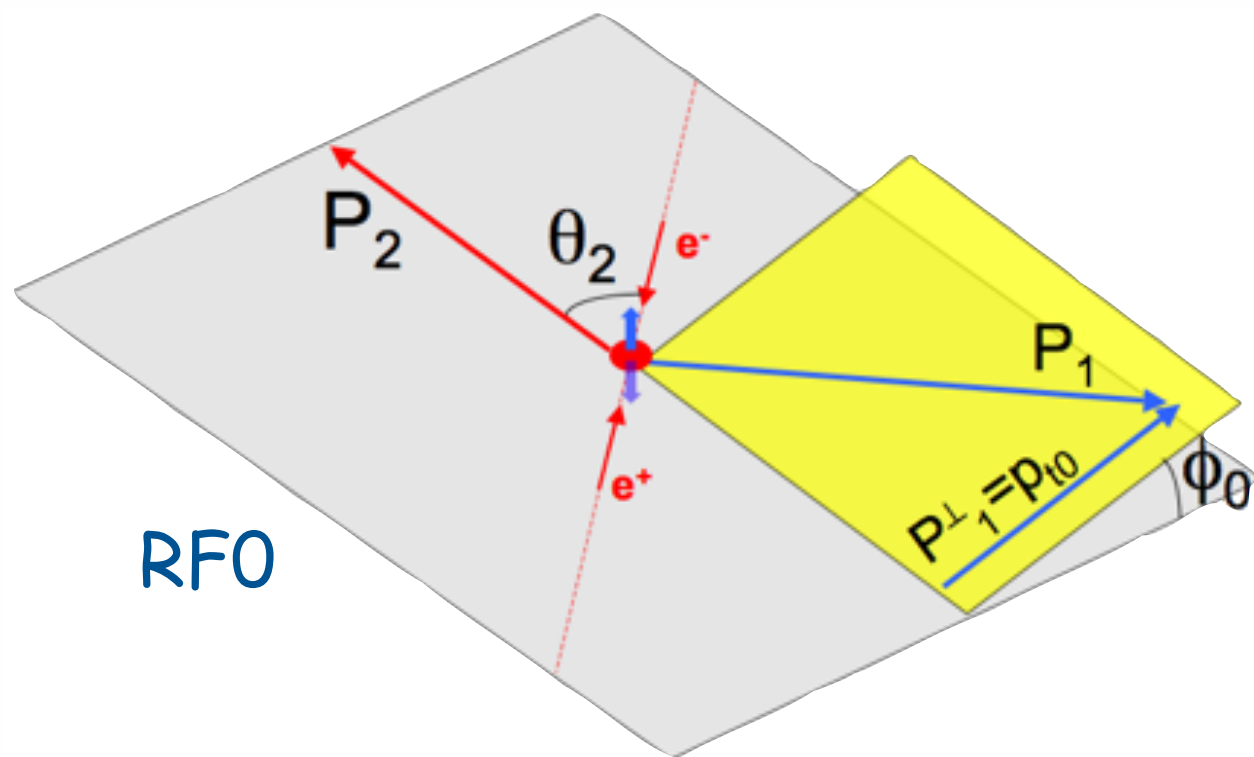


- experimental constraints on **individual z** restricts phase space of hadron pairs, however, not easy to avoid (detection requirements!)
- among others leads to mixing of partial-wave contributions [GS, QCDE'17]

polarization

hadron pairs: angular correlations

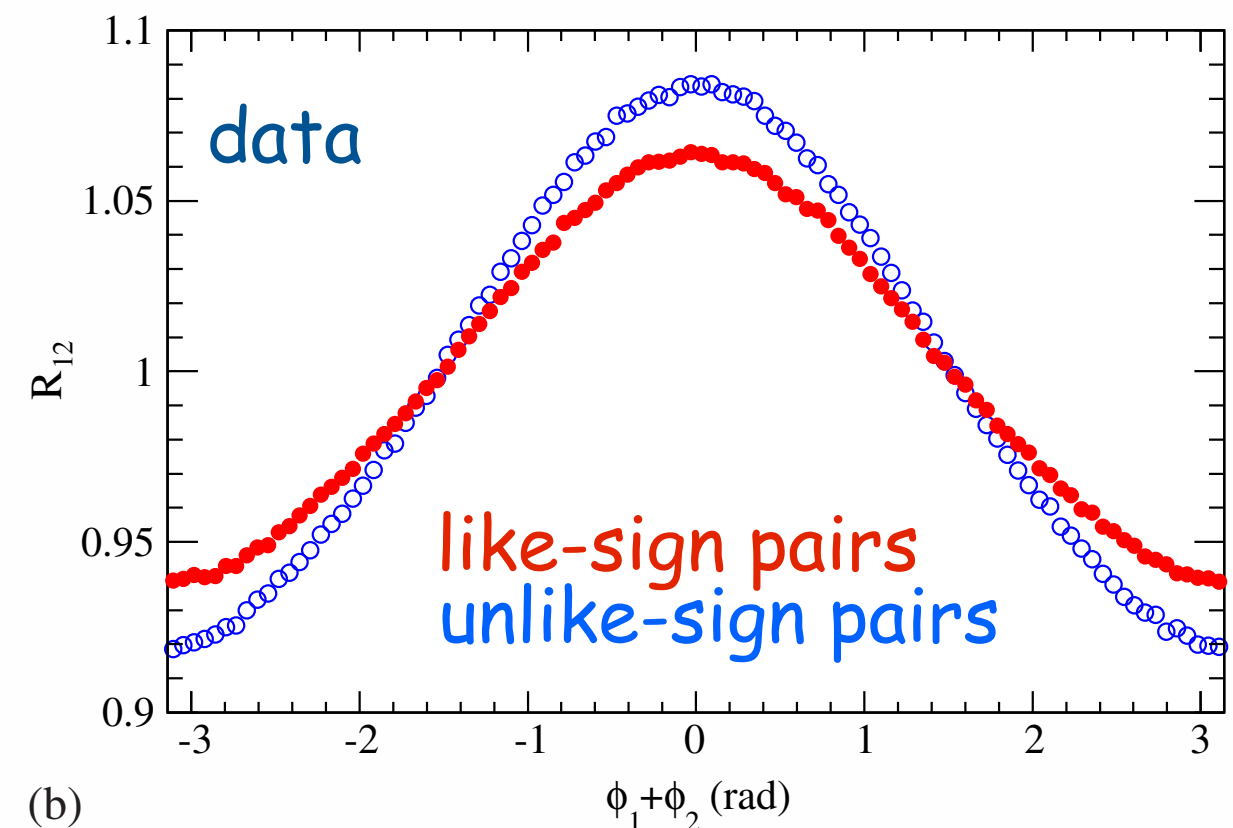
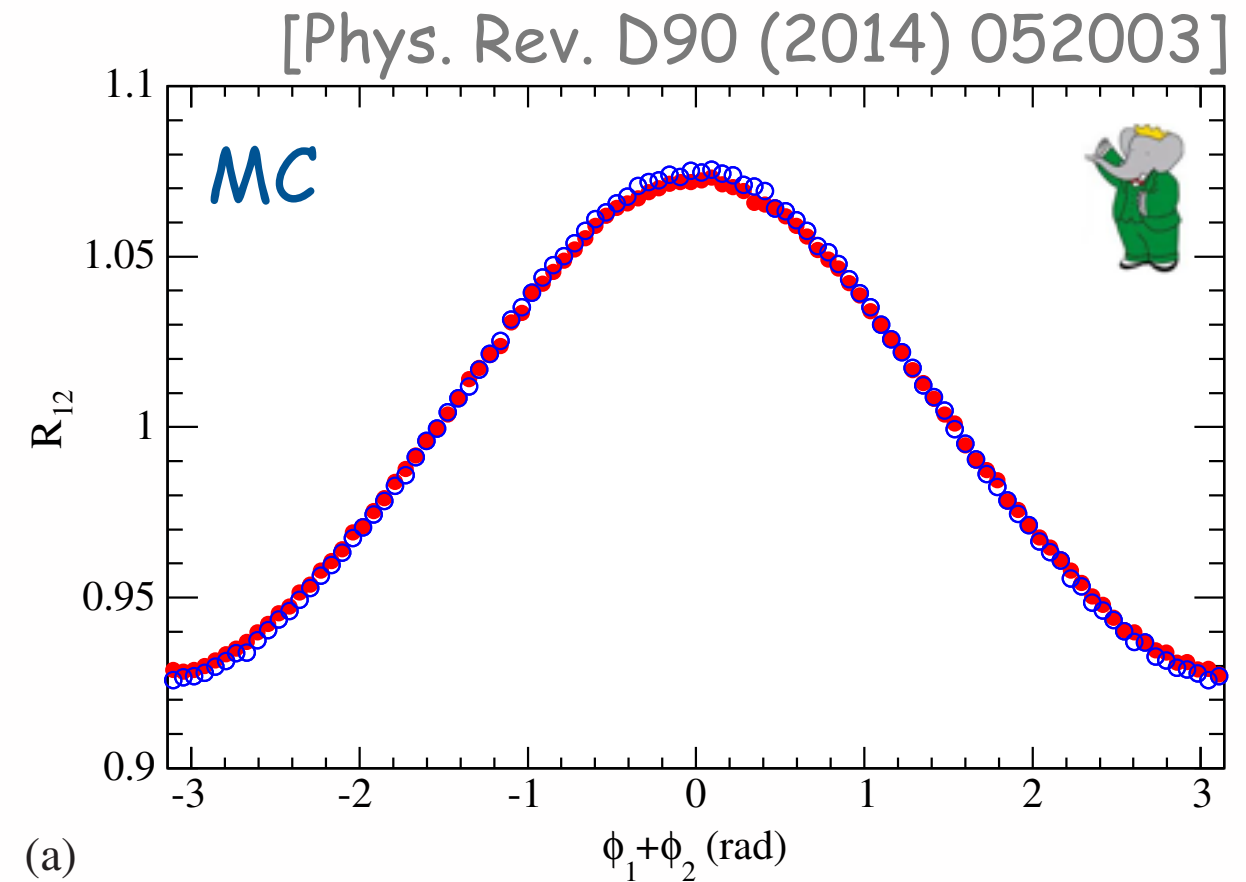
- angular correlations between nearly back-to-back hadrons used to tag transverse quark polarization \rightarrow Collins fragmentation functions
- RF0: one hadron as reference axis $\rightarrow \cos(2\phi_0)$ modulation
- RF12: thrust (or similar) axis $\rightarrow \cos(\phi_1 + \phi_2)$ modulation



- RF0 and RF12: different convolutions over transverse momenta
- debatable: MC used to "correct" thrust axis to $q\bar{q}$ axis

hadron pairs: angular correlations

- challenge: large modulations even without Collins effect (e.g., MC)

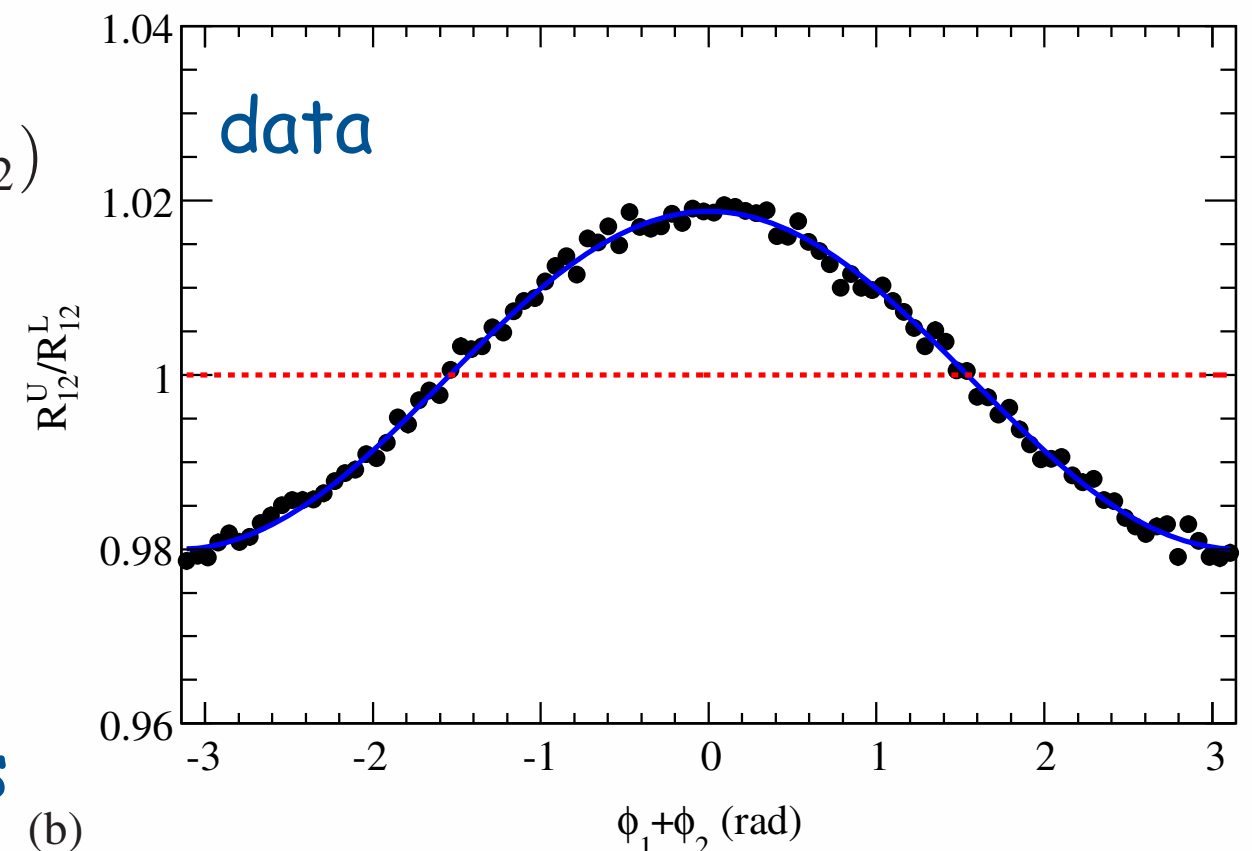
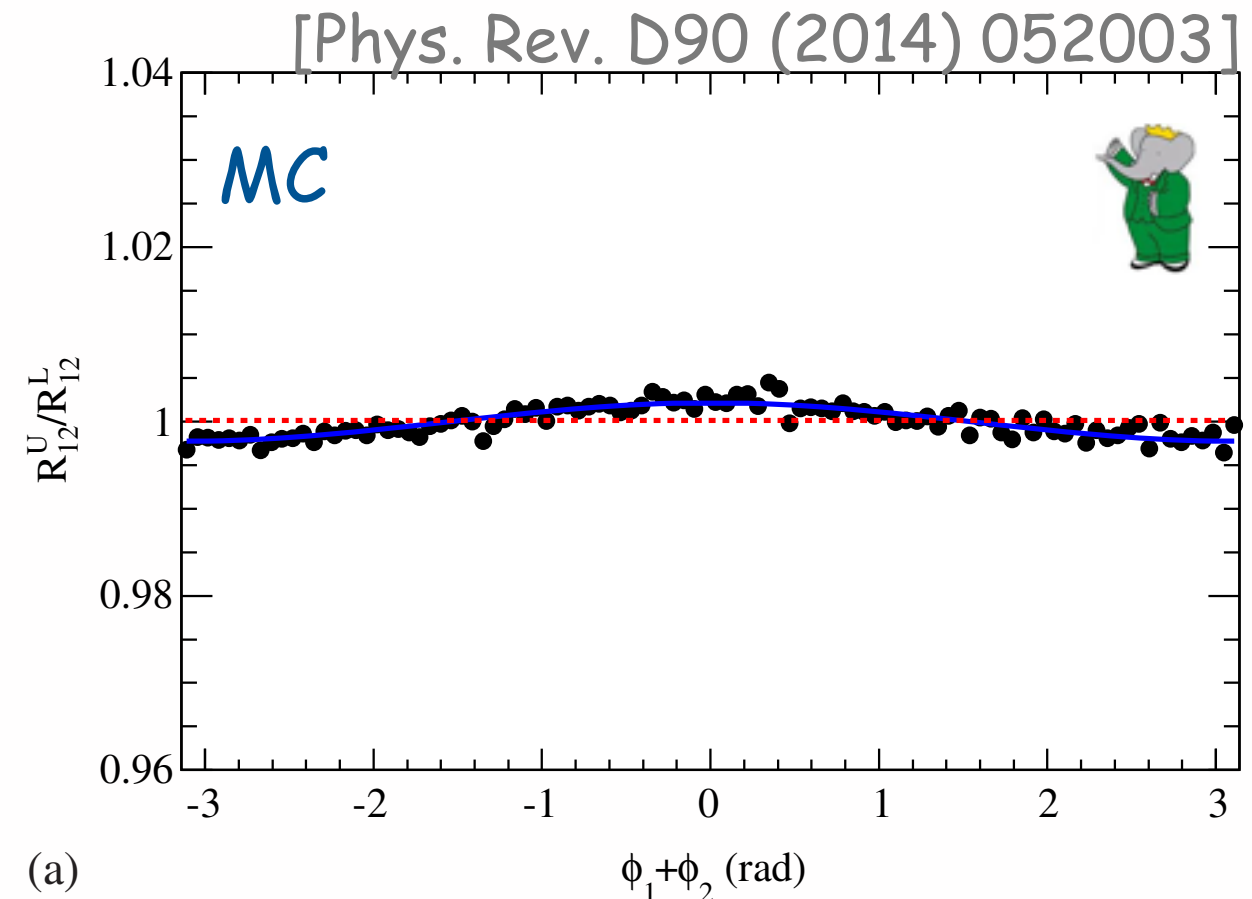


hadron pairs: angular correlations

- challenge: large modulations even without Collins effect (e.g., MC)
- construct double ratio of normalized-yield distributions R_{12} , e.g. unlike-/like-sign:

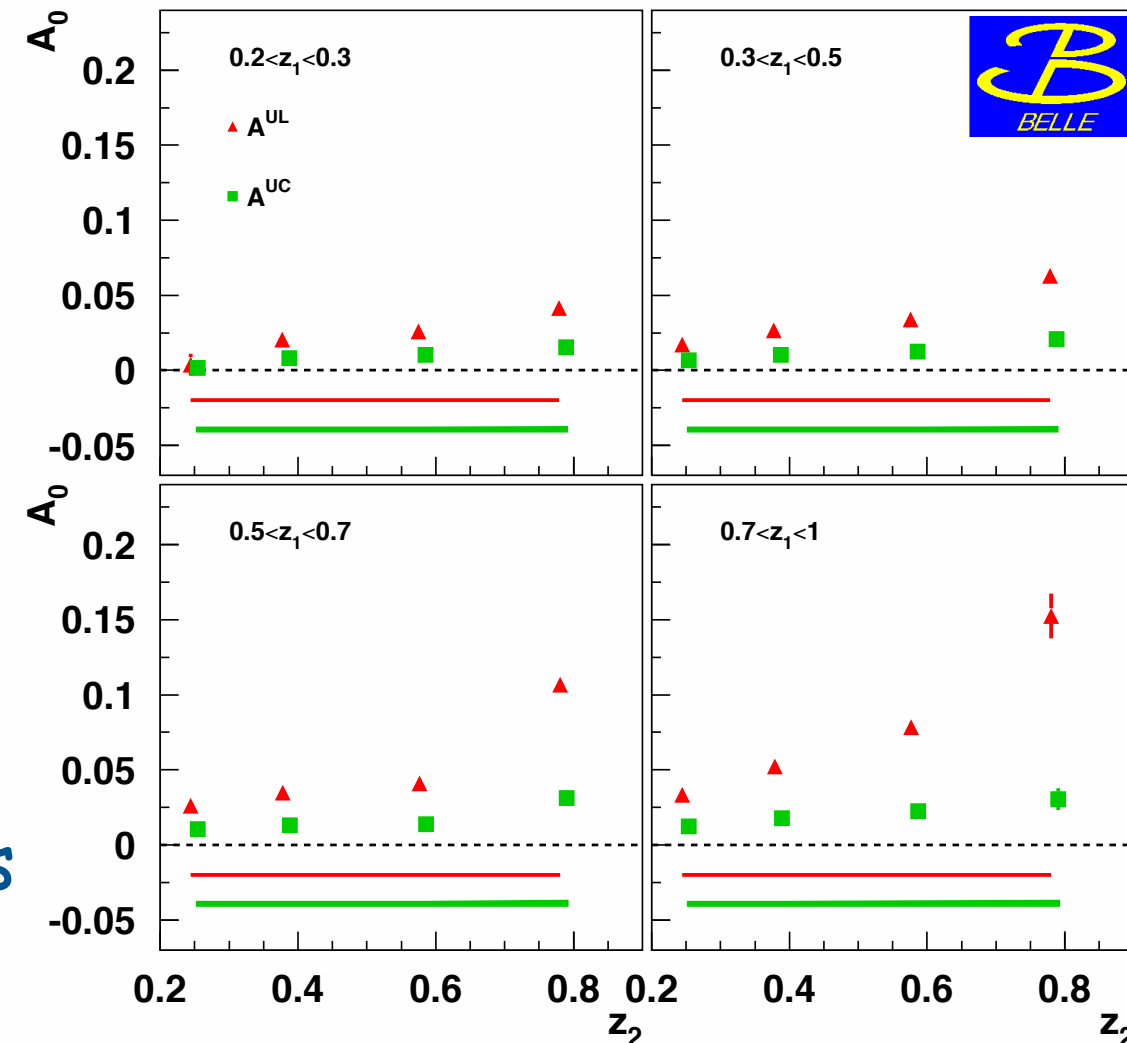
$$\begin{aligned} \frac{R_{12}^U}{R_{12}^L} &\simeq \frac{1 + \left\langle \frac{\sin^2 \theta_{\text{th}}}{1 + \cos^2 \theta_{\text{th}}} \right\rangle G^U \cos(\phi_1 + \phi_2)}{1 + \left\langle \frac{\sin^2 \theta_{\text{th}}}{1 + \cos^2 \theta_{\text{th}}} \right\rangle G^L \cos(\phi_1 + \phi_2)} \\ &\simeq 1 + \left\langle \frac{\sin^2 \theta_{\text{th}}}{1 + \cos^2 \theta_{\text{th}}} \right\rangle \{G^U - G^L\} \cos(\phi_1 + \phi_2) \end{aligned}$$

- suppresses flavor-independent sources of modulations
- $G^{U/L}$ specific combinations of FFs
- remaining MC asym.'s: systematics



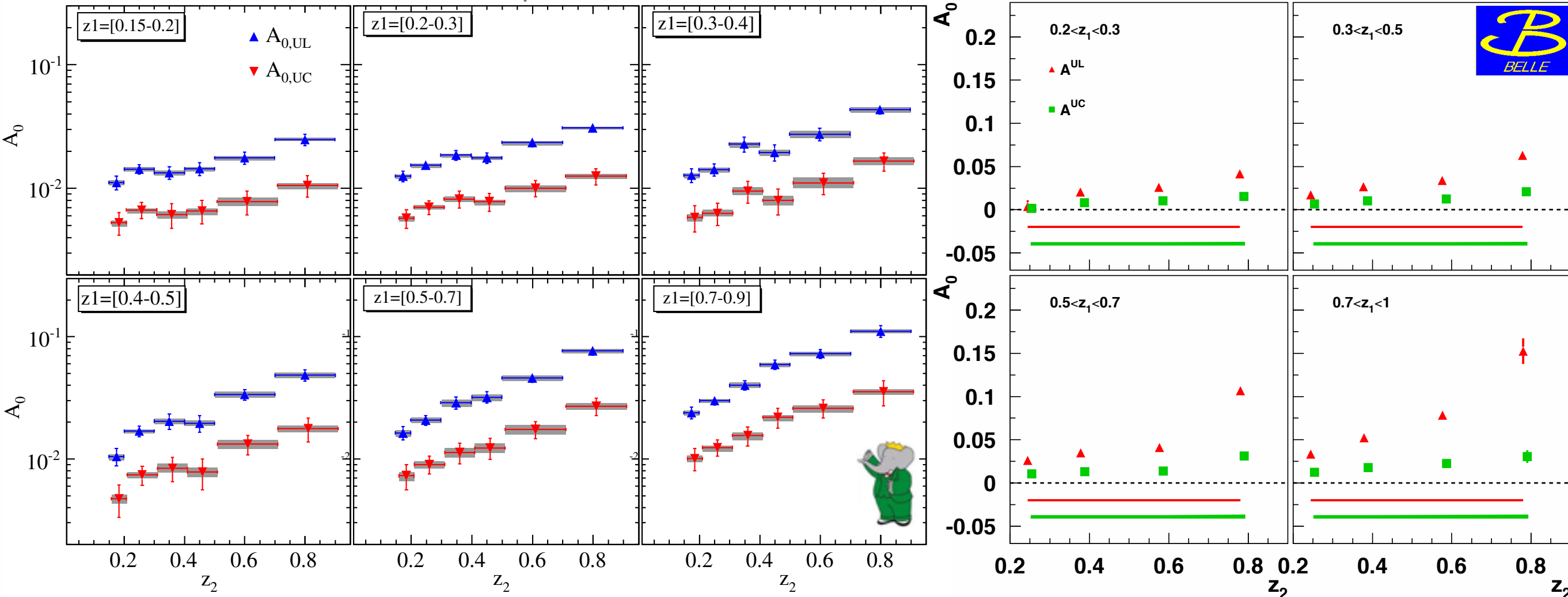
Collins asymmetries (RF0)

- first measurement of Collins asymmetries by Belle [PRL 96 (2006) 232002, PRD 78 (2008) 032011, PRD 86 (2012) 039905(E)]
- significant asymmetries rising with z
- used for first transversity and Collins FF extractions



Collins asymmetries (RF0)

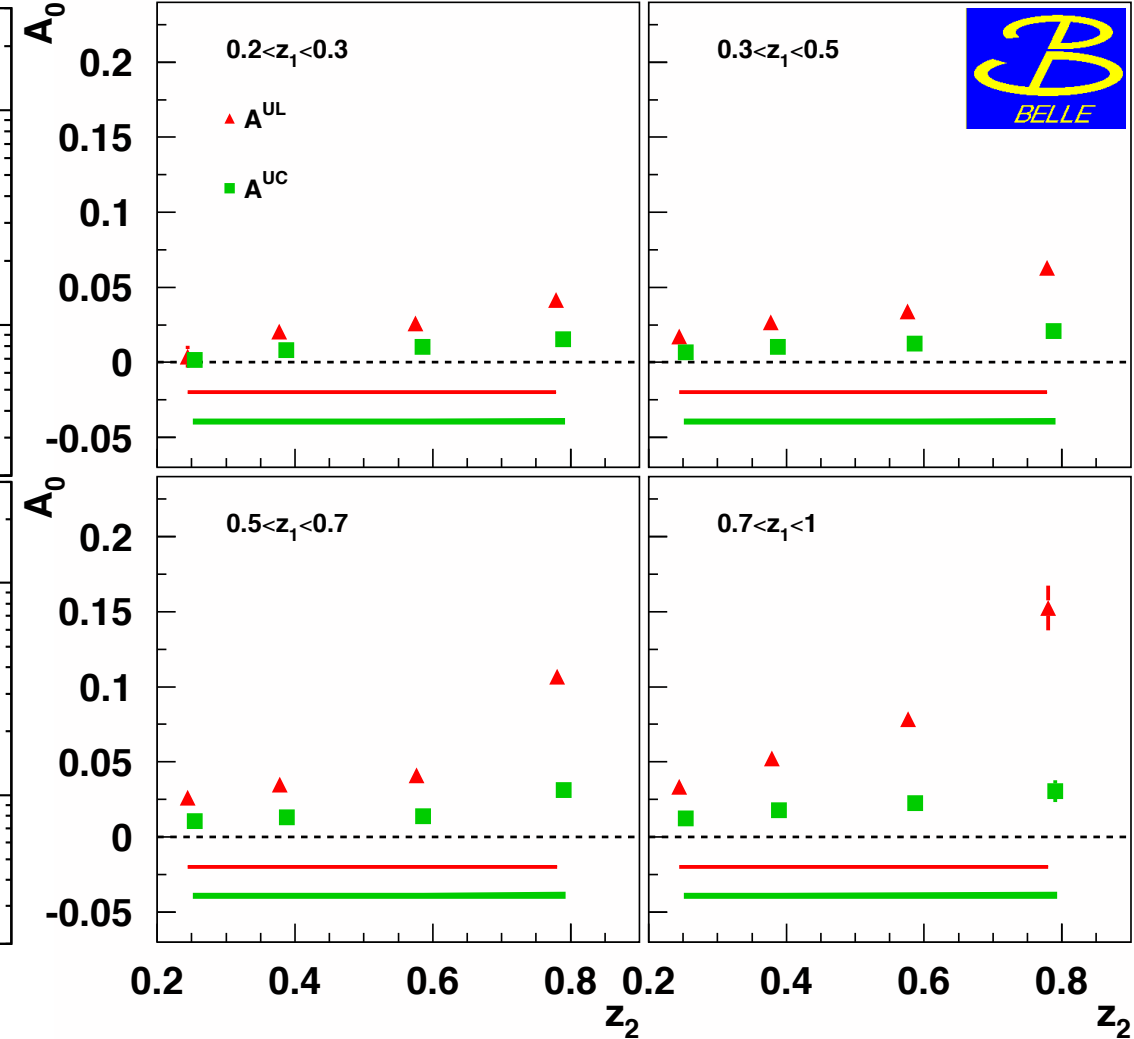
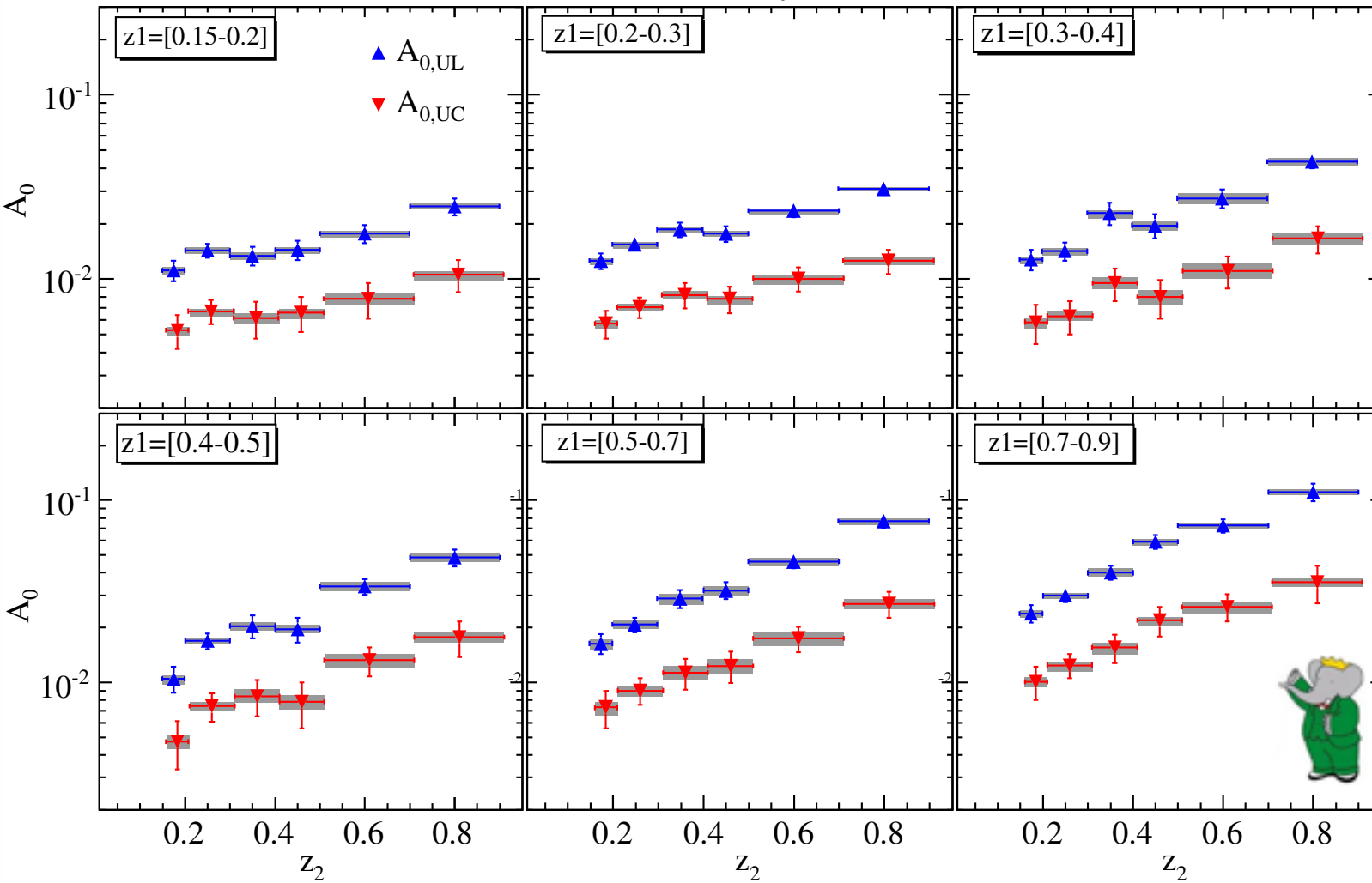
[Phys. Rev. D90 (2014) 052003]



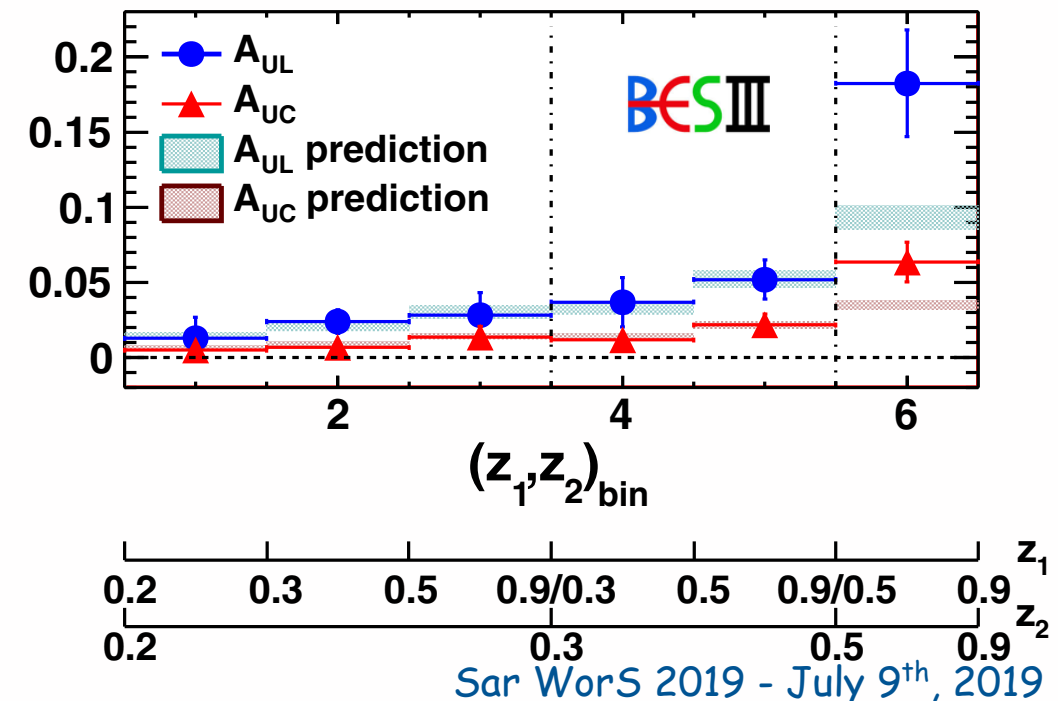
- BaBar results [PRD 90 (2014) 052003]
consistent with Belle

Collins asymmetries (RF0)

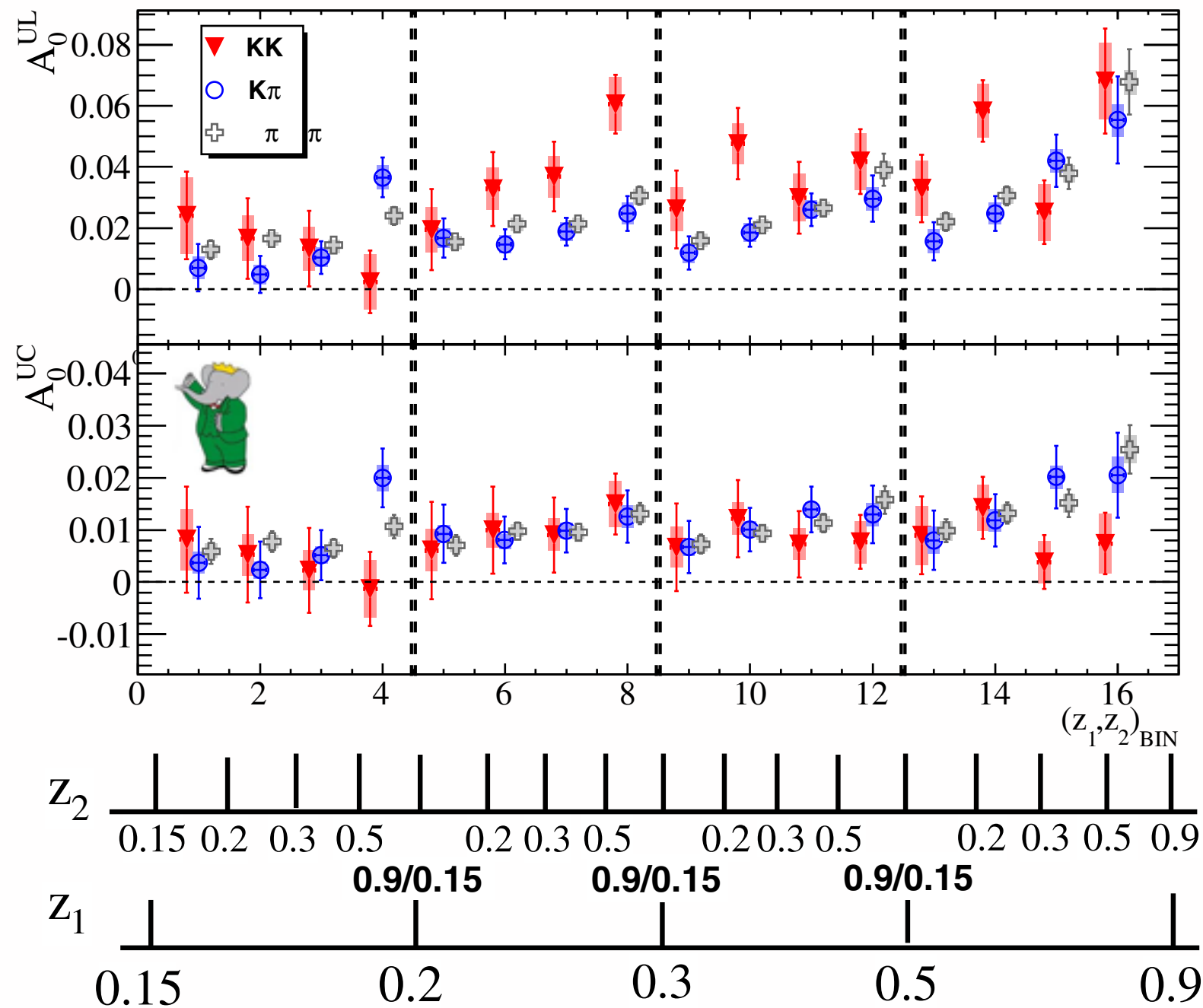
[Phys. Rev. D90 (2014) 052003]



- BaBar results [PRD 90 (2014) 052003] consistent with Belle
- BESIII [PRL 116 (2016) 042001] (at smaller s) consistent with TMD evolution [Z.-B. Kang et al., PRD 93 (2016) 014009]

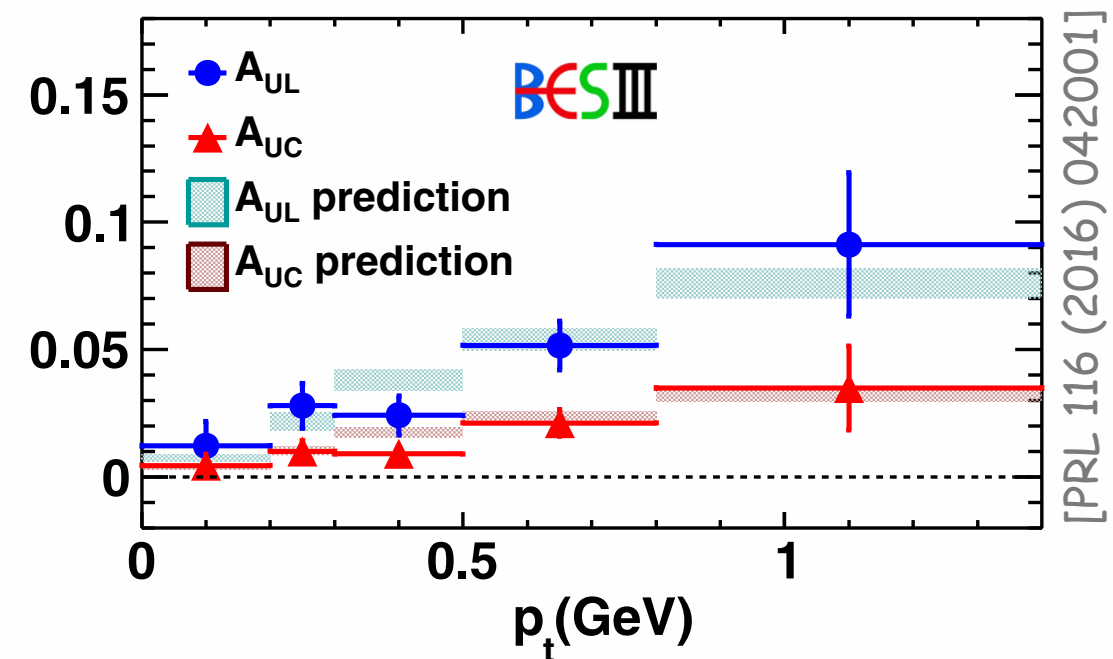
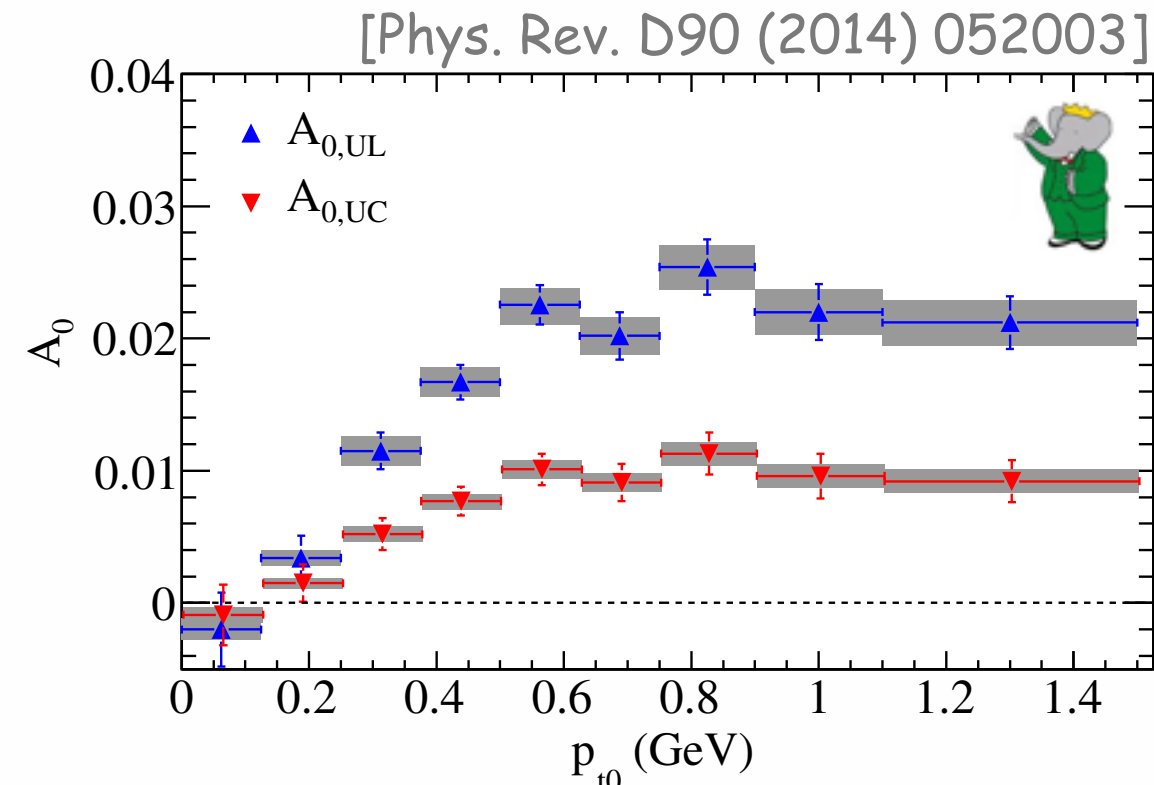
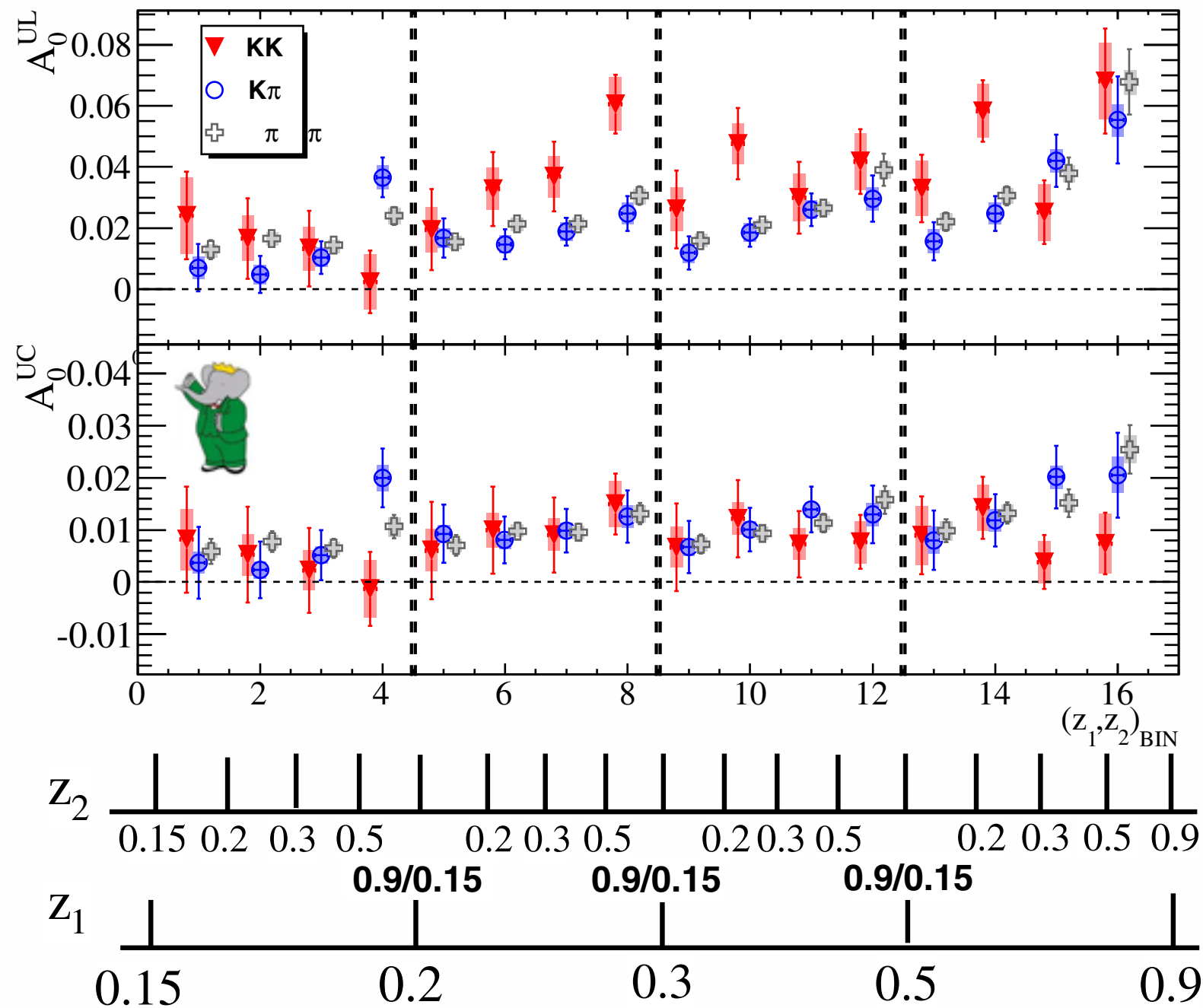


Collins asymmetries - going further



● even larger effects seen for kaon pairs

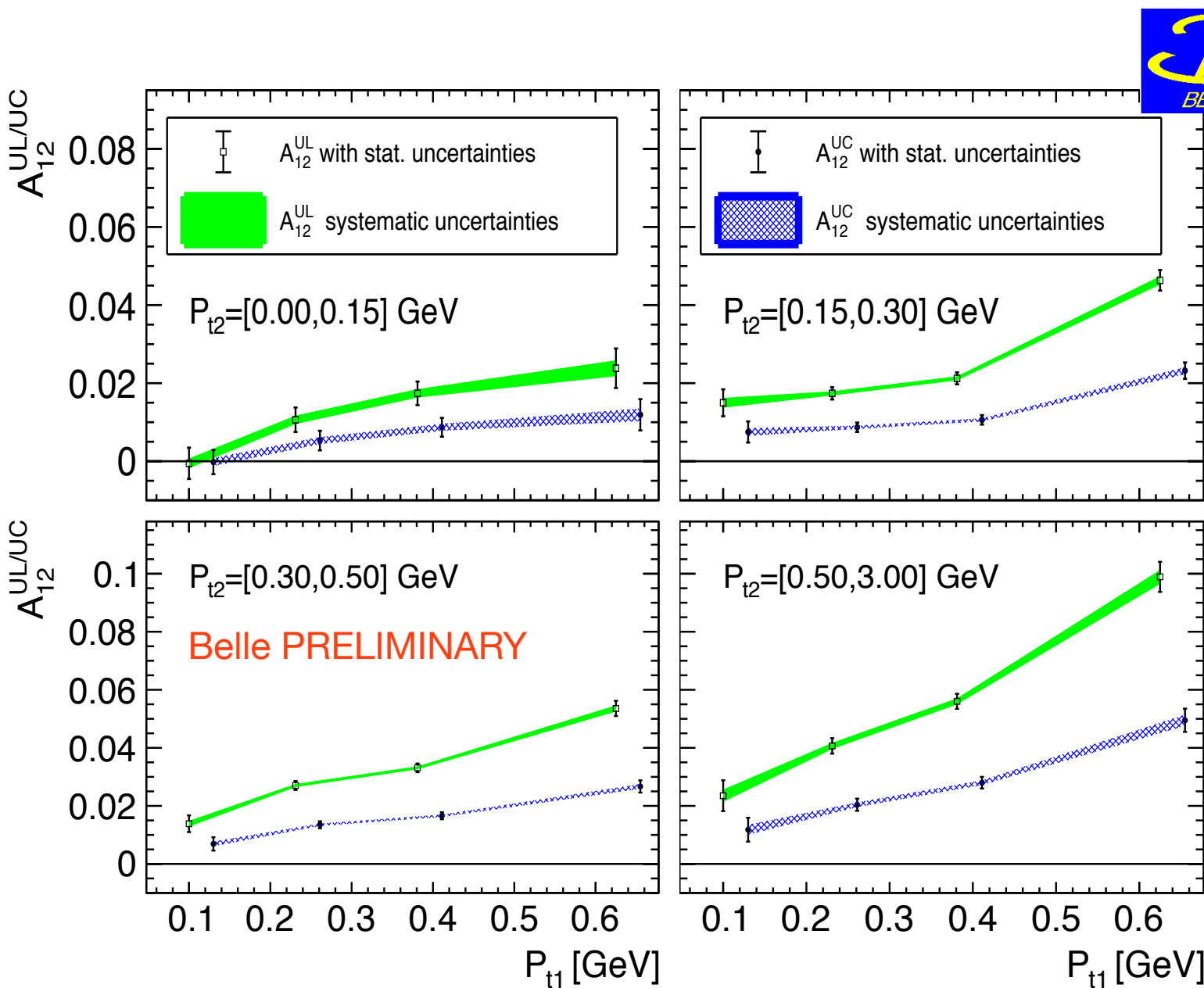
Collins asymmetries - going further



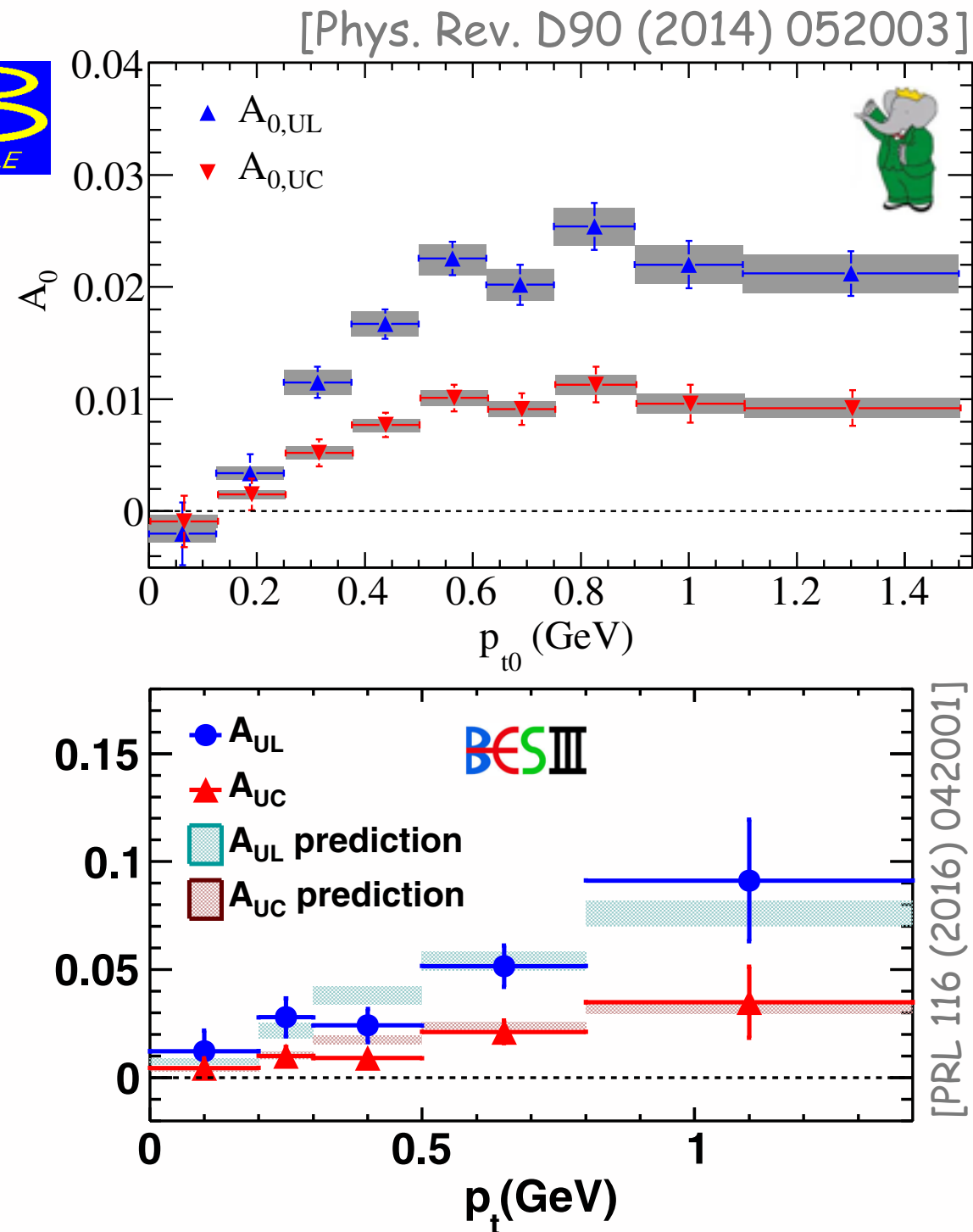
● even larger effects seen for kaon pairs

● p_T dependence for pions

Collins asymmetries - going further

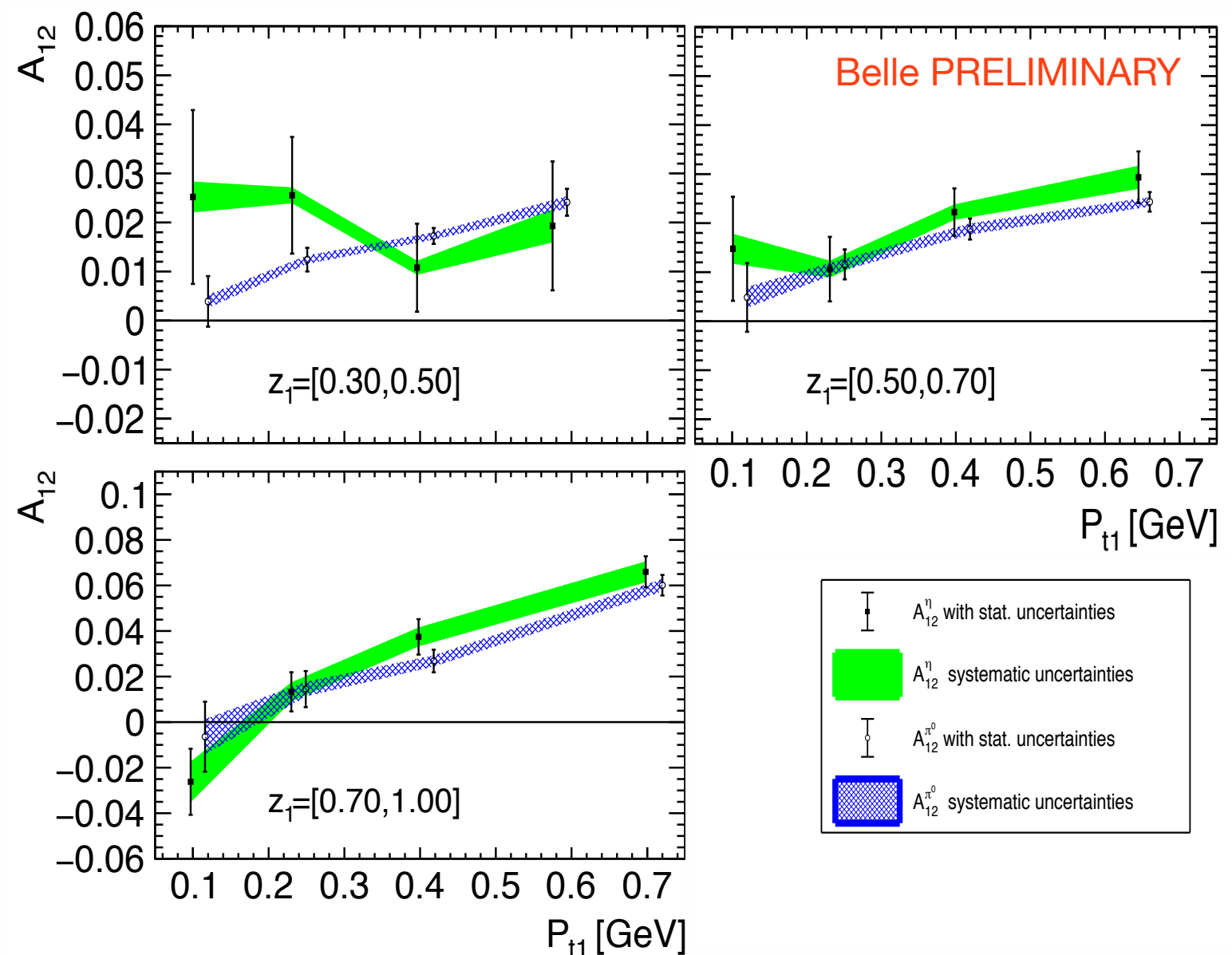
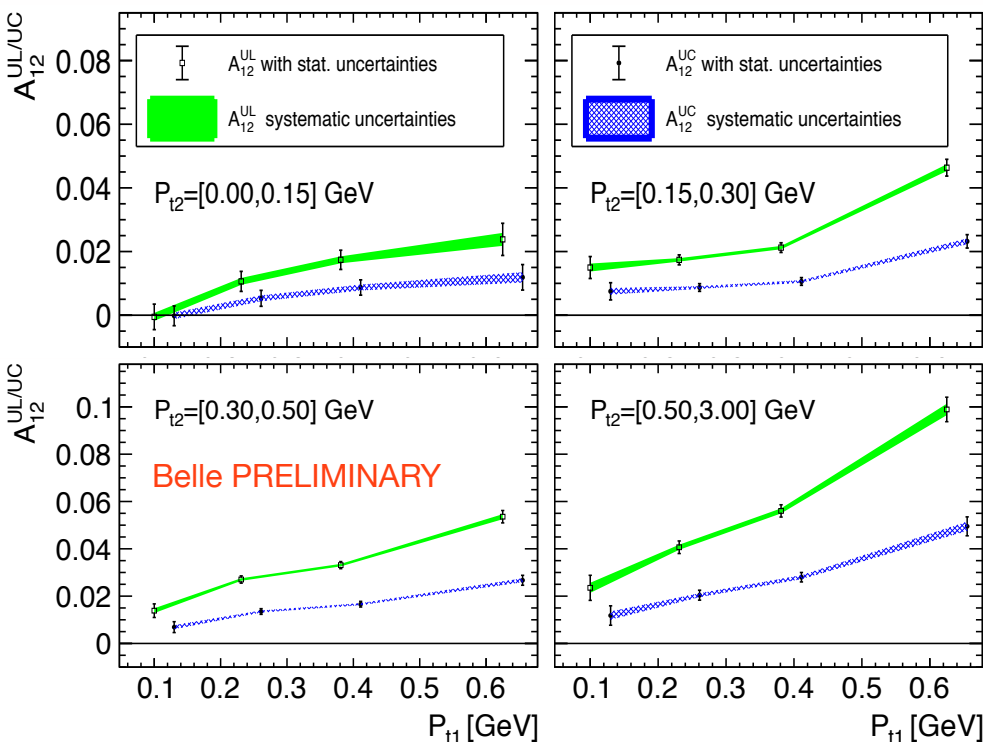


- now also from Belle for charged pions



- p_T dependence for pions

Collins asymmetries - going further



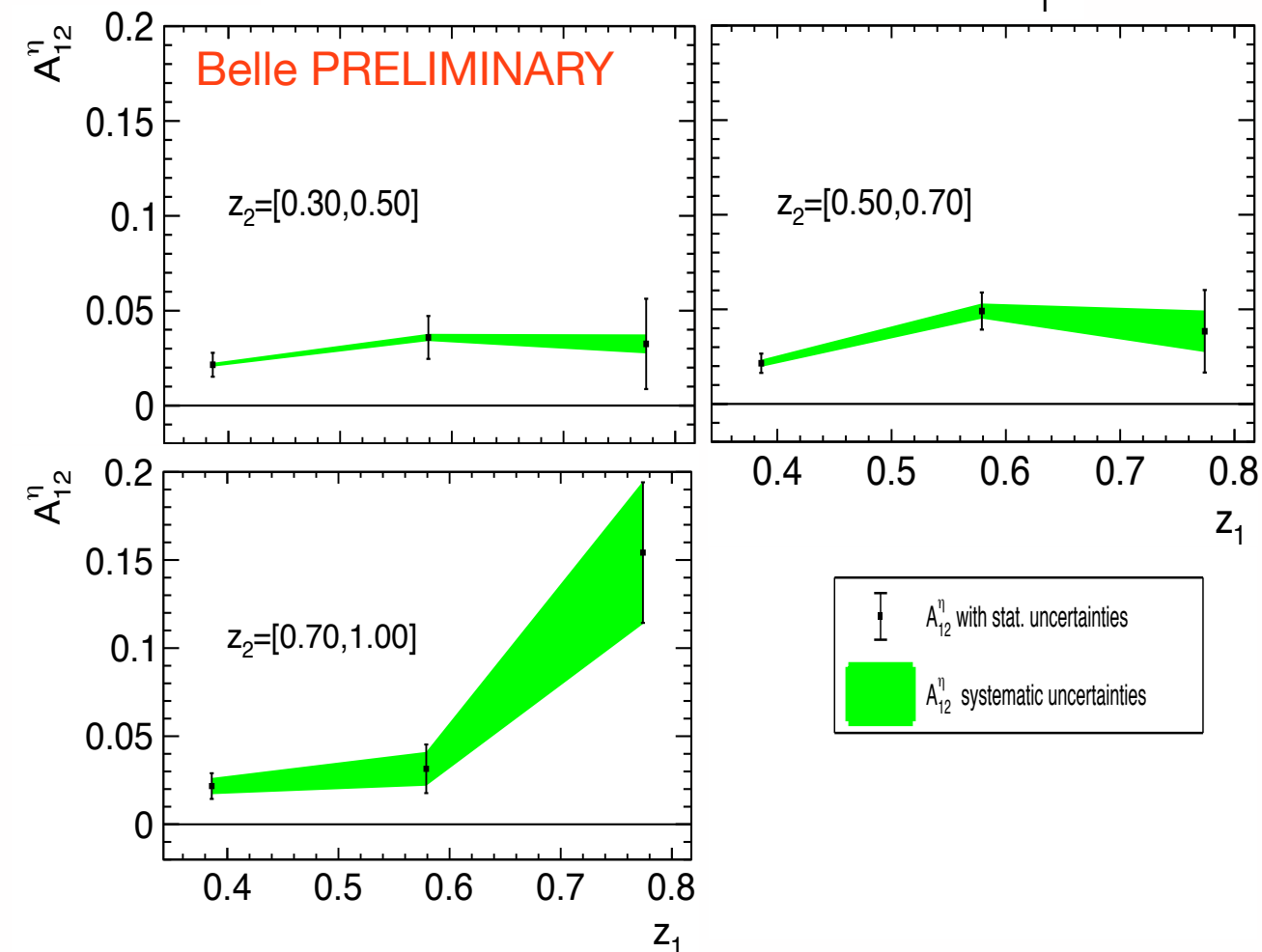
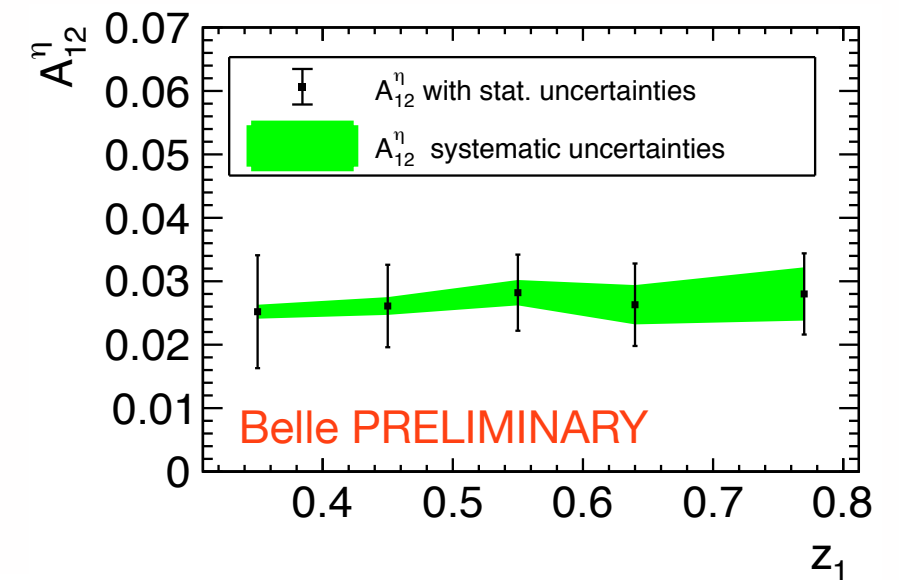
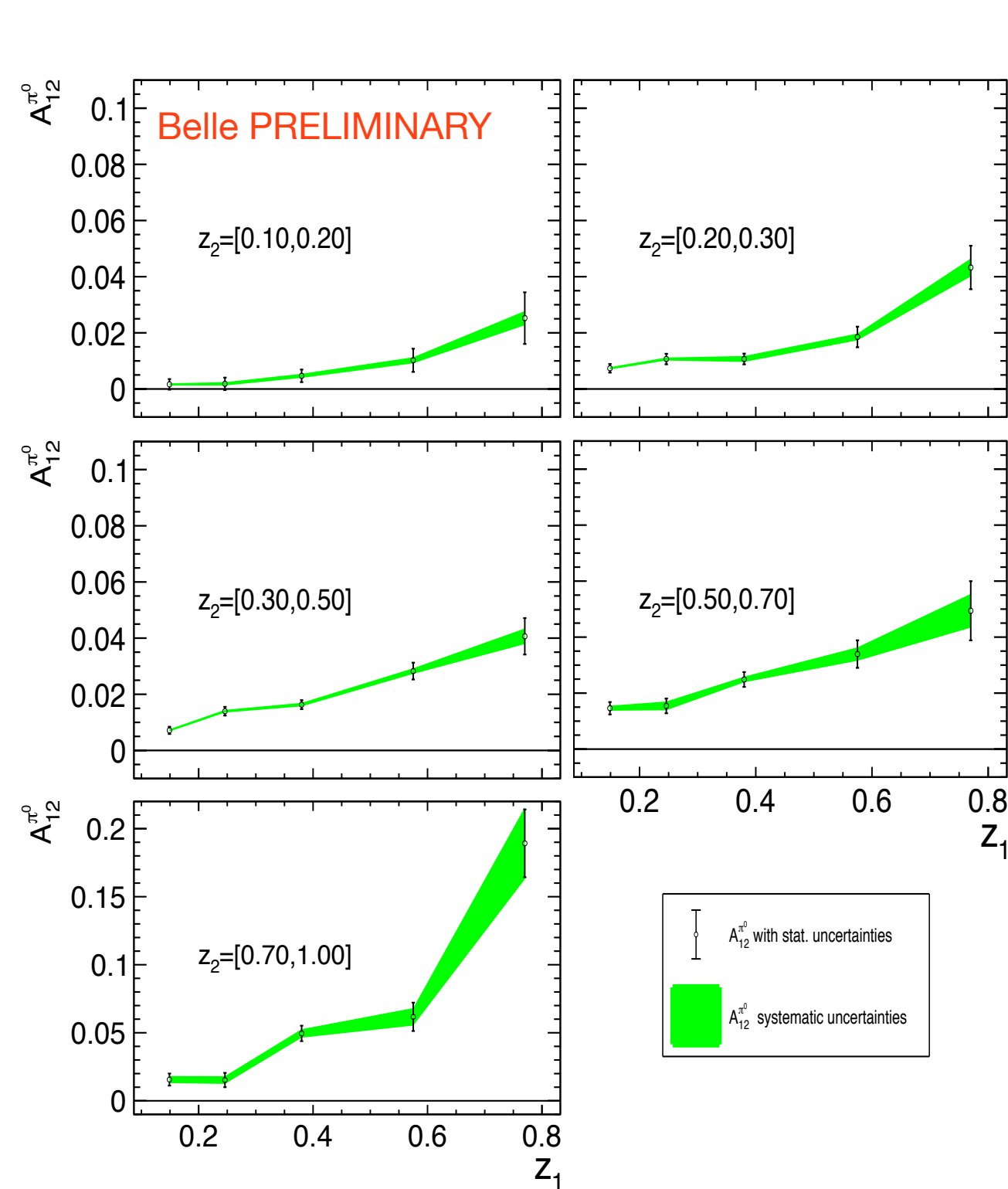
- now also from Belle for charged pions

$$R_{12}^{\pi^0} = \frac{R_{12}^{0\pm}}{R_{12}^L} = \frac{\pi^0\pi^+ + \pi^0\pi^-}{\pi^+\pi^+ + \pi^-\pi^-};$$

$$R_{12}^{\eta} = \frac{R_{12}^{\eta\pm}}{R_{12}^L} = \frac{\eta\pi^+ + \eta\pi^-}{\pi^+\pi^+ + \pi^-\pi^-}.$$

- ... and neutral pion and eta

Collins asymmetries - going further



- typical rise with z seen for neutral pion
- eta almost flat except for large- z region

Collins asymmetries - going further

$$R_{12}^{\pi^0} = \frac{R_{12}^{0\pm}}{R_{12}^L} \approx 1 + \cos(\phi_{12}) \frac{\sin^2(\theta)}{1 + \cos^2(\theta)}$$

$$\times \left\{ \frac{5(H_1^{\perp, fav} + H_1^{\perp, dis}) \otimes (H_1^{\perp, fav} + H_1^{\perp, dis}) + 4H_{1,s \rightarrow \pi}^{\perp, dis} \otimes H_{1,s \rightarrow \pi}^{\perp, dis}}{5(D_1^{fav} + D_1^{dis}) \otimes (D_1^{fav} + D_1^{dis}) + 4D_{1,s \rightarrow \pi}^{dis} \otimes D_{1,s \rightarrow \pi}^{dis}} \right.$$

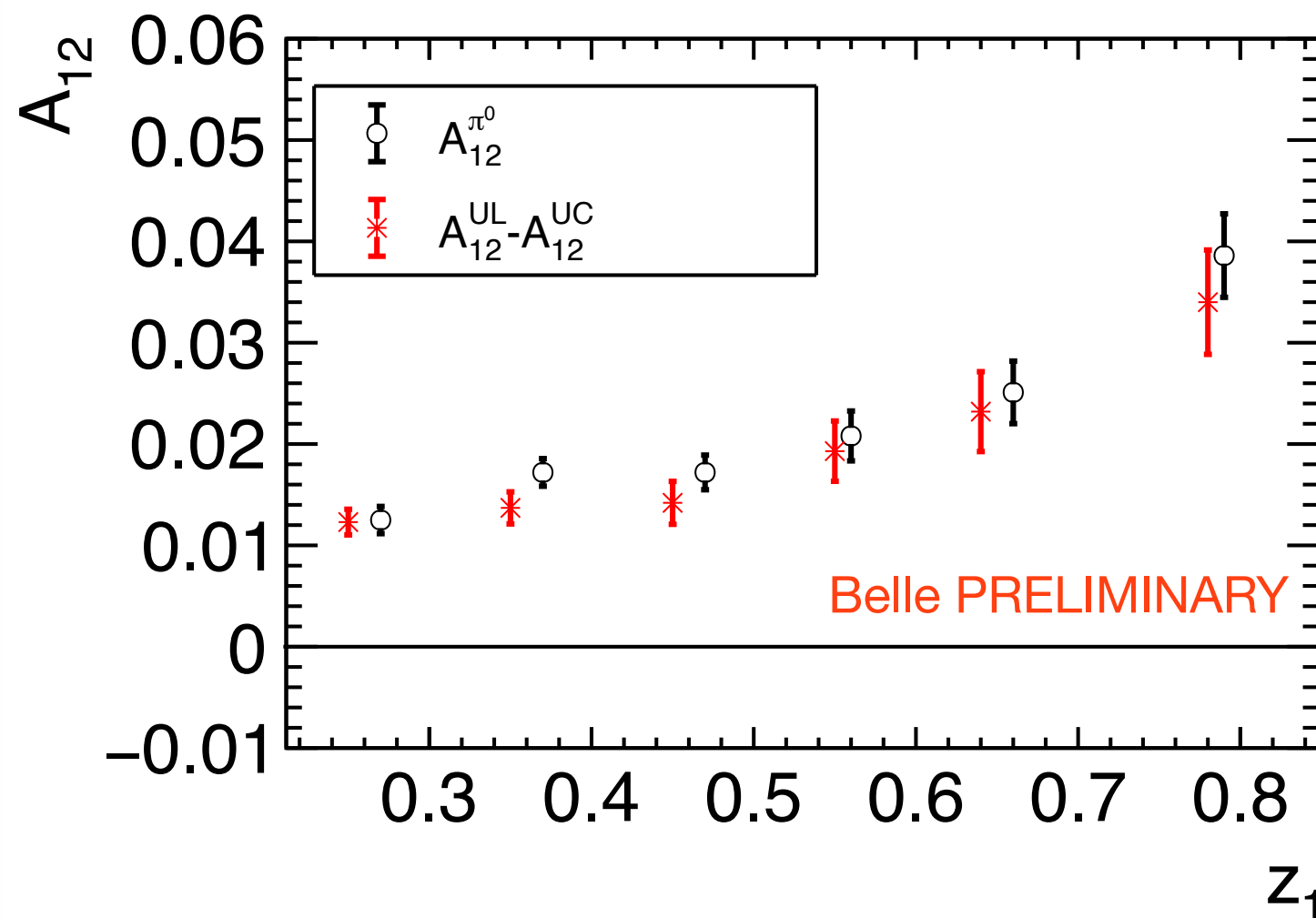
$$\left. - \frac{5(H_1^{\perp, fav} \otimes H_1^{\perp, dis} + H_1^{\perp, dis} \otimes H_1^{\perp, fav}) + 2H_{1,s \rightarrow \pi}^{\perp, dis} H_{1,s \rightarrow \pi}^{\perp, dis}}{5(D_1^{fav} \otimes D_1^{dis} + D_1^{dis} \otimes D_1^{fav}) + 2D_{1,s \rightarrow \pi}^{dis} \otimes D_{1,s \rightarrow \pi}^{dis}} \right\} \quad \left. \vphantom{\frac{5(H_1^{\perp, fav} + H_1^{\perp, dis}) \otimes (H_1^{\perp, fav} + H_1^{\perp, dis}) + 4H_{1,s \rightarrow \pi}^{\perp, dis} \otimes H_{1,s \rightarrow \pi}^{\perp, dis}}{5(D_1^{fav} + D_1^{dis}) \otimes (D_1^{fav} + D_1^{dis}) + 4D_{1,s \rightarrow \pi}^{dis} \otimes D_{1,s \rightarrow \pi}^{dis}}} \right\} \stackrel{\text{isospin}}{=} A_{12}^{UL} - A_{12}^{UC}$$

Collins asymmetries - going further

$$R_{12}^{\pi^0} = \frac{R_{12}^{0\pm}}{R_{12}^L} \approx 1 + \cos(\phi_{12}) \frac{\sin^2(\theta)}{1 + \cos^2(\theta)}$$

$$\times \left\{ \frac{5(H_1^{\perp, fav} + H_1^{\perp, dis}) \otimes (H_1^{\perp, fav} + H_1^{\perp, dis}) + 4H_{1,s \rightarrow \pi}^{\perp, dis} \otimes H_{1,s \rightarrow \pi}^{\perp, dis}}{5(D_1^{fav} + D_1^{dis}) \otimes (D_1^{fav} + D_1^{dis}) + 4D_{1,s \rightarrow \pi}^{dis} \otimes D_{1,s \rightarrow \pi}^{dis}} - \frac{5(H_1^{\perp, fav} \otimes H_1^{\perp, dis} + H_1^{\perp, dis} \otimes H_1^{\perp, fav}) + 2H_{1,s \rightarrow \pi}^{\perp, dis} H_{1,s \rightarrow \pi}^{\perp, dis}}{5(D_1^{fav} \otimes D_1^{dis} + D_1^{dis} \otimes D_1^{fav}) + 2D_{1,s \rightarrow \pi}^{dis} \otimes D_{1,s \rightarrow \pi}^{dis}} \right\}$$

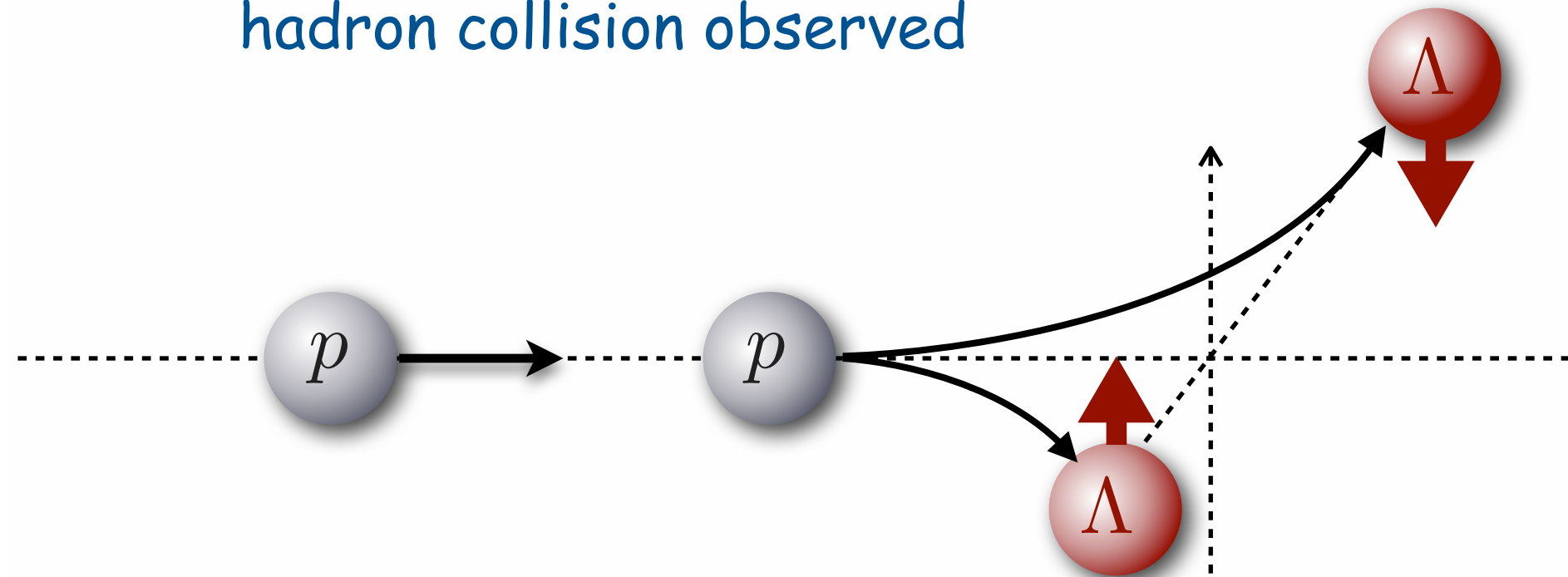
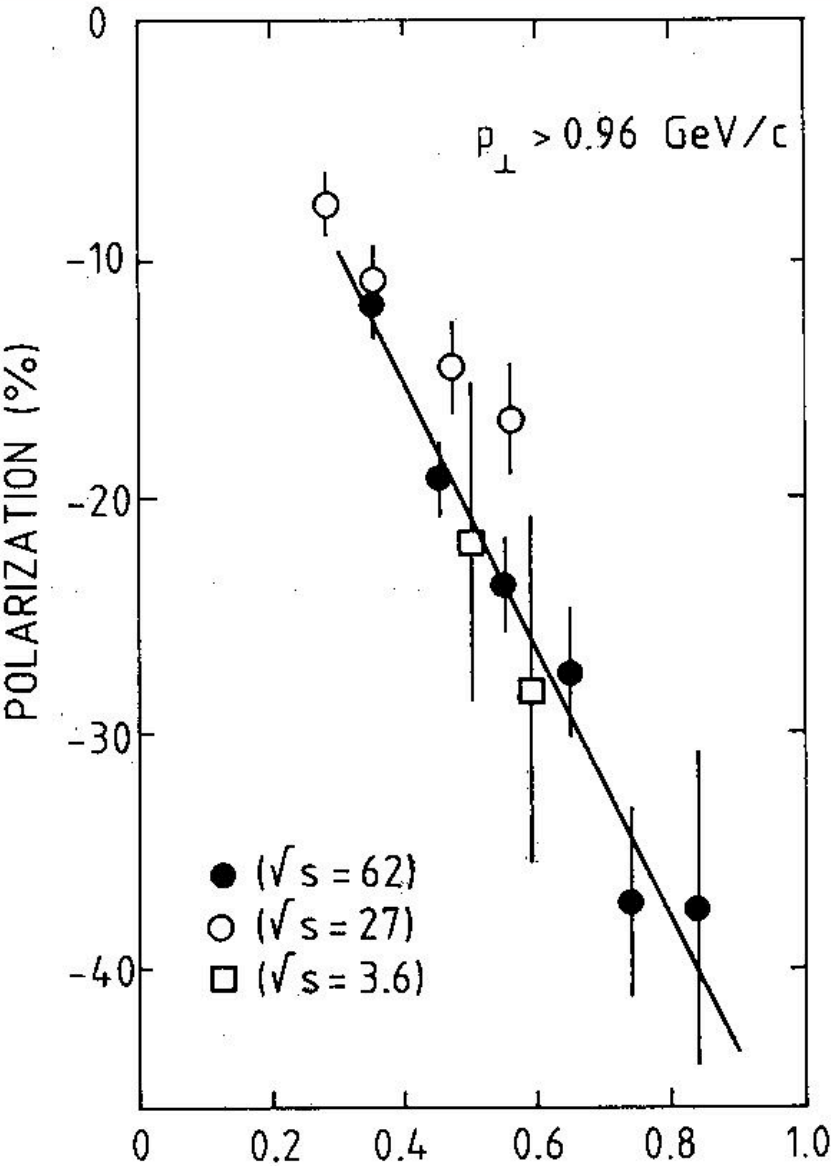
isospin $\underline{=}$ $A_{12}^{UL} - A_{12}^{UC}$



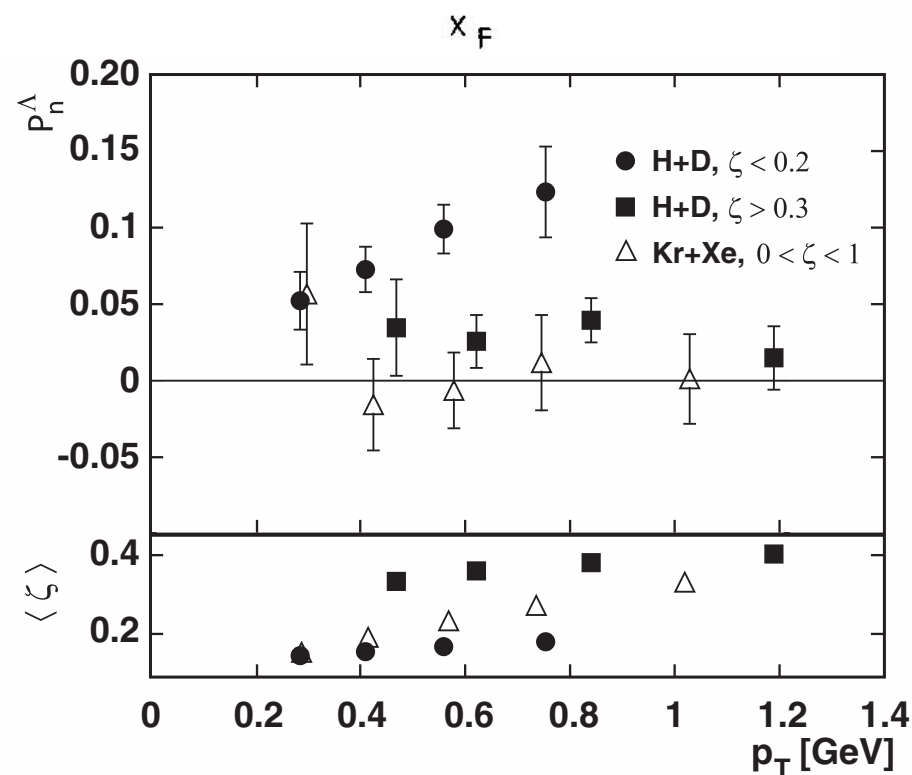
● consistency between neutral and charged pions

polarizing fragmentation

- large hyperon polarization in unpolarized hadron collision observed

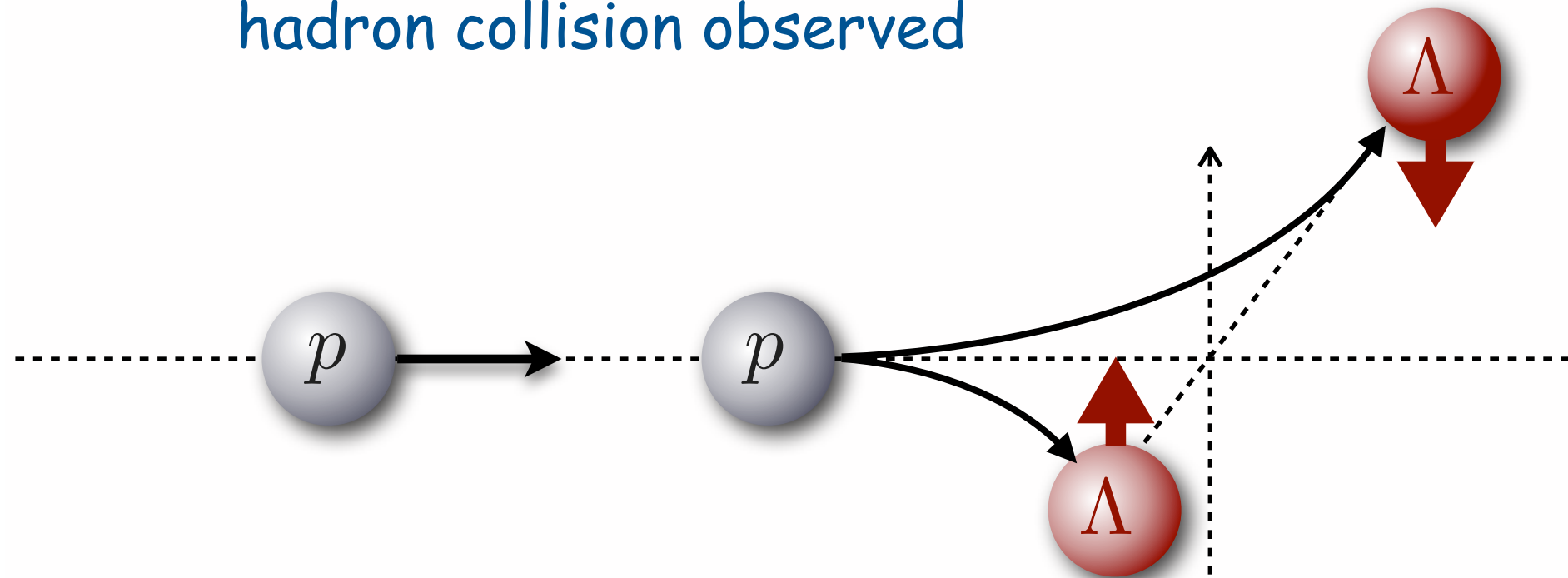
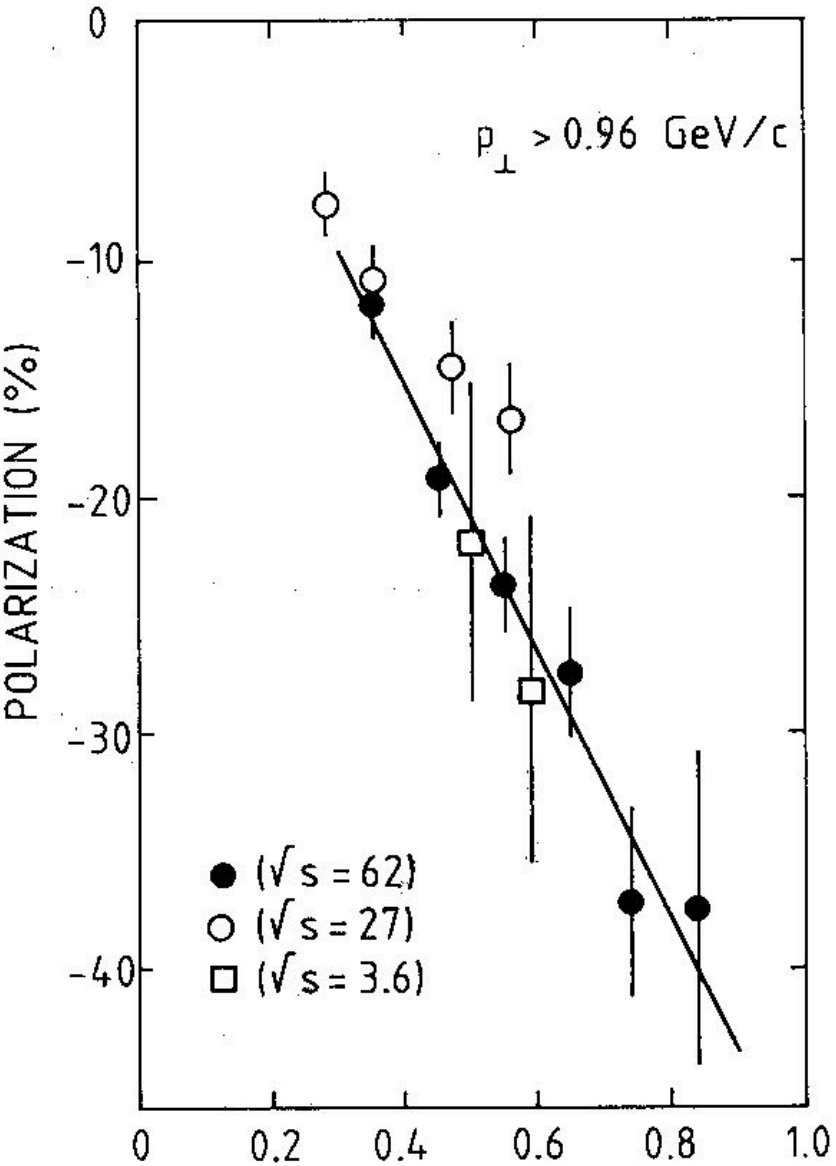


- ... as well as in inclusive lepto-production

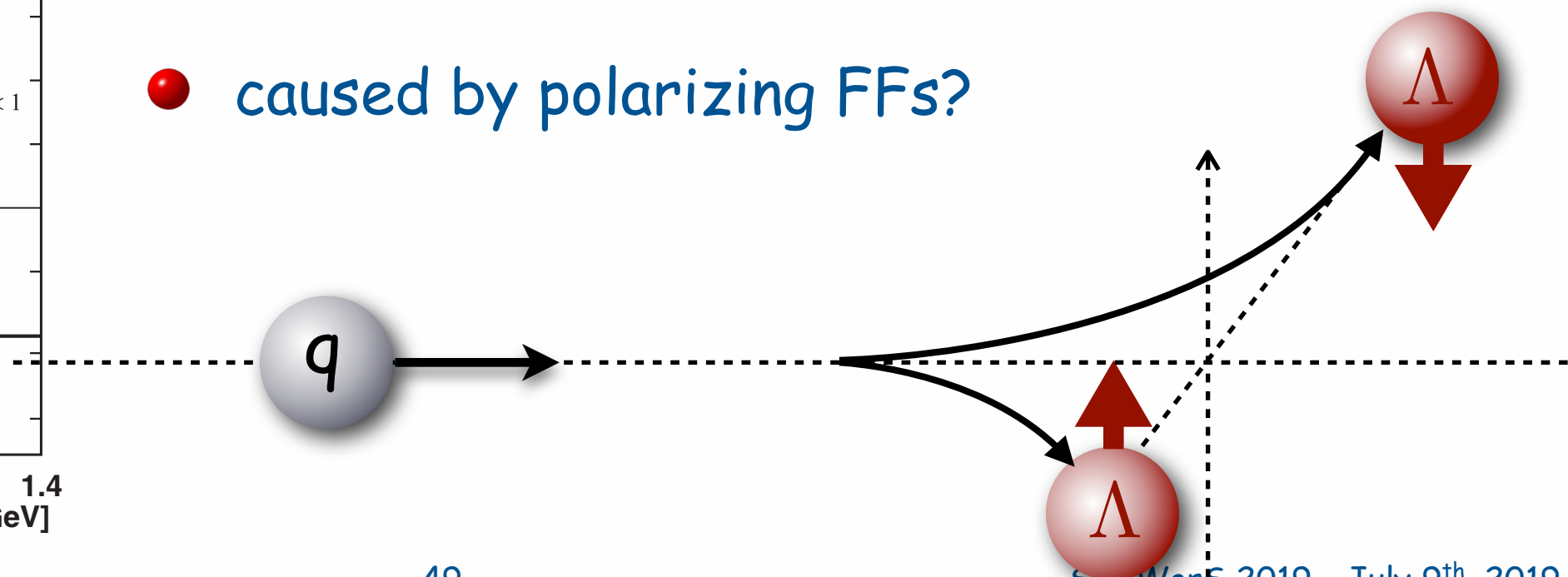
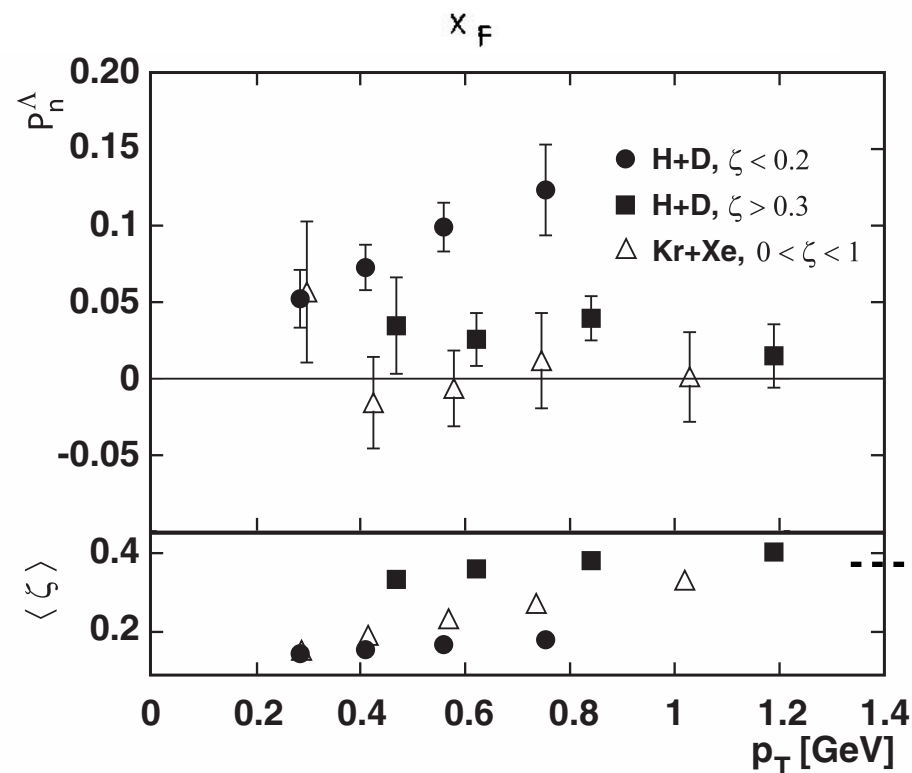


polarizing fragmentation

- large hyperon polarization in unpolarized hadron collision observed

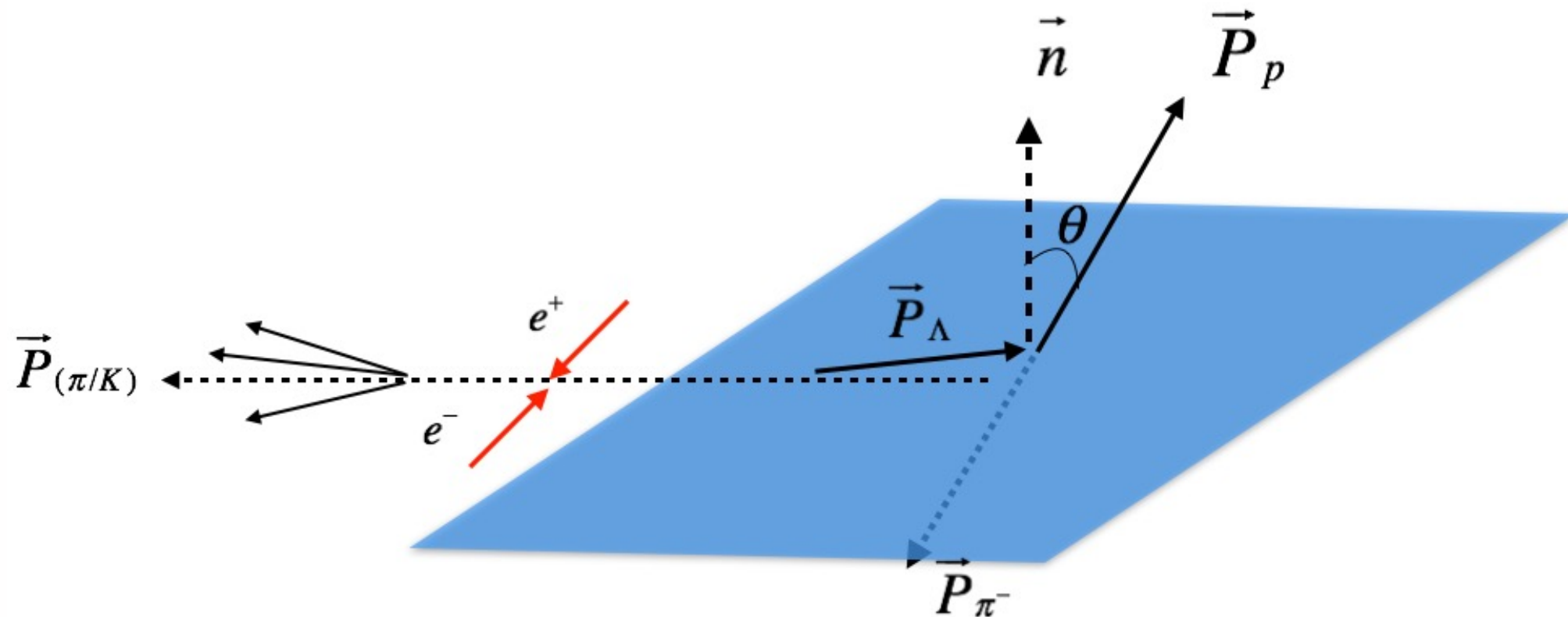


- ... as well as in inclusive lepto-production
- caused by polarizing FFs?



polarizing fragmentation function

- polarization measured normal to production plane, i.e. $\propto (\vec{P}_q \times \vec{P}_\Lambda)$

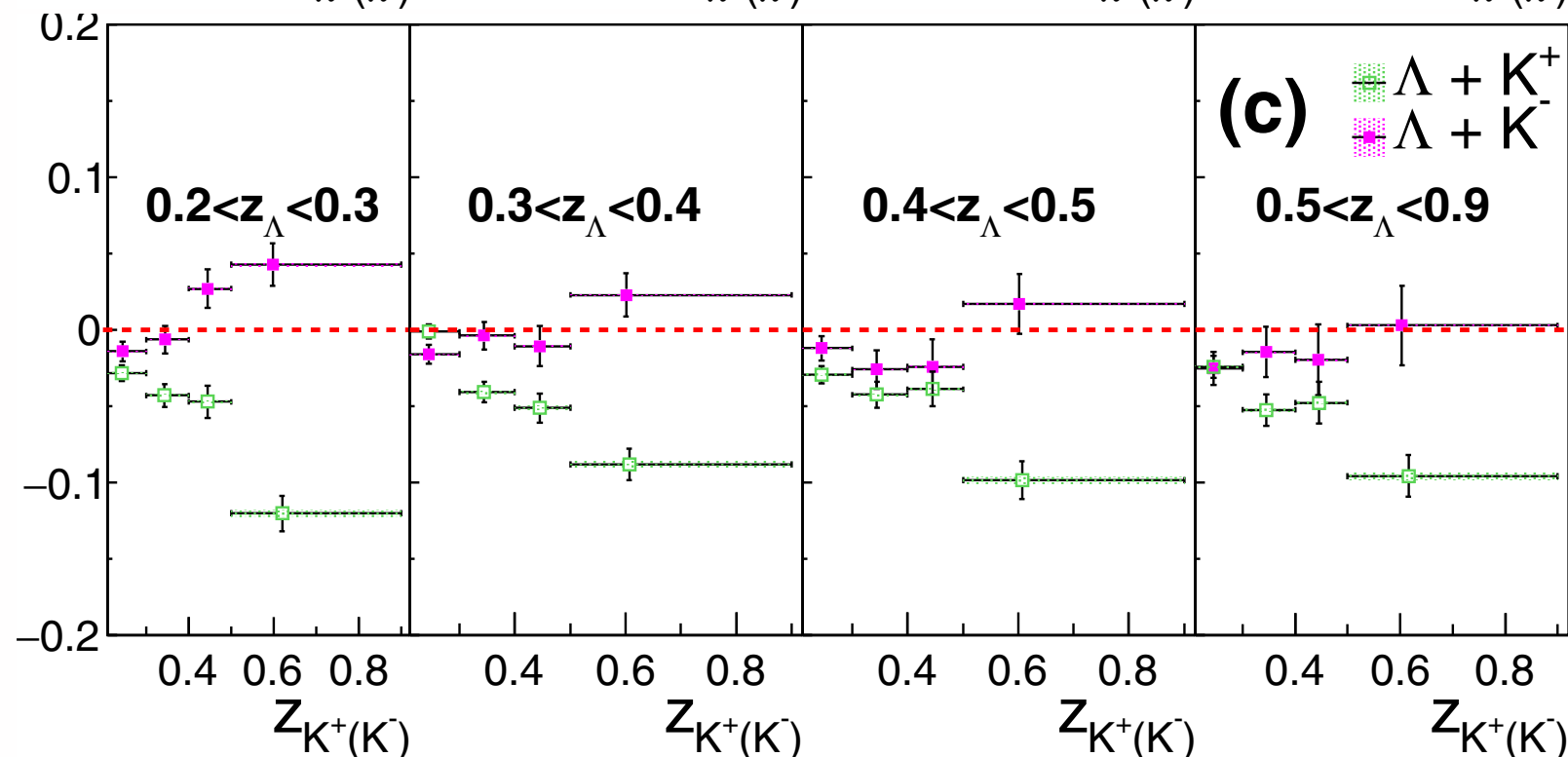
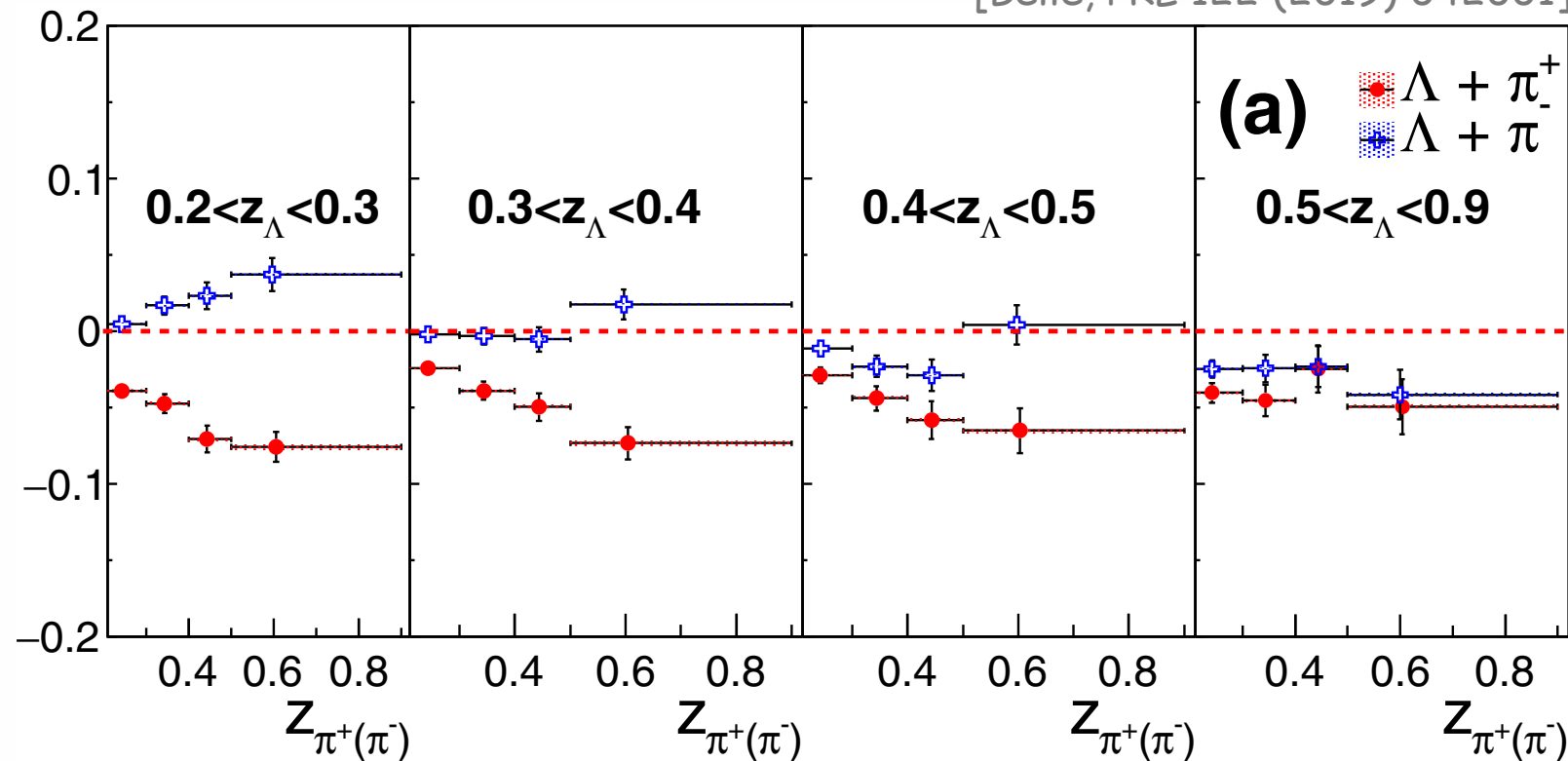


- reference axis to define transverse momentum:
 - "hadron frame" - use momentum direction of "back-to-back" hadron
 - "thrust frame" - use thrust axis
- use self-analyzing weak decay of Λ to determine polarization

polarizing fragmentation function

- flavor tagging through hadrons in opposite hemisphere:

[Belle, PRL 122 (2019) 042001]



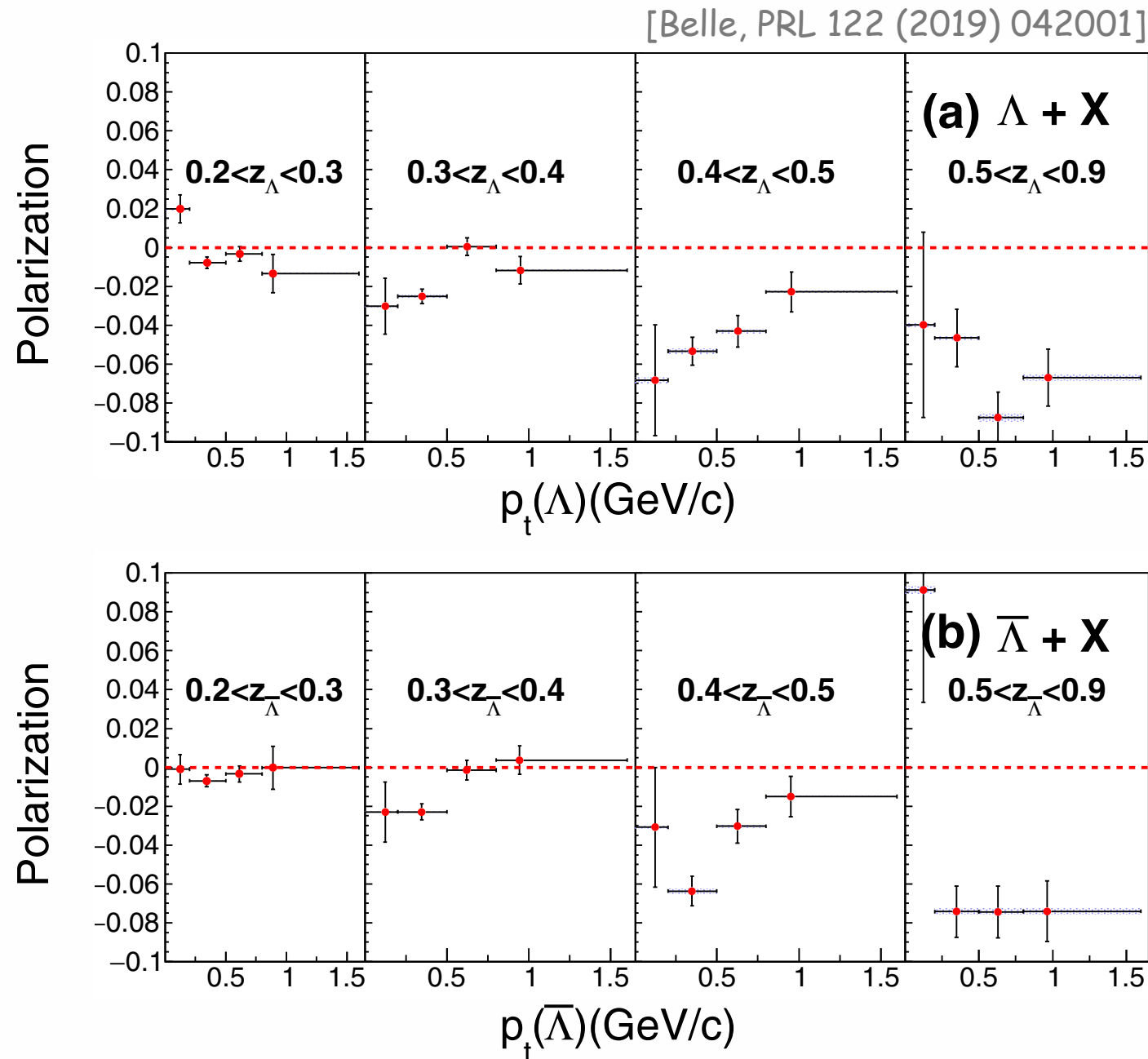
- large- z_h hadrons tag quark flavor more efficiently

➔ enlarges differences between oppositely charged hadrons

- MC-based quark-flavor decomposition in backup

polarizing fragmentation function

- polarization measured as function of z and p_t

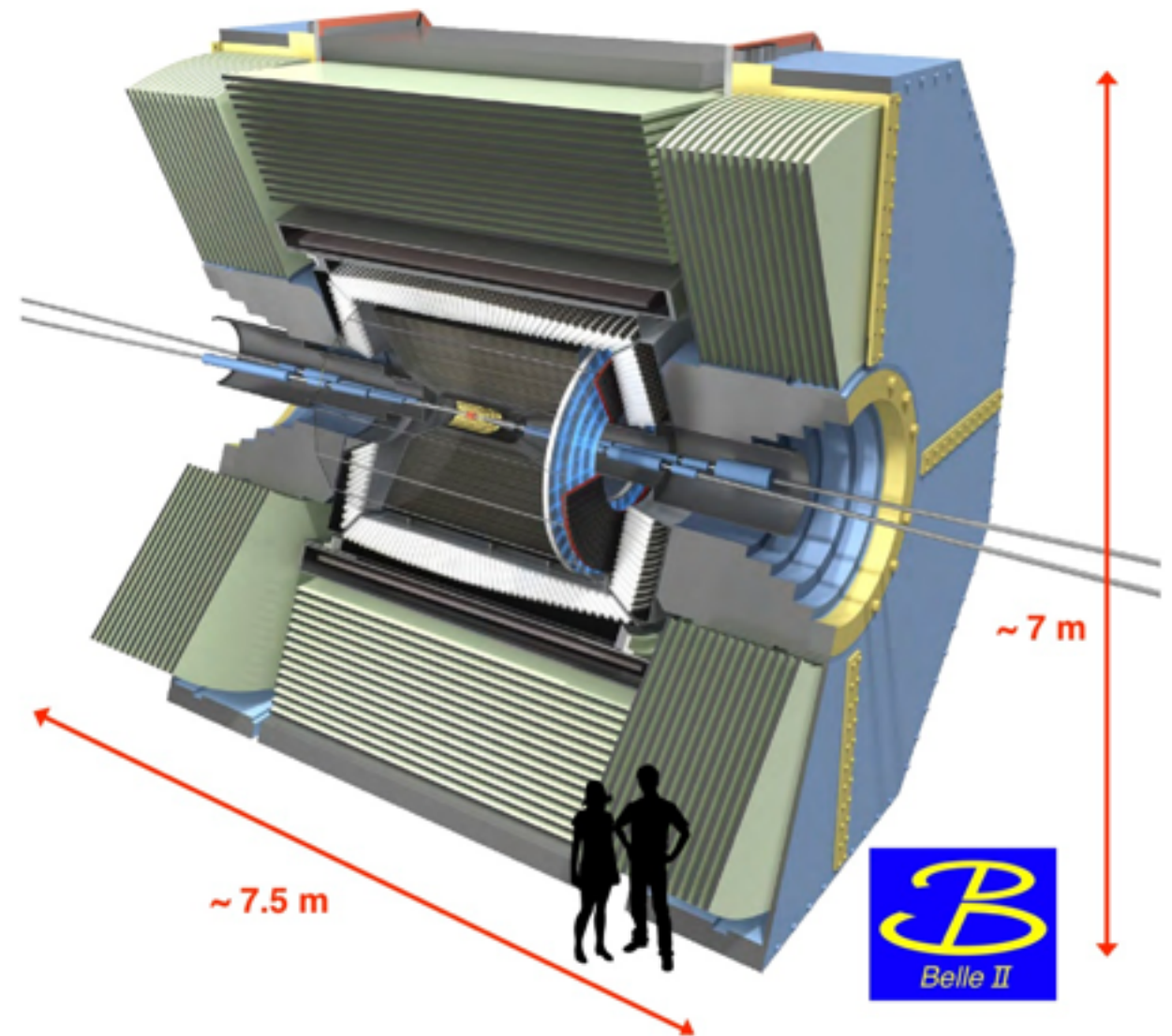
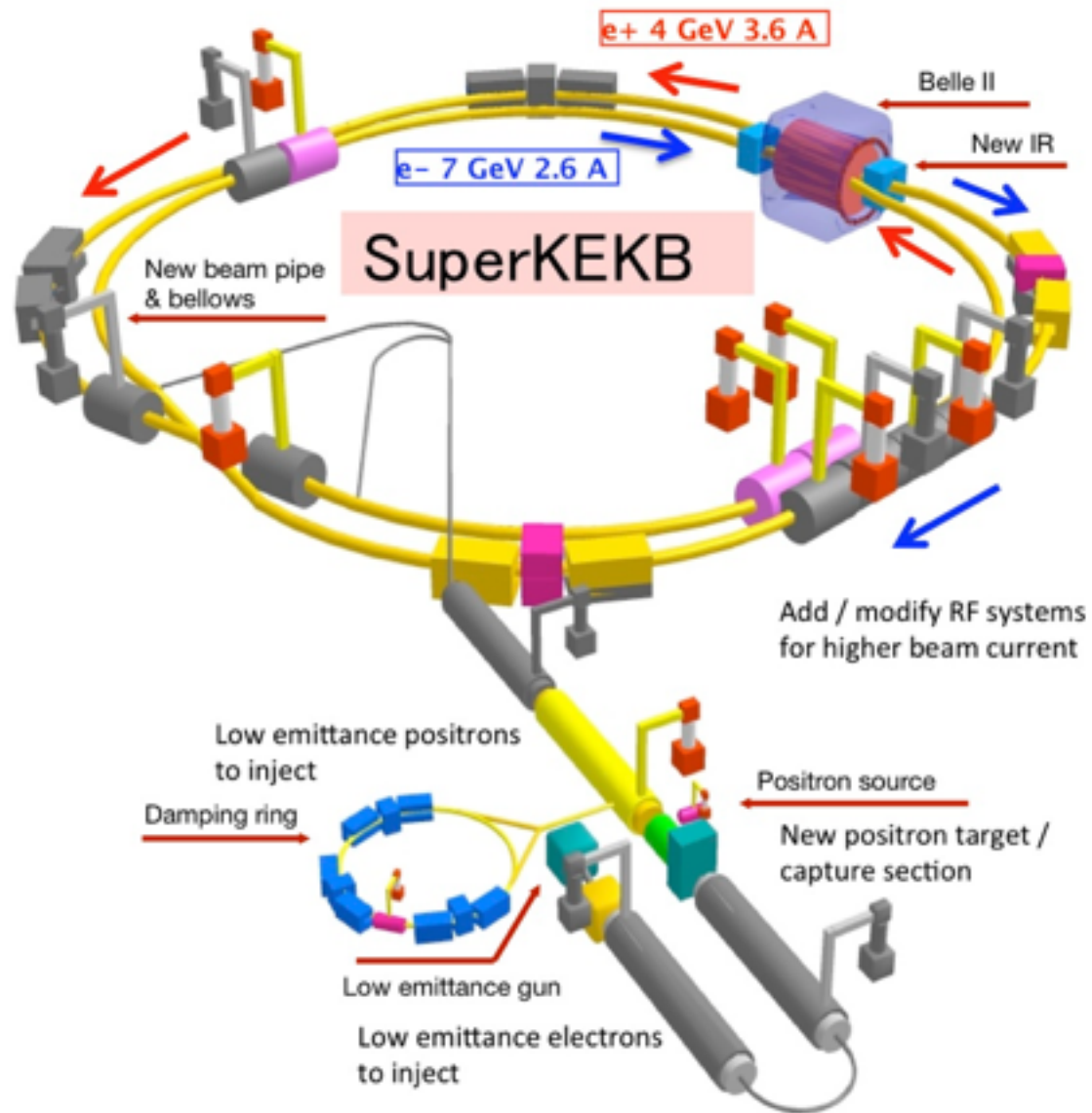


- strong dependence on both kinematics
- unexpected/surprising behavior for $p_t \rightarrow 0$

what to further expect (soon) from e^+e^-

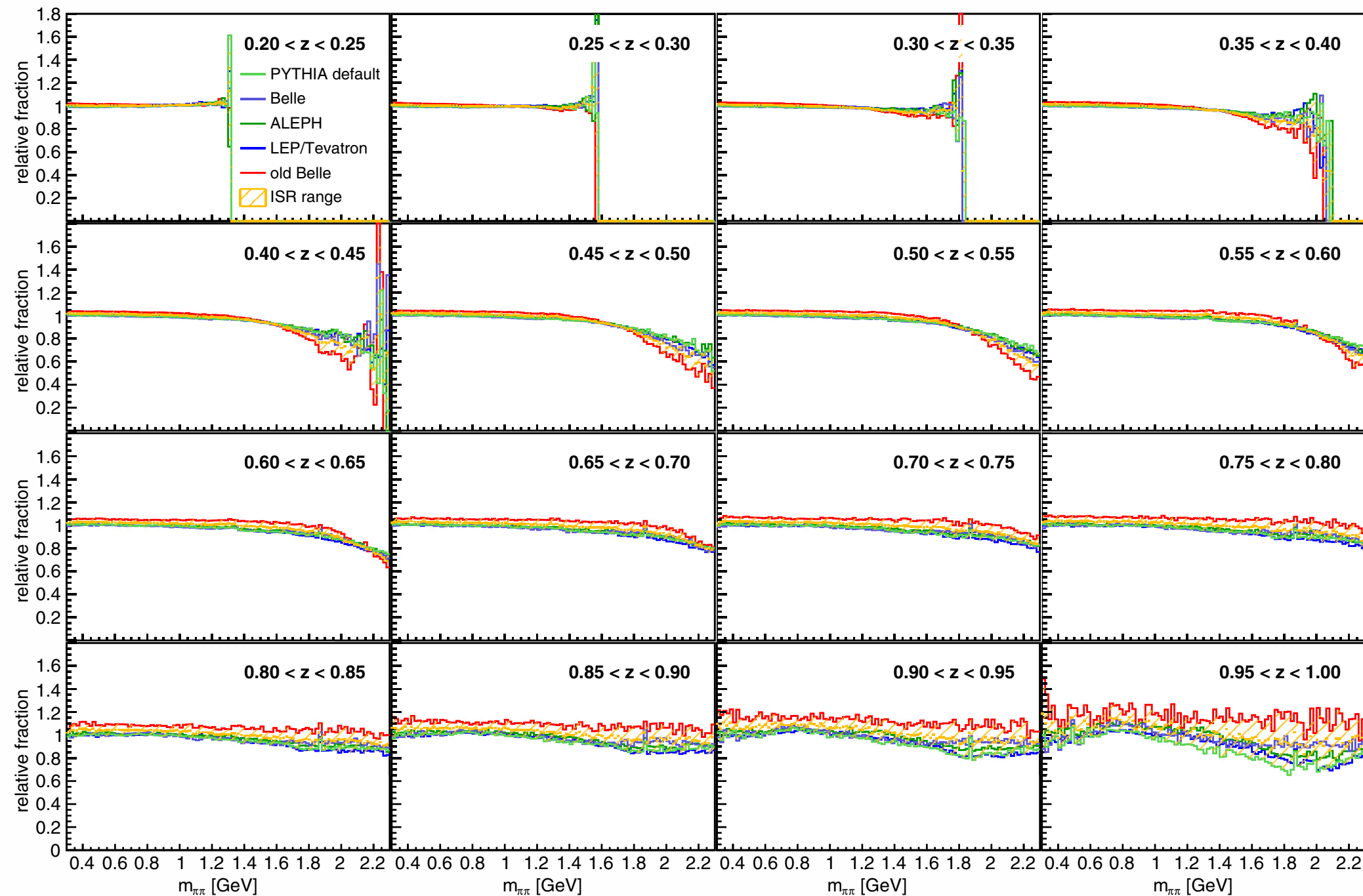
- single-hadron production
 - re-analysis with new ISR (Belle)
 - lower- s data (BESIII; Belle via ISR)
- k_T -dependent D_1 FFs (Belle, possibly BESIII)
 - nearly back-to-back hadrons
- Collins asymmetries:
 - neutral meson (pion and eta) incl. k_T dependence (Belle)
 - kaon and pion-kaon pairs as well as k_T dependence of Collins asymmetries (Belle, BESIII)
 - Collins asymmetries without double ratios (BaBar)
- di-hadron production (Belle)
 - collinear FF with new ISR and different momentum-fraction variable choice (Altarelli et al, NPB160 (1979) 301 & Mulders and Van Hulse)
 - helicity-dependent dihadron fragmentation function G_1^\perp ("jet handedness")

the future starts now



backup

ISR corrections - PRD96 (2017) 032005

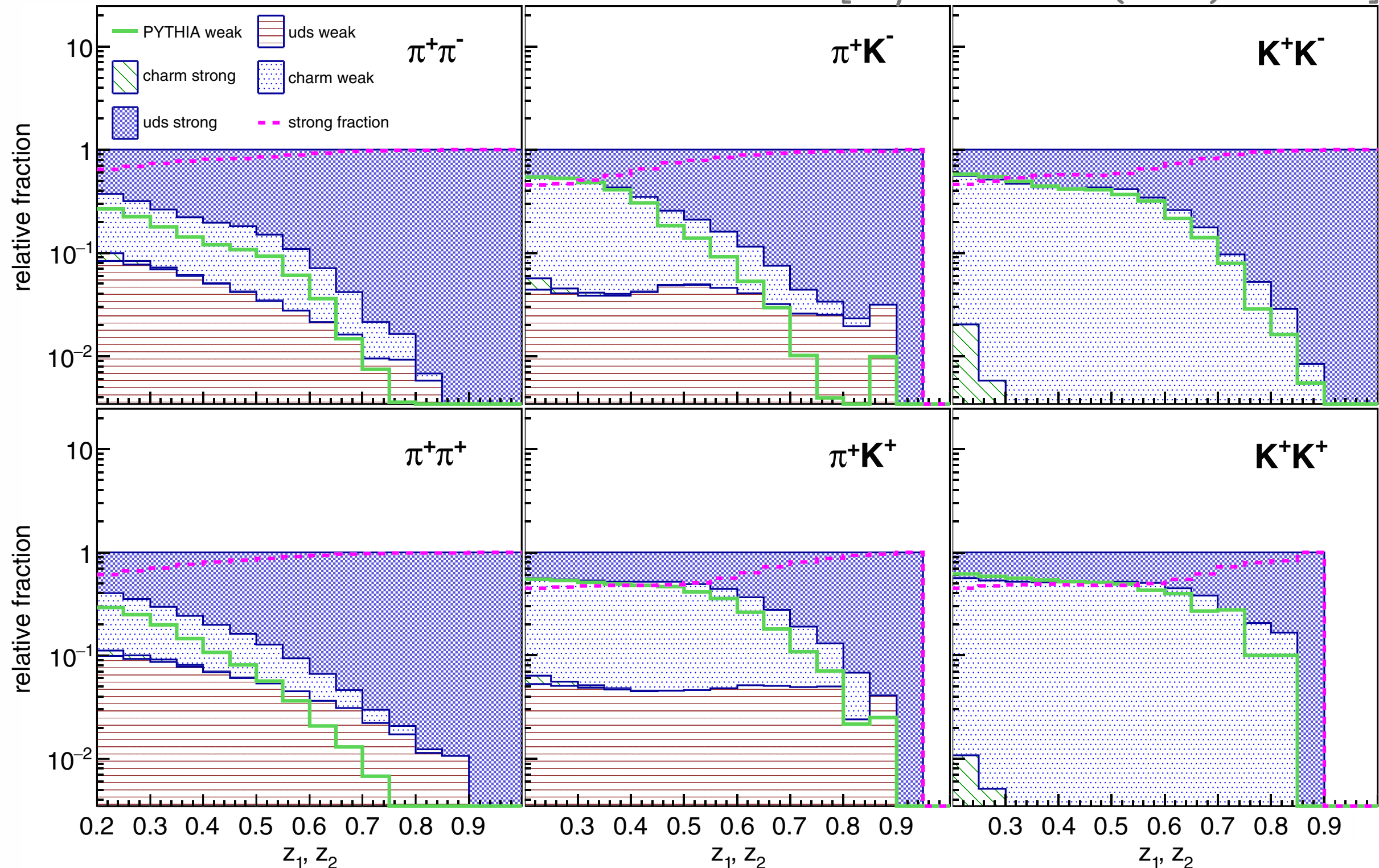


- non-ISR / ISR fractions based on PYTHIA switch MSTP(11)
[e.g., for dihadron invariant-mass analysis shown here, but discussed later]
- model dependence absorbed in variation of PYTHIA tunes in estimate of systematic uncertainty

hadron-pairs: weak-decay contributions

- not all hadrons originate from uds quarks but e.g., from D decay
- here only $z_1=z_2$ diagonal bins

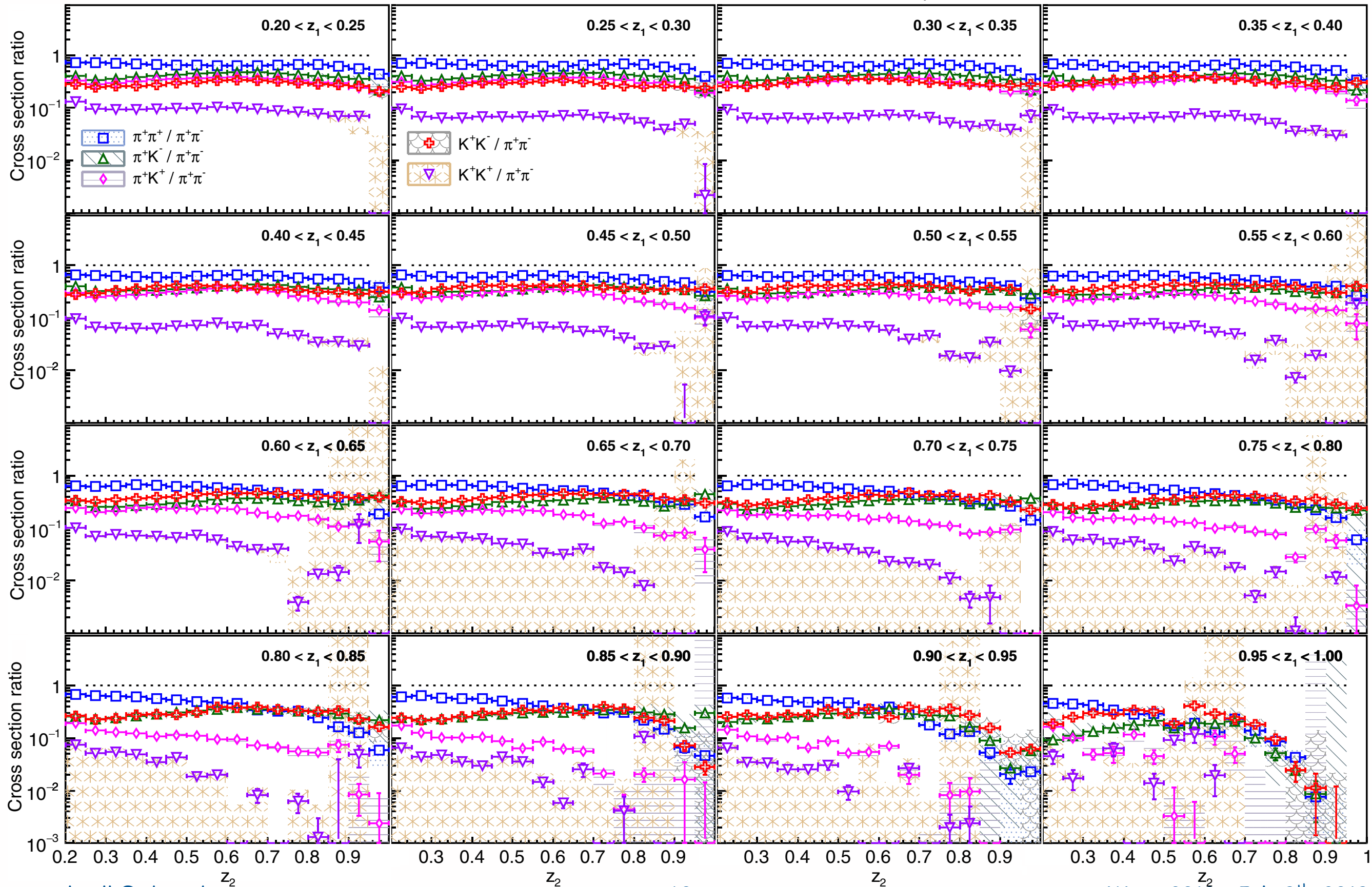
[Phys. Rev. D92 (2015) 092007]



no hemisphere selection

hadron-pair production

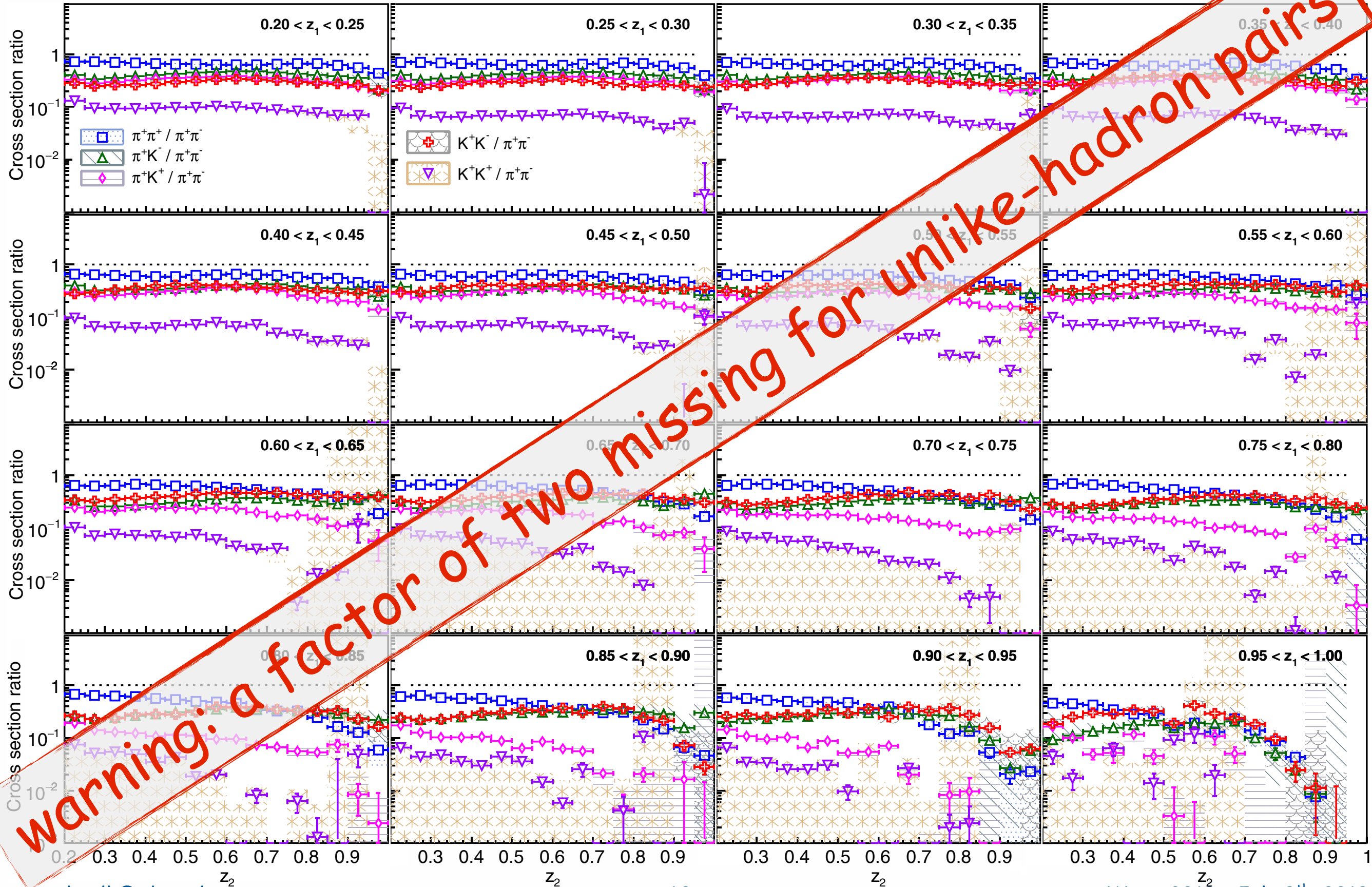
[Phys. Rev. D92 (2015) 092007]



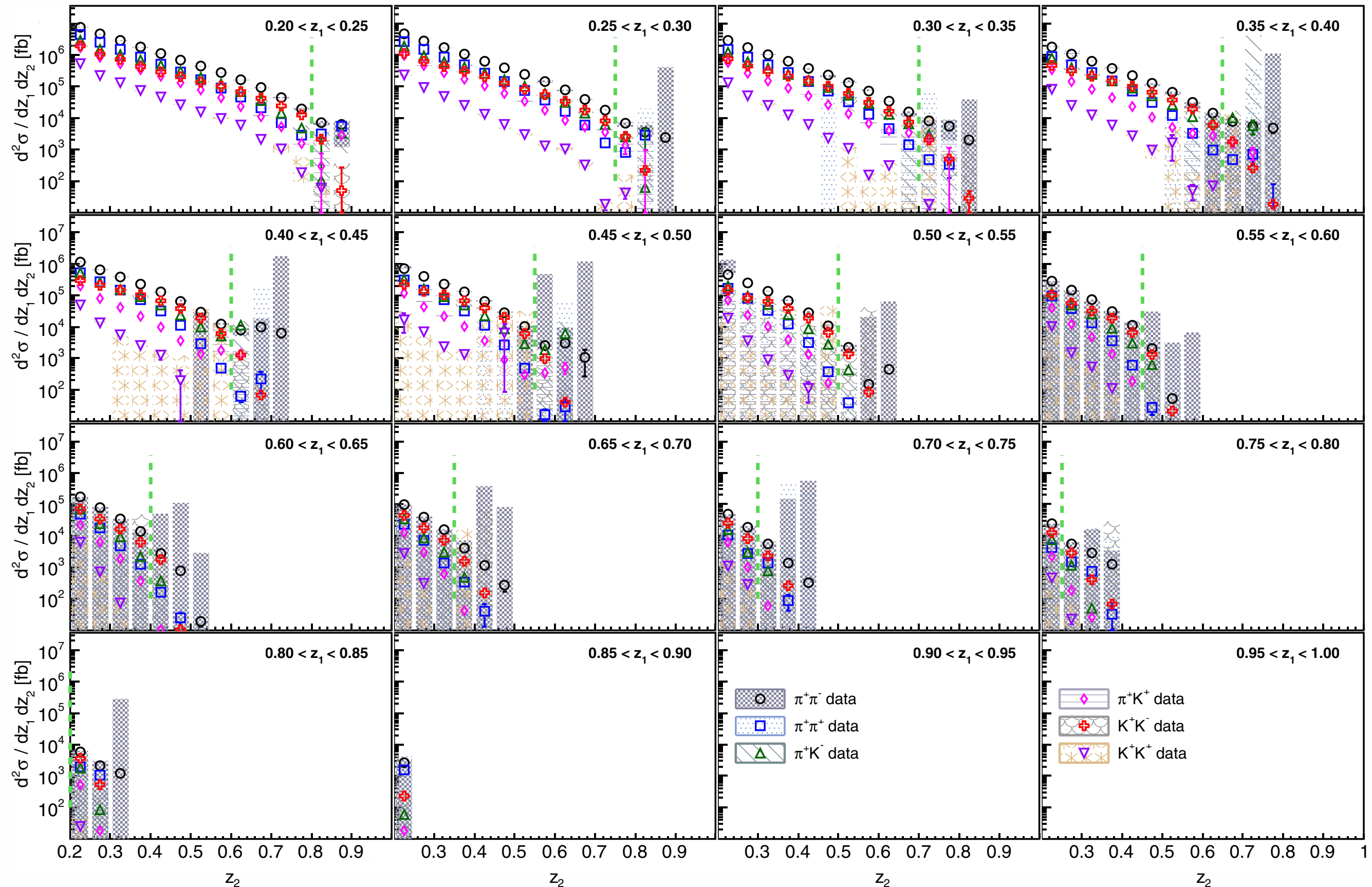
no hemisphere selection

hadron-pair production

[Phys. Rev. D92 (2015) 092007]

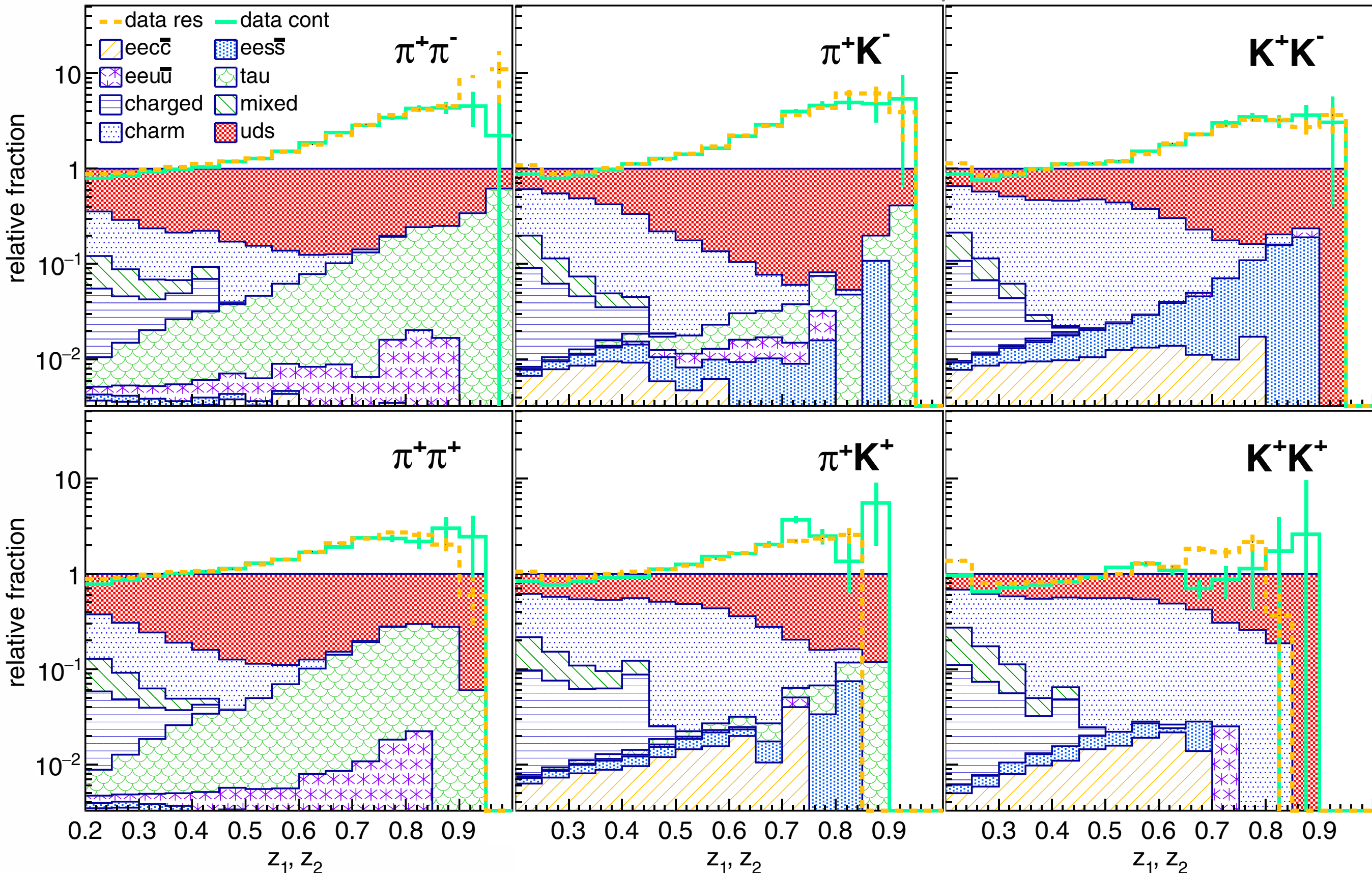


same-hemisphere hadron pairs



hadron-pairs: subprocess contributions

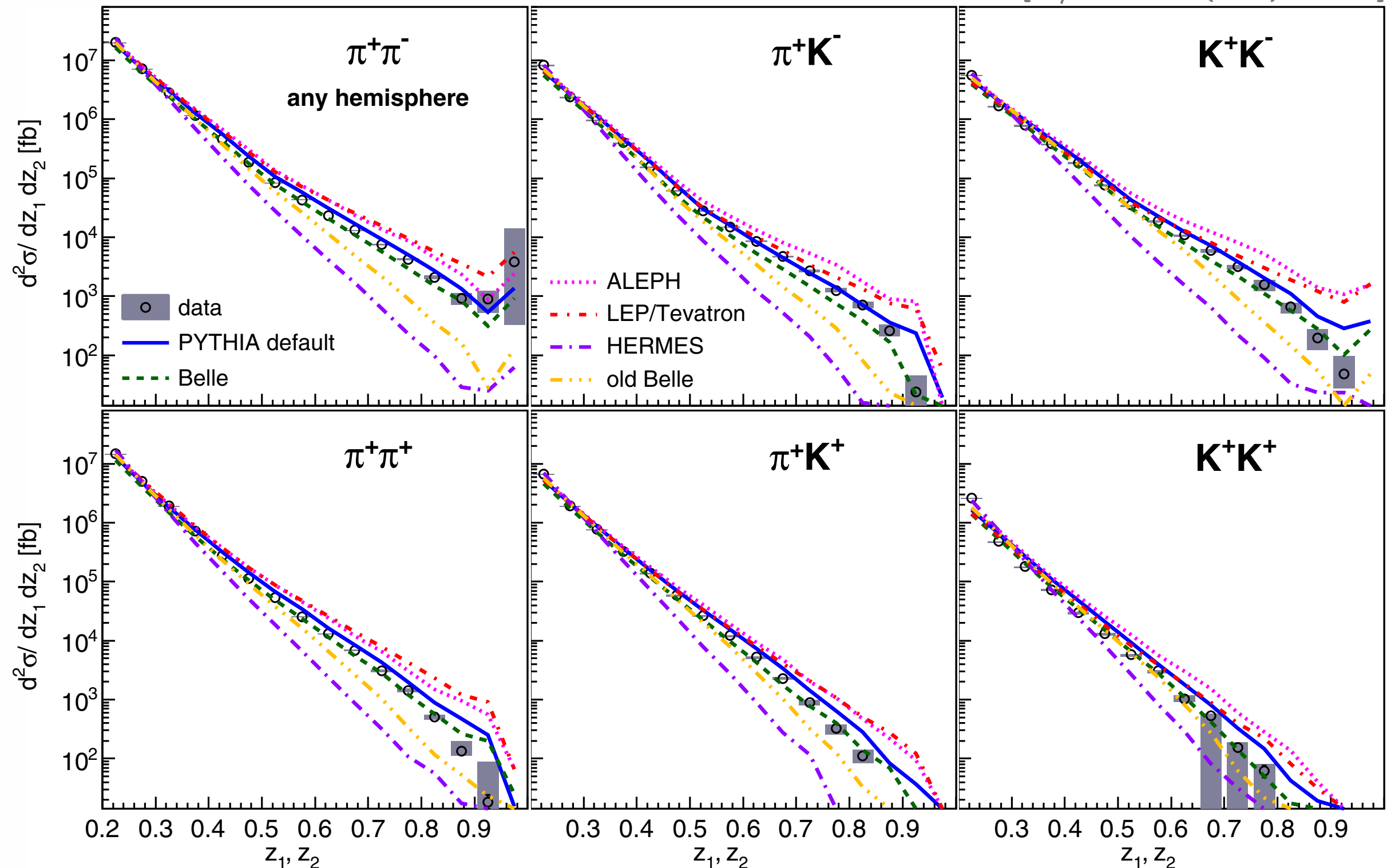
[Phys. Rev. D92 (2015) 092007]



hadron-pairs: comparison with PYTHIA

- generally good agreement at low z
- at large z only present Belle and PYTHIA default tunes satisfactory

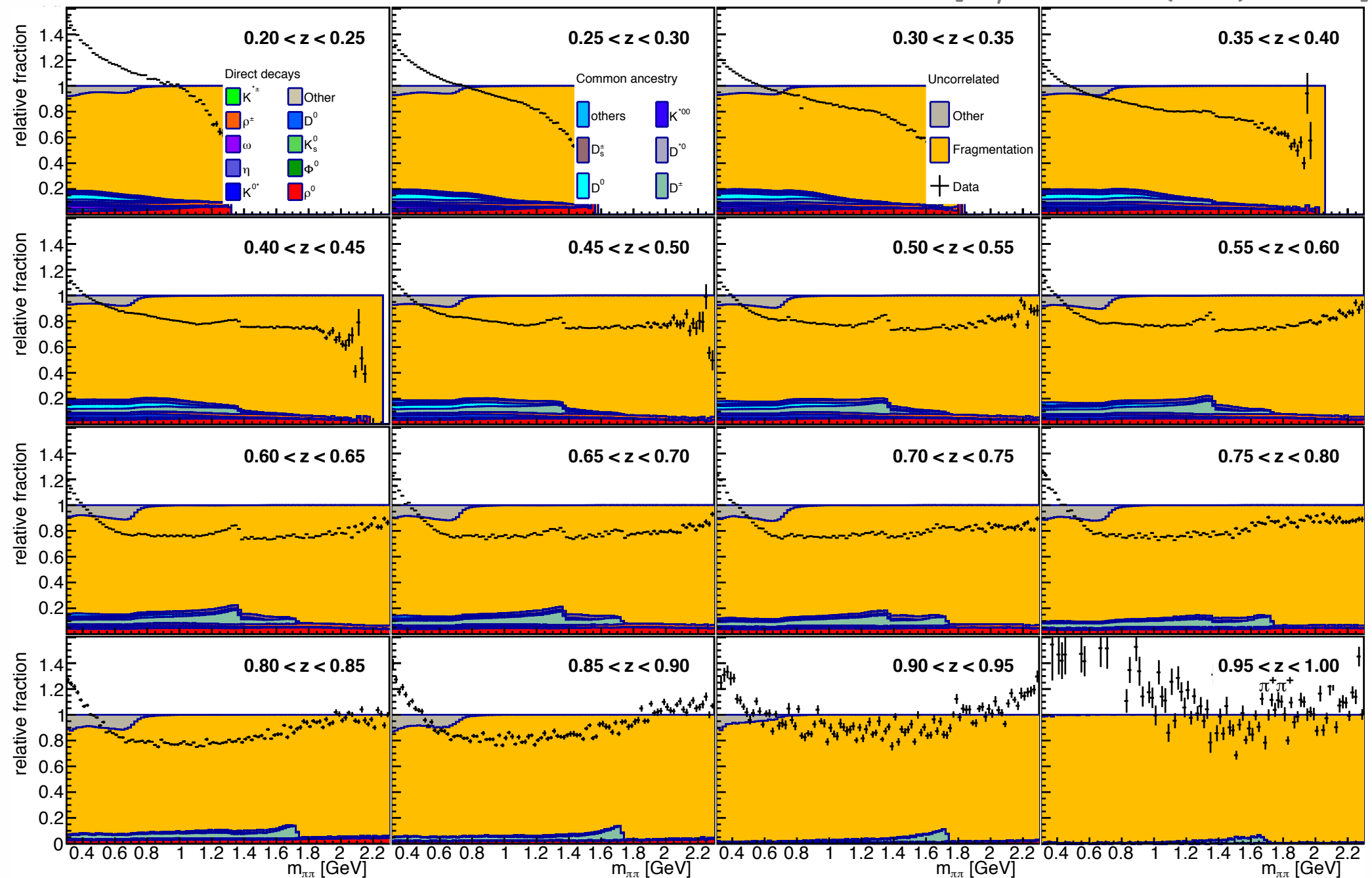
[Phys. Rev. D92 (2015) 092007]



same-hemisphere data: M_{h1h2} dependence

[Phys. Rev. D96 (2017) 032005]

like-sign
pion pairs



$T > 0.8$
 $z_{1,2} > 0.1$

- decomposition based on PYTHIA simulation
- though no strong resonance structure still clear MC/data discrepancy

quark-flavor contributions to Lambda prod.

- flavor tagging through opposite-hemisphere hadrons

