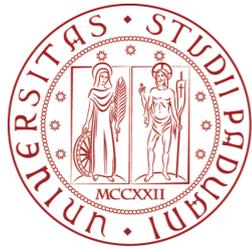


# Sub-GeV Dark Matter and the pre-BBN Universe



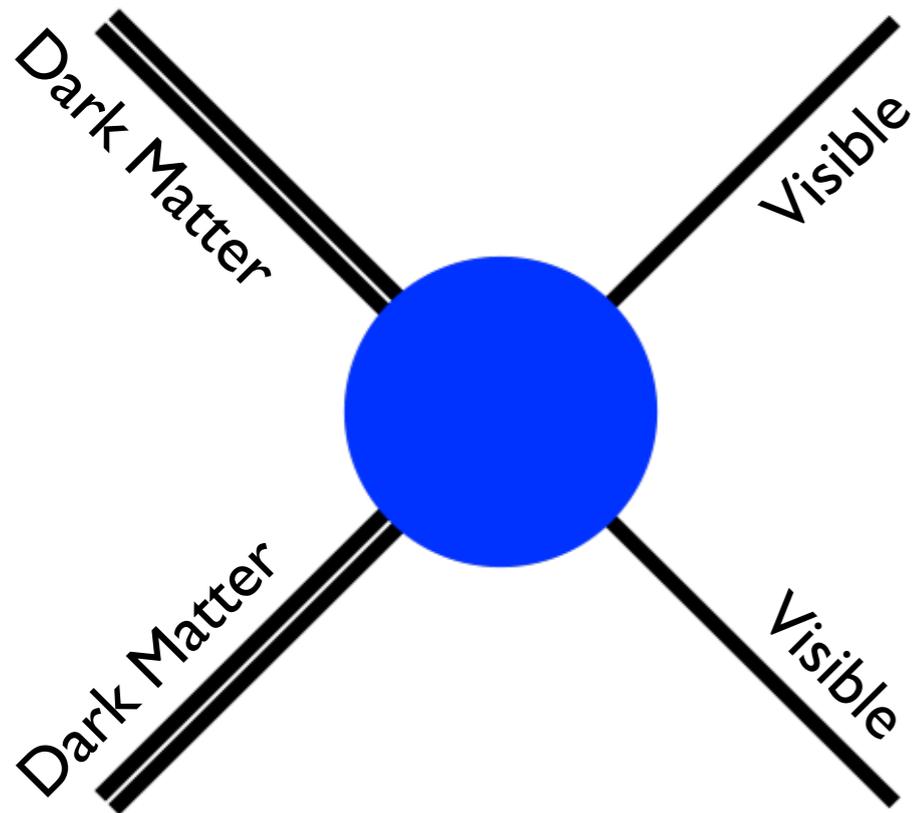
UNIVERSITÀ  
DEGLI STUDI  
DI PADOVA



Istituto Nazionale  
di Fisica Nucleare  
Sezione di Padova

**Francesco D'Eramo**

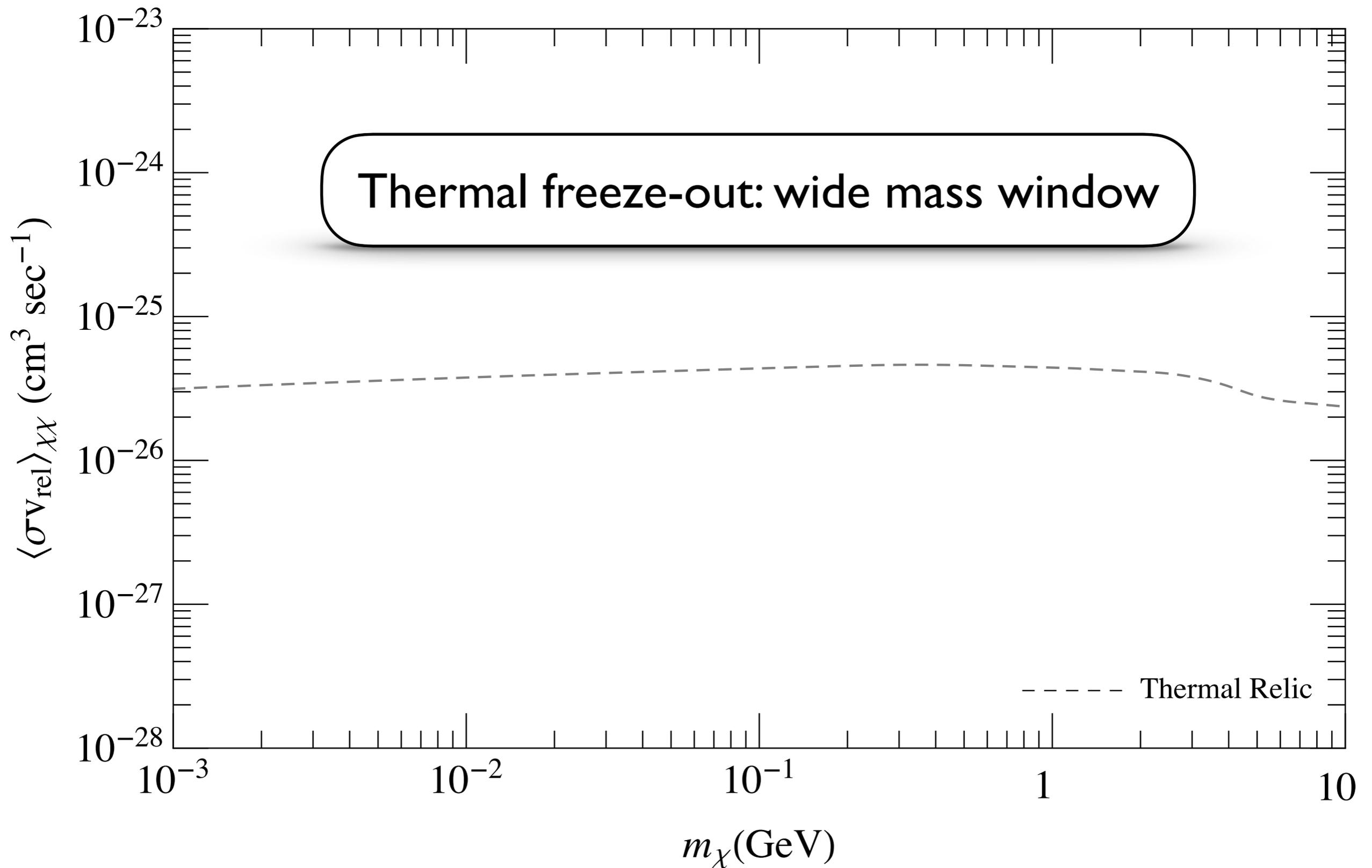
# What mass for the dark matter?



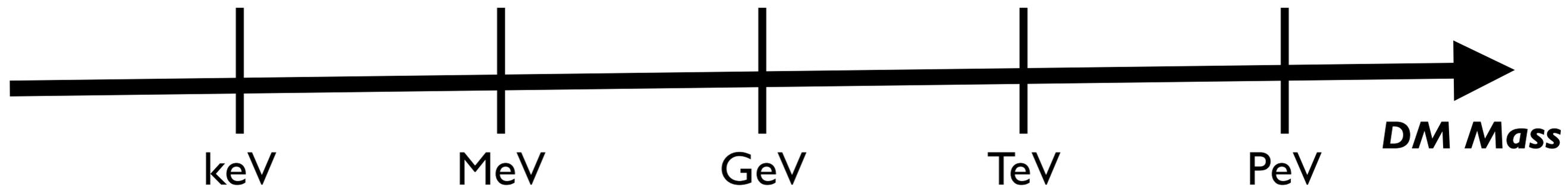
Relic density bounds  
annihilation rate at freeze-out

$$\langle \sigma v_{\text{rel}} \rangle \simeq 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$

# What mass for the dark matter?



# Bounds on the dark matter mass

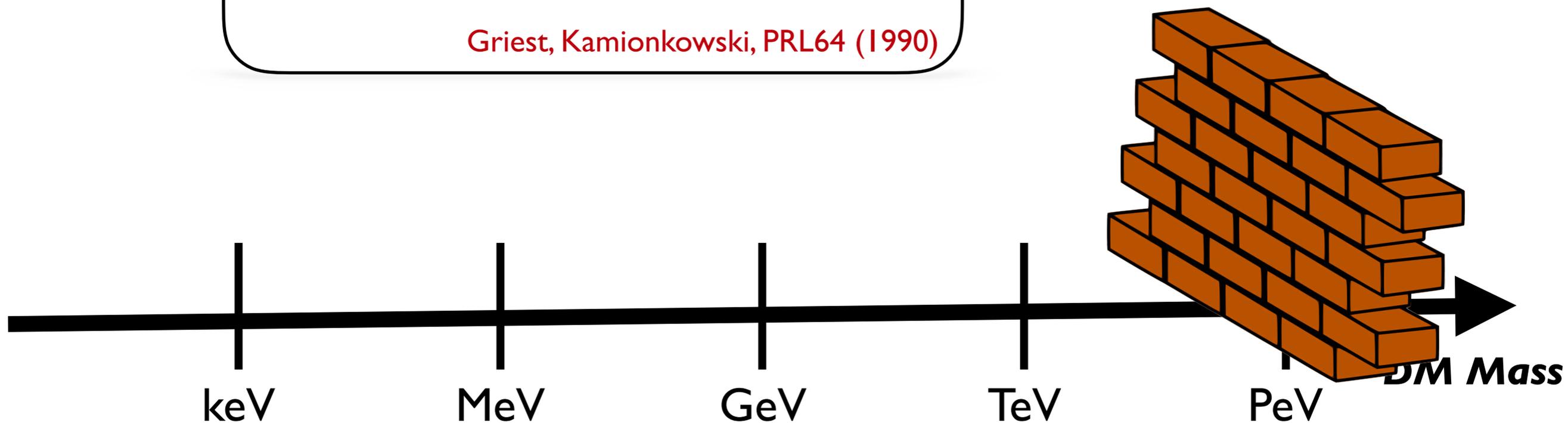


# Bounds on the dark matter mass

## Unitarity

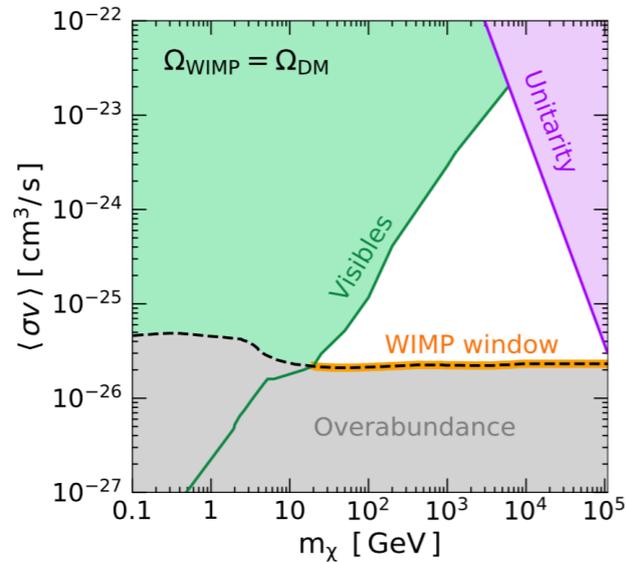
$$\sigma v \lesssim \frac{4\pi}{m_{\text{DM}}^2 v}$$

Griest, Kamionkowski, PRL64 (1990)



# Bounds on the dark matter mass

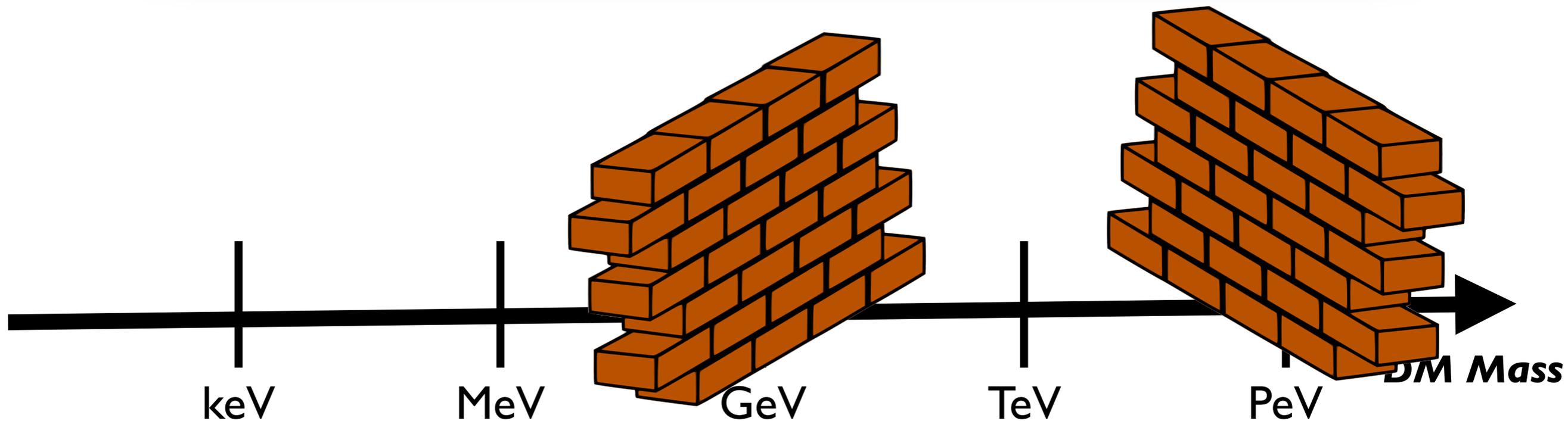
## The WIMP Window



Dark matter with s-wave annihilation to visible states

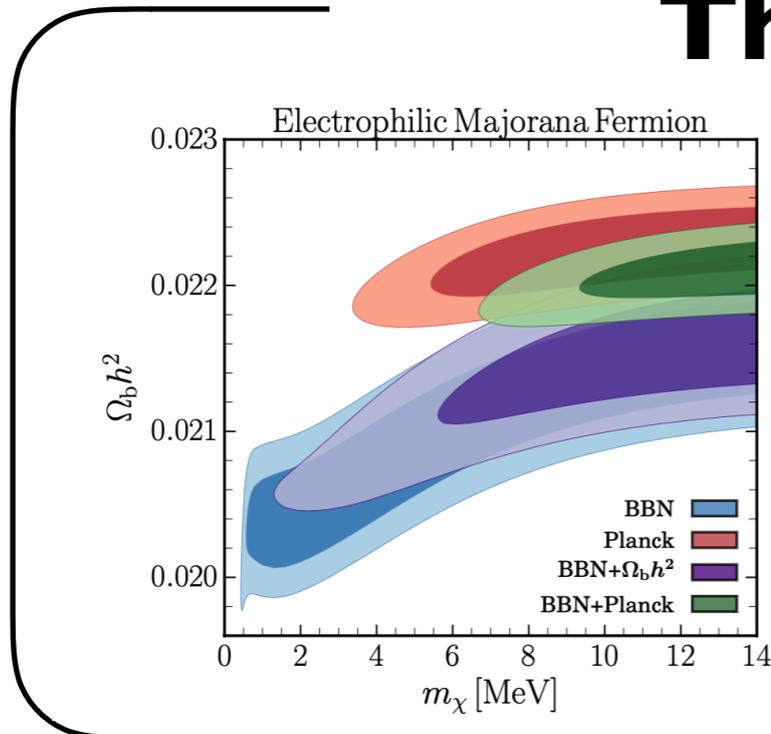
$$10 \text{ GeV} \lesssim m_{\text{WIMP}} \lesssim 100 \text{ TeV}$$

Leane, Slatyer, Beacom, Ng, PR D98 (2018)



# Bounds on the dark matter mass

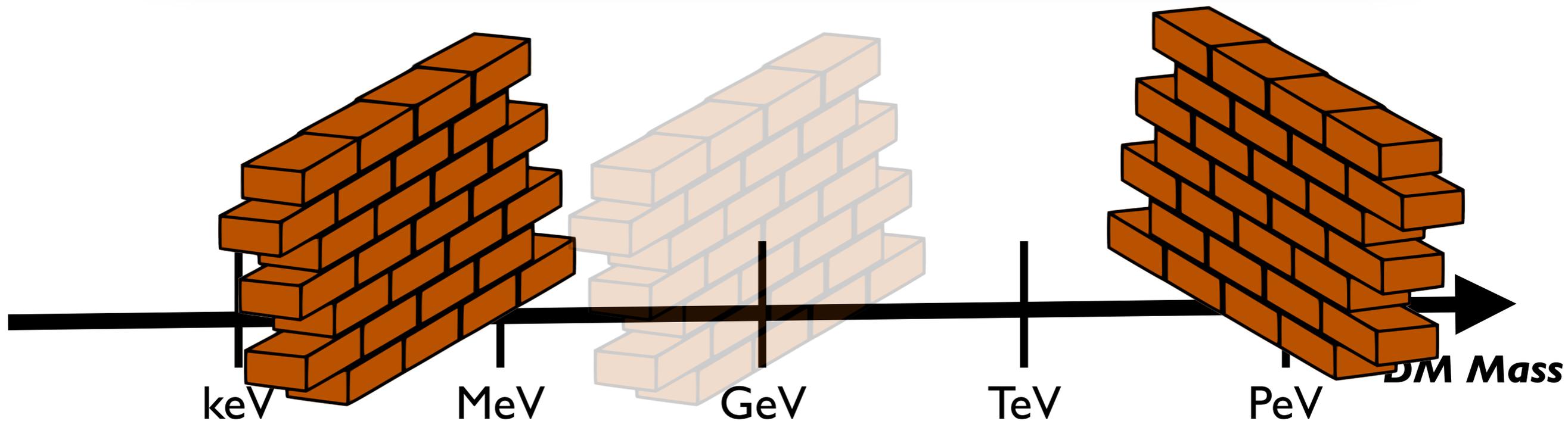
## The BBN Wall



Thermal relics with mass below MeV spoil BBN and CMB

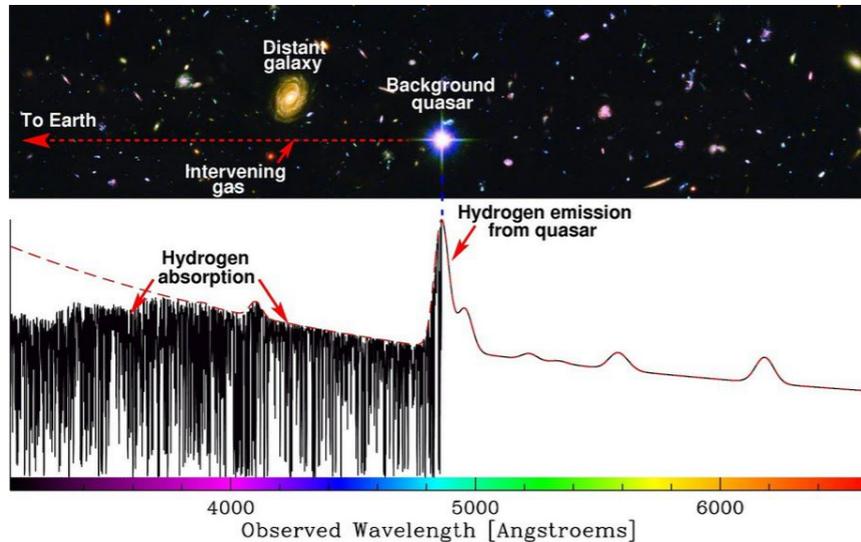
$$\text{MeV} \lesssim m_{\text{thermal}} \lesssim 100 \text{ TeV}$$

Sabti, Alvey, Escudero, Fairbairn, Blas, arXiv:1910.01649



# Bounds on the dark matter mass

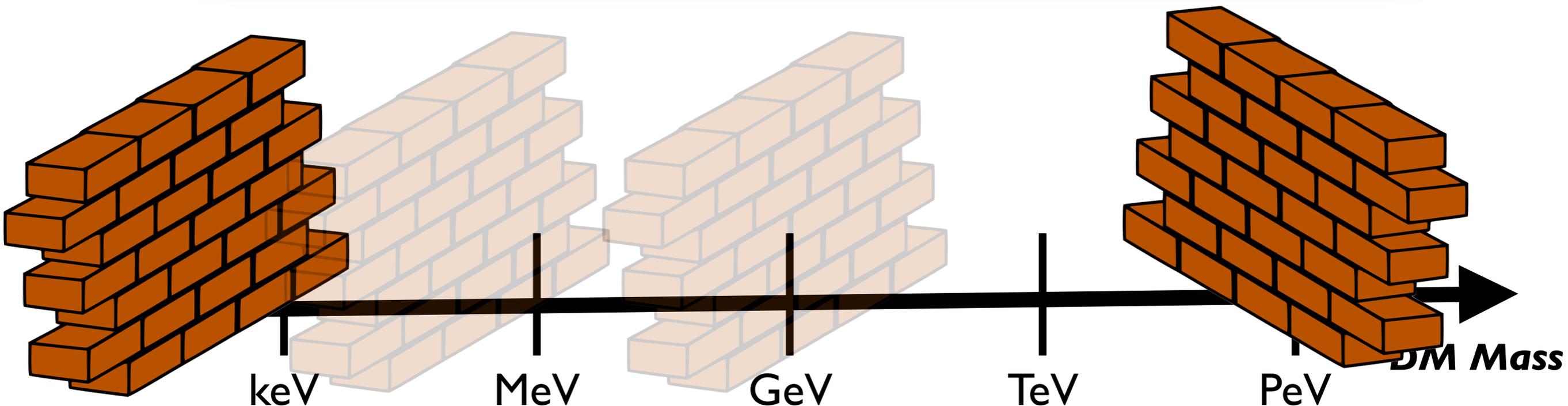
## Cannot be “too hot”



Bounds on the dark matter  
free-streaming length

$$m_{\text{warm}} \gtrsim \text{keV}$$

Viel, Becker, Bolton, Haehnelt, PRD88 (2013)

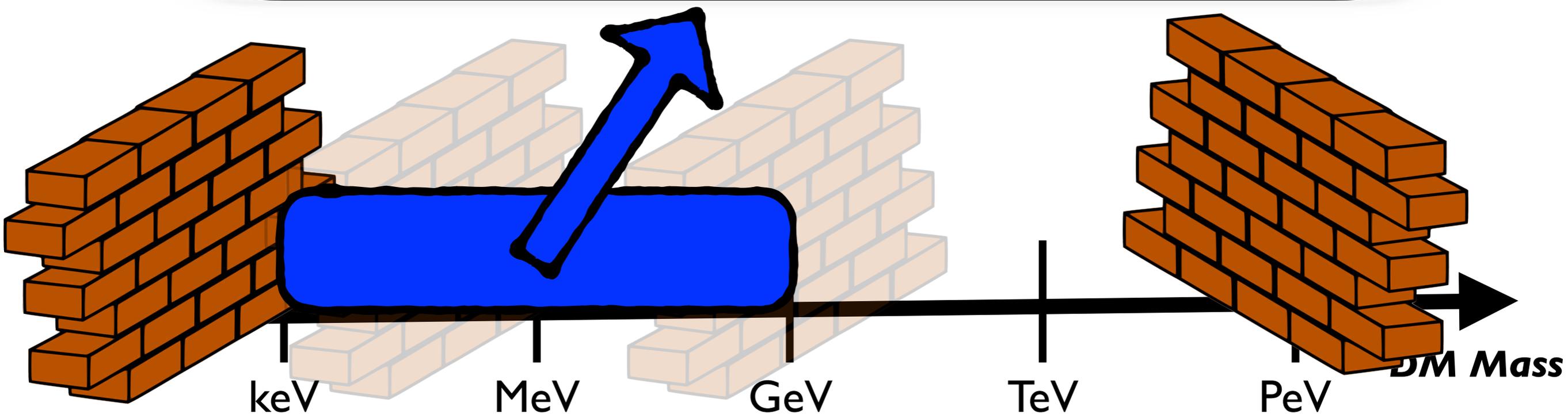


# The sub-GeV mass window

## Is it accessible?

Can we successfully produce sub-GeV dark matter in the early universe?

How do we test it today?



# The sub-GeV mass window

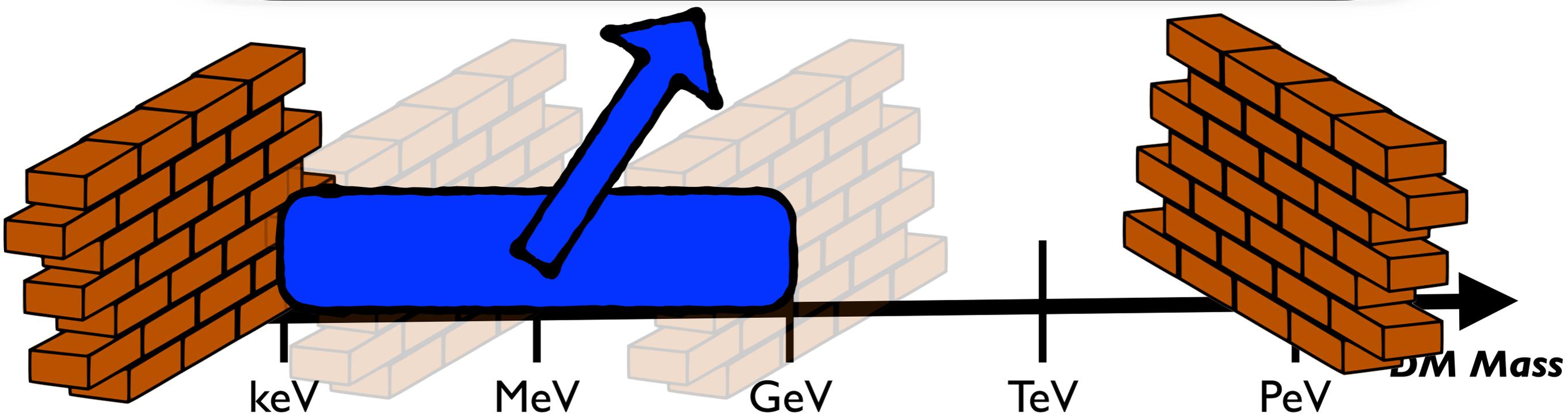
## Is it accessible?

Can we successfully produce sub-GeV dark matter in the early universe?

**YES!**

How do we test it today?

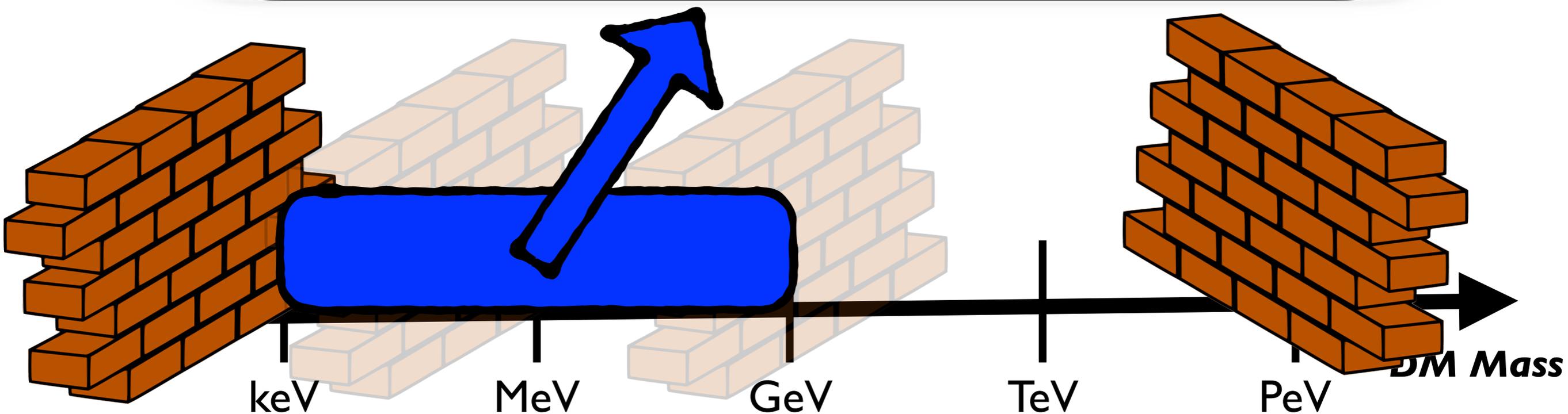
**DEPENDS!**



# The sub-GeV mass window

## Requirements

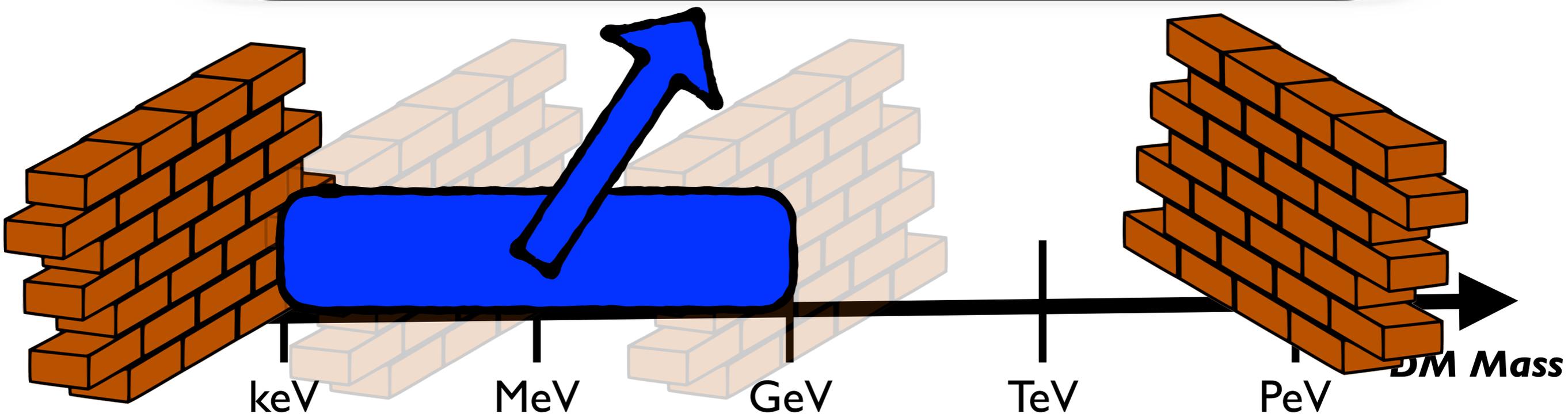
- It cannot annihilate via s-wave (GeV bound)
- It cannot be in thermal equilibrium at BBN (MeV bound)
- It has to talk to the visible sector (detectability)



# The sub-GeV mass window

## In this talk

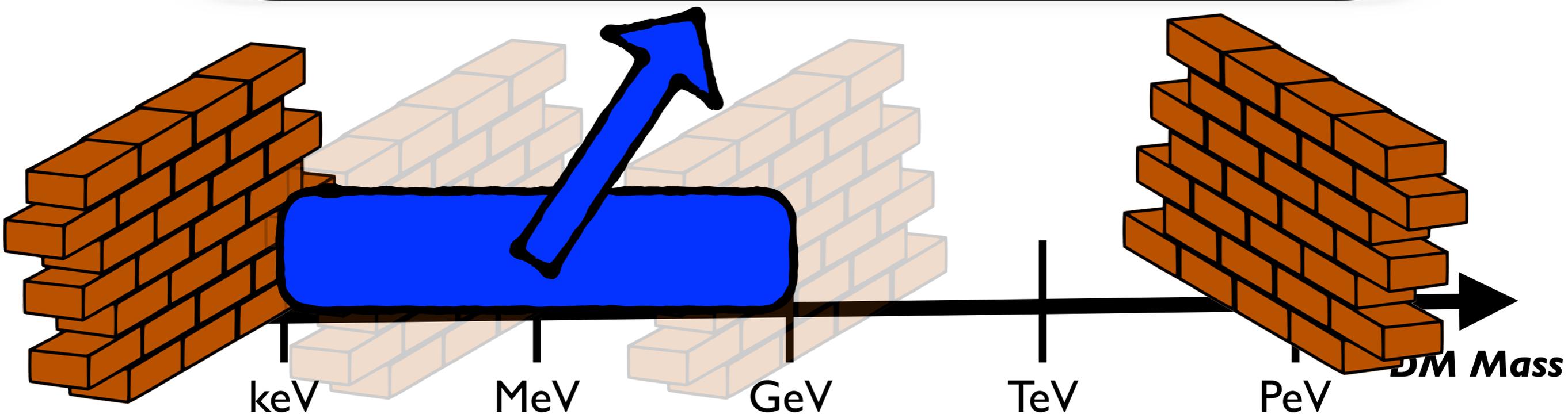
- Framework for sub-GeV non-thermal dark matter
- How to test it today
- What we can learn about our universe



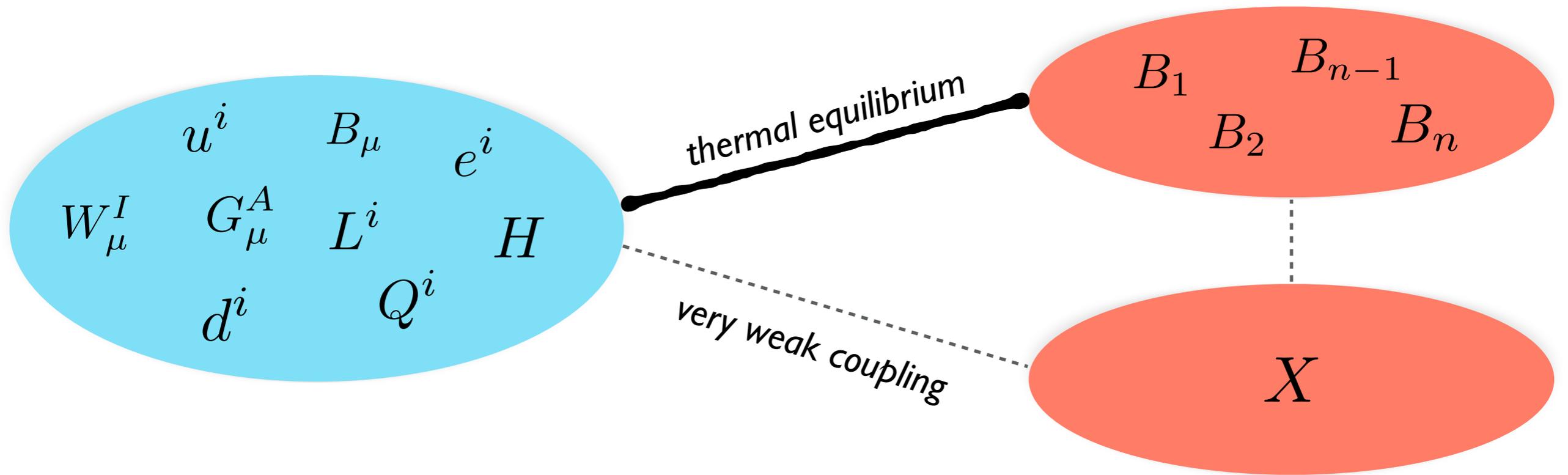
# The sub-GeV mass window



Other scenarios discussed by Berlin, D'Agnolo, Darmé, Kuflik in this workshop



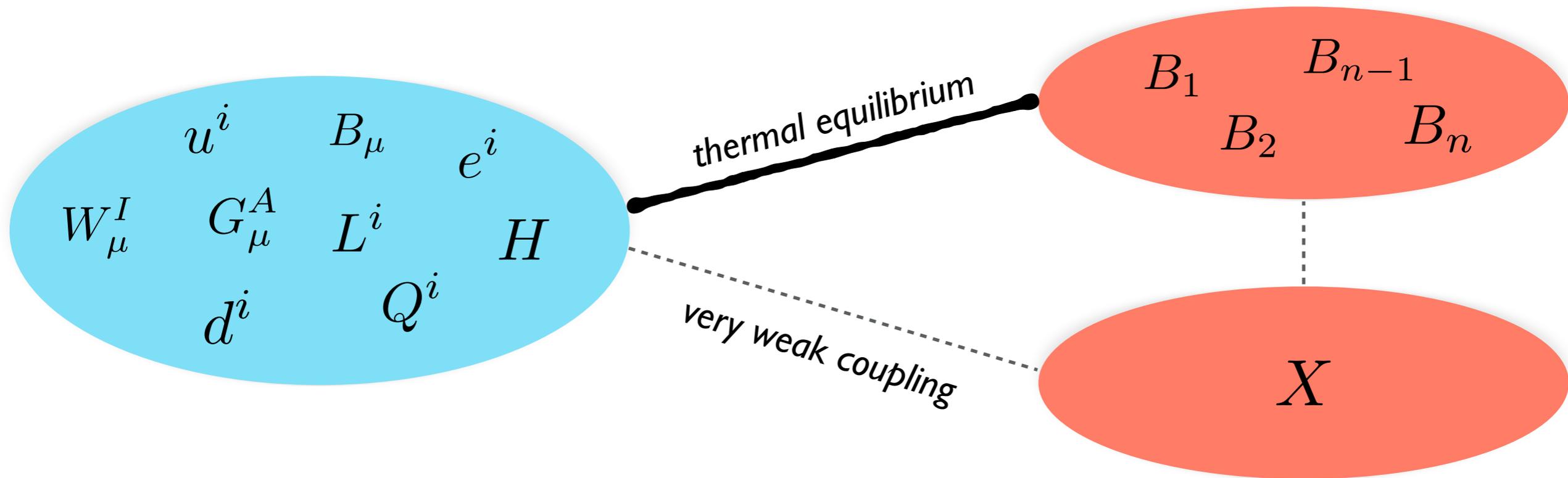
# FIMPs



$B_i$  and  $X$  odd under a  $Z_2$  symmetry  
For this talk:  $X$  is the lightest  $Z_2$ -odd particle

Hall, Jedamzik, March-Russell, West,  
JHEP 1003 (2010)

# FIMPs



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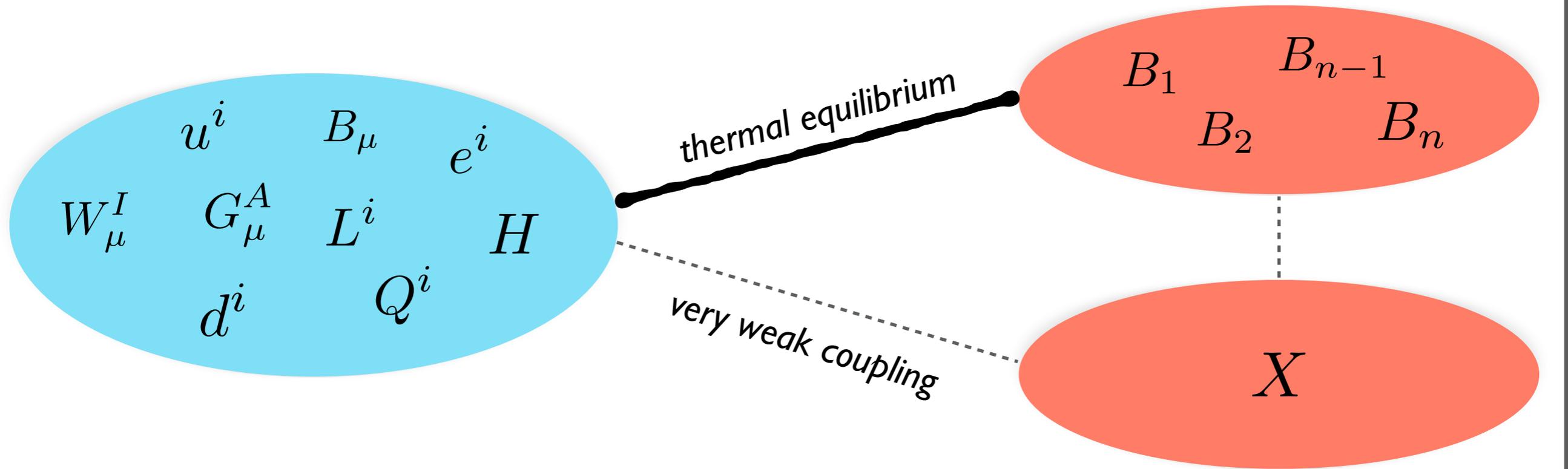
Hall, Jedamzik, March-Russell, West,  
JHEP 1003 (2010)



## Nightmare for dark matter searches

Testable by direct detection in some exceptions  
(see, e.g., Hambye et al., PRD98 (2018))

# FIMPs



$B_i$  and  $X$  odd under a  $Z_2$  symmetry  
For this talk:  $X$  is the lightest  $Z_2$ -odd particle

Hall, Jedamzik, March-Russell, West,  
JHEP 1003 (2010)

Production mechanism for  
FIMPs in the early universe?

# Freeze-in

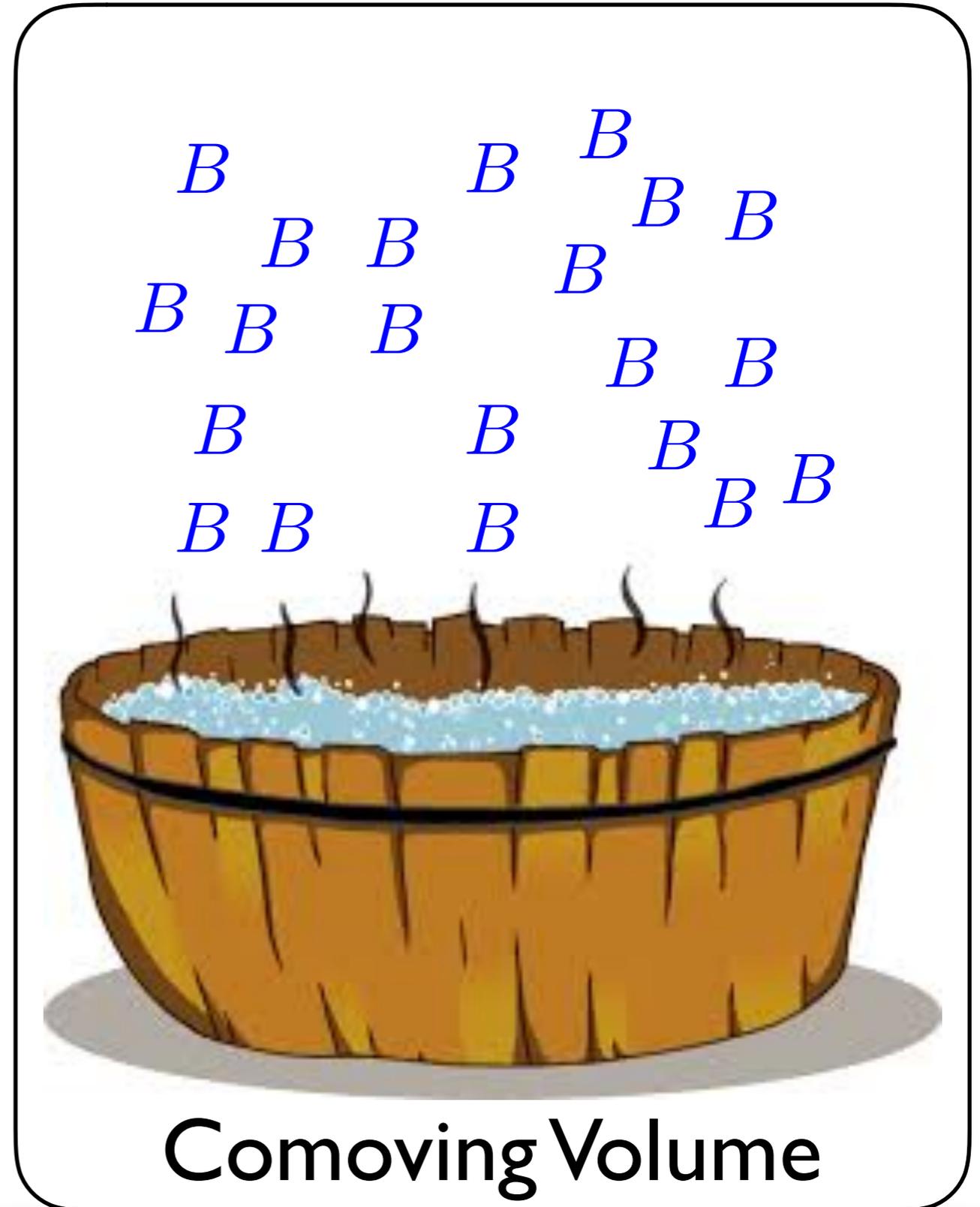
Bath particles collisions and/or  
decays dump  $X$  out of equilibrium

# Freeze-in

Bath particles collisions and/or decays dump  $X$  out of equilibrium

$T \gg$  bath particles mass:

after inflation, universe reheated without any dark matter particle in the bath  
(this is an assumption)

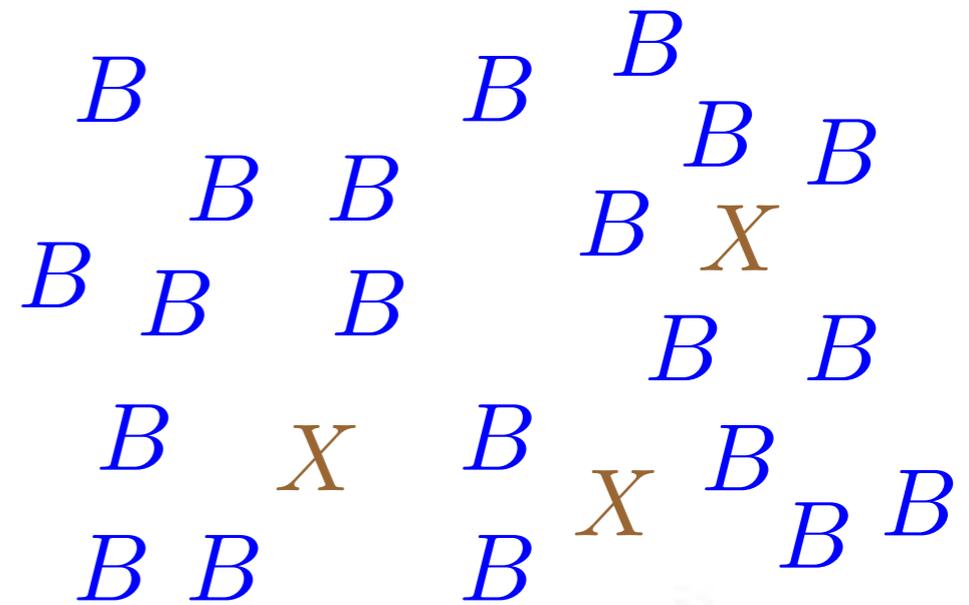


# Freeze-in

Bath particles collisions and/or decays dump  $X$  out of equilibrium

$T \sim$  freeze-in epoch:

dark matter particles dumped in the primordial plasma and then free-stream until the present time



Comoving Volume

# Freeze-in via decays

$$B \rightarrow \text{SM} + X$$

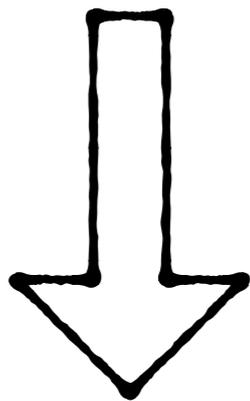
Bath particles decays produce  $X$  particle that will never thermalize

# Freeze-in via decays

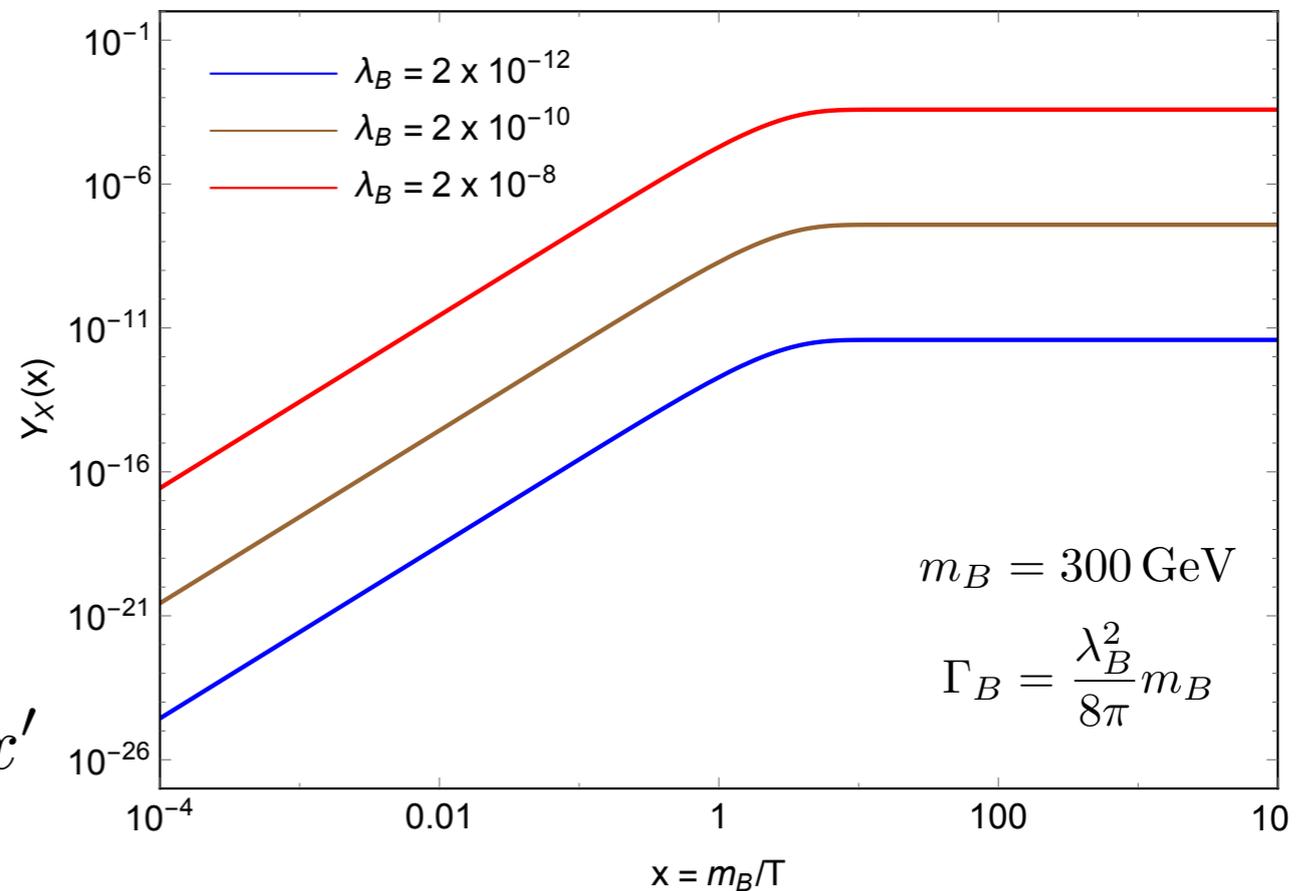


Bath particles decays produce  $X$  particle that will never thermalize

$$\frac{dn_X}{dt} = -3Hn_X + \Gamma_B n_B^{\text{eq}} \frac{K_1[m_B/T]}{K_2[m_B/T]}$$



$$Y_X(x) = g_B \frac{45}{4\pi^4} \Gamma_B \int_0^x \frac{x' K_1[x']}{g_{*s}(x') H(x')} dx'$$



# Freeze-in via decays

$$B \rightarrow \text{SM} + X$$

Bath particles decays produce  $X$  particle that will never thermalize

## **A collider signal?**

B's pair-produced at colliders, subsequent decays could be observed

# Freeze-in via decays

$$B \rightarrow \text{SM} + X$$

Bath particles decays produce  $X$  particle that will never thermalize

## A collider signal?

$B$ 's pair-produced at colliders, subsequent decays could be observed

Parameters “constrained” by relic density

$$\tau_B = \Gamma_B^{-1} \simeq 3.7 \times 10^8 \text{ cm} \left( \frac{m_X}{100 \text{ GeV}} \right) \left( \frac{300 \text{ GeV}}{m_B} \right)^2$$

# Freeze-in via decays

$B \rightarrow \text{SM} + X$       Bath particles decays produce  $X$   
particle that will never thermalize

## A collider signal?

**Hopeless?**  $B$  seems stable on the scale of experiments

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# Freeze-in via decays

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## A collider signal?

**Hopeless?**  $B$  seems stable on the scale of experiments

**CAVEAT:** result valid for a standard cosmological history

$$\tau_B = \Gamma_B^{-1} \simeq 3.7 \times 10^8 \text{ cm} \left( \frac{m_X}{100 \text{ GeV}} \right) \left( \frac{300 \text{ GeV}}{m_B} \right)^2$$

# Freeze-in via decays

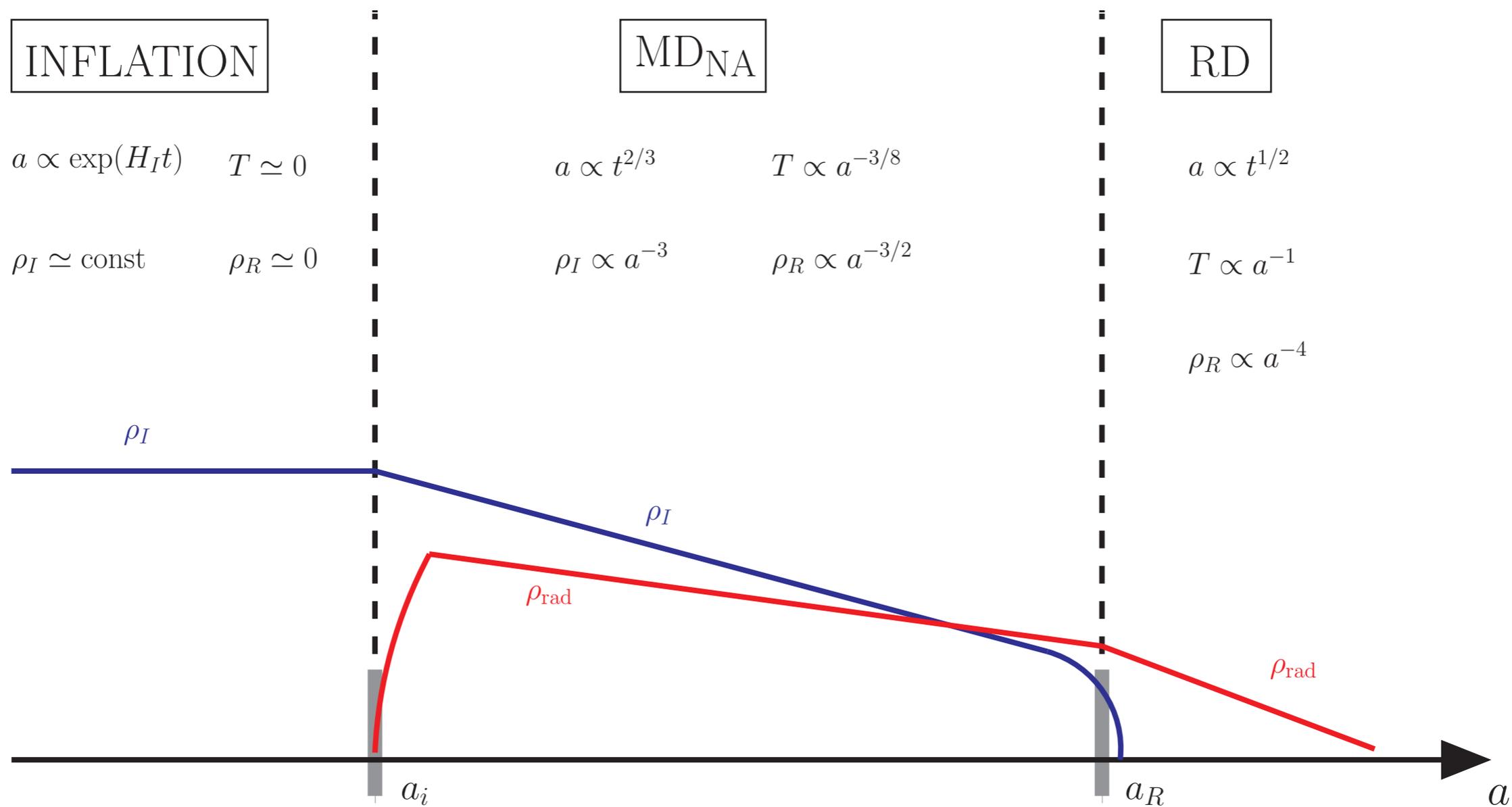
$B \rightarrow \text{SM} + X$       Both particles decays produce  $X$   
particle that will never thermalize

## A collider signal?

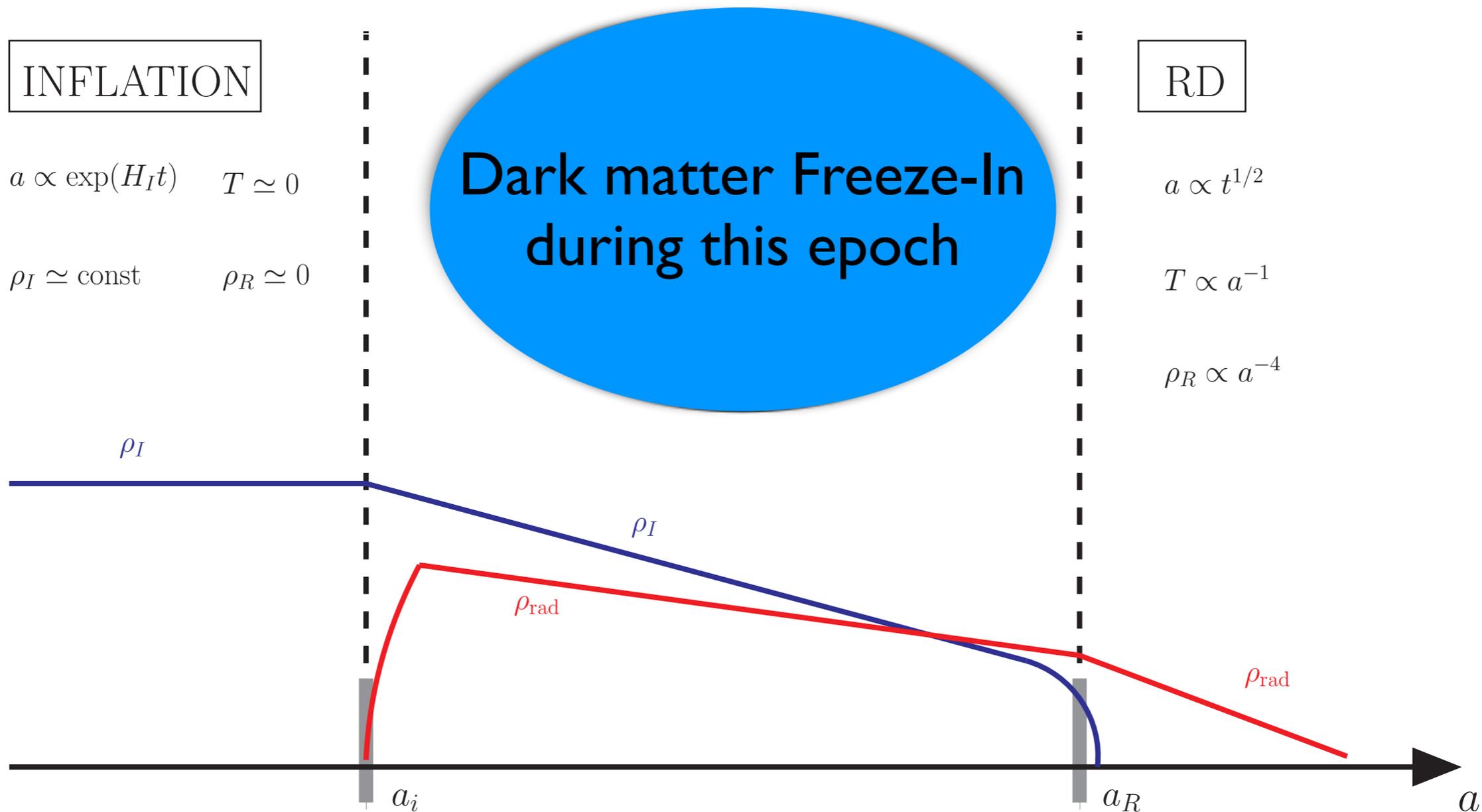
***Displaced events at collider could give information  
about the dark matter mass  
and the cosmological history of our universe***

$$\tau_B = \Gamma_B^{-1} \simeq 3.7 \times 10^8 \text{ cm} \left( \frac{m_X}{100 \text{ GeV}} \right) \left( \frac{300 \text{ GeV}}{m_B} \right)^2$$

# Case 1: FI and Early MD



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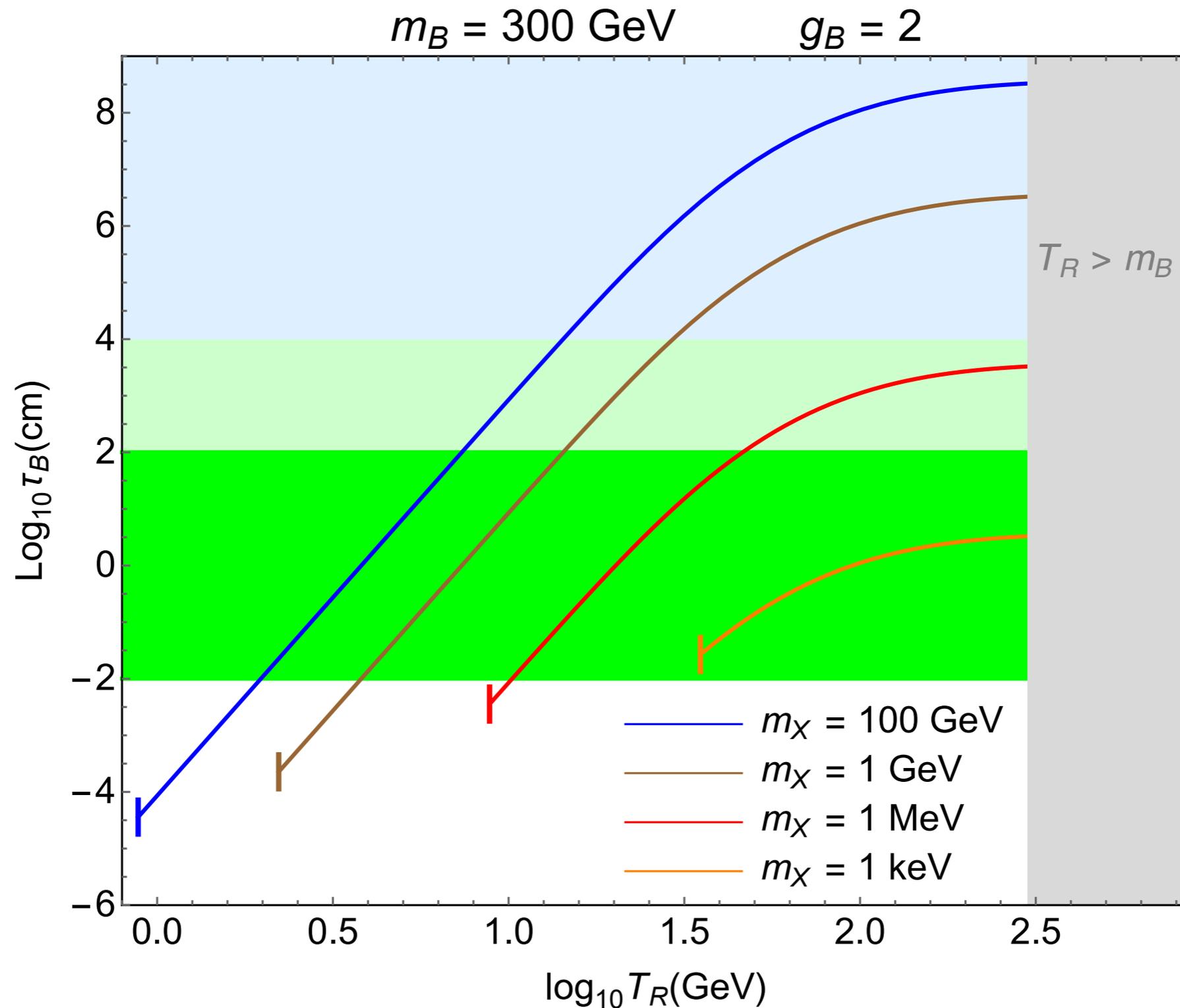


Decaying particle in the bath

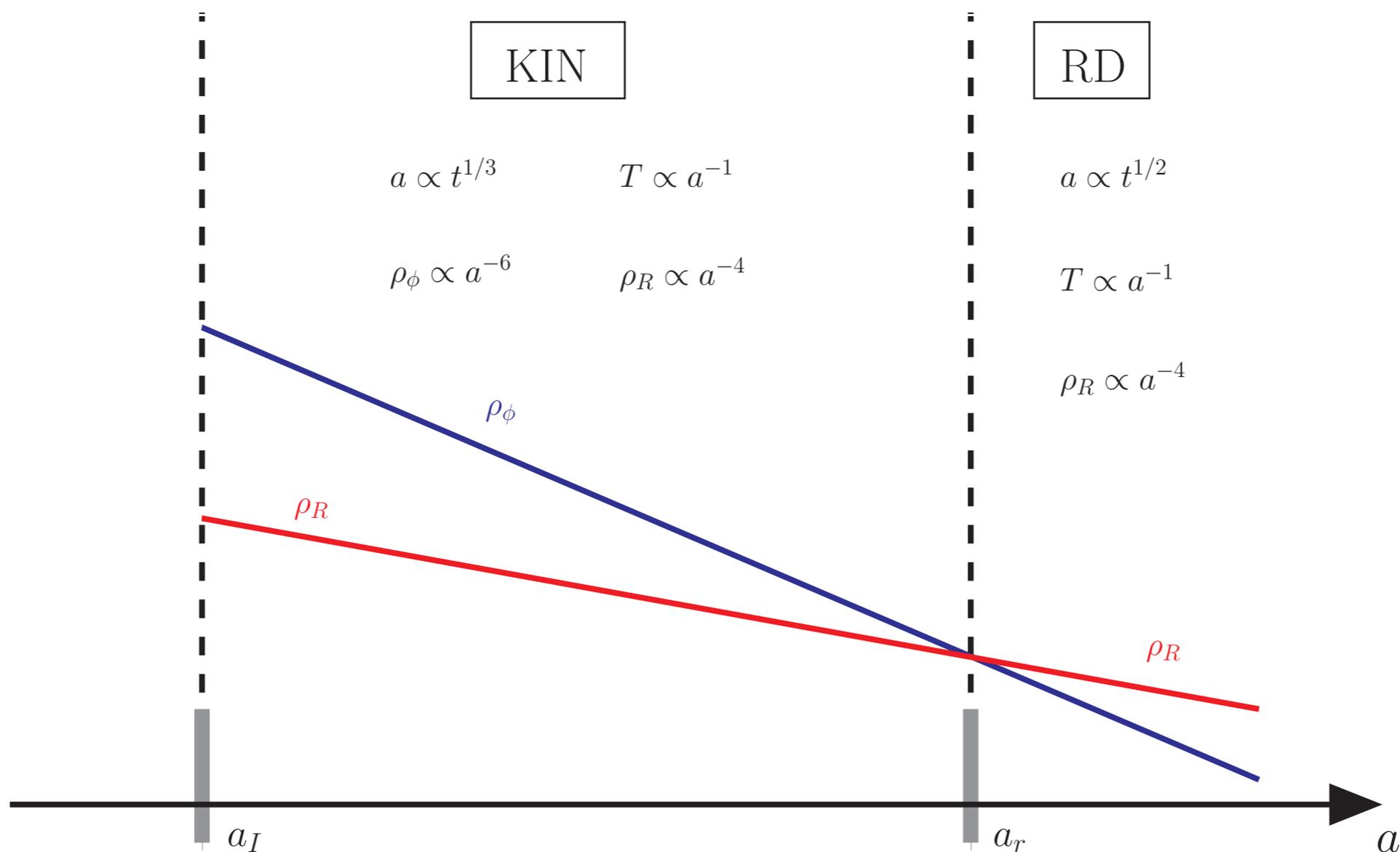
$$T_{\text{max}} \simeq (E_I T_R)^{1/2} \gtrsim m_B$$

Freeze-In density depends on the reheat temperature

# Case 1: FI and Early MD



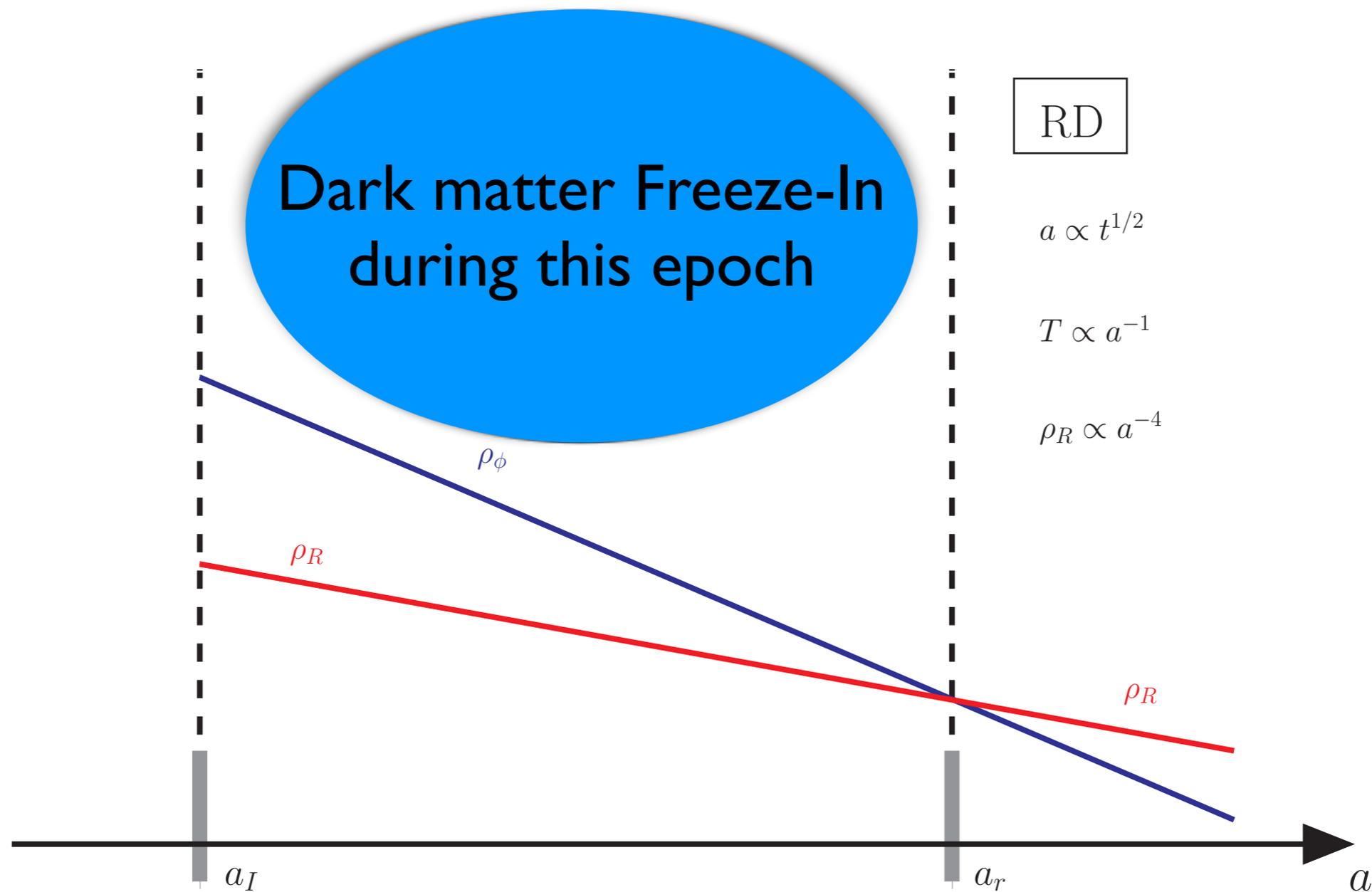
# Case 2: FI and Kination



$$\frac{p}{\rho} = w = \frac{\frac{\dot{\phi}^2}{2} - V(\phi)}{\frac{\dot{\phi}^2}{2} + V(\phi)}$$

Kination phase:  $w = 1$   
(Kinetic  $\gg$  Potential)

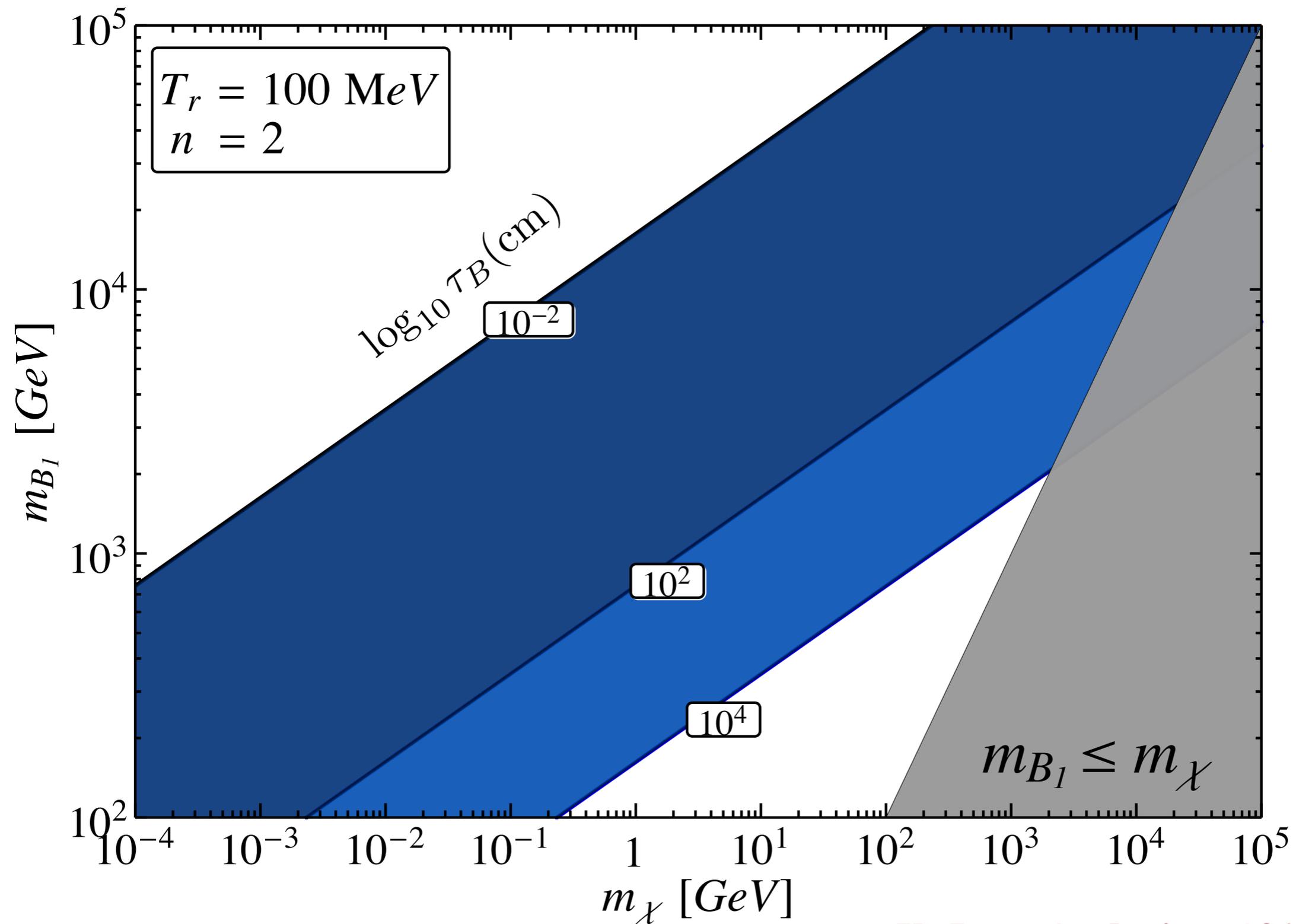
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# Case 2: FI and Kination



# A bottom-up approach

Classification of all possible operators mediating the decay

$$B \rightarrow \text{SM} + X$$

# A bottom-up approach

Classification of all possible operators mediating the decay

$$B \rightarrow \text{SM} + X$$

For each choice of the SM particle, we know the quantum numbers of B (X must be a gauge singlet)

We consider:  
 spin 0 and 1/2 DM  
 spin 0, 1/2 and 1 for B

Class for $A_{\text{SM}}$	Field for $A_{\text{SM}}$	Gauge Charges
$\psi_{\text{SM}}$	$Q_L^i$	$(\mathbf{3}, \mathbf{2})_{+1/6}$
	$u_R^i$	$(\mathbf{3}, \mathbf{1})_{+2/3}$
	$d_R^i$	$(\mathbf{3}, \mathbf{1})_{-1/3}$
	$E_L^i$	$(\mathbf{1}, \mathbf{2})_{-1/2}$
	$\bar{d}_R^i$	$(\mathbf{1}, \mathbf{1})_{-1}$
$F_{\mu\nu}$	$G_{\mu\nu}^A$	$(\mathbf{8}, \mathbf{1})_0$
	$W_{\mu\nu}^I$	$(\mathbf{1}, \mathbf{3})_0$
	$B_{\mu\nu}$	$(\mathbf{1}, \mathbf{1})_0$
$H$	$H$	$(\mathbf{1}, \mathbf{2})_{+1/2}$

# A bottom-up approach

Classification of all possible operators mediating the decay

$$B \rightarrow \text{SM} + X$$

$A_{\text{SM}}$	Spin $X$	Spin $B$	Interaction	Label
$\psi_{\text{SM}}$	0	1/2	$\overline{\psi_{\text{SM}}}\Psi_B\phi$	$\mathcal{F}_{\psi_{\text{SM}}\phi}$
	1/2	0	$\overline{\psi_{\text{SM}}}\chi\Phi_B$	$\mathcal{S}_{\psi_{\text{SM}}\chi}$
		1	$\overline{\psi_{\text{SM}}}\Gamma^\mu\chi V_B^\mu$	$\mathcal{V}_{\psi_{\text{SM}}\chi}$
$F_{\mu\nu}$	0	1	$V_B^{\mu\nu}F_{\mu\nu}\phi$	$\mathcal{V}_{F\phi}$
	1/2	1/2	$\overline{\psi_{\text{SM}}}\sigma_{\mu\nu}\chi F^{\mu\nu}$	$\mathcal{F}_{F\chi}$
$H$	0	0	$\Phi_B^\dagger H\phi$	$\mathcal{S}_{H\phi}$
		1	$V_B^\mu(c_\phi H\partial_\mu\phi + c_H\phi D_\mu H)$	$\mathcal{V}_{H\phi}$
	1/2	1/2	$\overline{\Psi_B}\chi H$	$\mathcal{F}_{H\chi}$

# A fermion DM example

Fermion dark matter with a scalar partner coupled to leptons

$$\bar{l} \chi \Phi_B$$

Also a Higgs portal operator allowed by all symmetries

$$\lambda_H H^\dagger H \Phi_B^\dagger \Phi_B$$

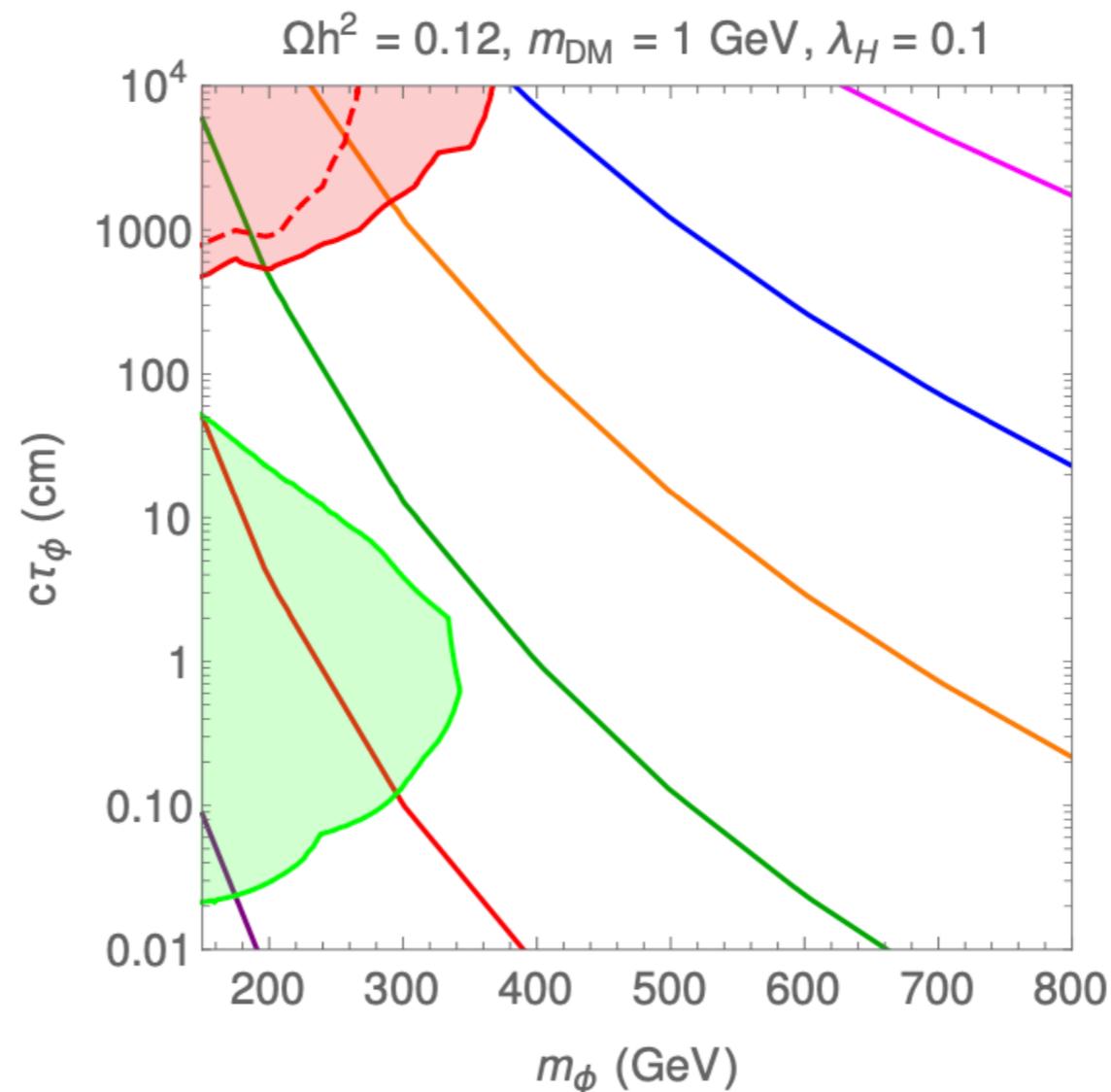
# A fermion DM example

Fermion dark matter with a scalar partner coupled to leptons

$$\bar{l} \chi \Phi_B$$

**Heavy Stable  
Charged Particles**

**Displaced Leptons**



**PRELIMINARY**

- $T_{\text{rh}} = 80 \text{ GeV}$
- $T_{\text{rh}} = 40 \text{ GeV}$
- $T_{\text{rh}} = 20 \text{ GeV}$
- $T_{\text{rh}} = 10 \text{ GeV}$
- $T_{\text{rh}} = 5 \text{ GeV}$
- $T_{\text{rh}} = 2 \text{ GeV}$

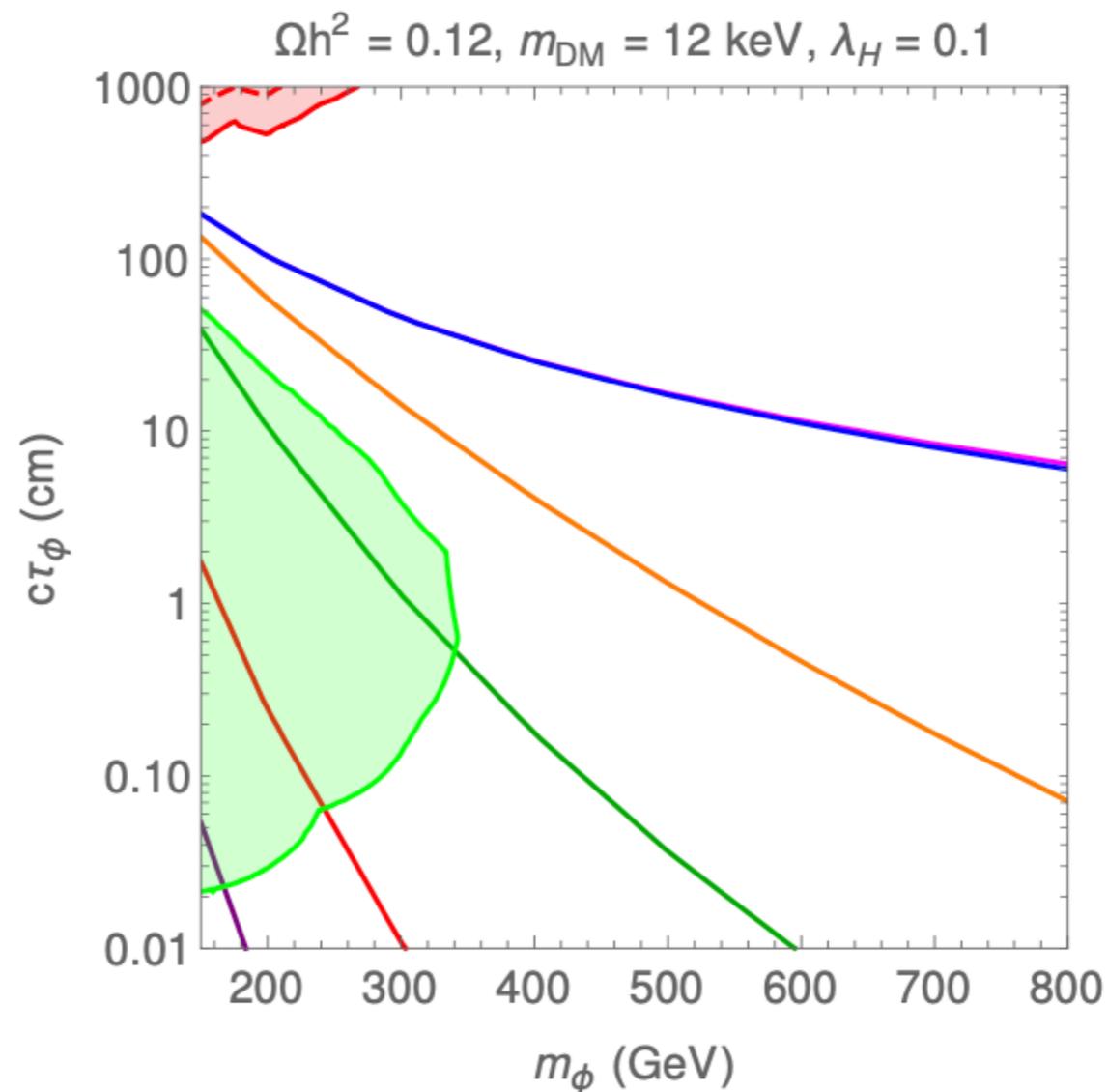
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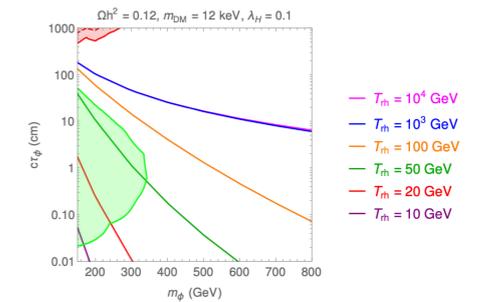
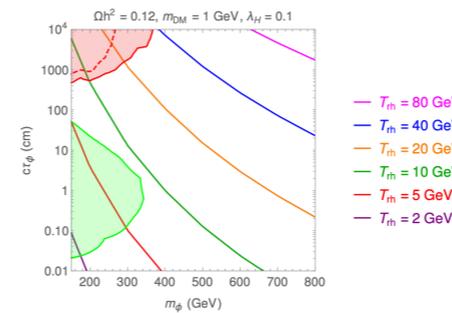
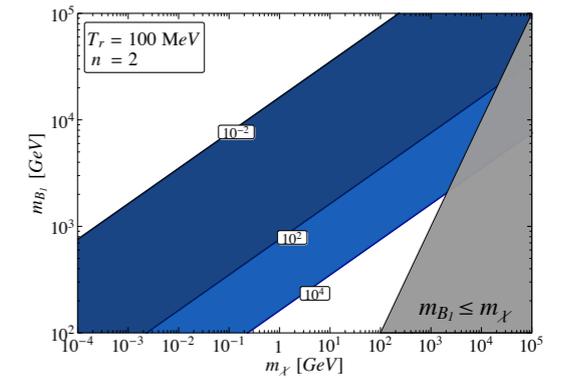
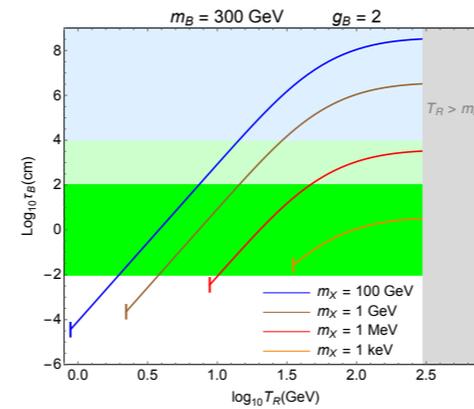


**PRELIMINARY**

- $T_{\text{rh}} = 10^4 \text{ GeV}$
- $T_{\text{rh}} = 10^3 \text{ GeV}$
- $T_{\text{rh}} = 100 \text{ GeV}$
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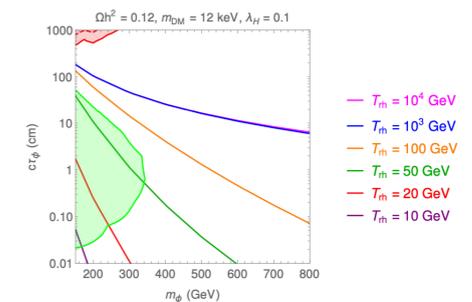
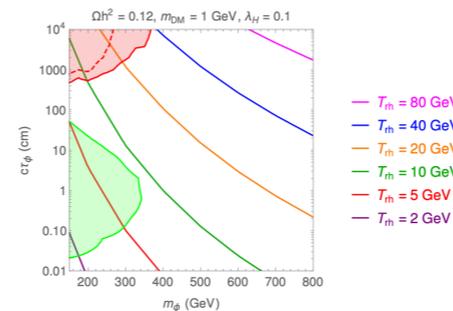
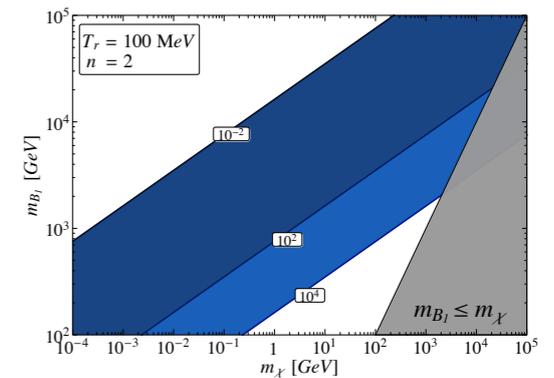
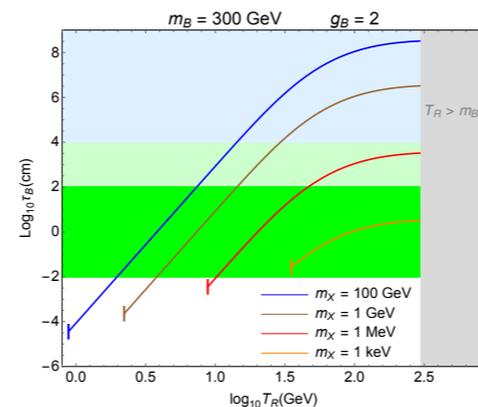
# Outlook

- Modified cosmologies and displaced signatures
- Bottom-up classification for displaced signatures @LHC



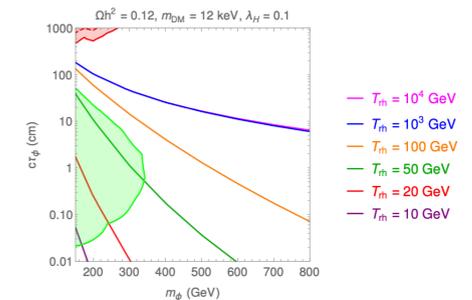
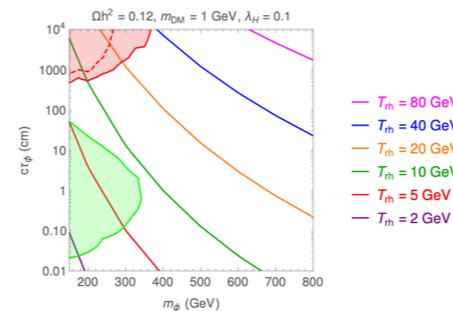
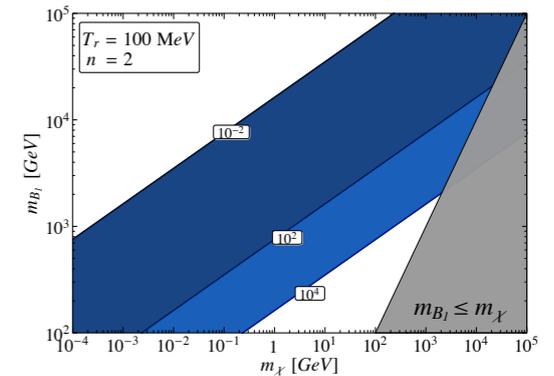
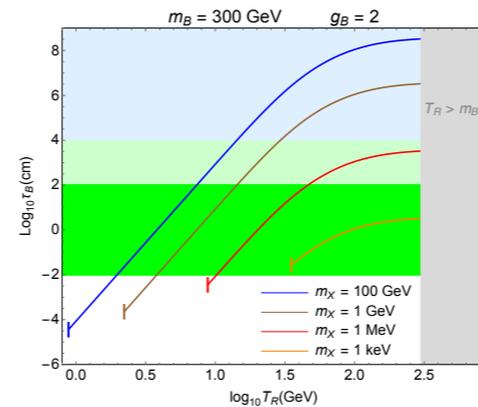
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**Thank you**