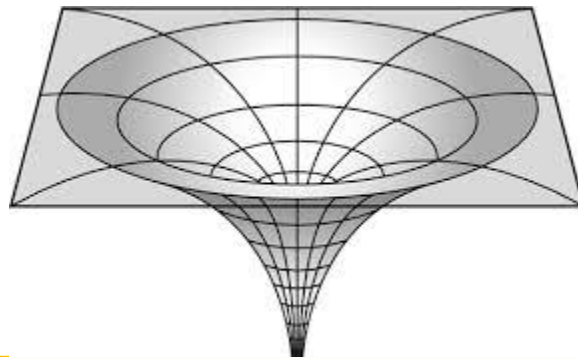


# BOOTSTRAPPING NEWTONIAN GRAVITY



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R. Casadio, M. Lenzi and O. Micu, Phys. Rev.D 98, 104016 (2018)  
R. Casadio, M. Lenzi and O. Micu, arXiv:1904.06752 [gr-qc]

# Introduction

- Motivations
- Bootstrapped approach
- Main results

# Motivations

- Gravitational collapse

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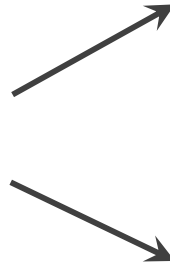
- Gravitational collapse



Singularity theorems

# Motivations

- Gravitational collapse



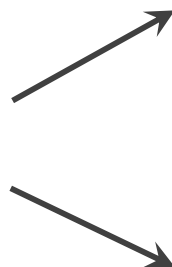
Singularity theorems

Buchdal limit

$$R > \frac{9}{4} G_{\text{N}} M \equiv R_{\text{BL}}$$

# Motivations

- Gravitational collapse



Singularity theorems

Buchdal limit

$$R > \frac{9}{4} G_{\text{N}} M \equiv R_{\text{BL}}$$

- Corpuscular model



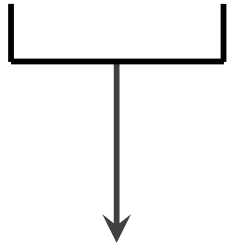
$$\lambda_{\text{G}} \sim R_{\text{H}} \quad N_{\text{G}} \sim \frac{M^2}{m_{\text{p}}^2} \sim \frac{R_{\text{H}}^2}{\ell_{\text{p}}^2}$$

# Bootstrapped approach

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# Effective field theory : bootstrap

$$L[V] = L_{\text{N}}[V]$$



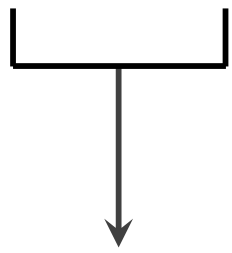
Newtonian part

$$\Delta V = 4\pi G_{\text{N}} \rho$$

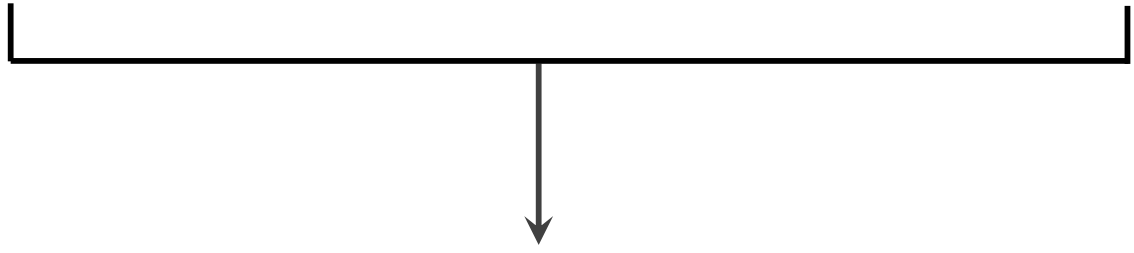


# Effective field theory : bootstrap

$$L[V] = L_N[V] - 4\pi \int_0^\infty r^2 dr [q_\Phi J_V V + q_\rho J_\rho (\rho + p)]$$



Newtonian part



Self-interaction terms

$$\Delta V = 4\pi G_N \rho$$

$$J_V \simeq \frac{dU_N}{d\mathcal{V}} = -\frac{[V'(r)]^2}{2\pi G_N}$$

$$J_\rho = -2V^2$$

# Effective field theory : bootstrap

$$L[V] = -4\pi \int_0^\infty r^2 dr \left[ \frac{(V')^2}{8\pi G_N} (1 - 4V) + (\rho + p) V (1 - 2V) \right]$$

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- Non-perturbative solution  $\longrightarrow$  BOOTSTRAP

# Main results

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# Homogeneous ball in vacuum

$$\rho = \frac{3 M_0}{4 \pi R^3} \Theta(R - r)$$

with

$$M_0 = 4 \pi \int_0^R r^2 dr \rho(r)$$

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Boundary conditions :

$$V'_{\text{in}}(0) = 0$$

$$V_{\text{in}}(R) = V_{\text{out}}(R) \equiv V_R$$



$$M_0 = M_0(M)$$

$$V'_{\text{in}}(R) = V'_{\text{out}}(R) \equiv V'_R$$

# Outer vacuum solution

$$\Delta V = \frac{2(V')^2}{1-4V} \quad \longrightarrow \quad V_c = \frac{1}{4} \left[ 1 - \left( 1 + \frac{6G_N M}{r} \right)^{2/3} \right]$$

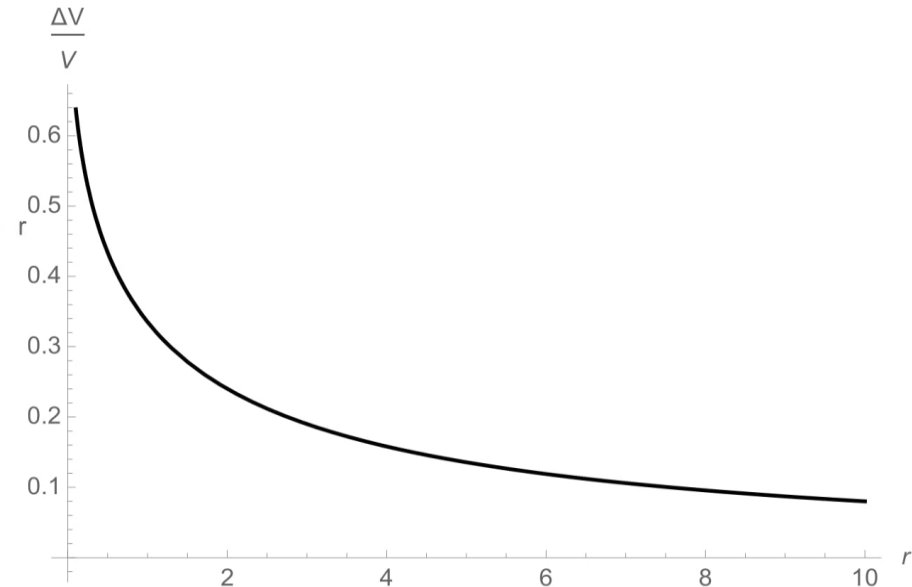
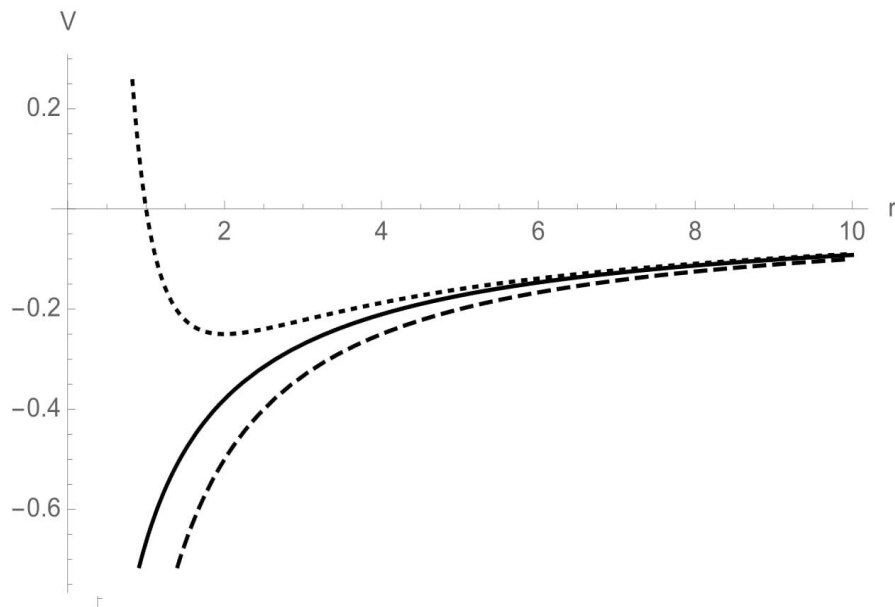
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- Large  $r$  expansion

$$V_c \underset{r \rightarrow \infty}{\simeq} -\frac{G_N M}{r} + \frac{G_N^2 M^2}{r^2}$$

- Always attractive potential



# Inner potential

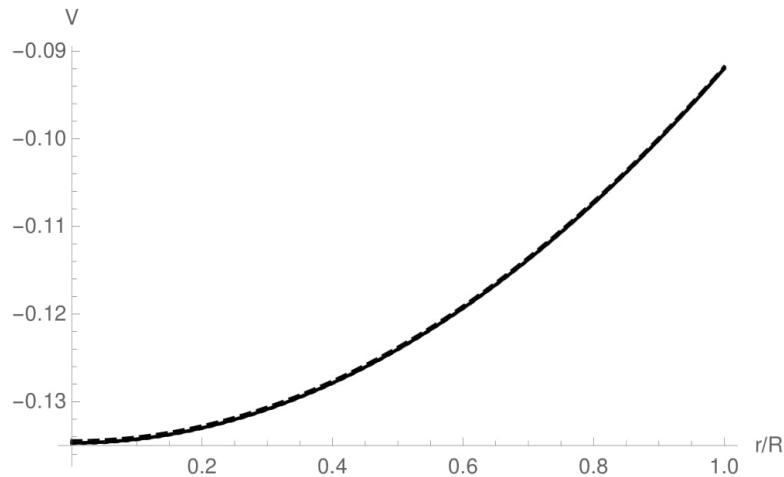
$$\Delta V = \frac{3 M_0}{R^3} e^{V_R - V} + \frac{2 (V')^2}{1 - 4V}$$

# Inner potential

$$\Delta V = \frac{3 M_0}{R^3} e^{V_R - V} + \frac{2 (V')^2}{1 - 4V}$$

Small/intermediate compactness

$$\frac{G_N M}{R} \ll 1 \quad \frac{G_N M}{R} \simeq 1$$



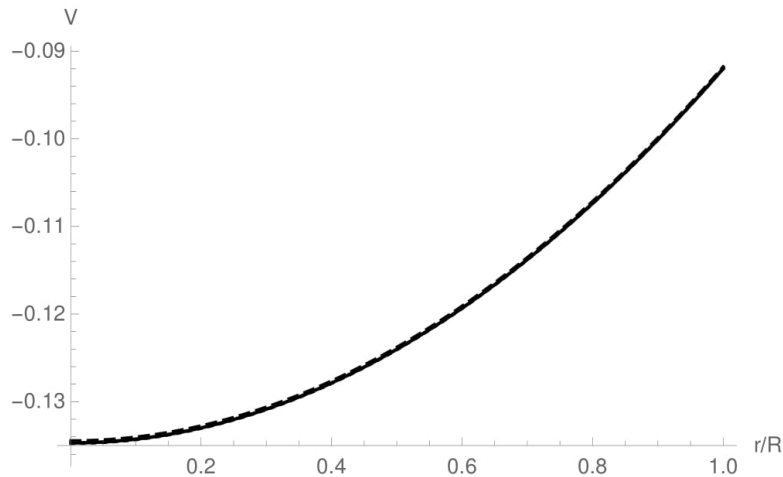
$$V_s = V_0 + \frac{G_N M_0}{2 R^3} e^{V_R - V_0} r^2$$

# Inner potential

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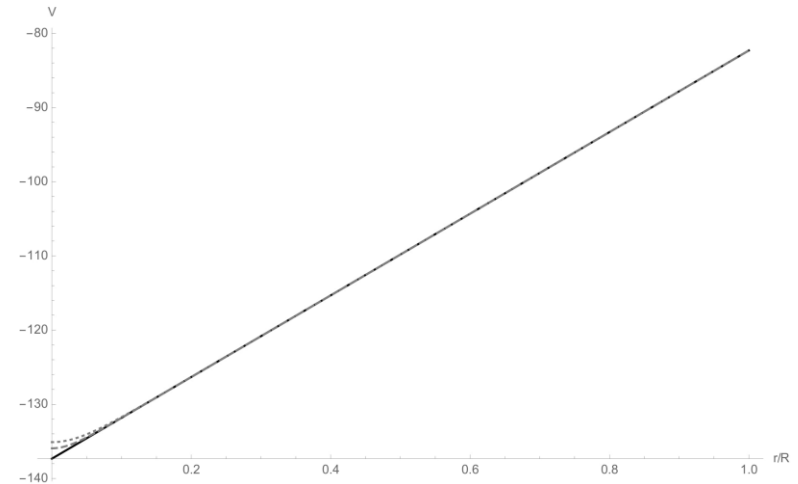
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High compactness

$$\frac{G_N M}{R} \gg 1$$



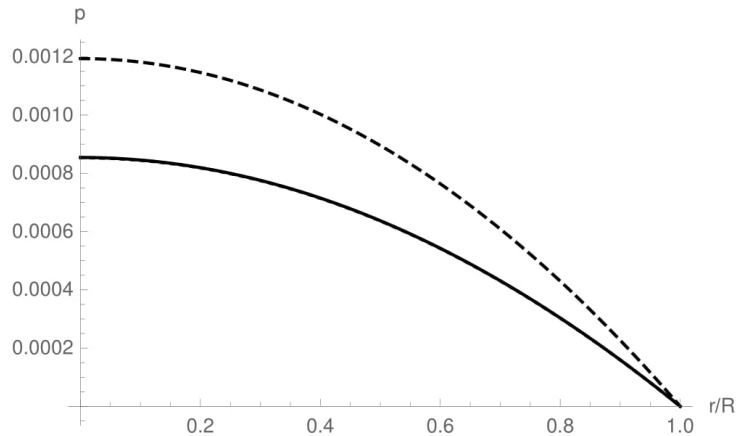
$$V_{\text{lin}} \simeq V_R + V'_R (r - R)$$

$$\frac{G_N M_0}{R} \sim \left( \frac{G_N M}{R} \right)^{2/3}$$

# Pressure

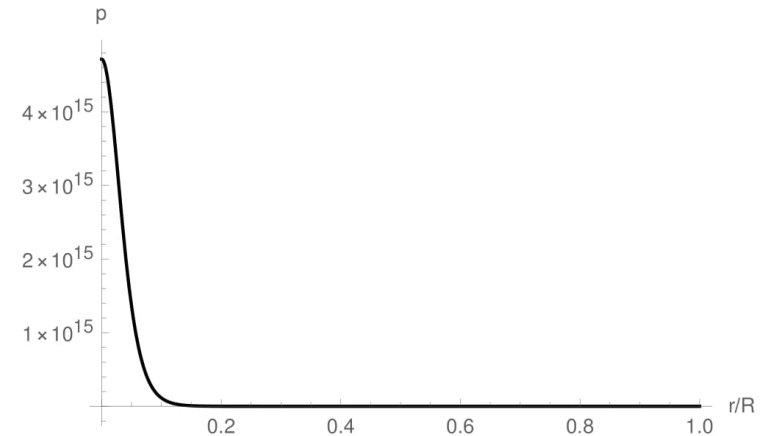
$$p = \rho_0 \left[ e^{V_R - V} - 1 \right]$$

Small/intermediate compactness



$$p \simeq \frac{3 G_N M^2 (R^2 - r^2)}{8 \pi R^6}$$

High compactness



$$p \simeq \frac{M^2 e^{\frac{1}{2} \left( \frac{G_N M}{\sqrt{6} R} \right)^{2/3} \left( 3 - \frac{5}{\tilde{C}} \right)}}{2 \pi \tilde{C}^2 (6 G_N M/R)^{2/3}} \left[ e^{\left( \frac{G_N M}{\sqrt{6} R} \right)^{2/3} \left( 1 - \frac{r}{R} \right)} - 1 \right]$$

# Horizon

- GR geometrical approach  $\longrightarrow$  Schwarzschild radius  $R_{\text{H}} \equiv 2 G_{\text{N}} M$

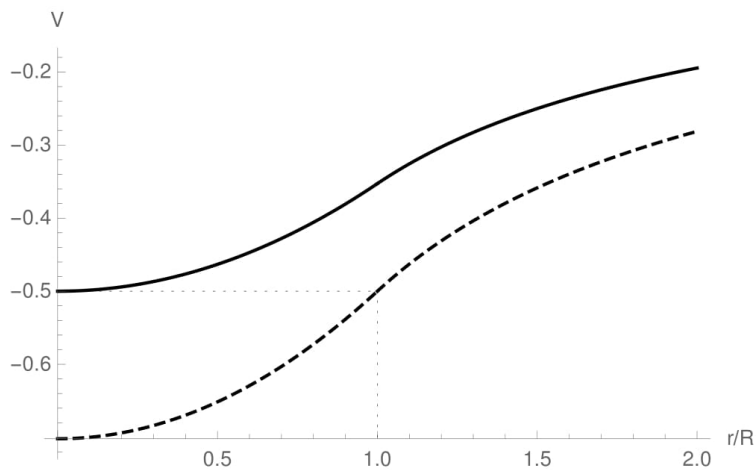


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$$\longrightarrow \begin{cases} \text{no horizon} & \text{for } G_{\text{N}} M/R \lesssim 0.46 \\ 0 < r_{\text{H}} \leq R \simeq 1.4 G_{\text{N}} M & \text{for } 0.46 \lesssim G_{\text{N}} M/R \leq 0.69 \\ r_{\text{H}} \simeq 1.4 G_{\text{N}} M & \text{for } G_{\text{N}} M/R \gtrsim 0.69 . \end{cases}$$

# Conclusions

- Investigation of the quantum features of this potential
- Connection with the corpuscular model