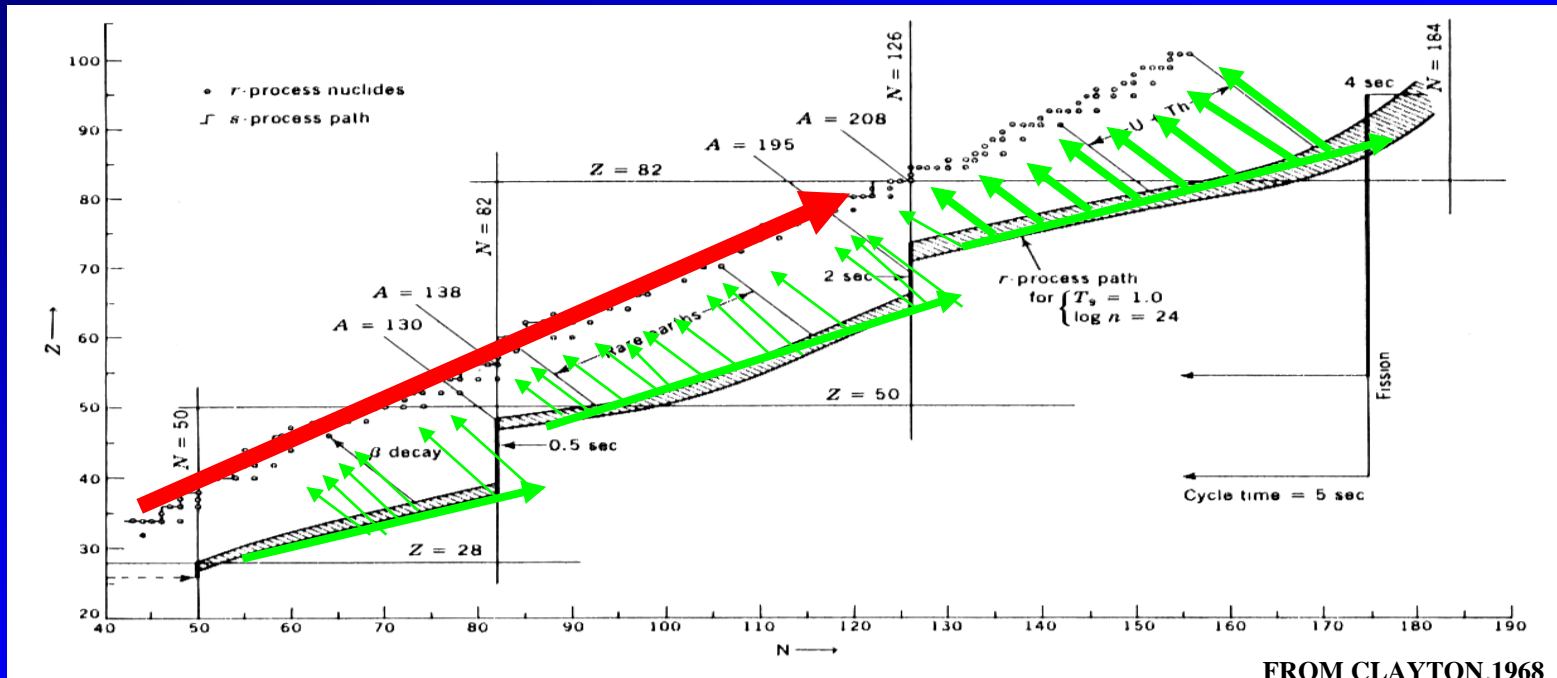


The importance of the $^{13}\text{C}(\alpha,n)^{16}\text{O}$ reaction in Asymptotic Giant Branch stars

Sergio Cristallo

INAF – Osservatorio Astronomico d'Abruzzo

INFN – Sezione di Perugia



***s*-process**

$$\tau_{\beta} \ll \tau_n$$

$$N_n \sim 10^7 \text{ n/cm}^3$$

***r*-process**

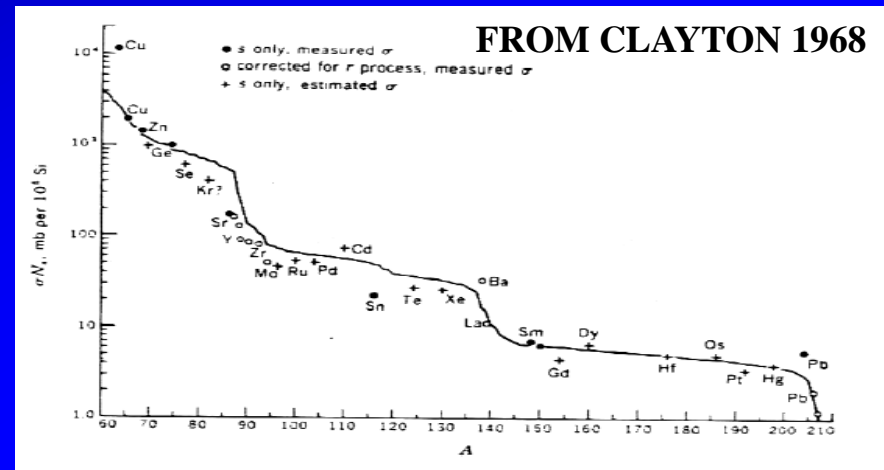
$$\tau_{\beta} \gg \tau_n$$

$$N_n > 10^{22} \text{ n/cm}^3$$



s-process

Easy to be reproduced with an exponential distribution of neutron exposures.



Moreover, since neutron flow reaches equilibrium between nuclei with magic neutron numbers, the product of the Maxwellian averaged stellar (n,γ) cross section of a nuclide, $\langle\sigma\rangle$, and its corresponding abundance, N_s , remains almost constant (the difference in the two product is much smaller than the magnitude of either one of them):

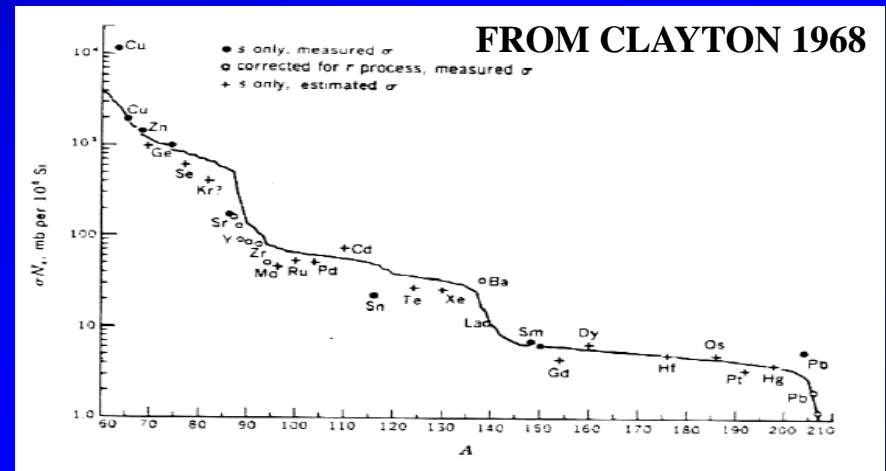
LOCAL APPROXIMATION

$$\langle\sigma\rangle_A N_A \approx \langle\sigma\rangle_{A+1} N_{A+1}$$



s-process

Easy to be reproduced with an exponential distribution of neutron exposures.



Moreover, since neutron flow reaches equilibrium between nuclei with magic neutron numbers, the product of the Maxwellian averaged stellar (n,γ) cross section of a nuclide, $\langle\sigma\rangle$, and its corresponding abundance, N_s , remains almost constant (the difference in the two product is much smaller than the magnitude of either one of them):

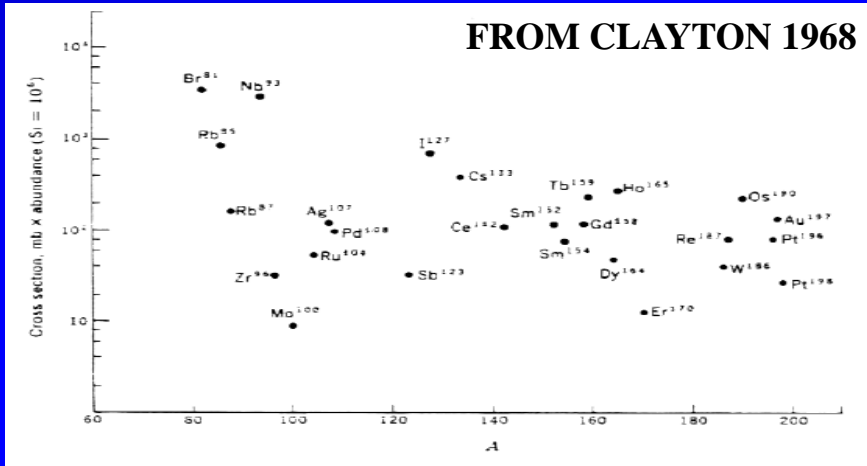
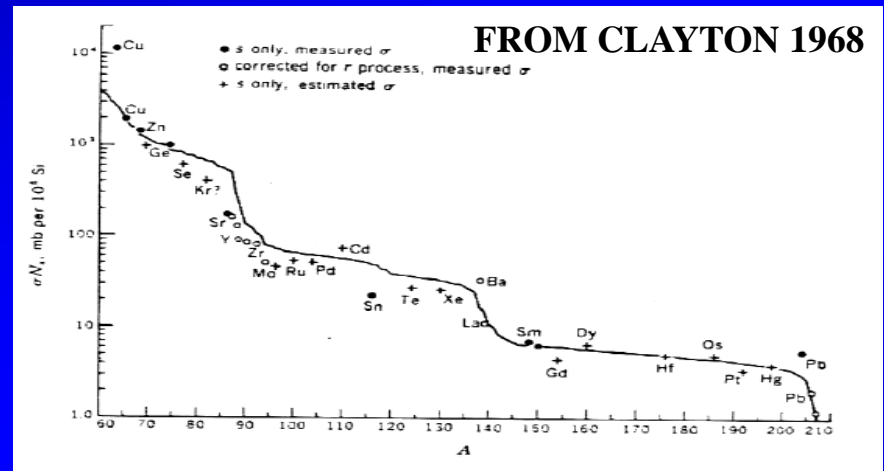
LOCAL APPROXIMATION

$$\langle\sigma\rangle_A N_A \approx \langle\sigma\rangle_{A+1} N_{A+1}$$



s-process

Easy to be reproduced with an exponential distribution of neutron exposures.

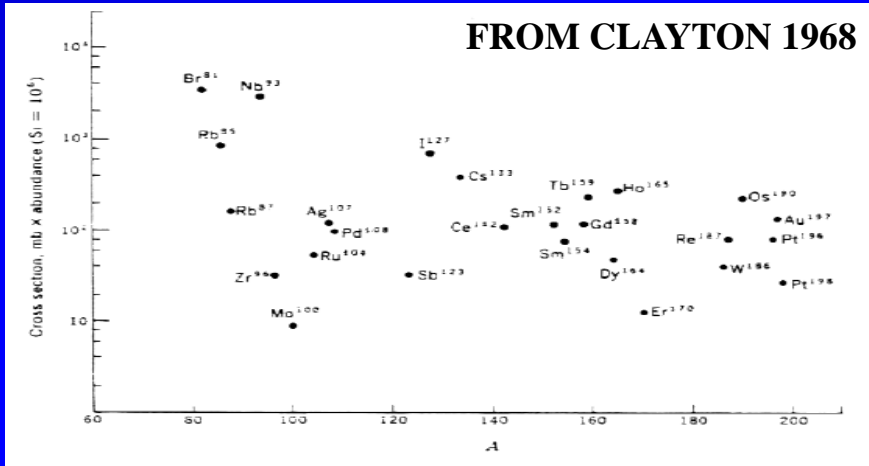
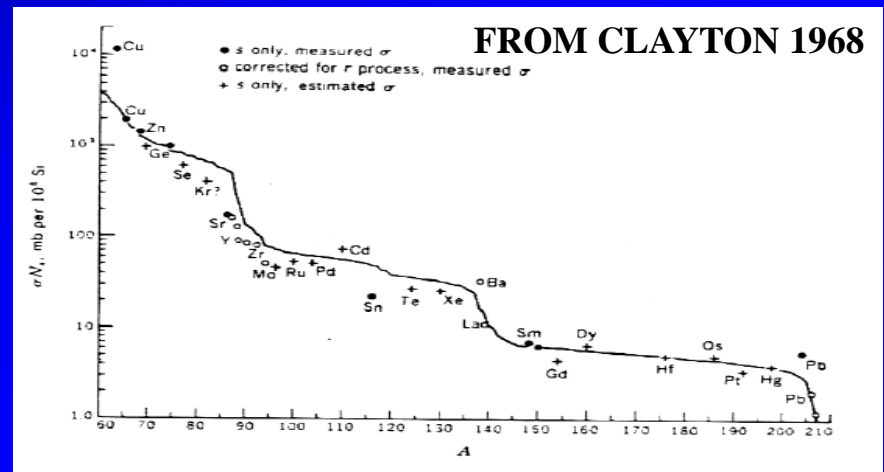


r-process



s-process

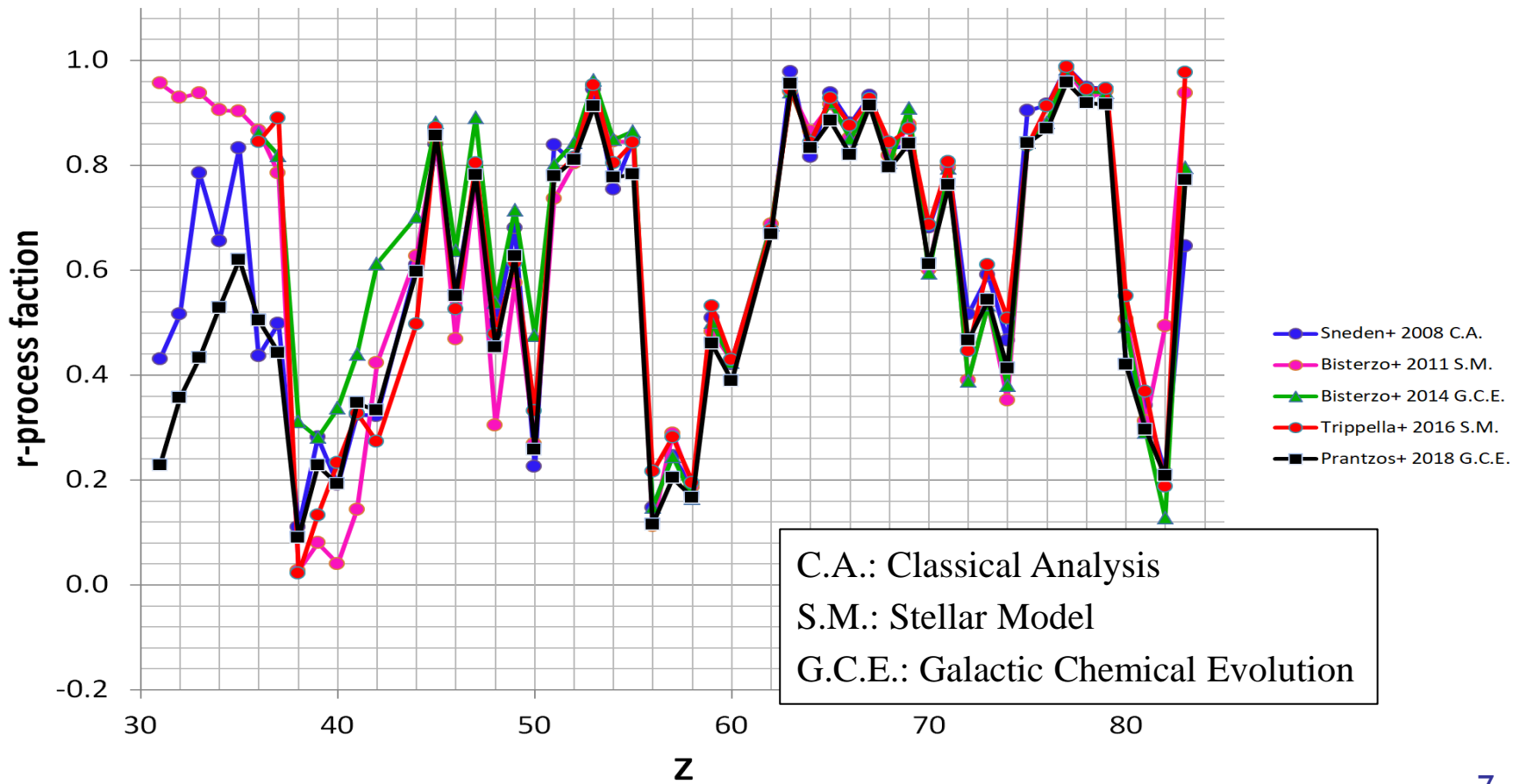
Easy to be reproduced with an exponential distribution of neutron exposures.



r-process

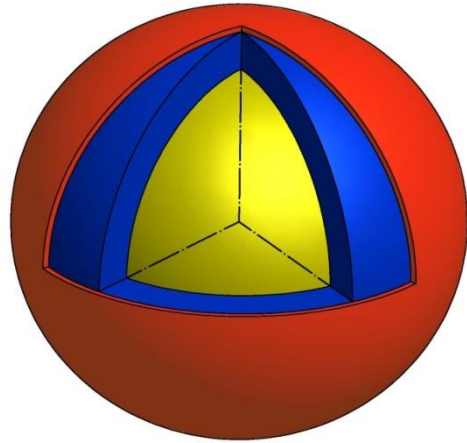
$$r = 1-s$$

r-process residuals from s-process studies



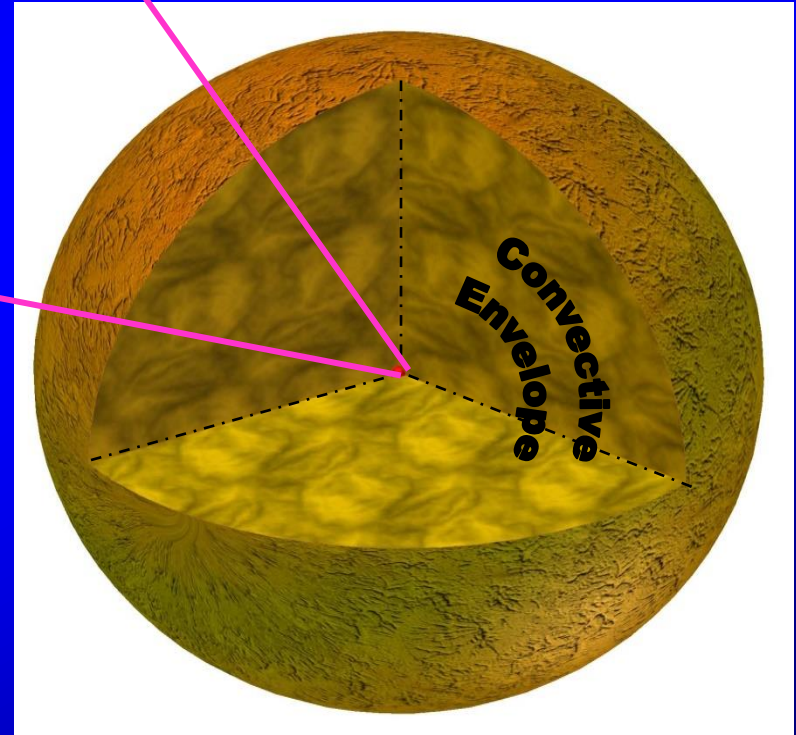
AGB structure

CO Core
He-shell
H-shell

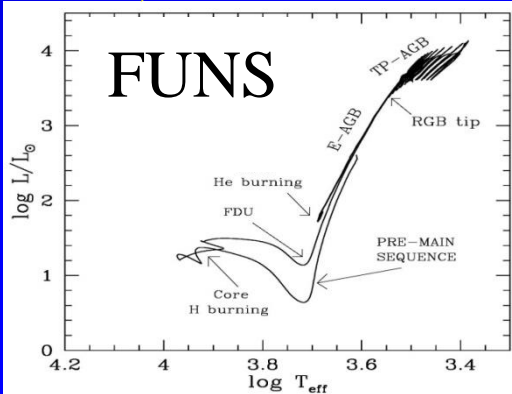


Earth radius ($\sim 10^{-2} R_{\text{SUN}}$)

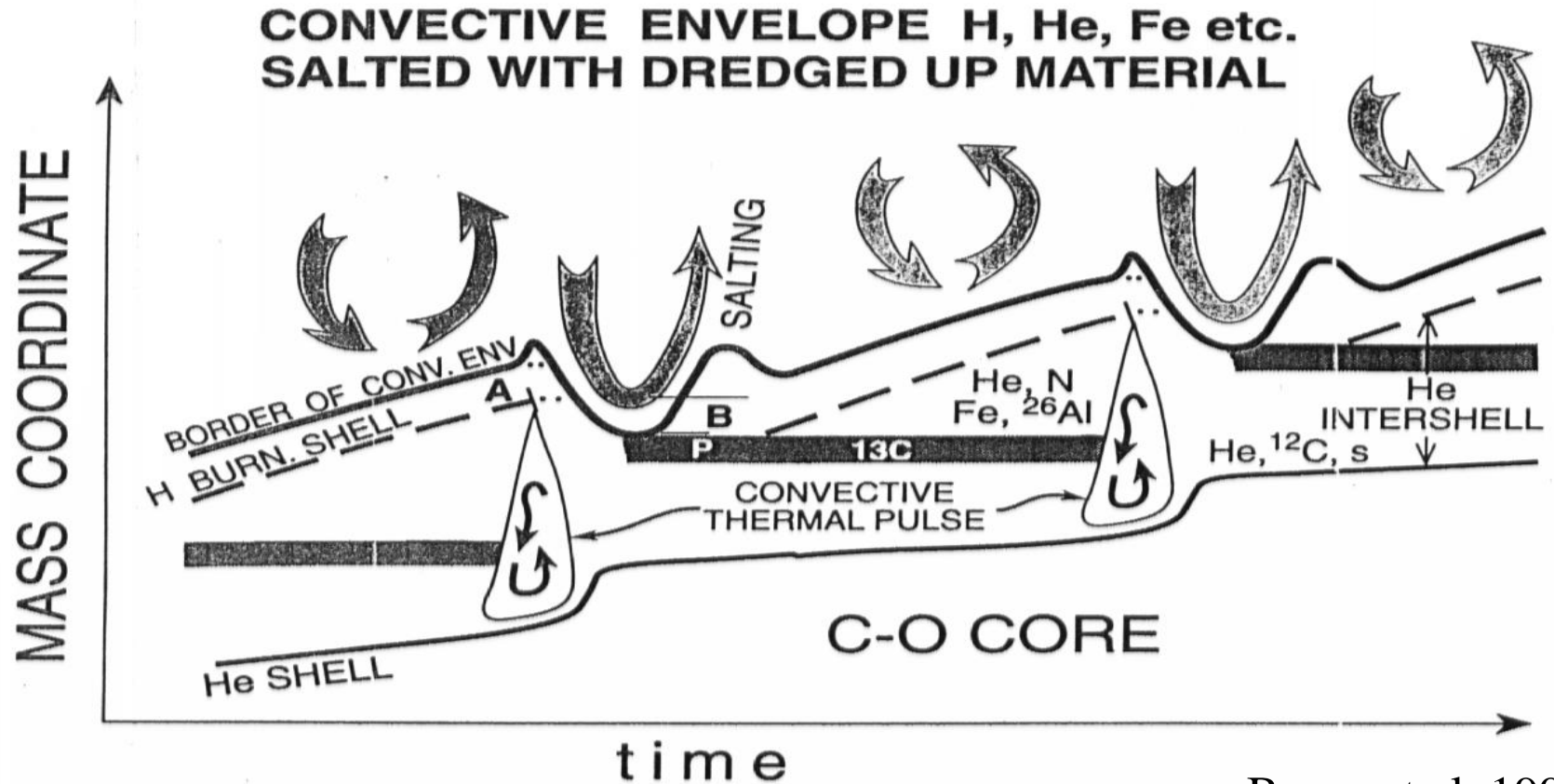
Earth-Sun
($\sim 200 R_{\text{SUN}}$)



Straniero, Gallino & Cristallo 2006

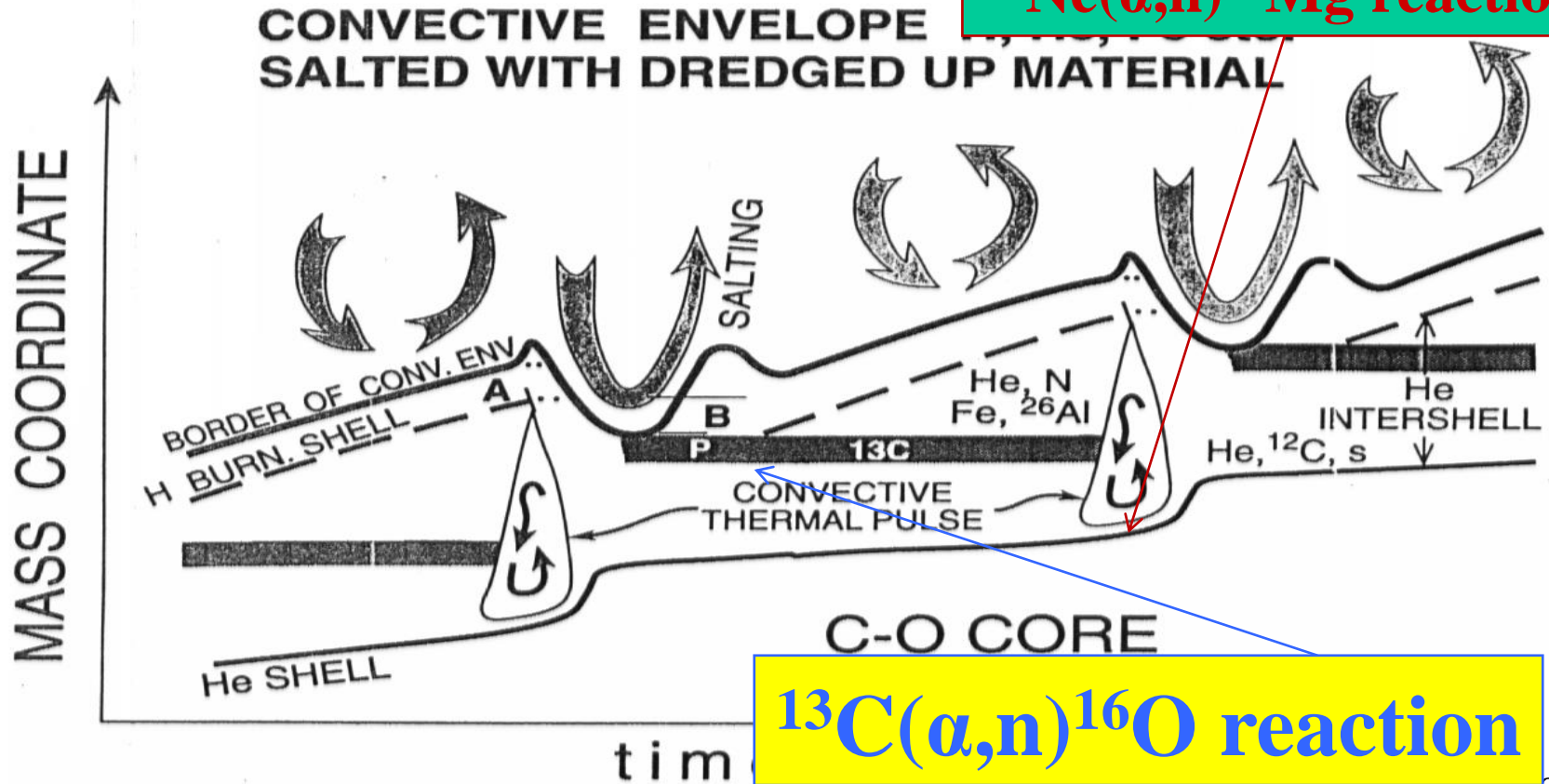


The s-process in AGB stars



The s-process in AGB stars

$^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction

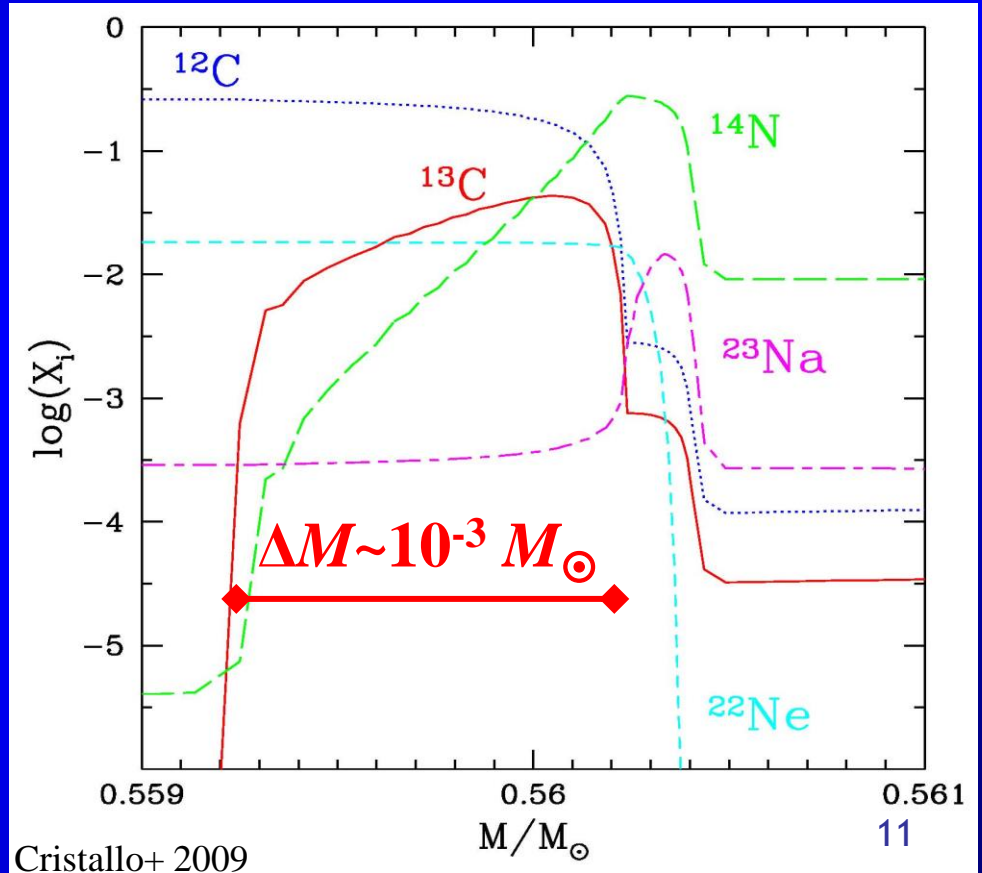


The formation of the ^{13}C pocket

^{13}C -pocket

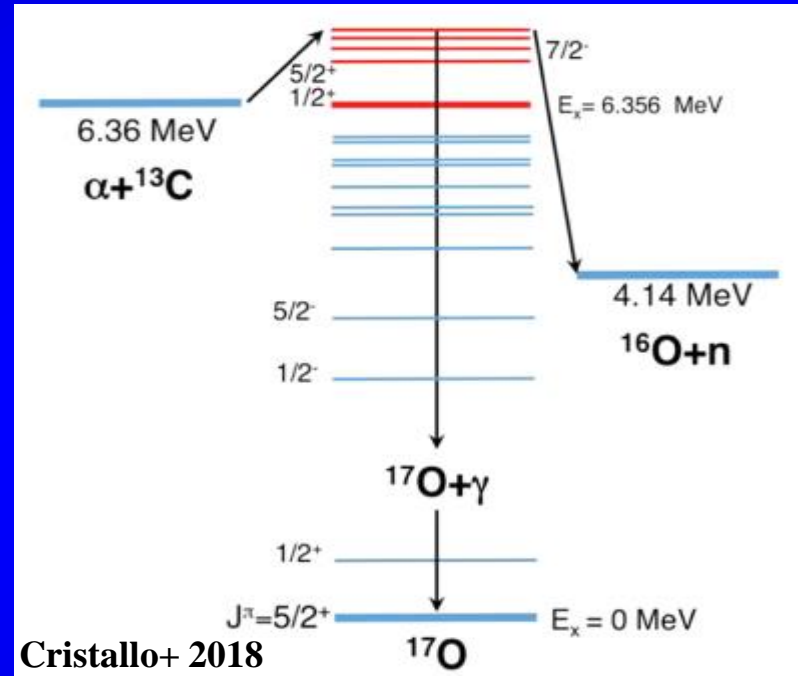
^{14}N -pocket

^{14}N strong neutron
poison via
 $^{14}\text{N}(n,p)^{14}\text{C}$ reaction



Measurements of the $^{13}\text{C}(\alpha, n)^{16}\text{O}$ reaction

- Trippella+ 2017
- Avila+ 2015
- La Cognata+ 2013
- Xu+ 2013
- La Cognata+ 2012
- Guo+ 2012
- Heil+ 2008
- Kubono+ 2003
- Angulo+ 1999
- Drotleff+ 1993
- ...



- S. Urlass: $^{16}\text{O}(n, \alpha)^{13}\text{C}$ (n_TOF)
- G.F. Ciani: $^{13}\text{C}(\alpha, n)^{16}\text{O}$ (LUNA)

Reference rate (most recent direct measure): Heil+ 2008

Upper case: *1.5 - *2

Lower case: *0.67 - *0.5

$$\sigma(E) = \frac{1}{E} S(E) e^{-2\pi\eta}$$

Substantial scatter of existing data, showing a broad (up to a factor of 2) range of absolute values for the astrophysical S-factor

Explored stellar models

- $M=1.5 M_{\text{SUN}}$ [Fe/H]=-0.15 Convective ^{13}C burning
- $M=3.0 M_{\text{SUN}}$ [Fe/H]=-0.15 s-process main component
- $M=4.0 M_{\text{SUN}}$ [Fe/H]=-2.15 Intermediate AGBs in GCs
- $M=1.3 M_{\text{SUN}}$ [Fe/H]=-2.85 Proton ingestions at low Z

$$[A/B] = \log(N_A/N_B)_{\text{STAR}} - \log(N_A/N_B)_{\text{SUN}}$$



The Importance of the $^{13}\text{C}(\alpha,n)^{16}\text{O}$ Reaction in Asymptotic Giant Branch Stars

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G. F. Ciani^{9,10}, F. Mingrone¹¹, L. Piersanti^{1,2} , and D. Vescovi^{2,10}

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⁴ Dipartimento di Fisica e Astronomia, Università di Bologna, Via Iriero 46, Bologna, Italy

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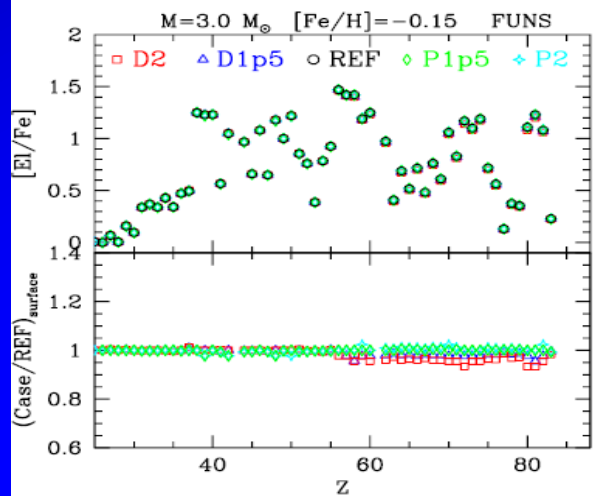
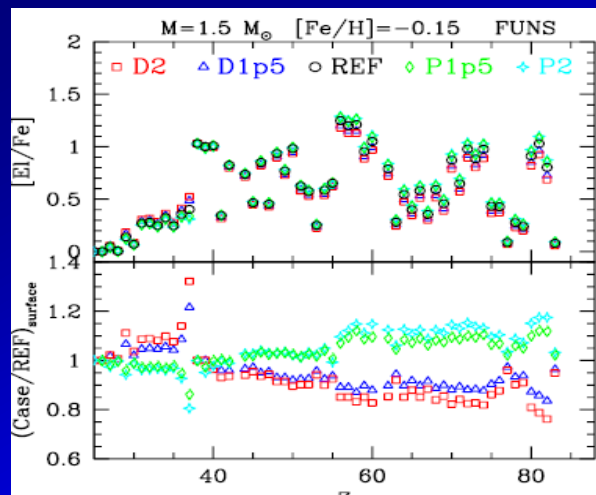
Explored stellar models

Showing a
absolute
factor

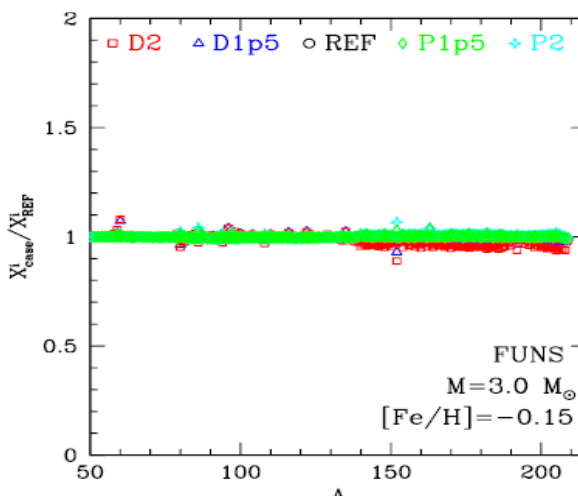
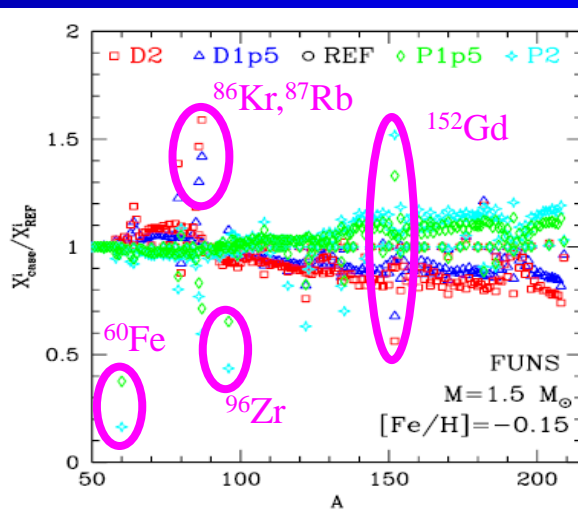
- $M=1.5 M_{\text{SUN}}$ $[\text{Fe}/\text{H}]=-0.15$ Convective ^{13}C burning
- $M=3.0 M_{\text{SUN}}$ $[\text{Fe}/\text{H}]=-0.15$ s-process main component
- $M=4.0 M_{\text{SUN}}$ $[\text{Fe}/\text{H}]=-2.15$ Intermediate AGBs in GCs
- $M=1.3 M_{\text{SUN}}$ $[\text{Fe}/\text{H}]=-2.85$ Proton ingestions at low Z

$$[A/B]=\log(N_A/N_B)_{\text{STAR}}-\log(N_A/N_B)_{\text{SUN}}$$

ELEMENTS



ISOTOPES

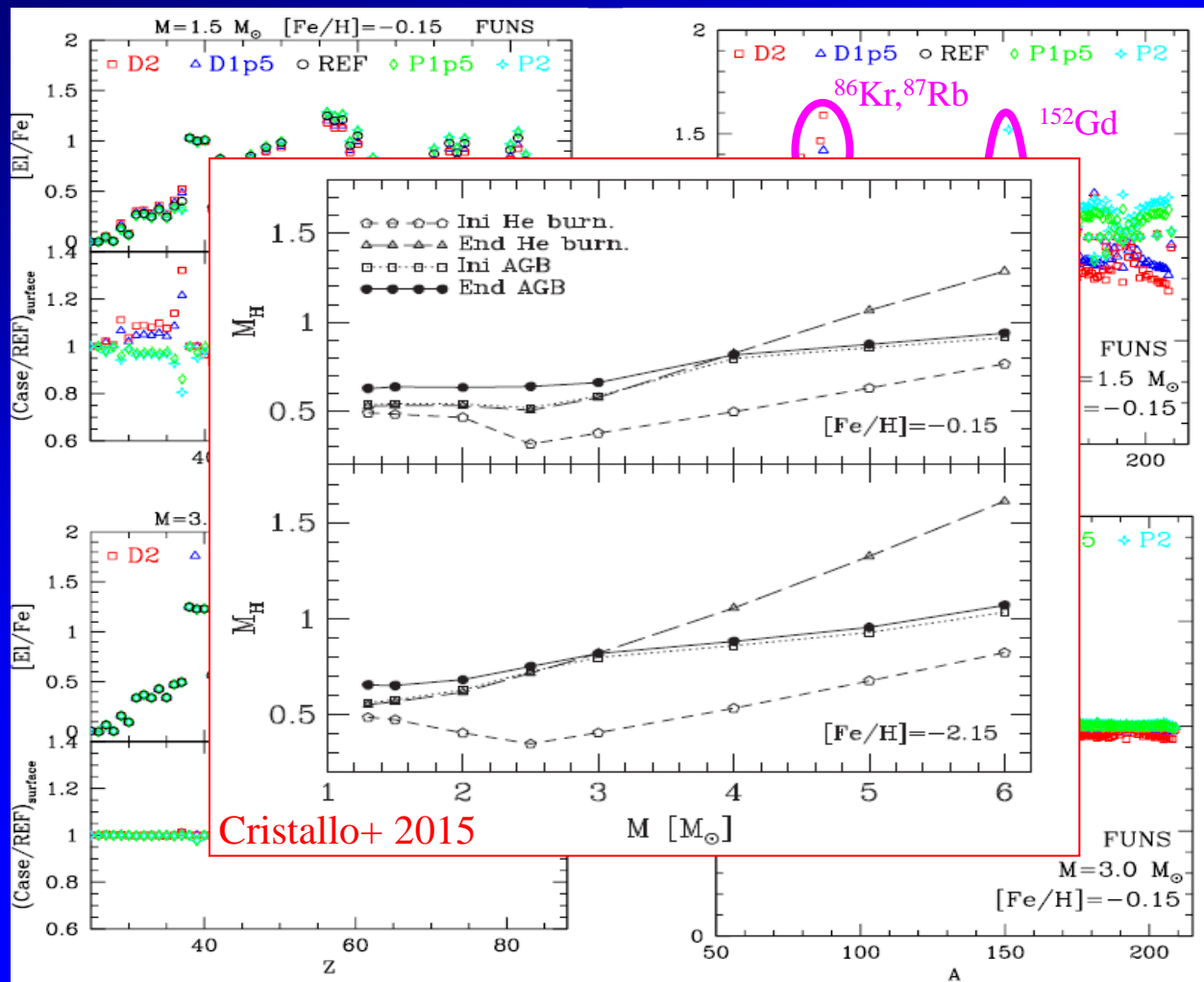


$M=1.5 M_{SUN}$
 $[Fe/H]=-0.15$

$M=3.0 M_{SUN}$
 $[Fe/H]=-0.15$

ELEMENTS

ISOTOPES

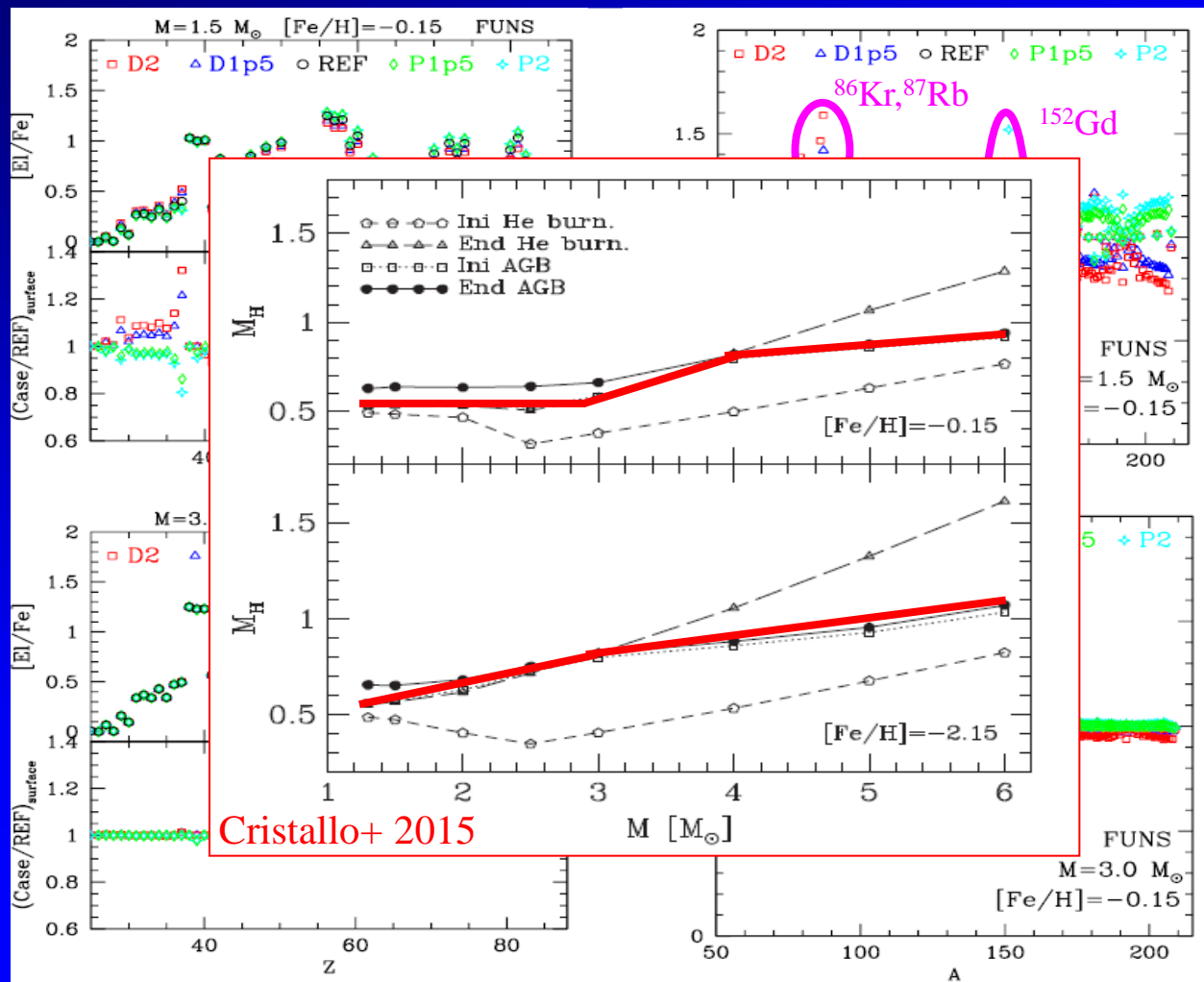


$M = 1.5 M_{\text{SUN}}$
 $[Fe/H] = -0.15$

$M = 3.0 M_{\text{SUN}}$
 $[Fe/H] = -0.15$

ELEMENTS

ISOTOPES

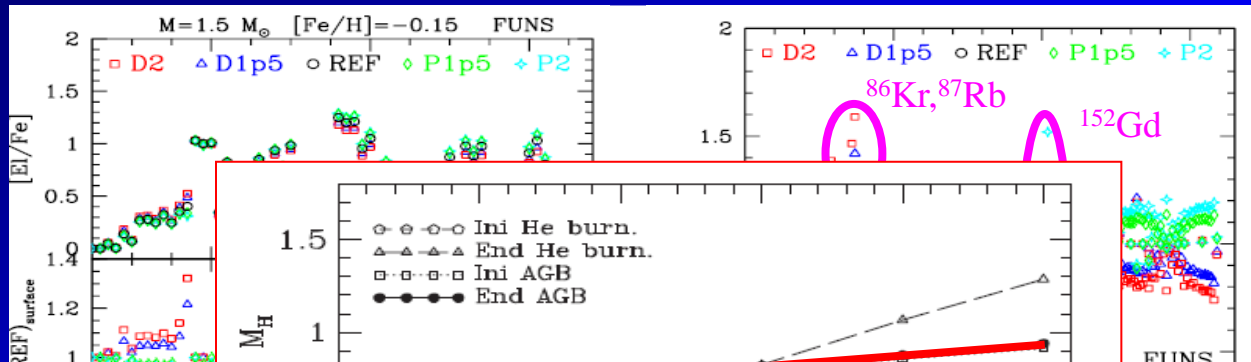


$M=1.5 M_{\text{SUN}}$
 $[Fe/H]=-0.15$

$M=3.0 M_{\text{SUN}}$
 $[Fe/H]=-0.15$

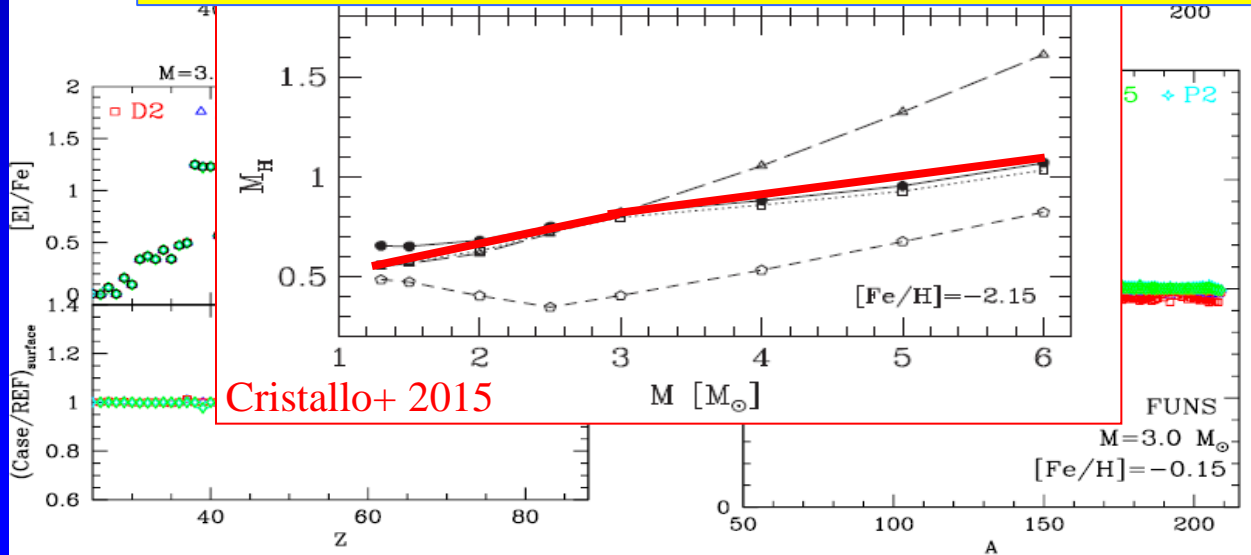
ELEMENTS

ISOTOPES



Same results with NEWTON code @ Perugia

$M=1.5 M_{\text{SUN}}$
 $[\text{Fe}/\text{H}]=-0.15$

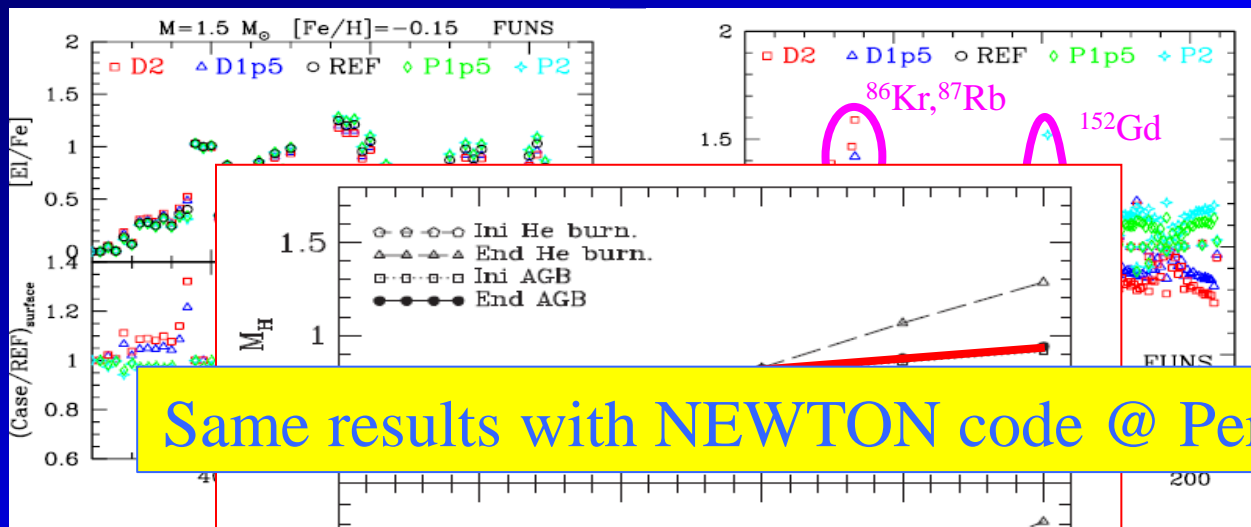


Cristallo+ 2015

$M=3.0 M_{\text{SUN}}$
 $[\text{Fe}/\text{H}]=-0.15$

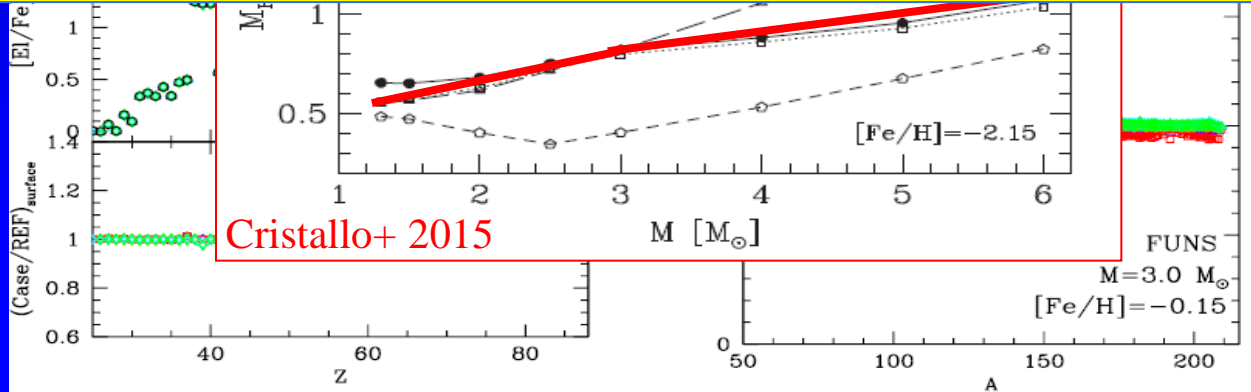
ELEMENTS

ISOTOPES



Same results with NEWTON code @ Perugia

We confirm the results by Guo+ 2012 and Trippella & La Cognata 2017

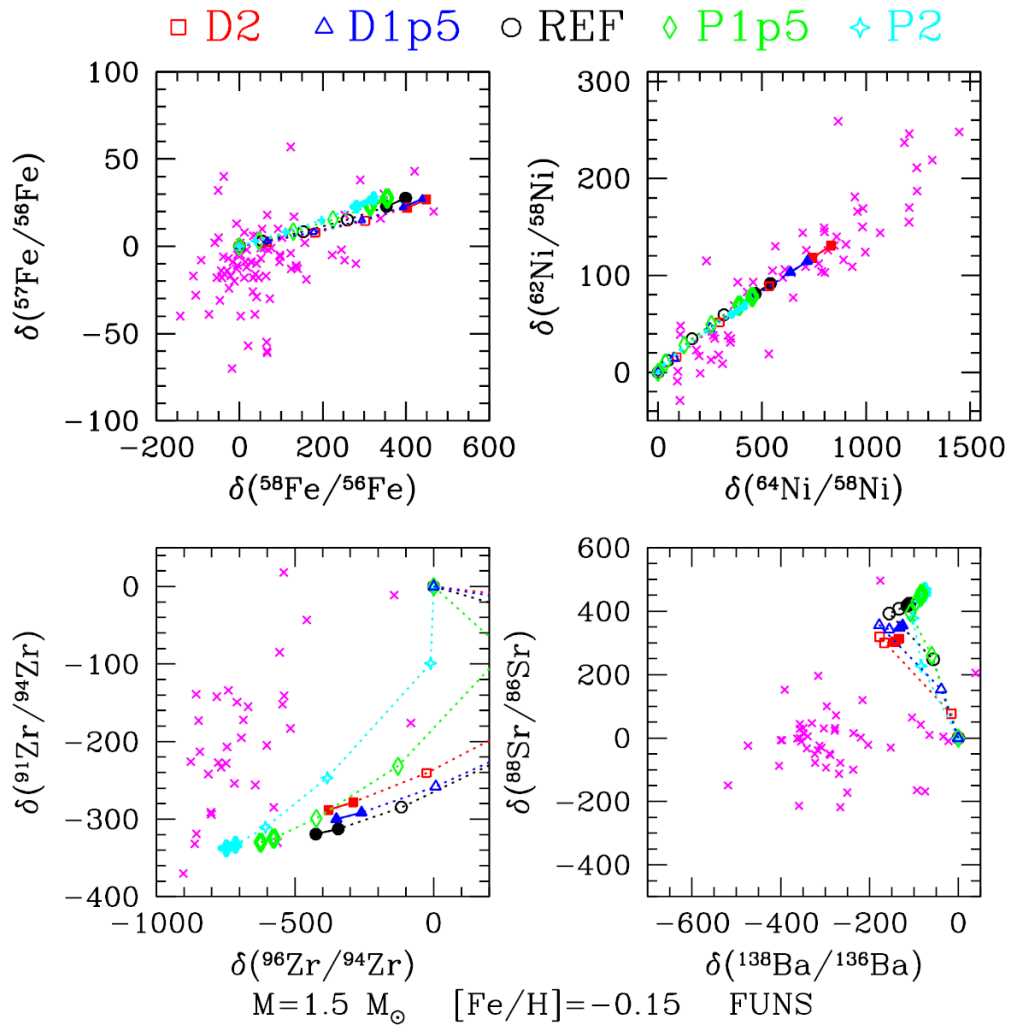


M=1.5 M_{SUN}
 [Fe/H]=-0.15

M=3.0 M_{SUN}
 [Fe/H]=-0.15

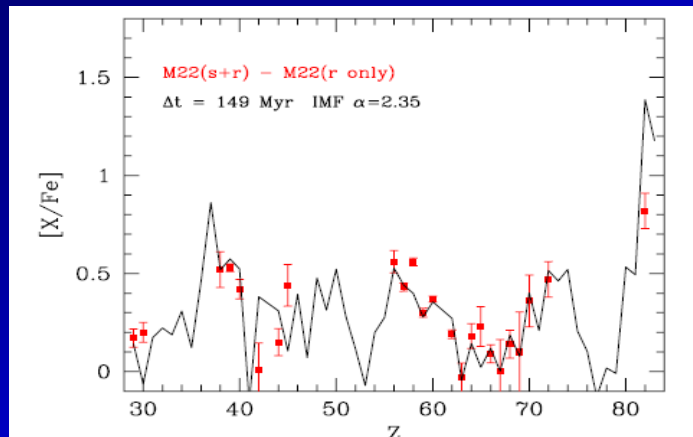
Isotopic anomalies in pre-solar SiC grains

$$\delta(X_i/X_j) \equiv [(X_i/X_j)_{\text{MEASURED}} / (X_i/X_j)_{\text{SUN}} - 1] \times 1000$$

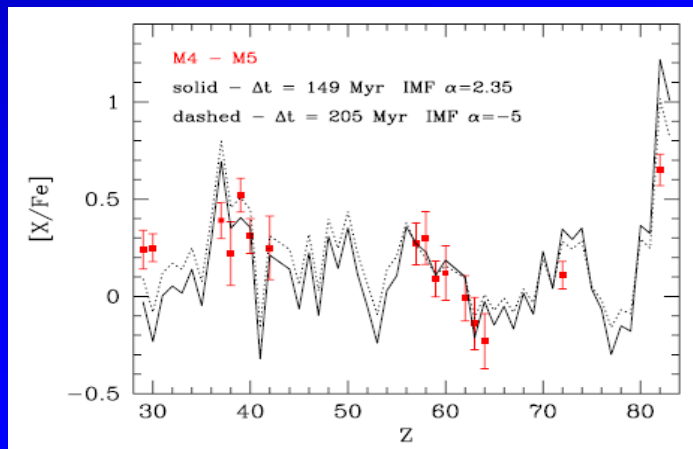


Schonbachler TALK
Davis TALK

s-rich Globular Clusters: the importance of intermediate AGBs

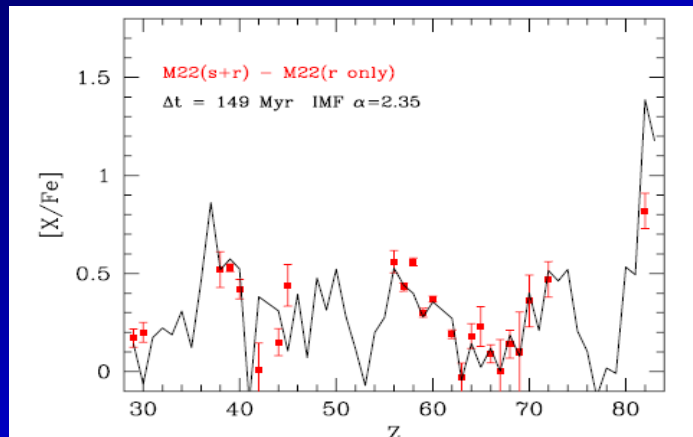


Straniero+ 2014 (Shingles+ 2014)

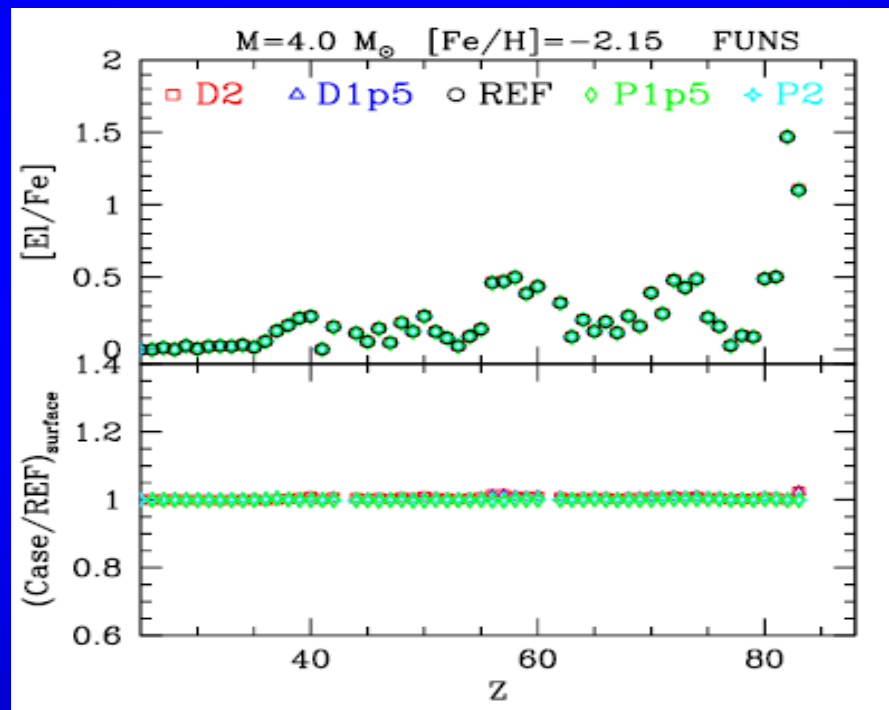
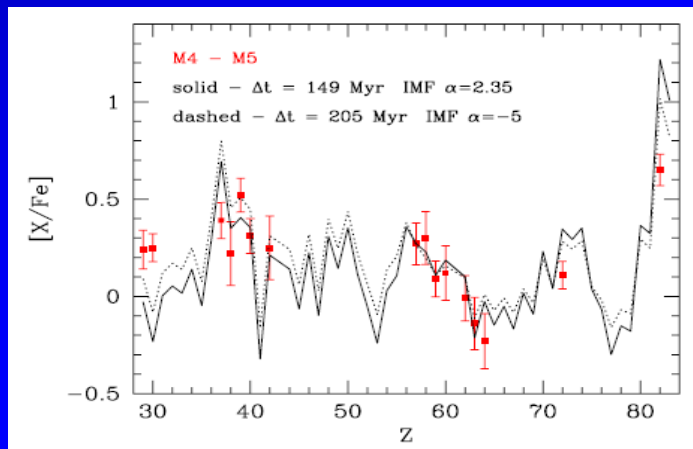


s-rich Globular Clusters: the importance of intermediate AGBs

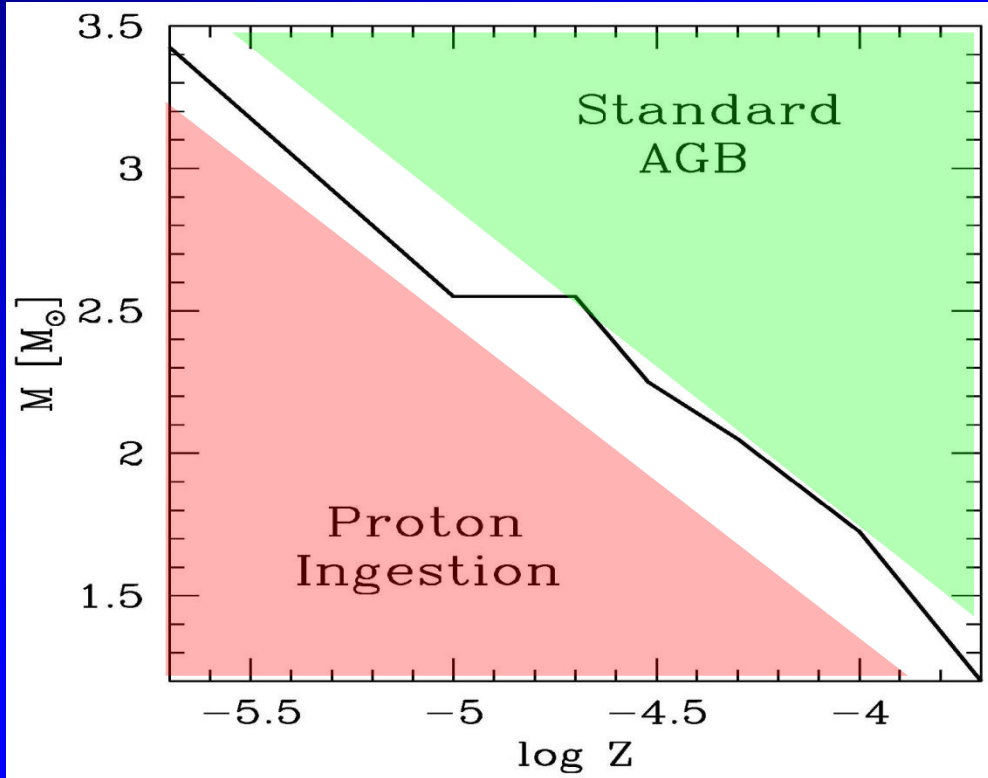
$M=4.0 M_{\text{SUN}}$ $[\text{Fe}/\text{H}]=-2.15$



Straniero+ 2014 (Shingles+ 2014)



Low-metallicity low-mass AGB stars



Hollowell+ 1990

Fujimoto+ 2000

Iwamoto+ 2004

Suda+ 2004

Campbell+ 2007

Cristallo+ 2009

Herwig+ 2014

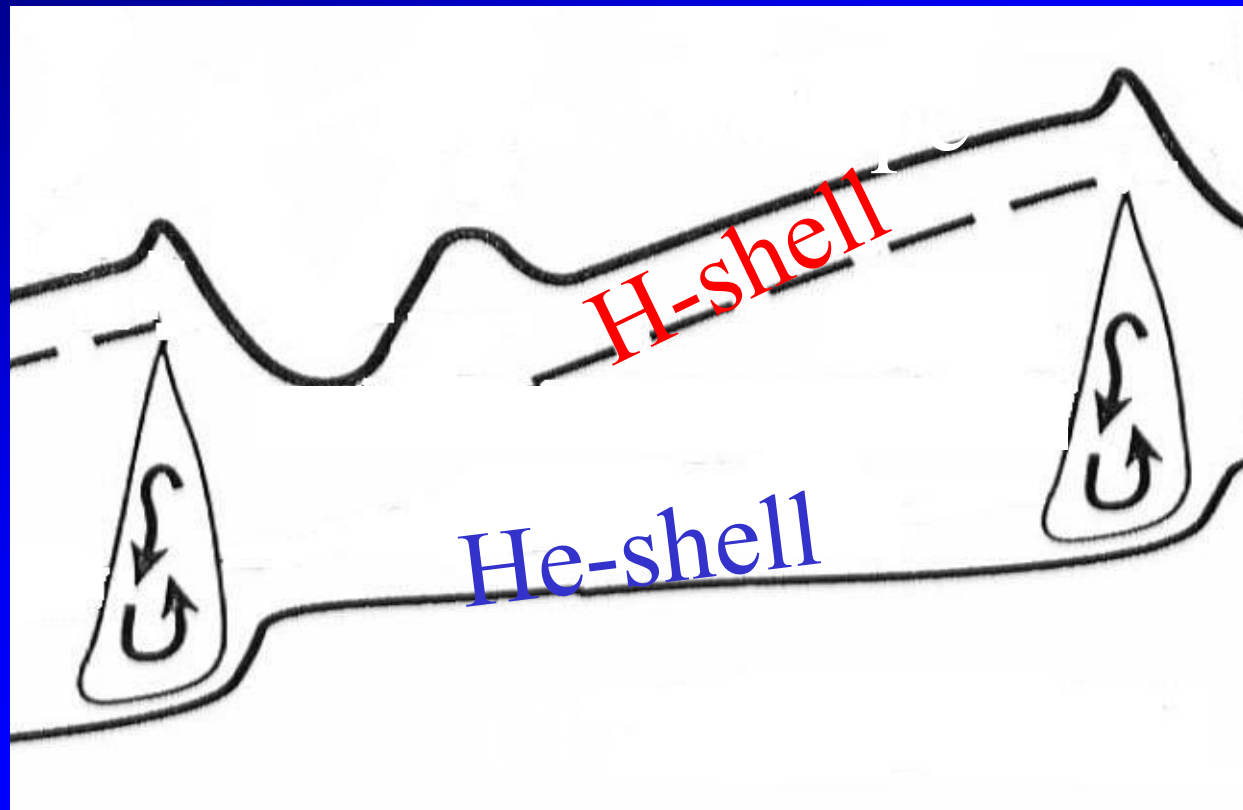
Dardelet+ 2015

Hempel+ 2016

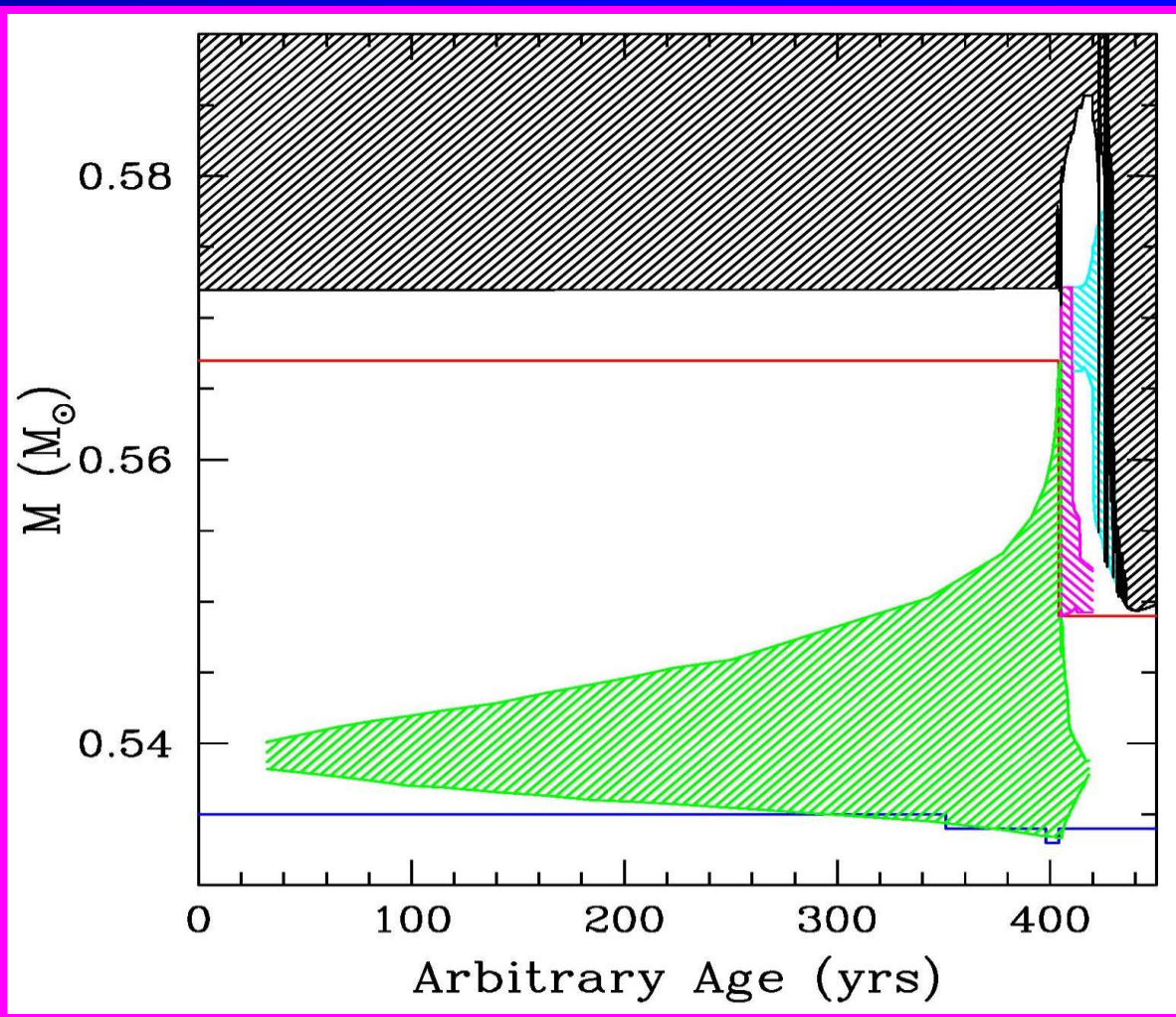
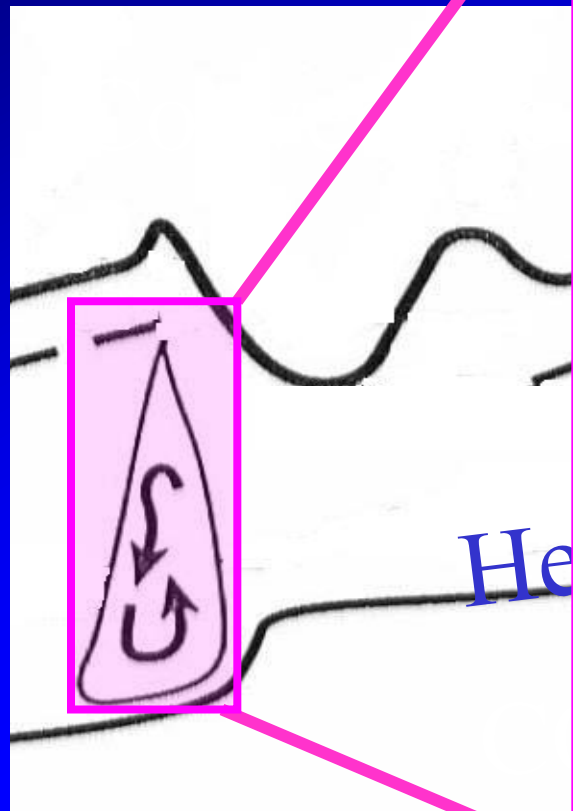
Cristallo+ 2016

$M=1.3 M_{\text{sun}}$

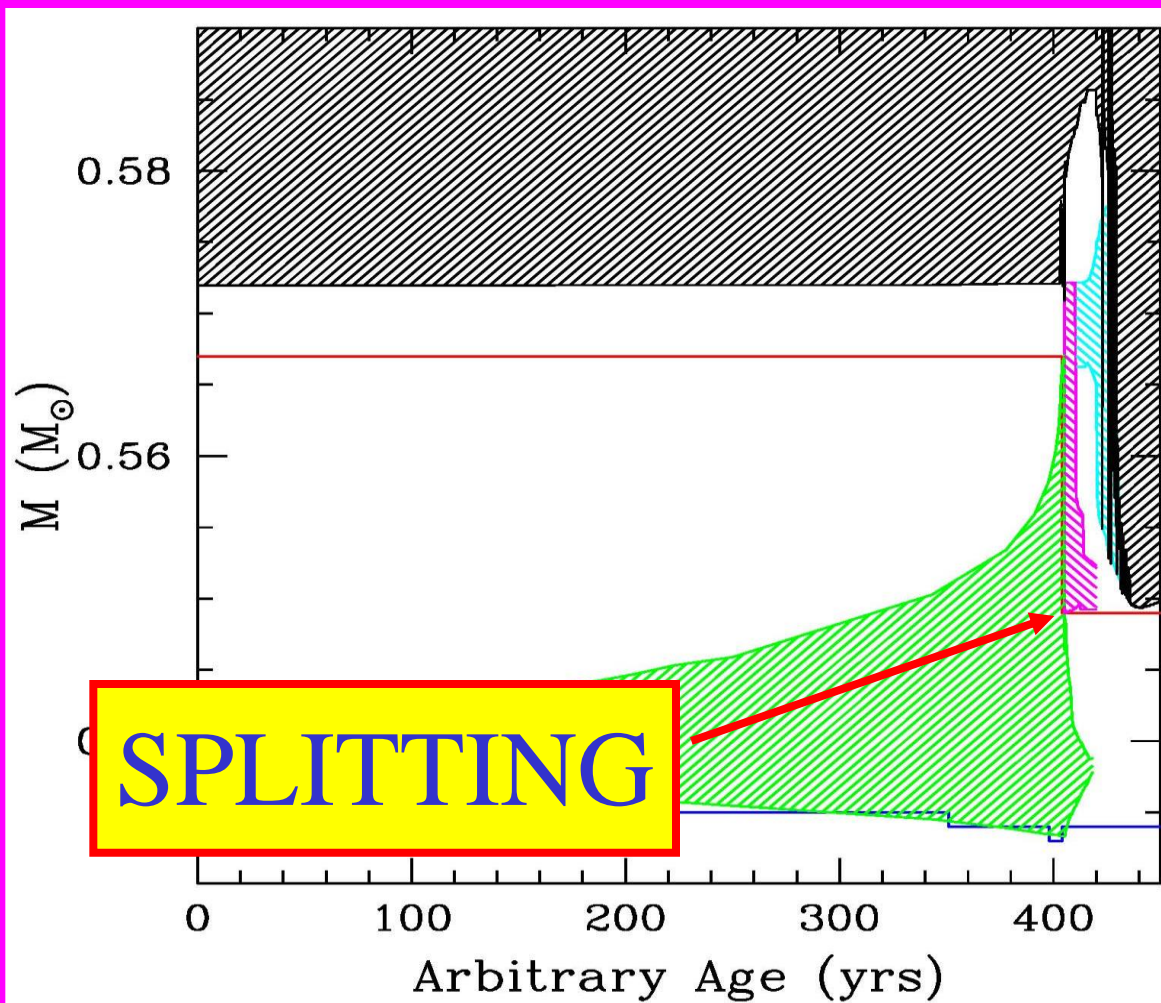
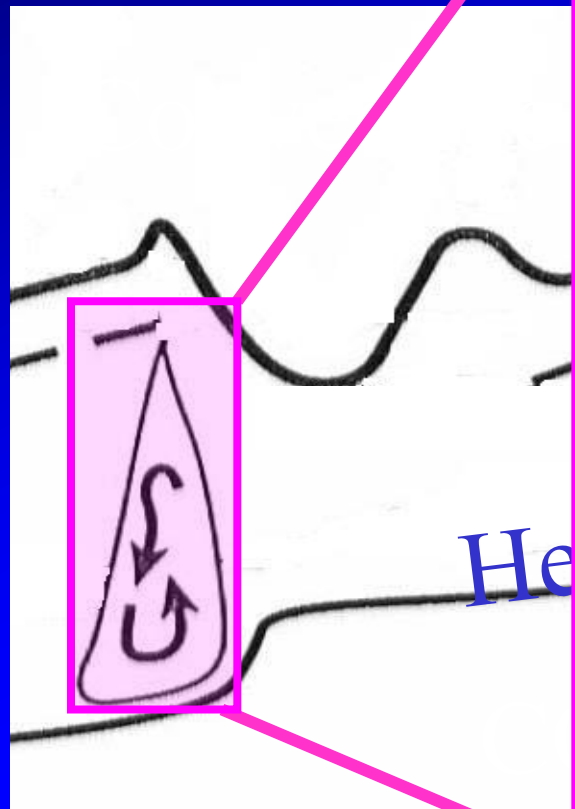
$[\text{Fe}/\text{H}]=-2.85$



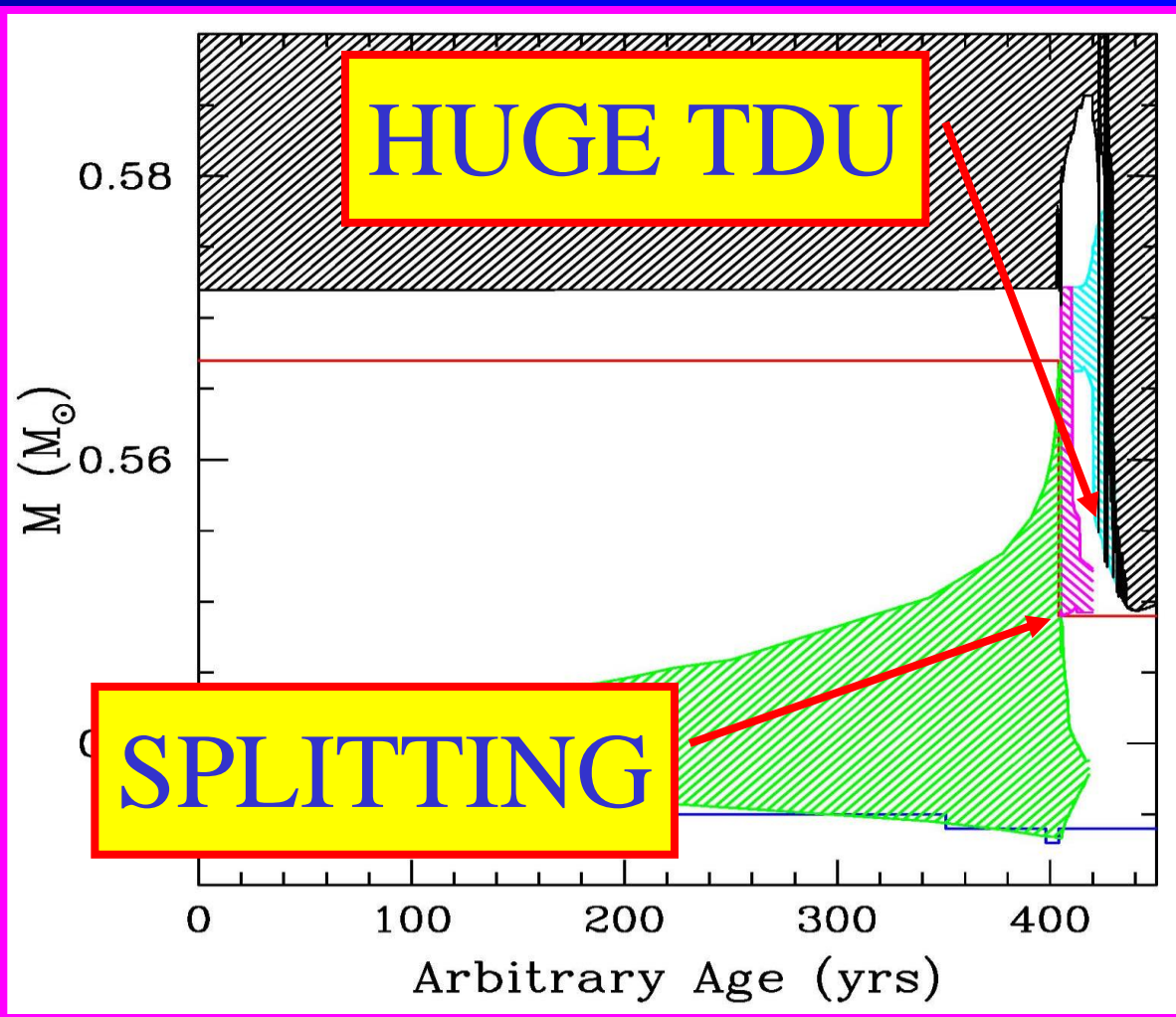
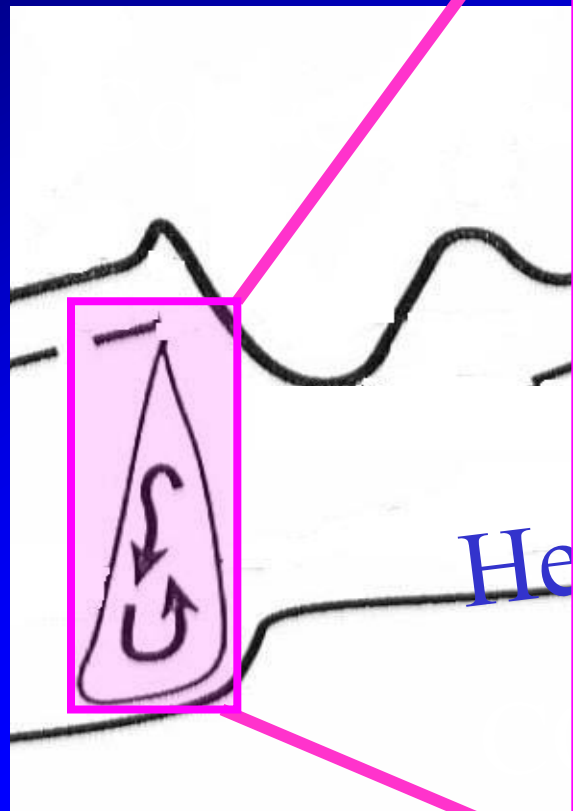
M=1.3 M_⊙



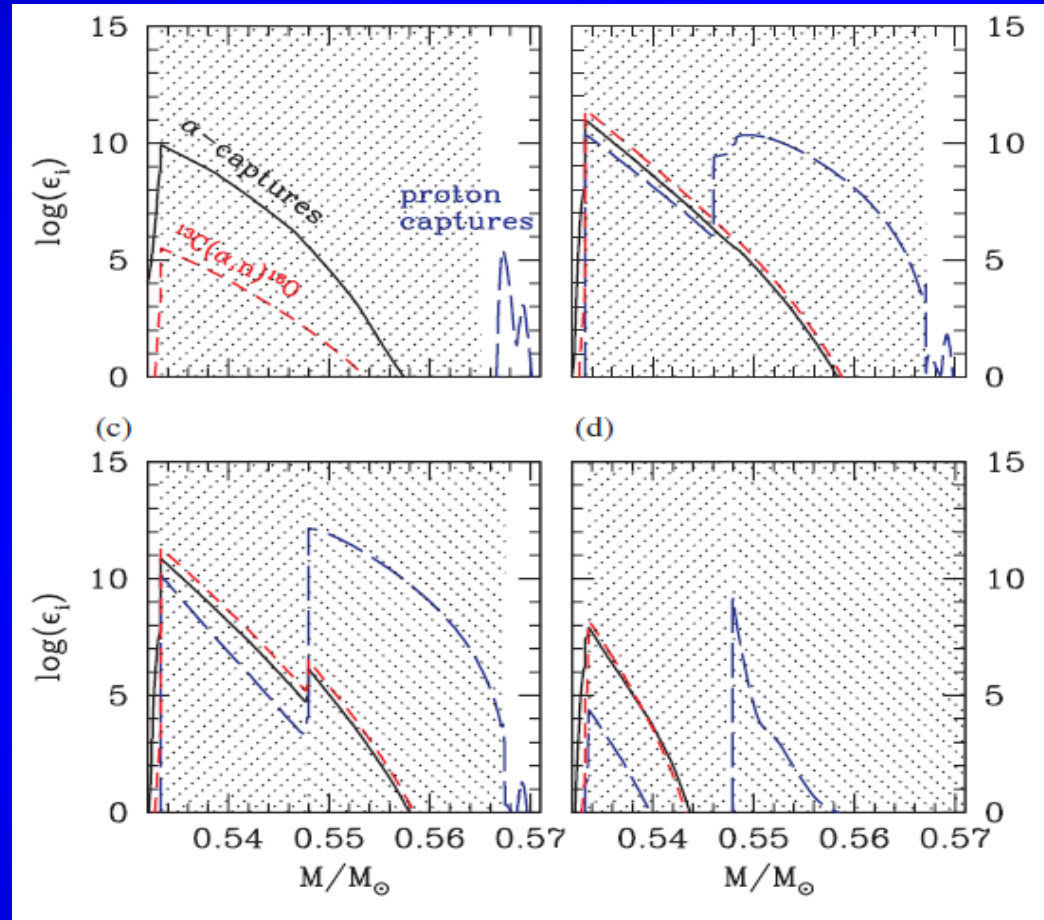
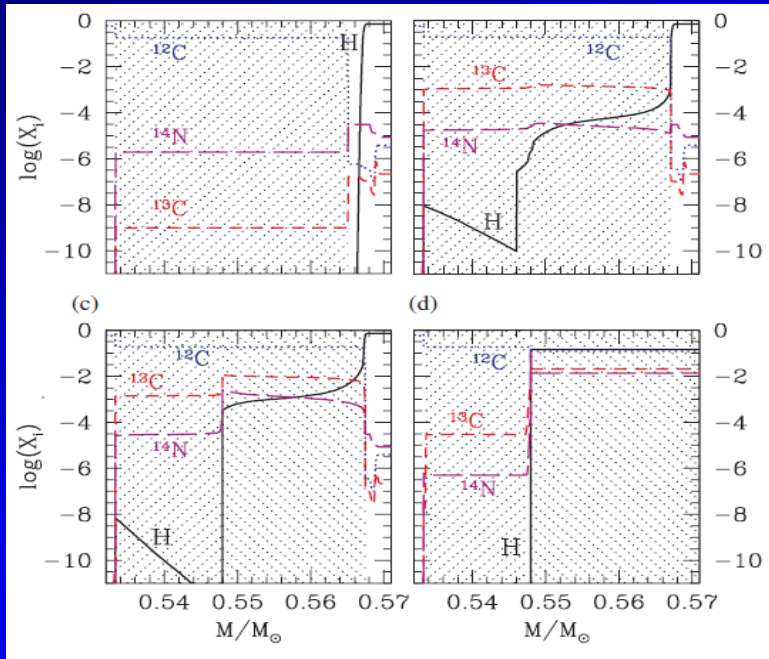
M=1.3 M_⊙



M=1.3 M_⊙



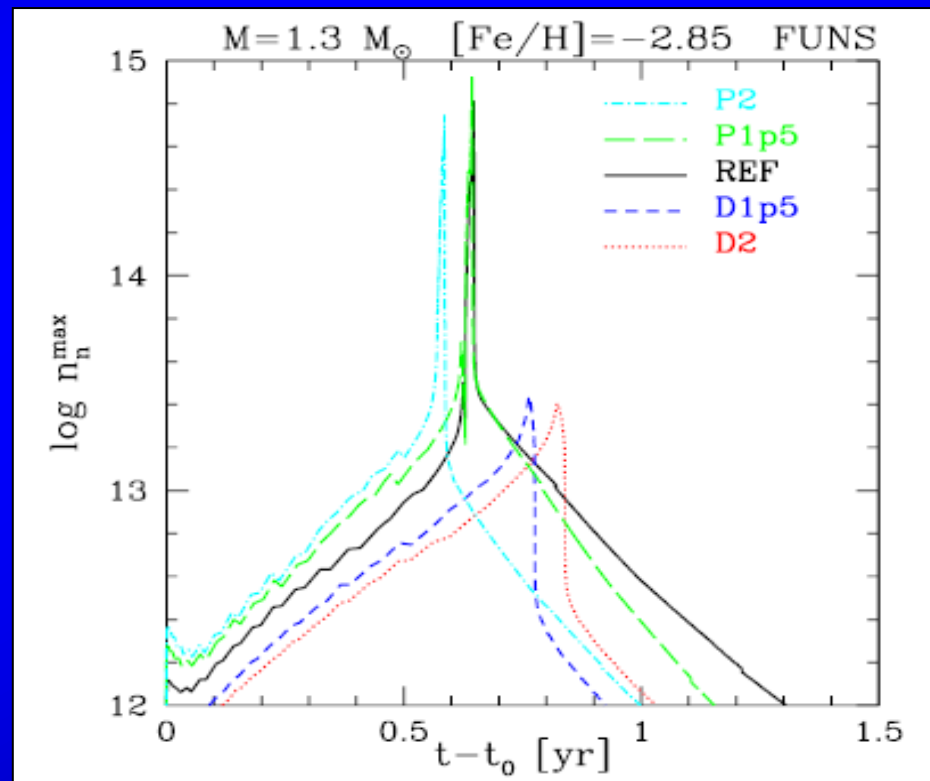
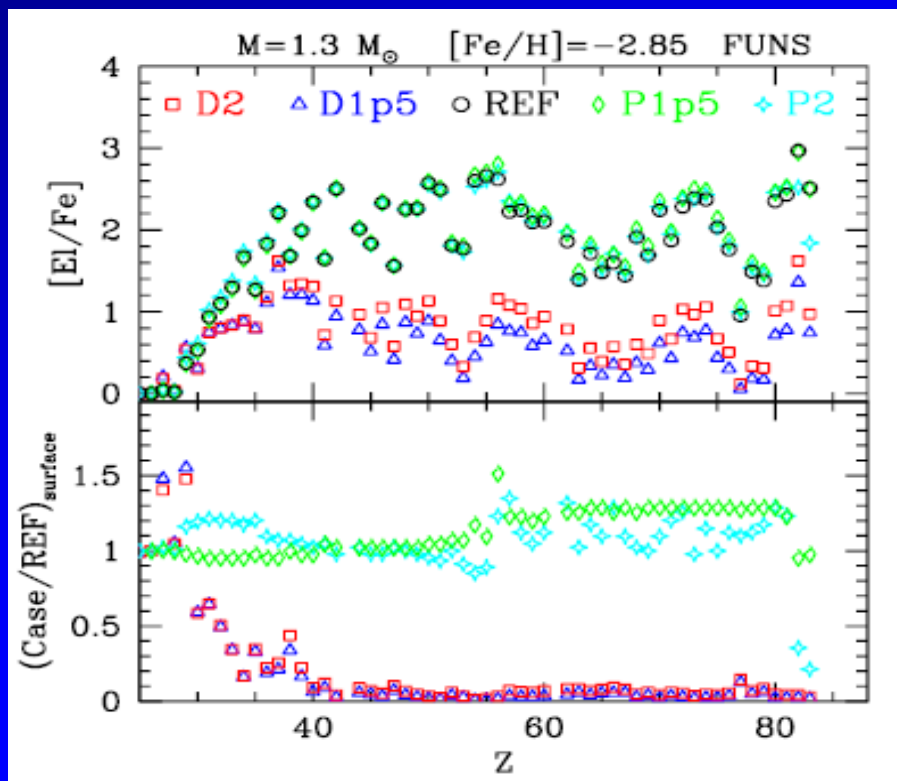
The importance of the $^{13}\text{C}(\alpha,n)^{16}\text{O}$ reaction (energetic point of view)



NECESSARY CONDITION:

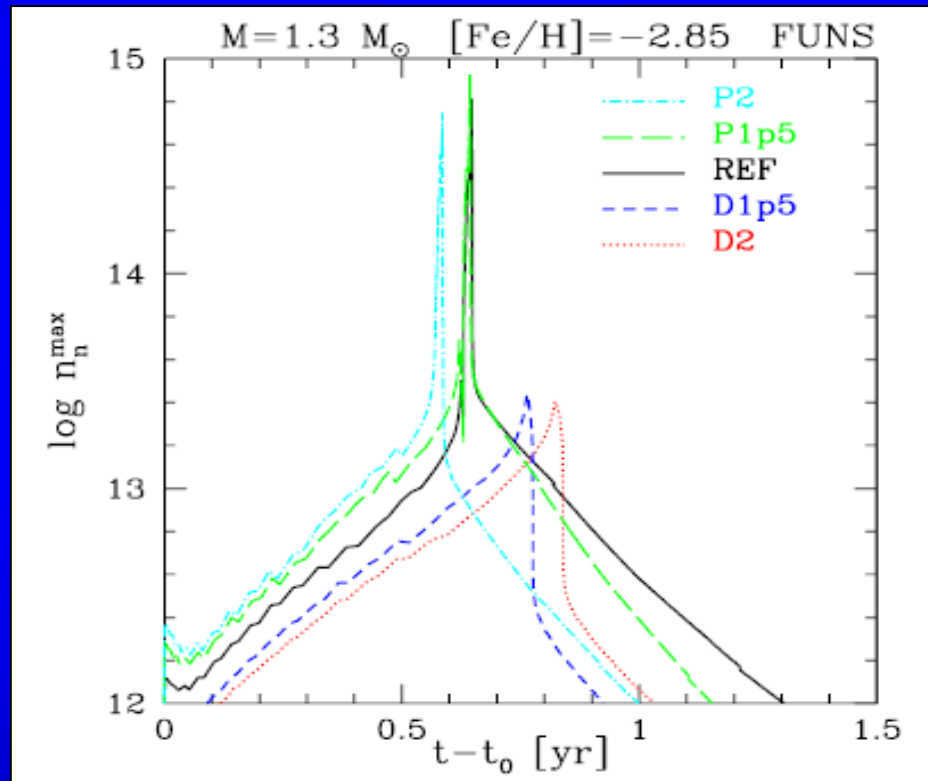
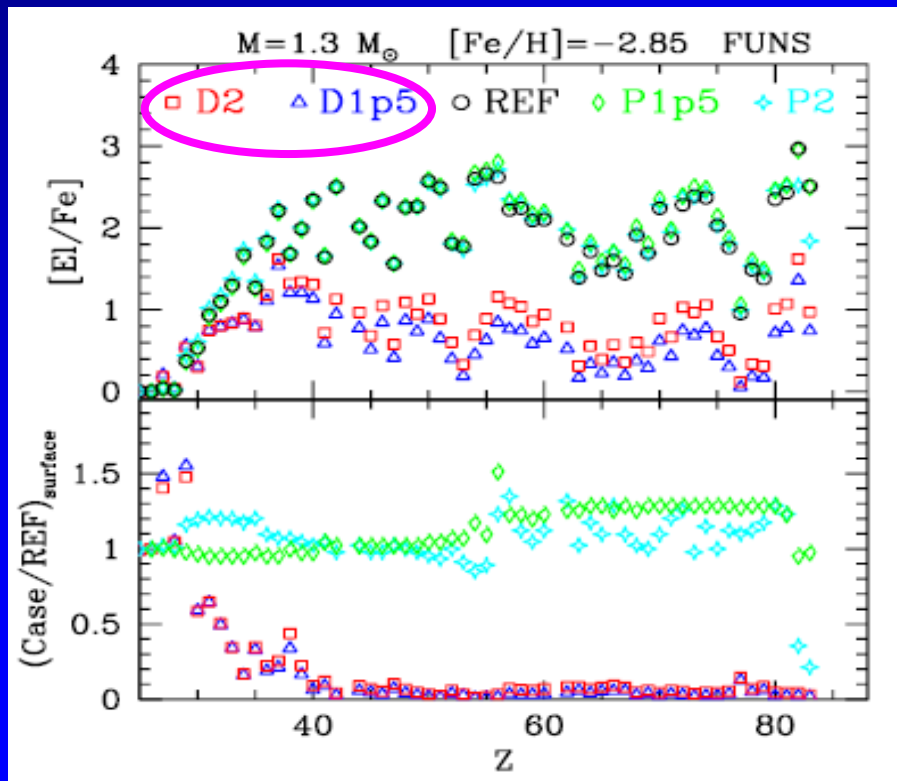
Full network coupled to physical evolution,
due to neutron captures energetics!!!

$M=1.3 M_{\text{SUN}}$ $[\text{Fe}/\text{H}]=-2.85$



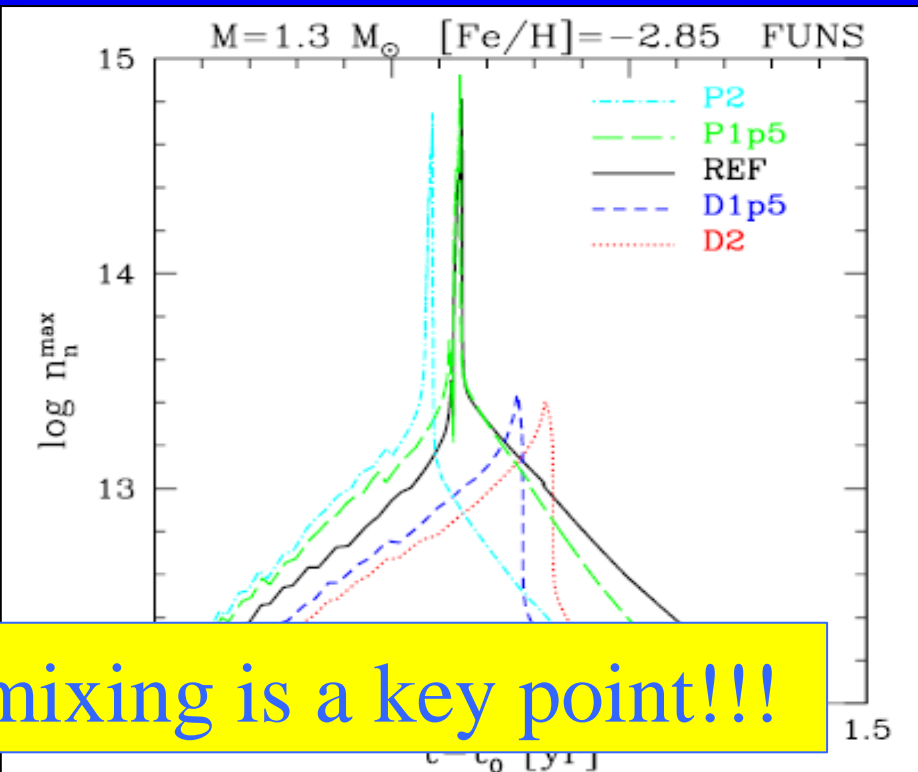
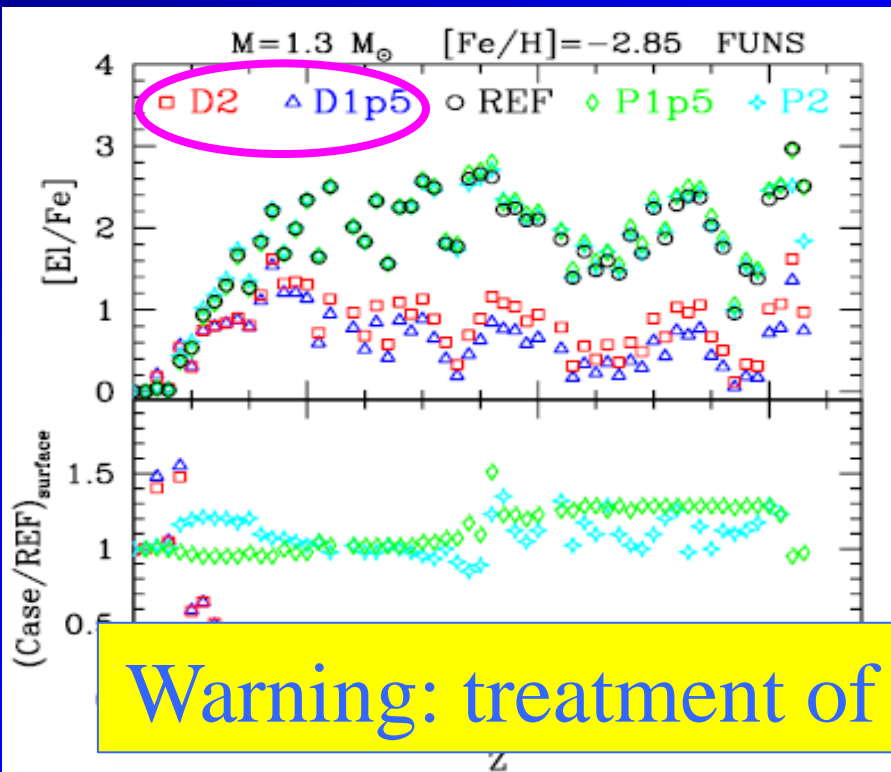
$M=1.3 M_{\text{SUN}}$ $[\text{Fe}/\text{H}]=-2.85$

Threshold effect: not enough ^{13}C mixed before shell splitting



$M=1.3 M_{\text{SUN}}$ $[\text{Fe}/\text{H}]=-2.85$

Threshold effect: not enough ^{13}C mixed before shell splitting



Warning: treatment of mixing is a key point!!!

CONCLUSIONS

- The $^{13}\text{C}(\alpha,n)^{16}\text{O}$ reaction rate is important in low mass AGBs ($M < 3.0 M_{\text{SUN}}$) at close-to-solar metallicities, because it determines how much ^{13}C burns in a convective environment;
- A variation of the $^{13}\text{C}(\alpha,n)^{16}\text{O}$ reaction rate does not change s-process abundances in more massive AGBs ($M > 3 M_{\text{SUN}}$), as well as in all masses at low metallicities;
- The $^{13}\text{C}(\alpha,n)^{16}\text{O}$ reaction could be important for low mass AGBs at very low metallicities, because in that case the convective ^{13}C burning (together with the subsequent neutron capture) affects the physical and chemical evolution of the model. However, the treatment of mixing is equally important.