

# Exercises

## Lecce Lectures on the Cosmological Bootstrap

1. Consider a massless scalar field in four-dimensional de Sitter space, with metric

$$ds^2 = \frac{-d\eta^2 + d\mathbf{x}^2}{(H\eta)^2},$$

where  $\eta$  is conformal time and  $H$  is the Hubble parameter.

- a. Determine the wave equation satisfied by the field. Show that for each Fourier mode you get solutions of the form

$$f_k(\eta) = \frac{H}{\sqrt{2k^3}}(1 + ik\eta)e^{-ik\eta},$$

and its complex conjugate.

- b. Promote  $\phi_{\mathbf{k}}$  to a quantum operator and write

$$\hat{\phi}_{\mathbf{k}}(\eta) = f_k(\eta)\hat{a}_{\mathbf{k}} + f_k^*(\eta)\hat{a}_{-\mathbf{k}}^\dagger.$$

Using the canonical commutation relation for  $\hat{\phi}$  and its conjugate momentum  $\hat{\pi}$ , find the commutator  $[\hat{a}_{\mathbf{k}}, \hat{a}_{\mathbf{k}'}^\dagger]$ . What is the physical meaning  $\hat{a}_{\mathbf{k}}$  and  $\hat{a}_{\mathbf{k}}^\dagger$ .

- c. Define the Bunch Davies vacuum through  $\hat{a}_{\mathbf{k}}|0\rangle = 0$ . Then compute

$$\langle 0 | \hat{\phi}_{\mathbf{k}}(\eta) \hat{\phi}_{\mathbf{k}'}(\eta) | 0 \rangle = \frac{2\pi^2}{k^3} P(k, \eta) \delta_D(\mathbf{k} + \mathbf{k}').$$

Determine the superhorizon limit of  $P(k, \eta)$ .

- d. By summing over all Fourier modes, calculate the equal time two-point function of the massless scalar in real space. What is the physical meaning of the IR divergence?
- e. Repeat the exercise for a conformally coupled scalar field, with  $m^2 = 2H^2$ .
- f. Photons are massless, but they are not produced during inflation. Why?

2. Add a  $\phi^3$  interaction to the above problem and compute the three-point function,

$$\langle 0 | \hat{\phi}_{\mathbf{k}_1} \hat{\phi}_{\mathbf{k}_2} \hat{\phi}_{\mathbf{k}_3} | 0 \rangle = \frac{(2\pi^2)^2}{(k_1 k_2 k_3)^2} B(k_1, k_2, k_3) \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3).$$

Understand the integration contour necessary for the computation. How do you kill off the oscillatory early time pieces? Understand that you need to rotate the contour into the imaginary direction.

*Warning:* This is not an easy problem. If you want a simpler warm-up problem, do the next problem first.

3. Consider a conformally coupled scalar field in a fixed de Sitter background.

a. Add a  $\phi^4$  interaction to the above problem and compute the four-point function,

$$\langle 0 | \hat{\phi}_{\mathbf{k}_1} \hat{\phi}_{\mathbf{k}_2} \hat{\phi}_{\mathbf{k}_3} \hat{\phi}_{\mathbf{k}_4} | 0 \rangle = F(k_1, k_2, k_3, k_4, k_I, k'_I) (2\pi)^4 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4),$$

where  $k_I \equiv |\mathbf{k}_1 + \mathbf{k}_2|$  and  $k'_I \equiv |\mathbf{k}_1 + \mathbf{k}_3|$ .

Consider only the “ $s$ -channel” permutations of the external momenta and write the result in terms of

$$u^{-1} \equiv \frac{k_1 + k_2}{k_I}, \quad v^{-1} \equiv \frac{k_3 + k_4}{k_I}.$$

Check that the result is conformally invariant, i.e. that the four-point function  $\hat{F}(u, v) \equiv k_I F$  satisfies

$$(\Delta_u - \Delta_v) \hat{F} = 0,$$

where  $\Delta_u \equiv u^2(1 - u^2)\partial_u^2 - 2u^3\partial_u$ .

b. Now, consider the interaction  $(\partial_\mu \phi)^4$  and compute the corresponding four-point function for “ $s$ -channel” permutations of the external momenta. Write the result in terms of  $u$  and  $v$ , and check conformal invariance. Show that the result can be written in the form

$$\hat{F} = (a_0 + a_1 \Delta_u + a_2 \Delta_u^2) \hat{F}_0,$$

where  $\hat{F}_0$  is the solution found in part a. and  $a_n$  are constants to be determined.

4. Consider a massless scalar field in an inflationary background. Add a  $(\partial_\mu \phi)^4$  interaction.

a. Evaluating one leg on the background leads to a cubic interaction for the inflationary fluctuations

$$\mathcal{L} \subset \frac{(\partial_\mu \phi)^4}{8\Lambda^4} \Rightarrow \mathcal{L}_{\text{int}} = \frac{\dot{\phi}}{4\Lambda^4} \dot{\varphi} (\partial_\mu \varphi)^2,$$

where  $\varphi(x) \equiv \phi(x) - \bar{\phi}(t)$ . Compute the three-point function associated with this interaction.

b. Re-derive this result from the cosmological bootstrap.