# Merging of fixed points in Chern-Simons-QCD<sub>3</sub>

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Pisa, 18-7-2019

based on:

arXiv:1812.01544 & arXiv:1902.05767, w/ Hrachya Khachatryan arXiv:19??.??, w/ Guillermo Arias-Tamargo and Diego Rodriguez-Gomez

#### Relativistic QFT's in 2 + 1 dimensions

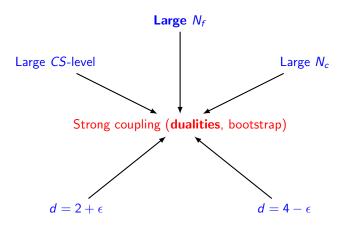
#### Interesting, for instance, for:

- theoretical laboratory for QFT's in general
- ▶ domain walls / boundary conditions of 4d QFT's (QCD)
- quantum phase transitions in condensed-matter
- Quantum Hall Effects
- ▶ graphene: multilayers, magic angles, ....
- emergent minimal supersymmetry

Lots of recent progress (last 3 years) in understanding dynamics, dualities.

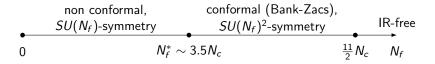
#### Dynamics of gauge theories in 2+1 dimensions

Five "simple" or weak-coupling limits:



#### Question: is there a CFT in the IR? Let's look at QCD<sub>4</sub>

In 4d QCD,  $SU(N_c)$  with  $N_f$  Dirac flavors, standard scenario:



Chiral symmetry breaking, triggered by a quartic fermion operator crossing marginality

$$\Delta[\Psi^4]=4$$
 in 4d: large anomalous dimension, strong coupling

$$\mathcal{L}_{int} = -\Psi^4 \quad \Rightarrow \quad \langle \Psi^2 
angle 
eq 0$$

 $\Psi^4$  is  $SU(N_f)^2$ -invariant but  $\langle \Psi^2 
angle$  breaks  $SU(N_f)^2 o SU(N_f)$ 

TWO REAL CFT'S "ANNIHILATE" INTO TWO COMPLEX CFT'S

#### A controlled example: bosonic QED in $4 - \epsilon$ dimensions

$$\mathcal{L} = \sum_{i=1}^{N_f} |D_{\mu}\phi_i|^2 + rac{1}{\mathrm{e}^2} F_{\mu
u}^2 - \lambda (\sum_{i=1}^{N_f} |\phi_i|^2)^2$$

One-loop beta functions:

$$\beta_{e^2} = -\epsilon e^2 + \frac{N_f}{6} e^4$$

$$\beta_{\lambda} = -\epsilon \lambda + \frac{N_f + 8}{6} \lambda^2 - 6e^2 \lambda + 18e^4$$

Two charged fixed points ( $\mathbb{CP}^{N_f-1}$ -model and tricritical QED):

$$(e^2)^* = \frac{6\epsilon}{N_f}$$
  $\lambda^* = \frac{3N_f + 54 \pm 3\sqrt{N_f^2 - 180N_f - 540}}{2N_f(N_f + 4)}\epsilon$ 

If  $N_f < 6 \left(15 + 4\sqrt{15}\right) \sim 182.95$  solutions are complex

Two loop completely changes qualitative behavior: fate in d=3 not clear Let us try to tackle this question using large  $N_f$  in d=3 instead!

## Large $N_f$ expansion: let's try $\mathcal{N}=1$ sQED

large  $N_f$  SB, Khachatryan 2018

Test a "mirror-like" duality for  $\mathcal{N} = 1$  sQED with 2 flavors

Large  $N_f$  works with  $\sim 10\%$  accuracy for  $2_f + 2_b$  complex flavors! (similar accuracy in O(N=8)-model or O(N=8)-Gross-Neveu model)

 $4 - \epsilon$  expansion SB. Benini 2018

#### Bosonic QED's at large $N_f$ in d=3

$$\begin{split} \mathcal{L}_{\textit{bQED}} &= \sum_{i=1}^{N_f} |(\partial_{\mu} + \textit{ia}_{\mu})\phi_i|^2 \\ \mathcal{L}_{\mathbb{CP}^{N_f-1} \text{model}} &= \sum_{i=1}^{N_f} |(\partial_{\mu} + \textit{ia}_{\mu})\phi_i|^2 + \sigma \sum_{i=1}^{N_f} |\phi_i|^2 \end{split}$$

Geometric series of bubble diagrams ightarrow propagators for  $a_{\mu}$  and  $\sigma$ 

Scaling dimension of simple operators (probably  $\sim$  accurate for  $N_f \ge 4$ ):

bQED (tricritical)	$\begin{split} & \Delta [\phi^* \phi_{SU(N_f)-adjoint}] = 1 - \frac{64}{3\pi^2 N_f} \\ & \Delta [ \phi ^2_{SU(N_f)-singlet}] = 1 + \frac{128}{3\pi^2 N_f} \\ & \Delta [\phi_i^* \phi_j^* \phi^k \phi^l - \text{traces}] = 2 - \frac{128}{3\pi^2 N_f} \\ & \Delta [ \phi ^4_{SU(N_f)-singlet}] = 2 + \frac{256}{3\pi^2 N_f} \end{split}$
$\mathbb{CP}^{N_f-1}$ model	$\begin{split} & \Delta [\phi^* \phi_{SU(N_f)-adjoint}] = 1 - \frac{48}{3\pi^2 N_f} \\ & \Delta [\phi_i^* \phi_j^* \phi^k \phi^l - \text{traces}] = 2 - \frac{48}{3\pi^2 N_f} \\ & \Delta [\sigma] = 2 - \frac{144}{3\pi^2 N_f} \\ & \Delta [\frac{-5 \mp \sqrt{37}}{12} \sigma^2 + F^{\mu\nu} F_{\mu\nu}] = 4 - \frac{32(4 \pm \sqrt{37})}{3\pi^2 N_f} \end{split}$

## Estimates of merging of $\mathbb{CP}^{F-1}$ and tricritical bQED: $N_{\mathfrak{E}}^* \sim 10$

SB, Khachatryan 2018

Three somewhat different ways give similar results:

 $|\phi|_{SU(N_f)-singlet}^4$  in tricritical QED crosses marginality:

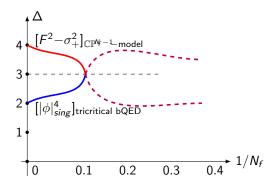
$$\Delta[|\phi|_{SU(N_f)-singlet}^4]_{bQED} = 3 \quad \rightarrow \quad N_f^* \sim \frac{256}{3\pi^2} \simeq 8.6$$

 $\sigma^2$  in  $\mathbb{CP}^{N_f-1}$  model crosses marginality:

$$\Delta [\sigma^2 + F^{\mu\nu} F_{\mu\nu}]_{\mathbb{CP}} = 3 \quad o \quad N_f^* \sim \frac{32(4 + \sqrt{37})}{3\pi^2} \simeq 10.9$$

 $|\phi|_{SU(N_f)-singlet}^4$  in tricritical QED meets  $\sigma$  in  $bC\mathbb{P}^{N_f-1}$  model:

$$\Delta[|\phi|^2_{singlet}]_{bQED} = 1 + \frac{128}{3\pi^2 N_f} = \Delta[\sigma]_{\mathbb{CP}} = 2 - \frac{144}{3\pi^2 N_f} \rightarrow N_f^* \sim 9.2$$



In particular for  $N_f = 2$  no  $2^{nd}$  order phase transition:

- ▶ important for celebrated Deconfined Quantum Critical point, (VBS ↔ Neel order transition)
- consistent with numerical simulations

Svistunov 2007

consistent with rigorous numerical bootstrap bounds

#### Chern-Simons-QCD<sub>3</sub>: $SU(N_c)_k$ with $N_f$ fermions

Regime  $k \geq \frac{N_f}{2}$ : Bosonization duality

$$SU(N_c)_k \text{ w}/N_f \psi' s \iff U(k + \frac{N_f}{2})_{-N_c} \text{ w}/N_f \phi' s$$
Aharony 2015

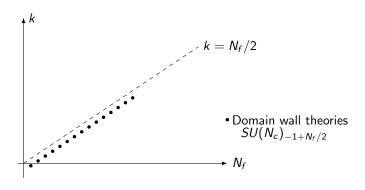
- massive phases match because of level-rank duality for TQFT's
- ▶ large- $N_c$  & large-k limit: exact checks  $\rightarrow$  CFT Aharony,Minwalla,...
- $ightharpoonup N_f = 1$ : holographic duality with Vasiliev higher-spin theory on  $AdS_4$

Regime  $k < \frac{N_f}{2}$ : possible region with "quantum-phase" and symmetry breaking

$$U(N_f) \rightarrow U(N_f/2 + k) \times U(N_f/2 - k)$$

Komargoski, Seiberg 2017 Armoni, Dumitrescu, Festuccia, Komargoski 2018 Argurio, Bertolini, Mignosa, Niro 2018

## Chern-Simons-QCD<sub>3</sub>: phase diagram



Domain walls of QCD<sub>4</sub> at  $\theta_{QCD}=\pi$ :  $SU(N_c)_{-1+N_f/2}$  w/  $N_f$   $\Psi$ 's Gaiotto, Komargoski, Seiberg 2017 Argurio, Bertolini, Bigazzi, Cotrone, Niro 2017

What can we say from large  $N_f$  expansion? Compute scaling dimensions of quartic operators and look for  $\Delta[\Psi^4]=3$ 

#### $CS-QCD_3$ at large $N_f$ : bilinear and quartic operators

Arias-Tamargo, SB, Rodriguez-Gomez

The bilinears,  $\bar{\Psi}_i \Psi^j$ , transform in the adjoint plus singlet of  $SU(N_f)$ :

$$\Delta[\Psi\Psi_{adjoint}] = 2 - \frac{64(N_c^2 - 1)}{3\pi^2 \left(1 + \left(\frac{8k}{\pi N_f}\right)^2\right) N_c N_f} + O(1/N_f^2)$$

$$\Delta[\Psi\Psi_{singlet}] = 2 - \frac{128(N_c^2 - 1) \left(2\left(\frac{8k}{\pi N_f}\right)^2 - 1\right)}{3\pi^2 \left(1 + \left(\frac{8k}{\pi N_f}\right)^2\right)^2 N_c N_f} + O(1/N_f^2)$$

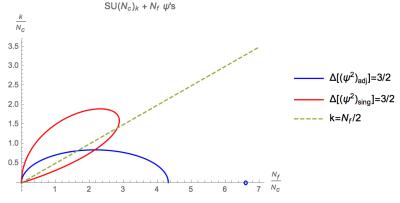
In principle we need to study mixing of four  $\Psi^4_{singlet}$  operators (WIP)

In the regime  $N_c>>1$ , only two double-trace quartic ops are into play:

$$\Delta[(\Psi\Psi_{adjoint})^2] \sim 2 \times \Delta[\Psi\Psi_{adjoint}] \sim 4 - \frac{c_1(k/N_f)N_c}{N_f} + O(N_c^2/N_f^2)$$

$$\Delta[(\Psi\Psi_{\textit{singlet}})^2] \sim 2 \times \Delta[\Psi\Psi_{\textit{singlet}}] \sim 4 - \frac{c_2(k/N_f)N_c}{N_f} + \textit{O}(N_c^2/N_f^2)$$

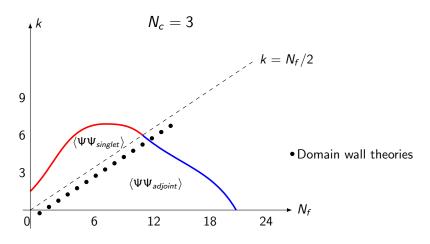
## Marginality crossing of the two double trace $\Psi^4_{sinolet}$



Inside blue curve: 
$$\Delta[(\Psi\Psi_{adjoint})^2] < 3 \Rightarrow \langle \Psi\Psi_{adjoint} \rangle \neq 0$$
  
Inside red curve:  $\Delta[(\Psi\Psi_{singlet})^2] < 3 \Rightarrow \langle \Psi\Psi_{singlet} \rangle \neq 0$ 

Blue dot: comes from  $\Delta[\Psi\Psi_{adjoint}]$  at  $O(1/N_f^2)$ , known for k=0 (important: confirms that merging region is non-empty)

## Qualitative picture and relation to domains walls



Conjecture: symmetry-breaking on the wall  $\leftrightarrow \chi\text{-symmetry-breaking}$  in the 3+1d bulk

#### **SUMMARY:**

non susy 2 + 1d gauge theories at strong coupling (small  $N_f$  & small k): often **no**  $2^{nd}$  **order phase transition** 

Large  $N_f$  techniques useful in d=3: minimal susy, QED and CS-QCD

Dualities for 1<sup>st</sup> order phase transitions less powerful (but complex CFT's)

#### **FUTURE DIRECTIONS:**

More general 4d gauge theories (quiver, rank-2 matter): cf. Poppitz talk

domains walls

3d dualities and IR scenarios (e.g. quantum phases)

4d IR scenarios (many different proposals, domain walls may help)

4d non susy dualities?

#### Numerical conformal bootstrap:

CFT's with intermediate number of flavors in 3d and 4d?