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# Light-cone Reduction

of Witten's open String Field Theory

based on JHEP04(2019)143 [arXiv:1901.08555]



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# Motivation

- We formulated cov. SFTs !!
  - Beautiful but complicated.
- We want more usable one as "a tool" ...
- As a tool, "light-cone SFT" may be better.
  - Easy to handle
  - Simple action even for closed strings
  - Quartic for open strings

#### Infinite # of vertices

$$S_{cov} = \frac{1}{2}\psi Q\psi + \frac{1}{3!}\psi^3 + \sum_{n>3}^{\infty} \frac{1}{n!}\psi^n$$

#### & quantum corrections



#### Just a Cubic vertex

 $S_{lc} = \frac{1}{2}\phi L^{lc}\phi + \frac{1}{3!}\phi^{3}$ 

#### without quantum correction

# Our ultimate goal

- Relating "light-cone SFT" with "covariant SFT".
  - People may expect that "light-cone theory" is a gauge fixed version of "covariant theory".
  - No explicit construction. . .

As a first step, we consider
 how the light-cone string field
 appears from covariant SFT

#### Lorentz covariant theory

Some gauge fixing?

#### **Light-cone theory**



#### Two differences

• "Field contents" (world-sheet variables) are different.

- Light-cone string field … Physical modes

$$\Psi_{lc} = \Psi_{lc}[X^{I}(\sigma), p^{+}, x^{-}]$$

- Covariant string field … (Phys + [gauge + unphys]) modes

$$\Psi_{cov} = \Psi_{cov}[X^{\mu}(\sigma), b(\sigma), c(\sigma)]$$
  
=  $\Psi_{cov}[X^{I}(\sigma), p^{+}, x^{-}; \partial X^{\pm}(\sigma), b(\sigma), c(\sigma)]$ 

**BRST quartets** 

 $\cdot$  "Neumann coefficients" (string-overlaps) are not equivalent.

LC type: 
$$\int dp^+ \begin{bmatrix} p_1^+ \\ p_2^+ \end{bmatrix} p_3^+$$
 Witten type:

## Sketch of light-cone & covariant theories

	Witten type	p-dependent overlap
Physical modes		Kaku-Kikkawa's LC theory
Physical modes + "∞ BRST quartets"	Witten theory	

So we can expect . . .

## Sketch of light-cone & covariant theories



Light-cone reduction (Removing the quartets)
 Field redefinition

# Sketch of light-cone & covariant theories



Light-cone reduction (Removing the quartets)
Field redefinition

# Our strategy

Covariant String Field
 = Light-cone String Field
 + BRST quartets

- Gauge fixing as making
   an effective field theory
  - Classical part = our LC action

 $S_{eff}[\psi_{lc}] = S[\psi_{lc}] + (\text{loop})$ 

 In terms of math...
 "Homological perturbation" for A<sub>∞</sub> / L<sub>∞</sub>

#### After a (linear) field redefinition...

$$\Psi_{\rm cov} = \underbrace{\Psi_{\rm lc}}_{\rm lc} + \Psi_{\rm quartet}$$

physical

gauge + unphys

#### Integrating out "gauge degrees"

$$\begin{split} \int \mathcal{D}[\psi] e^{-S[\psi]} &= \int \mathcal{D}[\psi_{lc}, \psi_g] e^{-S[\psi_{lc} + \psi_g]} \\ &= (\operatorname{Vol}_g) \cdot \int \mathcal{D}[\psi_{lc}] e^{-S_{eff}[\psi_{lc}]} \end{split}$$



# Our results

- Light-cone SFT from Witten's covariant SFT
  - A new type of theory
  - Different from the old Kaku-Kikkawa's one.

Amplitudes are exactly the same as Witten's SFT.
 (by construction)

Witten's SFT

 $S_{cov} = \frac{1}{2}\psi Q\psi + \frac{1}{3}\psi^3$ 

LC reduction

#### Light-cone SFT with A∞

$$S_{lc} = \frac{1}{2} \psi_{lc} c_0 L^{lc} \psi_{lc} + \frac{1}{3} \psi_{lc}^3 + \sum_{n>3} \frac{1}{n} \psi_{lc}^n$$



## Plan

- Introduction & Summary  $\checkmark$
- LC reduction of Witten's SFT

#### (a) Sketch of LC reduction

- Field redefinition
- Effective field theory

(b) More on the reduction- Homological perturbation& path-integrating-out



#### (a) Sketch of LC reduction

- Review of BRST quartets
- Find a field redefinition  $\Psi_{cov} = \mathscr{U}\Psi$  s.t.  $\Psi_{cov} = \underbrace{\Psi_{lc}}_{physical} + \underbrace{\Psi_{quartet}}_{unphys + gauge}$ - Integrating-out  $\Psi_{quartet}$  (= Homological Perturbation )

## Review: BRST quartets

•



$$L_{0} \equiv \underbrace{\frac{1}{2}p^{2} + \sum_{n \ge 1} a_{-n}^{I} a_{n}^{I} - 1}_{= L_{0}^{lc}} + \underbrace{\sum_{n \ne 0} \left[ n : c_{-n} b_{n} : - : a_{-n}^{+} a_{n}^{-} : \right]}_{= N}$$

 $N = \left[ d, \frac{1}{p^+} \sum_{n \neq 0} a_{-n}^+ b_n \right]$ 

- **N** counts the level of BRST quartets :  $\{a_{-n}^+, c_{-n}^-, ; b_n^-, a_n^-\}$
- Nilpotent operator  $d \equiv -p^+ \sum_{n \neq 0} c_{-n} a_n^-$  generates BRST transformations

$$\delta_{\rm B}(a_n^+) \equiv \left[\!\!\left[ d \, , \, a_n^+ \, \right]\!\!\right] = -n \left( p^+ c_n \right) \qquad \qquad \delta_{\rm B}\left( \frac{1}{p^+} b_{-n} \right) \equiv \left[\!\!\left[ d \, , \, \frac{1}{p^+} b_{-n} \, \right]\!\!\right] = a_{-n}^-$$

and **N** is a "*d*-exact operator".

Any "*d*-closed" state A satisfying N A = n A (n≠0) is "*d*-exact"

No **d**-cohomology except for **Ker**[**N**]!

## Light-cone & BRST

 $Q_{1} = \sum_{n \neq 0} c_{-n} \left[ \frac{1}{2} \sum_{k} a_{n-k}^{i} a_{k}^{i} - \sum_{k \neq 0} a_{n-k}^{+} a_{k}^{-} \right] - \frac{1}{2} \sum_{\substack{n \neq 0 \\ n \neq n \neq 0 \\ n \neq n \neq 0}} (m-n) : c_{-n} c_{-n} b_{n+m} :$   $Q \equiv c_{0} L_{0} + [Q_{non \ zero}] - b_{0} \sum_{n \neq 0} n : c_{-n} c_{n} :$   $= Q_{2}$   $= c_{0} (L_{0}^{lc} + N) + \left[ -p^{+} \sum_{\substack{n \neq 0 \\ n \neq 0 \\ = d}} c_{-n} a_{n}^{-} + Q_{1} - p^{-} \sum_{\substack{n \neq 0 \\ n \neq 0 \\ = Q_{3}}} c_{-n} a_{n}^{+} \right] + Q_{2}$   $= d + c_{0} (L_{0}^{lc} + N) + \sum_{k=1} Q_{k}$ 

• We find the similarity transformation :  $Q = \mathcal{U}^{-1} (c_o L^{lc} + d) \mathcal{U}$ 

$$\mathcal{U} \equiv \exp\left(-c_0 \frac{1}{p^+} \sum_{n \neq 0} a_{-n}^+ b_n\right) \exp\left(\frac{1}{p^+} \sum_{n \neq 0} \frac{1}{n} \left[\!\left[a_{-n}^+ b_n, Q_1 + \frac{1}{2} Q_2\right]\!\right]\right)$$

# Field redefinition

• LC dec. of BRST op.

 $Q = \mathcal{U}^{-1} \left( c_o L^{lc} + d \right) \mathcal{U}$ 

Field redefinition

 $\Psi \rightarrow \mathscr{U}\Psi = \underbrace{\Psi_{lc}}_{physical} + \underbrace{\Psi_{BRS\,quartet}}_{gauge}$ 

Light-cone Virasoro zero

$$L^{lc} = \frac{1}{2}p^2 + \sum_{n \ge 1} a^I_{-n} a^I_n - 1$$

BRST quartets :  $\{a_{-n}^+, c_{-n}^-, ; b_n^-, a_n^-\}$ 

$$d = -p^+ \sum_{n \neq 0} c_{-n} a_n^-$$

# Hodge type decomposition

- BRST quartets have no cohomology.
- We thus find a homotopy equivalence relation

 $dh + hd = 1 - \Pi$ 

- Actually...
  - $\Psi \rightarrow \mathscr{U}\Psi = \Pi \Psi + (dh + hd) \Psi$

 $= \underbrace{\Psi_{lc}}_{physical} + \underbrace{\Psi_{BRS\,quartet}}_{gauge}$ 

- "h" is a kind of propagator.

 $Q(bL^{-1}) + (bL^{-1})Q = 1 - e^{-\infty L}$ 



#### d : Quartets' Differential

$$d = -p^+ \sum_{n \neq 0} c_{-n} a_n^-$$

#### h: Homotopy contracting op.

$$h = \frac{1}{p^+ N} \sum_{n \neq 0} a_{-n}^+ b_n \left( 1 - \Pi \right)$$

#### Π: Projector onto Ker[N]

# Witten theory

String field redefinition

 $\Psi_{\rm cov} = \mathscr{U} \Psi$ 

- Witten theory becomes

$$S_{\text{Witten}} = \frac{1}{2} \Psi_{\text{cov}} \left( c_0 L^{lc} + d \right) \Psi_{\text{cov}} + \frac{1}{3} \left( \Psi_{\text{cov}} \right)^3$$

- Integrating-out all of the gauge & unphys degrees
  - = By HP lemma for  $A_{\infty}$

(We will see later...)



#### An "effective" theory

- Roughly, in terms of field theory...
  - Expanding the action  $\Psi_{cov} = \psi + c_0 \chi$

$$S[\psi + c_0 \chi] = \frac{1}{2} \langle \psi, c_0 K^{lc} \psi \rangle + \frac{1}{3} \langle \psi, m_2^{\text{cov}}(\psi, \psi) \rangle + \langle c_0 \chi, m_2^{\text{cov}}(\psi, \psi) \rangle + \langle c_0 \chi, d \psi \rangle + \langle \psi, m_2^{\text{cov}}(c_0 \chi, c_0 \chi) \rangle + \frac{1}{3} \langle c_0 \chi, m_2^{\text{cov}}(c_0 \chi, c_0 \chi) \rangle$$

- Solving the e.o.m. of gauge modes  $\chi = (dh + hd)\chi$ 

 $c_0 \left[ d\psi + m_2^{\text{cov}}(\psi, c_0\chi) + m_2^{\text{cov}}(c_0\chi, \psi) + m_2^{\text{cov}}(\psi, \psi) + m_2^{\text{cov}}(c_0\chi, c_0\chi) \right] = 0$ 

- LC SFT from covariant SFT
  - A new type of theory

$$S_{lc} = \frac{1}{2} \psi_{lc} c_0 L^{lc} \psi_{lc} + \frac{1}{3} \psi_{lc}^3 + \sum_{n>3} \frac{1}{n} \psi_{lc}^n$$



## (b) More on the reduction

- Homological Perturbation lemma & Integrating out  $\Psi_{unphys+gauge}$
- Application to Other types of reduction

## Review: HP lemma

For given "Standard situation"

$$h_{\bigcirc}(V,\partial_V) \stackrel{f}{\longleftrightarrow} (W,\partial_W)$$

 $1_V - g f = \partial_V h + h \, \partial_V$ 

& given "perturbation"

 $\Delta$  s.t.  $\left(\partial_V + \Delta\right)^2 = 0$ 

New Standard Situation

$$\widetilde{h}_{C} \left( V, \, \partial_{V} + \Delta \right) \stackrel{\widetilde{f}}{\underset{\widetilde{g}}{\longleftrightarrow}} \left( W, \, \widetilde{\partial}_{W} \right)$$
$$1_{V} - \widetilde{g} \, \widetilde{f} = \left( \partial_{V} + \Delta \right) \widetilde{h} + \widetilde{h} \, \left( \partial_{V} + \Delta \right)$$

V, W: vector space

- *∂v, ∂w* : nilpotent differential
- **f** and **g** preserve its cohomology (quasi-isomorphism)
- 1 and g f : homotopy equivalent



## Review: HP lemma

For given "Standard situation"

 $h_{\bigcirc}(V, \partial_V) \stackrel{f}{\longleftrightarrow} (W, \partial_W)$ 

 $1_V - g f = \partial_V h + h \, \partial_V$ 

& given "perturbation"

 $\Delta$  s.t.  $\left(\partial_V + \Delta\right)^2 = 0$ 

New Standard Situation

$$\widetilde{h}_{C}(V, \partial_{V} + \Delta) \xrightarrow{\widetilde{f}}_{\widetilde{g}} (W, \widetilde{\partial}_{W})$$

$$1_V - \widetilde{g}\,\widetilde{f} = (\partial_V + \Delta)\,\widetilde{h} + \widetilde{h}\,(\partial_V + \Delta)$$

V, W: vector space

*∂v, ∂w* : nilpotent differential

**f** and **g** preserve its cohomology (quasi-isomorphism)

1 and g f: homotopy equivalent

Input :

- Standard situation
- any perturbation

Output :

- New standard situation !!

## Review: HP lemma

For given "Standard situation"

 $h_{C}(V, \partial_{V}) \stackrel{f}{\longleftrightarrow} (W, \partial_{W}) \lt$ 

 $1_V - g f = \partial_V h + h \, \partial_V$ 

& given "perturbation"

 $\Delta$  s.t.  $\left(\partial_V + \Delta\right)^2 = 0$ 

**New** Standard Situation

$$\widetilde{h}_{C} \left( V, \partial_{V} + \Delta \right) \stackrel{\widetilde{f}}{\underset{\widetilde{g}}{\longleftrightarrow}} \left( W, \widetilde{\partial}_{W} \right)$$
$$1_{V} - \widetilde{g} \, \widetilde{f} = \left( \partial_{V} + \Delta \right) \widetilde{h} + \widetilde{h} \, \left( \partial_{V} + \Delta \right)$$

Roughly, in terms of QFT ...

- Homotopy equivalence -

We have a theory  $S_V = \frac{1}{2} \psi_v \partial_V \psi_v$ 

and its field-contents split

 $\phi_W$ 

 $\psi_v = (gf) \psi_v + (\partial_V h + h \partial_V) \psi_v$ 

no cohomology

- Standard situation -

By integrating-out  $(\partial_V h + h \partial_V) \psi_v$ ,  $\int \mathscr{D}[\psi_v] e^{-S_v[\psi_v]} \longrightarrow \int \mathscr{D}[\phi_w] e^{-S_w[\phi_w]}$  **S**<sub>V</sub> reduces to an equivalent one  $S_W = \frac{1}{2} \phi_w \partial_W \phi_w$ 

#### Explicit form of A<sub>∞</sub> vertices

• Start with "quartets"  $\{a_{-n}^+, c_{-n}^-; b_n, a_n^-\}_{n \neq 0}$ 

$$h_{C}(\mathcal{H}_{cov}, d) \stackrel{\pi}{\underset{\iota}{\longleftrightarrow}} \left( \underbrace{\operatorname{Ker}[N]}_{=\mathcal{H}_{lc}}, 0 \right) \quad \text{with} \quad 1_{\mathcal{H}_{cov}} - \underbrace{\iota \pi}_{=\Pi} = dh + hd$$

• Add "LC physics"  $c_0 L_0^{lc}$  as a perturbation

#### Example

• Removing BRST quartets :  $\{a_{-n}^+, c_{-n}^-; b_n, a_n^-\}$  except for "n=1"

Consider 
$$d_{[1]} \equiv -p^{+} \sum_{|n|>1} c_{-n} a_{n}^{-}$$

Then, we get string field theory which consists of tachyon,
 Covariant massless modes & transverse others

$$S[\phi] = \frac{1}{2} \langle \phi, \left[ c_0 L_0^{lc} + (d - d_{[1]}) \right] \phi \rangle + \sum_{n=2}^{\infty} \frac{1}{n+1} \langle \phi, \widetilde{m}_n(\underbrace{\phi, \dots, \phi}_n) \rangle$$

Gauge transformation of the string field

$$\delta\phi = \left[c_0 L_0^{lc} + (d - d_{[1]})\right]\lambda + \sum_{\text{cyclic}} \sum_{n=1}^{\infty} \widetilde{m}_{n+1}(\underbrace{\phi, \dots, \phi}_{n}, \lambda)$$

# Summary

- Our light-cone SFT has an A-infinity action.
- Different from Kaku Kikkawa's old LC theory.
- Consistent as a LC theory & No gauge degree.
- Our LC SFT has the same
   S-matrix as Witten's SFT.



	Witten type	p-dependent overlap	
Physical modes	LC theory w/ A∞	Kaku-Kikkawa's LC theory	
Physical modes + "∞ BRS quartets"	Witten theory	Some SFT	

	Witten type	p-dependent overlap
Physical modes	LC theory w/ A∞	Kaku-Kikkawa's LC theory
Physical modes + "∞ BRS quartets"	Witten theory	<b>a = p+ HIKKO</b> [Kugo-Zwiebach 92']
	Unitary (but almost cov.) ver. of <b>HIKKO's SFT</b>	

	Witten type	p-dependent overlap	
Physical modes	LC theory w/ A∞	Kaku-Kikkawa's LC theory	
Physical modes		5	
+ "∞ BRS quartets"	Witten theory	<b>a = p+ HIKKO</b> [Kugo-Zwiebach 92']	



		Witten type		p-dependent overlap
Physical modes		LC theory w/ A∞		Kaku-Kikkawa's LC theory
Physical modes + "∞ BRS quartets"		Witter	n theory	<b>a = p+ HIKKO</b> [Kugo-Zwiebach 92']
Physical modes + "(∞ <b>+1</b> ) BRS quartets"		<b>Osp(26,2</b> <b>extended</b> [ Siegel-Zwieba	2 <b>/2)-</b> theory ch, Kugo ]	

## Summary + a





# "Thank you !!"