

The 3D structure of the proton and its impact on high-energy precision measurements

giuseppe bozzi



A neither complete
nor exhaustive
review of something

giuseppe bozzi

(with the invaluable help of
Valerio Bertone and Miguel Echevarria)

Disclaimer

Just a few slides to stimulate discussion and try to anticipate possible future steps for our benchmark

Kind of bird-eye view of different formalisms, without many technical details and with some "dictionary" included

Please forgive and point out any omission/mistake/inaccuracy: slides are meant to be continuously (even real-time!) updated

SCET

SCET

- effective field theory: high energy d.o.f. integrated out, soft and collinear d.o.f. decouple

SCET

- effective field theory: high energy d.o.f. integrated out, soft and collinear d.o.f. decouple

$$A^\mu \rightarrow A_s^\mu + A_c^\mu + A_{\bar{c}}^\mu + \dots$$

$$\psi \rightarrow \psi_s + \psi_c + \psi_{\bar{c}} + \dots$$

$$\mathcal{L} = \mathcal{L}_s + \mathcal{L}_c + \mathcal{L}_{\bar{c}} + \dots$$

(DY)

$$\sigma \sim \text{Beam}(\mu, \nu) \otimes \text{Beam}(\mu, \nu) \otimes \text{Soft}(\mu, \nu) \otimes \text{Hard}(\mu, Q)$$

SCET

- effective field theory: high energy d.o.f. integrated out, soft and collinear d.o.f. decouple

$$A^\mu \rightarrow A_s^\mu + A_c^\mu + A_{\bar{c}}^\mu + \dots$$

$$\psi \rightarrow \psi_s + \psi_c + \psi_{\bar{c}} + \dots$$

$$\mathcal{L} = \mathcal{L}_s + \mathcal{L}_c + \mathcal{L}_{\bar{c}} + \dots$$

(DY) $\sigma \sim \text{Beam}(\mu, \nu) \otimes \text{Beam}(\mu, \nu) \otimes \text{Soft}(\mu, \nu) \otimes \text{Hard}(\mu, Q)$

- Beam function directly related to collinear PDF ($C \times f$)

SCET

- effective field theory: high energy d.o.f. integrated out, soft and collinear d.o.f. decouple

$$A^\mu \rightarrow A_s^\mu + A_c^\mu + A_{\bar{c}}^\mu + \dots$$

$$\psi \rightarrow \psi_s + \psi_c + \psi_{\bar{c}} + \dots$$

$$\mathcal{L} = \mathcal{L}_s + \mathcal{L}_c + \mathcal{L}_{\bar{c}} + \dots$$

$$(DY) \quad \sigma \sim \text{Beam}(\mu, \nu) \otimes \text{Beam}(\mu, \nu) \otimes \text{Soft}(\mu, \nu) \otimes \text{Hard}(\mu, Q)$$

- Beam function directly related to collinear PDF ($C \times f$)
- two non-physical scales for renormalisation of UV (μ) and rapidity (ν) divergences

SCET

- effective field theory: high energy d.o.f. integrated out, soft and collinear d.o.f. decouple

$$A^\mu \rightarrow A_s^\mu + A_c^\mu + A_{\bar{c}}^\mu + \dots$$

$$\psi \rightarrow \psi_s + \psi_c + \psi_{\bar{c}} + \dots$$

$$\mathcal{L} = \mathcal{L}_s + \mathcal{L}_c + \mathcal{L}_{\bar{c}} + \dots$$

$$(DY) \quad \sigma \sim \text{Beam}(\mu, \nu) \otimes \text{Beam}(\mu, \nu) \otimes \text{Soft}(\mu, \nu) \otimes \text{Hard}(\mu, Q)$$

- Beam function directly related to collinear PDF ($C \times f$)
- two non-physical scales for renormalisation of UV (μ) and rapidity (ν) divergences
- ν arises from distinguishing soft modes from collinear modes (connected by Lorentz boost), just as μ arises from distinguishing different virtualities

SCET

- effective field theory: high energy d.o.f. integrated out, soft and collinear d.o.f. decouple

$$A^\mu \rightarrow A_s^\mu + A_c^\mu + A_{\bar{c}}^\mu + \dots$$

$$\psi \rightarrow \psi_s + \psi_c + \psi_{\bar{c}} + \dots$$

$$\mathcal{L} = \mathcal{L}_s + \mathcal{L}_c + \mathcal{L}_{\bar{c}} + \dots$$

$$(DY) \quad \sigma \sim \text{Beam}(\mu, \nu) \otimes \text{Beam}(\mu, \nu) \otimes \text{Soft}(\mu, \nu) \otimes \text{Hard}(\mu, Q)$$

- Beam function directly related to collinear PDF ($C \times f$)
- two non-physical scales for renormalisation of UV (μ) and rapidity (ν) divergences
- ν arises from distinguishing soft modes from collinear modes (connected by Lorentz boost), just as μ arises from distinguishing different virtualities
- each function has its own RG evolution: $\frac{d \ln X}{d \ln \mu} = \Gamma_X$ with $X = B, H, S$ leading to resummed predictions and customary formula (next slide)

TMD

TMD

- 2D factorisation theorem (SIDIS, DY, $e^+e^- \rightarrow$ hadrons)

TMD

- 2D factorisation theorem (SIDIS, DY, $e^+e^- \rightarrow$ hadrons)

(DY)

$$\sigma \sim f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes \text{Hard}$$

TMD

- 2D factorisation theorem (SIDIS, DY, $e^+e^- \rightarrow$ hadrons)

$$(DY) \quad \sigma \sim f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes \text{Hard}$$

- same divergencies \rightarrow same RG evolutions as SCET

TMD

- 2D factorisation theorem (SIDIS, DY, $e^+e^- \rightarrow$ hadrons)

$$(DY) \quad \sigma \sim f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes \text{Hard}$$

- same divergencies \rightarrow same RG evolutions as SCET

$$f(x, k_T^2, \mu^2) \sim \text{Beam} \otimes \sqrt{\text{Soft}}$$

TMD

- 2D factorisation theorem (SIDIS, DY, $e^+e^- \rightarrow$ hadrons)

$$(DY) \quad \sigma \sim f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes \text{Hard}$$

- same divergencies \rightarrow same RG evolutions as SCET

$$f(x, k_T^2, \mu^2) \sim \text{Beam} \otimes \sqrt{\text{Soft}}$$

- For the Drell-Yan process, SCET and TMD are equivalent

TMD

- 2D factorisation theorem (SIDIS, DY, $e^+e^- \rightarrow$ hadrons)

$$(DY) \quad \sigma \sim f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes \text{Hard}$$

- same divergencies \rightarrow same RG evolutions as SCET

$$f(x, k_T^2, \mu^2) \sim \text{Beam} \otimes \sqrt{\text{Soft}}$$

- For the Drell-Yan process, SCET and TMD are equivalent

$$f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes \text{Hard} \sim \text{Beam} \otimes \text{Beam} \otimes \text{Soft} \otimes \text{Hard}$$

TMD

- 2D factorisation theorem (SIDIS, DY, $e^+e^- \rightarrow$ hadrons)

$$(DY) \quad \sigma \sim f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes \text{Hard}$$

- same divergencies \rightarrow same RG evolutions as SCET

$$f(x, k_T^2, \mu^2) \sim \text{Beam} \otimes \sqrt{\text{Soft}}$$

- For the Drell-Yan process, SCET and TMD are equivalent

$$f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes \text{Hard} \sim \text{Beam} \otimes \text{Beam} \otimes \text{Soft} \otimes \text{Hard}$$

- Common form (CSS)

TMD

- 2D factorisation theorem (SIDIS, DY, $e^+e^- \rightarrow$ hadrons)

$$(DY) \quad \sigma \sim f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes Hard$$

- same divergencies \rightarrow same RG evolutions as SCET

$$f(x, k_T^2, \mu^2) \sim Beam \otimes \sqrt{Soft}$$

- For the Drell-Yan process, SCET and TMD are equivalent

$$f(x, k_T^2, \mu^2) \otimes f(x, k_T^2, \mu^2) \otimes Hard \sim Beam \otimes Beam \otimes Soft \otimes Hard$$

- Common form (CSS)

$$\sigma \sim [C \otimes f(x)] Hard [C \otimes f(x)] \exp(S) \exp(S_{NP})$$

qt resummation

qt resummation

- 1D (collinear) factorisation theorem

qT resummation

- 1D (collinear) factorisation theorem

$$\sigma \sim f(x, \mu^2) \otimes f(x, \mu^2) \otimes \text{Hard}$$

QT resummation

- 1D (collinear) factorisation theorem

$$\sigma \sim f(x, \mu^2) \otimes f(x, \mu^2) \otimes \text{Hard}$$

- resummation (i.e. factorisation and exponentiation) of soft-gluon emissions to all orders

qT resummation

- 1D (collinear) factorisation theorem

$$\sigma \sim f(x, \mu^2) \otimes f(x, \mu^2) \otimes \text{Hard}$$

- resummation (i.e. factorisation and exponentiation) of soft-gluon emissions to all orders

- dynamical ($|M_{n\text{-gluons}}| \sim |M_{1\text{-gluon}}|^n$) and kinematical ($PS_{n\text{-gluons}} \sim \prod_n PS_{1\text{-gluon}}$) factorisation properties of QCD

qT resummation

- 1D (collinear) factorisation theorem

$$\sigma \sim f(x, \mu^2) \otimes f(x, \mu^2) \otimes \text{Hard}$$

- resummation (i.e. factorisation and exponentiation) of soft-gluon emissions to all orders
- dynamical ($|M_{n\text{-gluons}}| \sim |M_{1\text{-gluon}}|^n$) and kinematical ($PS_{n\text{-gluons}} \sim \prod_n PS_{1\text{-gluon}}$) factorisation properties of QCD
- identical formula as in the previous slide

qT resummation

- 1D (collinear) factorisation theorem

$$\sigma \sim f(x, \mu^2) \otimes f(x, \mu^2) \otimes \text{Hard}$$

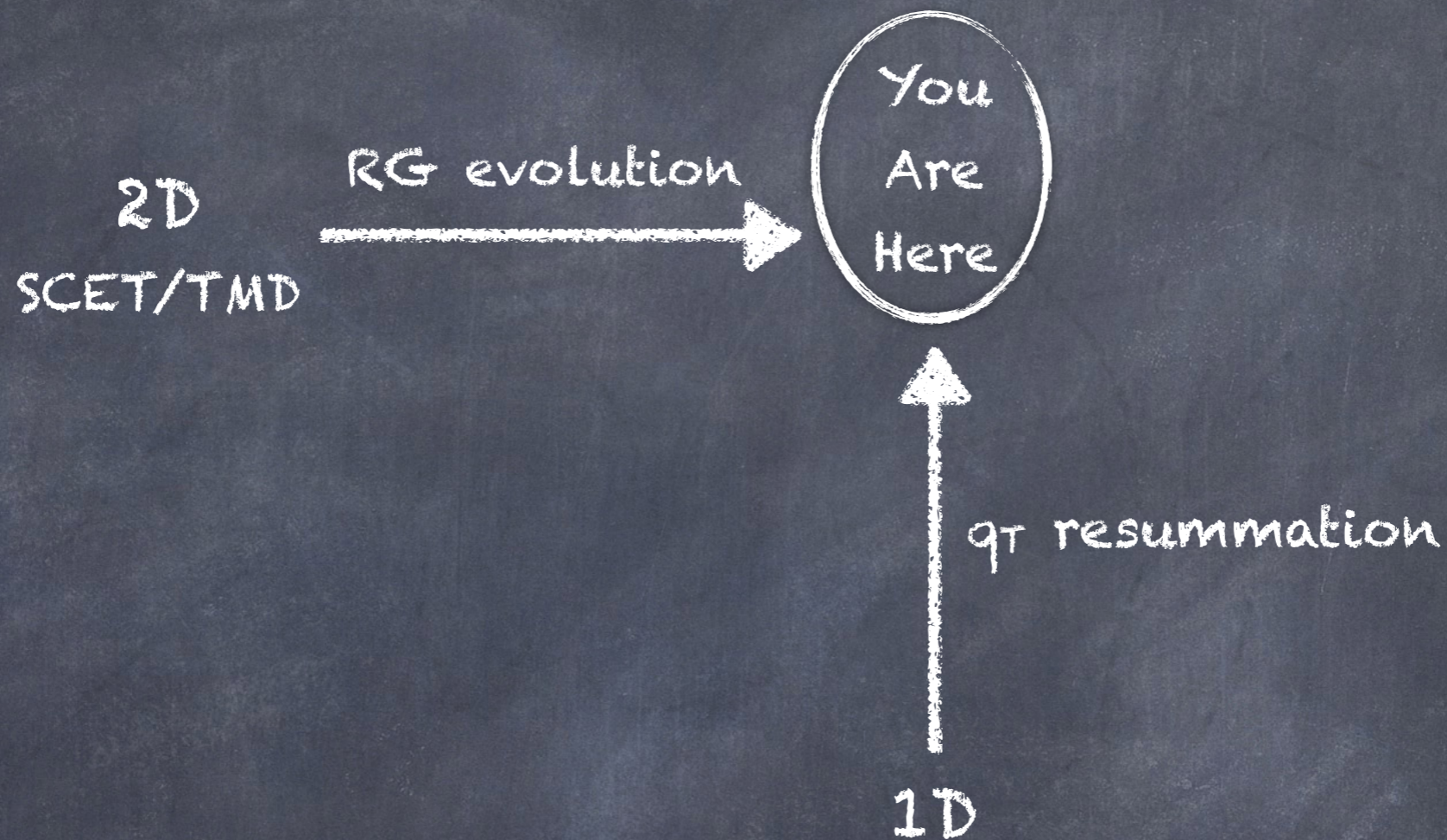
- resummation (i.e. factorisation and exponentiation) of soft-gluon emissions to all orders

- dynamical ($|M_{n\text{-gluons}}| \sim |M_{1\text{-gluon}}|^n$) and kinematical ($PS_{n\text{-gluons}} \sim \prod_n PS_{1\text{-gluon}}$) factorisation properties of QCD

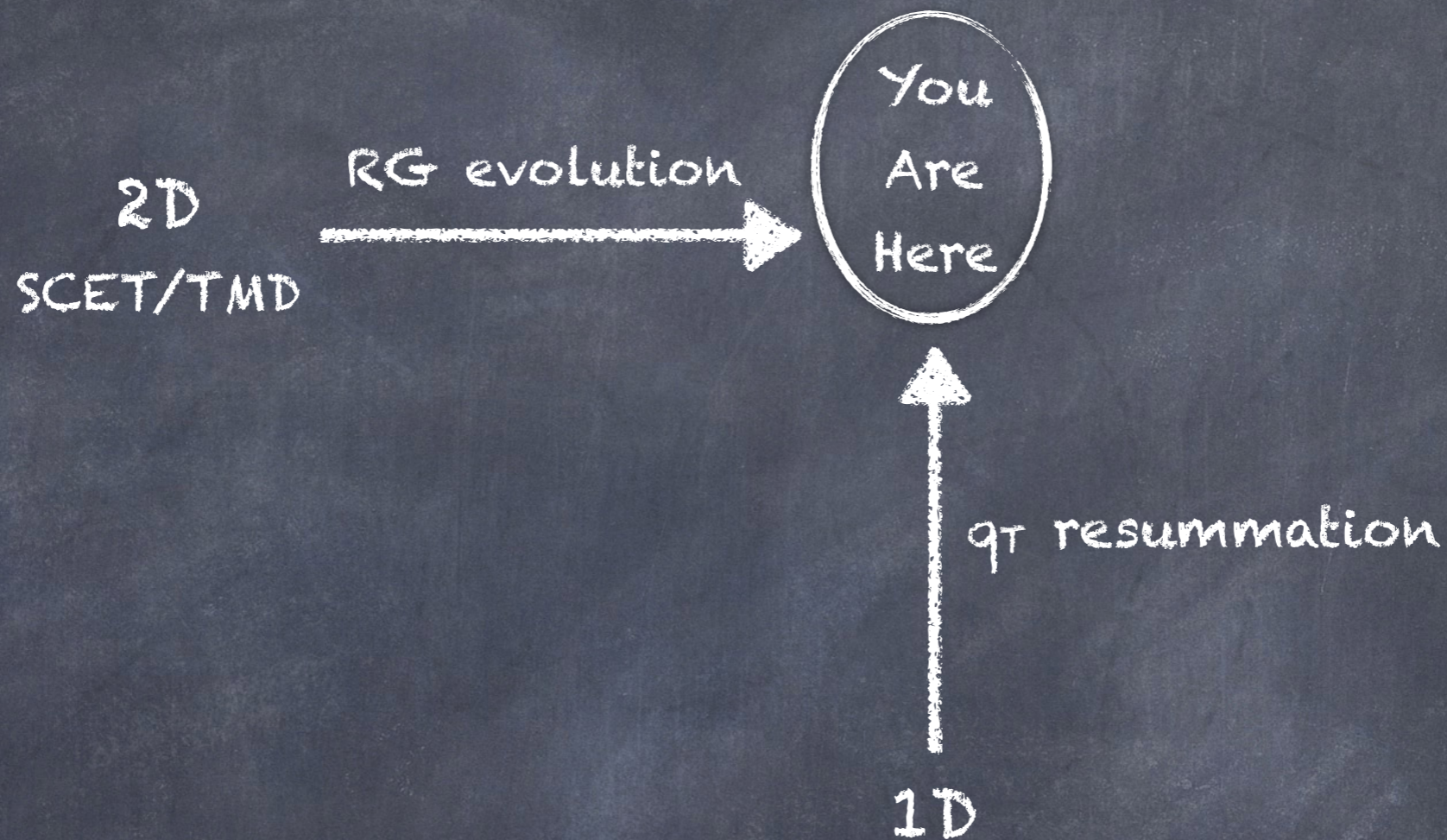
- identical formula as in the previous slide

$$\sigma \sim [C \otimes f(x)] \text{Hard} [C \otimes f(x)] \exp(S) \exp(S_{NP})$$

For hadroproduction of colourless final states:

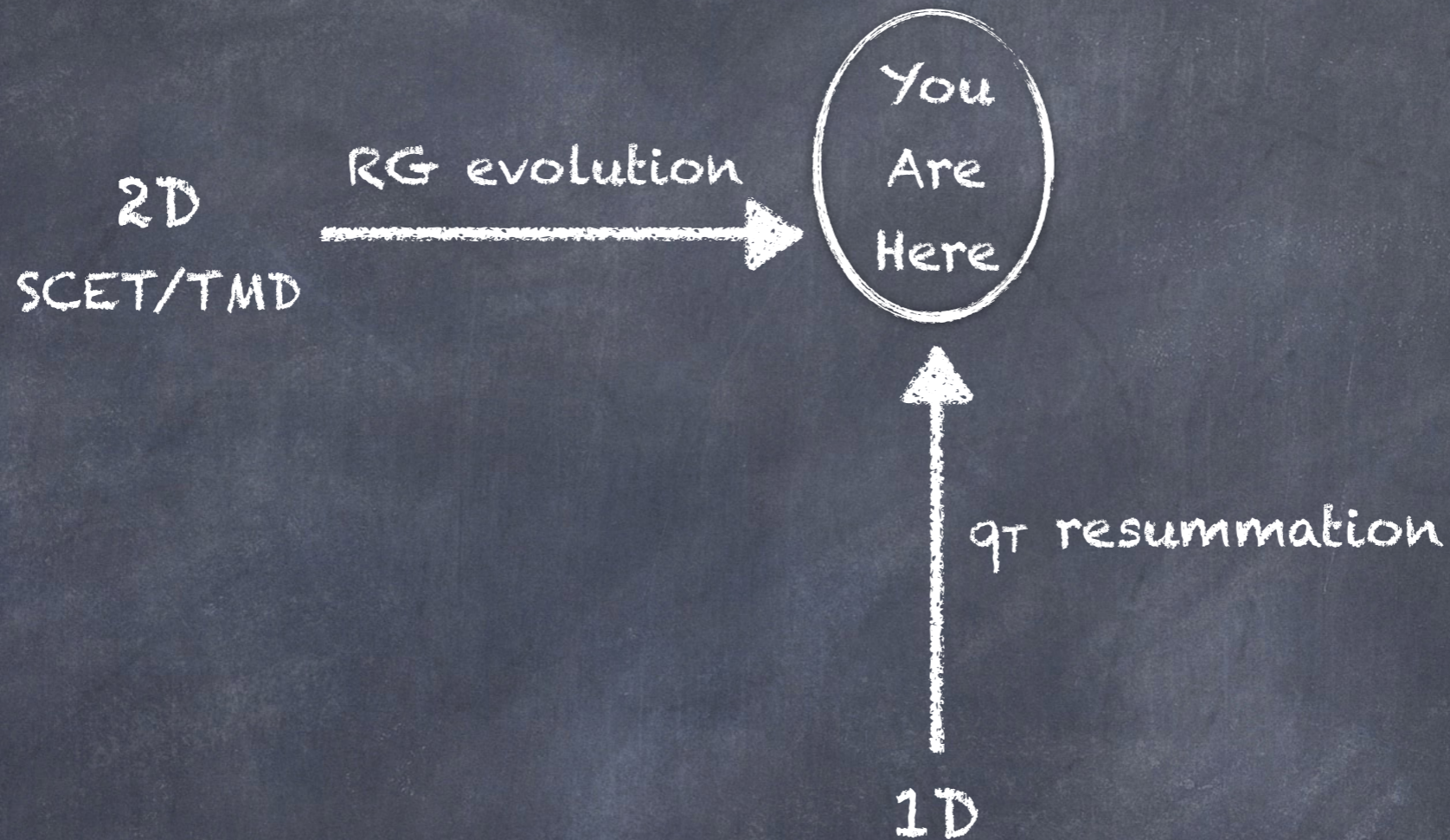


For hadroproduction of colourless final states:



$$\sigma \sim [C \otimes f(x)]_{\text{Hard}} [C \otimes f(x)] \exp(S) \exp(S_{NP})$$

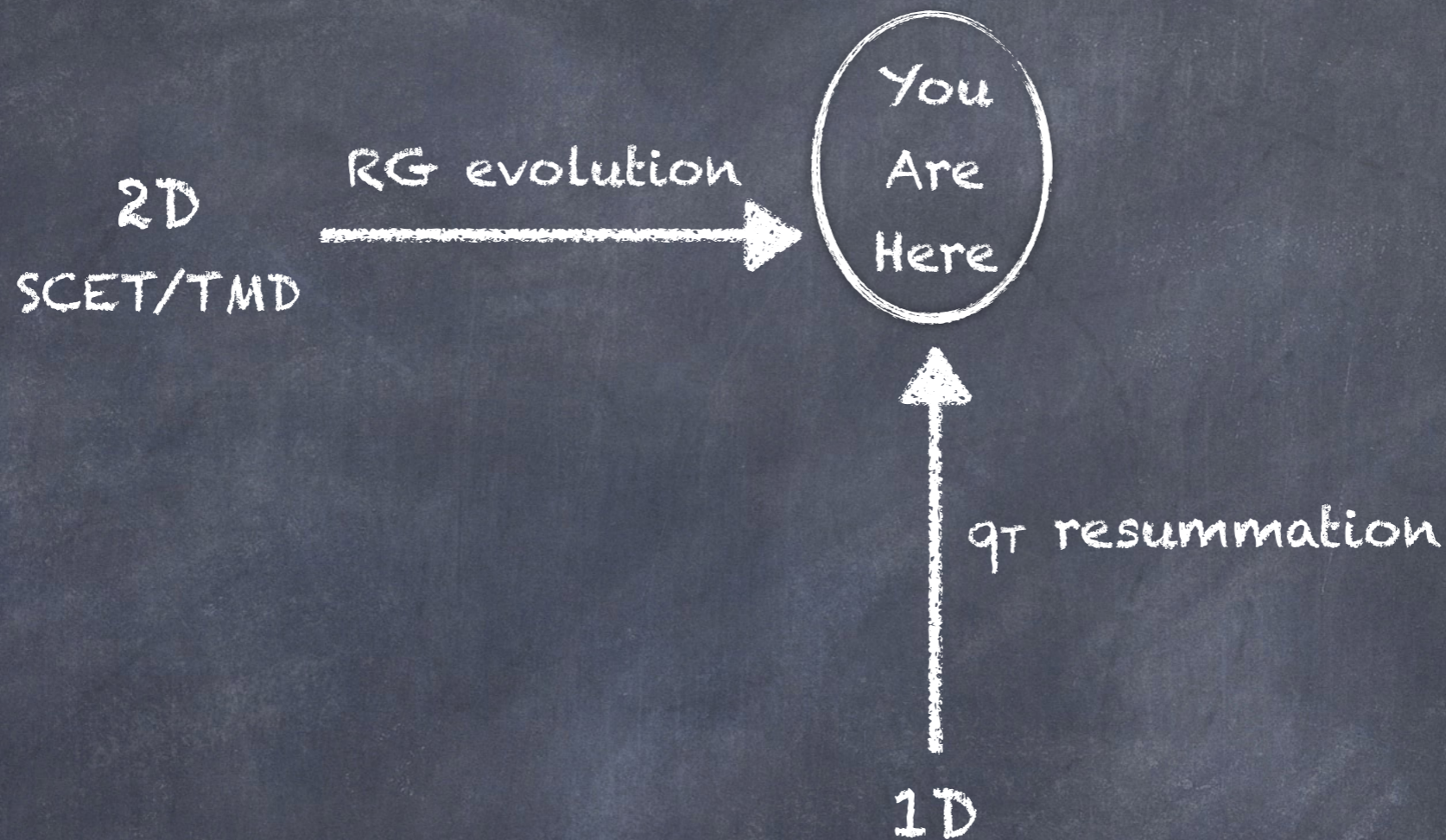
For hadroproduction of colourless final states:



$$\sigma \sim [C \otimes f(x)]_{Hard} [C \otimes f(x)] \exp(S) \exp(S_{NP})$$

1st order OPE
for TMD operator/
Collinear emissions
("matching"/"Wilson"/
"collinear"/"boundary")

For hadroproduction of colourless final states:

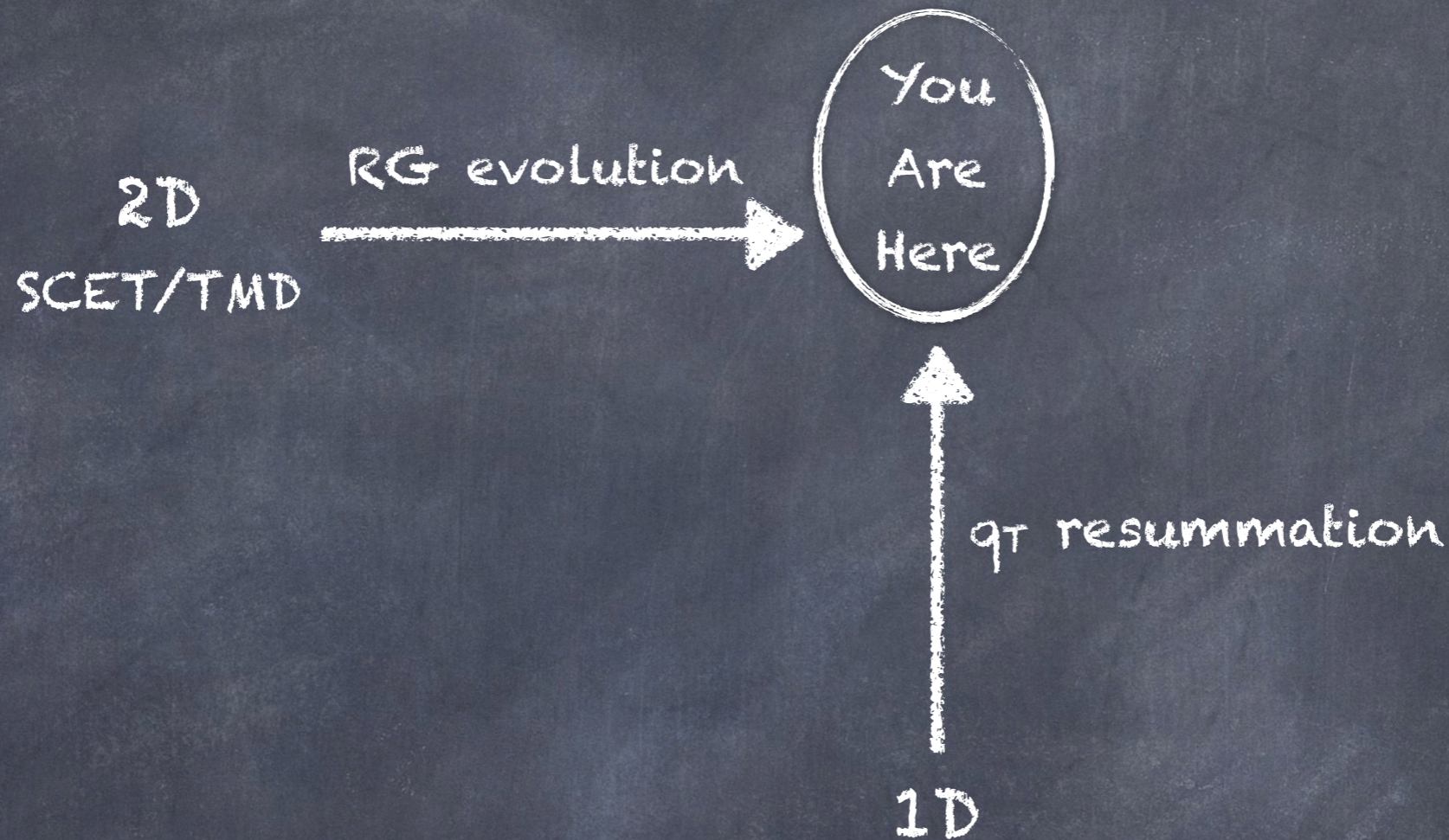


$$\sigma \sim [C \otimes f(x)]_{Hard} [C \otimes f(x)] \exp(S) \exp(S_{NP})$$

1st order OPE
for TMD operator/
Collinear emissions
("matching"/"Wilson"/
"collinear"/"boundary")

("Sudakov": soft emissions
contains "cusp"/"A"
and "non-cusp"/"B"
anomalous dimensions)

For hadroproduction of colourless final states:



$$\sigma \sim [C \otimes f(x)]_{Hard} [C \otimes f(x)] \exp(S) \exp(S_{NP})$$

1st order OPE
for TMD operator/
Collinear emissions
("matching"/"Wilson"/
"collinear"/"boundary")

("Sudakov": soft emissions
contains "cusp"/"A"
and "non-cusp"/"B"
anomalous dimensions)

necessary input:
RG evol. eq. needs initial condition
(to be determined from data)

Parlow Branching

Parton Branching

- parton-shower based: evolution equation with Sudakov factor denoting probability of no-(resolvable)branching

Parton Branching

- parton-shower based: evolution equation with Sudakov factor denoting probability of no-(resolvable)branching
- difference w.r.t. customary parton-shower: forward evolution (from hadron scale to hard scale) instead of backward evolution

Parton Branching

- parton-shower based: evolution equation with Sudakov factor denoting probability of no-(resolvable)branching
- difference w.r.t. customary parton-shower: forward evolution (from hadron scale to hard scale) instead of backward evolution
- angular-ordered emissions from initial parton \rightarrow non-ordered emissions give subleading logs

Parton Branching

- parton-shower based: evolution equation with Sudakov factor denoting probability of no-(resolvable)branching
- difference w.r.t. customary parton-shower: forward evolution (from hadron scale to hard scale) instead of backward evolution
- angular-ordered emissions from initial parton \rightarrow non-ordered emissions give subleading logs
- possible to prove formal equivalence with b-space formalism at various accuracies

Codes

- SCET: SCETLib, CuTe
- TMD: ResBos2, NangaParbat
- qT resummation: DYRes/DYTURBO, ReSolve
- shower-like: RadISH, PartonBranching

Codes

- SCET: SCETLib, CuTe
- TMD: ResBos2, NangaParbat
- qT resummation: DYRes/DYTURBO, ReSolve
- shower-like: RadISH, PartonBranching

Basic ingredients (A,B,C functions) common to all codes.

Codes

- SCET: SCETLib, CuTe
- TMD: ResBos2, NangaParbat
- qT resummation: DYRes/DYTURBO, ReSolve
- shower-like: RadISH, PartonBranching

Basic ingredients (A,B,C functions) common to all codes.

Main differences in:

Codes

- SCET: SCETLib, CuTe
- TMD: ResBos2, NangaParbat
- qt resummation: DYRes/DYTURBO, ReSolve
- shower-like: RadISH, PartonBranching

Basic ingredients (A,B,C functions) common to all codes.

Main differences in:

- working space for resummation/evolution (bT or qT)

Codes

- SCET: SCETLib, CuTe
- TMD: ResBos2, NangaParbat
- qT resummation: DYRes/DYTURBO, ReSolve
- shower-like: RadISH, PartonBranching

Basic ingredients (A,B,C functions) common to all codes.

Main differences in:

- working space for resummation/evolution (bT or qT)
- dealing with NP-physics (prescription/cutoff and intrinsic-kT)

Codes

- SCET: SCETLib, CuTe
- TMD: ResBos2, NangaParbat
- q_T resummation: DYRes/DYTURBO, ReSolve
- shower-like: RadISH, PartonBranching

Basic ingredients (A,B,C functions) common to all codes.

Main differences in:

- working space for resummation/evolution (b_T or q_T)
- dealing with NP-physics (prescription/cutoff and intrinsic- k_T)
- matching with fixed order at intermediate q_T

Differences

Differences

Differences

• NP-physics (1): avoiding the Landau pole

Differences

◉ NP-physics (1): avoiding the Landau pole

◉ qT-space: low-qT cutoff

Differences

◉ NP-physics (1): avoiding the Landau pole

◉ q_T-space: low-q_T cutoff

◉ b_T-space needs either a "freezing"/"saturation" of b (MANY CHOICES for "b*") or a "minimal prescription" (suitable integration path in complex-b plane)

Differences

⊙ NP-physics (1): avoiding the Landau pole

⊙ q_T-space: low-q_T cutoff

⊙ b_T-space needs either a "freezing"/"saturation" of b (MANY CHOICES for "b*") or a "minimal prescription" (suitable integration path in complex-b plane)

⊙ NP-physics (2): intrinsic-κ_T effects

Differences

- ◉ NP-physics (1): avoiding the Landau pole

 - ◉ q_T -space: low- q_T cutoff

 - ◉ b_T -space needs either a "freezing"/"saturation" of b (MANY CHOICES for " b_* ") or a "minimal prescription" (suitable integration path in complex- b plane)

- ◉ NP-physics (2): intrinsic- k_T effects

 - ◉ NP form factor to be determined from data, in principle kinematics- and flavour-dependent

Differences

- ◉ NP-physics (1): avoiding the Landau pole

 - ◉ q_T -space: low- q_T cutoff

 - ◉ b_T -space needs either a "freezing"/"saturation" of b (MANY CHOICES for " b_* ") or a "minimal prescription" (suitable integration path in complex- b plane)

- ◉ NP-physics (2): intrinsic- k_T effects

 - ◉ NP form factor to be determined from data, in principle kinematics- and flavour-dependent

- ◉ matching with fixed order at intermediate q_T

Differences

◉ NP-physics (1): avoiding the Landau pole

◉ q_T -space: low- q_T cutoff

◉ b_T -space needs either a "freezing"/"saturation" of b (MANY CHOICES for " b_* ") or a "minimal prescription" (suitable integration path in complex- b plane)

◉ NP-physics (2): intrinsic- k_T effects

◉ NP form factor to be determined from data, in principle kinematics- and flavour-dependent

◉ matching with fixed order at intermediate q_T

◉ multiplicative $\sigma_{res} \left[\frac{\sigma_{fix}}{\sigma_{res}} \right]_{expanded}$ or additive $\sigma_{res} + \sigma_{fix} - \sigma_{asy}$

Differences

◉ NP-physics (1): avoiding the Landau pole

◉ q_T -space: low- q_T cutoff

◉ b_T -space needs either a "freezing"/"saturation" of b (MANY CHOICES for " b_* ") or a "minimal prescription" (suitable integration path in complex- b plane)

◉ NP-physics (2): intrinsic- k_T effects

◉ NP form factor to be determined from data, in principle kinematics- and flavour-dependent

◉ matching with fixed order at intermediate q_T

◉ multiplicative $\sigma_{res} \left[\frac{\sigma_{fix}}{\sigma_{res}} \right]_{expanded}$ or additive $\sigma_{res} + \sigma_{fix} - \sigma_{asy}$

◉ damping function to switch off resummation/evolution (MANY choices)

Differences

◉ NP-physics (1): avoiding the Landau pole

◉ q_T -space: low- q_T cutoff

◉ b_T -space needs either a "freezing"/"saturation" of b (MANY CHOICES for " b_* ") or a "minimal prescription" (suitable integration path in complex- b plane)

◉ NP-physics (2): intrinsic- k_T effects

◉ NP form factor to be determined from data, in principle kinematics- and flavour-dependent

◉ matching with fixed order at intermediate q_T

◉ multiplicative $\sigma_{res} \left[\frac{\sigma_{fix}}{\sigma_{res}} \right]_{expanded}$ or additive $\sigma_{res} + \sigma_{fix} - \sigma_{asy}$

◉ damping function to switch off resummation/evolution (MANY choices)

◉ unitarity enforcing ($\int \frac{d\sigma}{dq_T} dq_T = \sigma$), i.e. modified logs (damping function and NP may spoil it!)

Differences

• NP-physics (1): avoiding the Landau pole

• q_T -space: low- q_T cutoff

• b_T -space needs either a "freezing"/"saturation" of b (MANY CHOICES for " b_* ") or a "minimal prescription" (suitable integration path in complex- b plane)

• NP-physics (2): intrinsic- k_T effects

• NP form factor to be determined from data, in principle kinematics- and flavour-dependent

• matching with fixed order at intermediate q_T

• multiplicative $\sigma_{res} \left[\frac{\sigma_{fix}}{\sigma_{res}} \right]_{expanded}$ or additive $\sigma_{res} + \sigma_{fix} - \sigma_{asy}$

• damping function to switch off resummation/evolution (MANY choices)

• unitarity enforcing ($\int \frac{d\sigma}{dq_T} dq_T = \sigma$), i.e. modified logs (damping function and NP may spoil it!)

• lepton cuts

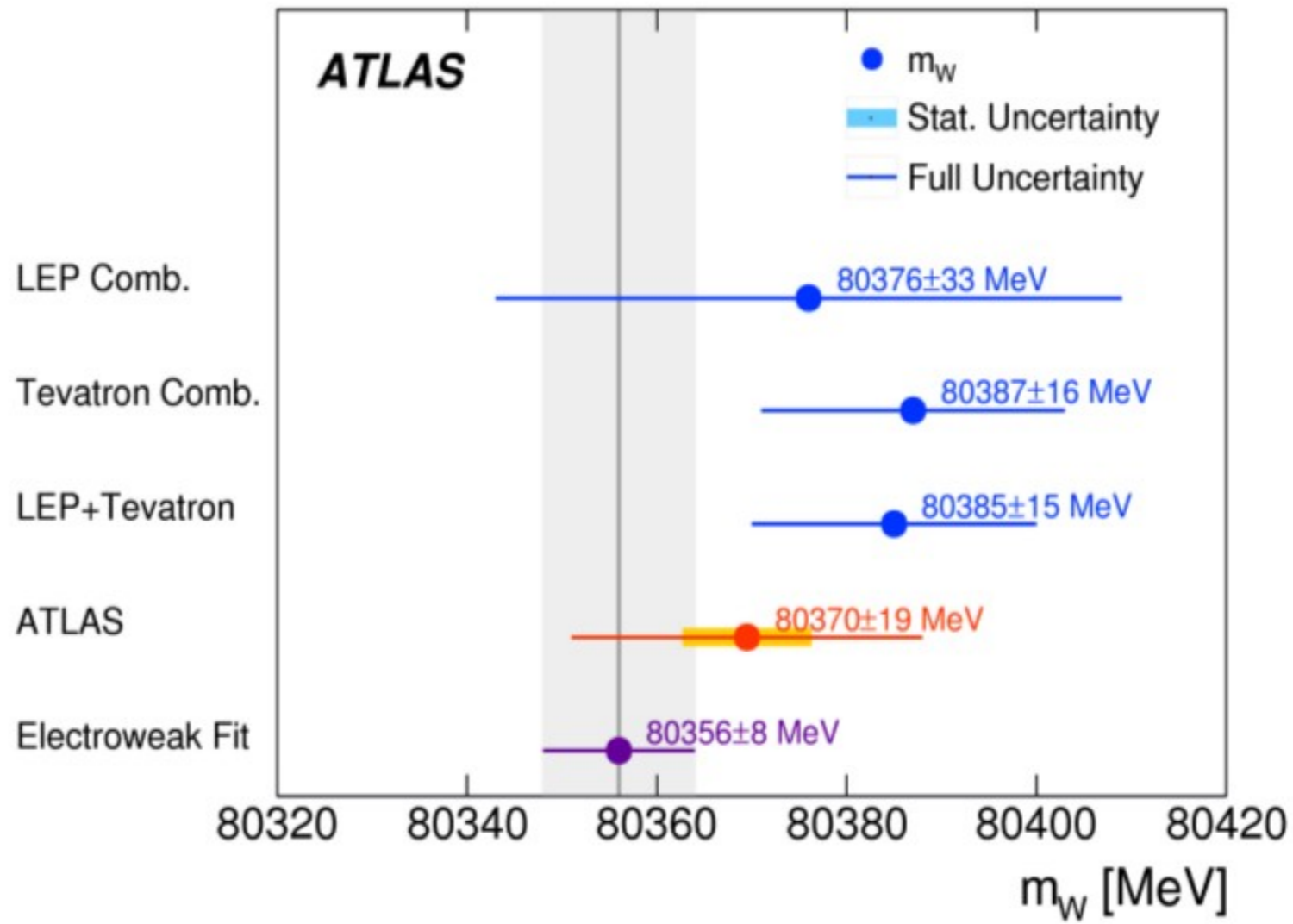
Impact on precision measurements at the LHC: the W mass case

in collaboration with:

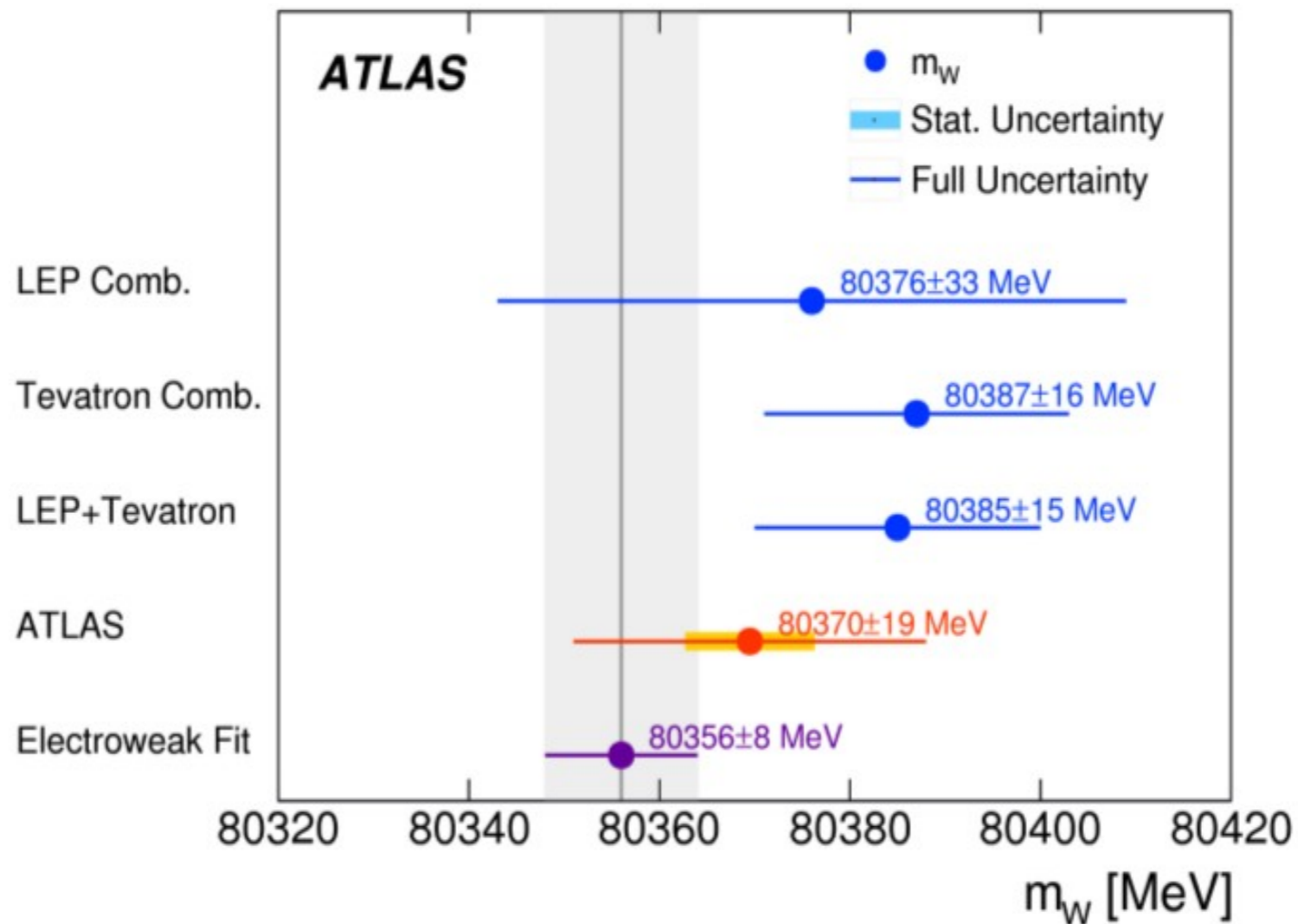
A. Bacchetta (Pavia), M. Radici (Pavia), A. Signori (Argonne)

arXiv:1807.02101 - Phys.Lett. B788 (2019) 542-545

The W mass

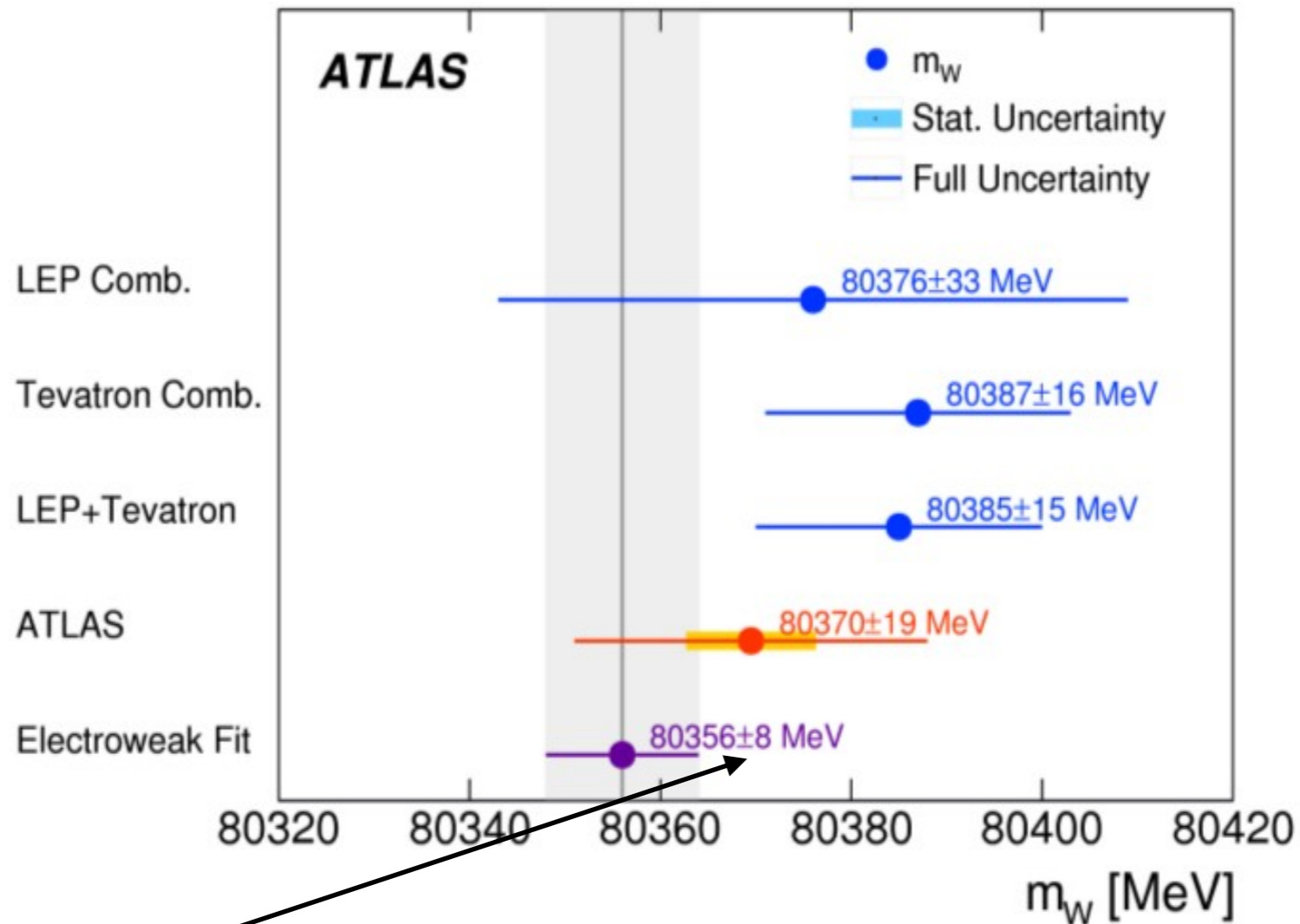


The W mass



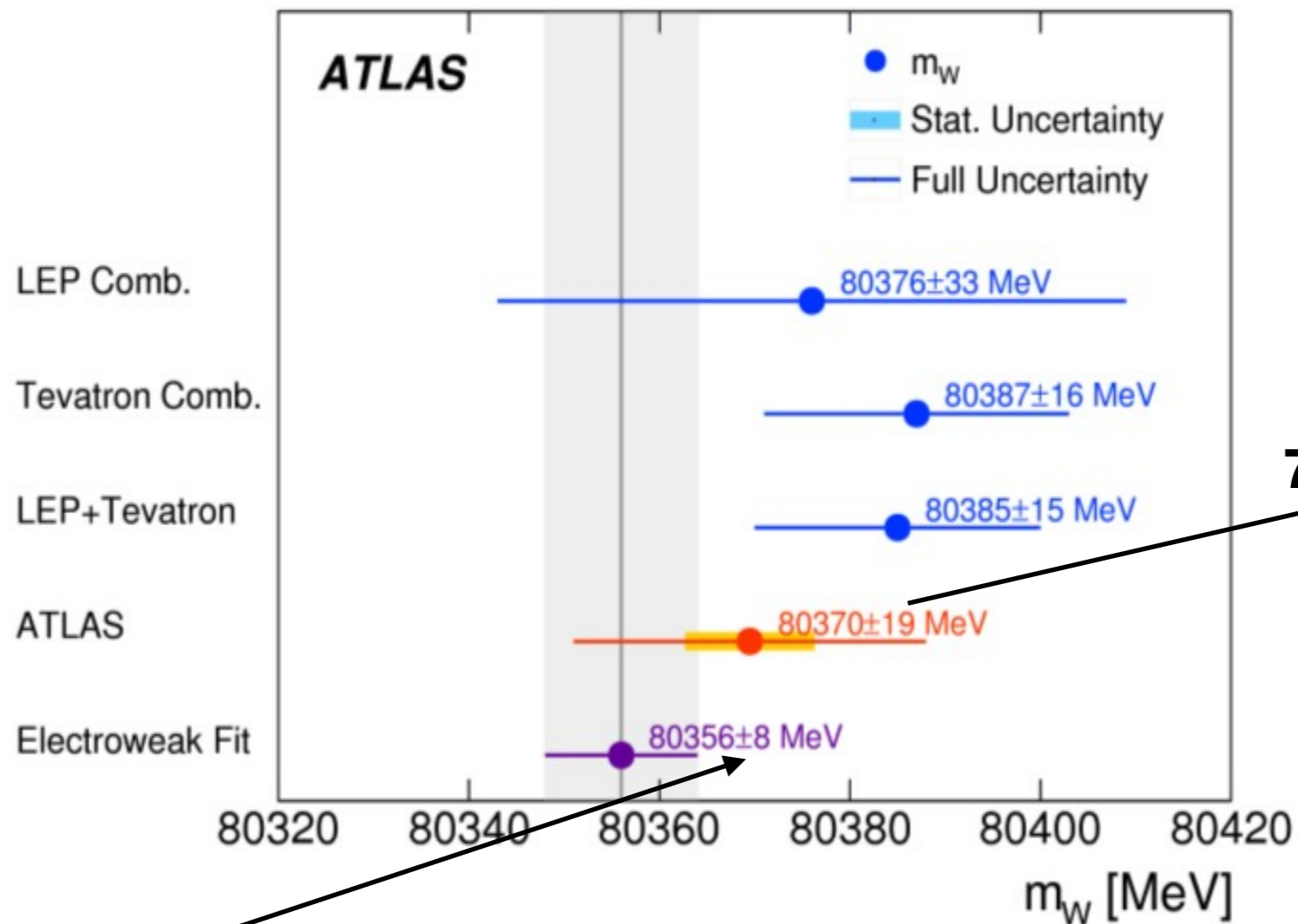
The determination of the W -boson mass from the global fit of the electroweak parameters has an uncertainty of 8 MeV, which sets a natural target for the precision of the experimental measurement of the mass of the W boson. The modelling uncertainties, which currently dominate the overall uncertainty on the m_W measurement presented in this note, need to be reduced in order to fully exploit the larger data samples available at centre-of-mass energies of 8 and 13 TeV. A better knowledge of the PDFs, as achievable with the inclusion in PDF fits of recent precise measurements of W - and Z -boson rapidity cross sections with the ATLAS detector [41], and improved QCD and electroweak predictions for Drell-Yan production, are therefore crucial for future measurements of the W -boson mass at the LHC. **ATLAS**, EPJC 78, 110 (2018)

The W mass



The determination of the W -boson mass from the global fit of the electroweak parameters has an uncertainty of 8 MeV, which sets a natural target for the precision of the experimental measurement of the mass of the W boson. The modelling uncertainties, which currently dominate the overall uncertainty on the m_W measurement presented in this note, need to be reduced in order to fully exploit the larger data samples available at centre-of-mass energies of 8 and 13 TeV. A better knowledge of the PDFs, as achievable with the inclusion in PDF fits of recent precise measurements of W - and Z -boson rapidity cross sections with the ATLAS detector [41], and improved QCD and electroweak predictions for Drell-Yan production, are therefore crucial for future measurements of the W -boson mass at the LHC. **ATLAS**, EPJC 78, 110 (2018)

The W mass



7 stat, 11 exp, 14 th

The determination of the W -boson mass from the global fit of the electroweak parameters has an uncertainty of 8 MeV, which sets a natural target for the precision of the experimental measurement of the mass of the W boson. The modelling uncertainties, which currently dominate the overall uncertainty on the m_W measurement presented in this note, need to be reduced in order to fully exploit the larger data samples available at centre-of-mass energies of 8 and 13 TeV. A better knowledge of the PDFs, as achievable with the inclusion in PDF fits of recent precise measurements of W - and Z -boson rapidity cross sections with the ATLAS detector [41], and improved QCD and electroweak predictions for Drell-Yan production, are therefore crucial for future measurements of the W -boson mass at the LHC. ATLAS, EPJC 78, 110 (2018)

The extraction of physical quantities

The extraction of physical quantities

Observables

- accessible via **counting experiments**: cross sections and asymmetries

The extraction of physical quantities

Observables

- accessible via **counting experiments**: cross sections and asymmetries

Pseudo-Observables

- functions of cross sections and symmetries
- **require a model** to be properly defined
 - M_Z at LEP as pole of the Breit-Wigner resonance factor
 - M_W at hadron colliders as fitting parameter of a *template fit* procedure

The extraction of physical quantities

Observables

- accessible via **counting experiments**: cross sections and asymmetries

Pseudo-Observables

- functions of cross sections and symmetries
- **require a model** to be properly defined
 - M_Z at LEP as pole of the Breit-Wigner resonance factor
 - M_W at hadron colliders as fitting parameter of a *template fit* procedure

Template fit

The extraction of physical quantities

Observables

- accessible via **counting experiments**: cross sections and asymmetries

Pseudo-Observables

- functions of cross sections and symmetries
- **require a model** to be properly defined
 - M_Z at LEP as pole of the Breit-Wigner resonance factor
 - M_W at hadron colliders as fitting parameter of a *template fit* procedure

Template fit

1. generate several histograms with highest available theoretical accuracy and best possible detector simulation, and let the fit parameter (e.g. M_W) vary in a range

The extraction of physical quantities

Observables

- accessible via **counting experiments**: cross sections and asymmetries

Pseudo-Observables

- functions of cross sections and symmetries
- **require a model** to be properly defined
 - M_Z at LEP as pole of the Breit-Wigner resonance factor
 - M_W at hadron colliders as fitting parameter of a *template fit* procedure

Template fit

1. generate several histograms with highest available theoretical accuracy and best possible detector simulation, and let the fit parameter (e.g. M_W) vary in a range
2. the histogram that best describes data selects the preferred (*i.e.* measured) M_W

The extraction of physical quantities

Observables

- accessible via **counting experiments**: cross sections and asymmetries

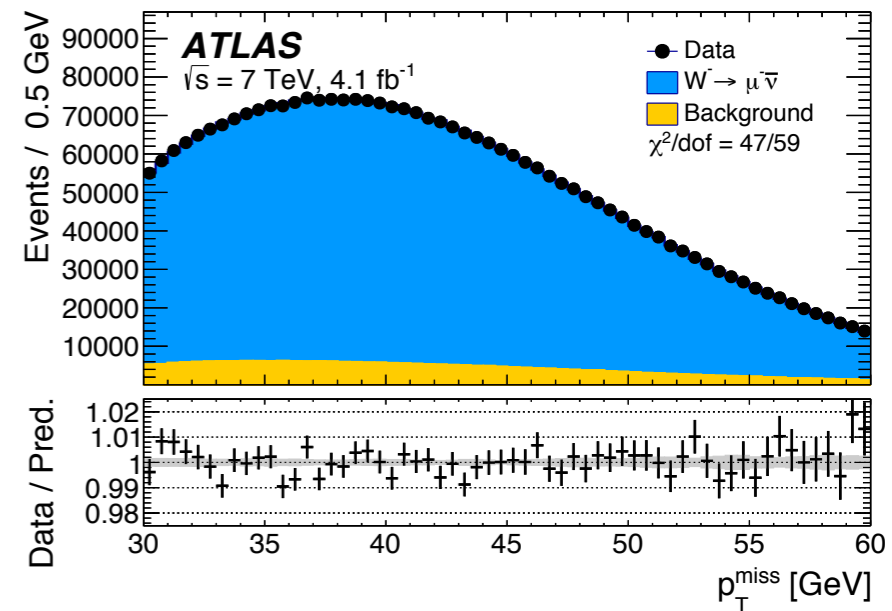
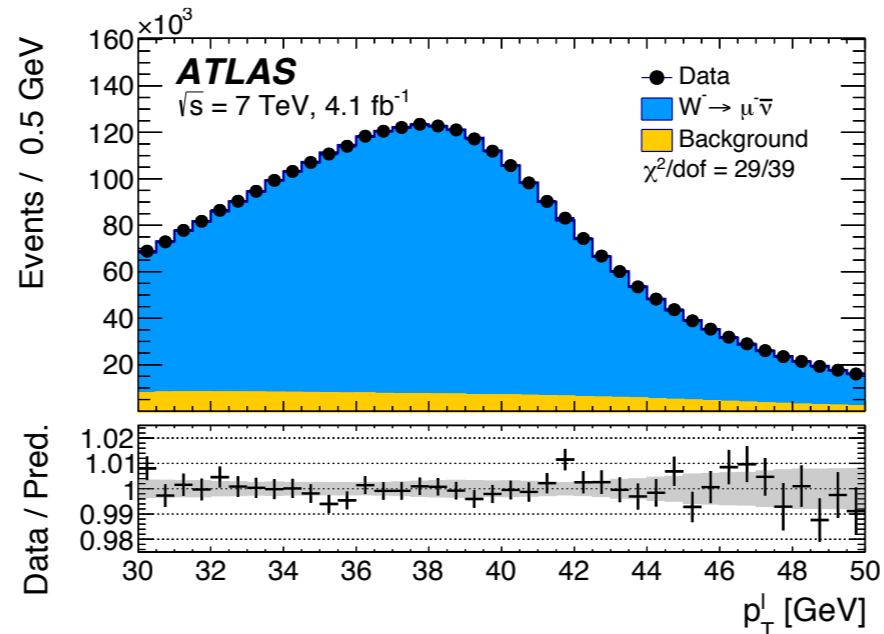
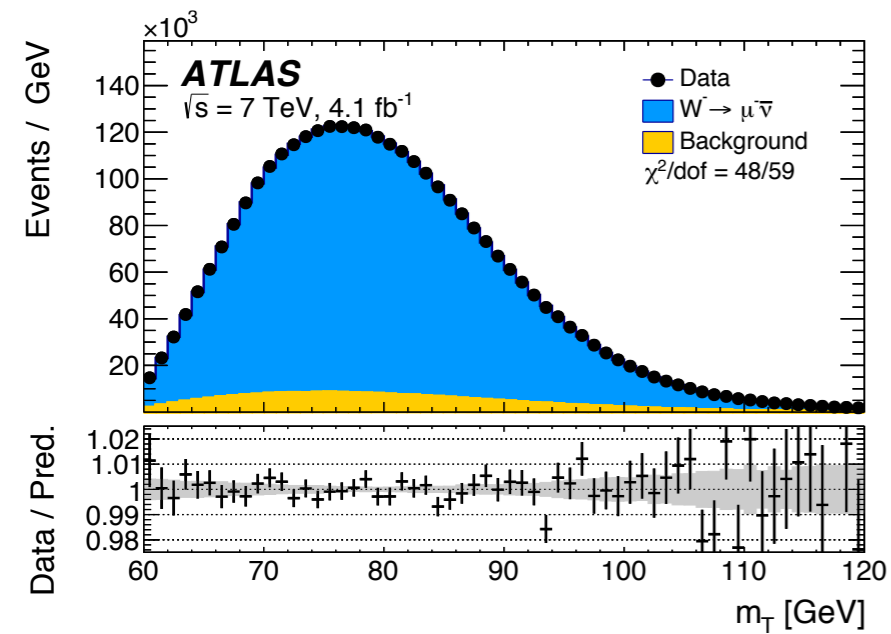
Pseudo-Observables

- functions of cross sections and symmetries
- **require a model** to be properly defined
 - M_Z at LEP as pole of the Breit-Wigner resonance factor
 - M_W at hadron colliders as fitting parameter of a *template fit* procedure

Template fit

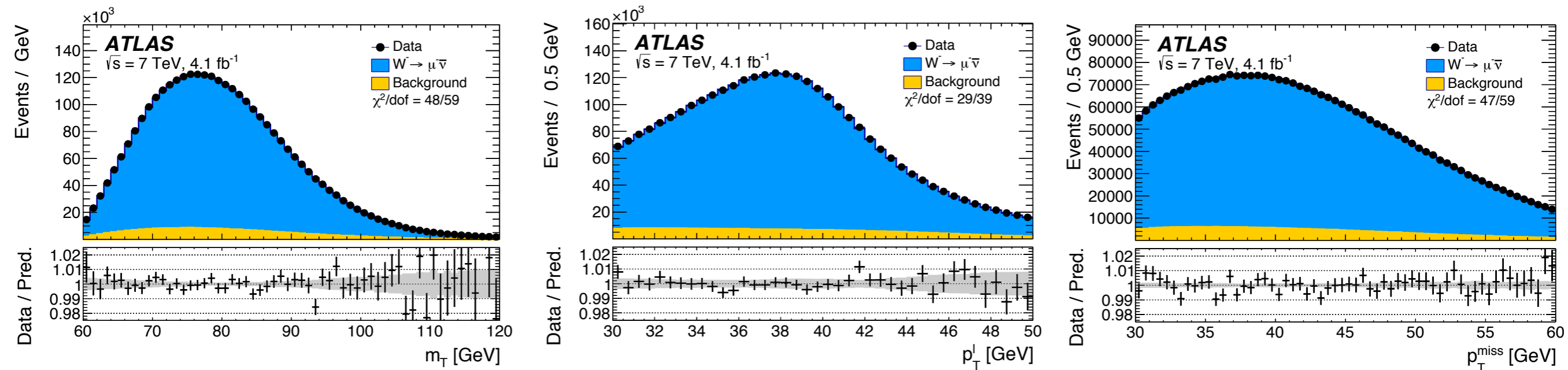
1. generate several histograms with highest available theoretical accuracy and best possible detector simulation, and let the fit parameter (e.g. M_W) vary in a range
 2. the histogram that best describes data selects the preferred (*i.e.* measured) M_W
- ➔ the result of the fit depends on the **hypotheses used to compute the templates** (PDFs, scales, non-perturbative, different prescriptions, ...)
 - ➔ these hypotheses **should be treated as theoretical systematic errors**

Observables and techniques

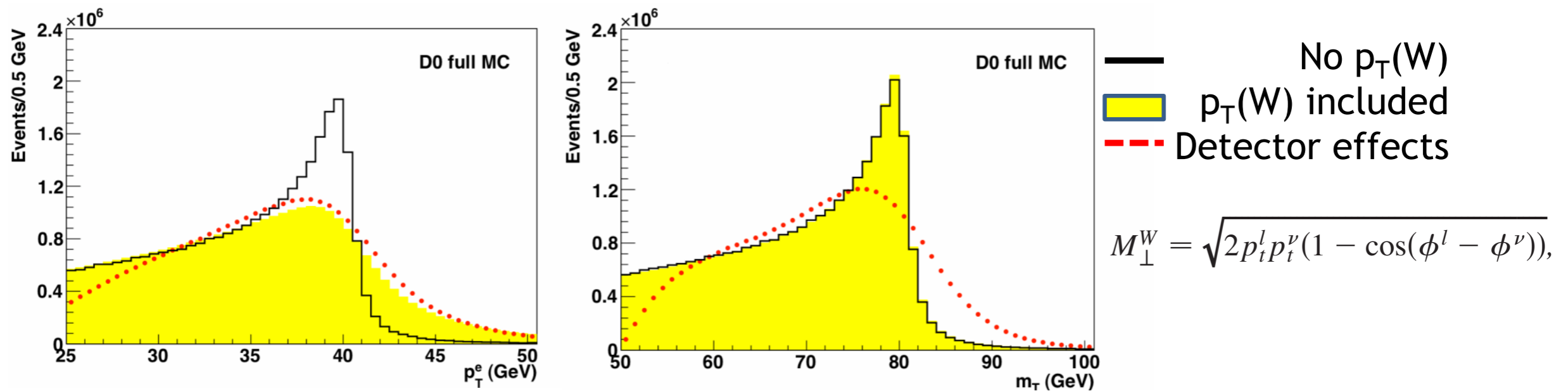


M_W extracted from the study of the **shape** of m_T , p_{Tl} , p_{Tmiss}
jacobian peak enhances sensitivity to M_W

Observables and techniques

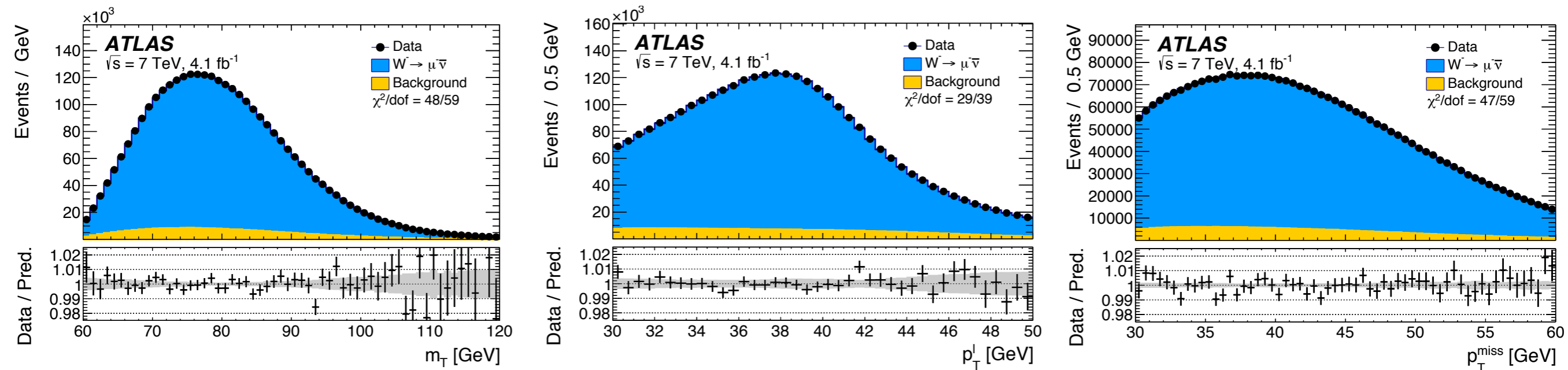


M_W extracted from the study of the **shape** of m_T , p_{Tl} , p_{Tmiss}
Jacobian peak enhances sensitivity to M_W

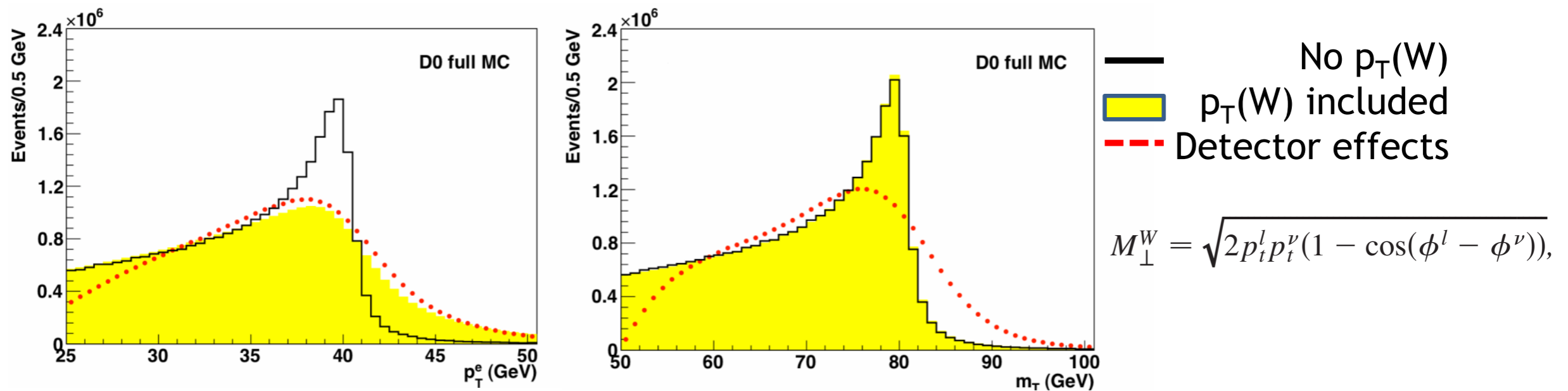


Transverse mass: **important** detector smearing effects, **weakly** sensitive to p_{TW} modelling
 Lepton p_T : **moderate** detector smearing effects, **extremely** sensitive to p_{TW} modelling

Observables and techniques



M_W extracted from the study of the **shape** of m_T , p_{Tl} , p_{Tmiss}
Jacobian peak enhances sensitivity to M_W

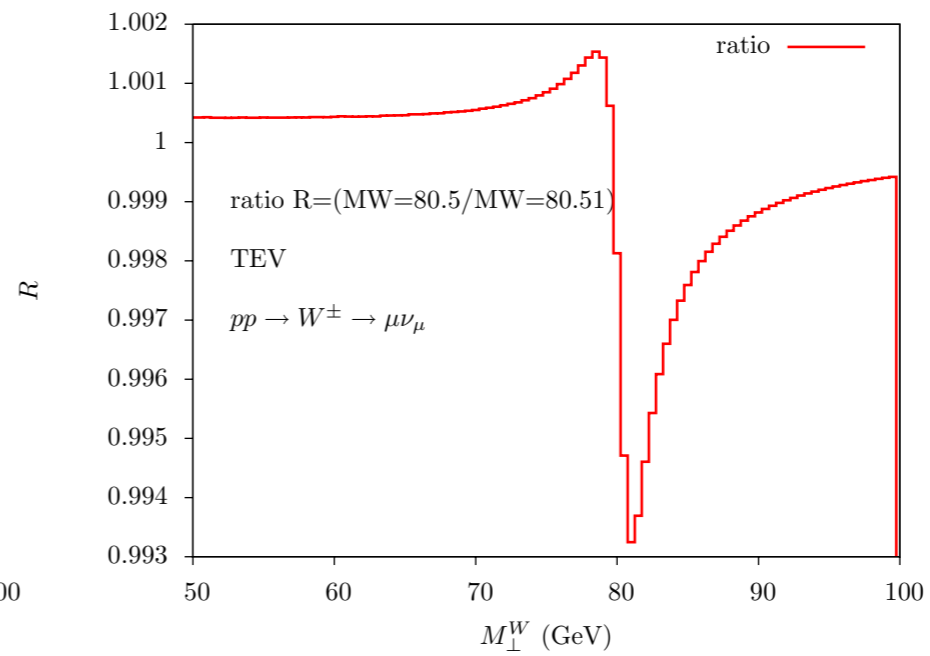
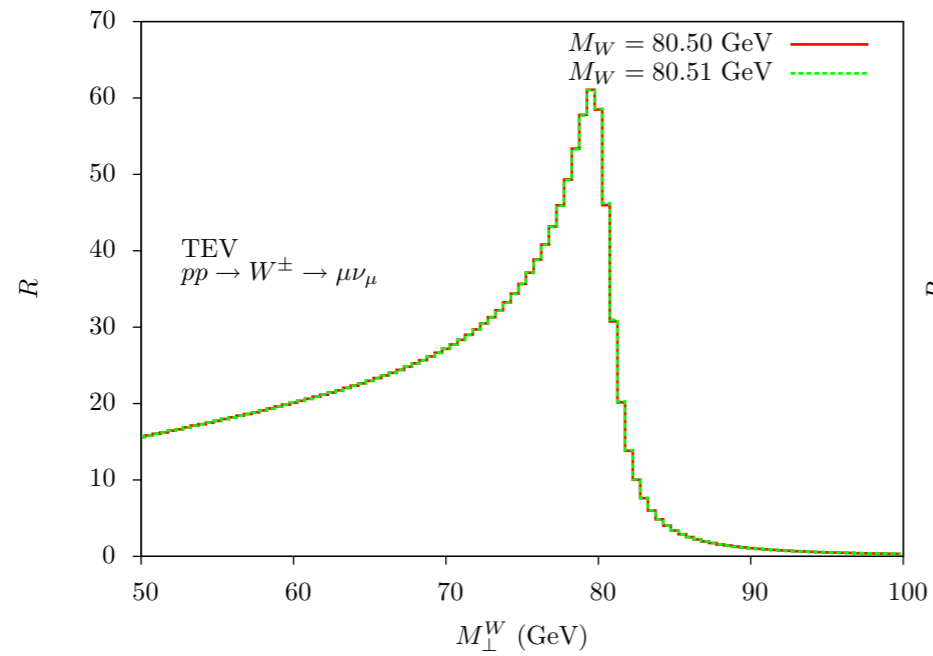


Transverse mass: **important** detector smearing effects, **weakly** sensitive to p_{TW} modelling
 Lepton p_T : **moderate** detector smearing effects, **extremely** sensitive to p_{TW} modelling
 p_{TW} modelling depends on flavour and all-order treatment of QCD corrections

Observables and techniques

Challenging shape measurement: a distortion at the **few per mille** level of the distributions yields a shift of **O(10 MeV)** of the M_W value

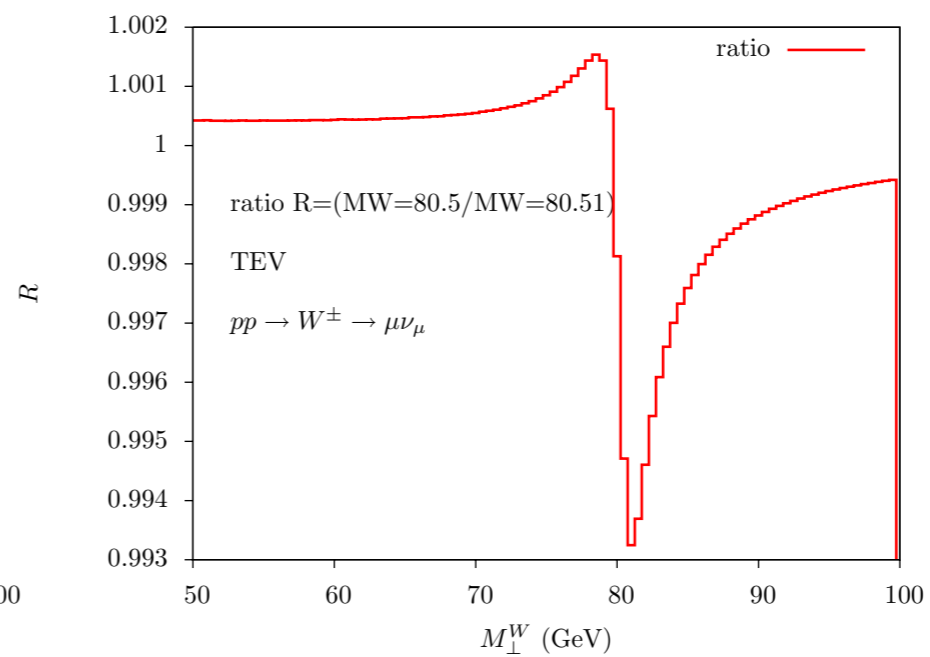
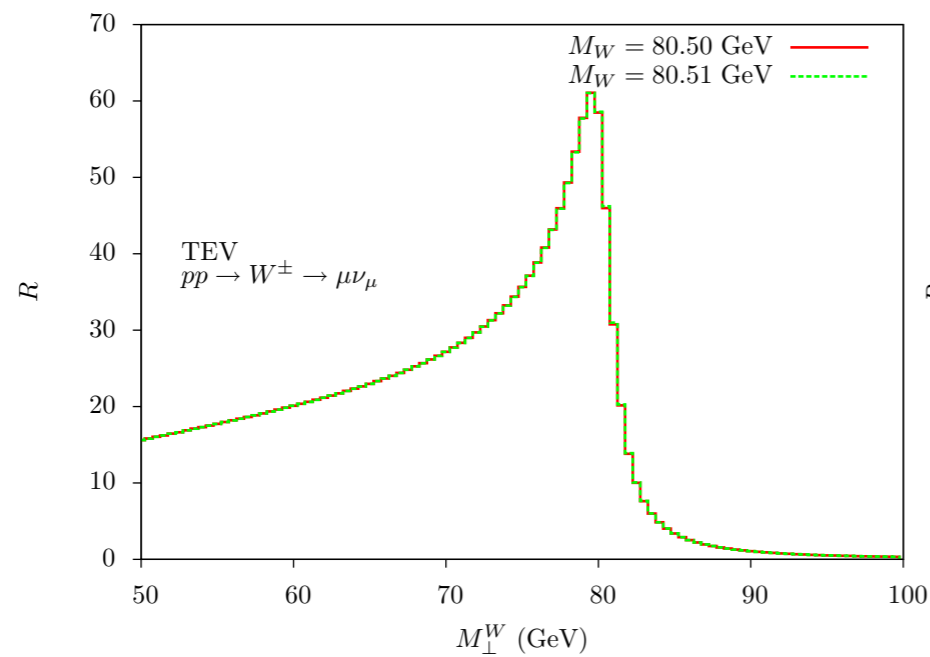
m_T



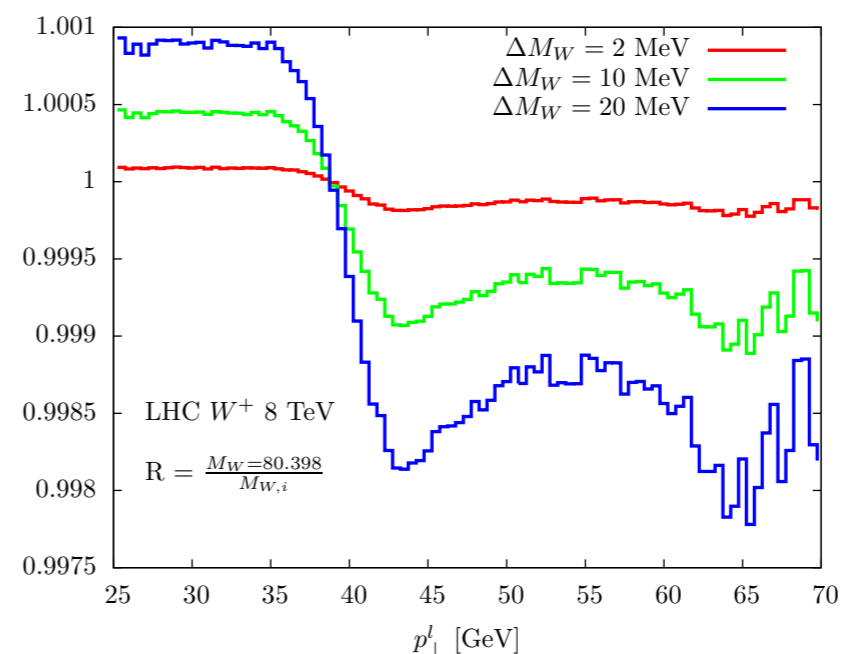
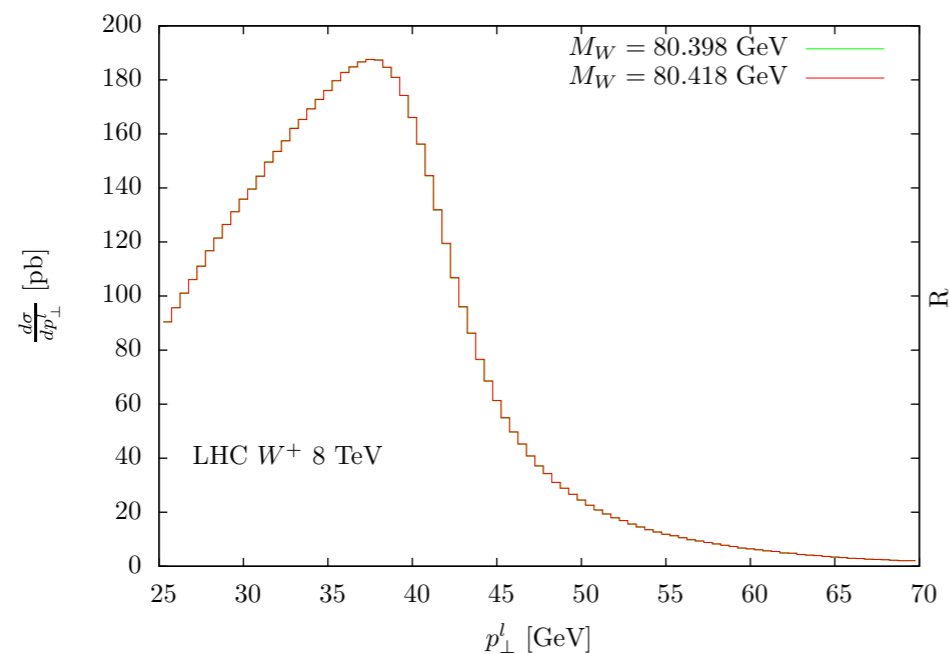
Observables and techniques

Challenging shape measurement: a distortion at the **few per mille** level of the distributions yields a shift of **O(10 MeV)** of the M_W value

m_T



p_{Tl}



Template-fit estimate of theoretical uncertainties (ex:PDF)

Carlone Calame, Montagna, Nicrosini, Treccani PRD 69 (2004)

Bozzi, Rojo, Vicini PRD 83 (2011)

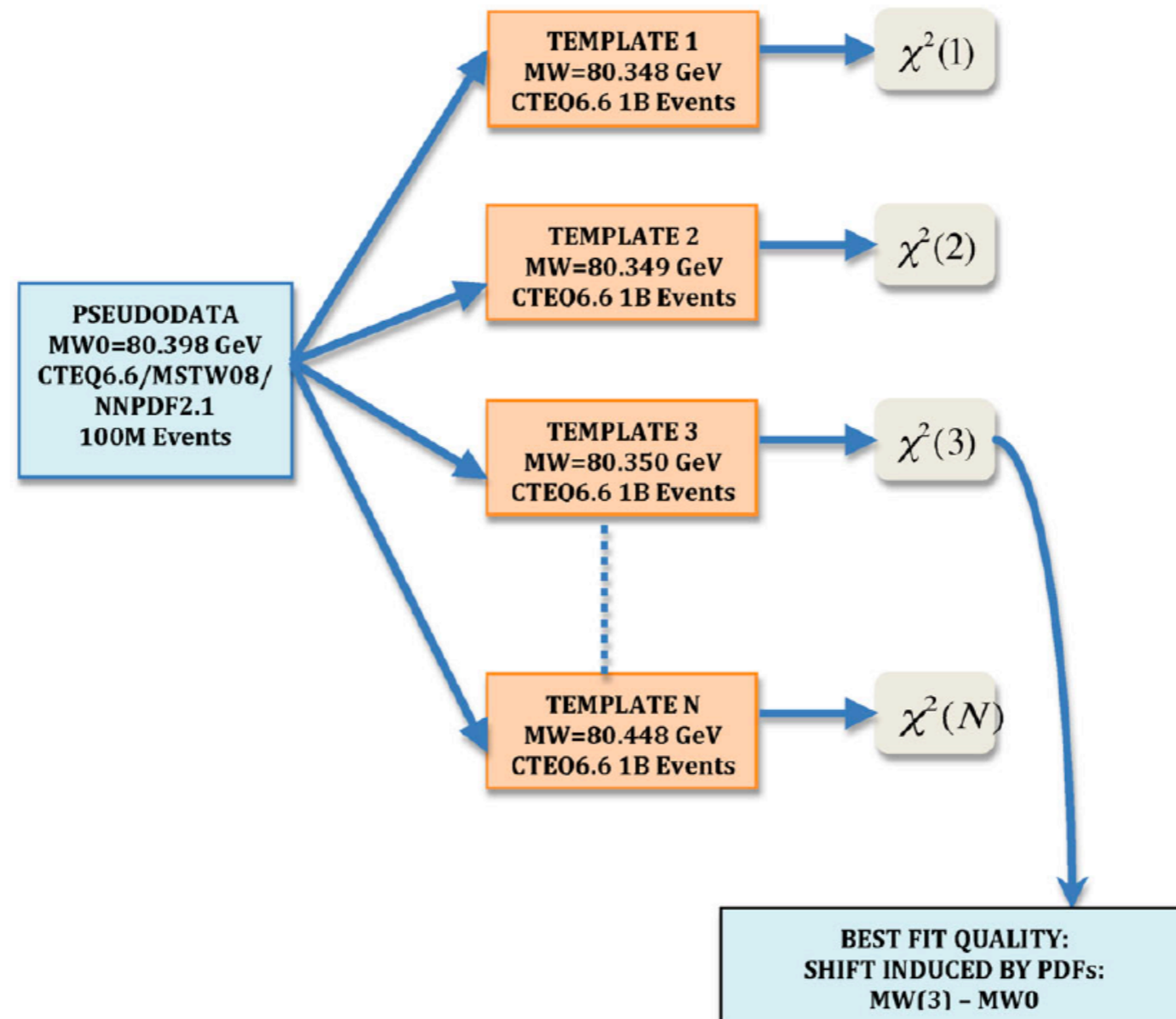
Bozzi, Citelli, Vicini PRD 91 (2015)

Bozzi, Citelli, Vesterinen, Vicini EPJC 75 (2015)

Template-fit estimate of theoretical uncertainties (ex:PDF)

Carloni Calame, Montagna, Nicosini, Treccani PRD 69 (2004)
Bozzi, Rojo, Vicini PRD 83 (2011)
Bozzi, Citelli, Vicini PRD 91 (2015)
Bozzi, Citelli, Vesterinen, Vicini EPJC 75 (2015)

- **pseudodata** with different PDF sets: low-statistics (100M) and fixed M_{W0}
- **templates** with a reference PDF set (CTEQ6.6): high-statistics (1B) and different M_W
- same code used to generate both pseudodata and templates → **only effect probed is the PDF one**



ρ_{TW} and the modelling of intrinsic k_T

p_{TW} and the modelling of intrinsic k_T

- $p_{Tl} \Leftrightarrow p_{TW} \Leftrightarrow$ QCD initial state radiation + intrinsic k_T

p_{TW} and the modelling of intrinsic k_T

- $p_{TI} \Leftrightarrow p_{TW} \Leftrightarrow$ QCD initial state radiation + intrinsic k_T
- Intrinsic k_T effects measured on Z data and used to predict W distributions, *assuming universality*

p_{TW} and the modelling of intrinsic k_T

- $p_{TI} \Leftrightarrow p_{TW} \Leftrightarrow$ QCD initial state radiation + intrinsic k_T
- Intrinsic k_T effects measured on Z data and used to predict W distributions, *assuming universality*

but

different flavour structure

*different phase space
available*

ρ_{TW} and the modelling of intrinsic k_T

- $\rho_{TI} \Leftrightarrow \rho_{TW} \Leftrightarrow$ QCD initial state radiation + intrinsic k_T
- Intrinsic k_T effects measured on Z data and used to predict W distributions, *assuming universality*

but

different flavour structure

different phase space available

—> *different Gaussian factors for different flavours*

$$f_1^{aNP}(b_T^2) \propto e^{-g_{NP}^a b_T^2}$$

~~Flavor and kinematic~~
dependent widths

p_{TW} and the modelling of intrinsic k_T

- $p_{TI} \Leftrightarrow p_{TW} \Leftrightarrow$ QCD initial state radiation + intrinsic k_T
- Intrinsic k_T effects measured on Z data and used to predict W distributions, *assuming universality*

but

different flavour structure

different phase space available

—> *different Gaussian factors for different flavours*

$$f_1^{aNP}(b_T^2) \propto e^{-g_{NP}^a b_T^2}$$

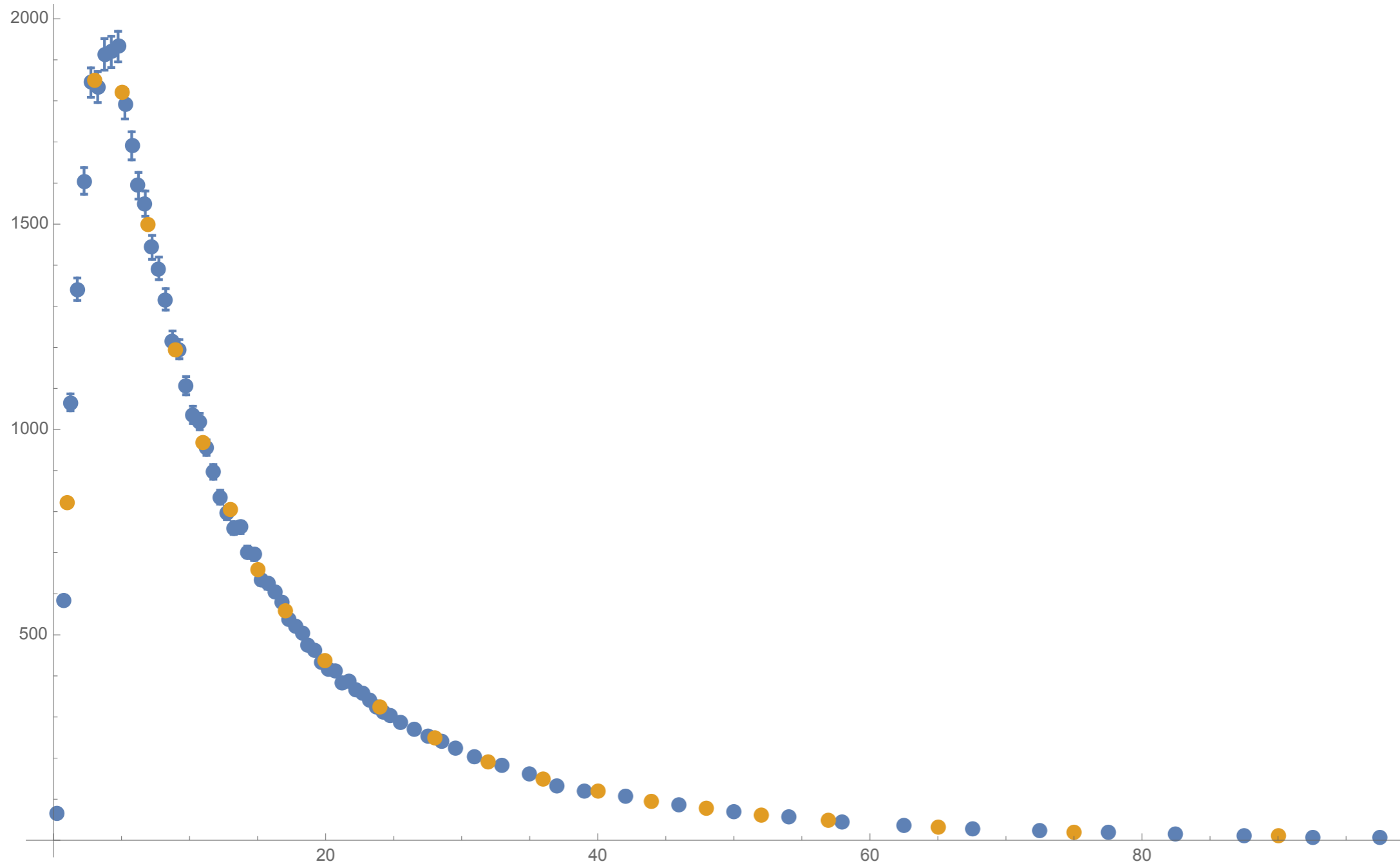
~~Flavor and kinematic dependent widths~~

We consider :

- **50 flavour-dependent sets** $\{g_{NP}^{u_v}, g_{NP}^{d_v}, g_{NP}^{u_s}, g_{NP}^{d_s}, g_{NP}^s\}$ with $g_{NP}^a \in [0.2, 0.6] \text{ GeV}^2$
- **1 flavour-independent set** with $g_{NP}^a = 0.4 \text{ GeV}^2$

“Z-equivalent” sets

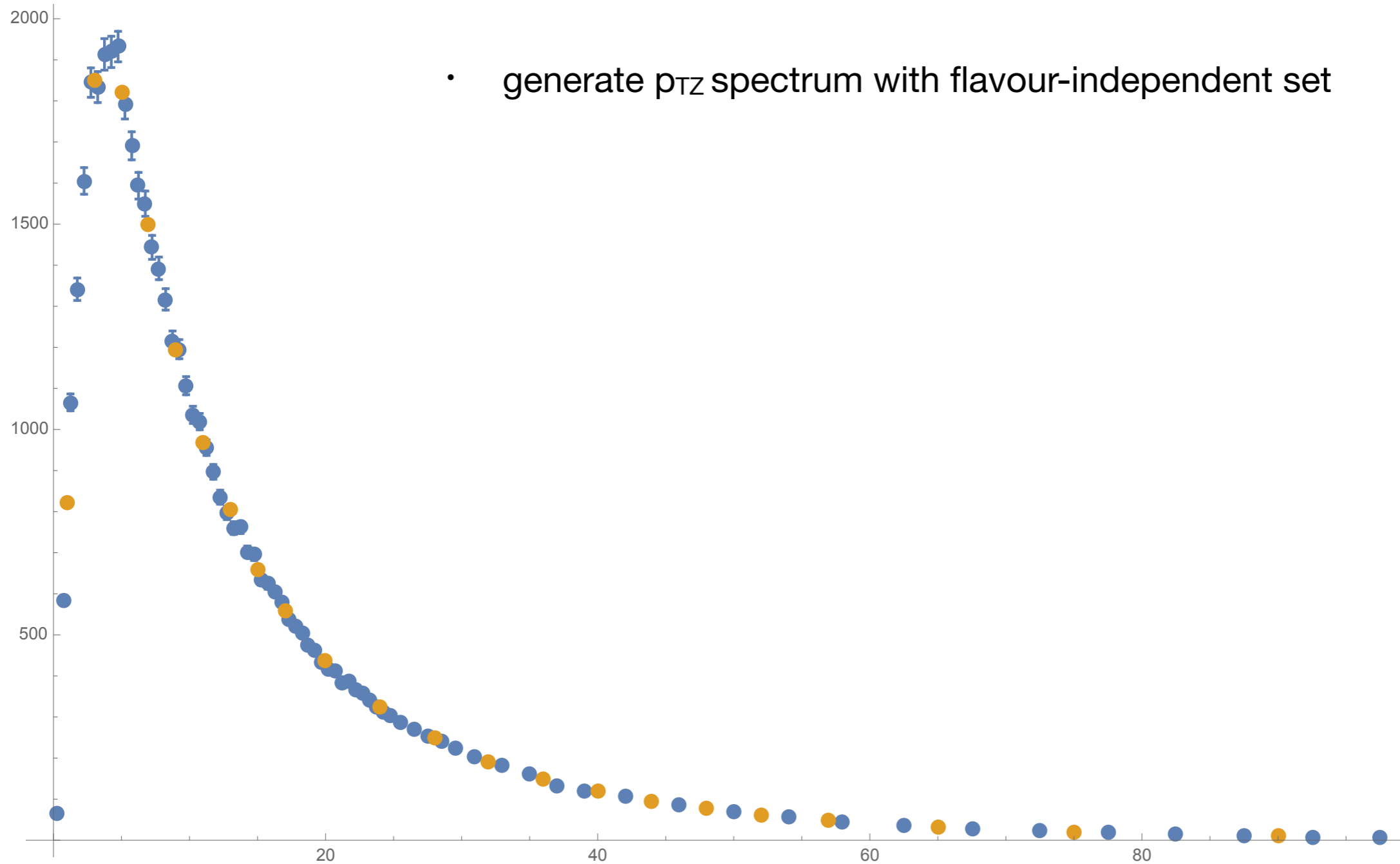
“Z-equivalent” sets



NLL+LO QCD curves obtained through a modified version of the **DYqT** code [Bozzi, Catani, deFlorian, Ferrera, Grazzini (2009,2011)]
(Tevatron 1.96 TeV & LHC 7 TeV)

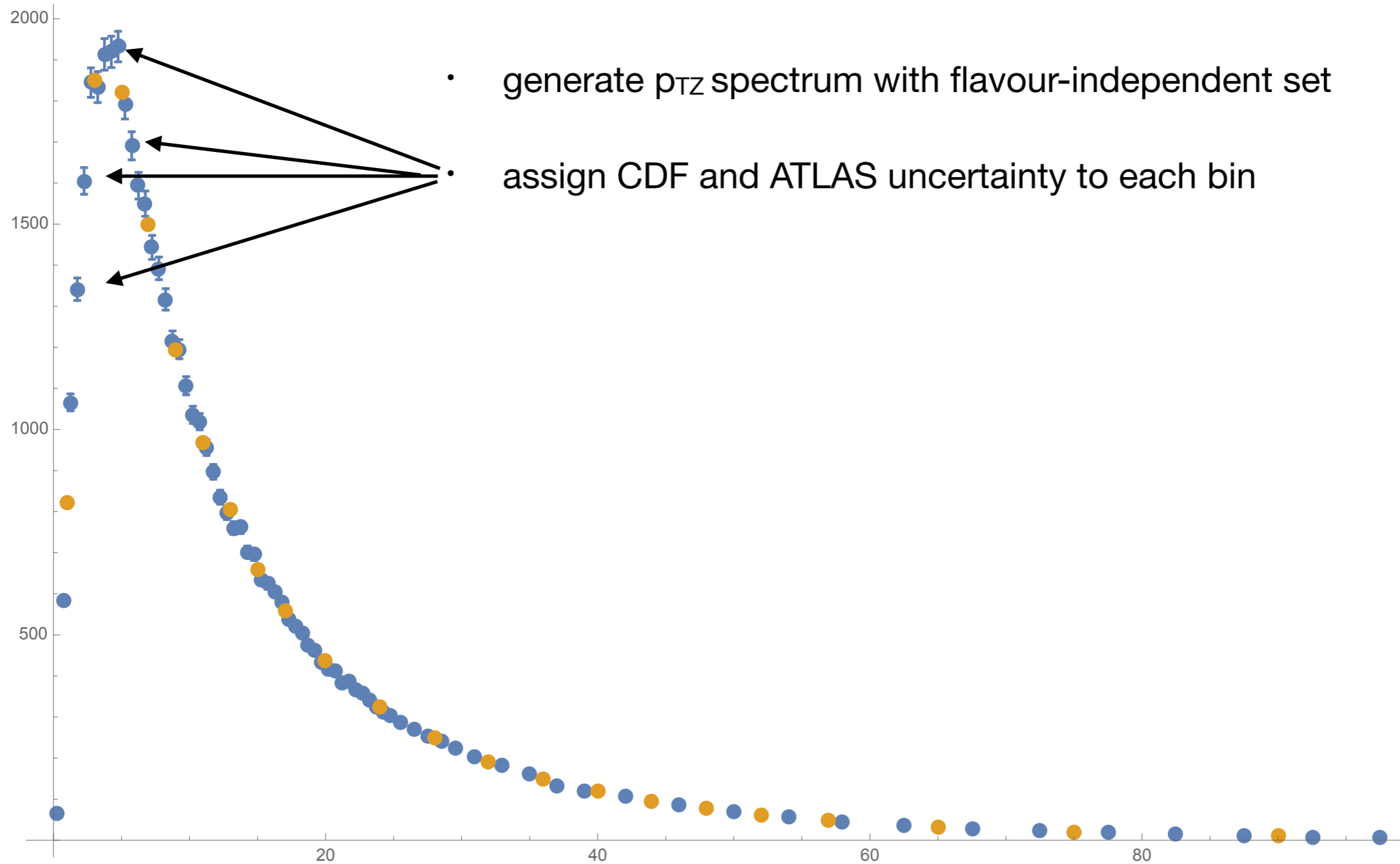
“Z-equivalent” sets

- generate p_{TZ} spectrum with flavour-independent set



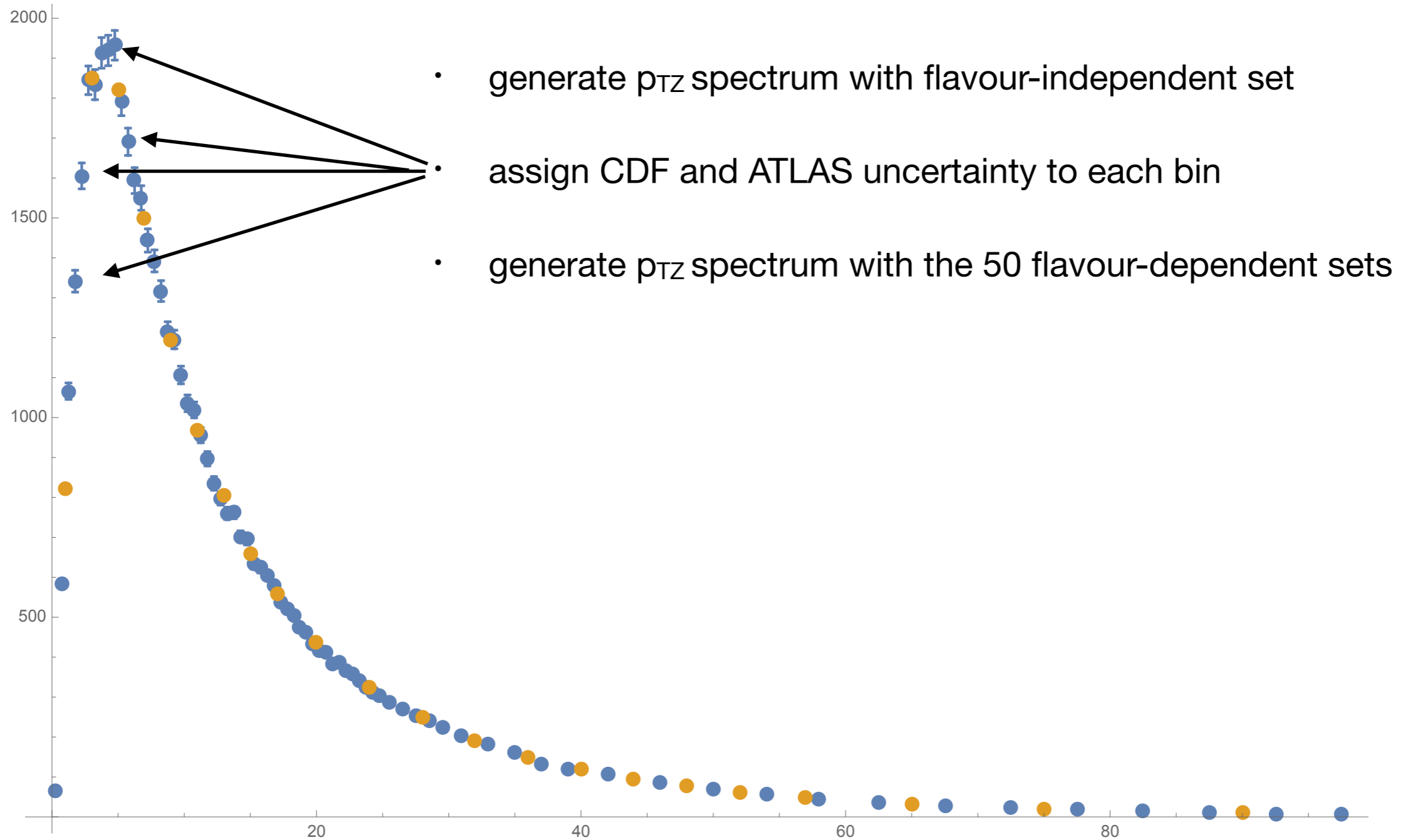
NLL+LO QCD curves obtained through a modified version of the **DYqT** code [Bozzi, Catani, deFlorian, Ferrera, Grazzini (2009,2011)]
(Tevatron 1.96 TeV & LHC 7 TeV)

“Z-equivalent” sets



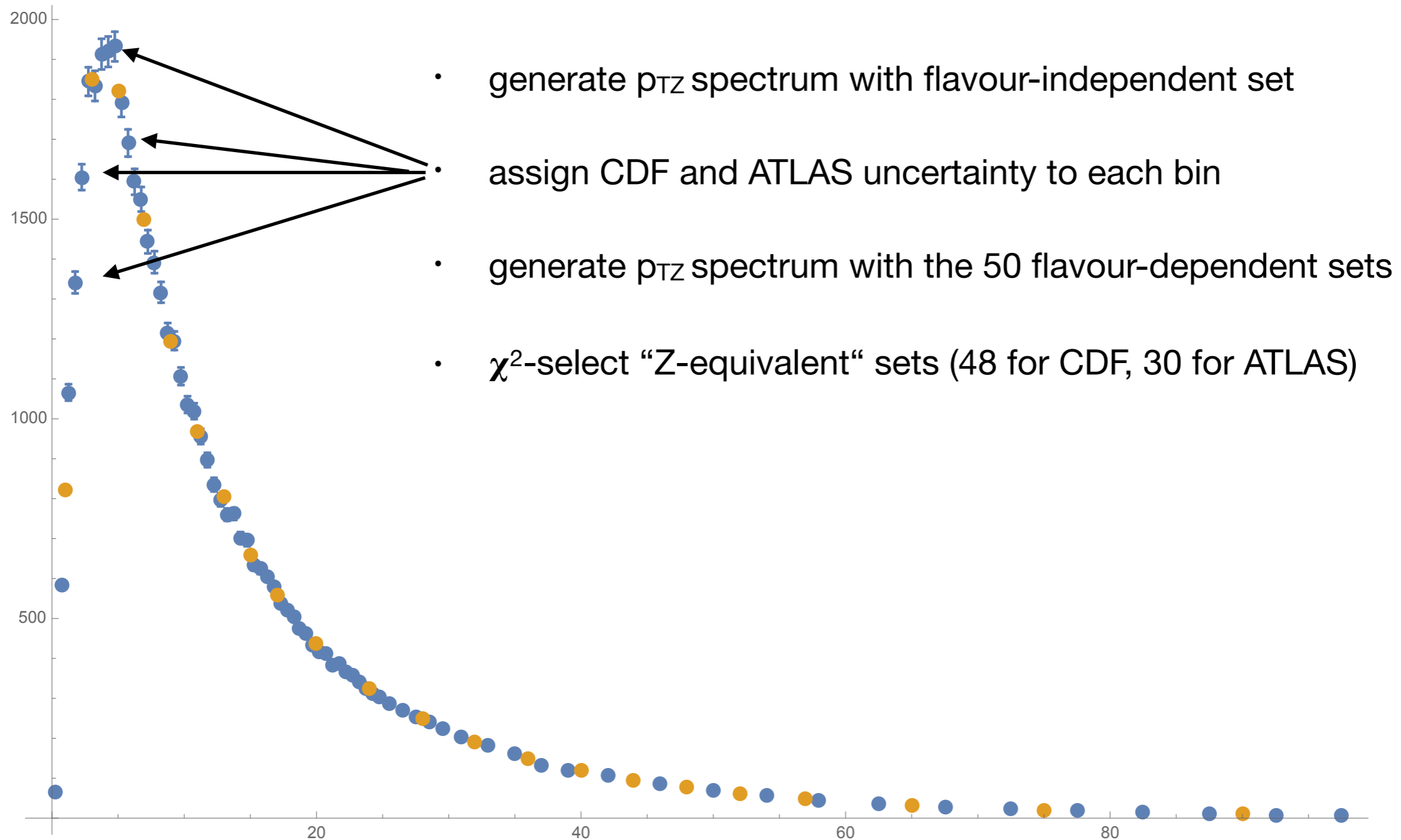
NLL+LO QCD curves obtained through a modified version of the **DYqT** code [Bozzi, Catani, deFlorian, Ferrera, Grazzini (2009,2011)]
(Tevatron 1.96 TeV & LHC 7 TeV)

“Z-equivalent” sets



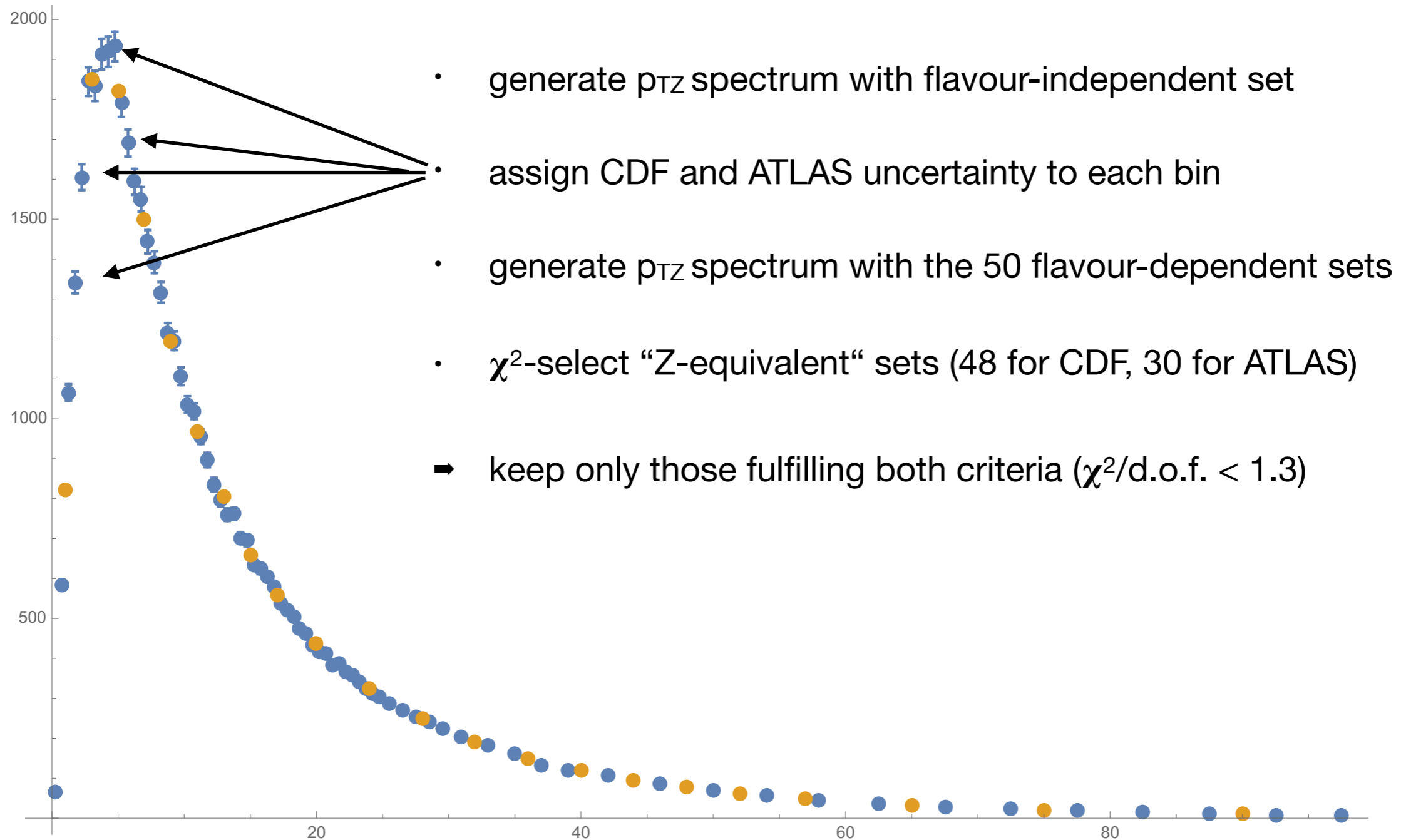
NLL+LO QCD curves obtained through a modified version of the **DYqT** code [Bozzi, Catani, deFlorian, Ferrera, Grazzini (2009,2011)]
(Tevatron 1.96 TeV & LHC 7 TeV)

“Z-equivalent” sets



NLL+LO QCD curves obtained through a modified version of the **DYqT** code [Bozzi, Catani, deFlorian, Ferrera, Grazzini (2009,2011)]
(Tevatron 1.96 TeV & LHC 7 TeV)

“Z-equivalent” sets



NLL+LO QCD curves obtained through a modified version of the **DYqT** code [Bozzi, Catani, deFlorian, Ferrera, Grazzini (2009,2011)]
(Tevatron 1.96 TeV & LHC 7 TeV)

Impact on the determination of M_w

Impact on the determination of M_W

- Take the “Z-equivalent” *flavour-dependent* parameter sets and compute *low-statistics* (135M) m_T, p_{Tl}, p_{Tn} distributions

Impact on the determination of M_W

- Take the “Z-equivalent” *flavour-dependent* parameter sets and compute *low-statistics* (135M) m_T, p_{Tl}, p_{Tn} distributions

➔ **pseudodata**

Impact on the determination of M_W

- Take the “Z-equivalent” *flavour-dependent* parameter sets and compute *low-statistics* (135M) m_T, p_{Tl}, p_{Tn} distributions

➔ **pseudodata**

- Take the *flavour-independent* parameter set and compute *high-statistics* (750M) m_T, p_{Tl}, p_{Tn} distributions for 30 different values of M_W

Impact on the determination of M_W

- Take the “Z-equivalent” *flavour-dependent* parameter sets and compute *low-statistics* (135M) m_T, p_{Tl}, p_{Tn} distributions

➔ **pseudodata**

- Take the *flavour-independent* parameter set and compute *high-statistics* (750M) m_T, p_{Tl}, p_{Tn} distributions for 30 different values of M_W

➔ **templates**

Impact on the determination of M_W

- Take the “Z-equivalent” *flavour-dependent* parameter sets and compute *low-statistics* (135M) $m_T, \rho_{Tl}, \rho_{Tn}$ distributions
 - ➔ **pseudodata**
- Take the *flavour-independent* parameter set and compute *high-statistics* (750M) $m_T, \rho_{Tl}, \rho_{Tn}$ distributions for 30 different values of M_W
 - ➔ **templates**
- **perform the template fit procedure and compute the shifts induced by flavour effects**

Impact on the determination of M_W

- Take the “Z-equivalent” *flavour-dependent* parameter sets and compute *low-statistics* (135M) $m_T, p_{T\ell}, p_{T\nu}$ distributions

➔ **pseudodata**

- Take the *flavour-independent* parameter set and compute *high-statistics* (750M) $m_T, p_{T\ell}, p_{T\nu}$ distributions for 30 different values of M_W

➔ **templates**

- **perform the template fit procedure and compute the shifts induced by flavour effects**

	ΔM_{W+}			ΔM_{W-}		
Set	m_T	$p_{T\ell}$	$p_{T\nu}$	m_T	$p_{T\ell}$	$p_{T\nu}$
1	0	-1	-2	-2	3	-3
2	0	-6	0	-2	0	-5
3	-1	9	0	-2	4	-10
4	0	0	-2	-2	-4	-10
5	0	4	1	-1	-3	-6
6	1	0	2	-1	4	-4
7	2	-1	2	-1	0	-8
8	0	2	8	1	7	8
9	0	4	-3	-1	0	7

TABLE I: ATLAS 7 TeV

	ΔM_{W+}			ΔM_{W-}		
Set	m_T	$p_{T\ell}$	$p_{T\nu}$	m_T	$p_{T\ell}$	$p_{T\nu}$
1	-1	-5	7	-1	-3	8
2	-1	-15	6	0	5	10
3	-1	1	8	-1	-7	5
4	-1	-15	6	0	-4	5
5	-1	-4	6	-1	-7	5
6	-1	-5	7	0	2	9
7	-1	-15	6	-1	-6	5
8	-1	0	8	0	3	10
9	-1	-7	7	0	4	10

TABLE II: LHCb 13 TeV

Set	u_ν	d_ν	u_s	d_s	s
1	0.34	0.26	0.46	0.59	0.32
2	0.34	0.46	0.56	0.32	0.51
3	0.55	0.34	0.33	0.55	0.30
4	0.53	0.49	0.37	0.22	0.52
5	0.42	0.38	0.29	0.57	0.27
6	0.40	0.52	0.46	0.54	0.21
7	0.22	0.21	0.40	0.46	0.49
8	0.53	0.31	0.59	0.54	0.33
9	0.46	0.46	0.58	0.40	0.28

NLL+LO QCD analysis obtained through a modified version of the **DYRes** code [Catani, deFlorian, Ferrera, Grazzini (2015)]

Statistical uncertainty: 2.5 MeV

Impact on the determination of M_W

- Take the “Z-equivalent” *flavour-dependent* parameter sets and compute *low-statistics* (135M) $m_T, p_{T\ell}, p_{T\nu}$ distributions

➔ **pseudodata**

- Take the *flavour-independent* parameter set and compute *high-statistics* (750M) $m_T, p_{T\ell}, p_{T\nu}$ distributions for 30 different values of M_W

➔ **templates**

- **perform the template fit procedure and compute the shifts induced by flavour effects**
- transverse mass: zero or few MeV shifts, generally favouring lower values for W^- (**preferred by EW fit**)

	ΔM_{W^+}			ΔM_{W^-}		
Set	m_T	$p_{T\ell}$	$p_{T\nu}$	m_T	$p_{T\ell}$	$p_{T\nu}$
1	0	-1	-2	-2	3	-3
2	0	-6	0	-2	0	-5
3	-1	9	0	-2	4	-10
4	0	0	-2	-2	-4	-10
5	0	4	1	-1	-3	-6
6	1	0	2	-1	4	-4
7	2	-1	2	-1	0	-8
8	0	2	8	1	7	8
9	0	4	-3	-1	0	7

TABLE I: ATLAS 7 TeV

	ΔM_{W^+}			ΔM_{W^-}		
Set	m_T	$p_{T\ell}$	$p_{T\nu}$	m_T	$p_{T\ell}$	$p_{T\nu}$
1	-1	-5	7	-1	-3	8
2	-1	-15	6	0	5	10
3	-1	1	8	-1	-7	5
4	-1	-15	6	0	-4	5
5	-1	-4	6	-1	-7	5
6	-1	-5	7	0	2	9
7	-1	-15	6	-1	-6	5
8	-1	0	8	0	3	10
9	-1	-7	7	0	4	10

TABLE II: LHCb 13 TeV

Set	u_ν	d_ν	u_s	d_s	s
1	0.34	0.26	0.46	0.59	0.32
2	0.34	0.46	0.56	0.32	0.51
3	0.55	0.34	0.33	0.55	0.30
4	0.53	0.49	0.37	0.22	0.52
5	0.42	0.38	0.29	0.57	0.27
6	0.40	0.52	0.46	0.54	0.21
7	0.22	0.21	0.40	0.46	0.49
8	0.53	0.31	0.59	0.54	0.33
9	0.46	0.46	0.58	0.40	0.28

NLL+LO QCD analysis obtained through a modified version of the **DYRes** code [Catani, deFlorian, Ferrera, Grazzini (2015)]

Statistical uncertainty: 2.5 MeV

Impact on the determination of M_W

- Take the “Z-equivalent” *flavour-dependent* parameter sets and compute *low-statistics* (135M) $m_T, p_{T\ell}, p_{T\nu}$ distributions

➔ **pseudodata**

- Take the *flavour-independent* parameter set and compute *high-statistics* (750M) $m_T, p_{T\ell}, p_{T\nu}$ distributions for 30 different values of M_W

➔ **templates**

- **perform the template fit procedure and compute the shifts induced by flavour effects**
- transverse mass: zero or few MeV shifts, generally favouring lower values for W^- (**preferred by EW fit**)
- lepton p_T : quite important shifts (envelope **up to 15 MeV**)

	ΔM_{W^+}			ΔM_{W^-}		
Set	m_T	$p_{T\ell}$	$p_{T\nu}$	m_T	$p_{T\ell}$	$p_{T\nu}$
1	0	-1	-2	-2	3	-3
2	0	-6	0	-2	0	-5
3	-1	9	0	-2	4	-10
4	0	0	-2	-2	-4	-10
5	0	4	1	-1	-3	-6
6	1	0	2	-1	4	-4
7	2	-1	2	-1	0	-8
8	0	2	8	1	7	8
9	0	4	-3	-1	0	7

TABLE I: ATLAS 7 TeV

	ΔM_{W^+}			ΔM_{W^-}		
Set	m_T	$p_{T\ell}$	$p_{T\nu}$	m_T	$p_{T\ell}$	$p_{T\nu}$
1	-1	-5	7	-1	-3	8
2	-1	-15	6	0	5	10
3	-1	1	8	-1	-7	5
4	-1	-15	6	0	-4	5
5	-1	-4	6	-1	-7	5
6	-1	-5	7	0	2	9
7	-1	-15	6	-1	-6	5
8	-1	0	8	0	3	10
9	-1	-7	7	0	4	10

TABLE II: LHCb 13 TeV

Set	u_ν	d_ν	u_s	d_s	s
1	0.34	0.26	0.46	0.59	0.32
2	0.34	0.46	0.56	0.32	0.51
3	0.55	0.34	0.33	0.55	0.30
4	0.53	0.49	0.37	0.22	0.52
5	0.42	0.38	0.29	0.57	0.27
6	0.40	0.52	0.46	0.54	0.21
7	0.22	0.21	0.40	0.46	0.49
8	0.53	0.31	0.59	0.54	0.33
9	0.46	0.46	0.58	0.40	0.28

NLL+LO QCD analysis obtained through a modified version of the **DYRes** code [Catani, deFlorian, Ferrera, Grazzini (2015)]

Statistical uncertainty: 2.5 MeV

Impact on the determination of M_W

- Take the “Z-equivalent” *flavour-dependent* parameter sets and compute *low-statistics* (135M) $m_T, p_{T\ell}, p_{T\nu}$ distributions

➔ **pseudodata**

- Take the *flavour-independent* parameter set and compute *high-statistics* (750M) $m_T, p_{T\ell}, p_{T\nu}$ distributions for 30 different values of M_W

➔ **templates**

- **perform the template fit procedure and compute the shifts induced by flavour effects**
- transverse mass: zero or few MeV shifts, generally favouring lower values for W^- (**preferred by EW fit**)
- lepton pt: quite important shifts (envelope **up to 15 MeV**)
- neutrino pt: same order of magnitude (or bigger) as lepton pt

	ΔM_{W^+}			ΔM_{W^-}		
Set	m_T	$p_{T\ell}$	$p_{T\nu}$	m_T	$p_{T\ell}$	$p_{T\nu}$
1	0	-1	-2	-2	3	-3
2	0	-6	0	-2	0	-5
3	-1	9	0	-2	4	-10
4	0	0	-2	-2	-4	-10
5	0	4	1	-1	-3	-6
6	1	0	2	-1	4	-4
7	2	-1	2	-1	0	-8
8	0	2	8	1	7	8
9	0	4	-3	-1	0	7

TABLE I: ATLAS 7 TeV

	ΔM_{W^+}			ΔM_{W^-}		
Set	m_T	$p_{T\ell}$	$p_{T\nu}$	m_T	$p_{T\ell}$	$p_{T\nu}$
1	-1	-5	7	-1	-3	8
2	-1	-15	6	0	5	10
3	-1	1	8	-1	-7	5
4	-1	-15	6	0	-4	5
5	-1	-4	6	-1	-7	5
6	-1	-5	7	0	2	9
7	-1	-15	6	-1	-6	5
8	-1	0	8	0	3	10
9	-1	-7	7	0	4	10

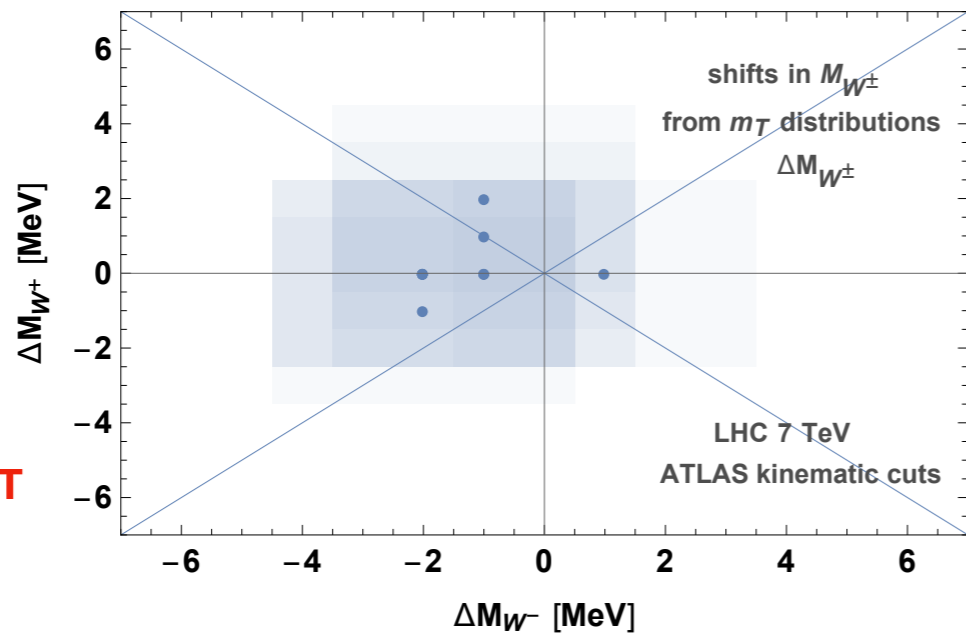
TABLE II: LHCb 13 TeV

Set	u_ν	d_ν	u_s	d_s	s
1	0.34	0.26	0.46	0.59	0.32
2	0.34	0.46	0.56	0.32	0.51
3	0.55	0.34	0.33	0.55	0.30
4	0.53	0.49	0.37	0.22	0.52
5	0.42	0.38	0.29	0.57	0.27
6	0.40	0.52	0.46	0.54	0.21
7	0.22	0.21	0.40	0.46	0.49
8	0.53	0.31	0.59	0.54	0.33
9	0.46	0.46	0.58	0.40	0.28

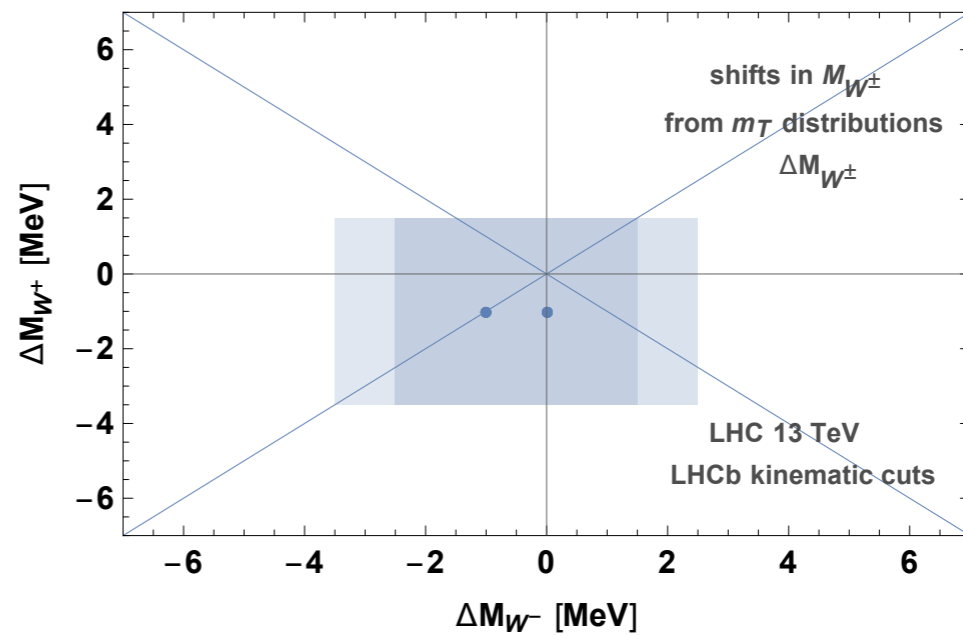
NLL+LO QCD analysis obtained through a modified version of the **DYRes** code [Catani, deFlorian, Ferrera, Grazzini (2015)]

Statistical uncertainty: 2.5 MeV

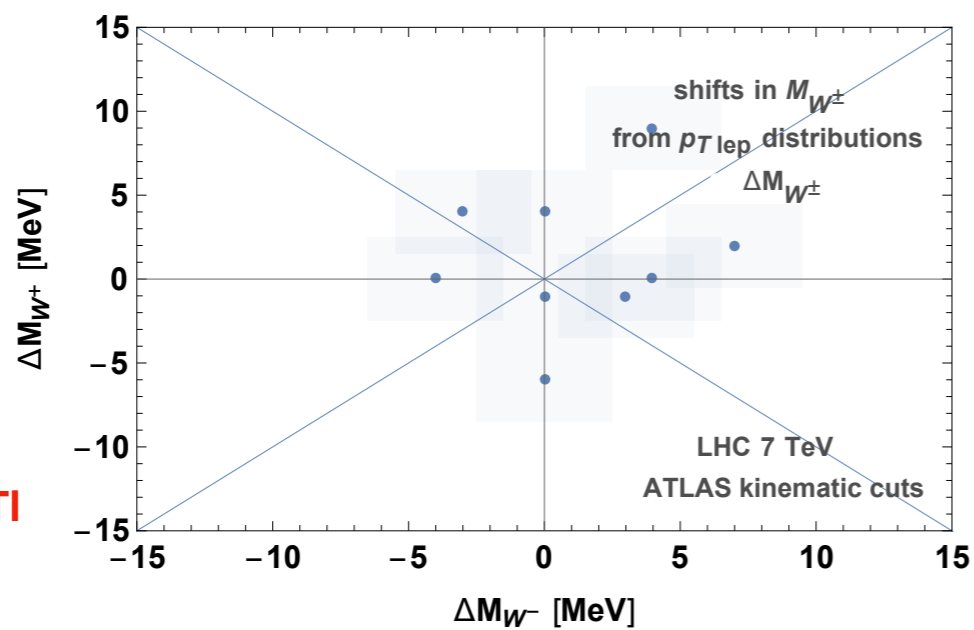
ATLAS m_T



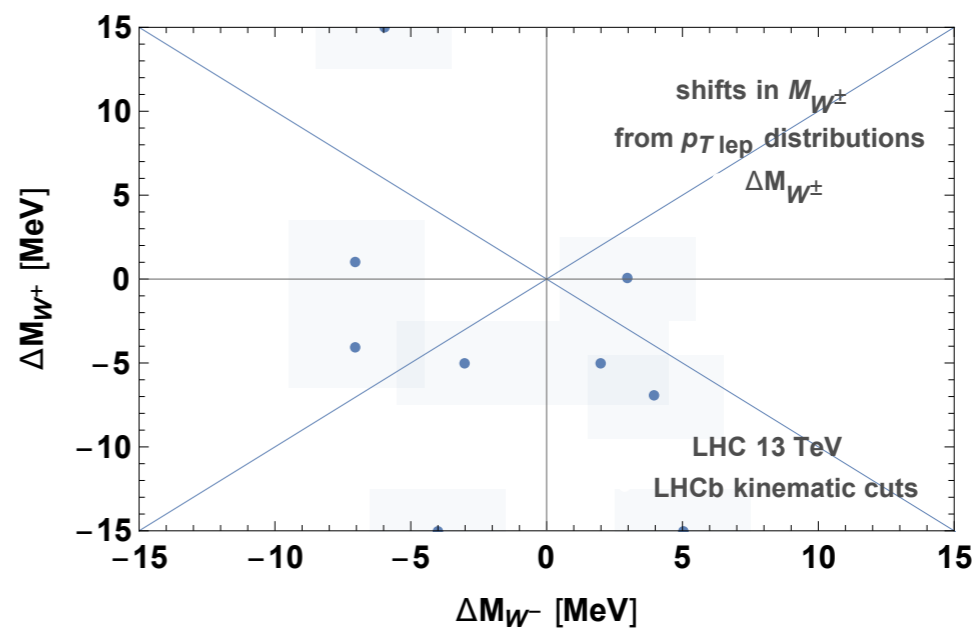
LHCb m_T



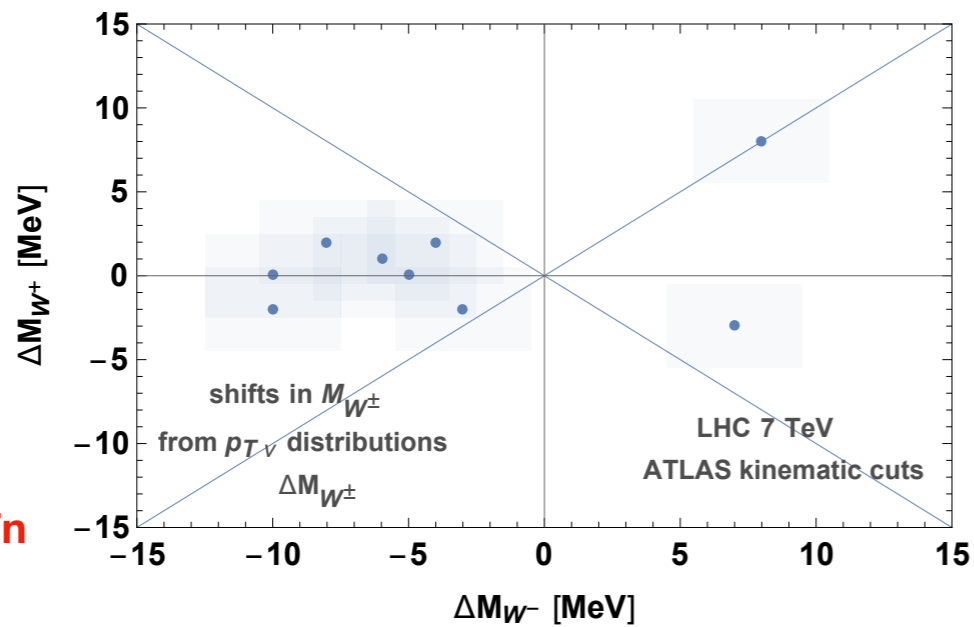
ATLAS p_{Tl}



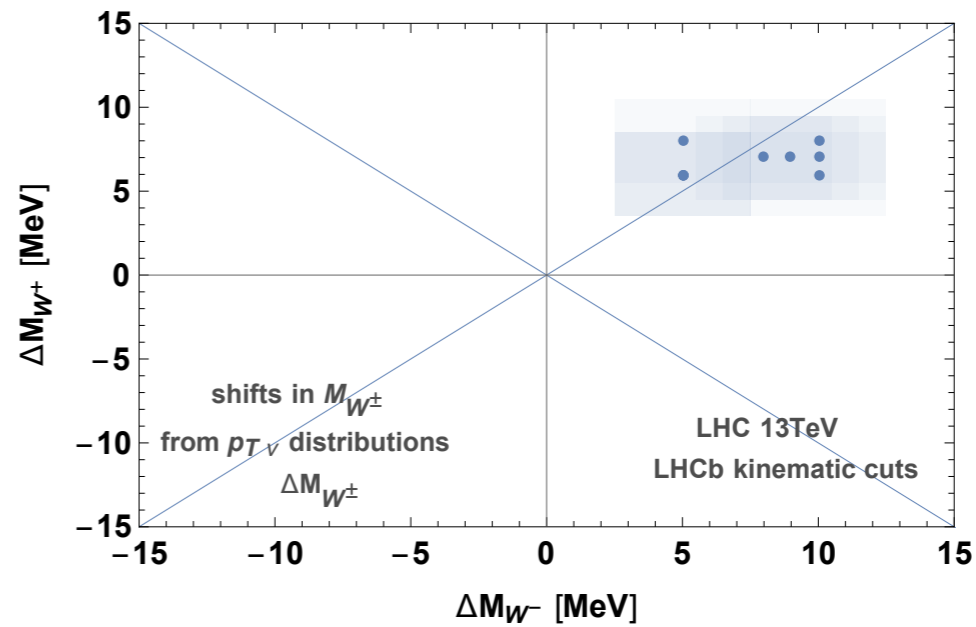
LHCb p_{Tl}



ATLAS p_{Tv}



LHCb p_{Tv}



Backup slides

Uncertainties on M_W due to p_{TW}

CDF

m_T fit uncertainties				p_T^ℓ fit uncertainties			
Source	$W \rightarrow \mu\nu$	$W \rightarrow e\nu$	Common	Source	$W \rightarrow \mu\nu$	$W \rightarrow e\nu$	Common
Lepton energy scale	7	10	5	Lepton energy scale	7	10	5
Lepton energy resolution	1	4	0	Lepton energy resolution	1	4	0
Lepton efficiency	0	0	0	Lepton efficiency	1	2	0
Lepton tower removal	2	3	2	Lepton tower removal	0	0	0
Recoil scale	5	5	5	Recoil scale	6	6	6
Recoil resolution	7	7	7	Recoil resolution	5	5	5
Backgrounds	3	4	0	Backgrounds	5	3	0
PDFs	10	10	10	PDFs	9	9	9
W boson p_T	3	3	3	W boson p_T	9	9	9
Photon radiation	4	4	4	Photon radiation	4	4	4
Statistical	16	19	0	Statistical	18	21	0
Total	23	26	15	Total	25	28	16

D0

Source	Section	m_T	p_T^ℓ	E_T
Experimental				
Electron Energy Scale	VII C4	16	17	16
Electron Energy Resolution	VII C5	2	2	3
Electron Shower Model	V C	4	6	7
Electron Energy Loss	VD	4	4	4
Recoil Model	VII D3	5	6	14
Electron Efficiencies	VII B10	1	3	5
Backgrounds	VIII	2	2	2
$\Sigma(\text{Experimental})$		18	20	24
W Production and Decay Model				
PDF	VIC	11	11	14
QED	VIB	7	7	9
Boson p_T	VIA	2	5	2
$\Sigma(\text{Model})$		13	14	17
Systematic Uncertainty (Experimental and Model)		22	24	29
W Boson Statistics	IX	13	14	15
Total Uncertainty		26	28	33

ATLAS

W -boson charge	W^+		W^-		Combined	
Kinematic distribution	p_T^ℓ	m_T	p_T^ℓ	m_T	p_T^ℓ	m_T
δm_W [MeV]						
Fixed-order PDF uncertainty	13.1	14.9	12.0	14.2	8.0	8.7
AZ tune	3.0	3.4	3.0	3.4	3.0	3.4
Charm-quark mass	1.2	1.5	1.2	1.5	1.2	1.5
Parton shower μ_F with heavy-flavour decorrelation	5.0	6.9	5.0	6.9	5.0	6.9
Parton shower PDF uncertainty	3.6	4.0	2.6	2.4	1.0	1.6
Angular coefficients	5.8	5.3	5.8	5.3	5.8	5.3
Total	15.9	18.1	14.8	17.2	11.6	12.9

Uncertainties on M_W due to p_{TW}

CDF

m_T fit uncertainties				p_T^ℓ fit uncertainties			
Source	$W \rightarrow \mu\nu$	$W \rightarrow e\nu$	Common	Source	$W \rightarrow \mu\nu$	$W \rightarrow e\nu$	Common
Lepton energy scale	7	10	5	Lepton energy scale	7	10	5
Lepton energy resolution	1	4	0	Lepton energy resolution	1	4	0
Lepton efficiency	0	0	0	Lepton efficiency	1	2	0
Lepton tower removal	2	3	2	Lepton tower removal	0	0	0
Recoil scale	5	5	5	Recoil scale	6	6	6
Recoil resolution	7	7	7	Recoil resolution	5	5	5
Backgrounds	3	4	0	Backgrounds	5	3	0
PDFs	10	10	10	PDFs	9	9	9
W boson p_T	3	3	3	W boson p_T	9	9	9
Photon radiation	4	4	4	Photon radiation	4	4	4
Statistical	16	19	0	Statistical	18	21	0
Total	23	26	15	Total	25	28	16

D0

Source	Section	m_T	p_T^ℓ	E_T
Experimental				
Electron Energy Scale	VII C4	16	17	16
Electron Energy Resolution	VII C5	2	2	3
Electron Shower Model	V C	4	6	7
Electron Energy Loss	VD	4	4	4
Recoil Model	VII D3	5	6	14
Electron Efficiencies	VII B10	1	3	5
Backgrounds	VIII	2	2	2
$\Sigma(\text{Experimental})$		18	20	24
W Production and Decay Model				
PDF	VIC	11	11	14
QED	VIB	7	7	9
Boson p_T	VIA	2	5	2
$\Sigma(\text{Model})$		13	14	17
Systematic Uncertainty (Experimental and Model)		22	24	29
W Boson Statistics	IX	13	14	15
Total Uncertainty		26	28	33

ATLAS

W -boson charge Kinematic distribution	W^+		W^-		Combined	
	p_T^ℓ	m_T	p_T^ℓ	m_T	p_T^ℓ	m_T
δm_W [MeV]						
Fixed-order PDF uncertainty	13.1	14.9	12.0	14.2	8.0	8.7
AZ tune	3.0	3.4	3.0	3.4	3.0	3.4
Charm-quark mass	1.2	1.5	1.2	1.5	1.2	1.5
Parton shower μ_F with heavy-flavour decorrelation	5.0	6.9	5.0	6.9	5.0	6.9
Parton shower PDF uncertainty	3.6	4.0	2.6	2.4	1.0	1.6
Angular coefficients	5.8	5.3	5.8	5.3	5.8	5.3
Total	15.9	18.1	14.8	17.2	11.6	12.9

Uncertainties on M_W due to p_{TW}

CDF

m_T fit uncertainties				p_T^ℓ fit uncertainties			
Source	$W \rightarrow \mu\nu$	$W \rightarrow e\nu$	Common	Source	$W \rightarrow \mu\nu$	$W \rightarrow e\nu$	Common
Lepton energy scale	7	10	5	Lepton energy scale	7	10	5
Lepton energy resolution	1	4	0	Lepton energy resolution	1	4	0
Lepton efficiency	0	0	0	Lepton efficiency	1	2	0
Lepton tower removal	2	3	2	Lepton tower removal	0	0	0
Recoil scale	5	5	5	Recoil scale	6	6	6
Recoil resolution	7	7	7	Recoil resolution	5	5	5
Backgrounds	3	4	0	Backgrounds	5	3	0
PDFs	10	10	10	PDFs	9	9	9
W boson p_T	3	3	3	W boson p_T	9	9	9
Photon radiation	4	4	4	Photon radiation	4	4	4
Statistical	16	19	0	Statistical	18	21	0
Total	23	26	15	Total	25	28	16

D0

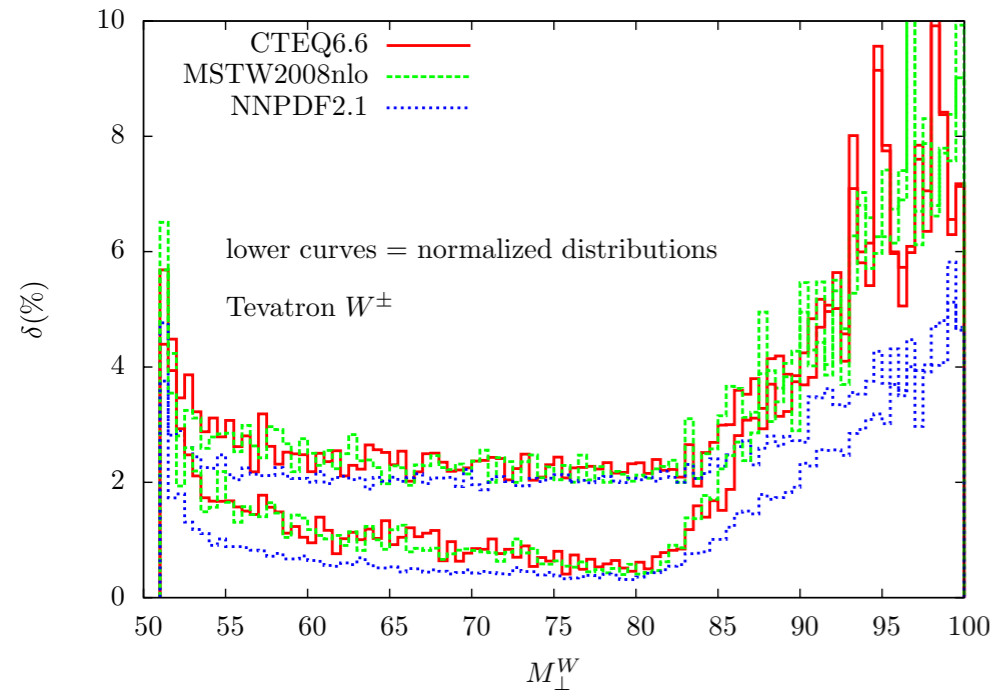
Source	Section	m_T	p_T^ℓ	E_T
Experimental				
Electron Energy Scale	VII C4	16	17	16
Electron Energy Resolution	VII C5	2	2	3
Electron Shower Model	V C	4	6	7
Electron Energy Loss	VD	4	4	4
Recoil Model	VII D3	5	6	14
Electron Efficiencies	VII B10	1	3	5
Backgrounds	VIII	2	2	2
$\Sigma(\text{Experimental})$		18	20	24
W Production and Decay Model				
PDF	VIC	11	11	14
QED	VIB	7	7	9
Boson p_T	VIA	2	5	2
$\Sigma(\text{Model})$		13	14	17
Systematic Uncertainty (Experimental and Model)		22	24	29
W Boson Statistics	IX	13	14	15
Total Uncertainty		26	28	33

ATLAS

W -boson charge Kinematic distribution	W^+		W^-		Combined	
	p_T^ℓ	m_T	p_T^ℓ	m_T	p_T^ℓ	m_T
δm_W [MeV]						
Fixed-order PDF uncertainty	13.1	14.9	12.0	14.2	8.0	8.7
AZ tune	3.0	3.4	3.0	3.4	3.0	3.4
Charm-quark mass	1.2	1.5	1.2	1.5	1.2	1.5
Parton shower μ_F with heavy-flavour decorrelation	5.0	6.9	5.0	6.9	5.0	6.9
Parton shower PDF uncertainty	3.6	4.0	2.6	2.4	1.0	1.6
Angular coefficients	5.8	5.3	5.8	5.3	5.8	5.3
Total	15.9	18.1	14.8	17.2	11.6	12.9

PDF effect on transverse mass

Bozzi, Rojo, Vicini PRD 83, 113008 (2011)

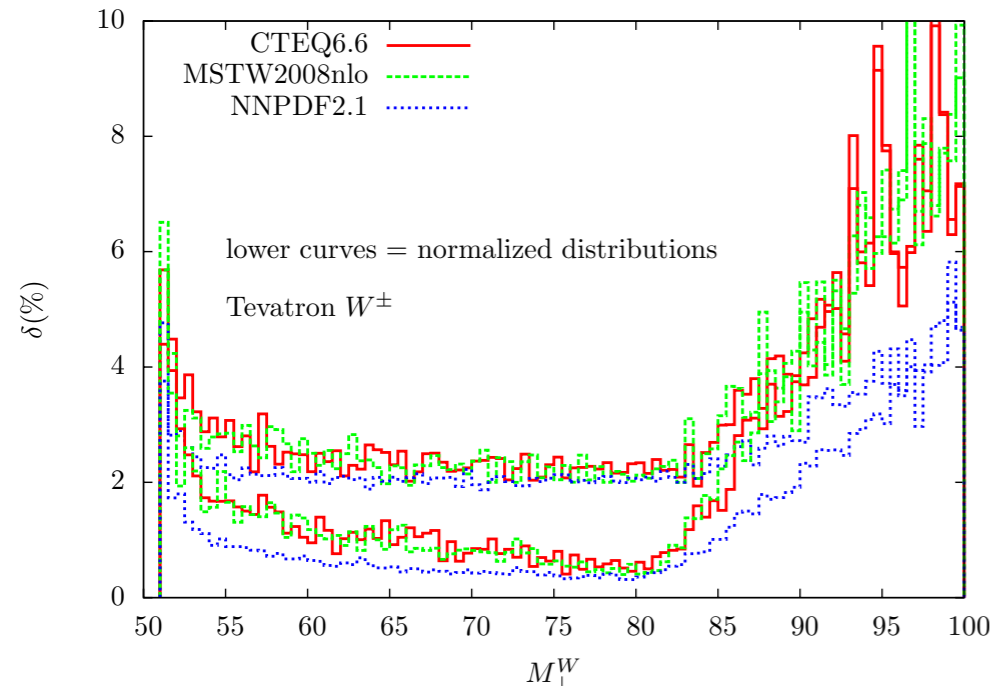


- **Normalised** distributions: reduced sensitivity to PDFs
- Ratio of (non-)normalised distributions w.r.t. to central PDF set
- Distributions obtained with **DYNNLO**

*in first approximation the PDF effects **factorise** w.r.t. all other theoretical and experimental factors*

PDF effect on transverse mass

Bozzi, Rojo, Vicini PRD 83, 113008 (2011)



- **Normalised** distributions: reduced sensitivity to PDFs
- Ratio of (non-)normalised distributions w.r.t. to central PDF set
- Distributions obtained with **DYNNLO**

*in first approximation the PDF effects **factorise** w.r.t. all other theoretical and experimental factors*

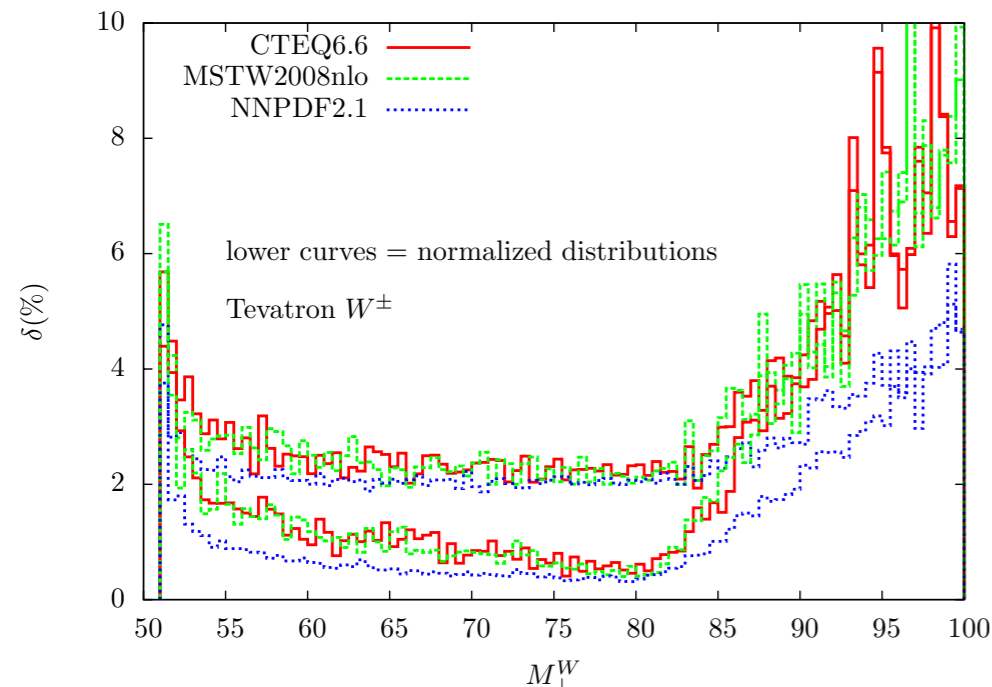
	CTEQ6.6		MSTW2008		NNPDF2.1		δ_{pdf}^{tot}
	$m_W \pm \delta_{pdf}$	$\langle \chi^2 \rangle$	$m_W \pm \delta_{pdf}$	$\langle \chi^2 \rangle$	$m_W \pm \delta_{pdf}$	$\langle \chi^2 \rangle$	
Tevatron, W^\pm	80.398 ± 0.004	1.42	80.398 ± 0.003	1.42	80.398 ± 0.003	1.30	4
LHC 7 TeV W^+	80.398 ± 0.004	1.22	80.404 ± 0.005	1.55	80.402 ± 0.003	1.35	8
LHC 7 TeV W^-	80.398 ± 0.004	1.22	80.400 ± 0.004	1.19	80.402 ± 0.004	1.78	6
LHC 14 TeV W^+	80.398 ± 0.003	1.34	80.402 ± 0.004	1.48	80.400 ± 0.003	1.41	6
LHC 14 TeV W^-	80.398 ± 0.004	1.44	80.404 ± 0.006	1.38	80.402 ± 0.004	1.57	8

PDF effect on transverse mass

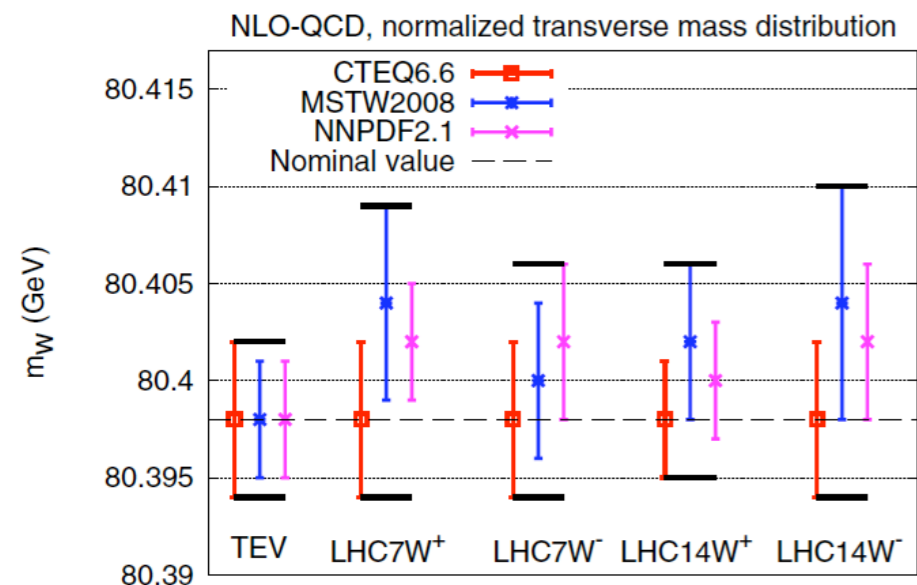
Bozzi, Rojo, Vicini PRD 83, 113008 (2011)

- **Normalised** distributions: reduced sensitivity to PDFs
- Ratio of (non-)normalised distributions w.r.t. to central PDF set
- Distributions obtained with **DYNNLO**

*in first approximation the PDF effects **factorise** w.r.t. all other theoretical and experimental factors*

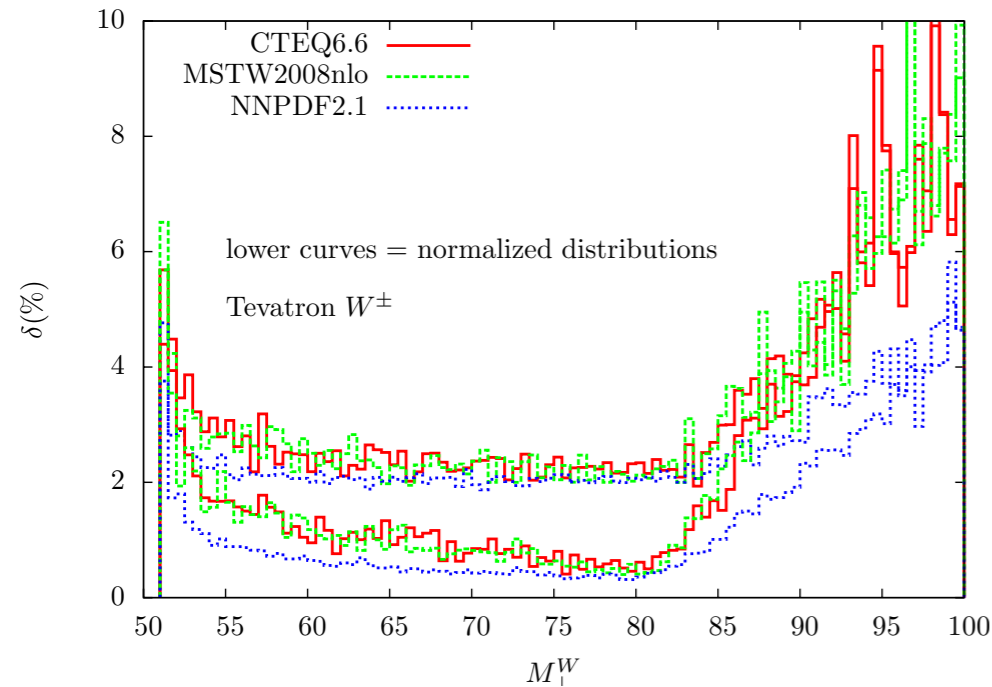


	CTEQ6.6		MSTW2008		NNPDF2.1		δ_{pdf}^{tot}
	$m_W \pm \delta_{pdf}$	$\langle \chi^2 \rangle$	$m_W \pm \delta_{pdf}$	$\langle \chi^2 \rangle$	$m_W \pm \delta_{pdf}$	$\langle \chi^2 \rangle$	
Tevatron, W^\pm	80.398 ± 0.004	1.42	80.398 ± 0.003	1.42	80.398 ± 0.003	1.30	4
LHC 7 TeV W^+	80.398 ± 0.004	1.22	80.404 ± 0.005	1.55	80.402 ± 0.003	1.35	8
LHC 7 TeV W^-	80.398 ± 0.004	1.22	80.400 ± 0.004	1.19	80.402 ± 0.004	1.78	6
LHC 14 TeV W^+	80.398 ± 0.003	1.34	80.402 ± 0.004	1.48	80.400 ± 0.003	1.41	6
LHC 14 TeV W^-	80.398 ± 0.004	1.44	80.404 ± 0.006	1.38	80.402 ± 0.004	1.57	8



PDF effect on transverse mass

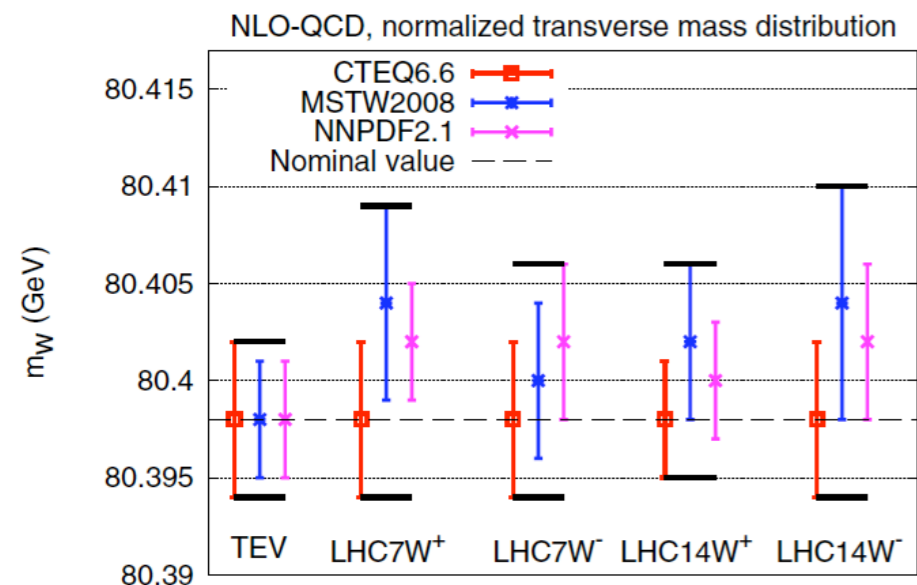
Bozzi, Rojo, Vicini PRD 83, 113008 (2011)



- **Normalised** distributions: reduced sensitivity to PDFs
- Ratio of (non-)normalised distributions w.r.t. to central PDF set
- Distributions obtained with **DYNNLO**

*in first approximation the PDF effects **factorise** w.r.t. all other theoretical and experimental factors*

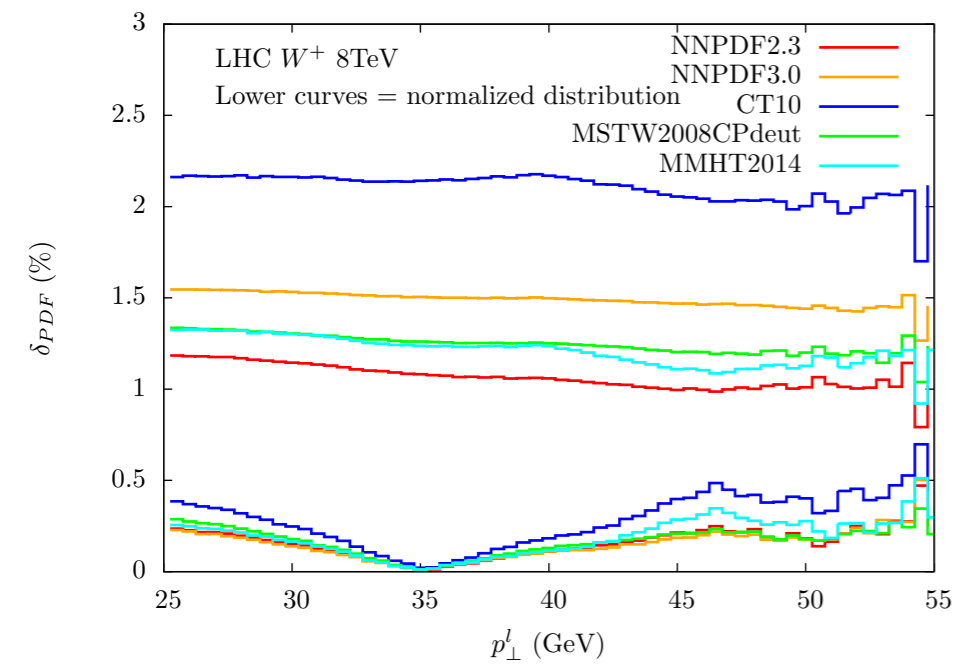
	CTEQ6.6		MSTW2008		NNPDF2.1		δ_{pdf}^{tot}
	$m_W \pm \delta_{pdf}$	$\langle \chi^2 \rangle$	$m_W \pm \delta_{pdf}$	$\langle \chi^2 \rangle$	$m_W \pm \delta_{pdf}$	$\langle \chi^2 \rangle$	
Tevatron, W^\pm	80.398 ± 0.004	1.42	80.398 ± 0.003	1.42	80.398 ± 0.003	1.30	4
LHC 7 TeV W^+	80.398 ± 0.004	1.22	80.404 ± 0.005	1.55	80.402 ± 0.003	1.35	8
LHC 7 TeV W^-	80.398 ± 0.004	1.22	80.400 ± 0.004	1.19	80.402 ± 0.004	1.78	6
LHC 14 TeV W^+	80.398 ± 0.003	1.34	80.402 ± 0.004	1.48	80.400 ± 0.003	1.41	6
LHC 14 TeV W^-	80.398 ± 0.004	1.44	80.404 ± 0.006	1.38	80.402 ± 0.004	1.57	8



- Accuracy of templates essential: highly demanding computing task!
- For transverse mass distribution, a **fixed-order NLO-QCD analysis is sufficient** to assess this PDF uncertainty
- PDF error is moderate at the Tevatron but also at the LHC

PDF effect on lepton p_T

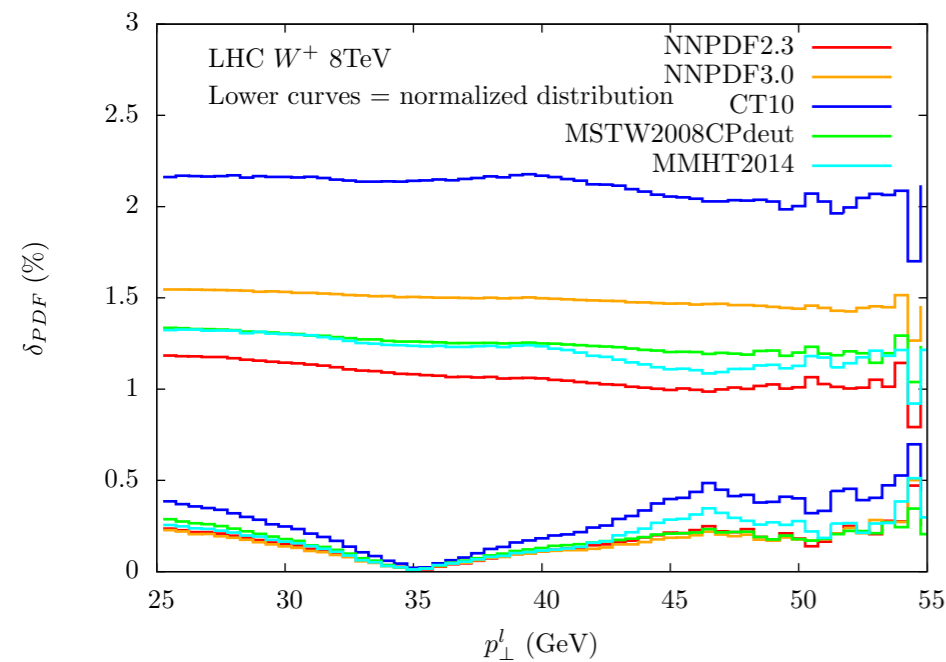
Bozzi, Citelli, Vicini PRD 91, 113005 (2015)



- **Conservative** estimate of the PDF uncertainty: **CC-DY channel alone**
- Distributions obtained with **POWHEG+PYTHIA 6.4**
- PDF uncertainty over relevant p_T range almost flat: O(2%)
- Uncertainty of normalised distributions: below the O(0.5%) level (but still sufficient to yield large M_W shifts)

PDF effect on lepton p_T

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

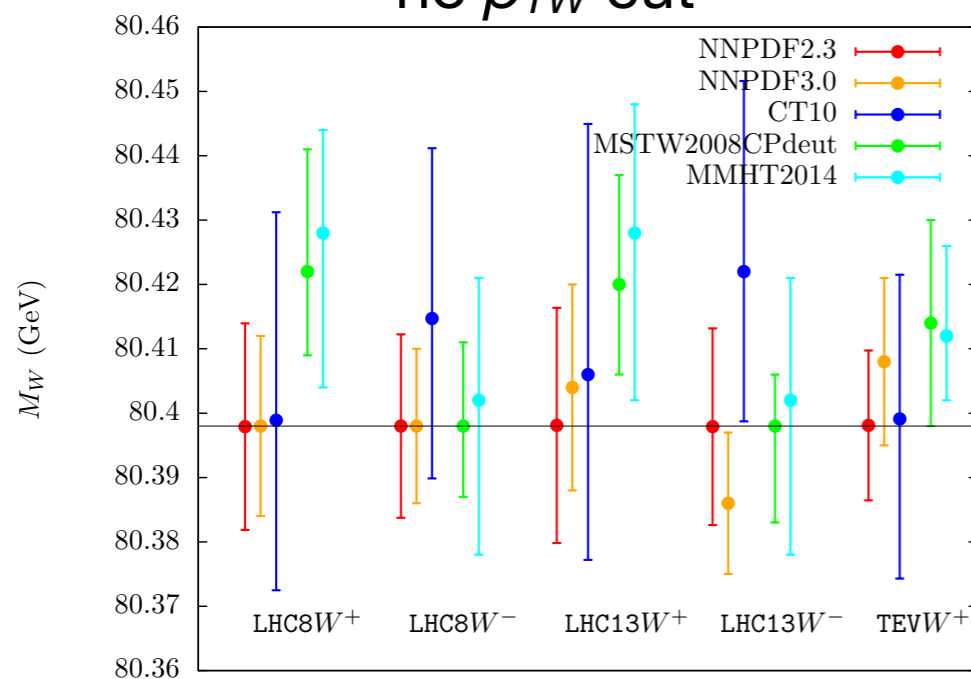


- **Conservative** estimate of the PDF uncertainty: **CC-DY channel alone**
- Distributions obtained with **POWHEG+PYTHIA 6.4**
- PDF uncertainty over relevant p_T range almost flat: $O(2\%)$
- Uncertainty of normalised distributions: below the $O(0.5\%)$ level (but still sufficient to yield large M_W shifts)

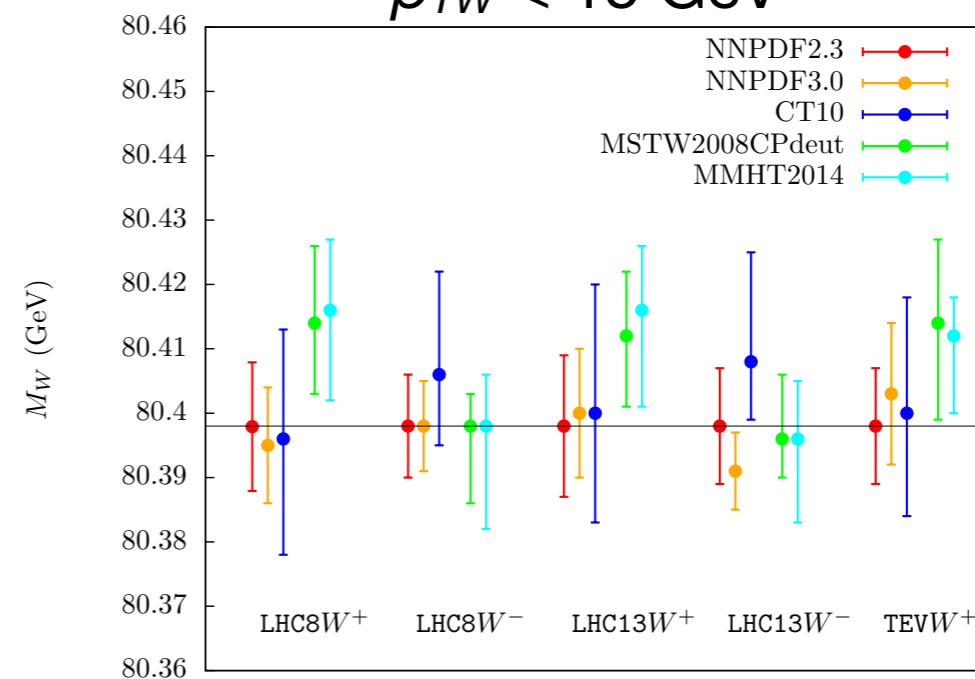
	no p_{\perp}^W cut		$p_{\perp}^W < 15$ GeV	
	δ_{PDF} (MeV)	Δ_{sets} (MeV)	δ_{PDF} (MeV)	Δ_{sets} (MeV)
Tevatron 1.96 TeV	27	16	21	15
LHC 8 TeV W^+	33	26	24	18
W^-	29	16	18	8
LHC 13 TeV W^+	34	22	20	14
W^-	34	24	18	12

- Individual PDF sets provide non-pessimistic estimates: $\Delta M_W \sim O(10$ MeV)
- Global envelope still shows large discrepancies of the central values
- p_{TW} cut is relevant

no p_{TW} cut



$p_{TW} < 15$ GeV



Acceptance cuts: interesting insights

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

normalized distributions			
cut on p_{\perp}^W	cut on $ \eta_l $	CT10	NNPDF3.0
inclusive	$ \eta_l < 2.5$	$80.400 + 0.032 - 0.027$	80.398 ± 0.014
$p_{\perp}^W < 20$ GeV	$ \eta_l < 2.5$	$80.396 + 0.027 - 0.020$	80.394 ± 0.012
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 10$ GeV	$ \eta_l < 2.5$	$80.392 + 0.015 - 0.012$	80.394 ± 0.007
$p_{\perp}^W < 15$ GeV	$ \eta_l < 1.0$	$80.400 + 0.032 - 0.021$	80.406 ± 0.017
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 15$ GeV	$ \eta_l < 4.9$	$80.400 + 0.009 - 0.004$	80.401 ± 0.003
$p_{\perp}^W < 15$ GeV	$1.0 < \eta_l < 2.5$	$80.392 + 0.025 - 0.018$	80.388 ± 0.012

Acceptance cuts: interesting insights

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

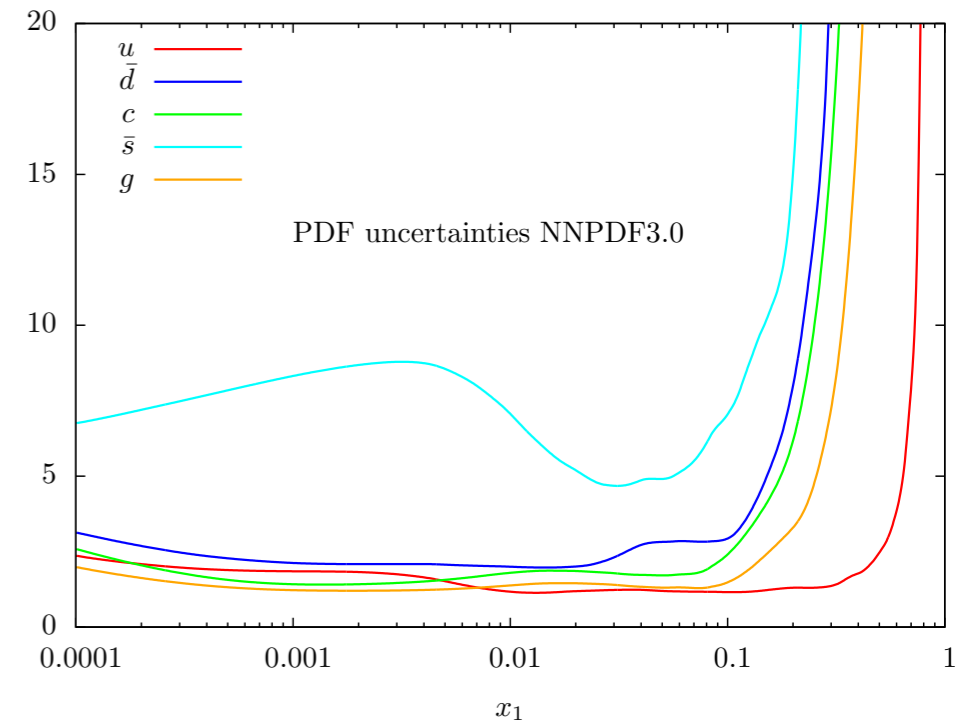
normalized distributions			
cut on p_{\perp}^W	cut on $ \eta_l $	CT10	NNPDF3.0
inclusive	$ \eta_l < 2.5$	$80.400 + 0.032 - 0.027$	80.398 ± 0.014
$p_{\perp}^W < 20$ GeV	$ \eta_l < 2.5$	$80.396 + 0.027 - 0.020$	80.394 ± 0.012
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 10$ GeV	$ \eta_l < 2.5$	$80.392 + 0.015 - 0.012$	80.394 ± 0.007
$p_{\perp}^W < 15$ GeV	$ \eta_l < 1.0$	$80.400 + 0.032 - 0.021$	80.406 ± 0.017
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 15$ GeV	$ \eta_l < 4.9$	$80.400 + 0.009 - 0.004$	80.401 ± 0.003
$p_{\perp}^W < 15$ GeV	$1.0 < \eta_l < 2.5$	$80.392 + 0.025 - 0.018$	80.388 ± 0.012

strong p_{TW} cut reduces M_W uncertainty

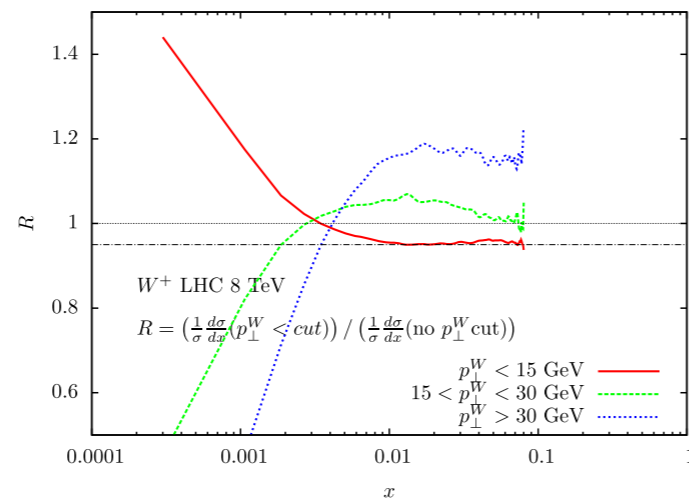
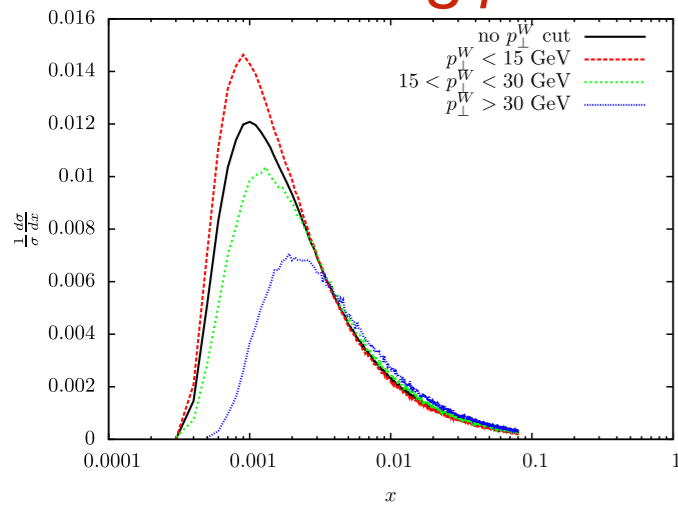
Acceptance cuts: interesting insights

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

normalized distributions			
cut on p_{\perp}^W	cut on $ \eta_l $	CT10	NNPDF3.0
inclusive	$ \eta_l < 2.5$	$80.400 + 0.032 - 0.027$	80.398 ± 0.014
$p_{\perp}^W < 20$ GeV	$ \eta_l < 2.5$	$80.396 + 0.027 - 0.020$	80.394 ± 0.012
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 10$ GeV	$ \eta_l < 2.5$	$80.392 + 0.015 - 0.012$	80.394 ± 0.007
$p_{\perp}^W < 15$ GeV	$ \eta_l < 1.0$	$80.400 + 0.032 - 0.021$	80.406 ± 0.017
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 15$ GeV	$ \eta_l < 4.9$	$80.400 + 0.009 - 0.004$	80.401 ± 0.003
$p_{\perp}^W < 15$ GeV	$1.0 < \eta_l < 2.5$	$80.392 + 0.025 - 0.018$	80.388 ± 0.012



strong p_{TW} cut reduces M_W uncertainty

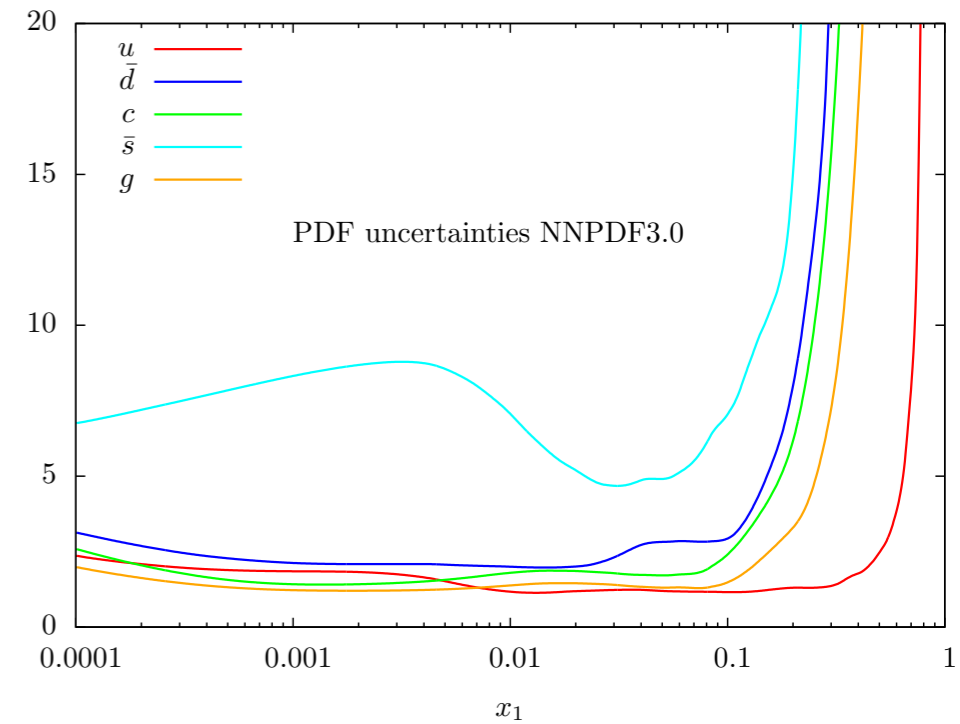


suppression of the large-x region

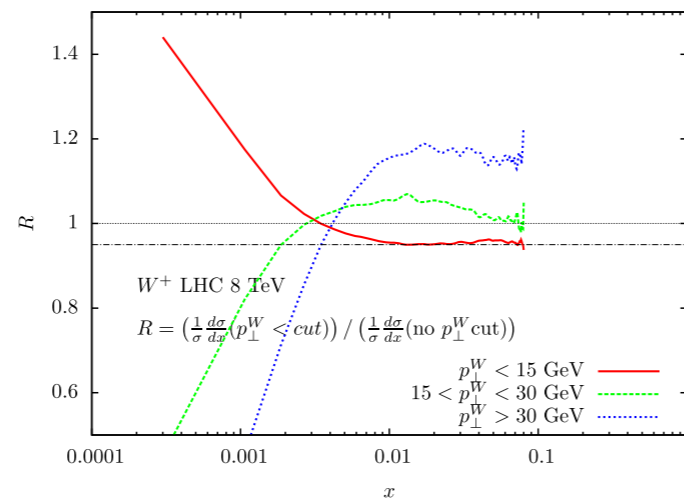
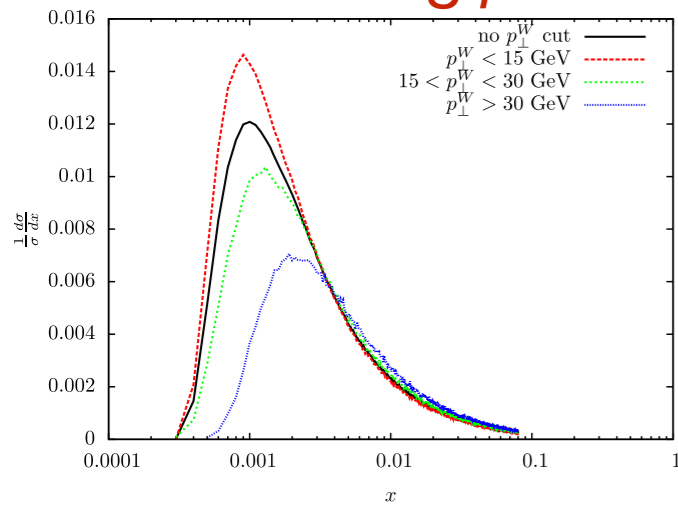
Acceptance cuts: interesting insights

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

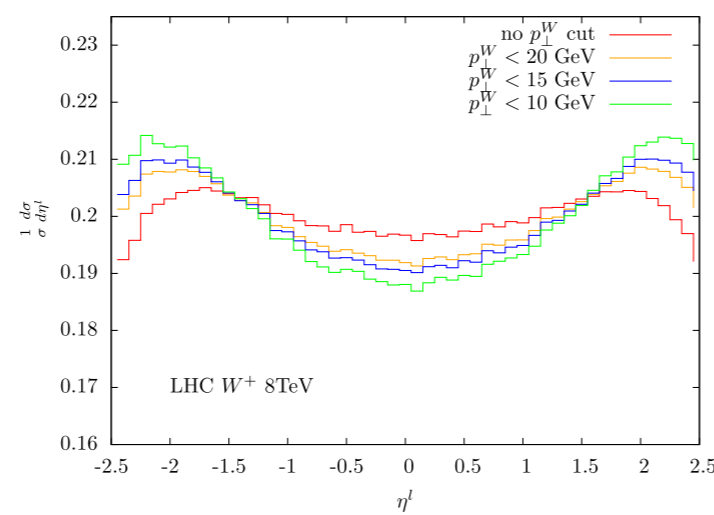
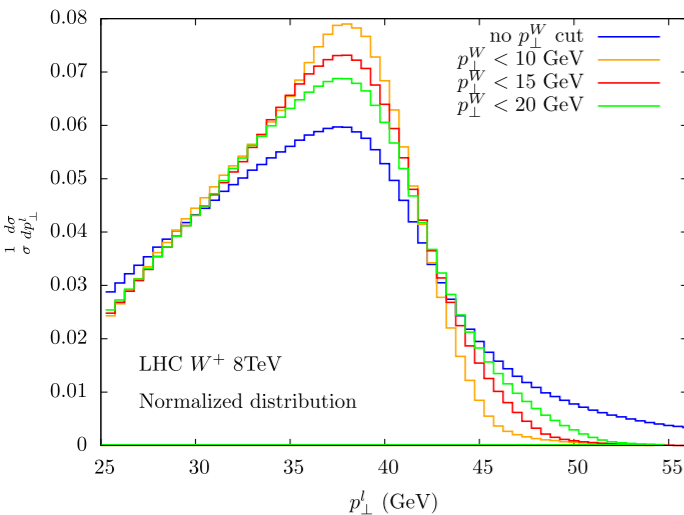
normalized distributions			
cut on p_{\perp}^W	cut on $ \eta_l$	CT10	NNPDF3.0
inclusive	$ \eta_l < 2.5$	$80.400 + 0.032 - 0.027$	80.398 ± 0.014
$p_{\perp}^W < 20$ GeV	$ \eta_l < 2.5$	$80.396 + 0.027 - 0.020$	80.394 ± 0.012
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 10$ GeV	$ \eta_l < 2.5$	$80.392 + 0.015 - 0.012$	80.394 ± 0.007
$p_{\perp}^W < 15$ GeV	$ \eta_l < 1.0$	$80.400 + 0.032 - 0.021$	80.406 ± 0.017
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 15$ GeV	$ \eta_l < 4.9$	$80.400 + 0.009 - 0.004$	80.401 ± 0.003
$p_{\perp}^W < 15$ GeV	$1.0 < \eta_l < 2.5$	$80.392 + 0.025 - 0.018$	80.388 ± 0.012



strong p_{TW} cut reduces M_W uncertainty



suppression of the large-x region



steeper shape of the p_{Tl} distribution

enhancement of high rapidity regions

Acceptance cuts: interesting insights

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

normalized distributions			
cut on p_{\perp}^W	cut on $ \eta_l $	CT10	NNPDF3.0
inclusive	$ \eta_l < 2.5$	$80.400 + 0.032 - 0.027$	80.398 ± 0.014
$p_{\perp}^W < 20$ GeV	$ \eta_l < 2.5$	$80.396 + 0.027 - 0.020$	80.394 ± 0.012
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 10$ GeV	$ \eta_l < 2.5$	$80.392 + 0.015 - 0.012$	80.394 ± 0.007
$p_{\perp}^W < 15$ GeV	$ \eta_l < 1.0$	$80.400 + 0.032 - 0.021$	80.406 ± 0.017
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 15$ GeV	$ \eta_l < 4.9$	$80.400 + 0.009 - 0.004$	80.401 ± 0.003
$p_{\perp}^W < 15$ GeV	$1.0 < \eta_l < 2.5$	$80.392 + 0.025 - 0.018$	80.388 ± 0.012

Acceptance cuts: interesting insights

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

normalized distributions			
cut on p_{\perp}^W	cut on $ \eta_l $	CT10	NNPDF3.0
inclusive	$ \eta_l < 2.5$	$80.400 + 0.032 - 0.027$	80.398 ± 0.014
$p_{\perp}^W < 20$ GeV	$ \eta_l < 2.5$	$80.396 + 0.027 - 0.020$	80.394 ± 0.012
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 10$ GeV	$ \eta_l < 2.5$	$80.392 + 0.015 - 0.012$	80.394 ± 0.007
$p_{\perp}^W < 15$ GeV	$ \eta_l < 1.0$	$80.400 + 0.032 - 0.021$	80.406 ± 0.017
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 15$ GeV	$ \eta_l < 4.9$	$80.400 + 0.009 - 0.004$	80.401 ± 0.003
$p_{\perp}^W < 15$ GeV	$1.0 < \eta_l < 2.5$	$80.392 + 0.025 - 0.018$	80.388 ± 0.012

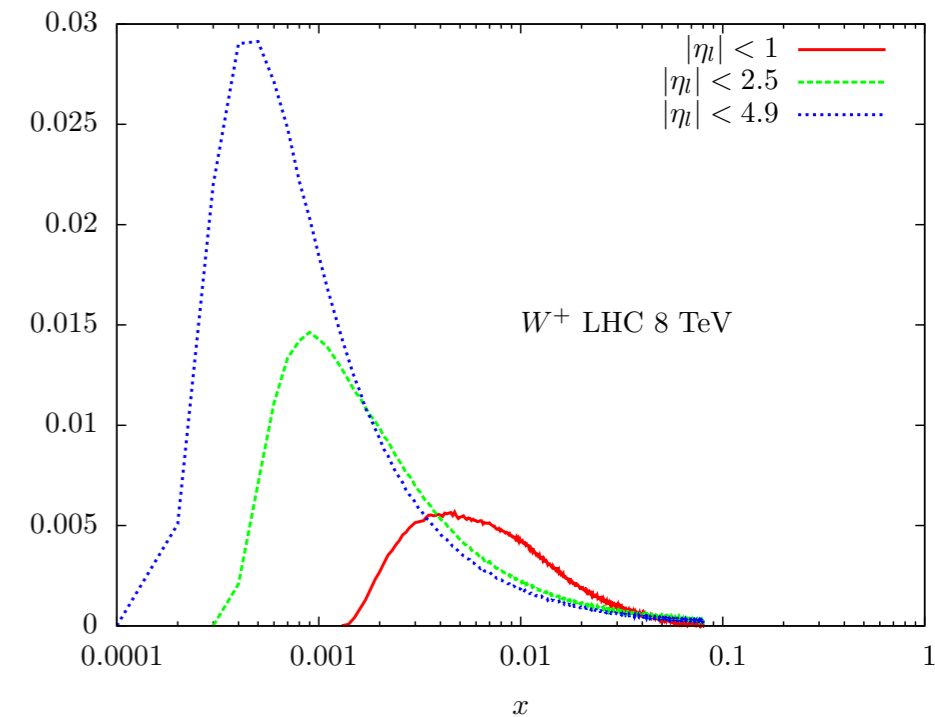
loose lepton pseudorapidity cut reduces M_W uncertainty

- uncertainties for ($\eta < 1$) and for ($1 < \eta < 2.5$)
are *separately larger* than for ($\eta < 2.5$)

Acceptance cuts: interesting insights

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

normalized distributions			
cut on p_{\perp}^W	cut on $ \eta_l $	CT10	NNPDF3.0
inclusive	$ \eta_l < 2.5$	$80.400 + 0.032 - 0.027$	80.398 ± 0.014
$p_{\perp}^W < 20$ GeV	$ \eta_l < 2.5$	$80.396 + 0.027 - 0.020$	80.394 ± 0.012
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 10$ GeV	$ \eta_l < 2.5$	$80.392 + 0.015 - 0.012$	80.394 ± 0.007
$p_{\perp}^W < 15$ GeV	$ \eta_l < 1.0$	$80.400 + 0.032 - 0.021$	80.406 ± 0.017
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 15$ GeV	$ \eta_l < 4.9$	$80.400 + 0.009 - 0.004$	80.401 ± 0.003
$p_{\perp}^W < 15$ GeV	$1.0 < \eta_l < 2.5$	$80.392 + 0.025 - 0.018$	80.388 ± 0.012



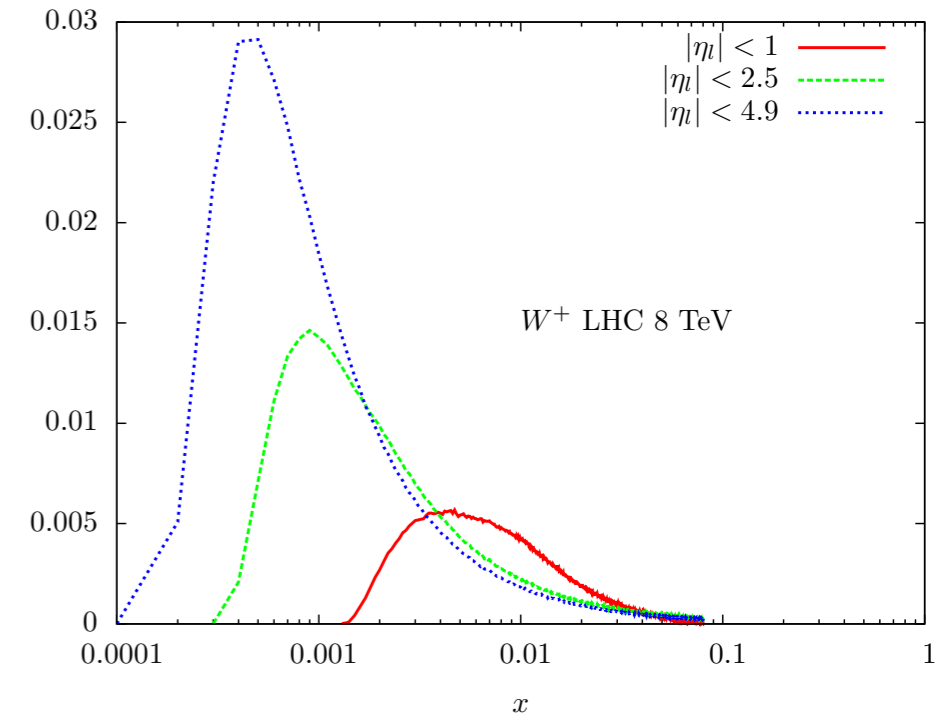
loose lepton pseudorapidity cut reduces M_W uncertainty

- uncertainties for ($\eta < 1$) and for ($1 < \eta < 2.5$) are *separately larger* than for ($\eta < 2.5$)
- normalized p_{Tl} distribution, integrated over whole rapidity range, does not depend on x

Acceptance cuts: interesting insights

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

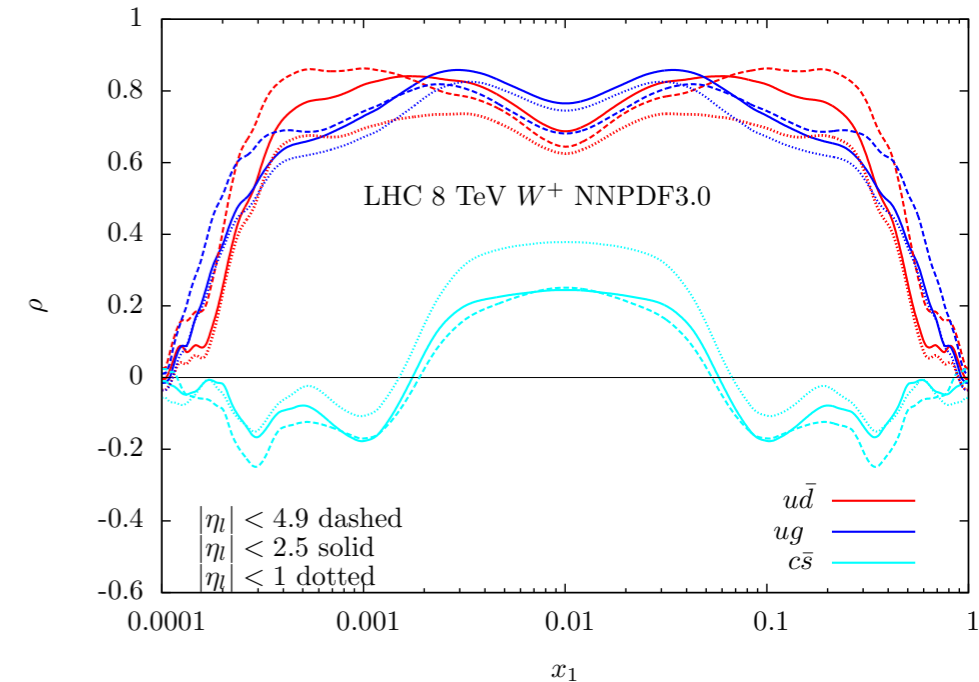
normalized distributions			
cut on p_{\perp}^W	cut on $ \eta_l $	CT10	NNPDF3.0
inclusive	$ \eta_l < 2.5$	$80.400 + 0.032 - 0.027$	80.398 ± 0.014
$p_{\perp}^W < 20$ GeV	$ \eta_l < 2.5$	$80.396 + 0.027 - 0.020$	80.394 ± 0.012
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 10$ GeV	$ \eta_l < 2.5$	$80.392 + 0.015 - 0.012$	80.394 ± 0.007
$p_{\perp}^W < 15$ GeV	$ \eta_l < 1.0$	$80.400 + 0.032 - 0.021$	80.406 ± 0.017
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 15$ GeV	$ \eta_l < 4.9$	$80.400 + 0.009 - 0.004$	80.401 ± 0.003
$p_{\perp}^W < 15$ GeV	$1.0 < \eta_l < 2.5$	$80.392 + 0.025 - 0.018$	80.388 ± 0.012



loose lepton pseudorapidity cut reduces M_W uncertainty

- uncertainties for ($\eta < 1$) and for ($1 < \eta < 2.5$) are *separately larger* than for ($\eta < 2.5$)
- normalized p_{Tl} distribution, integrated over whole rapidity range, does not depend on x

correlation of parton luminosities within the 40.5 GeV p_{Tl} bin

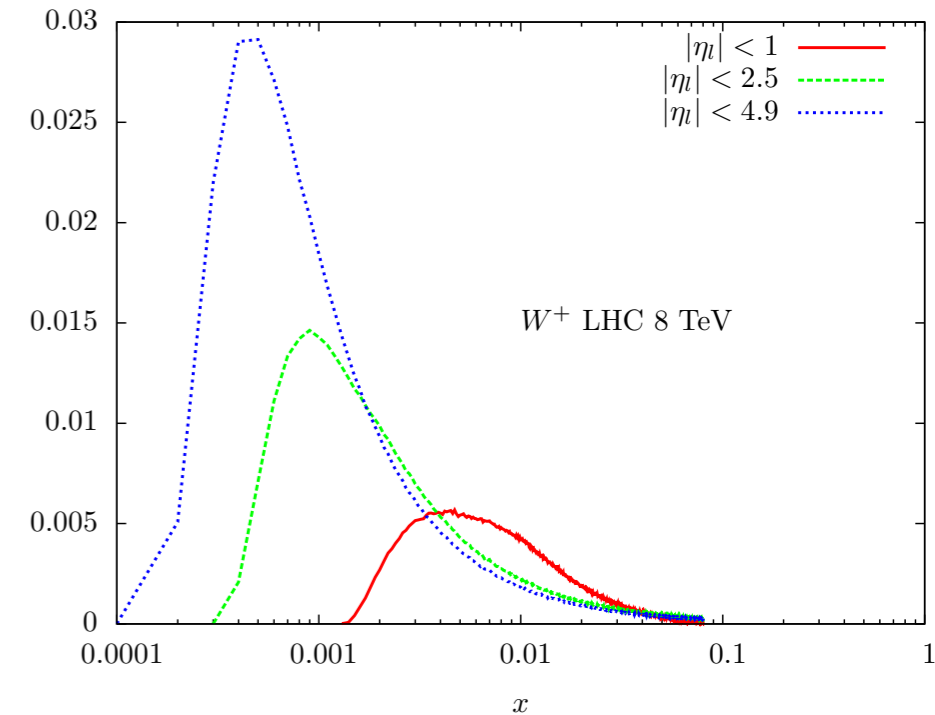


$$\rho(x, \tau) = \frac{\langle \mathcal{P}_{ij}(x, \tau) \frac{d\sigma}{dp_{\perp}^l} \rangle - \langle \mathcal{P}_{ij}(x, \tau) \rangle \langle \frac{d\sigma}{dp_{\perp}^l} \rangle}{\sigma_{\mathcal{P}_{ij}}^{\text{PDF}} \sigma_{d\sigma/dp_{\perp}^l}^{\text{PDF}}},$$

Acceptance cuts: interesting insights

Bozzi, Citelli, Vicini PRD 91, 113005 (2015)

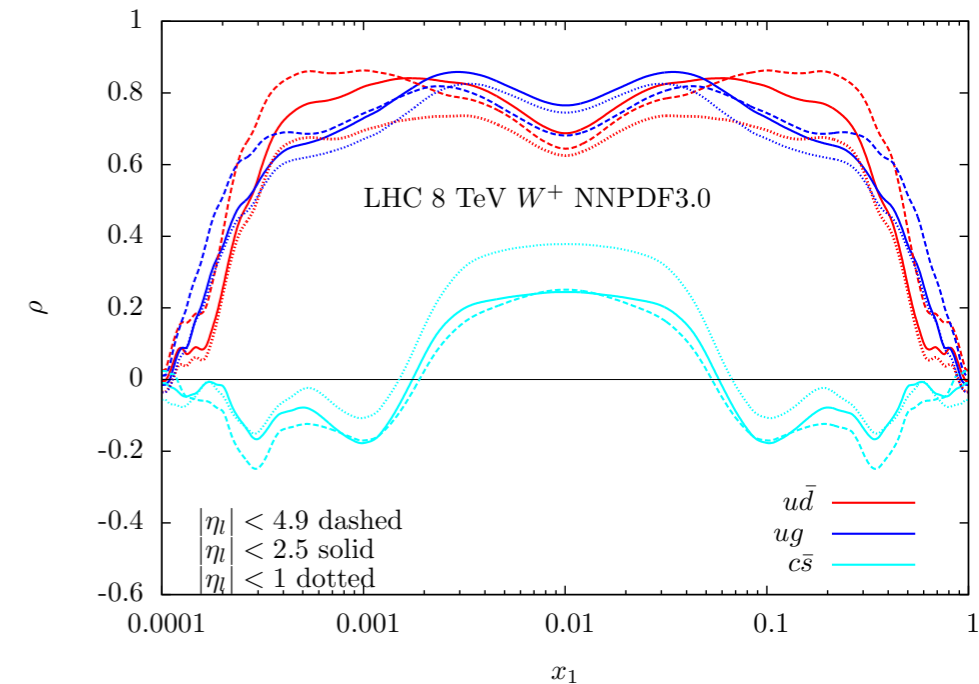
normalized distributions			
cut on p_{\perp}^W	cut on $ \eta_l $	CT10	NNPDF3.0
inclusive	$ \eta_l < 2.5$	$80.400 + 0.032 - 0.027$	80.398 ± 0.014
$p_{\perp}^W < 20$ GeV	$ \eta_l < 2.5$	$80.396 + 0.027 - 0.020$	80.394 ± 0.012
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 10$ GeV	$ \eta_l < 2.5$	$80.392 + 0.015 - 0.012$	80.394 ± 0.007
$p_{\perp}^W < 15$ GeV	$ \eta_l < 1.0$	$80.400 + 0.032 - 0.021$	80.406 ± 0.017
$p_{\perp}^W < 15$ GeV	$ \eta_l < 2.5$	$80.396 + 0.017 - 0.018$	80.395 ± 0.009
$p_{\perp}^W < 15$ GeV	$ \eta_l < 4.9$	$80.400 + 0.009 - 0.004$	80.401 ± 0.003
$p_{\perp}^W < 15$ GeV	$1.0 < \eta_l < 2.5$	$80.392 + 0.025 - 0.018$	80.388 ± 0.012



loose lepton pseudorapidity cut reduces M_W uncertainty

- uncertainties for ($\eta < 1$) and for ($1 < \eta < 2.5$) are *separately larger* than for ($\eta < 2.5$)
- normalized p_{Tl} distribution, integrated over whole rapidity range, does not depend on x
- PDF sum rules \rightarrow *non trivial compensations between different rapidity intervals among different flavours*

correlation of parton luminosities within the 40.5 GeV p_{Tl} bin



$$\rho(x, \tau) = \frac{\langle \mathcal{P}_{ij}(x, \tau) \frac{d\sigma}{dp_{\perp}^l} \rangle - \langle \mathcal{P}_{ij}(x, \tau) \rangle \langle \frac{d\sigma}{dp_{\perp}^l} \rangle}{\sigma_{\mathcal{P}_{ij}}^{\text{PDF}} \sigma_{d\sigma/dp_{\perp}^l}^{\text{PDF}}},$$

Choice of NP parameters

Choice of NP parameters

$$\frac{d\sigma}{dq_T} \sim \text{FT} \exp\{-g_{NP} b_T^2\}$$

Choice of NP parameters

$$\frac{d\sigma}{dq_T} \sim \text{FT} \exp\{-g_{NP} b_T^2\}$$



Fit to Z/γ^* Tevatron data: $g_{NP} \sim 0.8 \text{ GeV}^2$

[Guzzi, Nadolsky, Wang (2014)]

Choice of NP parameters

$$\frac{d\sigma}{dq_T} \sim \text{FT} \exp\{-g_{NP} b_T^2\} \longrightarrow \text{Fit to } Z/\gamma^* \text{ Tevatron data: } g_{NP} \sim 0.8 \text{ GeV}^2$$

[Guzzi, Nadolsky, Wang (2014)]

For each TMD: $0.4 \text{ GeV}^2 \sim g_{NP}^a \longrightarrow g_{evo} \ln\left(\frac{Q^2}{Q_0^2}\right) + g_a$

Choice of NP parameters

$$\frac{d\sigma}{dq_T} \sim \text{FT} \exp\{-g_{NP} b_T^2\} \longrightarrow \text{Fit to } Z/\gamma^* \text{ Tevatron data: } g_{NP} \sim 0.8 \text{ GeV}^2$$

[Guzzi, Nadolsky, Wang (2014)]

For each TMD: $0.4 \text{ GeV}^2 \sim g_{NP}^a \longrightarrow g_{evo} \ln\left(\frac{Q^2}{Q_0^2}\right) + g_a$

$$\text{Fit to SIDIS/DY/Z data: } g_{evo} \ln\left(\frac{Q^2}{Q_0^2}\right) \in [0.17, 0.39] \text{ GeV}^2$$

[Bacchetta, Delcarro, Pisano, Radici, Signori (2017)]

Choice of NP parameters

$$\frac{d\sigma}{dq_T} \sim \text{FT} \exp\{-g_{NP} b_T^2\} \longrightarrow \text{Fit to } Z/\gamma^* \text{ Tevatron data: } g_{NP} \sim 0.8 \text{ GeV}^2$$

[Guzzi, Nadolsky, Wang (2014)]

For each TMD: $0.4 \text{ GeV}^2 \sim g_{NP}^a \longrightarrow g_{evo} \ln\left(\frac{Q^2}{Q_0^2}\right) + g_a \longrightarrow \text{variation range for } g_a$

Fit to SIDIS/DY/Z data: $g_{evo} \ln\left(\frac{Q^2}{Q_0^2}\right) \in [0.17, 0.39] \text{ GeV}^2$

[Bacchetta, Delcarro, Pisano, Radici, Signori (2017)]

Choice of NP parameters

$$\frac{d\sigma}{dq_T} \sim \text{FT} \exp\{-g_{NP} b_T^2\} \longrightarrow \text{Fit to } Z/\gamma^* \text{ Tevatron data: } g_{NP} \sim 0.8 \text{ GeV}^2$$

[Guzzi, Nadolsky, Wang (2014)]

For each TMD: $0.4 \text{ GeV}^2 \sim g_{NP}^a \longrightarrow g_{evo} \ln\left(\frac{Q^2}{Q_0^2}\right) + g_a \longrightarrow \text{variation range for } g_a$

Fit to SIDIS/DY/Z data: $g_{evo} \ln\left(\frac{Q^2}{Q_0^2}\right) \in [0.17, 0.39] \text{ GeV}^2$

[Bacchetta, Delcarro, Pisano, Radici, Signori (2017)]

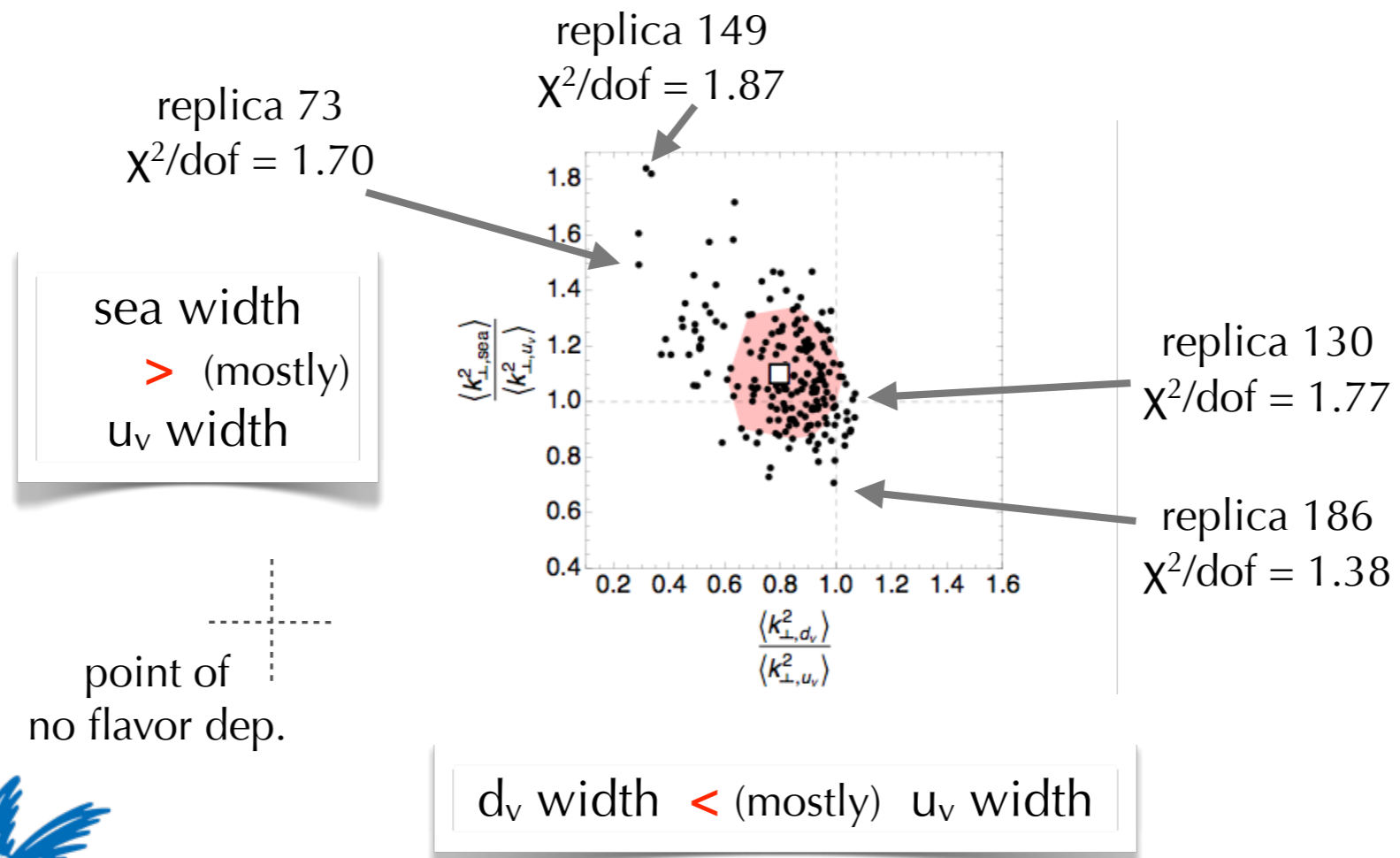
We consider :

- **50 flavour-dependent sets** $\{g_{NP}^{u_v}, g_{NP}^{d_v}, g_{NP}^{u_s}, g_{NP}^{d_s}, g_{NP}^s\}$ with $g_{NP}^a \in [0.2, 0.6] \text{ GeV}^2$
- **1 flavour-independent set** with $g_{NP}^a = 0.4 \text{ GeV}^2$

Extraction of parameters from SIDIS

Signori, Bacchetta, Radici, Schnell, JHEP 1311, 194 (2013)

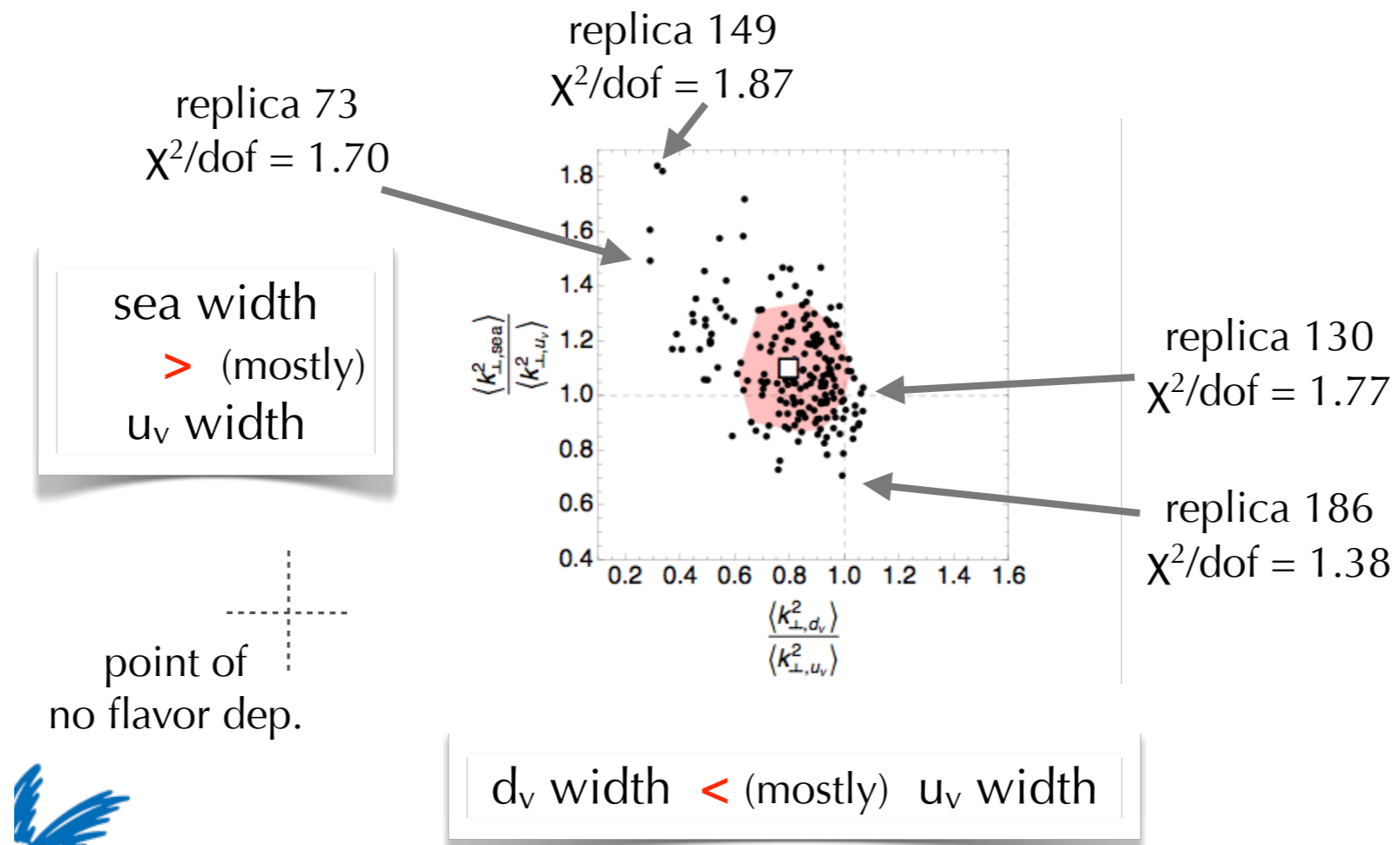
template fit on HERMES data: distribution of parameters



Extraction of parameters from SIDIS

Signori, Bacchetta, Radici, Schnell, JHEP 1311, 194 (2013)

template fit on HERMES data: distribution of parameters



On average, $sea > u_v > d_v$

flav-dep vs. flav-indep set

