Galactic Structure and Dark Matter Indirect Detection

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General outline

Indirect detection of dark matter: Theoretical motivation

The role of substructures

Subhalo population of our galaxy: -Prospects for detection in γ -rays -Multi-wavelength analysis

Indirect detection of dark matter: Theoretical motivation

Framework

The Standard Model of the Universe, as derived from data on large scale structures, distant supernovae, CMB, etc.

 $\Omega_{\text{tot}} \equiv \frac{\rho_{\gamma} + \rho_{\nu} + \rho_{b} + \rho_{DM} + \rho_{\Lambda}}{\rho_{c}} \sim 1$ value for a flat universe

predicts the existence of

- an unknown form of repulsive energy, or dark energy $\Omega_{\rm A}\sim 0.73$
- and an unknown type of non baryonic matter, or

DARK MATTER $\,\Omega_{\text{DM}}\sim0.23$

Framework The Standard Model of the Universe









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Framework The Standard Model of the Universe



 $\Omega_{\rm b} \sim 0.005 \text{ (galaxies)} \\ \sim 0.04 \text{ (BBN)} \\ \Omega_{\gamma} \sim 10^{-5} \\ 1.2 \cdot 10^{-3} < \Omega_{\gamma} < 1.5 \cdot 10^{-2} \\ \end{array}$



Dark because we haven't seen it (yet)

Matter because it interacts gravitationally building up the universe



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Micro to macro different particles have different free-streaming mass and damping scales and produce a different power spectrum of primordial density fluctuations



Dark because we haven't seen it (yet) Matter because it interacts gravitationally building up the universe

The evolution of DM halos Primordial density fluctuations grow and collapse in gravitationally bound structures which eventually virialize. 10-3 The assembly hystory depends on the particle. CDM proceeds via hierarchical merging, HDM via fragmentation. Baryons are captured in the dark matter potential well and form galaxies, clusters, etc.



Dark because we haven't seen it (yet) Matter because it interacts gravitationally building up the universe

Comparison with the data The observed large scale structure of the Universe compels the dark matter particle to be heavier than a few keV -> CDM scenario





Is the Neutralino the theoretical miracle?



BIG EXPERIMENTAL EFFORTS



x accelerators Many possible patterns for final states (jets and missing energy) $m_{\chi} < 1-2$ TeV **x** direct searches χ elastic diffusion on nuclei (nucleus recoil energy) 30 GeV < $m_{\chi} < 100-200$ GeV

- in close massive objects (Earth, Sun) \rightarrow neutrinos

x indirect searches ($\chi\chi$ annihilation) - in the Galactic halo and other compact objects, $\rightarrow \gamma_s$, v_s , antimatter

AND ACTIVE THEORETICAL MODELING



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Indirect detection of dark matter:

The role of substructures

Subhalo population of our galaxy: -Prospects for detection in γ -rays -Multi-wavelength analysis





Computing
$$\Phi_{\gamma} = \Phi_{\text{particle physics}} \times \Phi_{\text{cosmology}}$$



Pieri, Pizzella et al 2008 & Pieri, Bertone, Branchini 2008



Modeling the structure of dark matter halos

Halos form through a hierarchical process of successive mergers. The halo of our Galaxy will be self-similarly composed by: -a smoothly distributed component ($\rho^2_{DM(h)}$ single halo) -a number of virialized substructures ($\rho^2_{DM(subh)}$ all halos)



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N-body simulations study the smooth halo and the larger halos (M> $10^5 M_{sun}$).

Microphysics and theory of structure formation sets the mass of the smallest halo because there is no enough cpu power to simulate small halos from collapse till today.

Modeling the structure of dark matter halos from theory of structure formation (M< $10^5 M_{sun}$)

<u>Theory</u>: Damping of the primordial power spectrum due to CDM free streaming or acoustic oscillations after kinetic decoupling

Typical M_{min} for a WIMP = 10⁻⁶ M_{sun} 10⁻² 10⁻⁴ 10⁻⁴ 10⁻⁶ M_{sun}



patch

z=26

Diemand et al, 2005

 10^{-6}



Modeling the structure of dark matter halos from N-body simulations (M> 10⁵ M_{sun})



Via Lactea 2, Diemand et al

Aquarius, Springel et al

Modeling the structure of dark matter halos from N-body simulations (M> 10⁵ M_{sun})

Halo and subhalo profile shape and concentration



Concentration parameter (R_{vir}/r_s) has radial dependence higher concentration -> higher flux!



Modeling the structure of dark matter halos from N-body simulations (M> 10⁵ M_{sun})

Subhalo abundance and density distribution

Mass slope ~ M^{-2} f_{DM} (>10⁷ M_{sun}) ~ 11% f_{DM} (>10⁻⁶ M_{sun}) ~ 50%

Radial distribution ~ (1+R/rs)⁻¹ Mass slope ~ M^{-1.9} f_{DM} (>10⁷ M_{sun}) ~ 13% f_{DM} (>10⁻⁶ M_{sun}) ~ 25%

Radial distribution \sim Einasto α =0.67

Roche criterion sets the effect of tidal forces

Indirect detection of γ -rays: $\Phi_{\gamma} = \Phi_{\text{particle physics}} \times \Phi_{\text{cosmology}}$ $\Phi_{PP} = \frac{1}{4\pi} \frac{\sigma_{ann} v}{2m_{v}^{2}} \int_{F}^{m_{\chi}} \sum_{f} \frac{dN_{f}^{\gamma}}{dE_{v}} BR_{f}$ Step 1: Single subhalo contribution MW smooth $\Phi_{COSMO}^{halo}(M,R,r) \propto \int_{V.o.s.} dV \left| \frac{\rho_{DM}^{2} \left(M,c(M,R),r(d,V(\lambda',\theta',\phi'),\psi)\right)}{d^{2}} \right|$

> We created Monte Carlo simulations of the brightest and closest subhalos

Each source is characterized by its energy spectrum



Indirect detection of γ -rays: $\Phi_{\gamma} = \Phi_{\text{particle physics}} \times \Phi_{\text{cosmology}}$

$$\Phi_{PP} = \frac{1}{4\pi} \frac{\sigma_{ann} v}{2m_{\chi}^2} \int_{E_0}^{m_{\chi}} \sum_{f} \frac{dN_f^{\gamma}}{dE_{\gamma}} BR_f$$

Step 2: Integrated contribution of all the halos (sources) along the LOS

$$\Phi_{COSMO}(\psi, \Delta \Omega) \propto \int_{M} dM \int_{C} dc \iint_{\Delta \Omega} d\vartheta d\phi \int_{I.o.s} d\lambda \left[\rho_{sh}(M, R(R_{sun}, \lambda, \psi, \vartheta, \phi)) \cdot P(c) \cdot \Phi_{COSMO}^{halo}(M, c(M, R), r(\lambda, \lambda', \psi, \vartheta', \phi')) \right]$$

We modeled the LOS integral and integrated the signal over all sources

Indirect detection of γ -rays: $\Phi_{\gamma} = \Phi_{particle physics} \times \Phi_{cosmology}$ Step 3:Integrated contribution of
EXTRAGALACTIC halos and subhalos

Computing the cosmological γ -ray flux due to DM annihilation in halos...







The γ-ray sky Galactic and extragalactic: Smooth + subhalos

PHOTONS in 5 YEAR FERMI-LIKE OBSERVATION













Next step: Compare with Fermi MAP

first 3 months: 205 sources > 0.3 GeV > 10 σ

Many sources have been identified as pulsars or other astrophysical objects

Is any of the unidentified FERMI sources a DM subhalo?

Figure from D. Smith @ TANGO in Paris

Is it so "easy"?

Assume it is! We needs the energy spectrum! Will there be enough photons to get⁺it? + + Will it be enough? +

Is there any other clue that may point towards or exclude
+ the DM hypothesis? + +

Indirect detection of dark matter:

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The multiwavelength/multimessenger/multitarget approach

Φ = ParticlePhysics x Cosmology/Astrophysics x Transport



Slide: courtesy of M. Pato



ATIC

[HESS 2009]

PPB-BETS Kobayashi H.E.S.S.

Systematic error

Broken power-law fit

H.E.S.S. - low-energy analysis

Systematic error - low-energy analysis

10³

Energy (GeV)

a wealth of DM model building which is good independently on the nature of the excesses (that is probably astrophysical)

Possible Indirect Detection trigger: cosmic ray data



The radio sky GC, no subhalos

Assume a magnetic field

7.2 mGr<0.04pc</th>7.2 mG(r/0.04pc)-20.04pc<r<3.38pc</td>1 μGr>3.38pc

Compute synchtrotron power à la Bertone 2008

$$n_{e\pm}(\bar{x}, E_{e\pm}) = \frac{\sigma V}{2m_{DM}^2} \rho_{DM}^2(\bar{x}) \frac{N_{e\pm}(>E_{e\pm})}{b_{syn}(\bar{x}, E_{e\pm})}$$
$$v \frac{dW_{syn}}{dv} = \frac{\sigma V}{2m_{DM}^2} \int_{\Delta\Omega} d\Omega \int_{\log} ds \rho_{DM}^2(\bar{x}) E(\bar{x}, v) \frac{N_{e\pm}(>E_{e\pm})}{2}$$

The antimatter sky

Compute the number density à la Delahaye 2008

 $n_{CR}(t, \bar{x}, E_{CR}) = \frac{d^2 N_{CR}}{dV dE_{CR}}$

electrons and positrons



Compute fluxes and boosts à la Lavalle 2008

$$\phi_{CR,sm}(\mathsf{E}_{CR}) \propto < \sigma \mathsf{v} > \int_{\mathsf{E}_{CR}}^{\infty} d\mathsf{E} \frac{d\mathsf{N}_{CR}}{d\mathsf{E}} \int_{diff.zone}^{\mathsf{d}3} \bar{\mathsf{x}} \left(\frac{\rho_{sm}(\bar{\mathsf{x}})}{\rho_{sun}}\right)^2 G_{sun}^{CR}(\bar{\mathsf{x}},\lambda_{\mathsf{D}})$$

$$< \phi_{CR,cl} > (\mathsf{E}_{CR}) \propto < \sigma \mathsf{v} > \mathsf{N}_{cl} \int_{\mathsf{E}_{CR}}^{\infty} d\mathsf{E} \frac{d\mathsf{N}_{CR}}{d\mathsf{E}} \int_{diff.zone}^{\mathsf{d}3} \bar{\mathsf{x}} < \xi >_{\mathsf{M}} (\mathsf{R}) \frac{d\mathsf{P}_{\mathsf{V}}}{d\mathsf{V}} (\mathsf{R}) G_{sun}^{CR}(\bar{\mathsf{x}},\lambda_{\mathsf{D}}) = \mathsf{N}_{tot}^{sub} < \phi_{sub} > \mathsf{N}_{tot}^{Sub} < \phi_{sub} > \mathsf{N}_{tot}^{Sub} < \mathsf{N}_{to}^{Sub} < \mathsf{N}_{tot}^{Sub} < \mathsf{N}_$$

The antimatter sky



Compute fluxes and boosts à la Lavalle 2008

$$\phi_{CR,sm}(E_{CR}) \propto <\sigma v > \int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff.zone}^{d^{3}} d^{3}\bar{x} \left(\frac{\rho_{sm}(\bar{x})}{\rho_{sun}}\right)^{2} G_{sun}^{CR}(\bar{x},\lambda_{D})$$

$$<\phi_{CR,cl}>(E_{CR}) \propto <\sigma v>N_{cl}\int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff.zone} d^{3}\bar{x} <\xi>_{M} (R)\frac{dP_{V}}{dV}(R)G_{sun}^{CR}(\bar{x},\lambda_{D}) = N_{tot}^{sub} <\phi_{sub}>$$
LP, Lavalle, Bertone & Branchini 2009

The antimatter sky



Compute fluxes and boosts à la Lavalle 2008

$$\phi_{CR,sm}(E_{CR}) \propto <\sigma v > \int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff,zone}^{d^{3}} \overline{x} \left(\frac{\rho_{sm}(\overline{x})}{\rho_{sun}}\right)^{2} \mathcal{G}_{sun}^{CR}(\overline{x},\lambda_{D})$$

$$<\phi_{CR,cl}>(E_{CR}) \propto <\sigma_{V}>N_{cl}\int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff.zone} d^{3}\vec{x} <\xi>_{M} (R)\frac{dP_{V}}{dV}(R)G_{sun}^{CR}(\vec{x},\lambda_{D}) = N_{tot}^{sub} <\phi_{sub}>$$
LP, Lavalle, Bertone & Branchini 2009



It mimics an attractive force which arises when the two DM particles get close and are slow. Annihilation proceeds through the exchange of massive vector bosons $\sigma_{ann}v = S(\sigma_{ann}v)_{thermal}$

Particle Physics BF: Sommerfeld enhancement

Sommerfeld effect produces a local enhancement of the annihilation cross-section which depends on the DM velocity and mass, and does not touch the thermal value

$$\frac{1}{m_{DM}}\frac{d^{2}\psi(r)}{dr^{2}} = -m_{DM}\beta^{2}\psi(r) - \frac{\alpha}{r}e^{-m_{V}r}$$



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Particle Physics BF: Sommerfeld enhancement

Sommerfeld effect produces a local enhancement of the annihilation cross-section which depends on the DM velocity and mass, and does not touch the thermal value mass dependence (resonance)



Particle Physics BF and astrophysics BF: Sommerfeld enhancement and subhalos

Dwarf galaxies and galactic subhalos have low velocity dispersions, hence the Sommerfeld enhancement should be convolved with the sub-subhalo contribution

$$\sigma_{ann} v = (\sigma_{ann} v)_{thermal} S(\beta(r))$$

$$\Phi = T(E,d) \frac{(\sigma_{ann}v)_{thermal}}{2m_{\chi}^{2}} \int_{E_{0}}^{m_{\chi}} \sum_{f} \frac{dN_{f}^{\gamma}}{dE_{\gamma}} BR_{f} \int_{I.o.s} d\lambda S(\beta(M,r)) \rho_{DM}^{2} (M,c(M,r),R)$$

This holds for all annihilation products We can perform a multi-wavelength analysis to constrain models APPLYING BOOSTS TO BOTH Φ_{PP} AND Φ_{COSMO}



Pato, LP, Bertone 2009



Pato, LP, Bertone 2009



Conclusions

Detection of individual structures

In the best case scenario high mass halos are "detectable" Result poorly dependent on small mass extrapolation

Multi-wavelength constraints

Coherent prediction of signals from all annihilation products is now necessary in order to constrain (or discover) particle physics and cosmological DM models

Upcoming data

This is more than ever important in these years when data from satellites, Cherenkov Telescopes, accelerators and Direct detection are about to allow an unprecedented insight on the DM puzzle