Galactic Structure and Dark Matter Indirect Detection

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General outline

Indirect detection of dark matter: Theoretical motivation

The role of substructures

Subhalo population of our galaxy: -Prospects for detection in γ-rays -Multi-wavelength analysis

Indirect detection of dark matter: Theoretical motivation

Framework

The Standard Model of the Universe, as derived from data on large scale structures, distant supernovae, CMB, etc.

> \sim 1 $\Omega_{\text{tot}} \equiv \frac{\rho_{\gamma} + \rho_{\nu} + \rho_{b} + \rho_{\text{DM}} + \rho_{\Lambda}}{2}$ ρ_{c} value for a flat universe known particles unknown

predicts the existence of

- or dark energy $\, \Omega_{\! \vartriangle} \! \sim$ 0.73 - an unknown form of repulsive energy,
- -- and an unknown type of non baryonic matter, or

DARK MATTER $\Omega_{DM} \sim 0.23$

Framework The Standard Model of the Universe

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∼ 0.005 (galaxies) $\Omega_{\rm b}$ ∼ 0.04 (BBN) $1.2 \cdot 10^{-3} \cdot \Omega_v \cdot 1.5 \cdot 10^{-2}$

Framework The Standard Model of the Universe

Dark because we haven't seen it (yet)

Matter because it interacts gravitationally building up the universe

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of primordial density fluctuations

Dark because we haven't seen it (yet) **Matter** because it interacts gravitationally building up the universe

CDN

The evolution of DM halos Primordial density fluctuations grow and collapse in gravitationally bound structures which eventually virialize. 10^{-3} The assembly hystory depends on the particle. CDM proceeds via hierarchical merging, HDM via fragmentation. Baryons are captured in the dark matter potential well and form galaxies, clusters, etc.

Dark because we haven't seen it (yet) **Matter** because it interacts gravitationally building up the universe

Comparison with the data The observed large scale structure of the Universe compels the dark matter particle to be heavier than a few keV -> **CDM scenario**

Is the Neutralino the theoretical miracle?

BIG EXPERIMENTAL EFFORTS

x accelerators Many possible patterns for final states (jets and missing energy) m_x < 1-2 TeV **x** direct searches *χ* elastic diffusion on nuclei (nucleus recoil energy) 30 GeV < m_x < 100-200 GeV

 $-$ in close massive objects (Earth, Sun) \rightarrow neutrinos

x indirect searches ($\chi\chi$ annihilation) - in the Galactic halo and other compact objects, $\rightarrow \gamma_s$, v_s , antimatter

AND ACTIVE THEORETICAL MODELING

x accelerators Many possible patterns for final states (jets and missing energy) m_y< 1-2 TeV **x** direct searches *χ* elastic diffusion on nuclei (nucleus recoil energy) 30 GeV < m_x < 100-200 GeV

 $-$ in close massive objects (Earth, Sun) \rightarrow neutrinos

- in the Galactic halo and other compact objects, $\rightarrow \gamma_s$, γ_s , antimatter **×** indirect searches (χχ annihilation)

Indirect detection of dark matter:

The role of substructures

Subhalo population of our galaxy: -Prospects for detection in γ-rays -Multi-wavelength analysis

Computing
$$
\Phi_{\gamma} = \Phi_{\text{particle physics}} \times \Phi_{\text{cosmology}}
$$

Pieri, Pizzella et al 2008 & Pieri, Bertone, Branchini 2008

Modeling the structure of dark matter halos

Halos form through a hierarchical process of successive mergers. The halo of our Galaxy will be self-similarly composed by: -a smoothly distributed component (ρ^2 _{DM(h)} single halo) -a number of virialized substructures (ρ^2 _{DM(subh)} all halos)

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N-body simulations study the smooth halo and the larger halos ($M > 10^5$ M_{sun}).

Microphysics and theory of structure formation sets the mass of the smallest halo because there is no enough cpu power to simulate small halos from collapse till today.

Modeling the structure of dark matter halos from theory of structure formation ($M< 10^5 M_{sun}$)

Theory: Damping of the primordial power spectrum due to CDM free streaming or acoustic oscillations after kinetic decoupling

Modeling the structure of dark matter halos from N-body simulations ($M > 10^5$ M_{sun})

Via Lactea 2, Diemand et al Aquarius, Springel et al

Modeling the structure of dark matter halos from N-body simulations ($M > 10^5$ M_{sun})

→ **Halo and subhalo profile shape and concentration**

Concentration parameter (R_{vir}/r_s) has radial dependence higher concentration -> higher flux!

Modeling the structure of dark matter halos from N-body simulations ($M > 10^5$ M_{sun})

Subhalo abundance and density distribution

Mass slope ~ M-2 ${\sf f}_{\sf DM}$ (>10⁷ ${\sf M}_{\sf sun}$) ~ 11% f_DM (>10⁻⁶ M_sun) ~ 50%

Radial distribution $\sim (1 + R/r s)^{-1}$

Mass slope $\sim M^{-1.9}$ ${\sf f}_{\sf DM}$ (>10⁷ ${\sf M}_{\sf sun}$) ~ 13% ${\sf f}_{\sf DM}$ (>10⁻⁶ ${\sf M}_{\sf sun}$) ~ 25%

Radial distribution ~ Einasto α=0.67

Roche criterion sets the effect of tidal forces

Indirect detection of γ**-rays:** Φ_{γ} = Φ particle physics \times Φ cosmology $\Phi^\mathsf{halo}_{\mathcal{COSMO}}(\mathsf{M},\mathsf{R},\mathsf{r})\propto\;\int\mathsf{dV}$ V.o.s. $\int dV \left| \frac{\rho_{DM}^2(M,c(M,R),r(d,V(\lambda',\theta',\phi'),\psi))}{d^2} \right|$ d^2 \lceil \lfloor $\|$ \vert \vert . \rfloor $\overline{}$ $\overline{}$ $\overline{}$ **MW smooth Single subhalo contribution** $\Phi_{\sf PP} = \frac{1}{4}$ 4π σ_{ann} V $2m_\chi^2$ dN_f^{γ} $\frac{2}{f}$ dE $\frac{2}{f}$ \sum $\mathsf{E_{0}}$ m_{χ} $\int \sum \frac{dP}{dE}BR_f$ **Step 1:**

> **We created Monte Carlo simulations of the brightest and closest subhalos**

Each source is characterized by its energy spectrum

Indirect detection of γ**-rays:** Φ_{γ} = Φ particle physics \times Φ cosmology m_g

$$
\Phi_{\rm PP} = \frac{1}{4\pi} \frac{\sigma_{\rm ann} v}{2m_{\chi}^2} \int_{E_0}^{m_{\chi}} \sum_{f} \frac{dN_{f}^{\gamma}}{dE_{\gamma}} BR_{f}
$$

Integrated contribution of all the halos (sources) along the LOS Step 2:

$$
\Phi_{\text{COSMO}}(\psi, \Delta\Omega) \propto \int_{M} dM \int_{c} dc \iint_{\Delta\Omega} d\vartheta d\varphi \int_{l.o.s} d\lambda \left[\rho_{sh}(M, R(R_{sun}, \lambda, \psi, \vartheta, \varphi)) \cdot P(c) \cdot \Phi_{\text{COSMO}}^{halo}(M, c(M, R), r(\lambda, \lambda', \psi, \vartheta', \varphi')) \right]
$$

! **and integrated the signal over all sources We modeled the LOS integral**

Indirect detection of γ**-rays:** Φ_{γ} = Φ particle physics \times Φ cosmology **Integrated contribution of EXTRAGALACTIC halos and subhalos Step 3:**

Computing the cosmological γ -ray flux due to DM annihilation in halos...

The γ**-ray sky Galactic and extragalactic: Smooth + subhalos**

PHOTONS in 5 YEAR FERMI-LIKE OBSERVATION

Next step: Compare with Fermi MAP

first 3 months: 205 sources > 0.3 GeV > 10 σ

Many sources have been identified as pulsars or other astrophysical objects

Is any of the unidentified FERMI sources a DM subhalo?

Figure from D. Smith @ TANGO in Paris

Is it so "easy"?

Assume it is! We needs the energy spectrum! Will there be enough photons to get it? Will it be enough?

> $\overline{1}$ Is there any othen clue that may point towards or exclude + the DM hypothesis?

Indirect detection of dark matter:

The role of substructures

Subhalo population of our galaxy: -Prospects for detection in γ-rays -Multi-wavelength analysis

The multiwavelength/multimessenger/multitarget approach

Φ = ParticlePhysics x Cosmology/Astrophysics x Transport

Slide: courtesy of M. Pato

 $H.E.S.S.$

[HESS 2009]

Systematic error

Broken power-law fit

H.E.S.S. - low-energy analysis

Systematic error - low-energy analysis

 $10³$

Energy (GeV)

which is good independently on the nature of the excesses (that is probably astrophysical)

Possible Indirect Detection trigger:cosmic ray data

The radio sky GC, no subhalos

Assume a magnetic field

7.2 mG 7.2 mG(r/0.04pc)-2 0.04pc<r<3.38pc $1 \mu G$ r>3.38pc r<0.04pc

Compute synchtrotron power à la Bertone 2008

$$
n_{e_{\pm}}(\bar{x}, E_{e_{\pm}}) = \frac{\sigma v}{2m_{DM}^2} \rho_{DM}^2(\bar{x}) \frac{N_{e_{\pm}}(>E_{e_{\pm}})}{b_{syn}(\bar{x}, E_{e_{\pm}})}
$$

$$
v \frac{dW_{syn}}{dv} = \frac{\sigma v}{2m_{DM}^2} \int_{\Omega} d\Omega \int_{\Omega} ds \rho_{DM}^2(\bar{x}) E(\bar{x}, v) \frac{N_{e_{\pm}}(>E_{e_{\pm}})}{2}
$$

The antimatter sky

Compute the number density à la Delahaye 2008

$$
n_{_{CR}}(\mathbf{t}, \bar{\mathbf{x}}, \mathbf{E}_{_{CR}}) = \frac{d^2 N_{_{CR}}}{dV dE_{_{CR}}}
$$

electrons and positrons

$$
\frac{\partial n_{e+}}{\partial t} - K_{e+}(E_{e+})\nabla^2 n_{e+} - \frac{\partial}{\partial E_{e+}}(b(E_{e+})n_{e+}) = Q_{e+}(\vec{x}, E_{e+})
$$
\n
$$
\frac{\partial n_{\overline{p}}}{\partial t} - K_{\overline{p}}(T_{\overline{p}})\nabla^2 n_{\overline{p}} - \frac{\partial}{\partial z}(sgn(z)V_{c}n_{\overline{p}}) = Q_{\overline{p}}(\vec{x}, T_{\overline{p}}) - 2h\delta_{D}(z)\Gamma_{ann}^{p\overline{p}}(T_{\overline{p}})n_{\overline{p}}
$$
\n
$$
\frac{\partial n_{\overline{p}}}{\partial t} - K_{\overline{p}}(T_{\overline{p}})\nabla^2 n_{\overline{p}} - \frac{\partial}{\partial z}(sgn(z)V_{c}n_{\overline{p}}) = Q_{\overline{p}}(\vec{x}, T_{\overline{p}}) - 2h\delta_{D}(z)\Gamma_{ann}^{p\overline{p}}(T_{\overline{p}})n_{\overline{p}}
$$

Compute fluxes and boosts à la Lavalle 2008

!

$$
\varphi_{\text{CR,sm}}(E_{\text{CR}}) \propto <\sigma v> \int_{E_{\text{CR}}}^{\infty} dE \frac{dN_{\text{CR}}}{dE} \int_{\text{diff,zone}} d^3 \vec{x} \left(\frac{\rho_{\text{sm}}(\vec{x})}{\rho_{\text{sun}}}\right)^2 G_{\text{sun}}^{\text{CR}}(\vec{x}, \lambda_{\text{D}})
$$
\n
$$
<\varphi_{\text{CR,cl}}>(E_{\text{CR}}) \propto <\sigma v> N_{\text{cl}} \int_{E_{\text{CR}}}^{\infty} dE \frac{dN_{\text{CR}}}{dE} \int_{\text{diff,zone}} d^3 \vec{x} < \xi>_{\text{M}}(R) \frac{dP_{\text{V}}}{dV}(R) G_{\text{sun}}^{\text{CR}}(\vec{x}, \lambda_{\text{D}}) = N_{\text{tot}}^{\text{sub}} < \varphi_{\text{sub}}>
$$

The antimatter sky

Compute fluxes and boosts à la Lavalle 2008

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\varphi_{\text{CR,sm}}(E_{\text{CR}}) \propto <\sigma v> \int_{E_{\text{CR}}}^{\infty} dE \frac{dN_{\text{CR}}}{dE} \int_{diff,zone} d^3\vec{x} \left(\frac{\rho_{\text{sm}}(\vec{x})}{\rho_{\text{sun}}}\right)^2 G_{\text{sun}}^{\text{CR}}(\vec{x},\lambda_{\text{D}})
$$

$$
\langle \varphi_{CR,cl} \rangle (E_{CR}) \propto \langle \sigma v \rangle N_{cl} \int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff. zone} d^3 \vec{x} \langle \xi \rangle_{M} (R) \frac{dP_{V}}{dV} (R) G_{sun}^{CR} (\vec{x}, \lambda_{D}) = N_{tot}^{sub} \langle \varphi_{sub} \rangle
$$

LP, Lavalle, Bertone & Branchini 2009

The antimatter sky

Compute fluxes and boosts à la Lavalle 2008

$$
\varphi_{\text{CR,sm}}(E_{\text{CR}}) \propto <\sigma v> \int_{E_{\text{CR}}}^{\infty} dE \frac{dN_{\text{CR}}}{dE} \int_{diff.zone} d^3 \bar{x} \left(\frac{\rho_{\text{sm}}(\bar{x})}{\rho_{\text{sun}}} \right)^2 G_{\text{sun}}^{\text{CR}}(\bar{x}, \lambda_{\text{D}})
$$

$$
\langle \varphi_{CR,cl} \rangle \langle E_{CR} \rangle \propto \langle \sigma v \rangle \langle N_{cl} \int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff. zone} d^3\vec{x} \langle \xi \rangle_{M} (R) \frac{dP_{V}}{dV} (R) G_{sun}^{CR}(\vec{x}, \lambda_{D}) = N_{tot}^{sub} \langle \varphi_{sub} \rangle
$$

Annihilation proceeds through the exchange of massive vector bosons $\sum_{i=1}^{n}$ when the two DM particles get close and are slow. It mimics an attractive force which arises $\sigma_{\text{ann}}\mathsf{v} = \mathsf{S}(\sigma_{\text{ann}}\mathsf{v})_{\text{thermal}}$

Particle Physics BF: Sommerfeld enhancement

Sommerfeld effect produces a local enhancement of the annihilation cross-section which depends on the DM velocity and mass, and does not touch the thermal value

$$
\frac{1}{m_{DM}}\frac{d^2\psi(r)}{dr^2}=-m_{DM}\beta^2\psi(r)-\frac{\alpha}{r}e^{-m_Vr}
$$

Particle Physics BF: Sommerfeld enhancement

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$$
\frac{1}{m_{DM}}\frac{d^2\psi(r)}{dr^2}=-m_{DM}\beta^2\psi(r)-\frac{\alpha}{r}e^{-m_Vr}
$$

Particle Physics BF: Sommerfeld enhancement

mass dependence (resonance) Sommerfeld effect produces a local enhancement of the annihilation cross-section which depends on the DM velocity and mass, and does not touch the thermal value

Particle Physics BF and astrophysics BF: Sommerfeld enhancement and subhalos

Dwarf galaxies and galactic subhalos have low velocity dispersions, hence the Sommerfeld enhancement should be convolved with the sub-subhalo contribution

$$
\sigma_{\text{ann}} \mathsf{v} = (\sigma_{\text{ann}} \mathsf{v})_{\text{thermal}} \mathsf{S}(\beta(r))
$$

$$
\Phi = T(E,d)\frac{(\sigma_{\text{ann}}v)_{\text{thermal}}}{2m_{\chi}^2}\int\limits_{E_o}^{m_{\chi}}\sum_f \frac{dN_f^{\gamma}}{dE_{\gamma}}BR_f \int\limits_{\text{l.o.s}} d\lambda S(\beta(M,r))\rho_{\text{DM}}^2\Big(M,c(M,r),R\Big)
$$

 \mathbf{r} This holds for all annihilation products **We can perform a multi-wavelength analysis to constrain models APPLYING BOOSTS TO BOTH** Φ_{PP} **AND** Φ_{COSMO}

Pato, LP, Bertone 2009

Pato, LP, Bertone 2009

Conclusions

Detection of individual structures

In the best case scenario high mass halos are "detectable" Result poorly dependent on small mass extrapolation

Multi-wavelength constraints

Coherent prediction of signals from all annihilation products is now necessary in order to constrain (or discover) particle physics and cosmological DM models

Upcoming data

This is more than ever important in these years when data from satellites, Cherenkov Telescopes, accelerators and Direct detection are about to allow an unprecedented insight on the DM puzzle