Transverse Λ polarisation in e^+e^- collisions

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Introduction

- Transverse Lambda polarisation in high-energy inclusive processes: a long standing open issue
- Proton-proton data: large polarisation that cannot be explained within collinear factorisation in QCD at leading twist
- First phenomenological studies within a TMD scheme: Anselmino, Boer, D'Alesio, Murgia (2001)

Contents

- TMDs FF with Helicity Formalism
- $e^+e^- \rightarrow h_1(jet)X$
- $e^+e^- \rightarrow h^{\uparrow}_1 h_2 X$
- Phenomenolgy
- Preliminary Fits

New Datas in e^+e^- from Belle: 150 data points

Observation of Transverse $\Lambda/\bar{\Lambda}$ Hyperon Polarization in e^+e^- Annihilation at Belle

We report the first observation of the spontaneous polarization of Λ and $\overline{\Lambda}$ hyperons transverse to the production plane in e^+e^- annihilation, which is attributed to the effect arising from a polarizing fragmentation function. For inclusive $\Lambda/\overline{\Lambda}$ production, we also report results with subtracted feed-down contributions from Σ^0 and charm. This measurement uses a dataset of 800.4 fb⁻¹ collected by the Belle experiment at or near a center-of-mass energy of 10.58 GeV. We observe a significant polarization that rises with the fractional energy carried by the $\Lambda/\overline{\Lambda}$ hyperon.

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Benefits:

- No PDFs
- Cleaner process
- TMD Factorization proven

Helicity Formalism: TMD Fragmentation Functions for quarks

Quark Polarisation States determine the Hadron ones

$$\rho_{\lambda_h,\lambda'_h}^{h,S_h} \hat{D}_{h/q,s_q}(z,\mathbf{k}_{\perp h}) = \sum_{\lambda_q,\lambda'_q} \rho_{\lambda_q,\lambda'_q}^{q,s_q} \hat{D}_{\lambda_q,\lambda'_q}^{\lambda_h,\lambda'_h}(z,\mathbf{k}_{\perp h})$$
$$P_j^h \hat{D}_{h/q,s_q} = \hat{D}_{S_j/q,s_q}^h - \hat{D}_{-S_j/q,s_q}^h = \Delta \hat{D}_{S_j/s_q}^h.$$



Polarisation: $e^+e^- \rightarrow h_1(jet)X$

$$p_{\lambda_{h_1},\lambda'_{h_1}}^{h_1} \frac{d\sigma^{e^+e^- \to h_1(jet)X}}{d\cos\theta dz_1 d^2 p_{\perp 1}} = \sum_{q_c} \sum_{\{\lambda\}} \frac{1}{32\pi s} \frac{1}{4} \hat{M}_{\lambda_c \lambda_d,\lambda_a \lambda_b} \hat{M}^*_{\lambda'_c \lambda_d,\lambda_a \lambda_b} \hat{D}^{\lambda_{h_1},\lambda'_{h_1}}_{\lambda_c,\lambda'_c} (z_1, p_{\perp 1})$$



Polarisation: $e^+e^- \rightarrow h_1(jet)X$

Only two possibile cross sections

$$\frac{d\sigma^{e^+e^- \rightarrow h_1(jet)X}}{d\cos\theta dz_1 d^2 p_{\perp 1}} = \sum_{q_e} \frac{3e^4}{32\pi s} e_q^2 (1 + \cos^2\theta) \hat{D}_{h/q}(z_1, p_{\perp h_1})$$

$$P_Y^{h_1} \frac{d\sigma^{e^+e^- \rightarrow h_1(jet)X}}{d\cos\theta dz_1 d^2 p_{\perp 1}} = \sum_{q_e} \frac{3e^4}{32\pi s} e_q^2 (1 + \cos^2\theta) \Delta D_{S_Y/q}^{h_1}(z_1, p_{\perp h_1})$$

$$P_Y^{h_1} = \frac{\sum_{q_e} q_q^2 \Delta D_{S_Y/q}^{h_1}(z_1, p_{\perp h_1})}{\sum_{q} e_q^2 D_{h/q}(z_1, p_{\perp h_1})}$$

$$I \text{ is possible to extract}$$

$$\int_{0}^{0} \frac{1}{\sqrt{q_e^+ q_e^- q_e^+ q_e^+$$

Polarisation:
$$e^+e^- \rightarrow h^{\uparrow}_{1}h_{2}X$$

 $p_{\lambda_{n_{1}}\lambda_{n_{1}}}^{h_{1},S_{n_{1}}}p_{\lambda_{n_{2}}\lambda_{n_{2}}}^{h_{2},S_{n_{2}}}\frac{d\sigma^{e^+e^- \rightarrow h_{1}h_{2}X}}{d\cos\theta dz_{1}d^{2}\mathbf{p}_{\perp h_{1}}dz_{2}d^{2}\mathbf{p}_{\perp h_{1}}}$
 $=\sum_{q_{c}}\sum_{\chi_{1}}\frac{3}{32\pi s}\frac{1}{4}M_{\lambda_{\lambda}\lambda_{n,\lambda,\lambda,k}}M_{\lambda_{\lambda}\lambda_{d,\lambda,\lambda,k}}^{*}D_{\lambda_{c},\lambda_{c}}^{\lambda_{n_{1}},\lambda_{n_{1}}}(z_{1},\mathbf{p}_{\perp h_{1}})\hat{O}_{\lambda_{d},\lambda_{d}}^{\lambda_{n_{2}},\lambda_{h_{2}}}(z_{2},\mathbf{p}_{\perp h_{2}})}$
 $\hat{Y}_{h_{1}} = \frac{\hat{Y}_{h_{1}} \times \hat{Z}_{h_{1}}}{|\hat{q}_{1} \times P_{h_{1}}|}$
 $\hat{Z}_{h_{1}} = \frac{P_{h_{1}}}{|P_{h_{1}}|}$
 $\hat{Z}_{h_{1}} = \frac{P_{h_{1}}}{|P_{h_{1}}|}$
 $\hat{P}^{h_{1}} = P_{x}^{h_{1}}\hat{X}_{1} + P_{y}^{h_{1}}\hat{Y}_{1} + P_{z}^{h_{1}}\hat{Z}_{1}$
Kinematics
 $P_{1} = z_{p_{1}}q_{1} + p_{\perp 1}$
 $P_{2} = z_{p_{2}}q_{2} + p_{\perp 2}$



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Polarisation:
$$e^+e^- \rightarrow h^{\uparrow}_1 h_2 X$$

$$\frac{D_{h_1/q}D_{h_2/\bar{q}}}{\Delta^N D_{h_1/q^{\uparrow}}\Delta^N D_{h_2/\bar{q}^{\uparrow}}\cos(2\varphi_2 + \phi_1^{h_1})}$$

$$\begin{split} \Delta D^{h_1}_{S_Y/q} D_{h_2/\bar{q}} \\ d\sigma^{\uparrow_Y} - d\sigma^{\downarrow_Y} \\ \Delta^- D^{h_1}_{S_Y/s_T} \Delta^N D_{h_2/\bar{q}}^{\uparrow} \cos(2\varphi_2 + \phi_1^{h_1}) \end{split}$$

$$d\sigma^{\uparrow_X} - d\sigma^{\downarrow_X} \qquad \Delta D^{h_1}_{S_X/s_T} \Delta^N D_{h_2/\bar{q}^{\uparrow}} \sin(2\varphi_2 + \phi_1^{h_1})$$

In order to obtain the polarisation we firstly calculate the polarisation along the Y and X helicity axes:

$$\langle P_Y^{h_1} \rangle = \frac{\langle d\sigma^{\uparrow_Y} - d\sigma^{\downarrow_Y} \rangle}{\langle d\sigma \rangle}$$

$$\langle P_X^{h_1} \rangle = \frac{\langle d\sigma^{\uparrow_X} - d\sigma^{\downarrow_X} \rangle}{\langle d\sigma \rangle}$$

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$$\frac{d\sigma^{e^+e^- \to h_1 h_2 X}}{d\cos\theta dz_1 d^2 \mathbf{p}_{\perp h_1} dz_2 d^2 \mathbf{p}_{\perp h_2}}$$

$$l^2 p_{\perp 1} \longrightarrow dP_{1T} d\phi_1$$

 \rangle_p

$$d\sigma^{e^+e^- \to h_1 h_2 X}$$

 $d\cos\theta dz_1 dP_{1T} d\phi_1 dz_2 d^2 \mathbf{p}_{\perp h_2}$

We use for the FF the following parametrisations:

$$\begin{split} \Delta D_{S_Y/q}^h(z,p_{\perp}) &= \Delta D_{S_Y/q}^h(z)\sqrt{2e} \frac{p_{\perp}}{M_p} \frac{e^{-p_{\perp}^2/\langle p_{\perp}^2 \rangle}}{\pi \langle p_{\perp}^2 \rangle_h} \\ D_{h/q}(z,p_{\perp}) &= D_{h/q}(z) \frac{e^{-p_{\perp}^2/\langle p_{\perp}^2 \rangle_h}}{\pi \langle p_{\perp}^2 \rangle_h} \\ \frac{1}{\langle p_{\perp}^2 \rangle_p} &= \frac{1}{M_p^2} + \frac{1}{\langle p_{\perp}^2 \rangle_h} \end{split}$$

 $\int dP_{1T} d\phi_1 dp_{\perp 2} d\varphi_2$

We find an expression for the polarisation that depends only on the energy fractions z_1 and z_2

$$P^{h_1} \cdot \hat{n} :: \qquad \frac{\sum_q \Delta D^h_{S_Y/q}(z) D_{h_2/\bar{q}}(z_2)}{\sum_q D_{h_1/q}(z_1) D_{h_2/\bar{q}}(z_2)} \\ \times \quad \frac{\sqrt{2e\pi}}{2M_p} \frac{\langle p_\perp^2 \rangle_p^2}{\langle p_\perp^2 \rangle_1} \frac{z_2}{\sqrt{z_1^2 \langle p_\perp^2 \rangle_2 + z_2^2 \langle p_\perp^2 \rangle_p}}$$

Phenomenology

From datas we can extract different information, particularly:

- $\Lambda(jet)X$: Lambda polarising width $\langle p_{\perp}^2
 angle_p$
- $\Lambda \pi X$: Polarising FF (u,d)
- ΛkX : Polarising FF (u,s)

Fitted parameters:

Flav.	\mathcal{N}^p_q	$lpha_q$	eta_q	$\langle p_{\perp}^2 angle_p$
U	\mathcal{N}_{u}^{p}			$\langle p_{\perp}^2 \rangle_p$
d	\mathcal{N}^p_d			
S	\mathcal{N}^p_s	α_s		
sea	\mathcal{N}^p_{sea}		β_{sea}	

Polarising parametrization:
$\Delta D^h_{S_Y/q}(z) = \mathcal{N}^p_q(z) D_{h/q}(z)$
$\mathcal{N}_q^p(z) = \mathcal{N}_q^p z^{\alpha_q} (1-z)^{\beta_q} \frac{(\alpha_q + \beta_q)^{(\alpha_q + \beta_q)}}{\alpha_q^{\alpha_q} \beta_q^{\beta_q}}$

Since we did not have a set of FF that separates the Λ and $\overline{\Lambda}$:

$$D_q^{\Lambda} = D_q^{\Lambda^0} + D_q^{\bar{\Lambda}^0} \longrightarrow D_q^{\Lambda} = D_q^{\Lambda^0} + D_{\bar{q}}^{\Lambda^0}$$

$$D_{\bar{q}}^{\Lambda^{0}} = (1-z)^{\alpha} D_{q}^{\Lambda^{0}}$$
$$D_{q}^{\Lambda^{0}} = \frac{1}{1+(1-z)^{\alpha}} D_{q}^{\Lambda} \qquad \alpha = 1, 2$$
$$D_{\bar{q}}^{\Lambda^{0}} = \frac{(1-z)^{\alpha}}{1+(1-z)^{\alpha}} D_{q}^{\Lambda}$$

Global Fit

Polarisation: $\Lambda(jet)$



Global Fit



Datas cut for z > 0.5 : $\Lambda(jet)$, Λk

Global Fit



Datas cut for z > 0.5 : $\Lambda(jet)$, Λk

Fit: associated production only

Polarisation: $\Lambda\pi$



Datas cut for z > 0.5 : $\Lambda(jet)$, Λk

Fit: associated production only

Polarisation: Λk



Fit: associated production only (fixed gaussian width)

Polarization: $\Lambda\pi$ $\chi^2/dof = 1.5178$ z1=0.25, Λ π[±] z1=0.35, Λ π[±] 0.04 0.02 0.02 0.00 0.00 -0.02 H 0.02 -0.04 -0.04 -0.06 -0.06theo Λπ theo_Λπ⁴ theo Λπ⁻ theo Λπ⁻⁻ exp_Λπ⁺ $exp_\Lambda \pi^+$ ٠ -0.08 -0.08 exp_Λπ⁻ exp_Λπ⁻ 0.30 0.35 0.40 0.45 0.50 0.55 0.60 0.30 0.35 0.40 0.45 0.50 0.55 0.60 0.25 0.25 z1=0.45, ∧ n[±] z1=0.6, ∧ π[±] 0.02 -0.01 0.00 -0.02 -0.03 -0.02 ⊢ ଅ. ______ -0.04 -0.05 -0.06theo_Λπ⁺ theo Λπ* theo Λπ⁻ -0.06 theo Arr exp_лл* exp_Λπ⁺ ехр_Лл= exp_Λπ⁻ -0.08 -0.0 0.30 0.35 0.40 0.45 0.50 0.55 0.60 0.25 0.30 0.35 0.40 0.45 0.50 0.55 0.60 0.25 z2 z2

Datas cut for z > 0.5 : $\Lambda(jet)$, Λk

Fit: associated production only (fixed gaussian width)

Polarisation: Λk





Conclusions

• Global fit : $\chi^2/dof = 2.5683$

Lead to different width $\langle p_{\perp}^2
angle_p$

• Associated fit with no Lambda jet, fixing width: $\chi^2/dof = 1.5178$

• Associated fit with no Lambda jet: $\chi^2/dof = 1.499$

What's next?

- Different choices of the Gaussian width for the unpolarised and/or the polarising FF: z dependence, flavour dependence.
- Functional form of the polarising FF.
- Predictions for proton-proton and comparison with existing data and previous extractions

Grazie

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$$\begin{aligned} \frac{d\sigma^{e^+e^- \to h_1h_2X}}{d\cos\theta dz_1 d^2 p_{\perp 1} dz_2 d^2 p_{\perp 2}} & P_Y^{h_1} \frac{d\sigma^{e^+e^- \to h_1h_2X}}{d\cos\theta dz_1 d^2 p_{\perp 1} dz_2 d^2 p_{\perp 2}} \\ = \frac{6e^4 e_q^2}{64\pi \hat{s}} \bigg\{ D_{h_1/q}(z_1, p_{\perp 1}) D_{h_2/\bar{q}}(z_2, p_{\perp 2}) \Big(1 + \cos^2\theta \Big) & = \frac{6e^4 e_q^2}{64\pi \hat{s}} \bigg\{ \Delta D_{S_Y/q}^{h_1}(z_1, p_{\perp 1}) D_{h_2/\bar{q}}(z_2, p_{\perp 2}) \Big(1 + \cos^2\theta \Big) \\ + \frac{1}{4} \sin^2\theta \Delta^N D_{h_1/q^+}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}^+}(z_2, p_{\perp 2}) \cos(2\varphi_2 + \phi_1^{h_1}) \bigg\} & + \frac{1}{2} \sin^2\theta \Delta^- D_{S_Y/s_T}^{h_1}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}^+}(z_2, p_{\perp 2}) \cos(2\varphi_2 + \phi_1^{h_1}) \bigg\} \end{aligned}$$

$$P_X^{h_1} \frac{d\sigma^{e^+e^- \to h_1 h_2 X}}{d\cos\theta dz_1 d^2 p_{\perp 1} dz_2 d^2 p_{\perp 2}} = \frac{3e^4 e_q^2}{64\pi \hat{s}} \Delta D_{S_X/s_T}^{h_1}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}^{\uparrow}}(z_2, p_{\perp 2}) \sin^2\theta \sin(2\varphi_2 + \phi_1^{h_1})$$

$$P_{Z}^{h_{1}} \frac{d\sigma^{e^{+}e^{-} \to h_{1}h_{2}X}}{d\cos\theta dz_{1}d^{2}p_{\perp 1}dz_{2}d^{2}p_{\perp 2}}$$

= $\frac{3e^{4}e_{q}^{2}}{64\pi\hat{s}}\Delta D_{S_{Z}/s_{T}}^{h_{1}}(z_{1}, p_{\perp 1})\Delta^{N}D_{h_{2}/\bar{q}^{\uparrow}}(z_{2}, p_{\perp 2})\sin^{2}\theta\sin(2\varphi_{2} + \phi_{1}^{h_{1}})$

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TMD Formalism and Azimuthal Asymmetries

Generalization of PDF and FF including a Transverse Momentum Dependence.

Correlation Azimuthal distribution - Spin

PDF
$$\Delta^N f_{q/p^\uparrow}$$
 Sivers 1991
FF $\Delta^N D_{h/q^\uparrow}$ Collins 1993



Spin and TMD Effects

$$a^{\uparrow}b \to cd$$
$$a_N = \frac{d\sigma^{a^{\uparrow}b \to cd} - d\sigma^{a^{\downarrow}b \to cd}}{d\sigma^{a^{\uparrow}b \to cd} + d\sigma^{a^{\downarrow}b \to cd}}$$

 $a_N \propto \alpha_s rac{m}{\sqrt{s}} \simeq \alpha_s rac{m}{p_\perp}$ [Kane, Pumplin, Repko 1978]

Spin and transverse momentum effects were considered negligible.

But Experimental Data show $\,A_N\simeq 20\%$

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Collinear pQCD



 $p^{\uparrow}p \to \pi X$

 $A_N = \frac{d\sigma^{\uparrow} - d\sigma^{\downarrow}}{d\sigma^{\uparrow} + d\sigma^{\downarrow}}$

Cross Sections and Polarization States for $e^+e^- \rightarrow h^{\uparrow}_1 h_2 X$ $\rho_{\lambda_{h_1},\lambda_{h_1}'}^{h_1,S_{h_1}}\rho_{\lambda_{h_2},\lambda_{h_2}'}^{h_2,S_{h_2}}\frac{d\sigma^{e^+e^-\to h_1h_2X}}{d\cos\theta dz_1 d^2\mathbf{k}_{\perp h_1} dz_2 d^2\mathbf{k}_{\perp h_2}}$ $=\sum_{q_c}\sum_{\{\lambda\}}\frac{3}{32\pi s}\frac{1}{4}M_{\lambda_c\lambda_d,\lambda_a\lambda_b}M^*_{\lambda'_c\lambda'_d,\lambda_a\lambda_b}\hat{D}^{\lambda_{h_1},\lambda'_{h_1}}_{\lambda_c,\lambda'_c}(z_1,\mathbf{k}_{\perp h_1})\hat{O}^{\lambda_{h_2},\lambda'_{h_2}}_{\lambda_d,\lambda'_d}(z_2,\mathbf{k}_{\perp h_2})$ Cross Sections and Polarization States for $e^+e^- \rightarrow h_{\perp}^{\uparrow}h_2X$ $\langle P_Y^{h_1} \rangle = \frac{\langle d\sigma^{\uparrow_Y} - d\sigma^{\downarrow_Y} \rangle}{\langle d\sigma \rangle}$ $d\sigma^{nonpol.}$ $\Delta^{N} D_{h_{1}/q^{\uparrow}} \Delta^{N} D_{h_{2}/\bar{q}^{\uparrow}} \cos(\varphi_{h_{1}} + \varphi_{h_{2}})$ $d\sigma^{\uparrow_{Y}} - d\sigma^{\downarrow_{Y}} \qquad \frac{\Delta D_{S_{Y}/q}^{h_{1}} D_{h_{2}/\bar{q}}}{\Delta^{-} D_{S_{Y}/s_{T}}^{h_{1}/q} \Delta^{N} D_{h_{2}/\bar{q}^{\uparrow}} \cos(\varphi_{h_{1}} + \varphi_{h_{2}})}$ $\Delta \hat{D}^h_{S_Y/q}$ Polarizing FF

Observation of Transverse Λ Hyperon Polarization in e^+e^- Annihilation at Belle (2018)

Measurement for Λ polarization in:

- $e^+e^- \rightarrow \Lambda (\text{jet})X$
- $e^+e^- \rightarrow \Lambda \pi X$

Benefits:

- TMD factorization is valid;
- NO PDF
- Extraction Polarizing FF



Transverse Polarization predicted in the Helicity Frame of the Hadron

 P_Y^{Λ}

Measured along:

$$\hat{n} = \frac{\vec{P_{\Lambda}} \times \vec{P_{\pi}}}{|\vec{P_{\Lambda}} \times \vec{P_{\pi}}|}$$

FF properties

$$\begin{split} \hat{D}_{h/q}(z,\mathbf{k}_{\perp,h}) &= D_{h/q} = (D_{++}^{++} + D_{--}^{++}) \\ \bar{\hat{D}}_{h/q,s_{T}}(z,\mathbf{k}_{\perp,h}) &= \hat{D}_{h/q} + \frac{1}{2}\Delta\hat{D}_{h/q,s_{T}} \\ \Delta\hat{D}_{h/q,s_{T}}(z,\mathbf{k}_{\perp,h}) &= \Delta^{N}D_{h/q^{\dagger}}\sin(\phi_{s_{q}} - \phi_{h}) = 4ImD_{+-}^{++}\sin(\phi_{s_{q}} - \phi_{h}) \quad [Collins] \\ \Delta\hat{D}_{S_{Z}/s_{L}}^{h/q}(z,\mathbf{k}_{\perp,h}) &= \Delta D_{S_{Z}/s_{L}}^{h/q} = (D_{++}^{++} - D_{--}^{++}) \\ \Delta\hat{D}_{S_{Z}/s_{T}}^{h/q}(z,\mathbf{k}_{\perp,h}) &= \Delta D_{S_{Z}/s_{T}}^{h/q}\cos(\phi_{s_{q}} - \phi_{h}) = 2ReD_{+-}^{++}\cos(\phi_{s_{q}} - \phi_{h}) \\ \Delta\hat{D}_{S_{X}/s_{L}}^{h/q}(z,\mathbf{k}_{\perp,h}) &= \Delta D_{S_{X}/s_{L}}^{h/q} = 2ReD_{++}^{+-} \\ \Delta\hat{D}_{S_{X}/s_{T}}^{h/q}(z,\mathbf{k}_{\perp,h}) &= \Delta D_{S_{X}/s_{T}}^{h/q}\cos(\phi_{s_{q}} - \phi_{h}) = (D_{+-}^{+-} + D_{-+}^{+-})\cos(\phi_{s_{q}} - \phi_{h}) \\ \Delta\hat{D}_{S_{Y}/q}^{h}(z,\mathbf{k}_{\perp,h}) &= \Delta D_{S_{Y}/q}^{h/q} = -2ImD_{++}^{+-} \quad [Polarizing] \\ \Delta\hat{D}_{S_{Y}/s_{T}}^{h/q}(z,\mathbf{k}_{\perp,h}) &= \Delta\hat{D}_{S_{Y}/c}^{h/q} + \Delta^{-}\hat{D}_{S_{Y}/s_{T}}^{h/q} \\ \Delta^{-}\hat{D}_{S_{Y}/s_{T}}^{h/q}(z,\mathbf{k}_{\perp,h}) &= \Delta^{-}D_{S_{Y}/s_{T}}^{h/q}\sin(\phi_{s_{q}} - \phi_{h}) = (D_{+-}^{+-} - D_{-+}^{+-})\sin(\phi_{s_{q}} - \phi_{h}) \\ \end{split}$$

$$\begin{split} D_{++}^{++} &= D_{--}^{--} \\ D_{--}^{++} &= D_{++}^{--} \\ D_{+-}^{++} &= -D_{-+}^{--} \\ D_{-+}^{++} &= -D_{+-}^{--} \\ D_{++}^{+-} &= -D_{--}^{-+} \\ D_{++}^{+-} &= -D_{-+}^{+-} \\ D_{+-}^{+-} &= D_{-+}^{-+} \\ D_{+-}^{+-} &= D_{+-}^{-+} \\ D_{++}^{+-} &= (D_{-+}^{++})^* \\ D_{++}^{+-} &= (D_{-+}^{-+})^* \end{split}$$

$$P_T \cdot \hat{n} = P_X^{h_1} \cos \widetilde{\phi} + P_Y^{h_1} \sin \widetilde{\phi}$$

Helicity Scattering Amplitudes

$$M^0_{\lambda_c \lambda_d, \lambda_a \lambda_b}$$

$$q_1 = E(1,0,0,1) \qquad p_{e^+} = E(1,\sin\theta,0,\cos\theta)$$
$$q_2 = E(1,0,0,-1) \qquad p_{e^-} = E(1,-\sin\theta,0,-\cos\theta).$$



$$M^{0}_{+-,+-} = M^{0}_{-+,-+} = M^{0}_{2}$$
$$M^{0}_{+-,-+} = M^{0}_{-+,+-} = M^{0}_{3}$$

 $M_2^0 = -ie^2 e_q (1 - \cos \theta);$ $M_3^0 = -ie^2 e_q (1 + \cos \theta);$

Cross Sections

$$d\sigma^{unpol.}$$

$$= \frac{3\pi\alpha^{2}e_{q}^{2}}{2s} \left\{ (1+\cos^{2}\theta)D_{h_{1}/q}D_{h_{2}/\bar{q}} + \frac{1}{4}\sin^{2}\theta\Delta^{N}D_{h_{1}/q^{\uparrow}}\Delta^{N}D_{h_{2}/\bar{q}^{\uparrow}}\cos(\varphi_{h_{1}}+\varphi_{h_{2}}) \right\}$$

$$d\sigma^{\uparrow_{Y}} - d\sigma^{\downarrow_{Y}}$$

$$= \frac{3\pi\alpha^{2}e_{q}^{2}}{2s} \left\{ (1+\cos^{2}\theta)\Delta D_{S_{Y}/q}^{h_{1}}D_{h_{2}/\bar{q}} + \frac{1}{2}\sin^{2}\theta\Delta^{-}D_{S_{Y}/s_{T}}^{h_{1}/q}\Delta^{N}D_{h_{2}/\bar{q}^{\uparrow}}\cos(\varphi_{h_{1}}+\varphi_{h_{2}}) \right\}$$