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# HEP Theory

(part 2)



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**Key question for the future developments of HEP:**  
**Why don't we see the new physics we expected to be present around the TeV scale ?**

- **Is the mass scale beyond the LHC reach ?**
- **Is the mass scale within LHC's reach, but final states are elusive to the direct search ?**

These two scenarios are a priori equally likely, but they impact in different ways the future of HEP, and thus the assessment of the physics potential of possible future facilities

Readiness to address both scenarios is the best hedge for the field:

- *precision*
- *sensitivity (to elusive signatures)*
- *extended energy/mass reach*



## **Remark**

the discussion of the **future** in HEP must start from the understanding that there is no experiment/facility, proposed or conceivable, in the lab or in space, accelerator or non-accelerator driven, which can **guarantee discoveries** beyond the SM, and **answers** to the big questions of the field

The physics potential (the “case”) of a future facility for HEP should be weighed against criteria such as:

(1) the **guaranteed deliverables:**

- knowledge that will be acquired independently of possible discoveries (*the value of “measurements”*)

(2) the **exploration potential:**

- target broad and well justified BSM scenarios ... *but guarantee sensitivity to more exotic options*
- exploit both direct (large  $Q^2$ ) and indirect (precision) probes

(3) the potential to provide conclusive **yes/no answers** to relevant, broad questions.

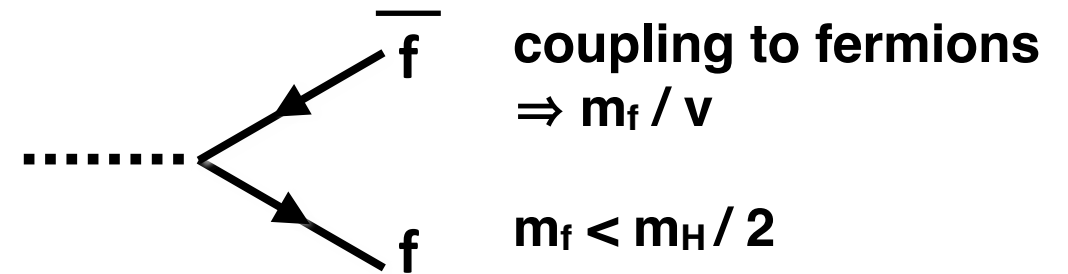
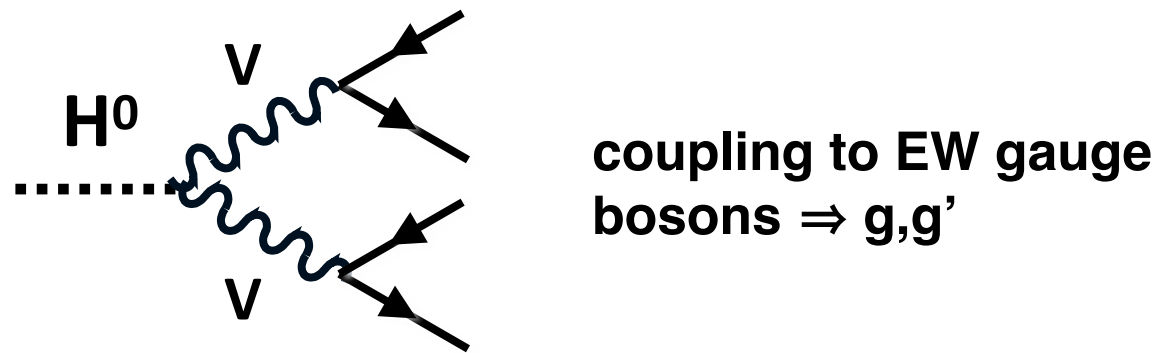


# What we want from a future collider

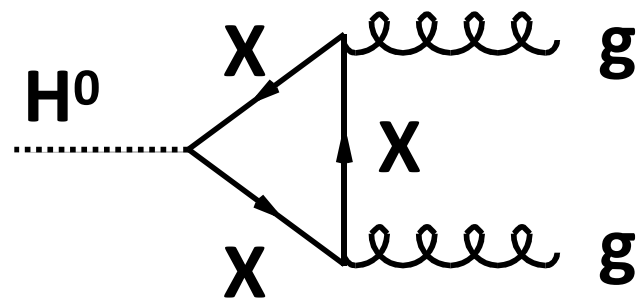
- Guaranteed deliverables:
  - study of Higgs and top quark properties, and exploration of EWSB phenomena, with the best possible **precision and sensitivity**
- Exploration potential:
  - exploit both direct (large  $Q^2$ ) and indirect (precision) probes
  - **enhanced mass reach for direct exploration**
    - *E.g. match the mass scales for new physics that could be exposed via indirect precision measurements in the EW and Higgs sector*
- Provide firm Yes/No answers to questions like:
  - is there a TeV-scale solution to the hierarchy problem?
  - is DM a thermal WIMP?
  - could the cosmological EW phase transition have been 1st order?
  - could baryogenesis have taken place during the EW phase transition?
  - could neutrino masses have their origin at the TeV scale?
  - ...

# Higgs observables: decay BRs

## Tree-level couplings

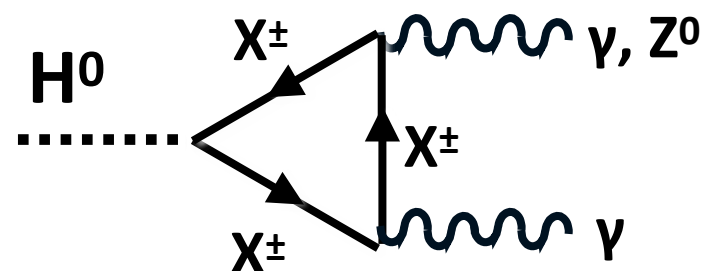


## Loop-level couplings



$X_{SM} = t, b, c$

$X_{BSM} = T, \text{stop}, \dots?$

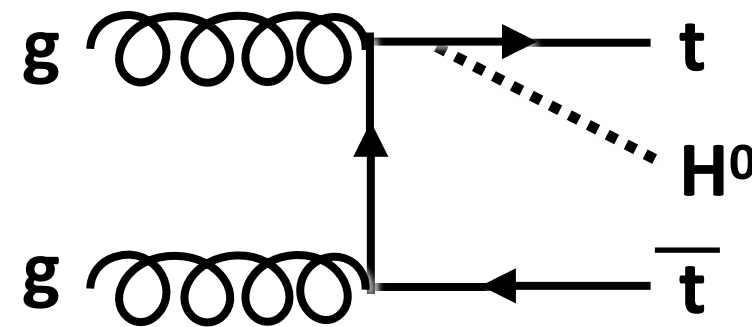
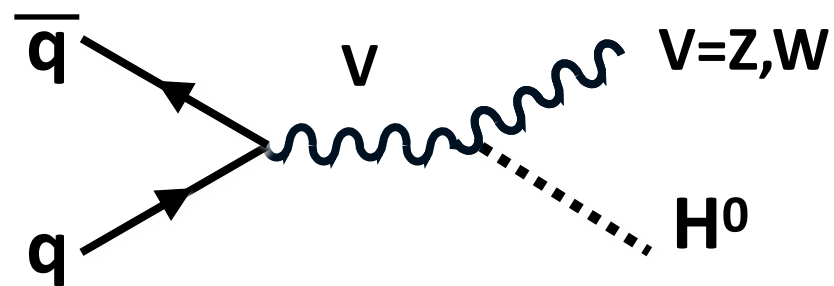
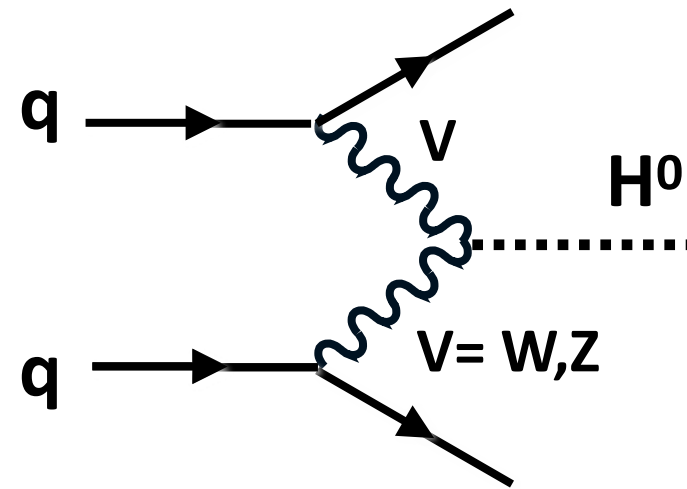
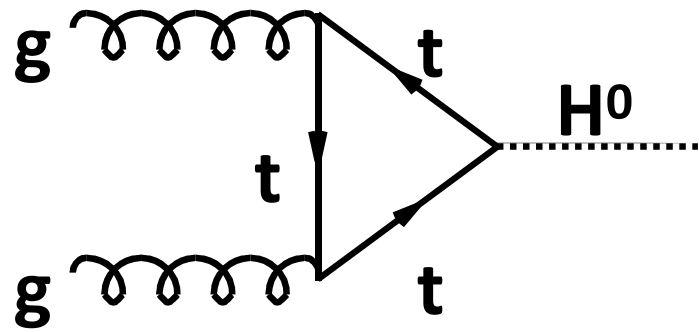


$X_{SM} = t, W^\pm$

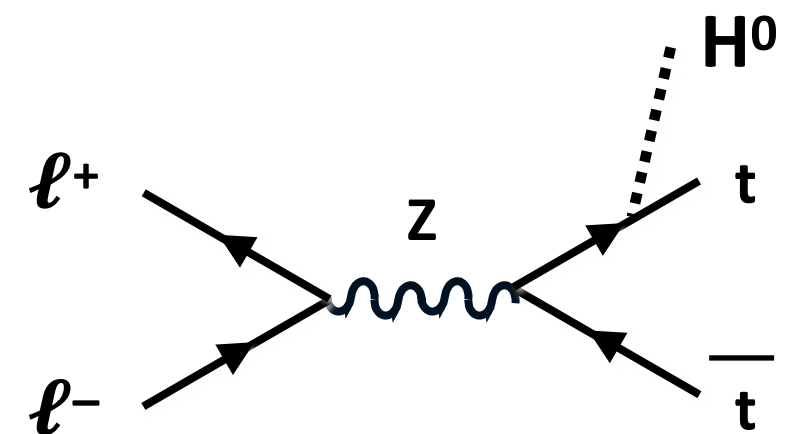
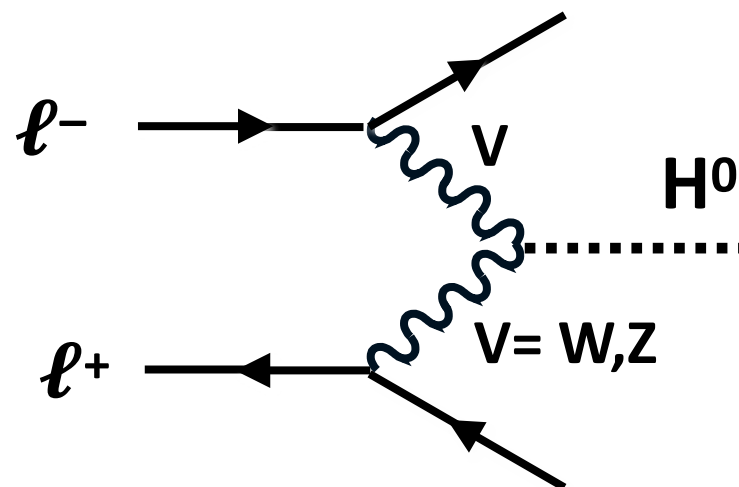
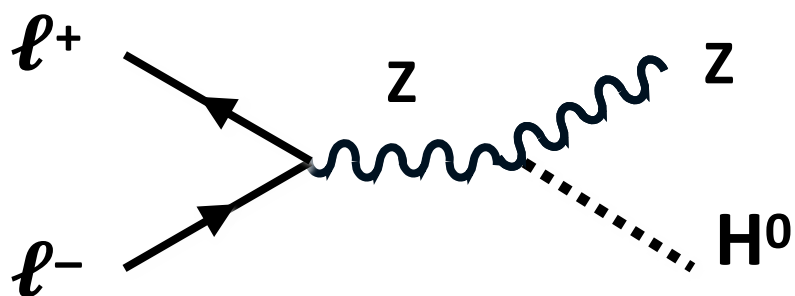
$X_{BSM} = T, \text{stop}, \text{chargino}, \dots?$

# Higgs observables: production rates

## Hadronic collisions



## Laptonic collisions





# Sensitivity of various Higgs couplings to examples of beyond-the-SM phenomena

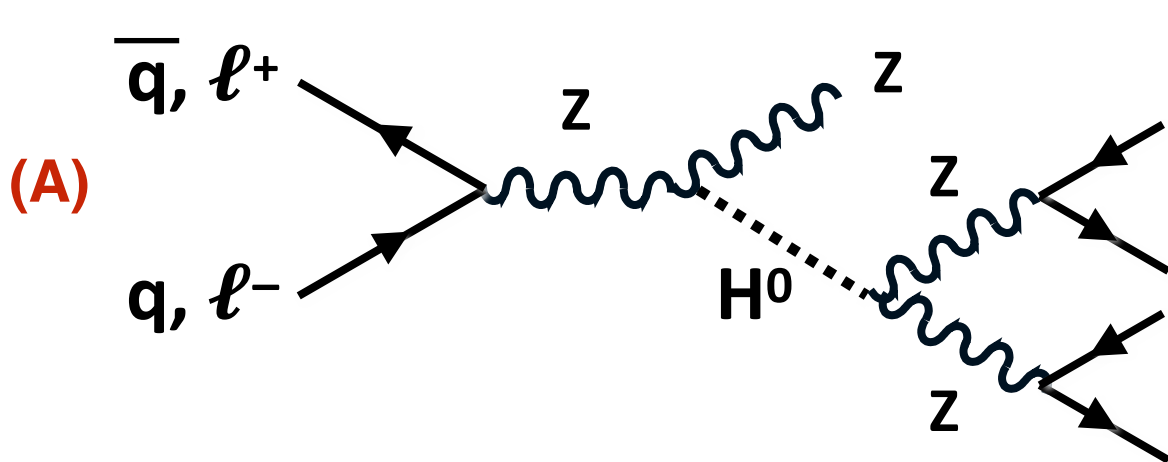
*arXiv:1310.8361*

Model	$\kappa_V$	$\kappa_b$	$\kappa_\gamma$
Singlet Mixing	$\sim 6\%$	$\sim 6\%$	$\sim 6\%$
2HDM	$\sim 1\%$	$\sim 10\%$	$\sim 1\%$
Decoupling MSSM	$\sim -0.0013\%$	$\sim 1.6\%$	$\sim -0.4\%$
Composite	$\sim -3\%$	$\sim -(3 - 9)\%$	$\sim -9\%$
Top Partner	$\sim -2\%$	$\sim -2\%$	$\sim +1\%$

**=> the goal should be (sub)percent precision!**

# Extracting couplings from measurements

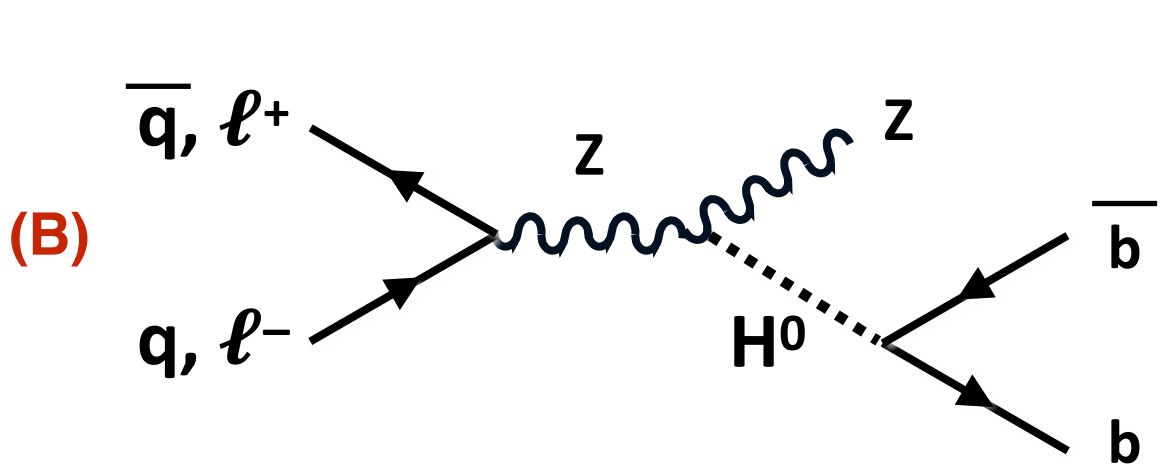
## Example



$$\sigma(pp\ell e \rightarrow ZH[\rightarrow ZZ^*]) \propto g_{HZZ}^2 \times \frac{g_{HZZ}^2}{\Gamma_H}$$

1 measurement, 2 parameters!

$B(H \rightarrow ZZ^*)$



$$\sigma(pp\ell e \rightarrow ZH[\rightarrow b\bar{b}]) \propto g_{HZZ}^2 \times \frac{g_{Hbb}^2}{\Gamma_H}$$

1 new measurement, but  
1 more parameter...

$B(H \rightarrow b\bar{b})$

... little progress, except we now know

$$\frac{g_{HZZ}^2}{g_{Hbb}^2} = \frac{\sigma_A}{\sigma_B}$$

Overall constraint:  $\sum_X B(H \rightarrow X) = 1$

Therefore:

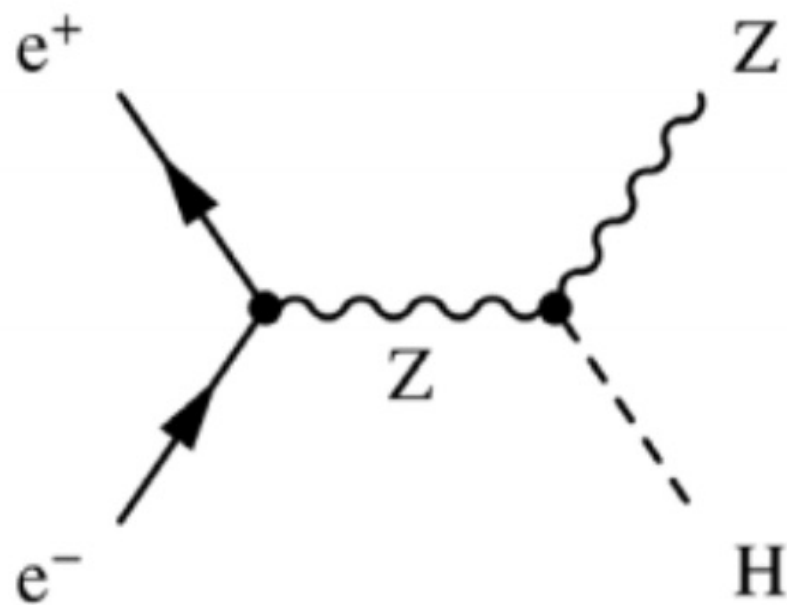
$$\sum_X \sigma \left( \begin{array}{c} \bar{q}, \ell^+ \\ q, \ell^- \end{array} \rightarrow \begin{array}{c} z \\ z \\ H^0 \end{array} \rightarrow X \right) = \sigma(ZH) \propto g_{HZZ}^2$$

How can we hope to detect ALL possible decays of the Higgs boson??

If the goal is to test its properties, we cannot make assumptions, and must be open to possible unexpected decays, possibly invisible, like  $H \rightarrow$  dark matter...

An  $\ell^+\ell^-$  collider provides the solution ....

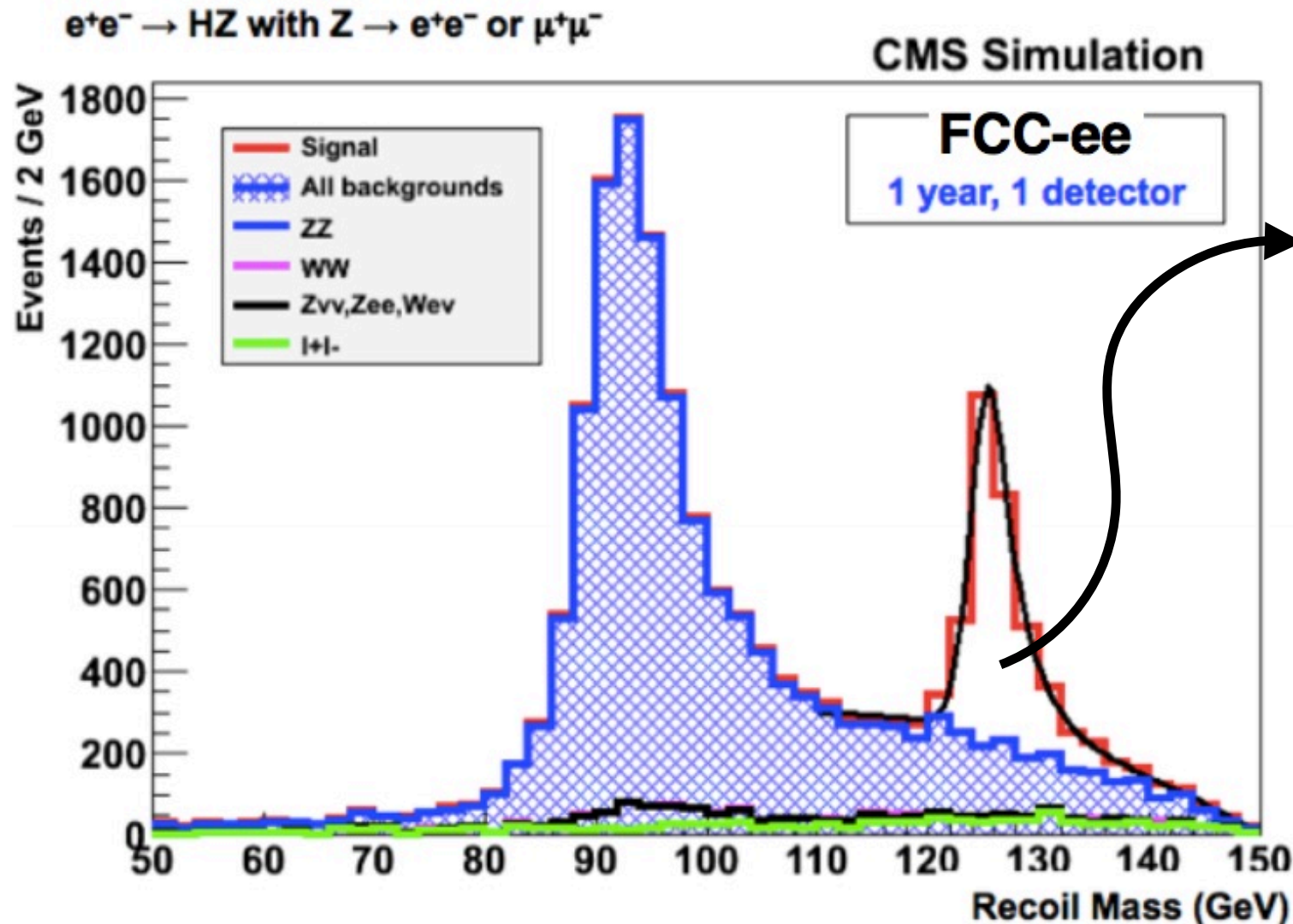




$$p(H) = p(e^-e^+) - p(Z)$$

$$\Rightarrow [ p(e^-e^+) - p(Z) ]^2 \text{ peaks at } m^2(H)$$

reconstruct Higgs events independently of the Higgs decay mode!

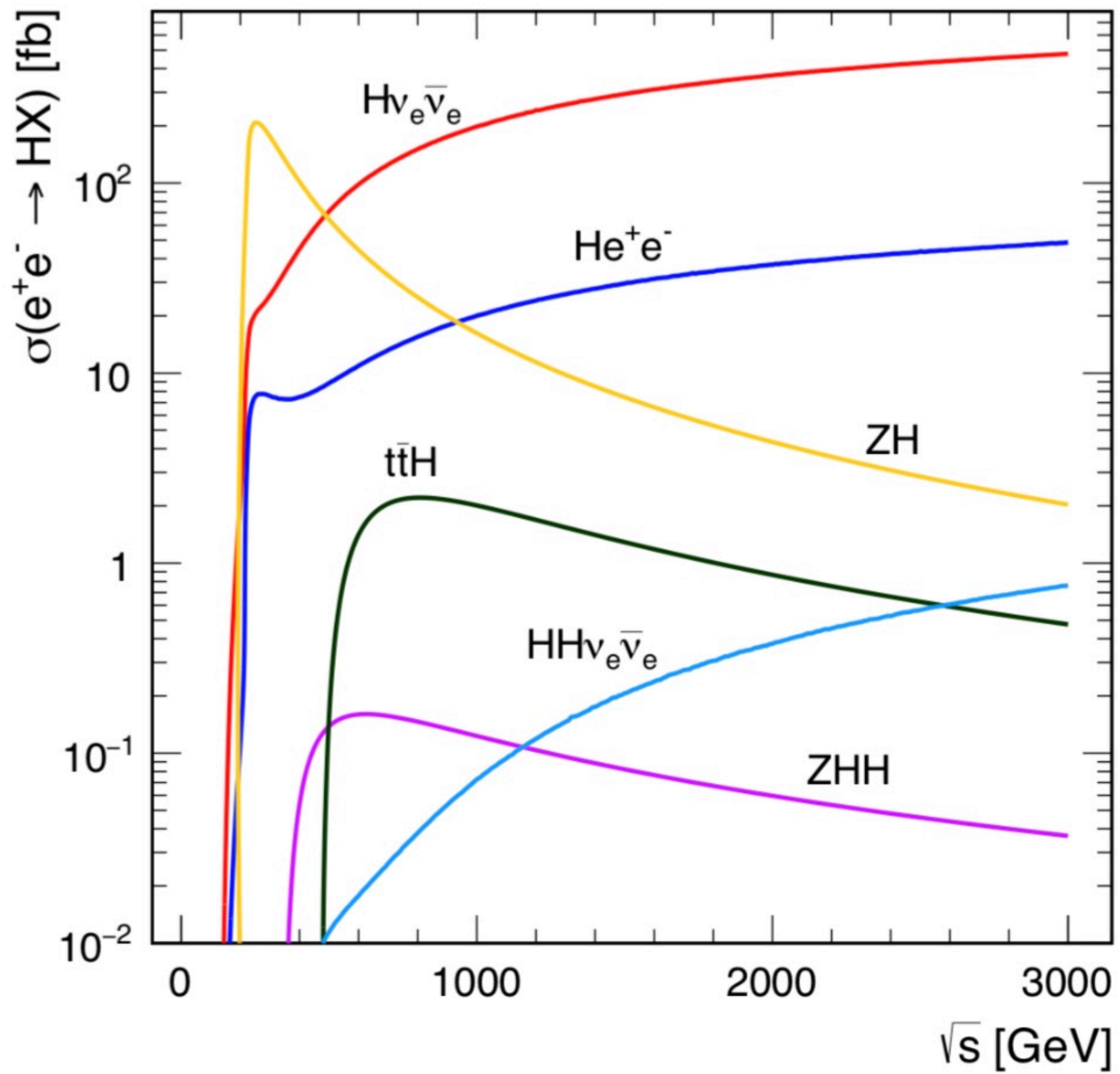


$$N(ZH) \propto \sigma(ZH) \propto g_{HZZ}^2$$

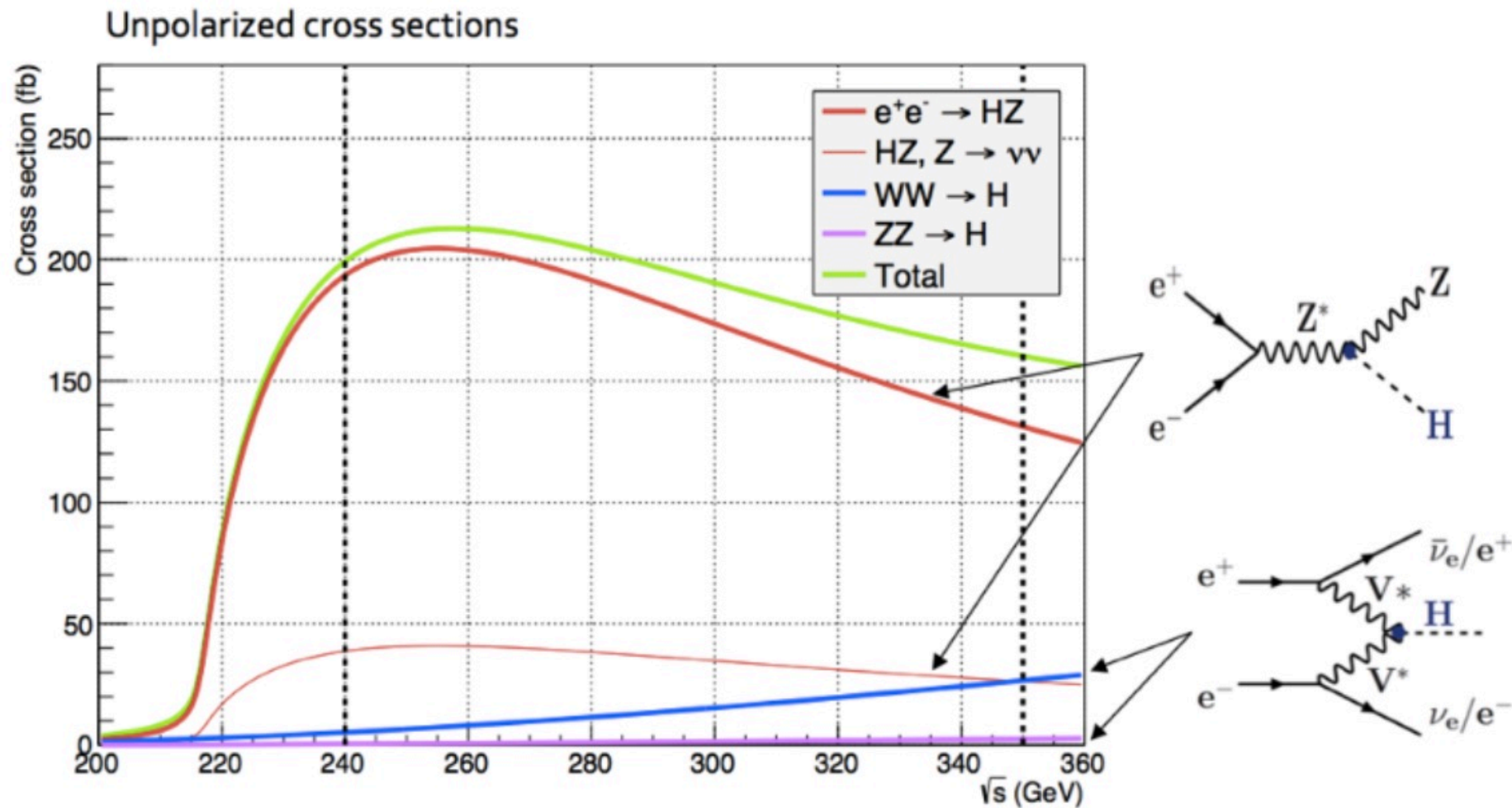
$$N(ZH[\rightarrow ZZ]) \propto \sigma(ZH) \times BR(H \rightarrow ZZ) \propto g_{HZZ}^2 \times g_{HZZ}^2 / \Gamma(H)$$

$\Rightarrow$  absolute measurement of width and couplings

$$m_{\text{recoil}} = \sqrt{ [ p(e^-e^+) - p(Z) ]^2 }$$



# FCC-ee



	FCC-ee 240 GeV	FCC-ee 350 GeV
<b>Total Integrated Luminosity (ab<sup>-1</sup>)</b>	<b>5</b>	<b>1.5</b>
<b># Higgs bosons from <math>e^+e^- \rightarrow HZ</math></b>	<b>1,000,000</b>	<b>200,000</b>
<b># Higgs bosons form fusion process</b>	<b>25,000</b>	<b>40,000</b>



# Higgs couplings: beyond the HL-LHC

Collider	HL-LHC	HL-LHC update	ILC <sub>250</sub>	CLIC <sub>380</sub>	LEP3 <sub>240</sub>	CEPC <sub>250</sub>	FCC-ee <sub>240+365</sub>		
Lumi (ab <sup>-1</sup> )	3	3	2	0.5	3	5	5 <sub>240</sub>	+1.5 <sub>365</sub>	+ HL-LHC
Years	25	25	15	7	6	7	3	+4	
$\delta\Gamma_H/\Gamma_H$ (%)	SM	50	3.6	6.3	3.6	2.6	2.7	<b>1.3</b>	1.1
$\delta g_{HZZ}/g_{HZZ}$ (%)	3.5	<b>1.5</b>	0.3	0.40	0.32	0.25	0.20	<b>0.17</b>	<b>0.16</b>
$\delta g_{HWW}/g_{HWW}$ (%)	3.5	<b>1.7</b>	1.7	0.8	1.7	1.2	1.3	<b>0.43</b>	<b>0.40</b>
$\delta g_{Hbb}/g_{Hbb}$ (%)	8.2	<b>3.7</b>	1.7	1.3	1.8	1.3	1.3	<b>0.61</b>	<b>0.56</b>
$\delta g_{Hcc}/g_{Hcc}$ (%)	SM	<b>SM</b>	2.3	4.1	2.3	1.8	1.7	<b>1.21</b>	<b>1.18</b>
$\delta g_{Hgg}/g_{Hgg}$ (%)	3.9	<b>2.5</b>	2.2	2.1	2.1	1.4	1.6	<b>1.01</b>	<b>0.90</b>
$\delta g_{H\tau\tau}/g_{H\tau\tau}$ (%)	6.5	<b>1.9</b>	1.9	2.7	1.9	1.4	1.4	<b>0.74</b>	<b>0.67</b>
$\delta g_{H\mu\mu}/g_{H\mu\mu}$ (%)	5.0	<b>4.3</b>	14.1	n.a.	12	6.2	10.1	<b>9.0</b>	<b>3.8</b>
$\delta g_{H\gamma\gamma}/g_{H\gamma\gamma}$ (%)	3.6	<b>1.8</b>	6.4	n.a.	6.1	4.7	4.8	<b>3.9</b>	<b>1.3</b>
$\delta g_{Htt}/g_{Htt}$ (%)	4.2	<b>3.4</b>	–	–	–	–	–	–	<b>3.1</b>
BR <sub>EXO</sub> (%)	SM	SM	< 1.7	< 3.0	< 1.6	< 1.2	< 1.2	< <b>1.0</b>	< <b>1.0</b>

**Table 1:** Relative statistical uncertainty on the Higgs boson couplings and total decay width, as expected from the FCC-ee data, and compared to those from HL-LHC and other  $e^+e^-$  colliders exploring the 240-to-380 GeV centre-of-mass energy range. All numbers indicate 68% CL intervals, except for the last line which gives the 95% CL sensitivity on the "exotic" branching fraction, accounting for final states that cannot be tagged as SM decays. The FCC-ee accuracies are subdivided in three categories: the first sub-column give the results of the model-independent fit expected with 5 ab<sup>-1</sup> at 240 GeV, the second sub-column in bold – directly comparable to the other collider fits – includes the additional 1.5 ab<sup>-1</sup> at  $\sqrt{s} = 365$  GeV, and the last sub-column shows the result of the combined fit with HL-LHC. The fit to the HL-LHC projections alone (first column) requires two additional assumptions to be made: here, the branching ratios into  $c\bar{c}$  and into exotic particles are set to their SM values.

\* M. Cepeda, S. Gori, P. J. Ilten, M. Kado, and F. Riva, (conveners), et al, *Higgs Physics at the HL-LHC and HE-LHC*, CERN-LPCC-2018-04, <https://cds.cern.ch/record/2650162>.

# Remarks and **key messages**

- Updated HL-LHC projections bring the coupling sensitivity to the few-% level. They are obtained by extrapolating **current** analysis strategies, and are informed by current experience plus robust assumptions about the performance of the phase-2 upgraded detectors in the high pile-up environment
  - Projections will improve as **new** analyses, allowed by higher statistics, will be considered
1. To significantly improve the expected HL-LHC results, future facilities must push Higgs couplings' precision to the sub-% level
  2. Event rates higher than what ee colliders can provide are needed to reach sub-% measurements of couplings such as  $H\gamma\gamma$ ,  $H\mu\mu$ ,  $HZ\gamma$ ,  $Htt$

# **The unique contributions of a 100 TeV pp collider to Higgs physics**

- Huge Higgs production rates:
  - access (very) rare decay modes
  - push to %-level Higgs self-coupling measurement
  - new opportunities to reduce syst uncertainties (TH & EXP) and push precision
- Large dynamic range for H production (in  $p_T^H$ ,  $m(H+X)$ , ...):
  - new opportunities for reduction of syst uncertainties (TH and EXP)
  - different hierarchy of production processes
  - develop indirect sensitivity to BSM effects at large  $Q^2$ , complementary to that emerging from precision studies (eg *decay BRs*) at  $Q \sim m_H$
- High energy reach
  - direct probes of BSM extensions of Higgs sector
    - SUSY Higgses
    - Higgs decays of heavy resonances
    - Higgs probes of the nature of EW phase transition
    - ...



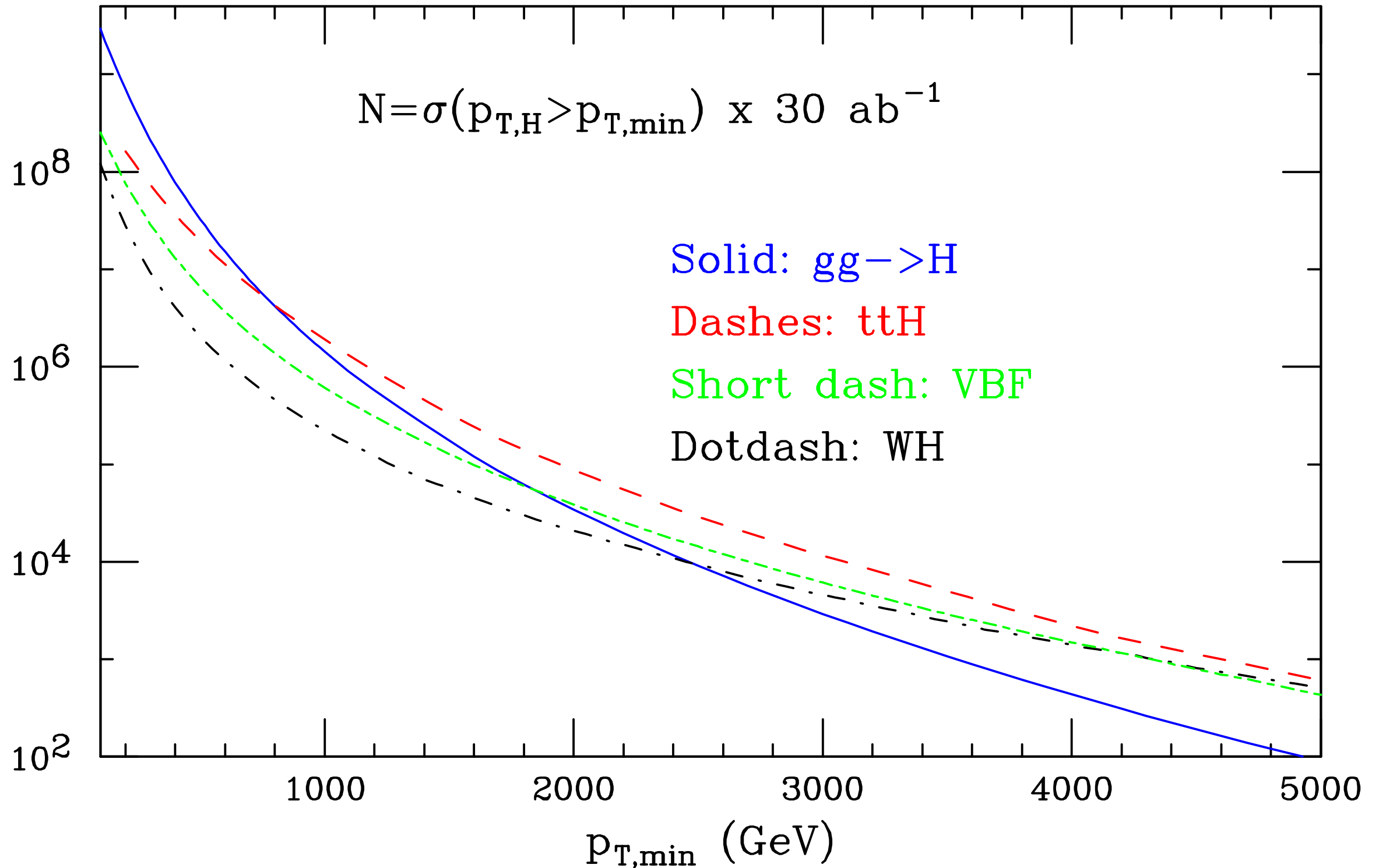
# SM Higgs: event rates in pp@100 TeV

	gg→H	VBF	WH	ZH	ttH	HH
$N_{100}$	24 x 10 <sup>9</sup>	2.1 x 10 <sup>9</sup>	4.6 x 10 <sup>8</sup>	3.3 x 10 <sup>8</sup>	9.6 x 10 <sup>8</sup>	3.6 x 10 <sup>7</sup>
$N_{100}/N_{14}$	180	170	100	110	530	390

$$N_{100} = \sigma_{100\text{TeV}} \times 30 \text{ ab}^{-1}$$

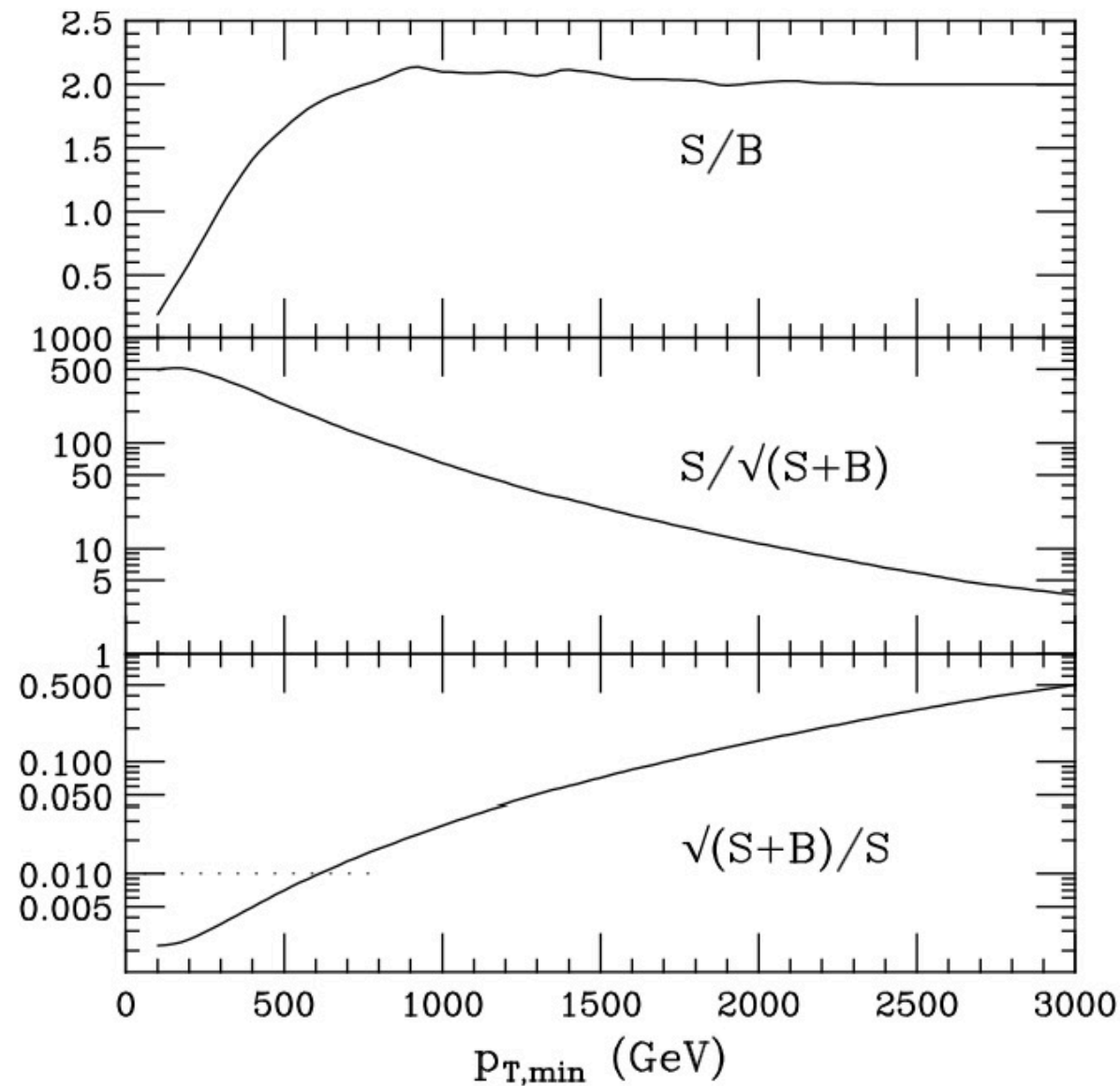
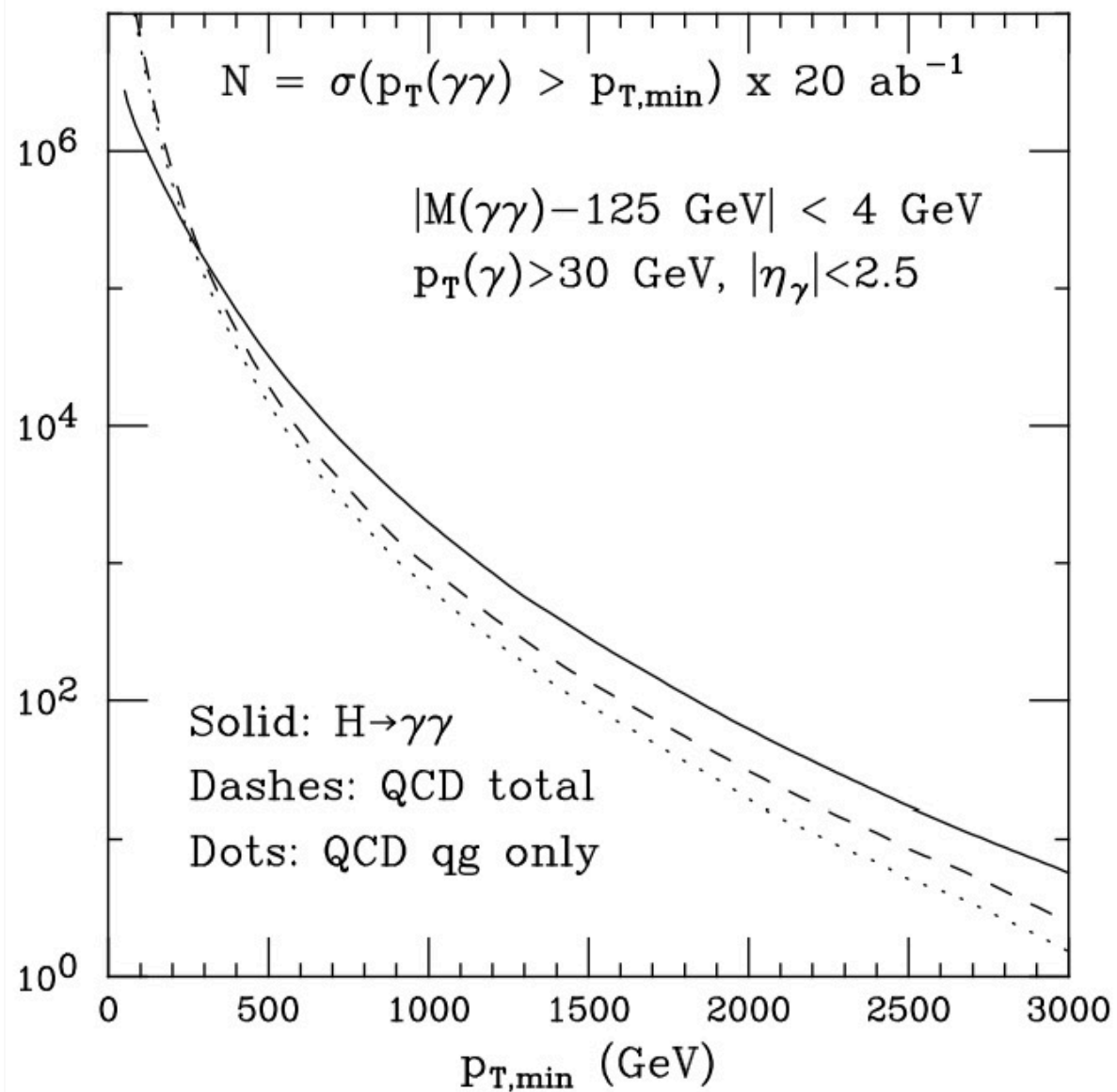
$$N_{14} = \sigma_{14\text{TeV}} \times 3 \text{ ab}^{-1}$$

# H at large $p_T$



- Hierarchy of production channels changes at large  $p_T(H)$ :
  - $\sigma(ttH) > \sigma(gg \rightarrow H)$  above 800 GeV
  - $\sigma(VBF) > \sigma(gg \rightarrow H)$  above 1800 GeV

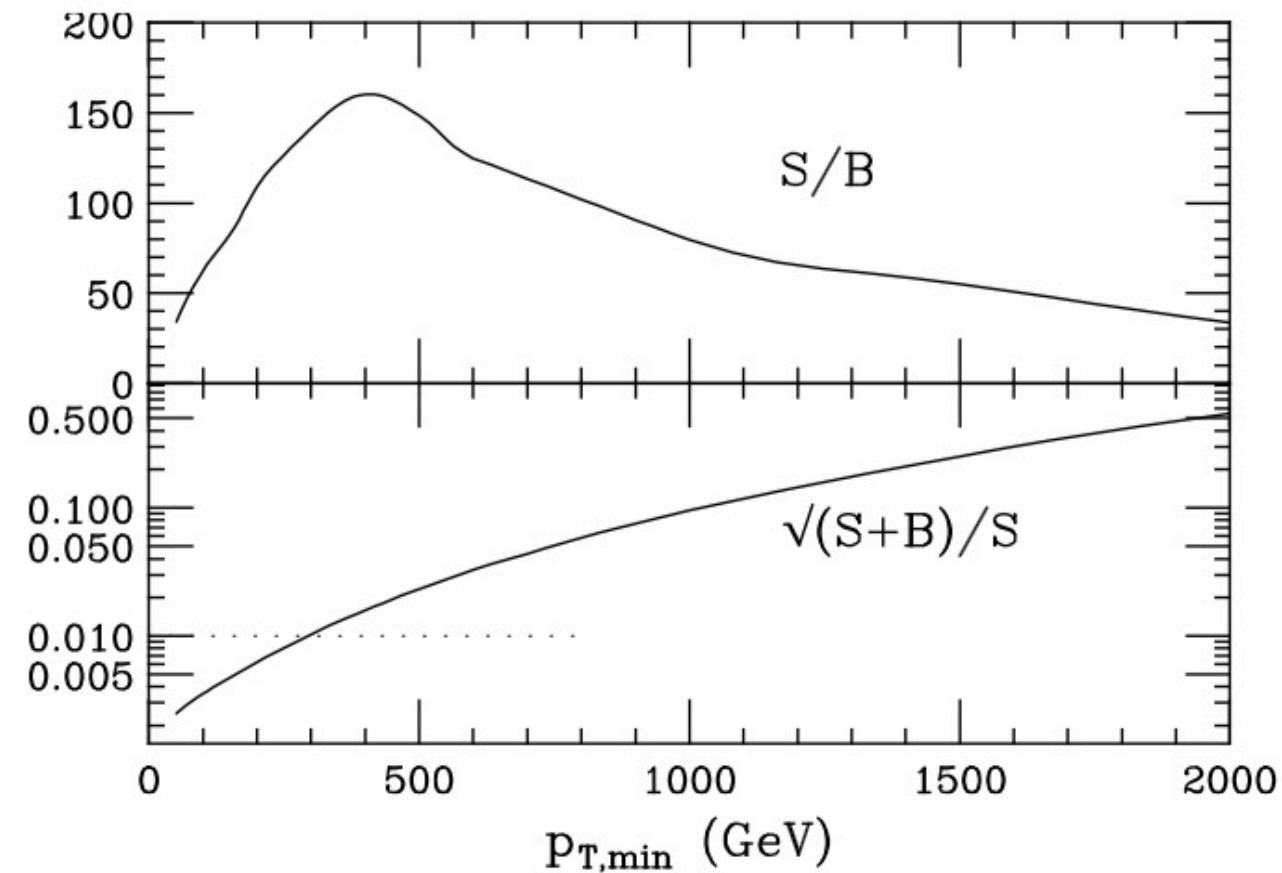
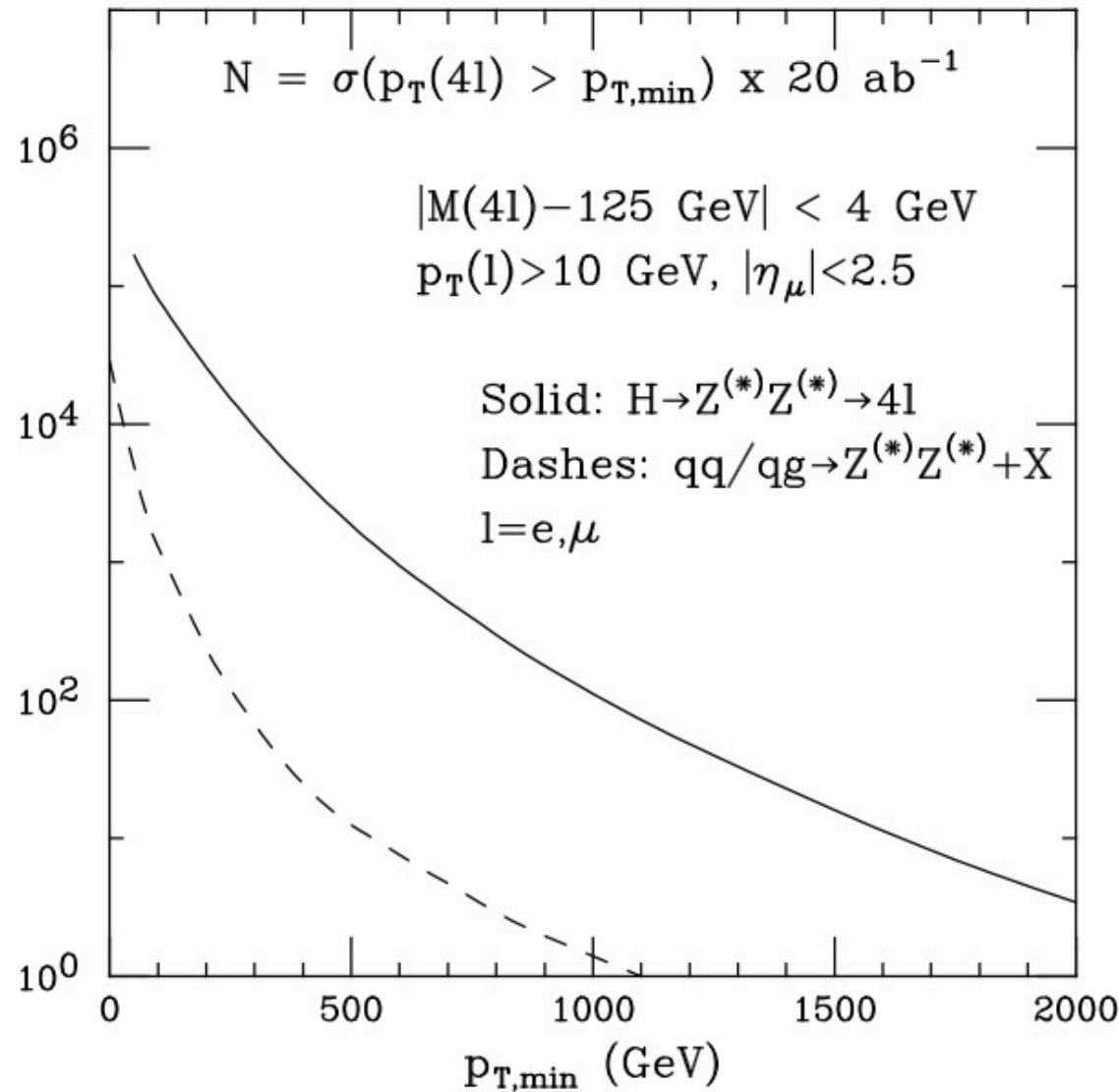
# $gg \rightarrow H \rightarrow \gamma\gamma$ at large $p_T$



- At LHC,  $S/B$  in the  $H \rightarrow \gamma\gamma$  channel is  $O(\text{few } \%)$
- At FCC, for  $p_T(H) > 300 \text{ GeV}$ ,  $S/B \sim 1$
- Potentially accurate probe of the  $H$   $p_T$  spectrum up to large  $p_T$

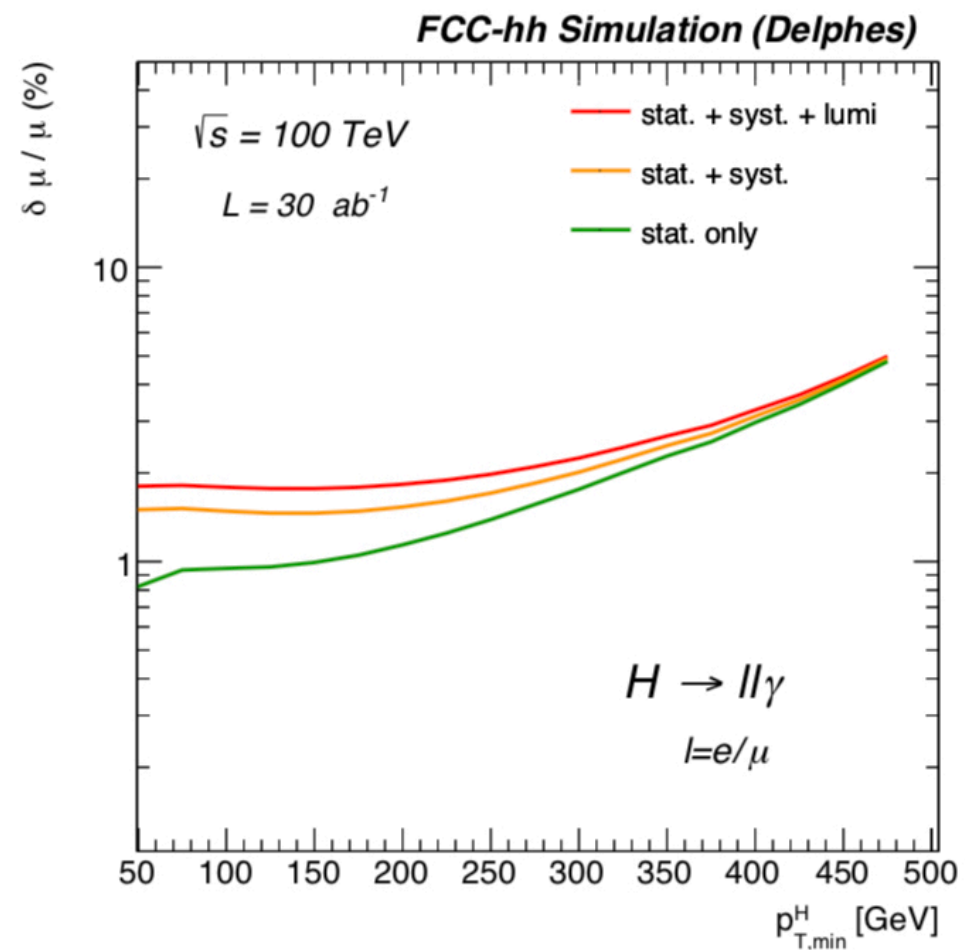
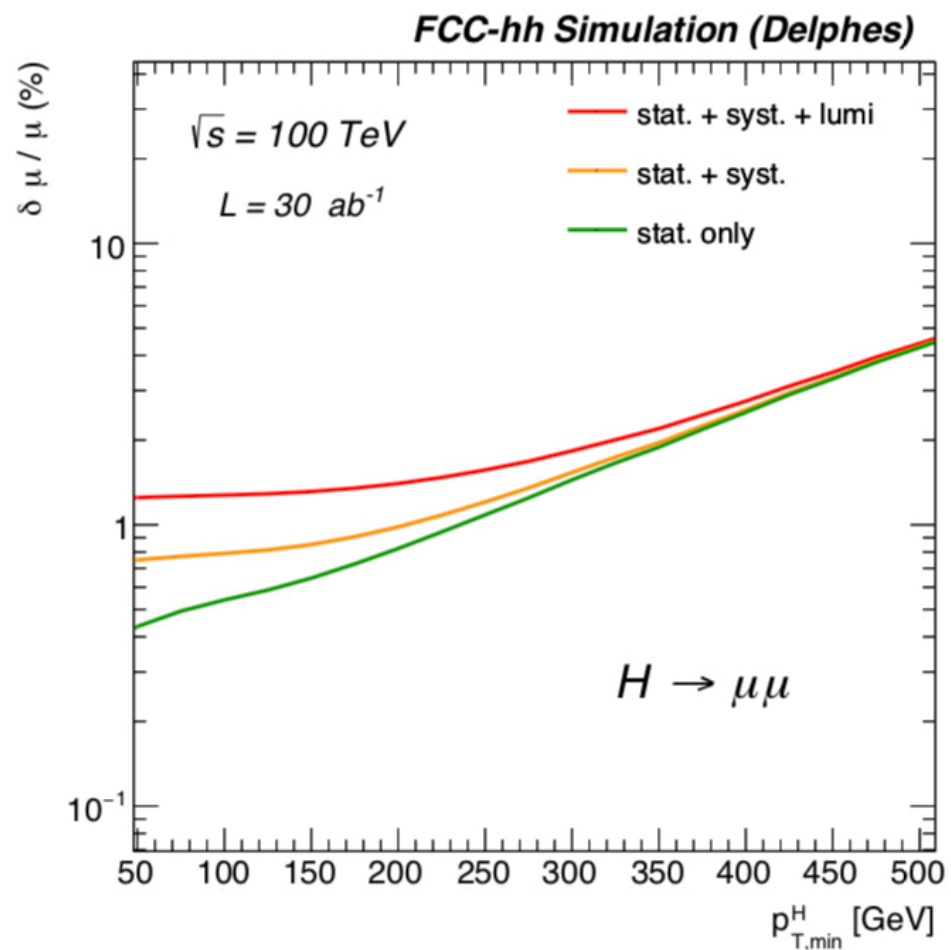
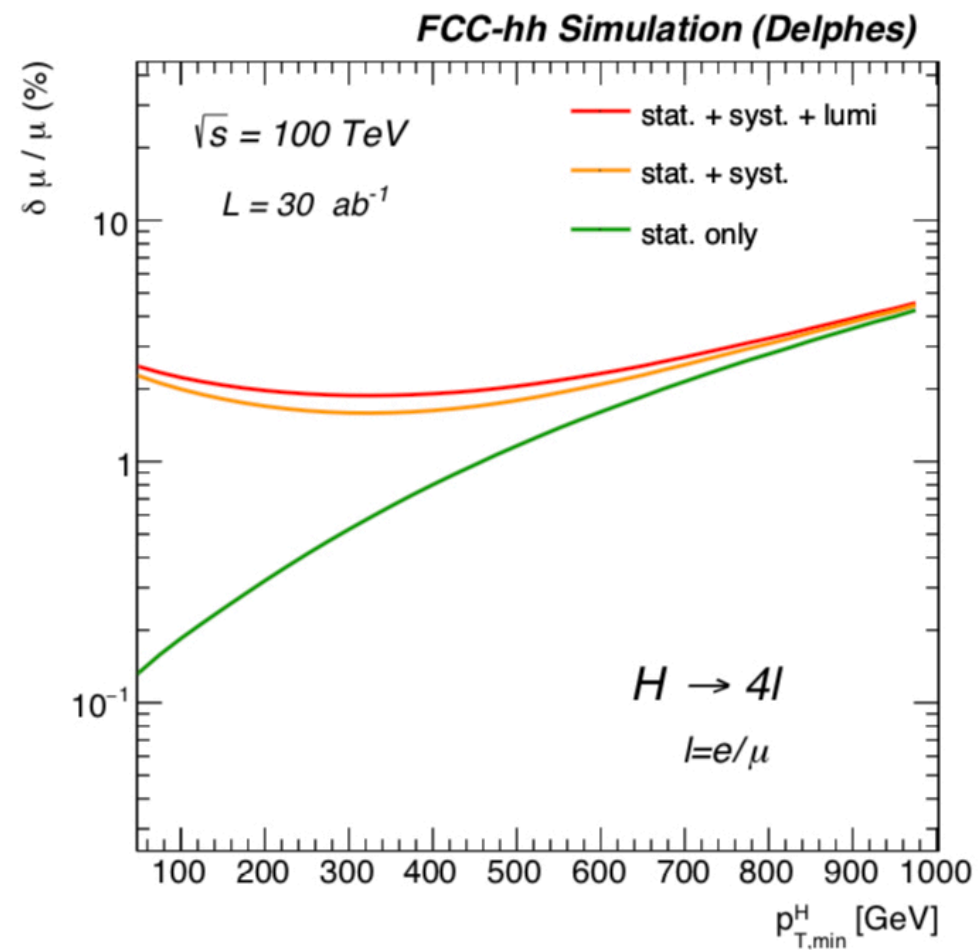
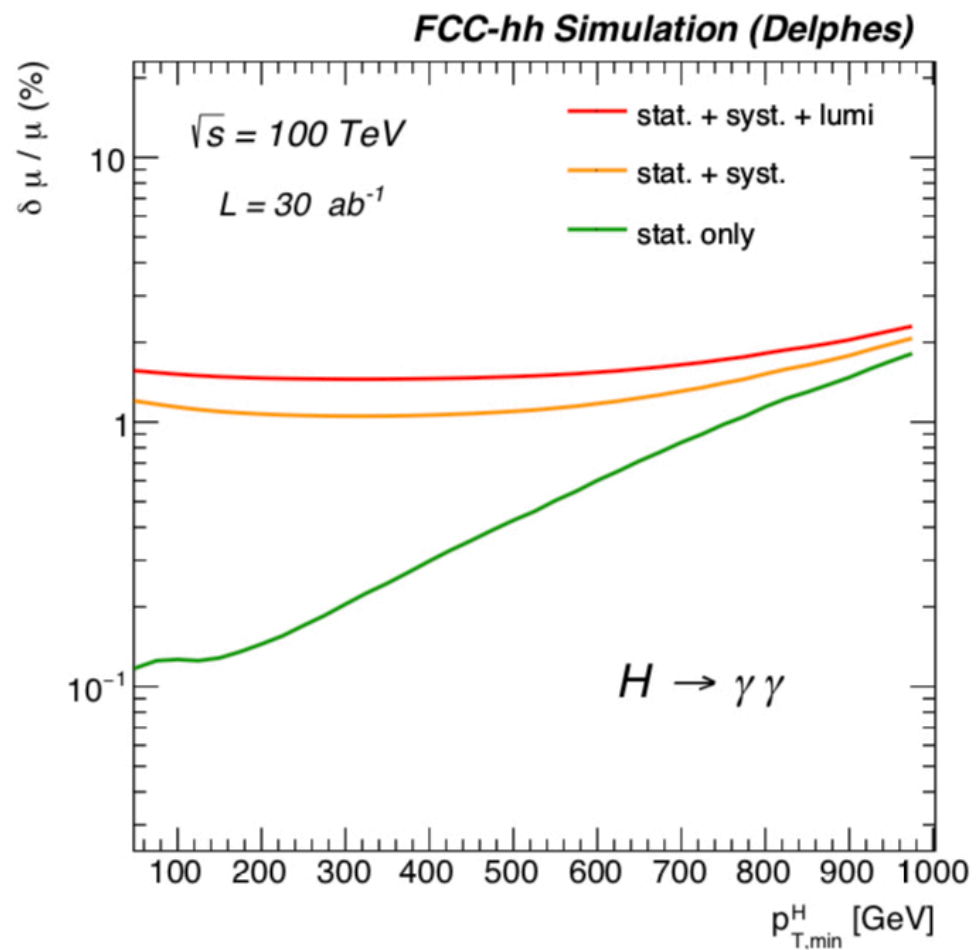
$p_{T,\min}$ (GeV)	$\delta_{\text{stat}}$
100	0.2%
400	0.5%
600	1%
1600	10%

# $gg \rightarrow H \rightarrow ZZ^* \rightarrow 4l$ at large $p_T$



- $S/B \sim 1$  for inclusive production at LHC
- Practically bg-free at large  $p_T$  at 100 TeV, maintaining large rates

$p_{T,\min}$ (GeV)	$\delta_{\text{stat}}$
100	0.3%
300	1%
1000	10%





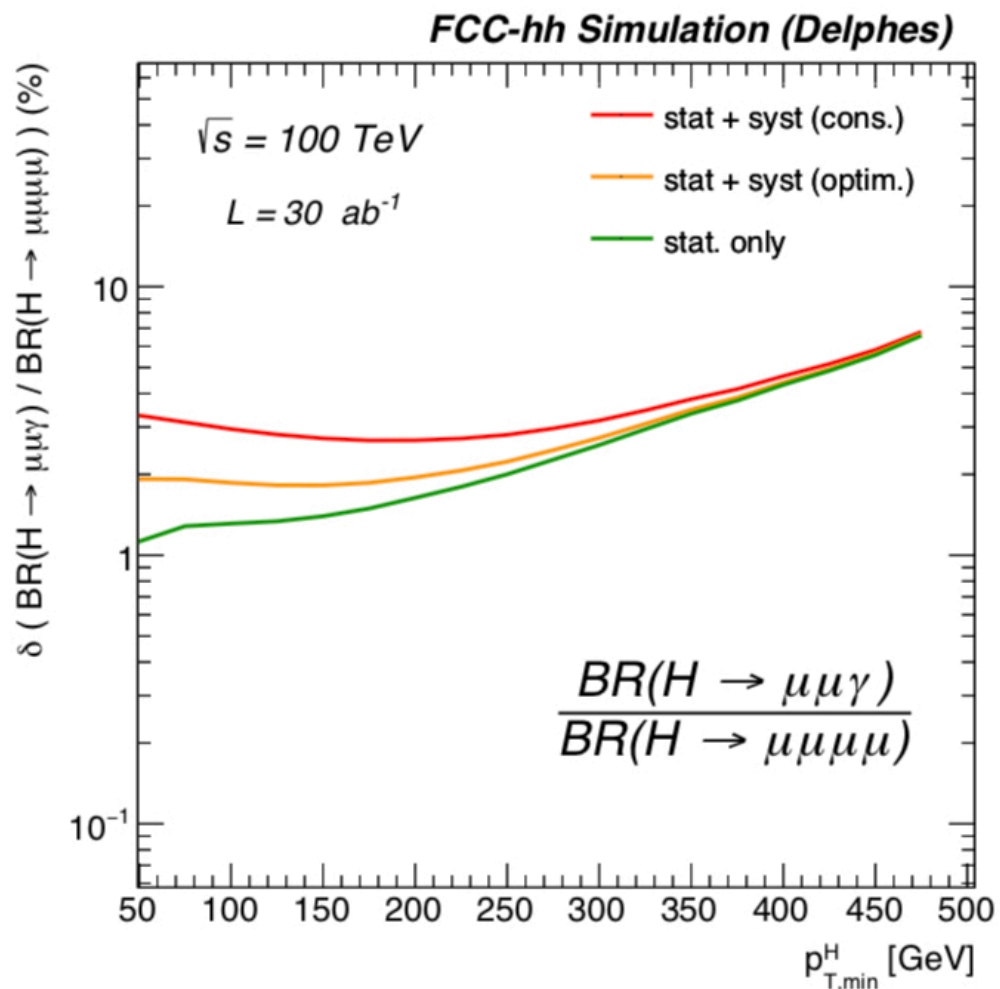
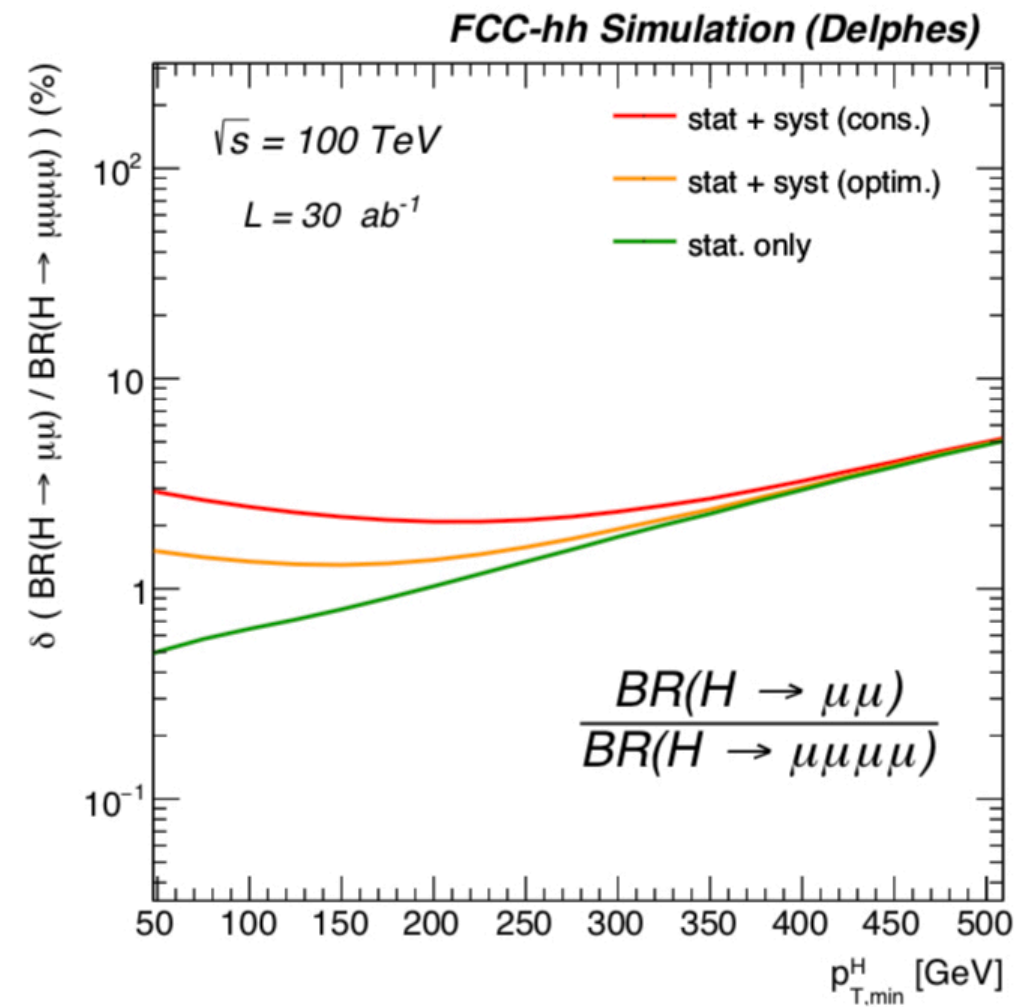
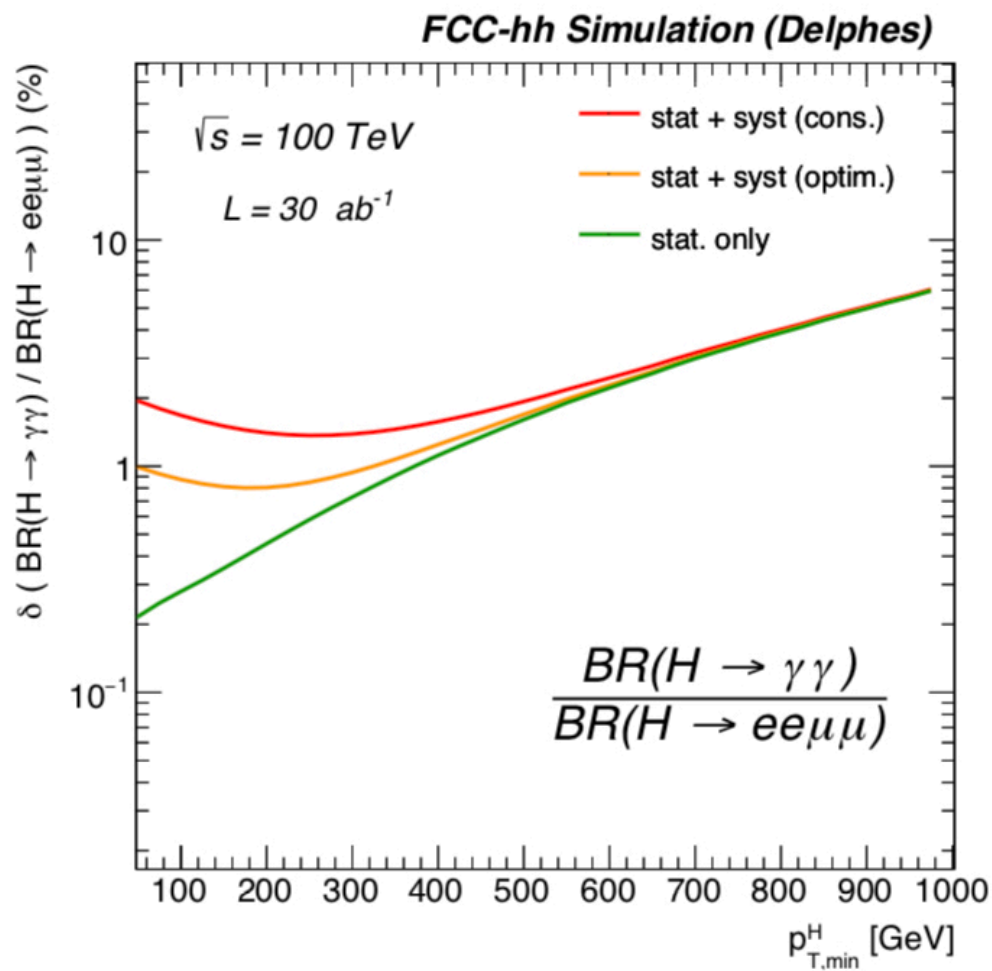


Table 4.4: Target precision for the parameters relative to the measurement of various Higgs decays, ratios thereof, and of the Higgs self-coupling  $\lambda$ . Notice that lagrangian couplings have a precision that is typically half that of what is shown here, since all rates and branching ratios depend quadratically on the couplings.

Observable	Parameter	Precision (stat)	Precision (stat+syst+lumi)
$\mu = \sigma(\text{H}) \times \text{B}(\text{H} \rightarrow \gamma\gamma)$	$\delta\mu/\mu$	0.1%	1.45%
$\mu = \sigma(\text{H}) \times \text{B}(\text{H} \rightarrow \mu\mu)$	$\delta\mu/\mu$	0.28%	1.22%
$\mu = \sigma(\text{H}) \times \text{B}(\text{H} \rightarrow 4\mu)$	$\delta\mu/\mu$	0.18%	1.85%
$\mu = \sigma(\text{H}) \times \text{B}(\text{H} \rightarrow \gamma\mu\mu)$	$\delta\mu/\mu$	0.55%	1.61%
$\mu = \sigma(\text{HH}) \times \text{B}(\text{H} \rightarrow \gamma\gamma) \text{B}(\text{H} \rightarrow \text{b}\bar{\text{b}})$	$\delta\lambda/\lambda$	5%	7.0%
$R = \text{B}(\text{H} \rightarrow \mu\mu) / \text{B}(\text{H} \rightarrow 4\mu)$	$\delta R/R$	0.33%	1.3%
$R = \text{B}(\text{H} \rightarrow \gamma\gamma) / \text{B}(\text{H} \rightarrow 2\text{e}2\mu)$	$\delta R/R$	0.17%	0.8%
$R = \text{B}(\text{H} \rightarrow \gamma\gamma) / \text{B}(\text{H} \rightarrow 2\mu)$	$\delta R/R$	0.29%	1.38%
$R = \text{B}(\text{H} \rightarrow \mu\mu\gamma) / \text{B}(\text{H} \rightarrow \mu\mu)$	$\delta R/R$	0.58%	1.82%
$R = \sigma(\text{t}\bar{\text{t}}\text{H}) \times \text{B}(\text{H} \rightarrow \text{b}\bar{\text{b}}) / \sigma(\text{t}\bar{\text{t}}\text{Z}) \times \text{B}(\text{Z} \rightarrow \text{b}\bar{\text{b}})$	$\delta R/R$	1.05%	1.9%
$B(\text{H} \rightarrow \text{invisible})$	$B@95\% \text{CL}$	$1 \times 10^{-4}$	$2.5 \times 10^{-4}$

## Importance of standalone precise “ratios-of-BRs” measurements:

- independent of  $\alpha_S$ ,  $m_b$ ,  $m_c$ ,  $\Gamma_{inv}$  systematics
- sensitive to BSM effects that typically influence BRs in different ways. Eg

$$\mathbf{BR(H \rightarrow \gamma\gamma) / BR(H \rightarrow ZZ^*)}$$

loop-level

tree-level

$$\mathbf{BR(H \rightarrow \mu\mu) / BR(H \rightarrow ZZ^*)}$$

2nd gen'n Yukawa

gauge coupling

$$\mathbf{BR(H \rightarrow \gamma\gamma) / BR(H \rightarrow Z\gamma)}$$

different EW charges in the loops of the two procs

$$\mathbf{BR(H \rightarrow inv) / BR(H \rightarrow \gamma\gamma)}$$

tree-level neutral

loop-level charged

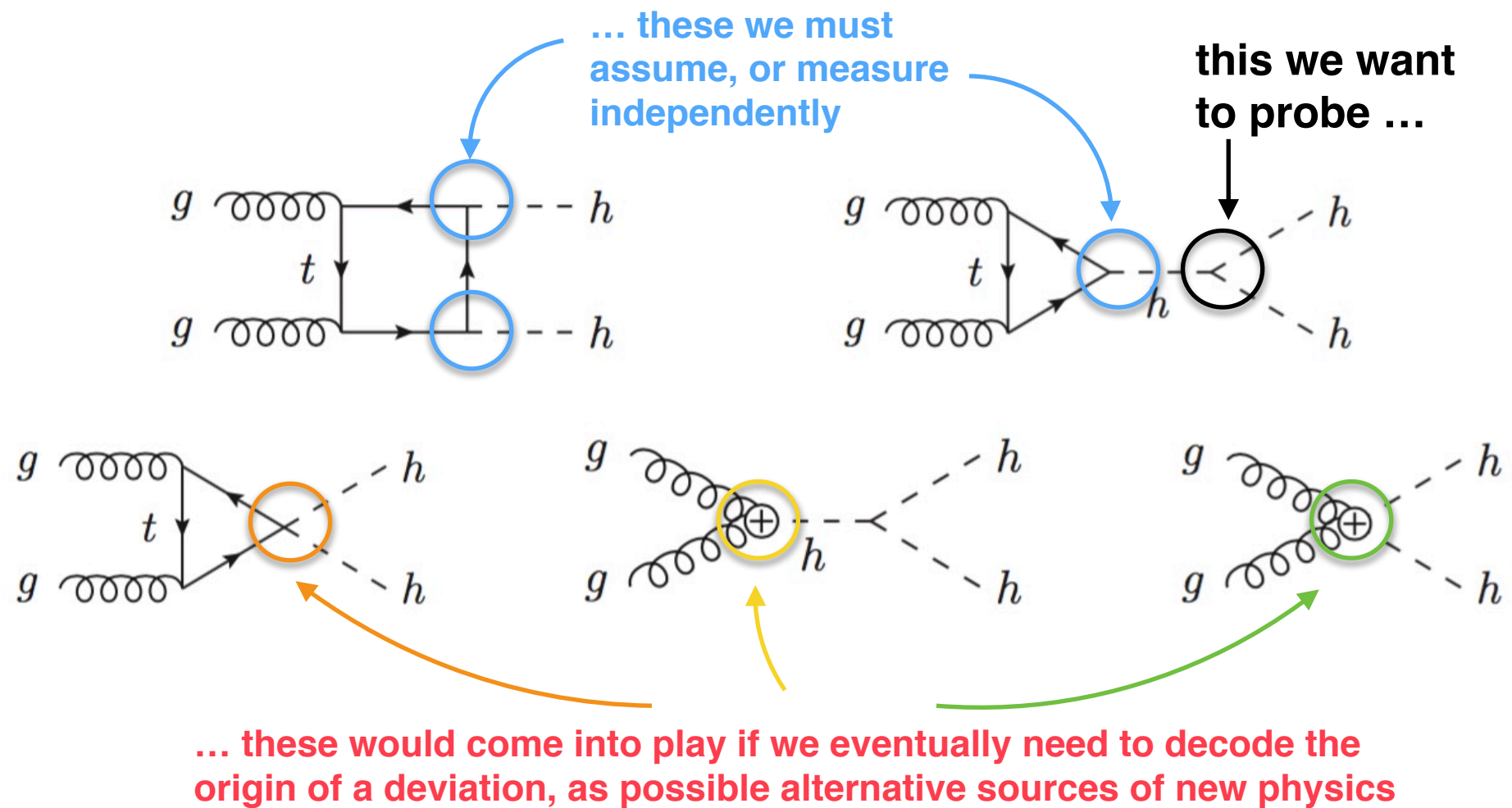
# Higgs couplings after FCC-ee / hh

	HL-LHC	FCC-ee	FCC-hh
$\delta\Gamma_H / \Gamma_H$ (%)	SM	<b>1.3</b>	tbd
$\delta g_{HZZ} / g_{HZZ}$ (%)	1.5	<b>0.17</b>	tbd
$\delta g_{HWW} / g_{HWW}$ (%)	1.7	<b>0.43</b>	tbd
$\delta g_{Hbb} / g_{Hbb}$ (%)	3.7	<b>0.61</b>	tbd
$\delta g_{Hcc} / g_{Hcc}$ (%)	~70	<b>1.21</b>	tbd
$\delta g_{Hgg} / g_{Hgg}$ (%)	2.5 (gg->H)	<b>1.01</b>	tbd
$\delta g_{H\tau\tau} / g_{H\tau\tau}$ (%)	1.9	<b>0.74</b>	tbd
$\delta g_{H\mu\mu} / g_{H\mu\mu}$ (%)	4.3	9.0	<b>0.65 (*)</b>
$\delta g_{H\gamma\gamma} / g_{H\gamma\gamma}$ (%)	1.8	3.9	<b>0.4 (*)</b>
$\delta g_{Htt} / g_{Htt}$ (%)	3.4	~10 (indirect)	<b>0.95 (**)</b>
$\delta g_{HZ\gamma} / g_{HZ\gamma}$ (%)	9.8	–	<b>0.9 (*)</b>
$\delta g_{HHH} / g_{HHH}$ (%)	50	~44 (indirect)	<b>6.5</b>
$BR_{\text{exo}}$ (95%CL)	$BR_{\text{inv}} < 2.5\%$	<b>&lt; 1%</b>	<b><math>BR_{\text{inv}} &lt; 0.025\%</math></b>

\* From BR ratios wrt  $B(H \rightarrow 4\text{lept})$  @ FCC-ee

\*\* From  $pp \rightarrow ttH$  /  $pp \rightarrow ttZ$ , using  $B(H \rightarrow bb)$  and  $ttZ$  EW coupling @ FCC-ee

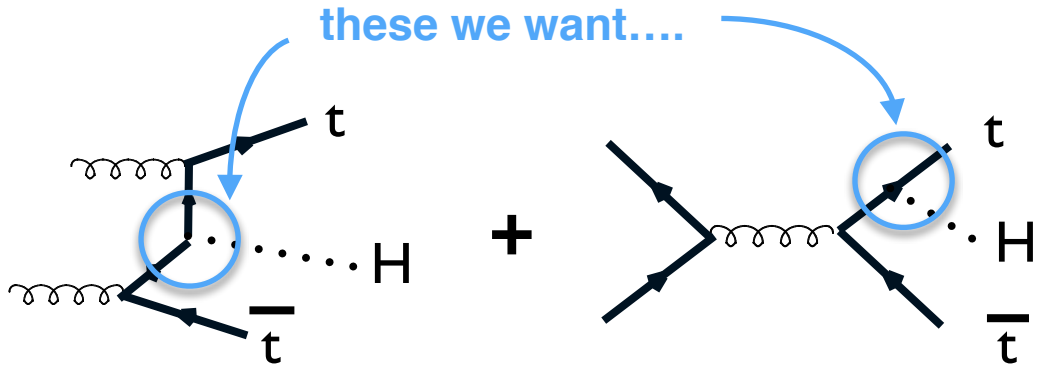
# Extracting Higgs self-coupling from $gg \rightarrow HH$ at FCC-hh



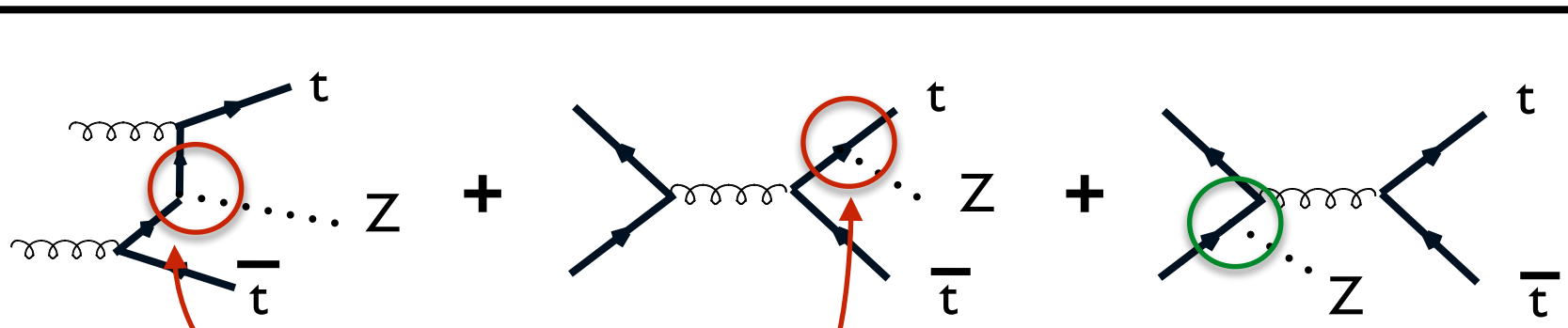


# Direct measurement of $ttH$ coupling: from $R_t = \sigma(ttH)/\sigma(ttZ)$

FCC-hh can measure  $R_t$  with  $\Delta R_t/R_t \sim 2\%$

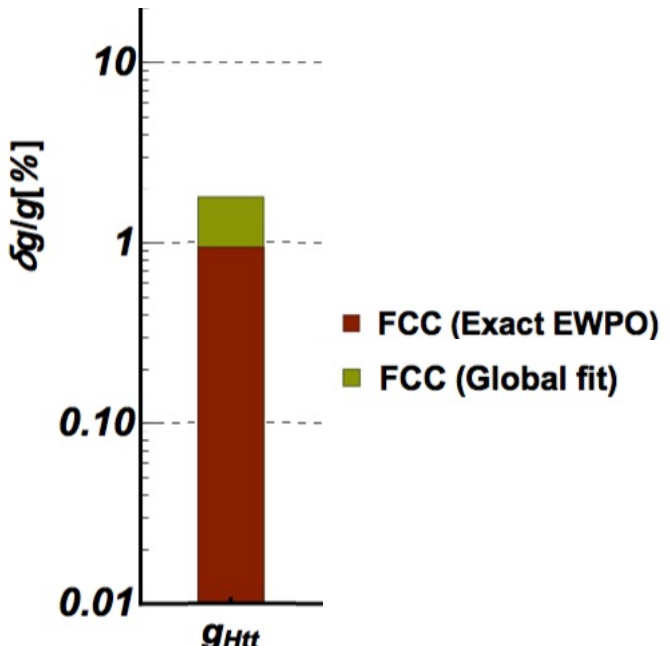
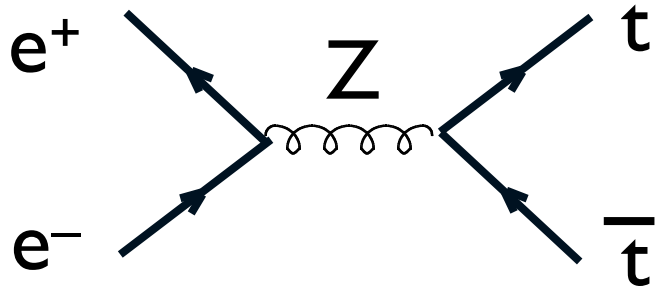


$R_t =$



this we must measure!

this we know (light quarks)



$\delta\lambda/\lambda=5\%$   
from  
 $gg \rightarrow HH$   
assuming  
SM inputs

$\delta\lambda/\lambda \sim 10\%$   
from global  
fit

# Higgs self-coupling, $gg \rightarrow HH$

From the detector performance studies:

Pheno-level studies:

	$bb\gamma\gamma$	$bbZZ[\rightarrow 4l]$	$bbWW[\rightarrow 2jlv]$	$4b+j$	$2b2\tau+j$
$\delta\kappa_\lambda$ (%)	6.5	14	40	30	8

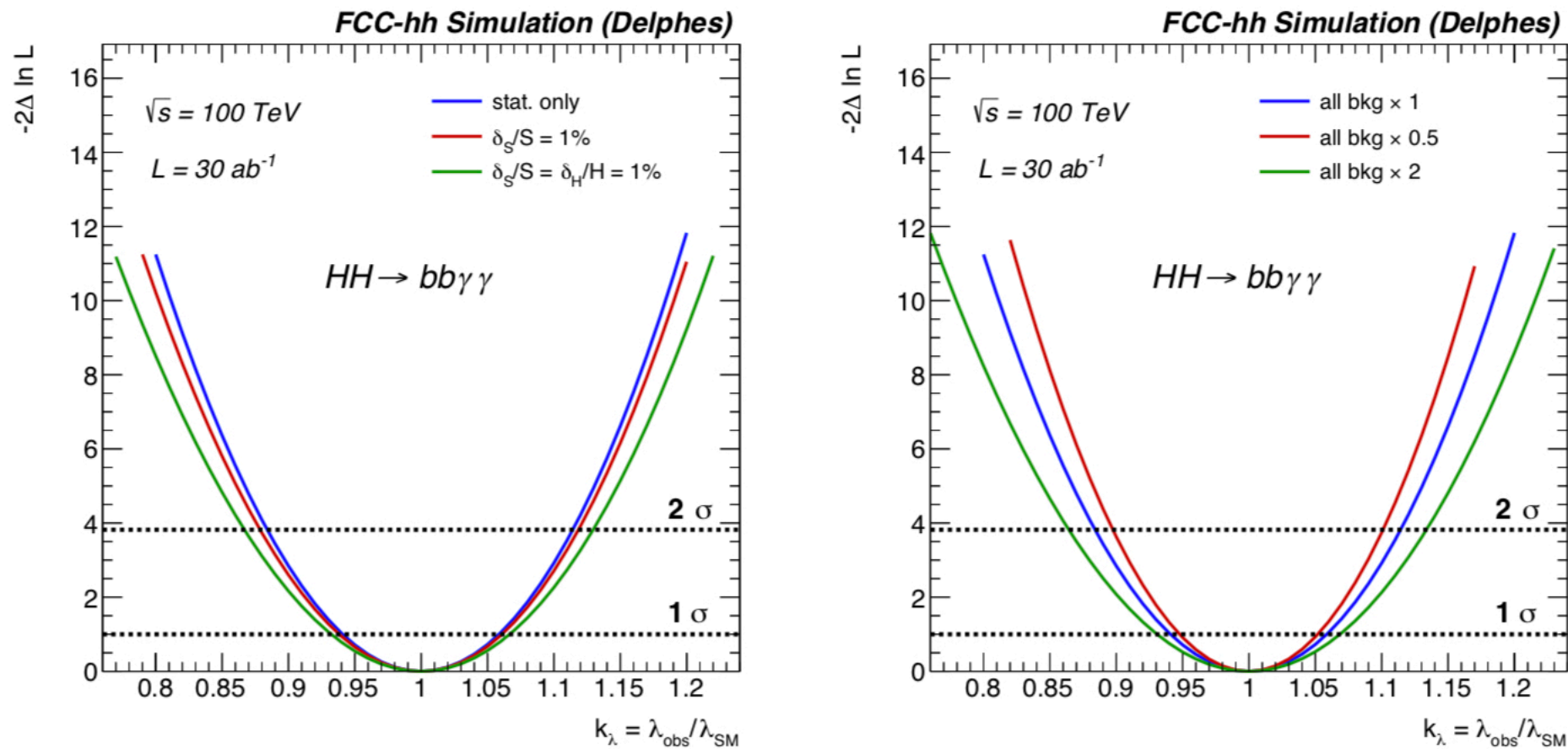
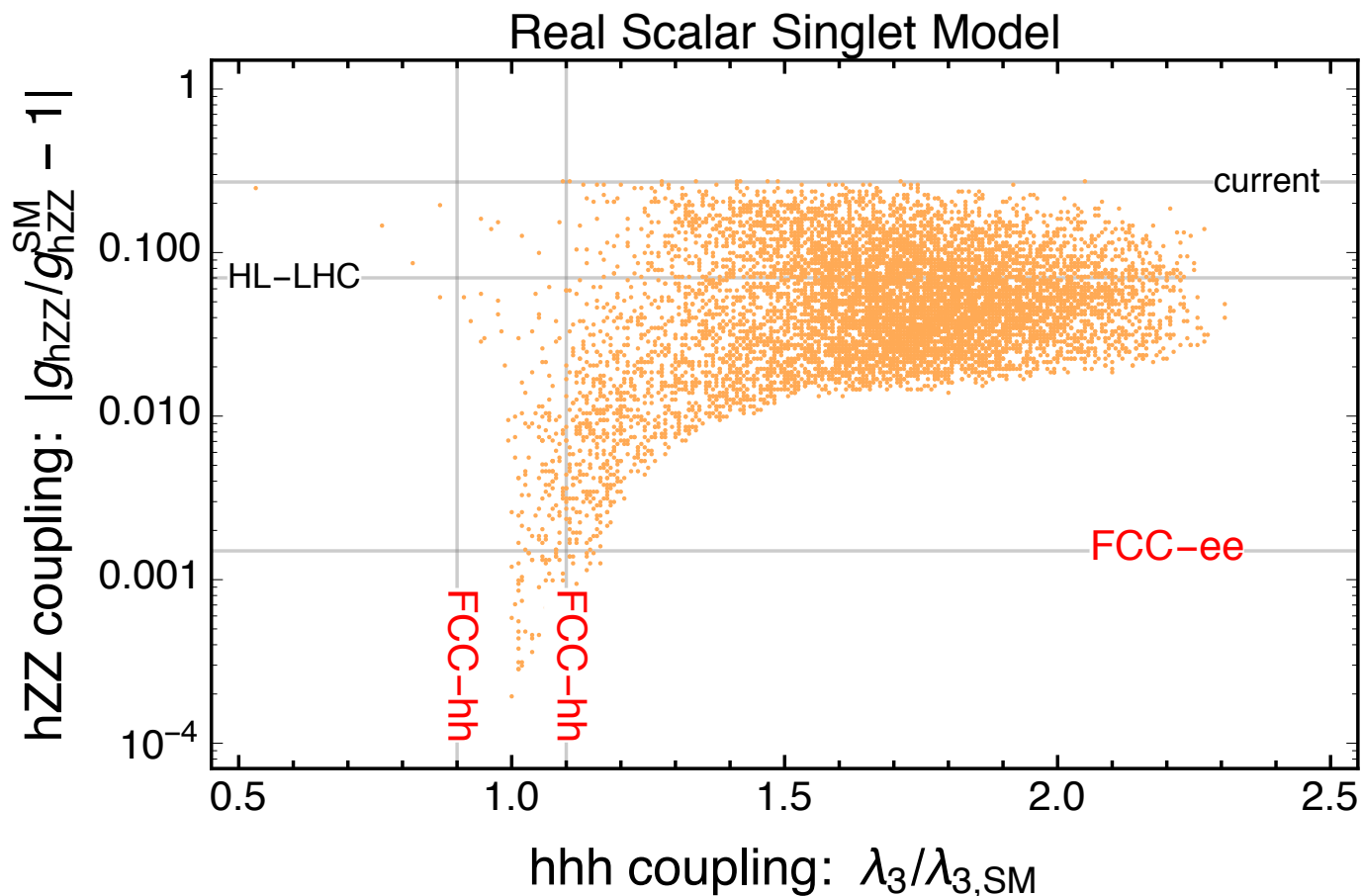


Figure 10.4: Expected precision on the Higgs self-coupling modifier  $\kappa_\lambda$  with no systematic uncertainties (only statistical), 1% signal uncertainty, 1% signal uncertainty together with 1% uncertainty on the Higgs backgrounds (left) and assuming respectively  $\times 1$ ,  $\times 2$ ,  $\times 0.5$  background yields (right).

# Example of precision targets: constraints on models with 1<sup>st</sup> order phase transition

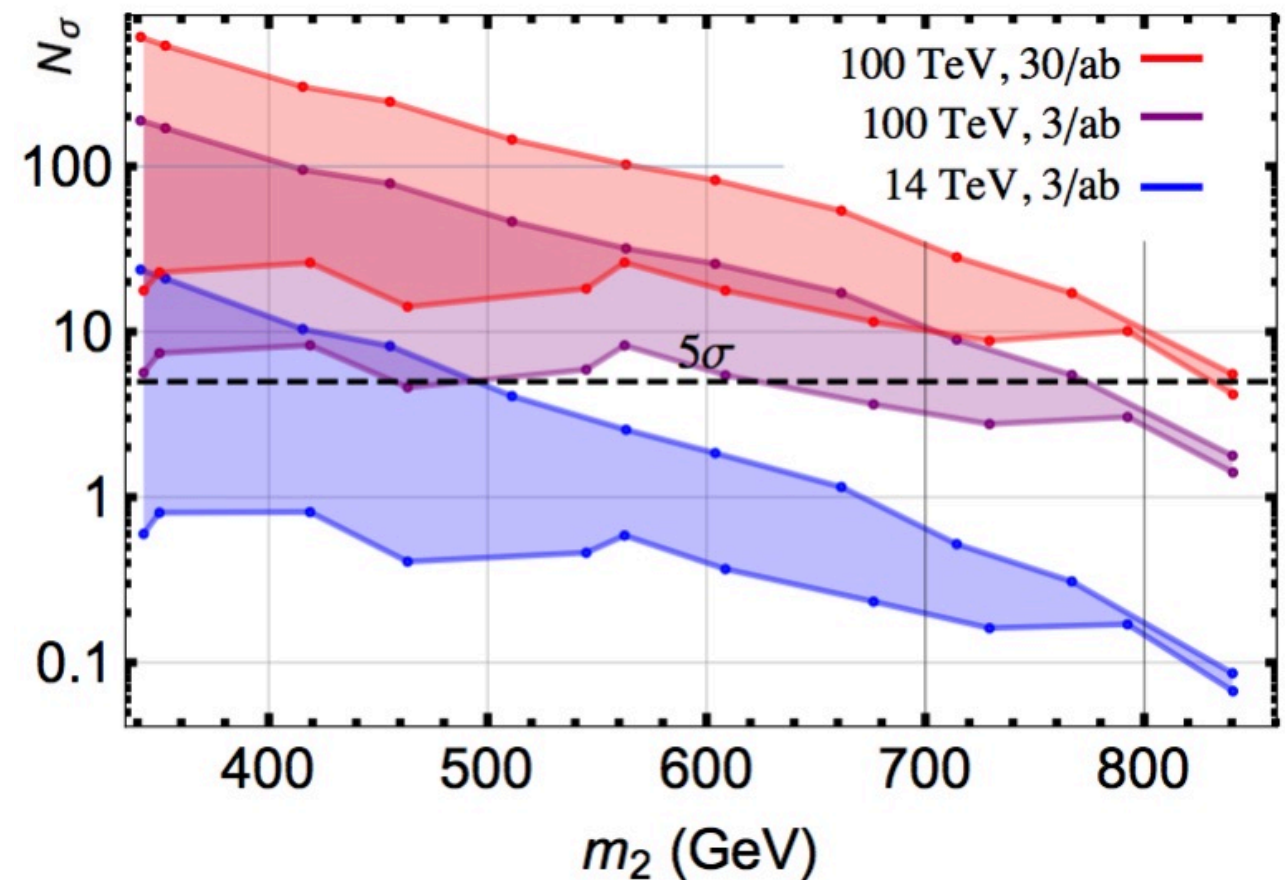
$$V(H, S) = -\mu^2 (H^\dagger H) + \lambda (H^\dagger H)^2 + \frac{a_1}{2} (H^\dagger H) S + \frac{a_2}{2} (H^\dagger H) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$$

Combined constraints from precision Higgs measurements at FCC-ee and FCC-hh



Parameter space scan for a singlet model extension of the Standard Model. The points indicate a first order phase transition.

Direct detection of extra Higgs states at FCC-hh



$h_2 \rightarrow h_1 h_1$  ( $b\bar{b}\gamma\gamma + 4\tau$ )  
( $h_2 \sim S, h_1 \sim H$ )

# Precision vs sensitivity

- We often talk about “**precise**” Higgs measurements. What we actually aim at is “**sensitive**” tests of the Higgs properties, where *sensitive* refers to the ability to reveal BSM behaviours.
- **Sensitivity** may not require extreme precision
  - Going after “sensitivity”, rather than *just* precision, opens itself new opportunities ...

# High- $Q^2$ observables : precision vs dynamic reach

$$L = L_{SM} + \frac{1}{\Lambda^2} \sum_k \mathcal{O}_k + \dots$$

$$O = | \langle f | L | i \rangle |^2 = O_{SM} [1 + O(\mu^2/\Lambda^2) + \dots]$$

For H decays, or inclusive production,  $\mu \sim O(v, m_H)$

$$\delta O \sim \left(\frac{v}{\Lambda}\right)^2 \sim 6\% \left(\frac{\text{TeV}}{\Lambda}\right)^2 \Rightarrow \text{precision probes large } \Lambda$$

$$\text{e.g. } \delta O = 1\% \Rightarrow \Lambda \sim 2.5 \text{ TeV}$$

For H production off-shell or with large momentum transfer  $Q$ ,  $\mu \sim O(Q)$

$$\delta O \sim \left(\frac{Q}{\Lambda}\right)^2 \Rightarrow \text{kinematic reach probes large } \Lambda$$

even if precision is “low”

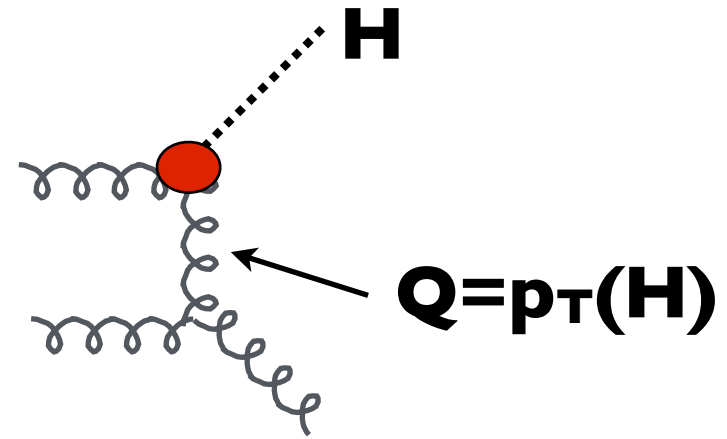
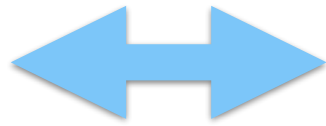
$$\text{e.g. } \delta O = 10\% \text{ at } Q = 1.5 \text{ TeV} \Rightarrow \Lambda \sim 5 \text{ TeV}$$

**Complementarity between precise measurements at ee collider and large- $Q$  studies at 100 TeV**

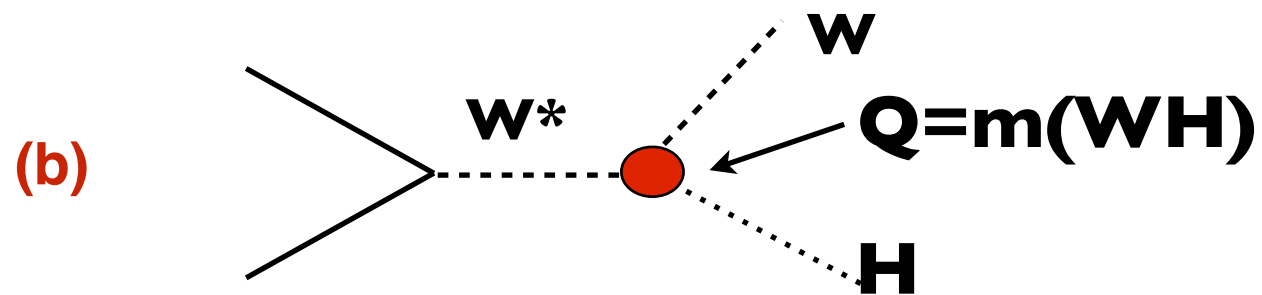
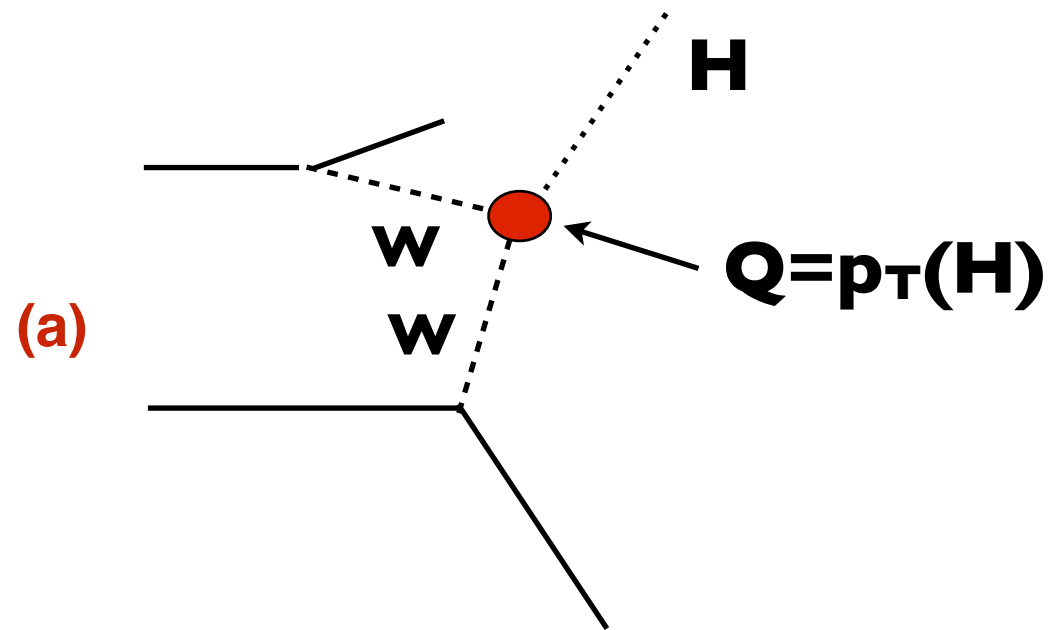
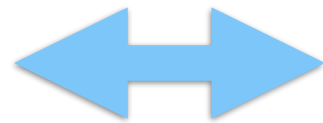


# Examples

$\delta BR(H \rightarrow gg)$



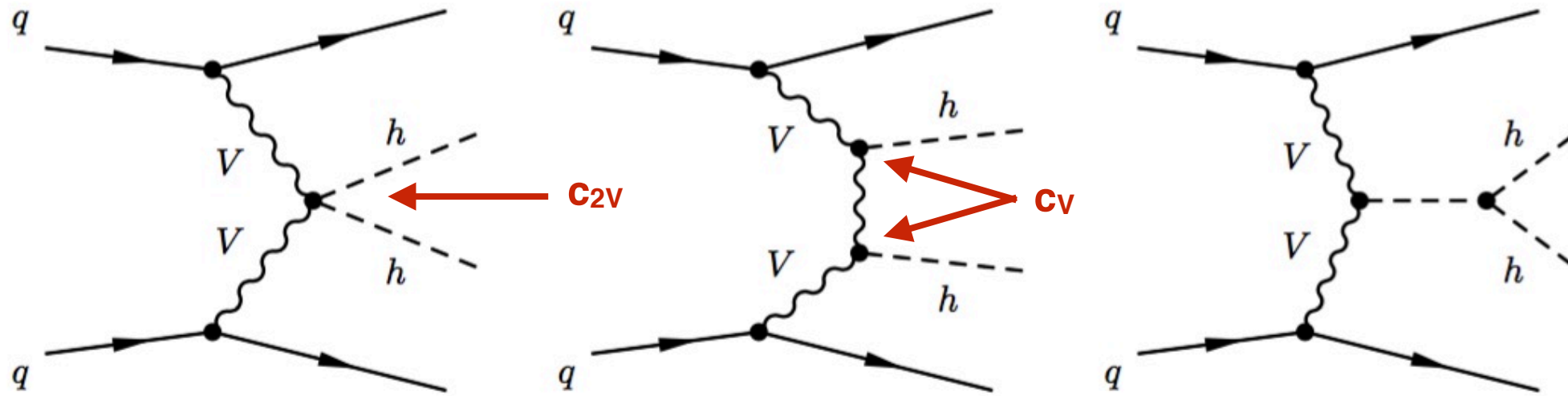
$\delta BR(H \rightarrow WW^*)$



$$L_{D=6} = \frac{ig c_W}{2 \Lambda^2} (H^\dagger \sigma^a D^\mu H) D^\nu V_{\mu\nu}^a$$

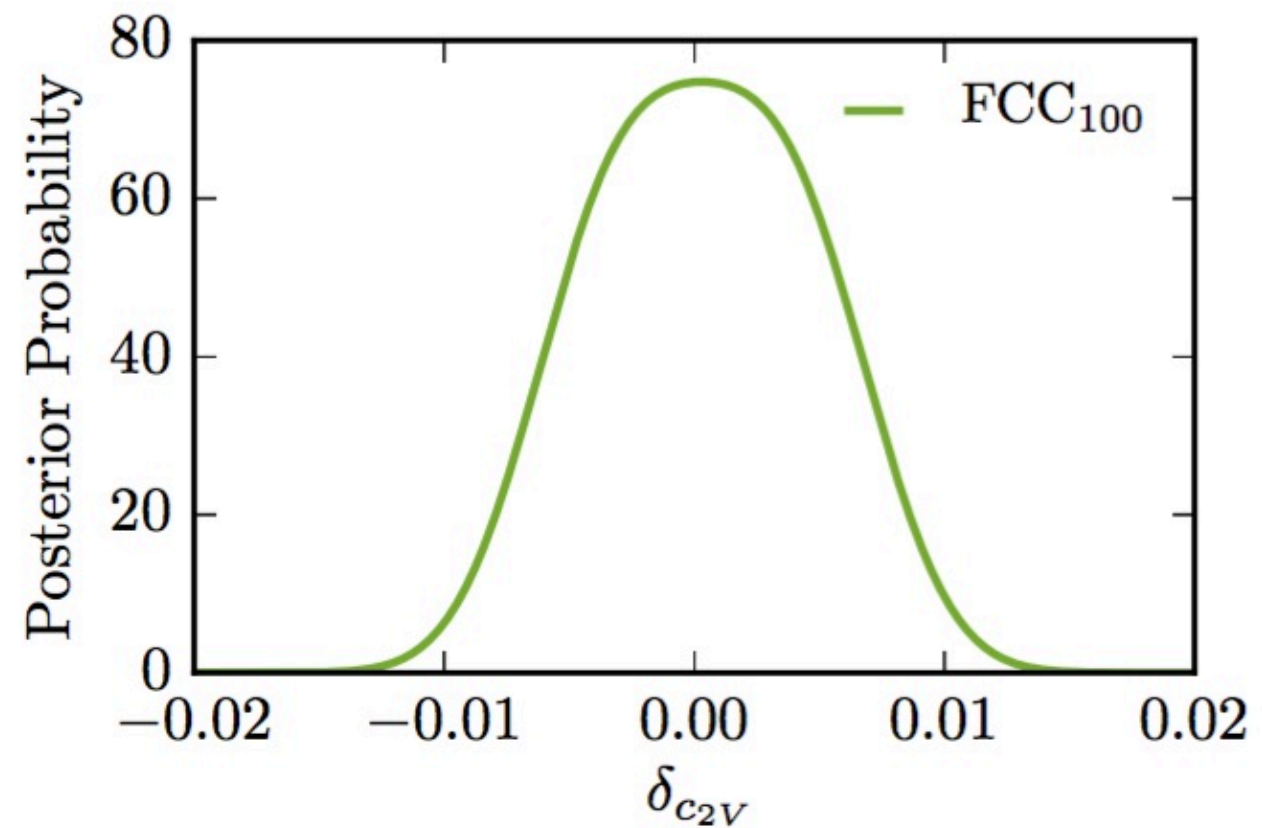
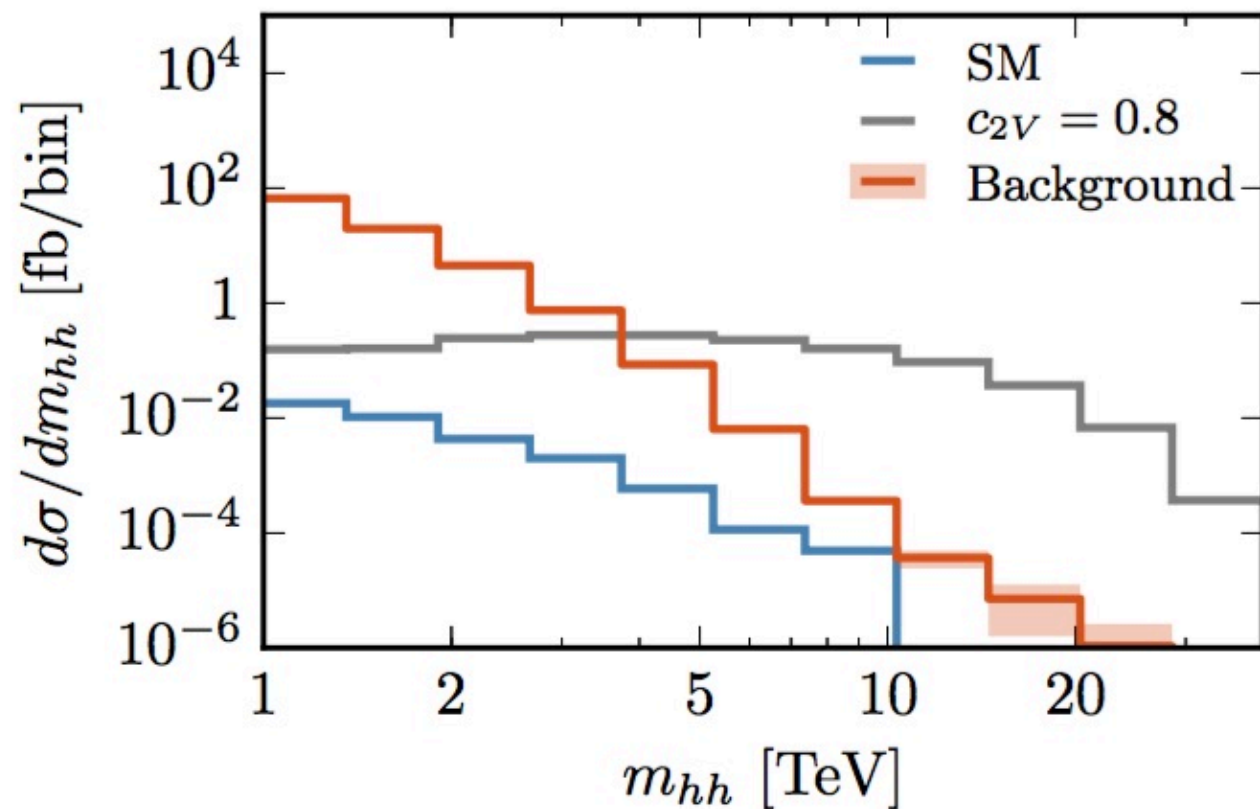
$$\frac{\sigma}{\sigma_{SM}} \sim \left( 1 + c_W \frac{\hat{s}}{\Lambda^2} \right)^2$$

# Example: high mass $VV \rightarrow HH$

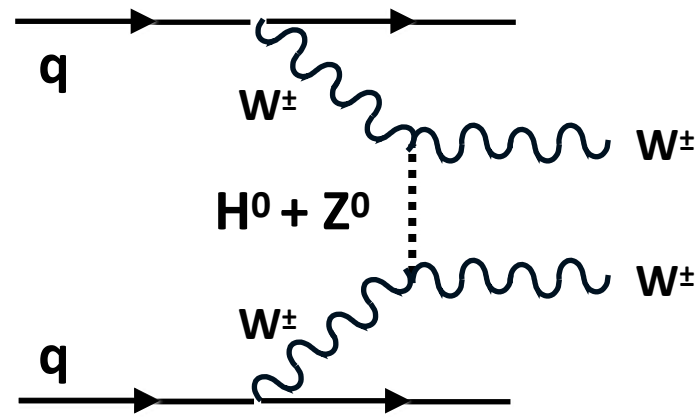


$$A(V_L V_L \rightarrow HH) \sim \frac{\hat{s}}{v^2} (c_{2V} - c_V^2) \cdot \text{where}$$

$$\begin{cases} c_V = g_{HVV}/g_{HVV}^{SM} \\ c_{2V} = g_{HHVV}/g_{HHVV}^{SM} \end{cases} \Rightarrow (c_{2V} - c_V^2)_{SM} = 0$$

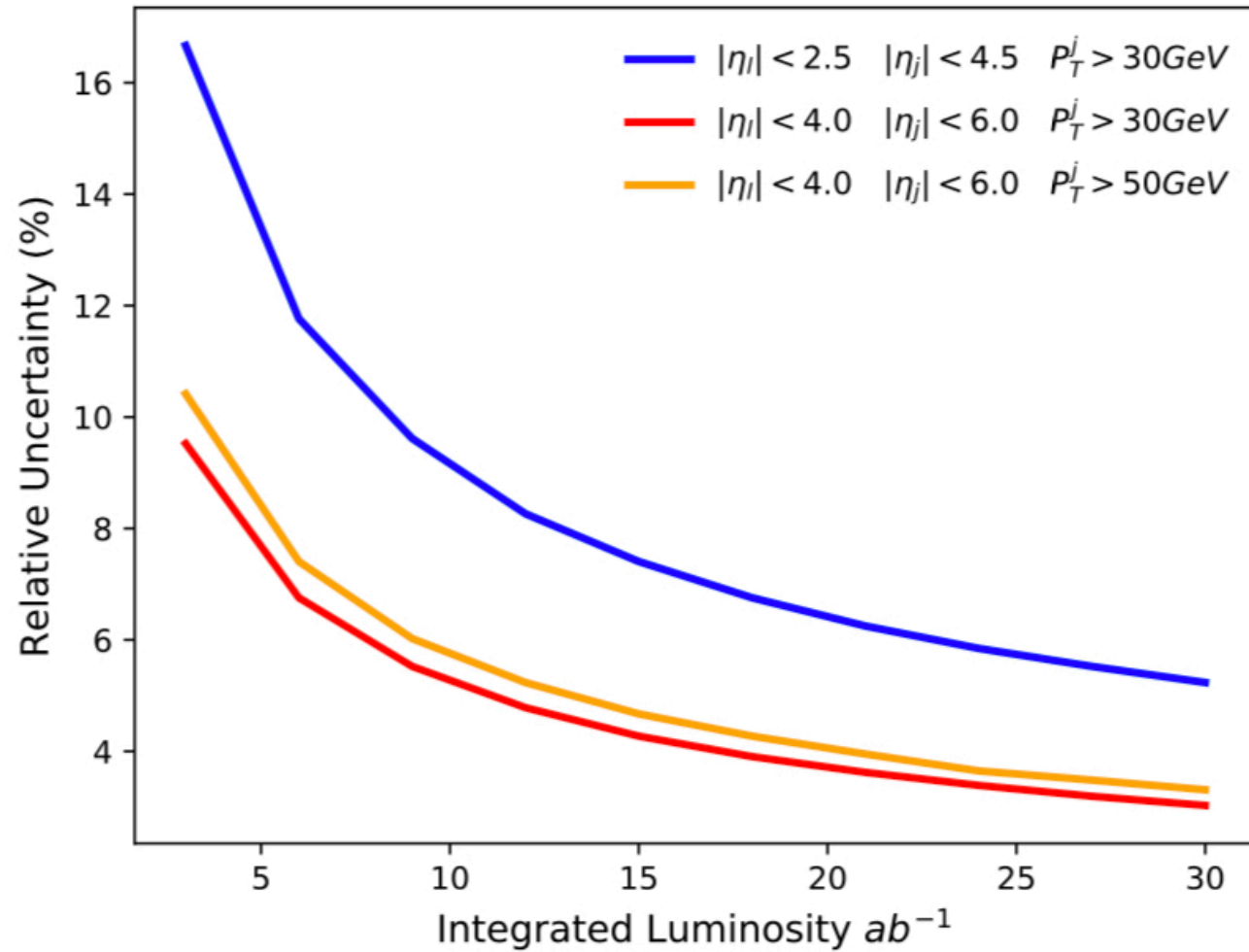


# $W_L W_L$ scattering



large  $m_{WW}$

VBS  $W_L W_L$  Same Sign Cross Uncertainty



FCC-hh Simulation (Delphes)

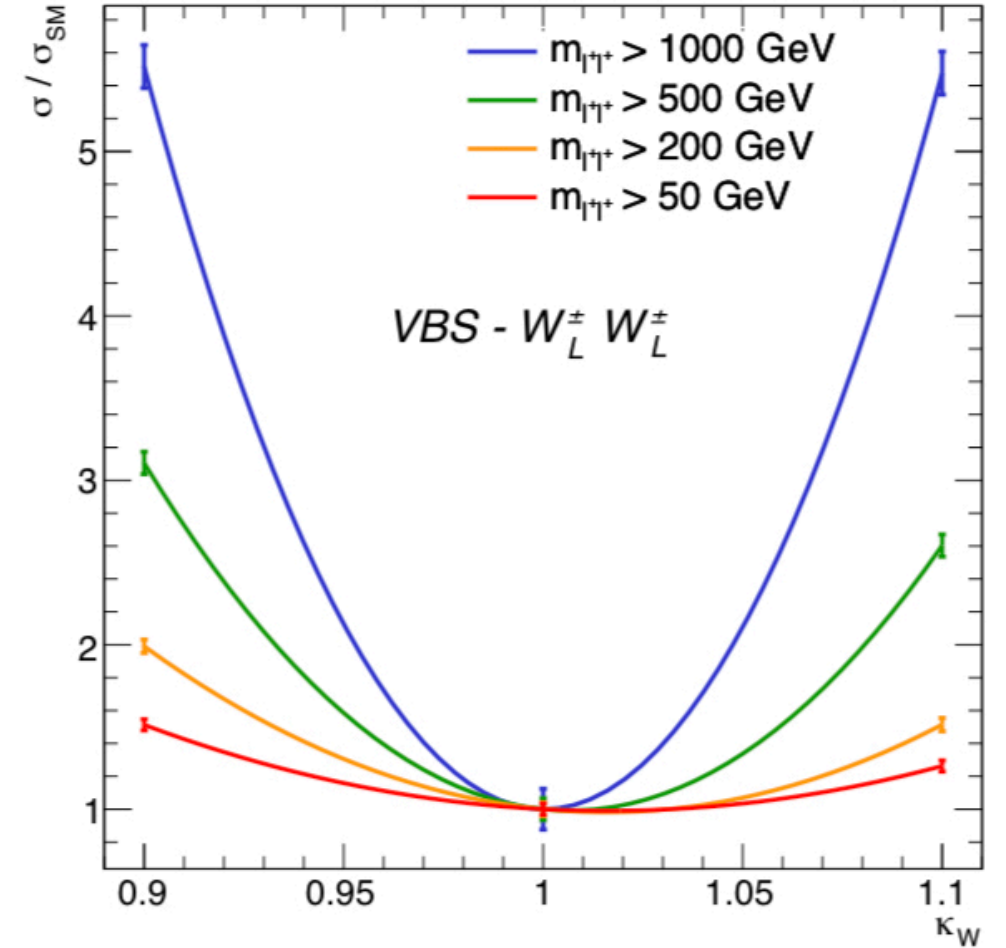


Table 4.5: Constraints on the HWW coupling modifier  $\kappa_W$  at 68% CL, obtained for various cuts on the di-lepton pair invariant mass in the  $W_L W_L \rightarrow HH$  process.

$m_{l+l^+}$ cut	$> 50\text{ GeV}$	$> 200\text{ GeV}$	$> 500\text{ GeV}$	$> 1000\text{ GeV}$
$\kappa_W \in$	[0.98, 1.05]	[0.99, 1.04]	[0.99, 1.03]	[0.98, 1.02]

$$\kappa_W = \frac{g_{HWW}}{g_{HWW}^{SM}}$$

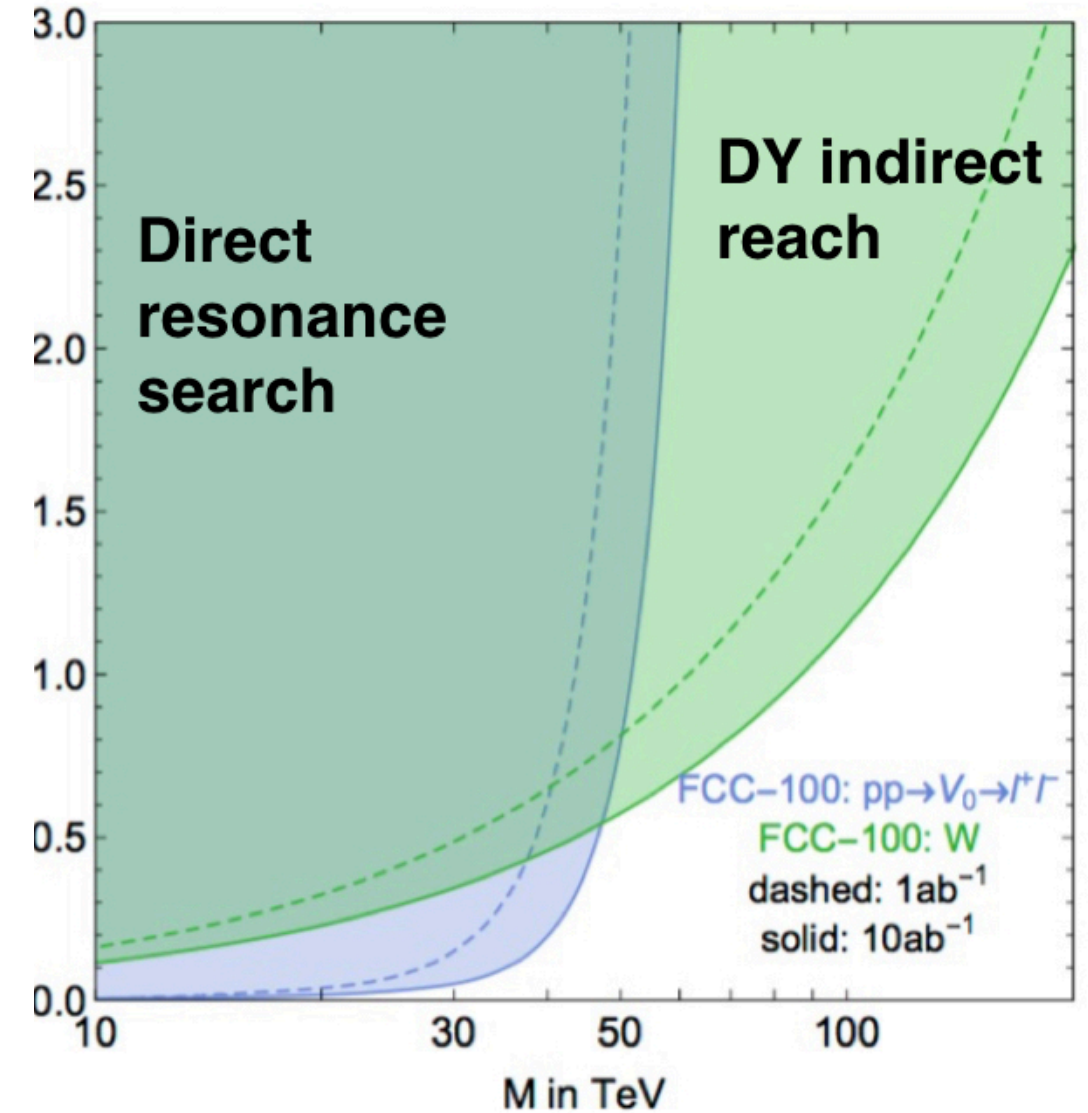
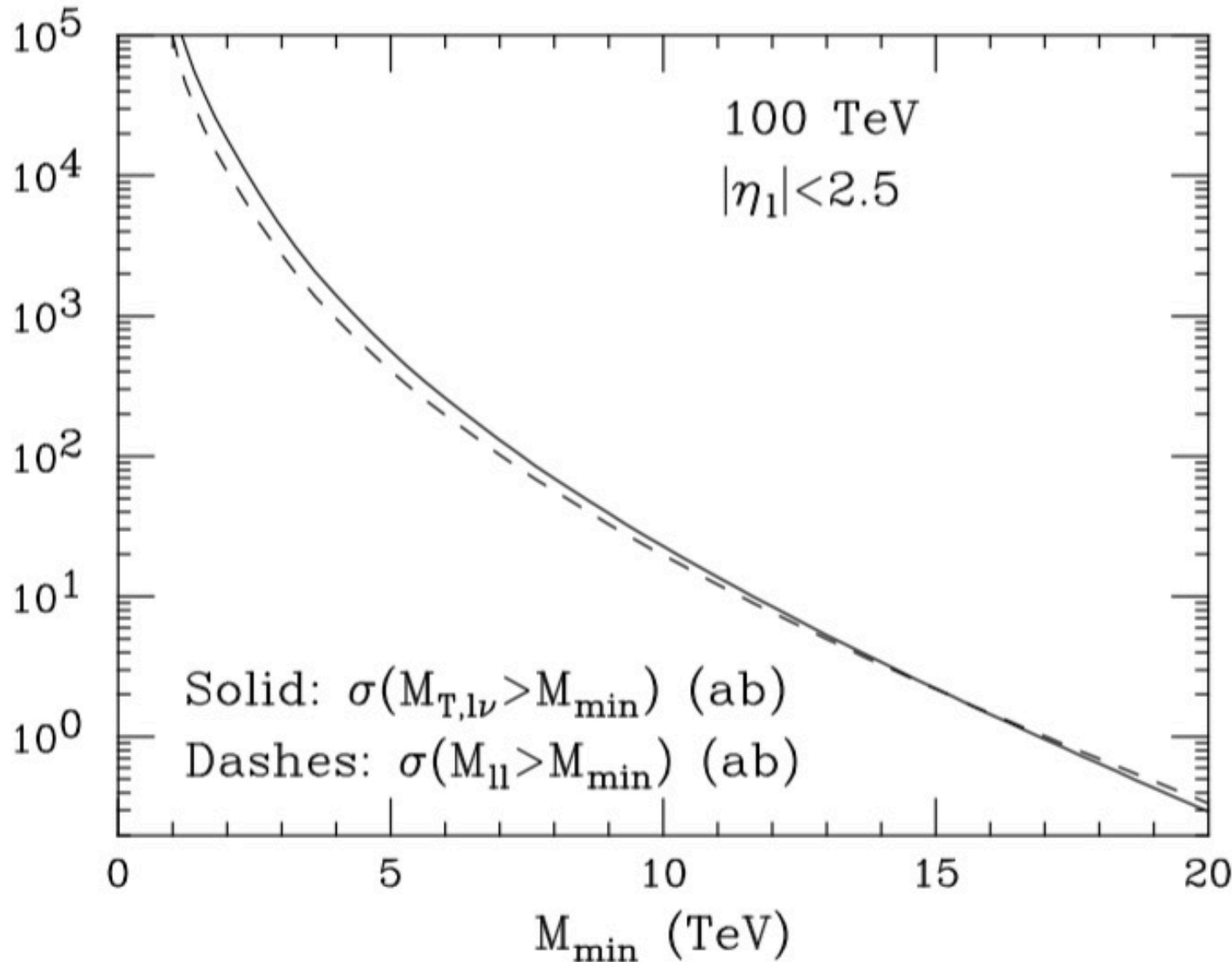
# Example: high mass DY

Constraints on Higher-dim op's

$$\hat{W} = -\frac{W}{4m_W^2} (D_\rho W_{\mu\nu}^a)^2, \quad \hat{Y} = -\frac{Y}{4m_W^2} (\partial_\rho B_{\mu\nu})^2$$

	LEP	LHC 13	FCC 100	ILC	TLEP	CEPC	ILC 500	CLIC 1	CLIC 3	
luminosity	$2 \times 10^7 Z$	0.3/ab	3/ab	10/ab	$10^9 Z$	$10^{12} Z$	$10^{10} Z$	3/ab	1/ab	1/ab
$W \times 10^4$	[-19, 3]	$\pm 0.7$	$\pm 0.45$	$\pm 0.02$	$\pm 4.2$	$\pm 1.2$	$\pm 3.6$	$\pm 0.3$	$\pm 0.5$	$\pm 0.15$
$Y \times 10^4$	[-17, 4]	$\pm 2.3$	$\pm 1.2$	$\pm 0.06$	$\pm 1.8$	$\pm 1.5$	$\pm 3.1$	$\pm 0.2$	$\sim \pm 0.5$	$\sim \pm 0.15$

$W / 4m_W^2 < 1 / (100 \text{ TeV})^2$



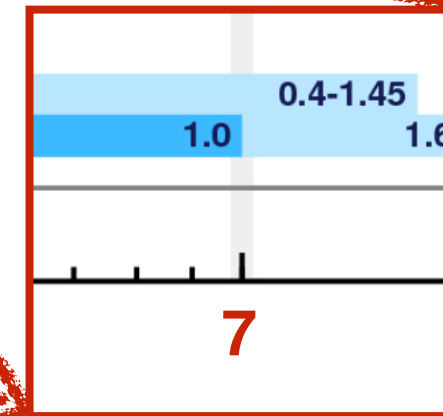
**Direct discovery reach:  
the power of 100 TeV**



Model	Signature	$\int \mathcal{L} dt \text{ [fb}^{-1}\text{]}$	Mass limit	Reference	
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 $e, \mu$ mono-jet	2-6 jets 1-3 jets $E_T^{\text{miss}}$ 36.1	$\tilde{q}$ [2x, 8x Degen.] $\tilde{q}$ [1x, 8x Degen.] 0.43 0.71 0.9 1.55 $m(\tilde{\chi}_1^0) < 100 \text{ GeV}$ $m(\tilde{q}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1712.02332 1711.03301
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0 $e, \mu$	2-6 jets $E_T^{\text{miss}}$ 36.1	$\tilde{g}$ $\tilde{g}$ Forbidden 0.95-1.6 2.0 $m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{\chi}_1^0) = 900 \text{ GeV}$	1712.02332 1712.02332
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	3 $e, \mu$ $ee, \mu\mu$	4 jets 2 jets $E_T^{\text{miss}}$ 36.1	$\tilde{g}$ $\tilde{g}$ 1.2 1.85 $m(\tilde{\chi}_1^0) < 800 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 50 \text{ GeV}$	1706.03731 1805.11381
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$	0 $e, \mu$ 3 $e, \mu$	7-11 jets 4 jets $E_T^{\text{miss}}$ 36.1	$\tilde{g}$ $\tilde{g}$ 0.98 1.8 $m(\tilde{\chi}_1^0) < 400 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200 \text{ GeV}$	1708.02794 1706.03731
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$	0-1 $e, \mu$ 3 $e, \mu$	3 $b$ 4 jets $E_T^{\text{miss}}$ 79.8 36.1	$\tilde{g}$ $\tilde{g}$ 1.25 2.25 $m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300 \text{ GeV}$	ATLAS-CONF-2018-041 1706.03731
3 <sup>rd</sup> gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/t\tilde{\chi}_1^\pm$	Multiple Multiple Multiple	36.1 36.1 36.1	$\tilde{b}_1$ $\tilde{b}_1$ $\tilde{b}_1$ Forbidden 0.9 Forbidden 0.58-0.82 Forbidden 0.7 $m(\tilde{\chi}_1^0) = 300 \text{ GeV}, \text{BR}(b\tilde{\chi}_1^0) = 1$ $m(\tilde{\chi}_1^0) = 300 \text{ GeV}, \text{BR}(b\tilde{\chi}_1^0) = \text{BR}(t\tilde{\chi}_1^\pm) = 0.5$ $m(\tilde{\chi}_1^0) = 200 \text{ GeV}, m(\tilde{\chi}_1^\pm) = 300 \text{ GeV}, \text{BR}(t\tilde{\chi}_1^\pm) = 1$	1708.09266, 1711.03301 1708.09266 1706.03731
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$	0 $e, \mu$	6 $b$ $E_T^{\text{miss}}$ 139	$\tilde{b}_1$ $\tilde{b}_1$ Forbidden 0.23-0.48 0.23-1.35 $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130 \text{ GeV}, m(\tilde{\chi}_1^0) = 100 \text{ GeV}$ $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130 \text{ GeV}, m(\tilde{\chi}_1^0) = 0 \text{ GeV}$	SUSY-2018-31 SUSY-2018-31
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$	0-2 $e, \mu$	0-2 jets/1-2 $b$ $E_T^{\text{miss}}$ 36.1	$\tilde{t}_1$ 1.0 $m(\tilde{\chi}_1^0) = 1 \text{ GeV}$	1506.08616, 1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1$ , Well-Tempered LSP	Multiple	36.1	$\tilde{t}_1$ 0.48-0.84 $m(\tilde{\chi}_1^0) = 150 \text{ GeV}, m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}, \tilde{t}_1 \approx \tilde{t}_L$	1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$	1 $\tau + 1 e, \mu, \tau$	2 jets/1 $b$ $E_T^{\text{miss}}$ 36.1	$\tilde{t}_1$ 1.16 $m(\tilde{\tau}_1) = 800 \text{ GeV}$	1803.10178
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$	0 $e, \mu$	2 $c$ $E_T^{\text{miss}}$ 36.1	$\tilde{t}_1$ $\tilde{t}_1$ 0.46 0.85 $m(\tilde{\chi}_1^0) = 0 \text{ GeV}$ $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 50 \text{ GeV}$ $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1805.01649 1805.01649 1711.03301
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$	0 $e, \mu$	mono-jet $E_T^{\text{miss}}$ 36.1	$\tilde{t}_1$ 0.43 $m(\tilde{\chi}_1^0) = 0 \text{ GeV}, m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 180 \text{ GeV}$	1706.03986
EW direct	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via WZ	2-3 $e, \mu$ $ee, \mu\mu$	$\geq 1$ $E_T^{\text{miss}}$ 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.17 0.6 $m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 10 \text{ GeV}$	1403.5294, 1806.02293 1712.08119
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via WW	2 $e, \mu$	$E_T^{\text{miss}}$ 139	$\tilde{\chi}_1^\pm$ 0.42 $m(\tilde{\chi}_1^0) = 0$	ATLAS-CONF-2019-008
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via Wh	0-1 $e, \mu$	2 $b$ $E_T^{\text{miss}}$ 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.68 $m(\tilde{\chi}_1^0) = 0$	1812.09432
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via $\tilde{\ell}_L/\tilde{\nu}$	2 $e, \mu$	$E_T^{\text{miss}}$ 139	$\tilde{\chi}_1^\pm$ 1.0 $m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	ATLAS-CONF-2019-008
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}_1\nu(\tau\tilde{\nu}), \tilde{\chi}_2^0 \rightarrow \tilde{\tau}_1\nu(\nu\tilde{\nu})$	2 $\tau$	$E_T^{\text{miss}}$ 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.76 0.22 $m(\tilde{\chi}_1^0) = 0, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 100 \text{ GeV}, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	1708.07875 1708.07875
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 $e, \mu$ 2 $e, \mu$	0 jets $\geq 1$ $E_T^{\text{miss}}$ 139 36.1	$\tilde{\ell}$ $\tilde{\ell}$ 0.7 0.18 $m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\ell}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	ATLAS-CONF-2019-008 1712.08119
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 $e, \mu$ 4 $e, \mu$	$\geq 3 b$ 0 jets $E_T^{\text{miss}}$ 36.1 36.1	$\tilde{H}$ $\tilde{H}$ 0.13-0.23 0.29-0.88 0.3 $\text{BR}(\tilde{\chi}_1^0 \rightarrow h\tilde{G}) = 1$ $\text{BR}(\tilde{\chi}_1^0 \rightarrow Z\tilde{G}) = 1$	1806.04030 1804.03602
Long-lived particles	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet $E_T^{\text{miss}}$ 36.1	$\tilde{\chi}_1^\pm$ $\tilde{\chi}_1^\pm$ 0.15 0.46 Pure Wino Pure Higgsino	1712.02118 ATL-PHYS-PUB-2017-019
	Stable $\tilde{g}$ R-hadron	Multiple	36.1	$\tilde{g}$ 2.0 $m(\tilde{\chi}_1^0) = 100 \text{ GeV}$	1902.01636, 1808.04095
	Metastable $\tilde{g}$ R-hadron, $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$	Multiple	36.1	$\tilde{g}$ [ $\tau(\tilde{g}) = 10 \text{ ns}, 0.2 \text{ ns}$ ] 2.05 2.4	1710.04901, 1808.04095
RPV	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\tau\mu/\mu\tau$	$e\mu, e\tau, \mu\tau$	3.2	$\tilde{\nu}_\tau$ 1.9 $\lambda'_{311} = 0.11, \lambda'_{132/133/233} = 0.07$	1607.08079
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\ell\nu\nu$	4 $e, \mu$	0 jets $E_T^{\text{miss}}$ 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ [ $\lambda'_{333} \neq 0, \lambda'_{12k} \neq 0$ ] 0.82 1.33 $m(\tilde{\chi}_1^0) = 100 \text{ GeV}$	1804.03602
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$	4-5 large-R jets	36.1	$\tilde{g}$ [ $m(\tilde{\chi}_1^0) = 200 \text{ GeV}, 1100 \text{ GeV}$ ] $\tilde{g}$ [ $\lambda'_{112} = 2e-4, 2e-5$ ] 1.05 1.3 1.9 2.0 Large $\lambda'_{112}$	1804.03568
	$\tilde{u}, \tilde{t} \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs$	Multiple	36.1	$\tilde{g}$ [ $\lambda'_{323} = 2e-4, 1e-2$ ] 0.55 1.05 $m(\tilde{\chi}_1^0) = 200 \text{ GeV}, \text{bino-like}$	ATLAS-CONF-2018-003
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$	2 jets + 2 $b$	36.7	$\tilde{t}_1$ [ $qq, bs$ ] 0.42 0.61 0.4-1.45 $m(\tilde{\chi}_1^0) = 200 \text{ GeV}, \text{bino-like}$	ATLAS-CONF-2018-003 1710.07171
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$	2 $e, \mu$ 1 $\mu$	2 $b$ DV 36.1 136	$\tilde{t}_1$ $\tilde{t}_1$ 1.0 1.6 [ $1e-10 < \lambda'_{23k} < 1e-8, 3e-10 < \lambda'_{23k} < 3e-9$ ] $\text{BR}(\tilde{t}_1 \rightarrow q\mu) = 100\%, \text{cos}\theta > 20\%$	1710.05544 ATLAS-CONF-2019-006	

10<sup>-1</sup> 1 Mass scale [TeV]

@14 TeV

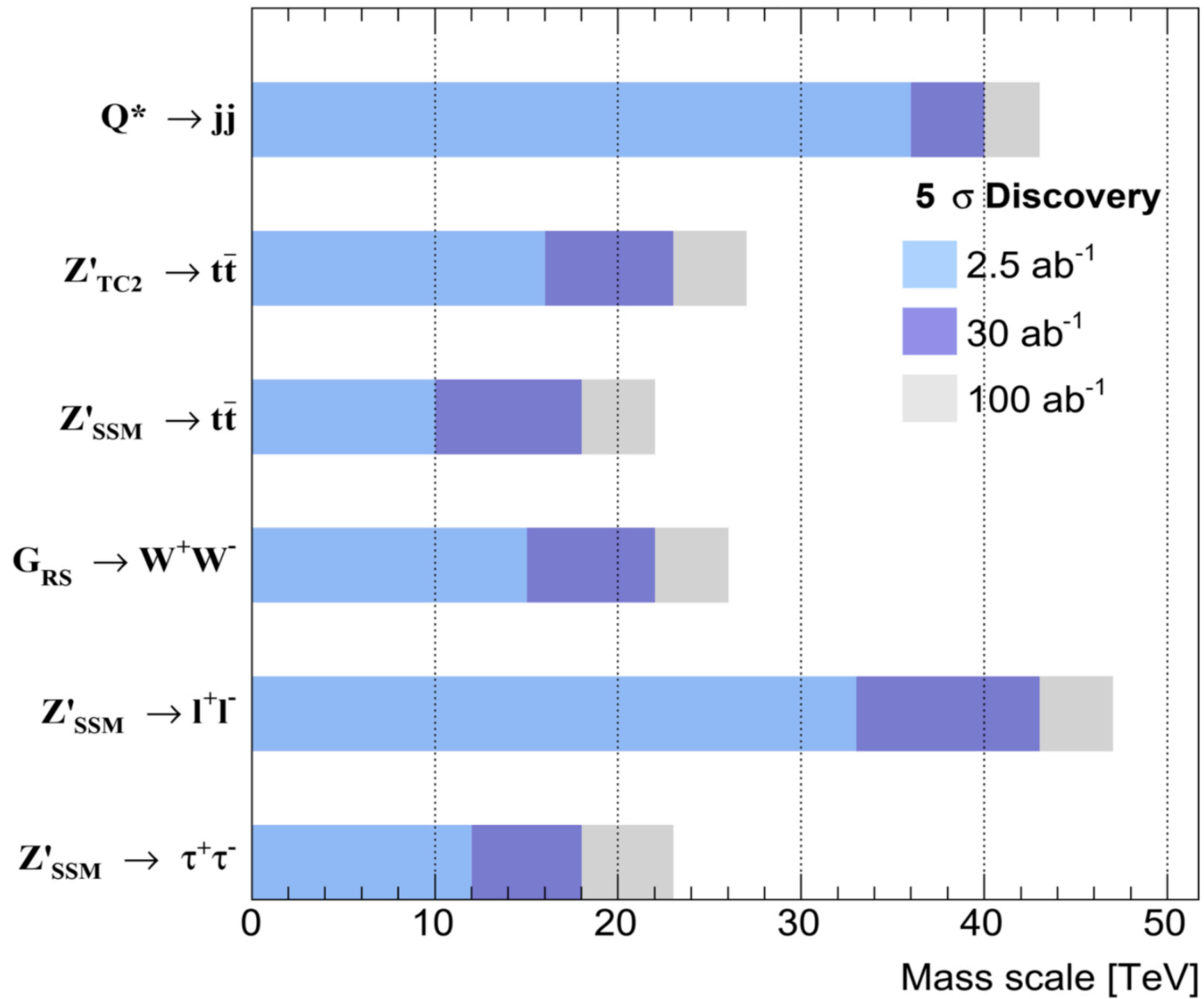


@100 TeV

\*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

# s-channel resonances

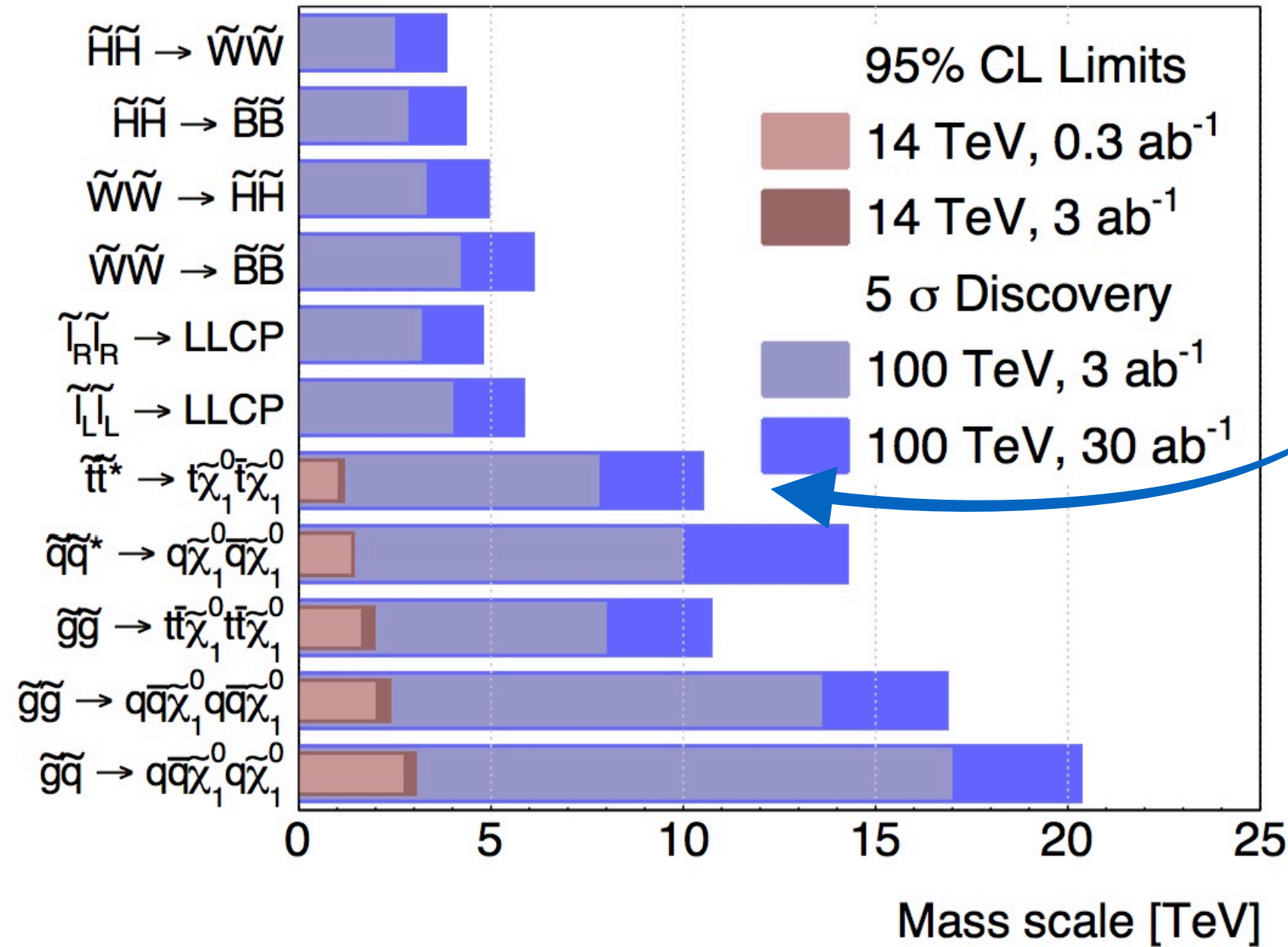
*FCC-hh Simulation (Delphes),  $\sqrt{s} = 100$  TeV*



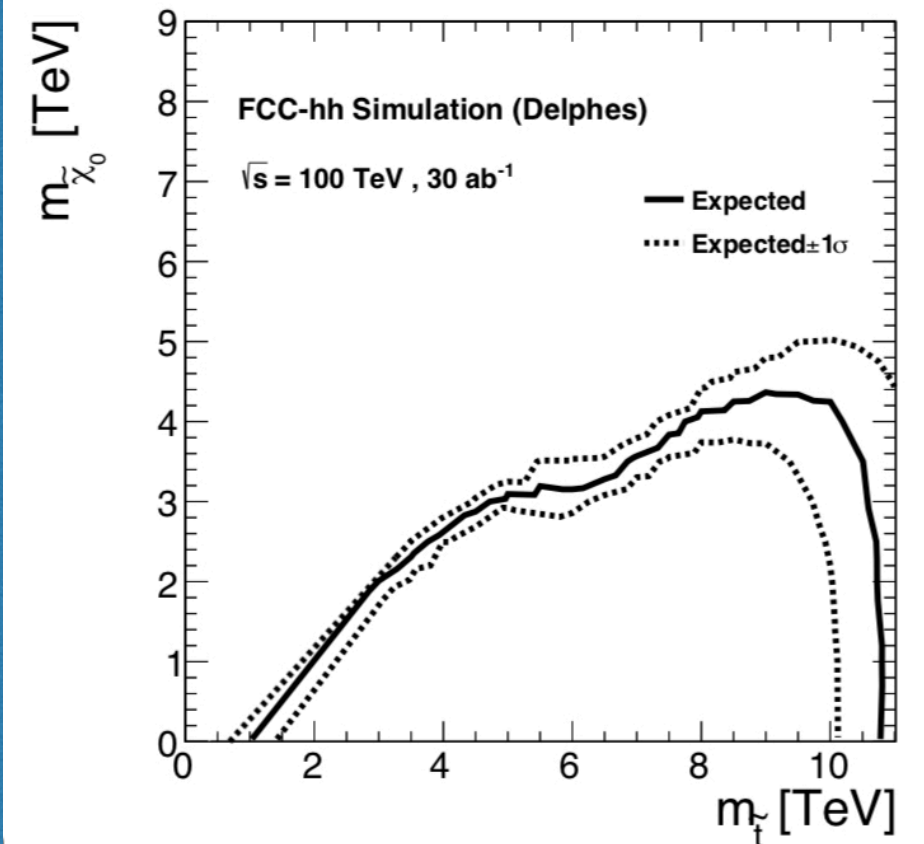
**FCC-hh reach  $\sim 6$  x HL-LHC reach**

# SUSY reach at 100 TeV

## Early phenomenology studies



## New detector performance studies



# WIMP DM theoretical constraints

For particles held in equilibrium by pair creation and annihilation processes, ( $\chi \chi \leftrightarrow \text{SM}$ )

$$\Omega_{\text{DM}} h^2 \sim \frac{10^9 \text{GeV}^{-1}}{M_{\text{pl}}} \frac{1}{\langle \sigma v \rangle}$$

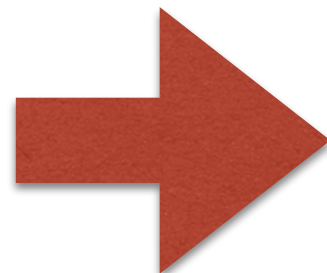
For a particle annihilating through processes which do not involve any larger mass scales:

$$\langle \sigma v \rangle \sim g_{\text{eff}}^4 / M_{\text{DM}}^2$$



$$\Omega_{\text{DM}} h^2 \sim 0.12 \times \left( \frac{M_{\text{DM}}}{2 \text{TeV}} \right)^2 \left( \frac{0.3}{g_{\text{eff}}} \right)^4$$

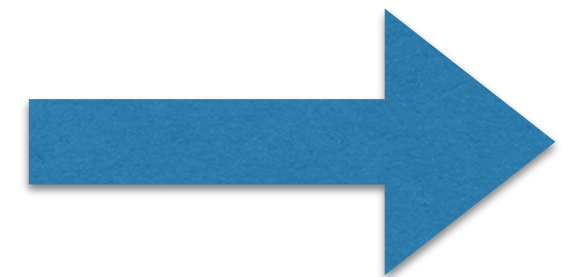
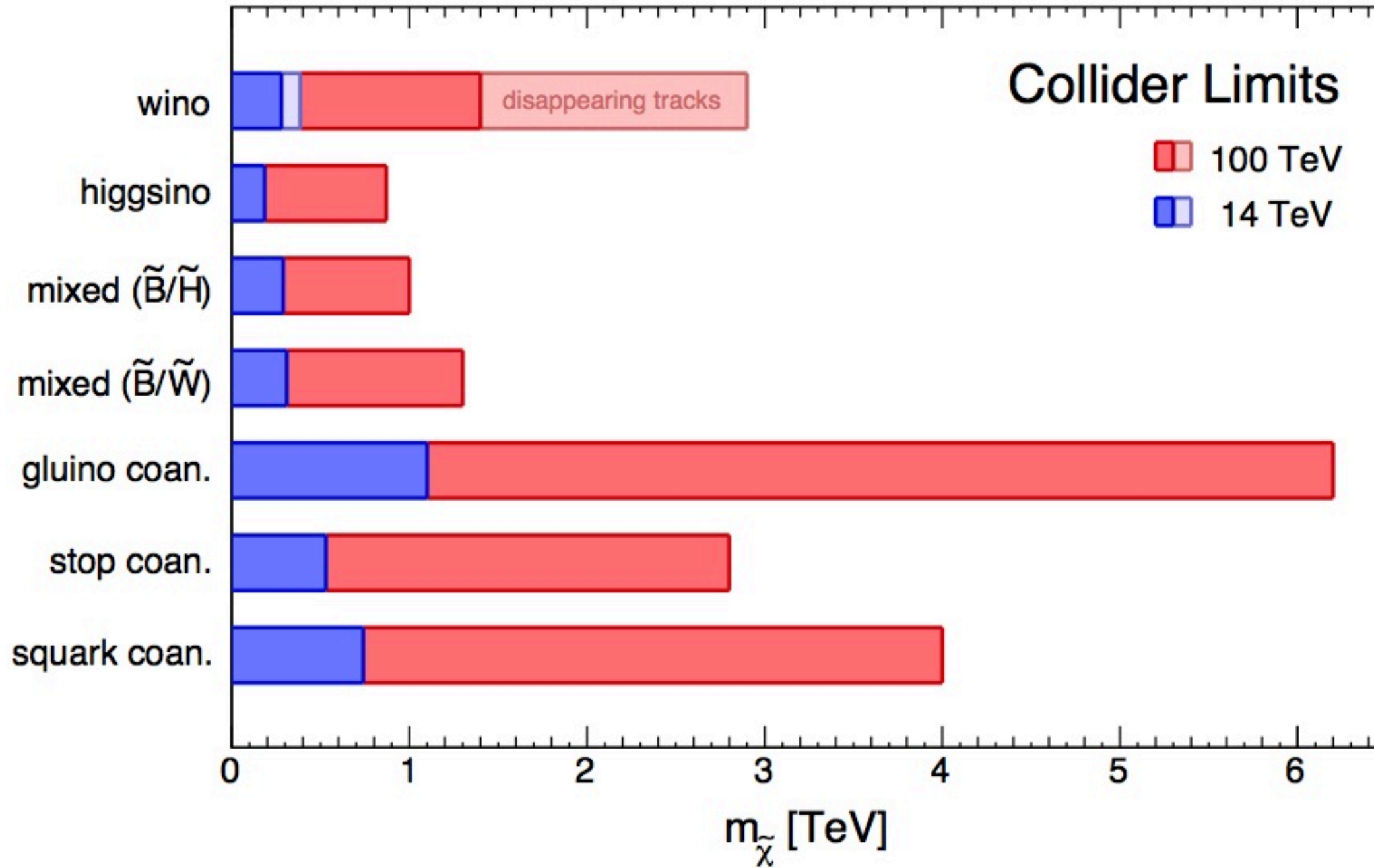
$$\Omega_{\text{wimp}} h^2 \lesssim 0.12$$



$$M_{\text{wimp}} \lesssim 2 \text{TeV} \left( \frac{g}{0.3} \right)^2$$

# DM reach at 100 TeV

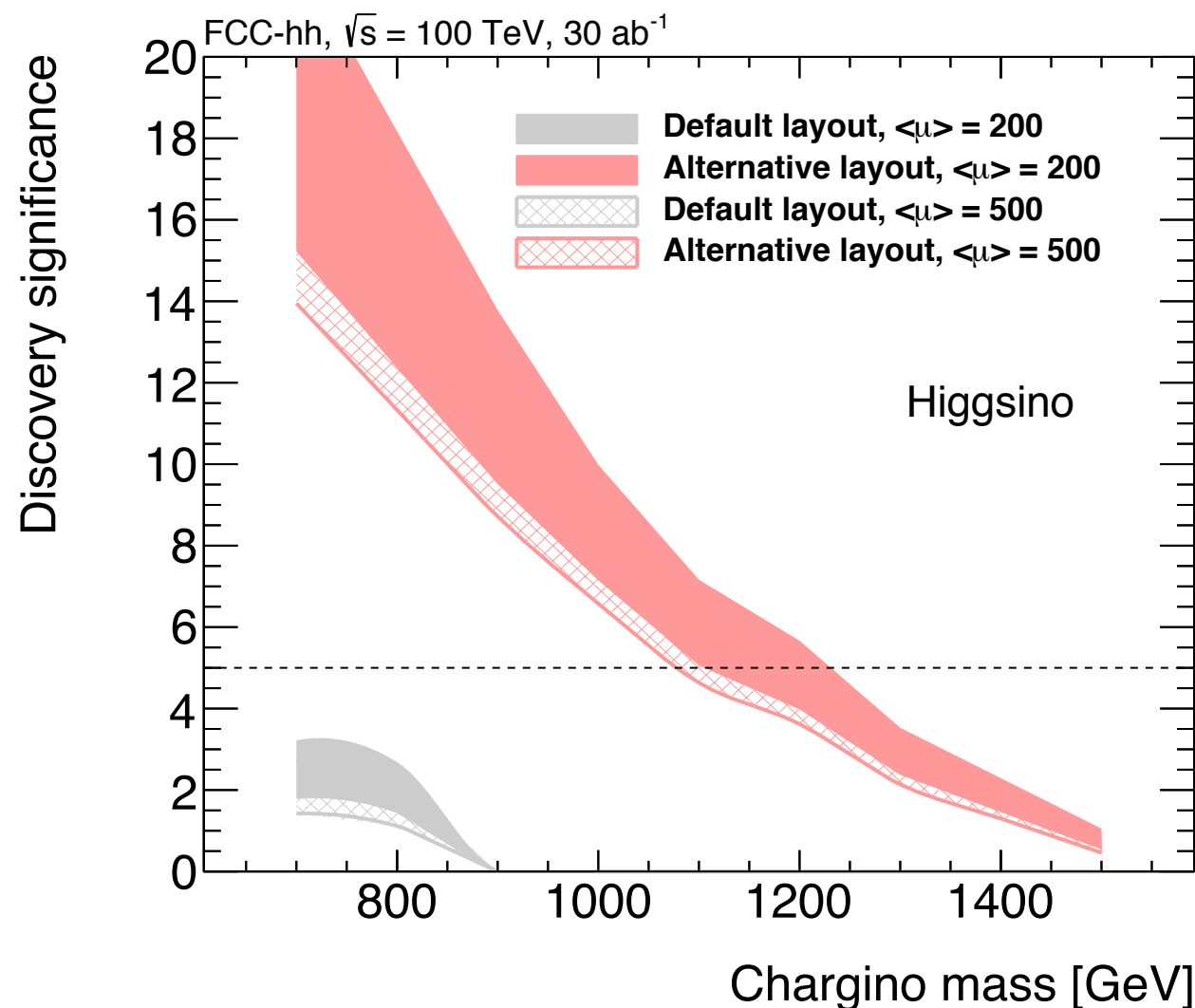
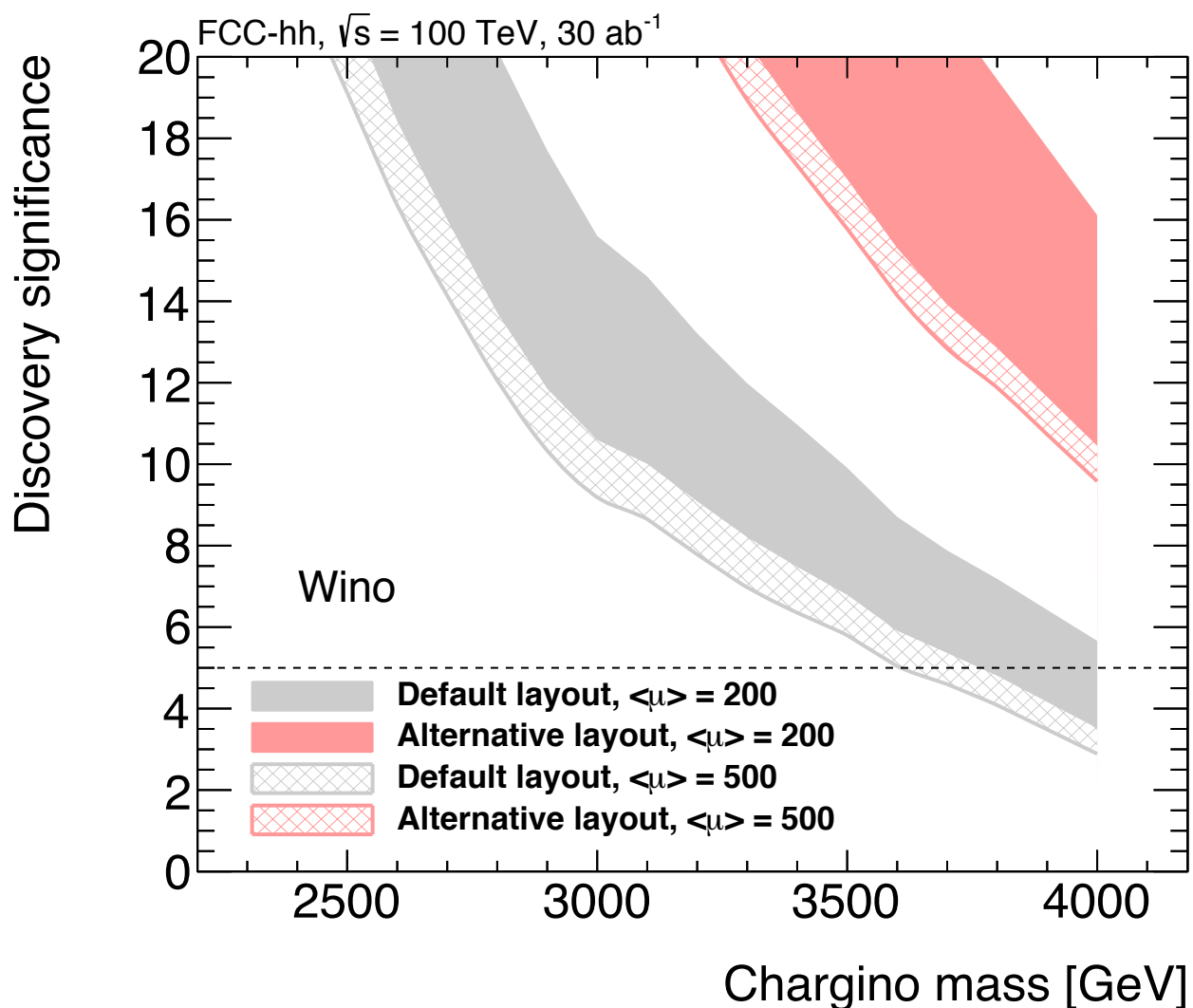
## Early phenomenology studies





New detector performance studies

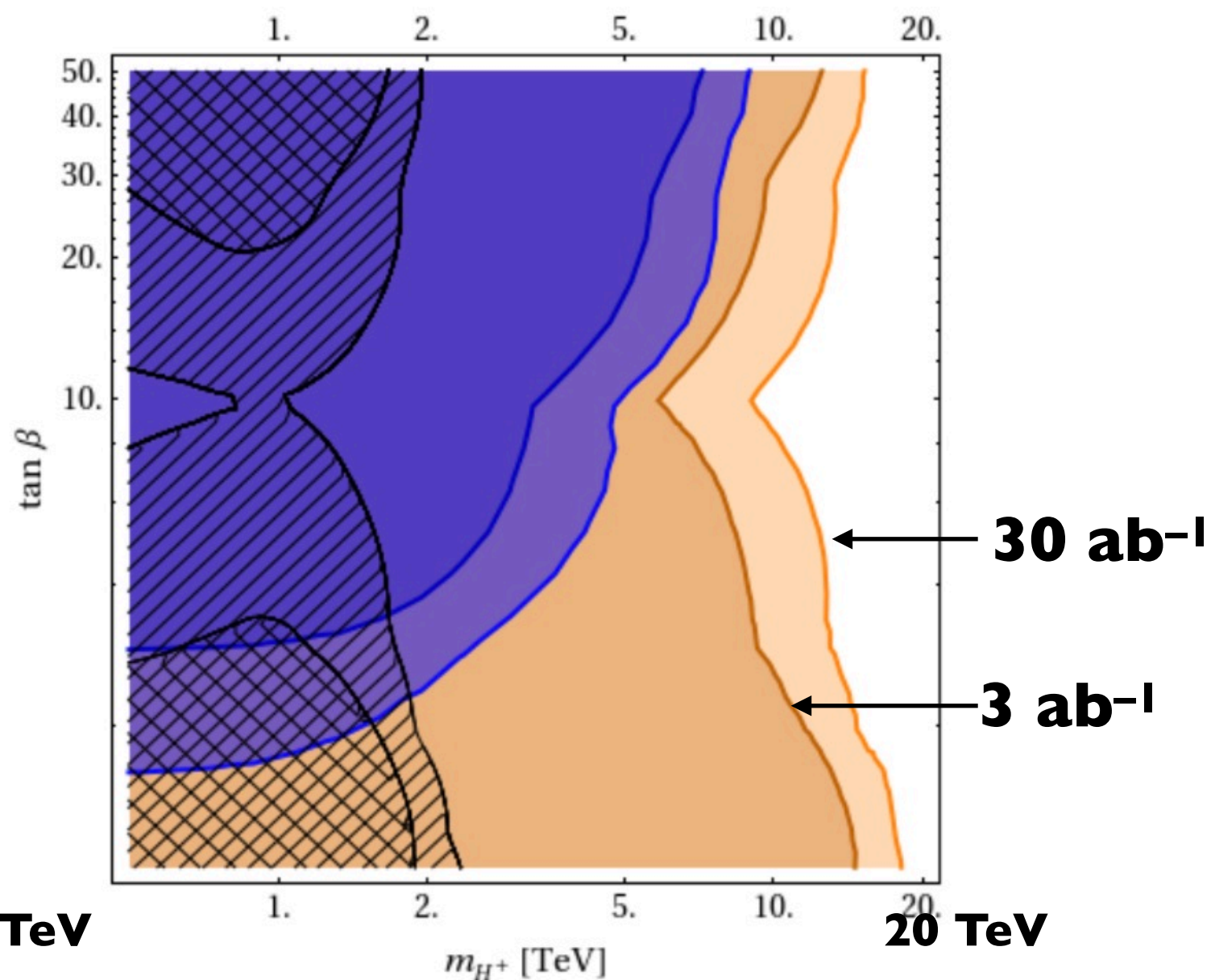
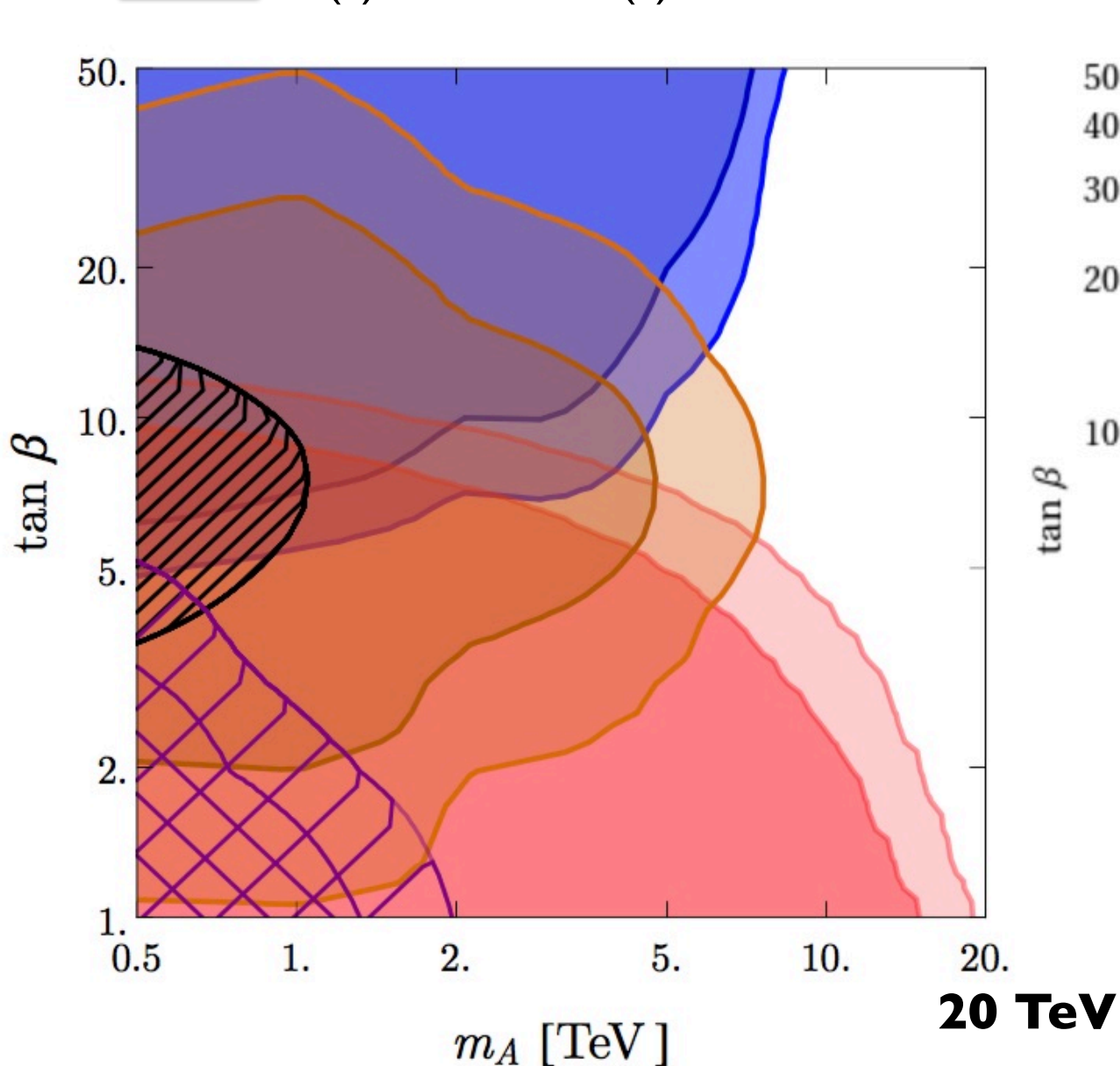
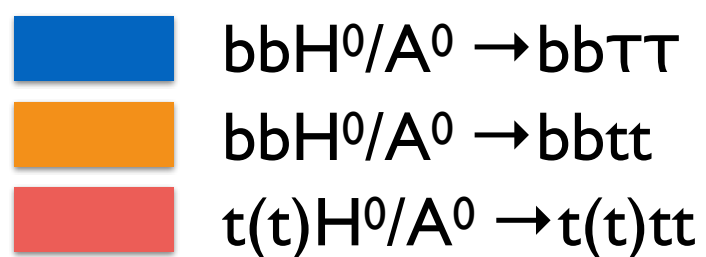
# Disappearing charged track analyses (at ~full pileup)



**=> coverage beyond the upper limit of the thermal WIMP mass range for both higgsinos and winos !!**

$$M_{wimp} \lesssim 2 \text{ TeV} \left( \frac{g}{0.3} \right)^2$$

# MSSM Higgs @ 100 TeV



N. Craig, J. Hajer, Y.-Y. Li, T. Liu, H. Zhang,  
arXiv:1605.08744

J. Hajer, Y.-Y. Li, T. Liu, and J. F. H. Shiu,  
arXiv:1504.07617

# Final remarks

- The study of the SM will not be complete until we clarify the nature of the Higgs mechanism and exhaust the exploration of phenomena at the TeV scale: many aspects are still obscure, many questions are still open.
- The combination of a versatile high-luminosity  $e^+e^-$  circular collider, with a follow-up pp collider in the 100 TeV range, appears like the ideal facility for the post-LHC era
  - *complementary and synergetic precision studies of EW, Higgs and top properties*
  - *energy reach to allow direct discoveries at the mass scales possibly revealed by the precision measurements*
  - *flavor factory at the Z pole, heavy ions and ep collisions: extremely diversified program => broad community engagement*