

Teoria – Alte Energie

progetti a Roma 1

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INFN, Rome 1 unit

Assisi, 16 Giugno 2019



High with respect to what?

My own answer: wrt to the “scale of the theory”:

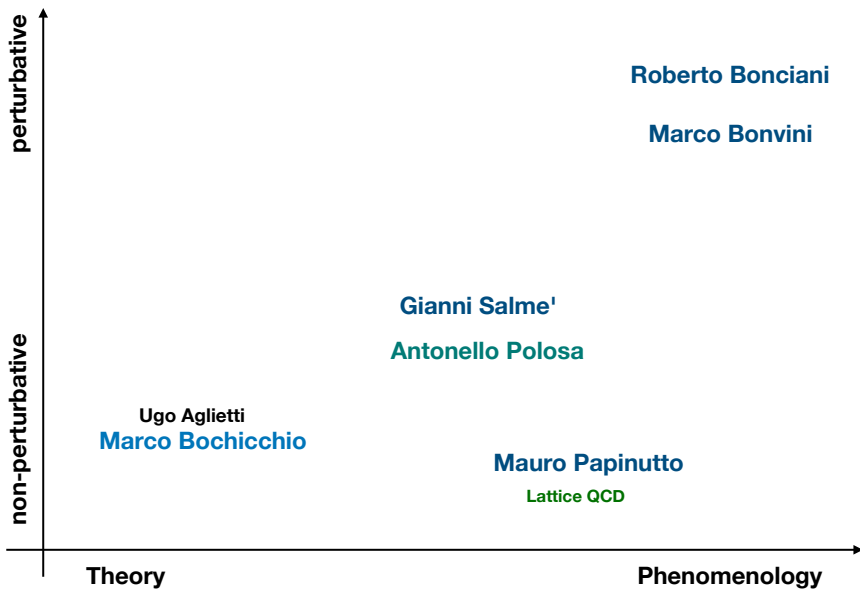
- QCD: $Q \gg \Lambda_{\text{QCD}} \sim 0.2 \text{ GeV}$
- EW: $Q \gg v \sim 250 \text{ GeV}$

This allows me to include the activity of more people in Rome

Flavour physics is excluded (tomorrow morning's session)

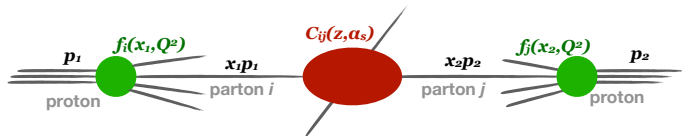
Overview of the theorists in Rome 1 + Sapienza





Collinear factorization theorem in QCD:

$$\frac{d\sigma}{dQ^2 dy dp_t \dots} \simeq \sum_{i,j=g,q} \int_0^1 dx_1 dx_2 f_i(x_1, Q^2) f_j(x_2, Q^2) C_{ij}(x_1, x_2, y, p_t, \dots, \alpha_s(Q^2))$$



- partonic cross sections $C_{ij}(\dots, \alpha_s)$ (observable-dependent, perturbative)
- parton distribution functions (PDFs) $f_i(x, Q^2)$ (universal, non-perturbative)

There exist generalizations (transverse momentum dependent factorization, double parton scattering, ...)

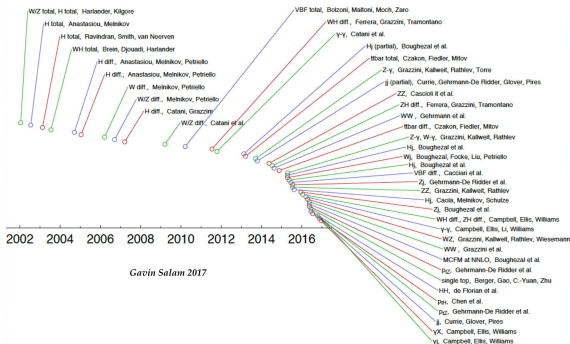
Perturbative computations of the partonic cross sections

$$C(\dots, \alpha_s) = \underbrace{C^{(0)}(\dots)}_{\text{LO}} + \underbrace{\alpha_s C^{(1)}(\dots)}_{\text{NLO}} + \underbrace{\alpha_s^2 C^{(2)}(\dots)}_{\text{NNLO}} + \underbrace{\alpha_s^3 C^{(3)}(\dots)}_{\text{N}^3\text{LO}} + \dots$$

LO: trivial

NLO revolution: automation (MG5_aMC@NLO and others, matching to parton showers)

NNLO revolution:

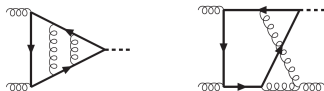


N³LO: inclusive total cross section for Higgs production in gluon fusion

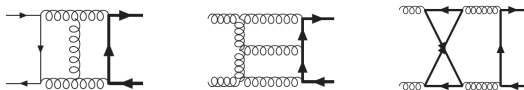
Analytic computations are faster and more powerful than numerical ones

Need to reduce to master integrals, and solve differential equations (sometimes elliptic)

- Higgs production: mass corrections to total XS at NNLO or to $H + j$ at NLO (sensitive to new physics in loop at medium/high p_t)



- top pair production at NNLO (asymmetries and top properties, constraints on the gluon PDF at large x)



- di-jet production at NNLO (determination of α_s and constraints on PDFs at medium-large x)



In general, perturbative coefficients contain logarithms of dimensionless ratios

$$L = \log(\text{something})$$

Sometimes, they are logarithmically enhanced:

$$\begin{aligned} C(\dots, \alpha_s) = & a_0 \\ & + \alpha_s [a_1 L + b_1] \\ & + \alpha_s^2 [a_2 L^2 + b_2 L + c_2] \\ & + \alpha_s^3 [a_3 L^3 + b_3 L^2 + c_3 L + d_3] \\ & + \alpha_s^4 [a_4 L^4 + b_4 L^3 + c_4 L^2 + d_4 L + e_4] \\ & + \dots \end{aligned}$$

If/when $\alpha_s L \sim 1$ the fixed-order expansion is no longer predictive!

Resum the logs, and convert to a “logarithmic-order” expansion:

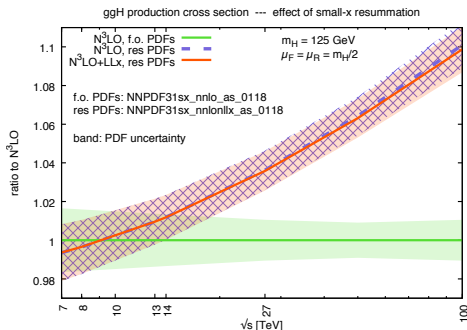
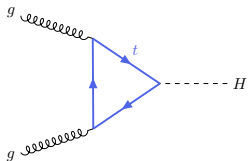
$$g_{\text{LL}}(\alpha_s L) + \alpha_s g_{\text{NLL}}(\alpha_s L) + \alpha_s^2 g_{\text{NNLL}}(\alpha_s L) + \alpha_s^3 g_{\text{N}^3\text{LL}}(\alpha_s L) + \dots$$

Leading log (LL), next-to-leading log (NLL), next-to-next-to-leading log (NNLL)...

Past and current activities:

- threshold logarithms
- small p_t logarithms
- collinear logarithms in presence of massive (heavy) quarks
- small- x (high-energy) resummation

Applications e.g. to Higgs physics



M.Bonvini: also interested in theoretical uncertainties

... with DGLAP evolution

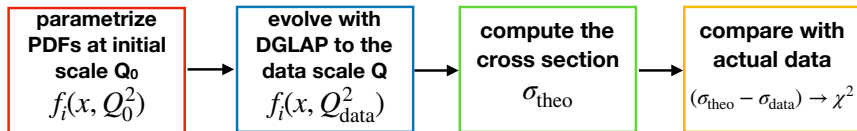
$$Q^2 \frac{d}{dQ^2} f_i(x, Q^2) = \sum_{j=g,q} \int_x^1 \frac{dz}{z} P_{ij}(z, \alpha_s(Q^2)) f_j\left(\frac{x}{z}, Q^2\right)$$

Splitting functions $P_{ij}(x, \alpha_s)$ (universal, perturbative)

PDFs at a given scale Q_0 + DGLAP evolution \rightarrow PDFs at any scale Q

$f_i(x, Q_0^2)$ at low $Q_0 \sim 1$ GeV is non-perturbative: how to determine?

Standard approach: determine PDFs by comparison with (many) data



Activity [M.Bonvini]: include all-order resummations in PDF fits

Non-perturbative problem \rightarrow numerical simulations on a discretized spacetime
 However, the field-theoretic definition of PDFs involves light-cone distances:

$$f_q(x, Q^2) = \int \frac{d\xi}{4\pi} e^{-ix\xi P^+} \langle P | \bar{\psi}_q(\xi) \not{n} U_n(\xi, 0) \psi_q(0) | P \rangle \quad \xi^2 = 0$$

but in lattice QCD simulation the spacetime is Euclidean, where light-cone separations are only trivial ($\xi = 0$) \rightarrow **PDFs cannot be computed in lattice QCD!**

Need to work directly in Minkowski space

A possible approach makes use of the Bethe-Salpeter equation to study a relativistic bound state, which can account for non-perturbative effects

$$G_{4pt}(k, q; p) \sim \frac{\chi(k, p)\chi(q, p)}{2\omega(p^0 - \omega + i\epsilon)} \quad \chi(k, p) = G_0(k, p) \int d^4k' I(k, k', p)\chi(k', p)$$

Working in the Euclidean is simpler, but staying in Minkowski, though more difficult, is now achievable thanks to the so-called Light-Front framework

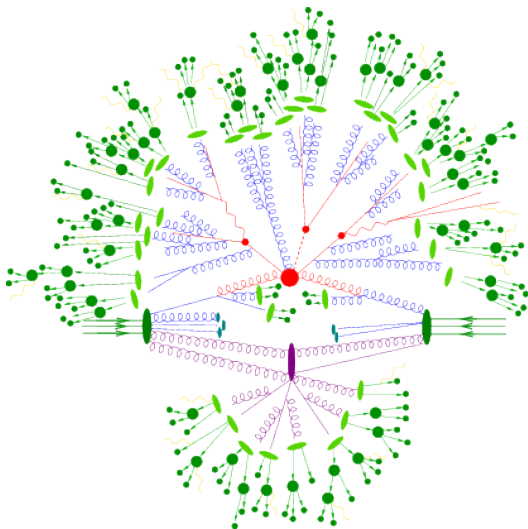
Obtaining the PDFs is the ultimate goal of such a business
 (transverse momentum dependent PDFs could be obtained similarly)

From initial state to final state

So far:
PDFs and hard scattering

Afterwards:
parton shower

Finally:
hadronization and decays



Belle, BES, BaBar, LHCb,... have found a large number of exotic resonances, which appear to be composed of 4 or 5 quarks

Interpretation as bound states of di-quarks explains:

- mixing
- degeneracies
- decay rates
- spin interactions

Dynamical description of how QCD produces these bound states

State	m (MeV)	Γ (MeV)	J^{PC}	Process (mode)	References
$X(3872)$	3871.69 ± 0.17	< 1.2	1^{++}	$B \rightarrow K(\pi^+ \pi^- J/\psi)$ $p\bar{p} \rightarrow (\pi^+ \pi^- J/\psi) + \dots$ $e^+e^- \rightarrow \gamma(\pi^+ \pi^- J/\psi)$ $B \rightarrow K(\omega J/\psi)$ $B \rightarrow K(D^{*0} \bar{D}^0)$ $B \rightarrow K(\gamma J/\psi)$ and $B \rightarrow K(\gamma \psi(2S))$	Belle [10,32], BaBar [36], LHCb [34,73] CDF [31,74,75], D0 [76] BES III [77] Belle [78], BaBar [33] Belle [38,79], BaBar [37] Belle [29], BaBar [30], LHCb [40]
$Z_c(3900)^+$	3888.7 ± 3.4	35 ± 7	1^+	$e^+e^- \rightarrow (J/\psi \pi^+) \pi^-$ $e^+e^- \rightarrow (D\bar{D}^*)^+ \pi^-$	Belle [43], BES III [55] BES III [56]
$X(3915)$	3915.6 ± 3.1	28 ± 10	$0/2^{++}$	$B \rightarrow K(\omega J/\psi)$ $e^+e^- \rightarrow e^+e^-(\omega J/\psi)$ $e^+e^- \rightarrow J/\psi(DD^*)$ $e^+e^- \rightarrow J/\psi(\dots)$	Belle [80], BaBar [33] Belle [81], BaBar [82] Belle [83] Belle [84]
$Y(4008)$	3891 ± 42	255 ± 42	1^{--}	$e^+e^- \rightarrow \gamma(\pi^+ \pi^- J/\psi)$	Belle [42,43]
$Z_c(4050)^+$	4051^{+24}_{-43}	82^{+51}_{-55}	?	$B \rightarrow K(\pi^+ \chi_{c1}(1P))$	Belle [53], BaBar [54]
$X(4050)^+$	4054 ± 3	45	?	$e^+e^- \rightarrow (\pi^+ \psi(2S)) \pi^-$	Belle [85]
$Y(4140)$	4143.4 ± 3.0	15^{+11}_{-7}	$?^{2+}$	$B \rightarrow K(\phi J/\psi)$	CDF [72], D0 [86]
$X(4160)$	4156^{+29}_{-25}	139^{+113}_{-65}	$?^{2+}$	$e^+e^- \rightarrow J/\psi(D\bar{D}^*)$	Belle [83]
$Z_c(4200)^+$	4196^{+35}_{-32}	370^{+99}_{-140}	?	$B \rightarrow K(\pi^+ J/\psi)$	Belle [87]
$Z_c(4250)^+$	4248^{+185}_{-45}	177^{+321}_{-72}	?	$B \rightarrow K(\pi^+ \chi_{c1}(1P))$	Belle [53], BaBar [54]
$Y(4260)$	4263 ± 5	108 ± 14	1^{--}	$e^+e^- \rightarrow \gamma(\pi^+ \pi^- J/\psi)$ $e^+e^- \rightarrow (\pi^+ \pi^- J/\psi)$ $e^+e^- \rightarrow (\pi^0 \pi^0 J/\psi)$	BaBar [41,88], CLEO [89], Belle [42,43] CLEO [47], BES III [56] CLEO [47]
$X(4350)$	$4350.6^{+4.6}_{-5.1}$	$13.3^{+18.4}_{-10.0}$	$?^{2+}$	$e^+e^- \rightarrow e^+e^-(\phi J/\psi)$	Belle [90]
$Y(4360)$	4361 ± 13	74 ± 18	1^{--}	$e^+e^- \rightarrow \gamma(\pi^+ \pi^- \psi(2S))$	BaBar [44], Belle [45,85]
$Z_c(4430)^+$	4485^{+36}_{-25}	200^{+80}_{-58}	1^+	$B \rightarrow K(\pi^+ \psi(2S))$ $B \rightarrow K(\pi^+ J/\psi)$	Belle [49,51,52], BaBar [50], LHCb [21] Belle [87], BaBar [50]
$X(4630)$	4634^{+9}_{-11}	92^{+41}_{-32}	1^{--}	$e^+e^- \rightarrow \gamma(\Lambda_c^+ \Lambda_c^-)$	Belle [91]
$Y(4660)$	4664 ± 12	48 ± 15	1^{--}	$e^+e^- \rightarrow \gamma(\pi^+ \pi^- \psi(2S))$	Belle [45]
$Z_c(10610)^+$	10607.2 ± 2.0	18.4 ± 2.4	1^+	$e^+e^- \rightarrow (b\bar{b} \pi^+) \pi^-$	Belle [20]
$Z_b(10610)^0$	$10609 \pm 4 \pm 4$	NA	1^{+-}	$e^+e^- \rightarrow (\Upsilon(2, 3S) \pi^0) \pi^0$	Belle [23]
$Z_b(10650)^+$	10652.2 ± 1.5	11.5 ± 2.2	1^+	$e^+e^- \rightarrow (b\bar{b} \pi^+) \pi^-$	Belle [20]
$Y_b(10888)$	10888.4 ± 3.0	$30.7^{+5.9}_{-3.7}$	1^{--}	$e^+e^- \rightarrow (\pi^+ \pi^- \Upsilon(nS))$	Belle [60,62]

A.Polosa: also carbon nanotubes and graphene as directional detectors of light Dark Matter and effective theories of liquid helium for Dark Matter detectors

Find solution of the large- N limit of QCD

- theory of quark and gluons \rightarrow theory of glueballs and mesons
- strong coupling $\alpha_s \rightarrow$ weak coupling $1/N$

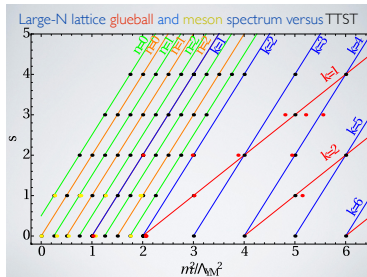
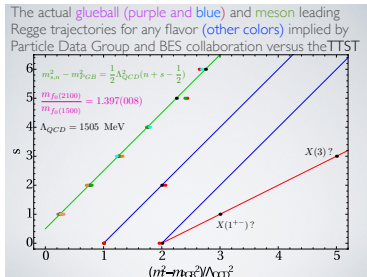
The solution must possess two fundamental properties:

- asymptotic freedom
- renormalizability (renormalization group)

Canonical string theory (glueballs:closed mesons:open strings) cannot be a solution

\rightarrow **Non-canonical string solutions** (twistor strings)

It is the only presently known string theory that may be consistent, both for the ultraviolet asymptotics and the spectrum, with the fundamental principles of QCD

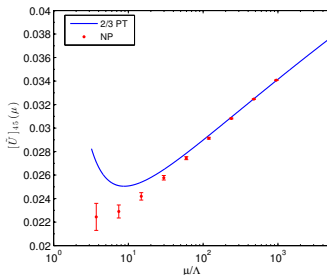


In Rome1 + Rome3 a strong lattice QCD group, mostly interested in flavour physics
I assume it will be covered tomorrow morning

Important: non-perturbative renormalization

[M.Papinutto]

At low scales computing Wilson coefficients perturbatively can miss important higher-order and non-perturbative effects



Future project: lattice simulations of large- N QCD

[M.Papinutto+M.Bochicchio]

String theory is a consistent quantum theory that contains gravity
General relativity arises in string theory as a Wilsonian effective theory

Local activity:

- understanding the general structure of the low energy theories in four dimensions (new results have been obtained in this directions in the past couple of years — fluxes and exotic branes)
- how to obtain inflation and a vacuum with a positive cosmological constant
- a microscopic understanding of the Bekenstein-Hawking entropy of black holes

Many theoretic ideas were born in this field of research, like large extra dimensions, holography, techniques to compute scattering amplitudes, ... and in this sense this line of research is very fruitful although not in direct contact with experiments

Two classes of motivations for BSM

- internal consistency issues of the SM (hierarchy, naturalness, strong CP, ...)
- experimental facts (dark matter, baryon asymmetry, neutrino masses, ...)

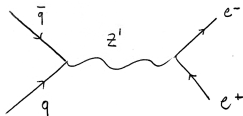
Two approaches to BSM

- model building (to address some of the above issues)
- phenomenology (either of a single model, or using SM-EFT)

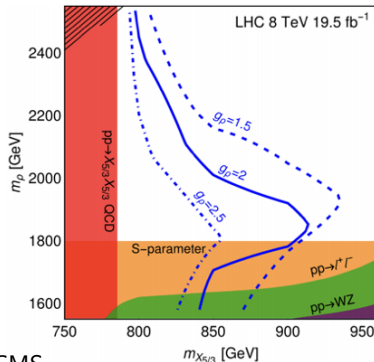
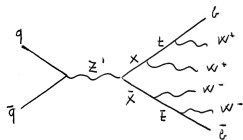
Activities and expertise:

- (non-minimal) composite Higgs [hierarchy/naturalness]
- (non-minimal) SUSY [hierarchy/naturalness]
- WIMPs [DM + hierarchy/naturalness]
- hidden sector [DM + hierarchy/naturalness]
- technicolor [DM + neutrino mass]
-
- and of course flavour...

Conventional signature, $\sim g_w/g_\rho$ (suppressed)



Non-conventional signature, $\sim g_\rho$ (enhanced)



This work triggered an experimental analysis in CMS

Current interests:

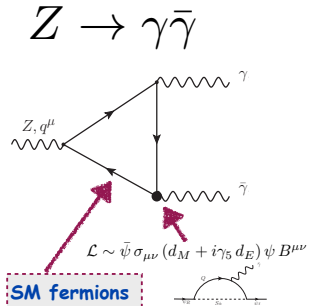
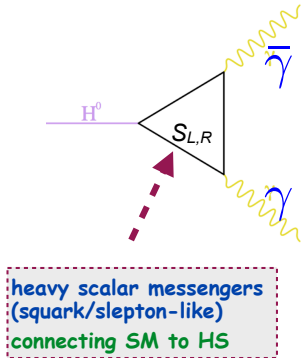
- BSM models for dark matter and/or baryogenesis and/or neutrino masses, with interests to theories with first-order phase transitions leading to gravitational waves
- SM-EFT in EW and top sectors, with interests on the recent flavour anomalies

Member of CMS as theory collaborator, participating to light ($\lesssim 30$ GeV) scalar searches

A model of dark fermions with an unbroken $U(1)$ symmetry to explain

- dark matter
- flavour hierarchies

Mono-photon signatures in Higgs and Z decays



Also new Higgs signatures at e^+e^- colliders

Several options:

- e^+e^- : ILC, CLIC (linear), FCC-ee (circular)
- $e^\pm p$: LHeC, FCC-eh
- pp : HE-LHC, FCC-hh
- $\mu^+\mu^-$: muon collider

[B.Mele, D.Barducci]

[M.Bonvini]

[B.Mele, D.Barducci, M.Bonvini]

[B.Mele]

$e^\pm p$ collider very important for PDF determination
(functional to pp colliders)

Also very interesting physics cases!
[slide from Uta Klein]

VBF Higgs Production in ep (top) and pp (bottom)

ep: Higgs production in ep comes uniquely from either CC or NC DIS via VBF

Clean bb final state, $S/B > 1$
e-h Cross Calibration for Precision ep
Clean, precise reconstruction and easy distinction of ZZH and WWH without pile-up:
 < 0.1 @ LHeC up to 1 @ FCCeh events

VBF: Small theoretical uncertainties!

pp: Higgs production in pp comes predominantly from $gg \rightarrow H$:
high rates crucial for rare decays
 LHC VBF cross section about 200 fb (about as large as at the LHeC).

Pile-up in pp at $5 \cdot 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ is 150 @ 25ns
FCC-hh: pile-up 500-1000
 S/B very small for bb

Final Precision in pp needs accurate $N^3\text{LO}$ PDFs & α_s

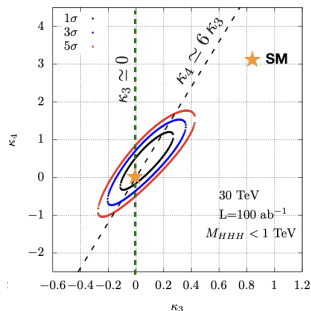
Can reach very high energies $\sqrt{s} \sim 10 \div 30$ TeV (negligible synchrotron radiation), but high luminosity is a challenge

Super clean: $\mu^+ \mu^- \rightarrow f \bar{f}$ with $m_f \lesssim \sqrt{s}/2$

Higgs self-interaction measurements at a new level of precision

$$\mathcal{L} = -\frac{1}{2} m_h^2 h^2 - \lambda_3 \frac{m_h^2}{2v} h^3 - \lambda_4 \frac{m_h^2}{8v^2} h^4$$

$$\lambda_3^{SM} = \lambda_4^{SM} = 1$$



For instance the quartic coupling λ_4 (assuming $\lambda_3 = 1$)

- FCC-hh at $\sqrt{s} = 100$ TeV for 30/ab: $\lambda_4 \in [-4, 16]$ at 68% c.l.
- muon coll at $\sqrt{s} = 30$ TeV for 100/ab: $\lambda_4 \in [0.8, 1.5]$ at 68% c.l.

Pros of the group:

- we cover various expertises
- various (inter)national connections

Cons of the group:

- we are quite disconnected/isolated locally
- very few students and post-docs

Suggestions/requests:

- improve communication among us, perhaps with a common room, a better seminar organization (e.g. with food after seminars), journal club(s), ...
- from M.Bochicchio: access to CPU hours at CINECA (issues related to assignments to iniziative specifiche)