



Università degli Studi di Padova

Constraining Dark Energy parameters with needlets

Bianca De Caro PUMA22 - 27/09/2022

Outline

- Integrated Sachs Wolfe effect and cross-correlation;
- Needlets-based estimator;
- ISW detection and Dark energy parameter estimation with needlets;
- Conclusion and future prospects.

Integrated Sachs-Wolfe effect

- The ISW is due to the **evolution** of the gravitational potential Ψ after recombination;
- Gravitational potential **decays** between the era of transition (from matter to dark energy dominated universe), giving rise to temperature **anisotropy** in the CMB;
- ISW happens at $z \leq 1$ and at extremely large scale $\ell \leq 30$;
- It is related to the **Dark Energy** parameters.

- At low ℓ , CMB signal is dominated by **cosmic variance**;
- ISW anisotropy is suppressed with respect to primary anisotropy (before recombination);
- The most direct way to detect this effect it to **cross-correlate** CMB photons with the Large Scale Structure (LSS) at low ℓ ;
- The angular power spectrum in the harmonics space C_{ℓ} is the most common **observable**.

• Cross-correlation angular power spectrum at low ℓ is dependent on

cosmological parameter
$$\Omega_{\Lambda}=1-$$

$$\Omega_{\rm m}$$
 and $w = \frac{p}{\rho}$ of DE;

• Ω_{Λ} changes the **amplitude** of the power spectrum;



Other estimators?

- Harmonic estimator exploits orthogonal properties in harmonic space;
- It possess the **minimal** variance in case of no mask, spatially uniform noise;
- Building new estimators to deal with a more realistic scenario.

Needlets

- Needlet-based estimator of the cross-correlation power spectrum between CMB photons and LSS;
- Needlets are a form of **spherical wavelet**;
- Needlets are a convolution of spherical harmonics and a suitably chosen window function;
- Needlets are localized both on real space and on harmonic space;
- Needlets are also used for testing the presence of systematics.







Needlets – CMBxG power spectrum

- Validation of the needlet estimator:
 - Calculation of the auto- and cross- angular power spectra for CMB temperature and density of galaxies from CAMB;
 - Simulation of Healpix maps for the correlated T and G from the C'_es, in the case of full-sky and zero noise;
 - Extraction of the needlets **coefficients** β_j for the cross-correlation power spectrum from maps;
 - Computation of the **covariance** matrix from simulations;

Needlets - ISW

- From the analysis of these coefficients, it is possible to predict the detectability of ISW;
- As test of significance, we compute the **chi** square between β_j of a single realization and the mean of all the realizations;
- We found $\chi^2 = 15.6$, so for 12 d.o.f. we can exclude the null hypothesis (no ISW) with **79%** confidence.



Needlets – Dark Energy parameter estimation

- We produced a **grid** of spectra for different values of Ω_{Λ} ;
- In each point we evaluated a Gaussian **likelihood** for a single realization, with mean equal to β_i in that point and covariance from the simulations;
- We **sampled** from the normalized likelihood in order to find the **posterior** distribution for Ω_{Λ} .
- We calculated the **median** of the distribution and confidence interval at **68%** confidence level.

Needlets – Dark Energy parameter estimation

- We repeated that procedure for β_j from different realizations and for 30 and 60 values of Ω_{Λ} ;
- We compared our result with the
 Planck best-fit value taken as fiducial;
- We found that the estimator is unbiased;



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Conclusion and future prospective

- We found that it is possible to **detect** the ISW from a needlets-based estimator of the cross-correlation power spectrum between CMB temperature and galaxies at low ℓ ;
- We also demonstrate that it is possible to **constrain** Dark Energy parameters from the β_j spectrum;
- Verification with masked sky and catalogues of galaxies in order to forecast some results for the Euclid survey.
- Comparison with other estimators.

Backup

Late ISW

- The late ISW occurs at late times, $z \lesssim 1$
- It is restricted to extremely large scale $l \lesssim 30$
- It depends on the cosmological parameters related to the Dark Energy
- The most direct way to detect this effect it to cross correlate CMB photons with the Large Scale Structure (LSS) at low redshifts.
- Galaxies and clusters are tracers of the dark matter overdensity δ_m related to the evolution of the gravitational potential Ψ trough the Poisson equation:

$$\nabla^2 \Psi = 4\pi G a^2 \rho_m \delta_m(\vec{x}, a)$$

Integrated Sachs-Wolfe effect

$$\Theta_{\ell}^{ISW} = -2 \int_0^{\eta_0} d\eta \, e^{-\tau} \psi'(k,\eta) j_{\ell}[k(\eta_0 - \eta)]$$

- The ISW is due to the evolution of the gravitational potential Ψ after recombination;
- The ISW at late time is due to the decay of gravitational potential. This gives rise to temperature anisotropy in the CMB spectrum during the era of transition between a matter dominated universe to a dark energy dominated one;
- ISW happens at $z \leq 1$ and it is restricted to extremely large scale $l \leq 30$;
- It depends on the cosmological parameters related to the Dark Energy.

• The most direct way to detect this effect it to cross correlate CMB photons with the Large Scale Structure (LSS) at low. The cross-correlation power spectrum C_{ℓ}^{TG} is:

$$C_{\ell} = \frac{2}{\pi} \int_0^{\infty} dk k^2 P_M(k) \int_0^{\infty} dz \, W^{ISW}(z) \int_0^{\infty} dz' \, W^g(z')$$

- $P_M(k)$ is the power spectrum of the matter overdensity today;
- $W^g(z)$ is the distribution of galaxies/clusters in a redshift interval;
- $W^{ISW}(z)$ is the contribution from the CMB temperature anisotropies due to the ISW:

$$W^{ISW}(z) = -3 \frac{\Omega_m H_0^2}{k} \frac{1}{D_+(a)} e^{-\tau} \frac{d}{dz} \left[\frac{D_+(a)}{a(z)} \right] j_\ell(k\chi(z))$$

 It depends on the cosmological parameters of the Dark Energy.



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Research activity - Needlets

- I am currently working on a needlet-based estimator of the cross-correlation power spectrum between CMB photons and LSS;
- Needlets are a form of spherical wavelet widely study for data analysis of power spectra. Wavelets are useful to handle with rapidly changeable signals; $\psi_{jk}(\hat{\mathbf{n}}) = \sqrt{\lambda_{jk}} \sum_{\ell=[B^{j-1}]}^{[B^{j+1}]} b\left(\frac{\ell}{B^{j}}\right) \sum_{m=-\ell}^{\ell} Y_{\ell m}^{*}(\hat{\mathbf{n}}) Y_{\ell m}(\xi_{jk})$
- The parameter B defines the localization in the harmonic space, since $\ell \in [B^{j-1}, B^{j+1}]$, and in the real spaces;
- Needlets-based estimator: $\beta_{jk} = \sqrt{\lambda_{jk}} \sum_{\ell=[B^{j-1}]}^{[B^{j+1}]} b\left(\frac{\ell}{B^j}\right) \sum_{m=-\ell}^{\ell} x_{\ell m} Y_{\ell m}(\xi_{jk}) \qquad \hat{\beta}_j^{XY} = \frac{1}{N_{\text{pix}}} \sum_k \beta_{jk}^X \beta_{jk}^Y$
- Mean and variance: $\langle \hat{\beta}_{j}^{XY} \rangle \equiv \beta_{j}^{XY} = \sum_{\ell} \frac{2\ell+1}{4\pi} b^2 \left(\frac{\ell}{B^j}\right) C_{\ell}^{XY} \quad (\Delta \beta_{j}^{XY})^2 \equiv \operatorname{Var}[\hat{\beta}_{j}^{XY}] = \sum_{\ell} \frac{2\ell+1}{16\pi^2} b^4 \left(\frac{\ell}{B^j}\right) \left[(C_{\ell}^{XY})^2 + C_{\ell}^{XX} C_{\ell}^{YY} \right]$
- Covariance: $\operatorname{Cov}_{jj'} \equiv \operatorname{Cov}[\hat{\beta}_j, \hat{\beta}_{j'}] = \langle (\hat{\beta}_j \langle \hat{\beta}_j \rangle_{\mathrm{MC}}) (\hat{\beta}_{j'} \langle \hat{\beta}_{j'} \rangle_{\mathrm{MC}}) \rangle_{\mathrm{MC}}$

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Research activity – Needlets (real space)

B = 1.74 , j = 8





Research activity – Needlets (real space)

B = 1.5 , j = 8





Needlets – Dark Energy parameter estimation

- We repeated that procedure for β_j from different realizations and for 30 and 60 values of Ω_{Λ} ;
- We compared our result with the Planck best-fit value taken as fiducial;
- We found that the estimator is **unbiased**;
- We found that the accuracy on the constrains on the parameter improves as the points of the grid are increasing.

