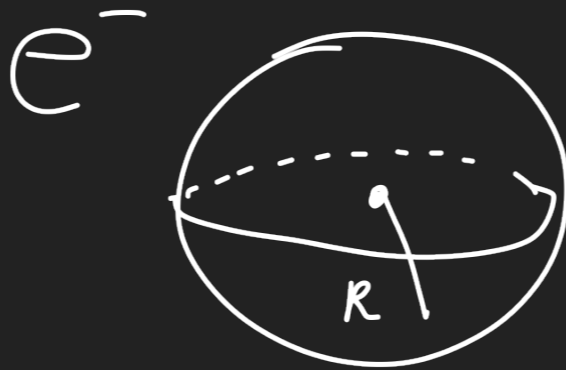


AD POLOSA

VIS. TEOR. 1

NEW PHYSICS AT ~1 MEV



(EM. CONTRIBUTION TO THE MASS OF e)

$$E_{st} \lesssim \frac{e \cdot e}{R} \lesssim mc^2 \text{ (self-energy)}$$

$$R \gtrsim \frac{e^2}{mc^2} = \hbar c \frac{\alpha}{0.5 \text{ MeV}} \approx 3 \text{ fm}$$

AS BIG AS A NUCLEUS WITH $A=27$

MAYBE IN THE QUANTUM THEORY CANCELLATIONS ARE AT WORK AS $R \rightarrow 0$.

NEW PHYSICS AT ~ 1 MEV

IF ONE ELECTRON ALONE IS PRESENT, THE PROB. OF FINDING A CHARGE DENSITY AT TWO DIFFERENT POINTS SIMULTANEOUSLY (DISTANT BY $|\vec{\xi}|$) IS ZERO.

$$\left\langle \int \rho(\vec{r} + \vec{\xi}/2) \rho(\vec{r} - \vec{\xi}/2) d^3r \right\rangle \propto \delta^3(\vec{\xi})$$

THUS

$$E_{st} \rightarrow \infty$$

NEW PHYSICS AT ~1 MEV

BELOW SOME 'CUT-OFF' DISTANCE λ THE
'POSITRON THEORY' HAS TO BE USED.

$$\left(\begin{array}{l} \text{EM. CONTRIBUTION} \\ \text{TO THE MASS OF } e \end{array} \right) E_{st} = \lim_{R \rightarrow 0} \alpha m c^2 \ln \left(\frac{\lambda_c}{R} \right) \quad (\text{Weinkopf})_{39}$$

THIS DISTANCE IS SET AT $\approx \lambda_c \approx \frac{1}{m_e c}$
— THE MASS OF THE POSITRON, THE NEW
PARTICLE OF QED

NEW PHYSICS AT ~1 MEV

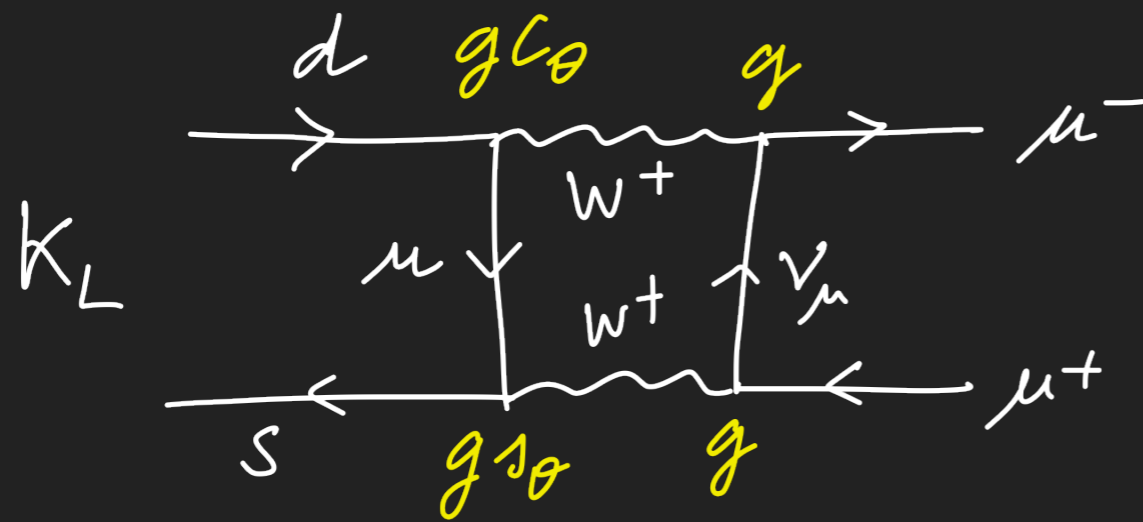
$$\propto mc^2 \ln\left(\frac{\lambda_c}{R}\right) \lesssim mc^2$$

only need that

$$R \gtrsim \lambda_c e^{-137} \quad (\text{instead of } R \gtrsim 3\text{fm})$$

QED RESTORES UNNATURALNESS:
NO MYSTERIOUS 'CANCELLATIONS' BUT A
NEW PHYSICS W/ e^- , e^+ & γ .

NEW PHYSICS AT ~1 GEV



(Ioffe & Shabalin)
'68

$$\int^{\Lambda} \frac{\left(\frac{k^\mu k^\nu}{m_W^2}\right)^2}{(k^2 - m_W^2)^2} \cdot \frac{k+m}{(k^2 - m^2)} \cdot \frac{k}{k^2} d^4 k \sim \Lambda^2$$

POWER COUNTING $\sim G_{\neq}^2 \Lambda^2$

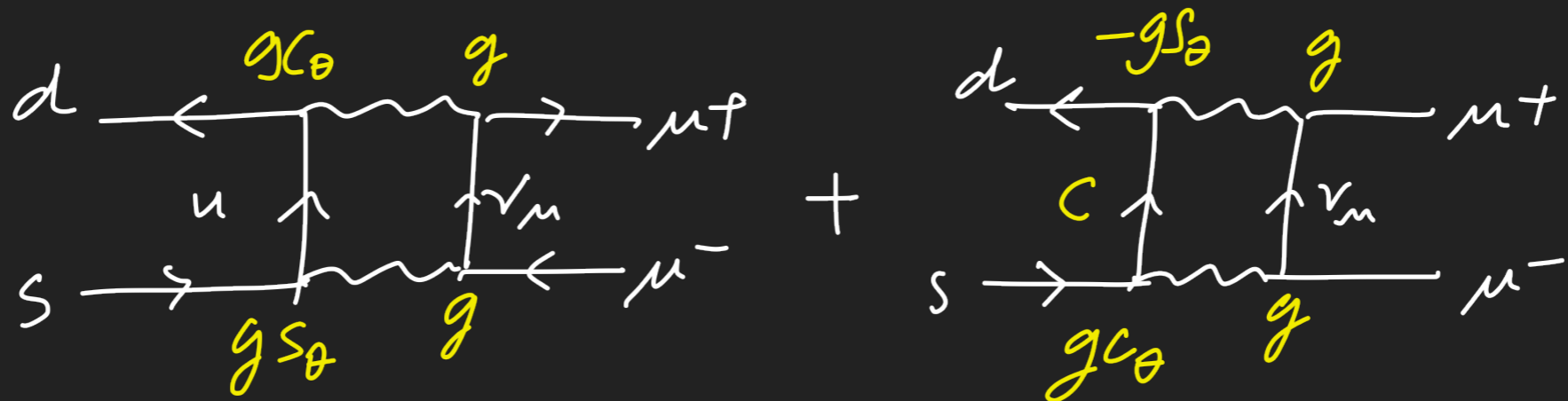
CFR w/ $\mathcal{B}(K_L \rightarrow \mu^+ \mu^-) \sim 7 \times 10^{-9}$

$\Lambda \lesssim 3 \text{ GeV}$

Does this 'cutoff' also correspond to the existence of a new pvt?

NEW PHYSICS AT ~ 1 GEV

GIM ('70) INTRODUCES A DYNAMICAL MECH.
TO SUPPRESS $K_L \rightarrow \mu^+ \mu^-$



THE DIFFERENCE IN MASS $m_c - m_u$ MAKES THE
CANCELLATION INCOMPLETE.

$$\text{Amp}(K_L \rightarrow \mu^+ \mu^-) \propto G^2 S_\theta C_\theta (m_c^2 - m_u^2)$$

A CONVERGENT RESULT

$$m_c \simeq 2-3 \text{ GeV}$$

NEW PHYSICS AT ~1 TEV

$$\left(\text{EW. CONTRIBUTION TO THE HIGGS MASS} \right)^2 \propto G_F f(m_W, m_Z, m_t, m_H) \Lambda^2 \lesssim m_H^2$$

(the constant is $3/4\pi^2\sqrt{2}$, $f = 4m_t^2 - m_Z^2 - m_W^2 - m_H^2$)

THIS REQUIRES $\Lambda \lesssim 0.5 \text{ TeV}$

Does this 'cutoff' ALSO CORRESPOND TO THE EXISTENCE OF A NEW PARTICLE/NEW PHYSICS @ 1 TeV?

NEW PHYSICS AT ~1 TEV

HIGGS MASS $m_H \sim G_F^{-1/2}$ SHOULD BE CLOSE TO THE
MAX. ENERGY ALLOWED BY THEORY $M_{PL} \sim G_N^{-1/2}$

$$M_{PL} = \sqrt{\frac{\hbar c}{G_N}}$$

$$\frac{G_N}{\hbar c} = 6.7 \times 10^{-39} \left(\frac{\text{GeV}}{c^2} \right)^{-2}$$

$$\frac{G_F}{(\hbar c)^3} = 1.16 \times 10^{-5} \text{GeV}^{-2}$$

$$\Sigma \quad \frac{G_N}{G_F} \frac{\hbar^2}{c^2} = 5.7 \times 10^{-34} \neq 1 !!$$

ANTHROPIC

IN THE VAST MAJORITY OF UNIVERSES

$$G_F / G_N \sim 1$$

BUT THOSE HAVEN'T THE RIGHT PROPERTIES
TO DEVELOP OBSERVERS.

WIMPS



$n(t)$ = number density of χ

Rate of decrease in the number of LEFT-OVER PARTICLES
IN A "COMOVING" VOLUME

$$\underbrace{n(t) a^3(t)}_{\#(t) \text{ cm}^{-3}} \times \underbrace{n(t) \langle v \sigma \rangle_T}_{\text{sec}^{-1}} =$$

$$= n^2(t) a^3(t) \langle v \sigma \rangle_T$$

AT EQUILIBRIUM

$$= n_{\text{EQ}}^2 a^3(t) \langle v \sigma \rangle_T$$

WIMPS

$$\frac{d}{dt} (n(t) a^3(t)) = - (n^2(t) - n_{\text{EQ}}^2) a^3(t) \langle v \sigma \rangle_T$$

$$\frac{d}{dt} (n a^3) = \left(\frac{dn}{dt} \right) a^3 + n 3 a^2 \cdot \dot{a}$$

$$\boxed{\frac{dn}{dt} = - 3Hn - (n^2 - n_{\text{EQ}}^2) \langle v \sigma \rangle_T}$$

dilution from expanding universe

$$H(t) = \frac{\dot{a}(t)}{a(t)}$$

$$H(\text{today}) = H_0 = 100 h \frac{\text{km}}{\text{Mpc}} \text{sec}^{-1} \quad \text{if } h = 0.72$$

WIMPS AND FREEZE-OUT

$$\frac{dn}{dt} = -3Hn - \underbrace{(n^2 - n_{\text{EQ}}^2)}_{\text{dilution from expanding universe}} \langle \sigma v \rangle_T$$

If $T < m_\chi$, creation term is Gibbs suppressed and at

$$n \langle \sigma v \rangle = H$$

we have

$$\frac{dn}{dt} \propto -Hn$$

$$\ln n \propto -Ht$$

$$n(t) \propto e^{-Ht}$$

FREEZE-OUT

WIMPS

$$\frac{d}{dt} (n(t) a^3(t)) = - (n^2(t) - n_{\text{EQ}}^2) a^3(t) < v \sigma_T$$

$$\downarrow t \rightarrow T$$

$$\frac{d}{dx} n(x) = B [n^2(x) - n_{\text{eq}}^2(x)]$$

$$x = T/m_\chi, \quad n = n/T^3, \quad B = \# \langle v \sigma_T \rangle m_\chi$$

$$\langle v \sigma_T \rangle \sim G_F^2 m_\chi^2 f$$

SUPPOSE DM HAS WEAK INTERACTIONS

WIMPS

$$\frac{d}{dx} \mu(x) = \beta [\mu^2(x) - \mu_{\text{eq}}^2(x)]$$

THE PRESENT ENERGY DENSITY OF LEFT-OVER PARTS. IS

$$m_\chi \mu(0) T_0^3 = \rho_\chi$$

REQUIRE

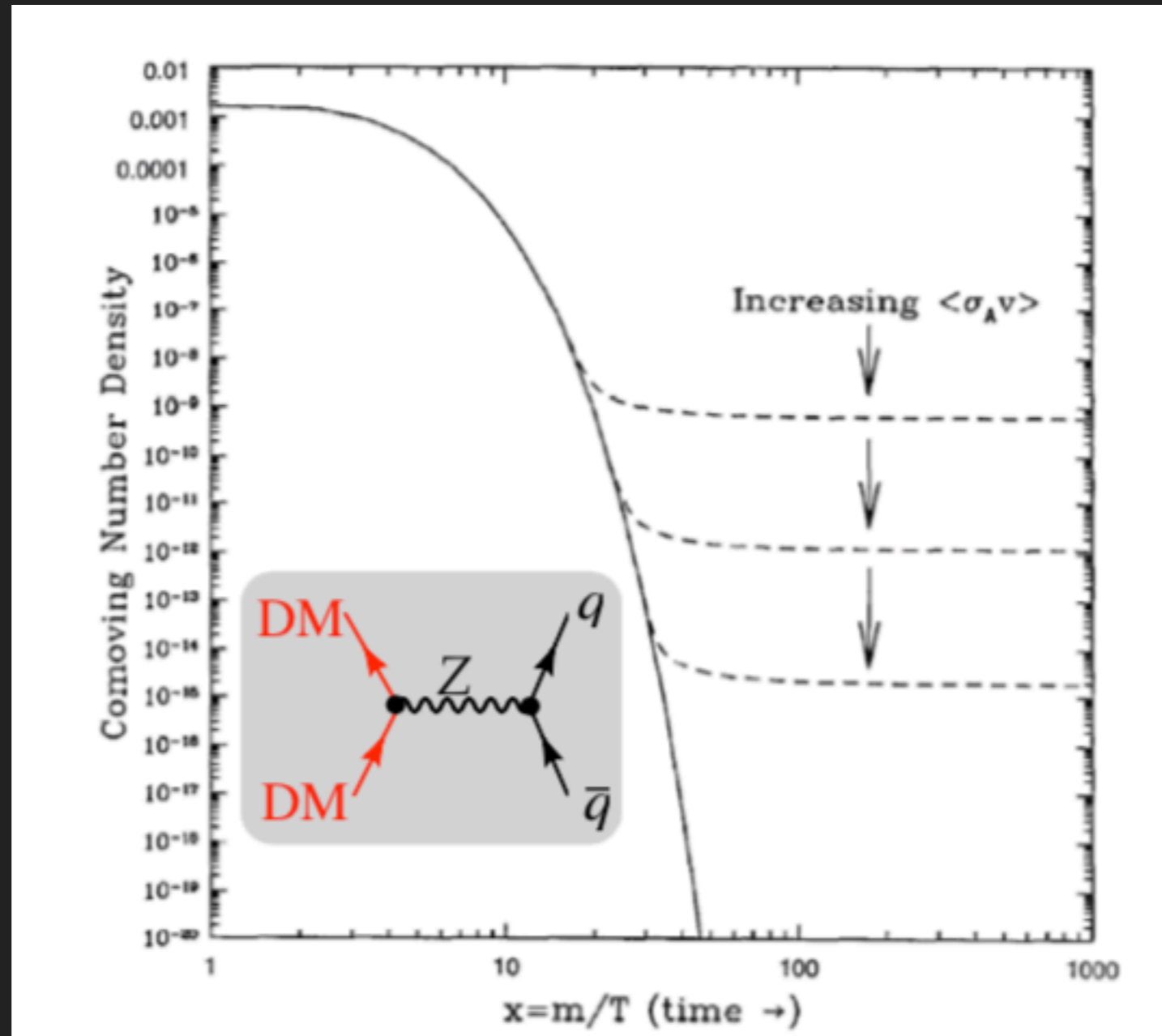
$$1) \quad \Omega_\chi = \frac{8\pi G_N \rho_\chi}{3H_0^2} \approx \Omega_M$$

$$2) \quad \text{Solve finding } \mu(0) \approx 6\beta^{-1}$$

$$m_\chi \left(\frac{f}{\sqrt{N}} \right)^{1/2} \approx 3.7 \left(\frac{\Omega_M h^2}{0.15} \right)^{-1/2} \text{ GeV}$$

WIMPS & FREEZE-OUT

Larger cross section, later freeze-out, LOWER number density (today)



Correct DM density for $\langle v\sigma \rangle \simeq 3 \times 10^{-26} \frac{\text{cm}^3}{\text{sec}} \simeq \frac{1}{(20 \text{ TeV})^2}$

Dark sector interactions & mass close to the weak scale

UNNATURAL HIGGS & WIMPS

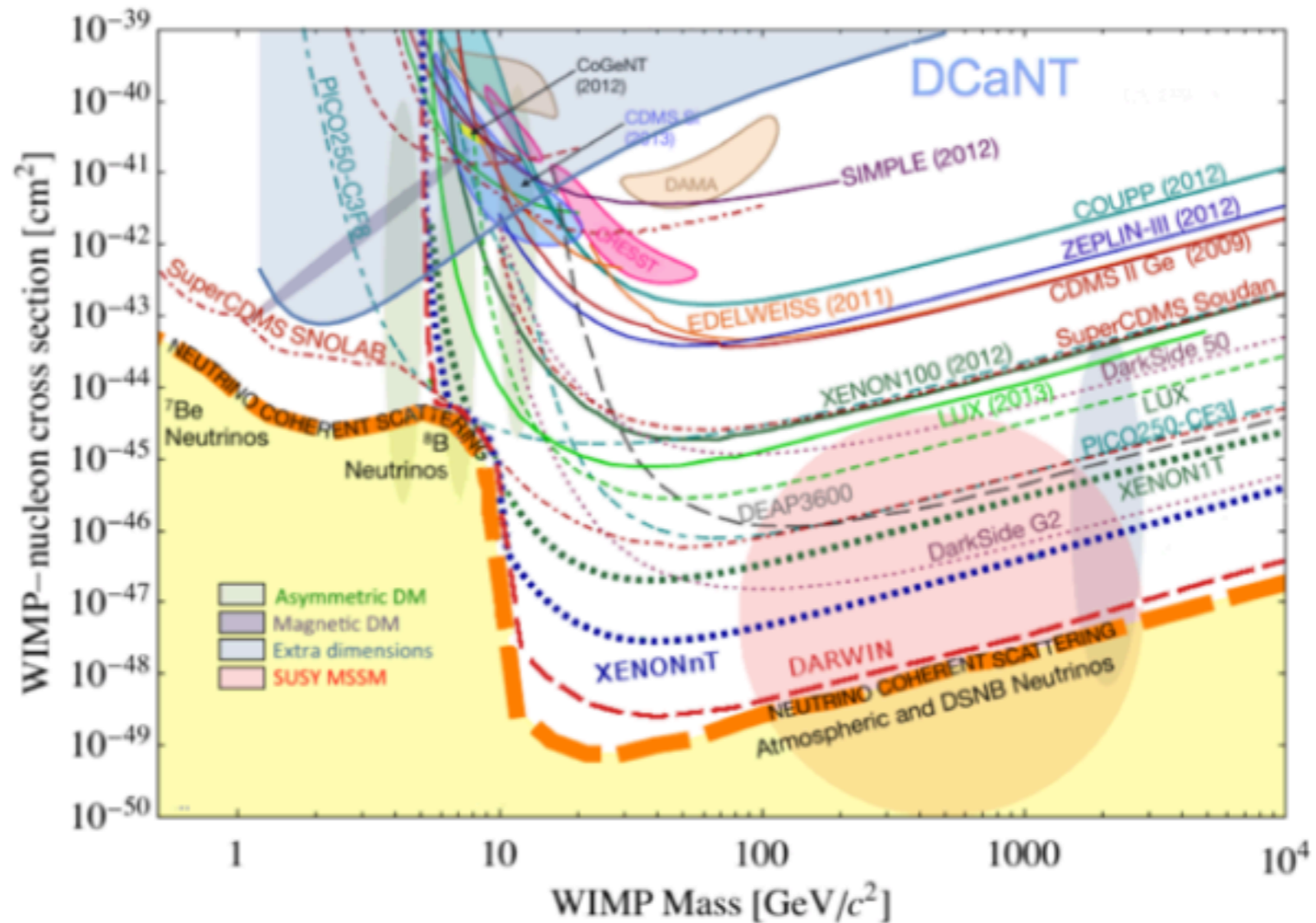
WIMPS COULD (HAVE BEEN?) BE UNDERSTOOD
WITH THE SAME PHYSICS NEEDED TO CURE THE
"UNNATURAL HIGGS"

SENSITIVITY

Directionality gives a better control on backgrounds.

Z tree

W loop



Exposure $0.4 \text{ * Kg * 1 year}$ – output ions @ 1 keV

Capparelli et al. Phys. Dark Univ. 9-10 (2015) 24, ibid. Phys. Dark Univ. 11 (2016) 79;

DARK MATTER LANDSCAPE



Superfluid

Semicond.

WIMP

He

SuperCDMS

Xenon

Sensei, DAMIC

DarkSide

LZ

Carbon nanotubes
& Graphene

← Axions
& like.

ADMX

meV

eV

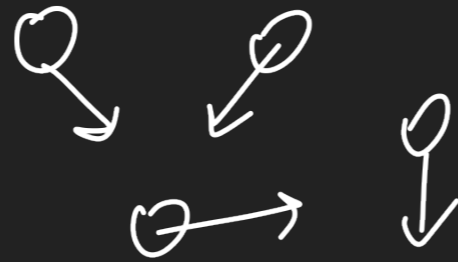
keV



↳ resolutions required by exp -

DARK MATTER LANDSCAPE

WIMPS



(Scalars)

think in terms of waves



$$\lambda = \frac{\hbar}{m_\chi v}$$

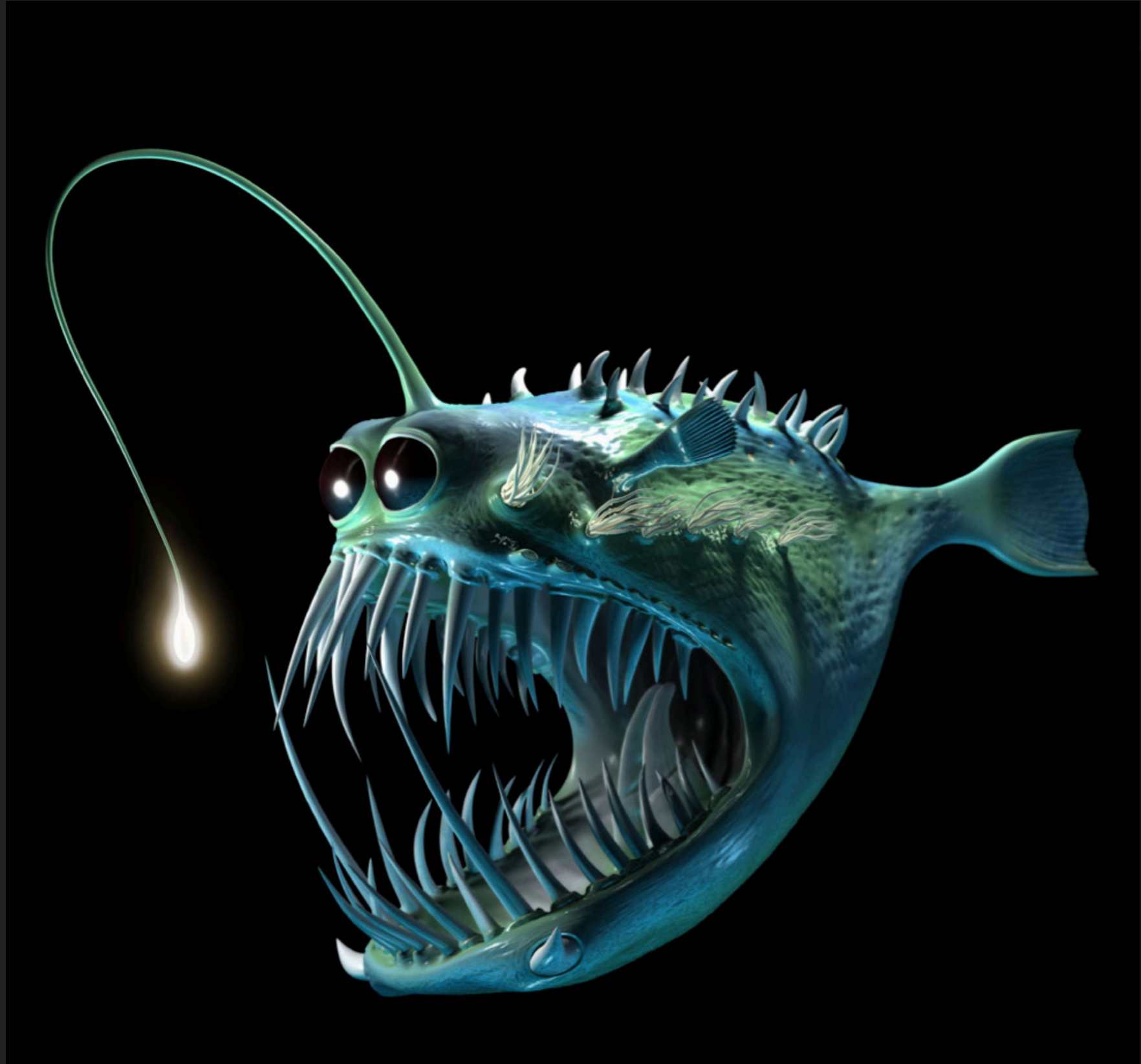
decrease m_χ



(Scalar waves)

$$\phi(t) \approx \phi_0 \cos \omega_\chi t$$

$$\omega_\chi \approx \frac{m_\chi c^2}{\hbar}; \quad \phi_0 = \sqrt{\rho_\chi} / m_\chi$$



BACKUP

$$\frac{d}{dt} (n(t) a^3(t)) = - (n^2(t) - n_{\text{EQ}}^2) a^3(t) \langle v \sigma \rangle_T$$

If $m_\chi \ll T$ the creation term is "on" and n_{EQ} is a solution of the eq. Thus

$$n_{\text{EQ}} a^3 \sim \text{const}$$

$$n_{\text{EQ}} \sim \frac{1}{a^3} \sim T^3$$

$$S(T) a^3 \sim \text{const.}$$

ADIABATIC EXPANSION

Indeed the number density in a rel. gas at thermal equilibrium is $\sim T^3$.

WIMPS

$$\frac{d}{dt} (n(t) a^3(t)) = - (n^2(t) - n_{\text{EQ}}^2) a^3(t) \langle v\sigma \rangle_T$$

WHEN T DROPS $T \ll m_\chi$, THE CREATION TERM IS "OFF" (GIBBS SUPPRESSED) AND WE HAVE

$$\frac{d}{dt} n a^3 = - n^2(t) a^3(t) \langle v\sigma \rangle_T$$

$$n(t) a^3(t) = \frac{n(t_1) a^3(t_1)}{1 + n(t_1) a^3(t_1) \int_{t_1}^t \frac{\langle v\sigma \rangle_T}{a^3(t')} dt'}$$

If the integral converges then are LEFT OVER PARTS.