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Binary compact objects as sources of gravitational waves

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OUTLINE:

1. The formation of compact objects from stellar evolution and supernova explosions

2. Isolated formation of binary compact objects with EXERCISES on binary compact objects

3. Dynamical formation of binary compact objects



What have astrophysicists learned from O1 + O2 detections?

- 1. Binary neutron star (BNS) mergers are associated with a gorgeous electromagnetic emission (Abbott+ 2017 on GW170817)
- 1. Binary black holes (BBHs) exist

(Tutukov & Yungelson 1973; Thorne 1987; Schutz 1989)

- **2. BBHs can merge in a Hubble time**
- 3. Massive black holes (BHs) exist i.e. stellar BHs with mass >20 M⊙ (Heger et al. 2003; MM et al. 2009, 2010; Belczynski+ 2010)

BHs in X-ray binaries < 20 M☉ (Ozel+ 2010) Most models of BH demography do not predict massive BH



LIGO-Virgo | Frank Elavsky | Northwestern



(BSE code, Hurley+ 2002)



From Spera, MM & Bressan 2015, MNRAS, 451, 4086

See also MM+ 2009, MNRAS, 395, L71; MM+ 2010, MNRAS, 408, 234; Belczynski+ 2010, ApJ, 714, 1217; Fryer+ 2012, ApJ, 749, 91; MM+ 2013, MNRAS, 429, 2298; Belczynski+ 2016, A&A, 594, 97; Spera & MM 2017, MNRAS, 470, 4739

Two critical ingredients:

1) PROGENITOR STAR EVOLUTION (STELLAR WINDS)

2) SUPERNOVA (SN) EXPLOSION





Winds ejected by Eta Carinae (HST, credits: NASA)

Chandra + HST + Spitzer Image of the SN remnant Cassiopeia A

1. The formation of compact objects: stellar winds

Massive stars (>30 Msun) might lose >50% mass by winds Stellar wind models underwent major upgrade in last ~10 yr (Vink+ 2001, 2005, 2011; see Vink+ 2016 for a short review)

Photons in atmosphere of a star couple with ions

 \rightarrow transfer linear momentum to the ions and unbind them

Coupling through resonant METAL LINES (especially Fe lines)

 \rightarrow MASS LOSS DEPENDS ON METALLICITY



How do we define metallicity in astrophysics? Metallicity in astrophysics is NOT same as chemistry

Metals in Astro: every element heavier than Helium

Measured with *Z* = FRACTION of elements heavier than He

X + Y + Z = 1.0

If M = total mass of system

 $X = m_p / M$ $Y = m_{He} / M$ $Z = \sum_i m_i / M$

Cosmological values: Sun values: *X* ~ 0.75, *Y* ~ 0.25, *Z* ~ 0 *X* ~ 0.73, *Y* ~ 0.25, *Z* ~ 0.02

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Metallicity dependence less important when STAR is CLOSE to electron-scattering EDDINGTON LIMIT (RADIATION PRESSURE dominates)



1. The formation of compact objects: stellar winds



Models from PARSEC stellar evolution code (Bressan+ 2012; Tang+ 2014; Chen, Bressan+ 2015)

Pre-supernova mass of a star is very important because affects the outcome of the SUPERNOVA



When Fe core forms in a massive (> 8 Msun) star

- 1) Fe-group atoms (Ni-62, Fe-58, Fe-56) have maximum binding energy: no more energy released by fusion
 → core starts collapsing because pressure drops
- 2) electron degeneracy pressure tries to stop collapse but if core mass > Chandrasekhar mass (~1.4 Msun) electron + proton capture removes electrons <u>→ electron pressure decreases</u>



- → COLLAPSE to NUCLEAR DENSITY, where <u>neutron degeneracy pressure</u> stops collapse
- → PROTO-NEUTRON STAR FORMS

Collapse of the core to nuclear density produces BOUNCE SHOCK

Fraction of binding energy of core (Eb,c ~10⁵³ erg) is converted into thermal energy (mostly of neutrinos)



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SHOCK MUST REVERSE COLLAPSE OF OUTER LAYERS

But density must be sufficiently high that neutrinos interact, otherwise neutrinos leak away without transferring energy → SHOCK MIGHT STALL

 \rightarrow SN FAILS

WHAT CAN REVIVE THE SHOCK?

STANDARD MODEL: CONVECTIVE ENGINE

Fryer 2014, http://pos.sissa.it/archive/conferences/237/004/FRAPWS2014_004.pdf



DEVELOPMENT OF CONVECTIVE BUBBLES

HELPS ENERGY FLUX TO REACH OUTER

ν

CONVECTIVE ENGINE:

LAYERS: SHOCK IS REVIVED

ν

ν

Collapsed core neutron pressure supported (proto NS)

Supernova shock stops anyway if BOUND MASS is too LARGE (Fryer 1999; Fryer & Kalogera 2001)

Back-of-the-envelope calculation to connect direct collapse and pre-supernova mass:

$$E_{\rm SN} = \frac{G M_{\rm env} \left(M_{\rm env} + M_{\rm core}\right)}{R_{\rm env}} \sim 1 \,\rm Msun}$$

Star cannot explode if
envelope binding energy
> SN energy
$$M_{\rm env} \sim 50 \,M_{\odot} \left(\frac{E_{\rm SN}}{10^{51} \rm erg}\right)^{1/2} \left(\frac{R_{\rm env}}{10 \,R_{\odot}}\right)^{1/2}$$

If M_{fin}>50 Msun this SN fails and star collapses to a BH

>

CRITERIA FOR COLLAPSE TO A REMNANT

depends on the "compactness" of the inner layers of the star

1. MASS OF CARBON-OXYGEN CORE If Mco > 8 – 12 Msun SN FAILS (Fryer+ 1999, 2012; Belczynski+ 2010)

2. COMPACTNESS

3. TWO-PARAMETER CRITERION

Core-collapse (CC) SN depends on the "compactness" of the inner layers

COMPACTNESS (= ratio between mass and radius) of a given portion of the stellar core at the onset of collapse

(O'Connor & Ott 2011)

$$\xi_M \equiv \frac{M/M_{\odot}}{R(M)/1000 \,\mathrm{km}}$$

 $M = 2.5 M_{\odot}$ is usually adopted

Star collapses if $\xi_{2.5} > 0.2$

(Ugliano+ 2012; Horiuchi+ 2012)

Figure from O'Connor & Ott 2011



Core-collapse (CC) SN depends on the "compactness" of the inner layers

Compactness correlates well with mass of CO core

 \rightarrow compactness > 0.2 corresponds to CO core > 8 M \odot



Core-collapse (CC) SN depends on the enclosed mass (M₄) and mass gradient (μ_4) at a dimensionless entropy per nucleon s = 4

$$M_4 = m(s=4)/M_{\odot}$$
 $\mu_4 = \left[\frac{dm/M_{\odot}}{dr/1000 \,\mathrm{km}}\right]_{s=4}$



Ertl et al. 2016

Core-collapse (CC) SN depends on the enclosed mass (M₄) and mass gradient (μ_4) at a dimensionless entropy per nucleon s = 4



ISLANDS OF DIRECT COLLAPSE AND SN EXPLOSION

Concluding remark: MANY MODELS of core-collapse SN EXPLOSION – REMNANT MASS CONNECTION BUT IF THE STAR IS VERY MASSIVE (>40 M☉) THEY GIVE SIMILAR RESULT

CC SN depends on the "fallback" of the outer layers of the star:

How much material falls back to the proto-NS after the SN

Barely constrained – depends on explosion energy, angular momentum, progenitor's mass/metallicity



Heger 2003

CC SN depends on the rapidity of the explosion:

(e.g. Fryer+ 2012; Fryer 2014)



From Fryer 2014, http://pos.sissa.it/archive/conferences/237/004/FRAPWS2014_004.pdf

PAIR-INSTABILITY SUPERNOVAE (PISNe)

If star is very massive, Helium core mass > 64 M_{\odot} \rightarrow central temperature > 7 x 10⁸ K \rightarrow efficient production of γ -ray radiation in core

 \rightarrow γ -ray photons scattering atomic nuclei produce electron-positron pairs (1 Mev)

The missing pressure of γ-ray photons produces dramatic collapse during O burning, without Fe core

→ high-Temperature collapse ignites all remaining species

 \rightarrow an explosion is induced that leaves NO remnant

Ober, El Eid & Fricke 1983; Bond, Arnett & Carr 1984; Heger et al. 2003; Woosley, Blinnikov & Heger 2007

positron

electron

nucleus

PULSATIONAL PAIR INSTABILITY (PPI)

If star is quite massive, 64 M₀> Helium core mass > 32 M₀ \rightarrow some production of γ -ray radiation in core

→ γ-ray photons scattering atomic nuclei produce electron-positron pairs (1 Mev)

The missing pressure of γ -ray photons produces contraction during O burning, without Fe core

- → enhancement of nuclear reaction restores pressure
- → star gains equilibrium after one or more oscillations
- \rightarrow oscillations enhance mass loss and final mass is lower

Barkat, Rakavy & Sack 1967; Woosley, Blinnikov & Heger 2007; Yoshida et al. 2016; Woosley 2017

positron

electron

nucleus

Very complicated. However, as rule of thumb (MM+ 2009, 2013):





Heger et al. (2003)



My cartoon from Heger et al. (2003)

What about intermediate metallicities between 0 and solar?

- more difficult because stellar winds are uncertain
- importance of final mass: pre-supernova mass of the star (when CO core built)



Spera, MM, Bressan 2015

Remnant mass follows same trend as final mass \rightarrow stellar winds are crucial



Spera, MM, Bressan 2015

Importance of supernova model for "LOW" STAR MASSES (<40 Mo)



Evolution of very massive stars still uncertain → stellar winds are Eddington-limited rather than metallicity dependent



Spera & MM 2017

Role of pulsational pair-instability and pair-instability supernovae (still missing in most models)



Spera & MM 2017

2. Binary compact objects

LIGO – Virgo observe compact object BINARIES

How do binary black holes or binary neutron stars or BH – NS form?

1) PRIMORDIAL BINARIES or ISOLATED BINARIES: two stars form from same cloud and evolve into two black holes (BHs) gravitationally bound





2) DYNAMICAL BINARIES: BBH forms and/or evolves by dynamical processes
PRIMORDIAL BINARIES or ISOLATED BINARIES:

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MOST MASSIVE STARS ARE IN BINARY SYSTEMS (e.g. Sana et al. 2012; Moe & Di Stefano 2017)

NOT SO EASY: Many evolutionary processes can affect the binary

SN kick wind mass transfer Roche lobe mass transfer common envelope

tidal evolution magnetic braking orbital evolution gravitational wave decay

Mass transfer

Two stars in a binary might exchange mass

- **1.** wind mass transfer
- 2. Roche lobe overflow
- 3. common envelope

If two stars exchange mass (and some mass is lost from the system), the final mass of the black holes will be completely different from two single stars



Credits: ESO/L. Calçada/M. Kornmesser/S.E. de Mink

Mass transfer



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Roche lobes

Equipotential surfaces in a binary system

For all *q* > 0 exists a 8-shaped critical potential surface

The connecting point is L1 (inner Lagrangian point)

The two lobes are called ROCHE LOBES

Approximation assuming 2 lobes are perfect circles (accurate to ~ 1 %)

$$\frac{r_1}{a} = \frac{0.49 \, q^{2/3}}{0.6 \, q^{2/3} + \ln\left(1 + q^{1/3}\right)}$$

where a = semi-major axis $q = m_1 / m_2$ Eggleton 1983

Roche lobes

If a star fills its Roche lobe (:= if star radius is equal or larger than Roche lobe),

matter flows without energy change into the other star → MASS TRANSFER

called ROCHE LOBE OVERFLOW



Star that fills the RL: DONOR Star that receives (accretes) mass: ACCRETOR



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Note: accreting material gains angular momentum from Coriolis force and can form an accretion DISK around the accretor

Common envelope in binaries:

If mass transfer becomes unstable (e.g. both stars fill Roche lobe), COMMON ENVELOPE (CE) phase = Two stars, one envelope



Two massive stars initially underfilling Roche lobe

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Mass transfer becomes unstable: CE phase



Drag by the envelope leads the two cores to spiral in

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Probably the least understood process in binary evolution Four STAGES (with different physics):

1. loss of COROTATION: instable mass transfer prevents the envelope to co-rotate with the core NOT YET MODELLED SELF-CONSISTENTLY (Ivanova et al. 2013)

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 fast SPIRAL IN: two cores spiral in – they lose kinetic energy by drag with the gas and heat the gaseous envelope – on dynamical time scale (~100d) – SIMULATED IN 3D (Ricker & Taam 2008, 2012; Passy et al. 2012; Ohlmann+ 2016)

2. Binary compact objects



From Ohlmann et al. 2016, ApJ, 816, L9

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- 4. MERGER of the cores or EJECTION of ENVELOPE

SEE IVANOVA ET AL. 2013, A&ARv, 21, 59 for a review

Common envelope in binaries:

Most used analytic formalism ($\alpha\lambda$, Webbink 1984) does not capture physics. In its version by Hurley+ (2002, MNRAS, 329, 897) the $\alpha\lambda$ formalism is:

1. initial binding energy of envelope (λ = free parameter, geometrical factor)

$$E_{\text{bind,i}} = -\frac{G}{\lambda} \left(\frac{M_1 M_{\text{env,1}}}{r_1} + \frac{M_2 M_{\text{env,2}}}{r_2} \right)$$

2. orbital energy of the cores

$$E_{\rm orb} = -\frac{1}{2} \, \frac{G \, M_{\rm c,1} \, M_{\rm c,2}}{a}$$

3. change of orbital energy needed to unbind the envelope:

$$E_{\text{bind,i}} = \Delta E_{\text{orb}} = \alpha \left(E_{\text{orb,f}} - E_{\text{orb,i}} \right)$$

 α is second free parameter (energy removal efficiency)

Common envelope in binaries:

4. if
$$a_{
m f} < (r_{
m c,1} + r_{
m c,2})$$

or $r_{\mathrm{c,i}} < r_{\mathrm{L,i}}$

i.e. any of the two cores fills Roche lobe before envelope ejection THEN the cores merge (Hurley+ 2002, MNRAS, 329, 897)

PROBLEM IS: HOW TO CONSTRAIN $\alpha \,$ and $\,\lambda \,$?

Observations of WD binaries, NS binaries, SNIa, now gravitational wave events,

Common envelope in binaries:



updated version of BSE (MM+ 2017, Giacobbo+ 2018)

Alternative to common envelope:

chemically homogeneous evolution

(Marchant+ 2016; Mandel & de Mink 2016; de Mink & Mandel 2016)

BASIC IDEA:

- if stars are chemically homogeneous, their radii are smaller
- \rightarrow close binaries avoid common envelope and premature merger

To be chemically homogeneous, stars need to ROTATE fast

OVERCONTACT BINARIES (Marchant+ 2016):

Metal-poor fast rotating stars may OVERFILL ROCHE LOBE WITHOUT ENTERING COMMON ENVELOPE



Why?

Star rotation induces chemical mixing

Chemical mixing prevents star radius from growing significantly (efficient only if star is metal poor)

Predictions of this model:

- * nearly equal-mass BH-BH
- * BH masses ~25 60, 130 230 Msun increasing with decreasing metallicity (no low-mass BHs!)
- * aligned spins unless SN reset them

Supernova kicks and BH binaries:

A massive-star binary can become a BH-BH binary only if it is not unbound by SN kicks

WHY KICKS?

- * asymmetry in mass ejection during core collapse
- * asymmetry in neutrino emission during core collapse
- compact object

* symmetric mass loss in a binary: breaks the binary only if pre-SN mass > companion mass (Blaauw mechanism, Blaauw 1961)

Supernova kicks and BH binaries:

SN kicks for NSs constrained from velocity of PULSARS



Supernova kicks and BH binaries:

Hobbs+ (2005): 3-D velocity distribution of pulsars obtained from the observed 2-D distributions of pulsars

 \rightarrow Maxwellian distribution with sigma ~ 265 km/s



Supernova kicks and BH binaries:

High (>100 km/s) velocity kicks for NSs (with caveats!)

WHAT ABOUT BHs?

No reliable methods to measure. Then people assume

1. conservation of linear momentum

$$v_{\rm kick, BH} = \frac{m_{\rm NS}}{m_{\rm BH}} v_{\rm kick, NS}$$

2. BHs formed without SN (failed or direct collapse) get NO KICK + kick modulated by FALLBACK

$$v_{\rm kick, BH} = (1 - f_{\rm fb}) v_{\rm kick, NS}$$

<u>Isolated binary evolution</u> <u>summary:</u>

- * possible Roche lobe
- * 1st BH formation
- * Common envelope BH – giant crucial to shrink the binary from >>100 Rsun to <100 Rsun
- * If binary survives common envelope, formation of second BH
- * BH BH merger

cartoon from MM2018



2. Binary compact objects: EXERCISES

Binary evolution studied via POPULATION SYNTHESIS CODES:

- * include models of stellar evolution in a simplified way
- * include prescriptions for supernova explosions
- * include treatment of binary evolution processes
- * based on a Monte Carlo approach (direct integration would be too expensive)

Examples of population-synthesis codes

BINARY_C (Izzard+ 2000)

BPASS (Eldridge+ 2017)

BSE/MOBSE (Hurley+ 2002; Mapelli+ 2017; Giacobbo+ 2018)

COMBINE (Kruckow+ 2018)

Seba (Portegies Zwart+ 2001; Mapelli+ 2013)

SEVN (Spera+ 2015, 2019; Spera & Mapelli 2017)

StarTrack (Belczynski+ 2007, 2010)

https://drive.google.com/open?id=1YIXgf_Q8qXO8hp5jEcNXGIn3_F5jdo1P

Enter in folder BBH

File bbh_binary.txt contains the data of a simulation of a massive binary

The simulation was done with MOBSE (Giacobbo et al. 2018, https://ui.adsabs.harvard.edu/abs/2018MNRAS.474.2959G/)

The first line explains the content. In particular

col.1: time [Myr] col.3: mass star1 [Msun] col.6: log10(R1) [Rsun]

col.9: mass star2 [Msun] col.12: log10(R2) [Rsun]

col.14: semi-major axis [Rsun] col.15: eccentricity

Plot the following quantities:

- **1.** Masses of the two stars as a function of time
- 2. Radii of the two stars and their Roche lobe radii as a function of time
- 3. Semi-major axis as a function of time

QUESTIONS YOU WANT TO TRY TO ANSWER WITH THESE PLOTS

- * When do the two stars become black holes?
- * Does the binary undergo Roche lobe overflow?
- * Is there a common-envelope episode?
- * Does the binary merges by gravitational-wave emission?

Note that in the directory BBH you find also the file bbh_binary_log.txt which contains a useful log file of the simulation



* When do the two stars become black holes?



* Does the binary undergo Roche lobe overflow?* Is there a common-envelope episode?



* Does the binary merges by gravitational-wave emission?



when orbital sep. increases RL expands

EXERCISE 2: properties of a compact binary sample

https://drive.google.com/open?id=1YIXgf_Q8qXO8hp5jEcNXGIn3_F5jdo1P

Enter in folder A5

The folder contains several files, which are catalogues of merging binary black holes, binary neutron stars and black hole – neutron star binaries

simulated with MOBSE (simulation named α 5 in Giacobbo & MM 2018, https://ui.adsabs.harvard.edu/abs/2018MNRAS.480.2011G/)

File names are of typedata_BHBs_*.txtfor binary black holesdata_BHNS_*.txtfor black hole – neutron star binariesdata_DNSs_*.txtfor binary neutron stars

The number before .txt indicates the METALLICITY of the PROGENITOR STAR Z = 0.02, 0.016, 0.012, 0.008, 0.006, 0.004, 0.002, 0.0016, 0.0012, 0.0008. 0.0004 and 0.0002

The first line of each file contains just two numbers: the total stellar mass which has been simulated and the number of merging compact binaries in the catalogue (ignore this line)

Each of the other lines is a single compact binary
https://drive.google.com/open?id=1YIXgf_Q8qXO8hp5jEcNXGIn3_F5jdo1P

Description of the columns:

col. 1: identifier of the binary

col.2 : ZAMS mass of stellar progenitor 1 [Msun]

col. 3: ZAMS mass of stellar progenitor 2 [Msun]

col. 4: mass of compact object 1 [Msun] NOTE that compact object 1 is not necessarily more massive than compact object 2, it is just the compact object that forms from stellar progenitor 1

col. 5: mass of compact object 2 [Msun]

col. 6: mass of compact object 1 + mass of compact object 2 [Msun]

col. 7: delay time [Gyr]

col. 8: semi-major axis [Rsun] when the binary becomes double degenerates

col. 9: eccentricity when the binary becomes double degenerates

https://drive.google.com/open?id=1YIXgf_Q8qXO8hp5jEcNXGIn3_F5jdo1P

1. Make a histogram plot of the delay times of files

data_BHBs_0.0002.txt data_BHNS_0.0002.txt data_DNSs_0.0002.txt

WHAT IS THE TYPICAL TREND OF DELAY TIMES? DO YOU SEE DIFFERENCES BETWEEN NEUTRON STARS AND BLACK HOLES?

2. Make a histogram plot of the masses of the black holes in binary black holes and black hole – neutron star binaries for several different metallicities (it is not needed that you plot all of them)

WHAT IS THE MAXIMUM BLACK HOLE MASS? HOW DOES IT DEPEND ON METALLICITY?



WHAT IS THE TYPICAL TREND OF DELAY TIMES? DO YOU SEE DIFFERENCES BETWEEN NEUTRON STARS AND BLACK HOLES?

Black holes in BBHs:



WHAT IS THE MAXIMUM BLACK HOLE MASS? HOW DOES IT DEPEND ON METALLICITY?

Black holes in BHNS systems:



WHAT IS THE MAXIMUM BLACK HOLE MASS? HOW DOES IT DEPEND ON METALLICITY?

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2) DYNAMICAL BINARIES: BBH forms and/or evolves by dynamical processes

3. The dynamics of black hole (BH) binaries:

DYNAMICS is IMPORTANT ONLY IF

n > 10³ stars pc⁻³

i.e. only in dense star clusters, where encounters are common

BUT massive stars (compact-object progenitors) form in star clusters

(Lada & Lada 2003; Weidner & Kroupa 2006; Weidner, Kroupa & Bonnell 2010; Gvaramadze et al. 2012; see Portegies Zwart+ 2010 for a review)



R136 in the LMC

There are many different flavours of star clusters



Globular clusters

- ✓ Formed mainly 12 Gyr ago
- ✓ Single-age stars
- ✓ Long lived
- ✓ Very massive $(10^{4-6} M_{\odot})$

There are many different flavours of star clusters



Nuclear star clusters

- ✓ At center of galaxies
- Prolonged star formation still ongoing (3 Myr – 12 Gyr ago)
- ✓ Long lived
- ✓ Very massive (>10⁶ M☉)
- Sometimes coexist with super-massive black hole (eg in the Milky Way)

There are many different flavours of star clusters



Open clusters

- ✓ Age from few Myr to several Gyr
- ✓ Single-age stars
- ✓ Not so long lived: when they die they release stellar content in the field
 → building blocks of field
- ✓ Lower mass $(10^{2-5} M\odot)$

There are many different flavours of star clusters



Young star clusters

- ✓ Young (<100 Myr)</p>
- Not so long lived: when they die they release stellar content in the field
 - \rightarrow building blocks of field
- ✓ Spread of masses (>10^{2 - 5} M☉)
- ✓ Are the NURSERY of massive stars

There are many different flavours of star clusters



Young star clusters

A large fraction of what we call "field binaries" might have formed in young star clusters

What processes happen in star clusters which cannot happen in the field?



3. The dynamics of stellar BH binaries: **3-body encounters**

Binaries have a energy reservoir (internal energy)

$$E_{int} = \frac{1}{2}\,\mu\,v^2 - \frac{G\,m_1\,m_2}{r}$$

where m_1 and m_2 are the mass of the primary and secondary member of the binary, μ is the reduced mass (:= $m_1 m_2/(m_1+m_2)$), r and v are the relative separation and velocity.

$$E_{int} = -\frac{G\,m_1\,m_2}{2\,a} = -E_b$$

THE ENERGY RESERVOIR of BINARIES can be EXCHANGED with stars during a 3-BODY INTERACTION, i.e. an interaction between a binary and a single star



3. The dynamics of stellar BH binaries: FLYBYs



In a flyby, the star acquires kinetic energy from the binary

- \rightarrow the binary shrinks
- → shorter coalescence time

3. The dynamics of stellar BH binaries: FLYBYs



Hills 1992, AJ, 103, 1955; Sigurdsson & Hernquist 1993, Nature, 364, 423; Portegies Zwart & McMillan 2000, ApJ, 528, L17; Aarseth 2012, MNRAS, 422, 841; Breen & Heggie 2013, MNRAS, 432, 2779; MM+ 2013, MNRAS, 429, 2298; Ziosi+ 2014, MNRAS, 441, 3703; Rodriguez+ 2015, PhRvL, 115, 1101; Rodriguez+ 2016, PhRvD, 93, 4029; MM 2016, MNRAS, 459, 3432; Banerjee 2017, MNRAS, 467, 524 and many others **3.** The dynamics of stellar BH binaries: EXCHANGEs



Exchanges bring BHs in binaries

BHs are FAVOURED BY EXCHANGES BECAUSE THEY ARE MASSIVE!

BH born from single star in the field never acquires a companion BH born from single star in a cluster likely acquires companion from dynamics

NEUTRON STARs (NSs) are lighter \rightarrow **Dynamics is less important for NSs**

3. The dynamics of stellar BH binaries: EXCHANGEs

Credits: Aaron Geller @Northwestern



ciera.northwestern.edu/Research/visualizations/videos/Binary+single.mp4 ciera.northwestern.edu/Research/visualizations/videos/Binary+singleex.mp4 ciera.northwestern.edu/Research/visualizations/videos/Triple+binary.mp4 **3.** The dynamics of stellar BH binaries: EXCHANGEs



>90% BH-BH binaries in young star clusters form by exchange (Ziosi, MM+ 2014, MNRAS, 441, 3703)

EXCHANGES FAVOUR THE FORMATION of BH-BH BINARIES WITH

- * THE MOST MASSIVE BHs
- * HIGH ECCENTRICITY
- * MISALIGNED BH SPINS

3. The dynamics of stellar BH binaries: MASSEs



Di Carlo et al. 2019, arXiv:1901.00863 see also Banerjee+ 2010; Ziosi+ 2014; MM 2016; Kimpson+ 2016; Banerjee 2017, 2018; Rastello+ 2018; Kumamoto+ 2018

3. The dynamics of stellar BH binaries: ECCENTRICITY



Ziosi, MM+ 2014, MNRAS, 441, 3703; Rodriguez+ 2015, Phys. Review Letter, 115, 1101; Hurley+ 2016, PASA, 33, 36; Askar+ 2017, MNRAS, 464, L36; Banerjee 2017, MNRAS, 467, 524 and many others





Spins of BBHs formed by exchange are ISOTROPICALLY distributed

Spins of BBHs formed from isolated binaries can be misaligned by SN kicks, but most remain aligned (especially massive binaries)

3. The dynamics of stellar BH binaries: intermediate-mass BHs (IMBHs)

Formalism by Miller & Hamilton (2002) to form IMBHs (100 – 10'000 Msun)

In a old cluster stellar BHs can grow in mass because of repeated mergers with the companion triggered by 3-body encounters



3. The dynamics of stellar BH binaries: wrap up

Dynamical binary evolution summary:

- * no need for Roche lobe or common envelope (but might happen)
- * exchanges build up more massive black hole binaries
- hardening by three-body encounters favours the binary shrinking
- * BH BH merger

cartoon from MM 2018, https://arxiv.org/abs/1809.09130



3. The dynamics of stellar BH binaries: wrap up



THANK YOU

