IFAE 2019

Incontri di Fisica delle Alte Energie

Napoli, 8 – 10 Aprile

Misure di luminosità ad ATLAS e CMS in run pp a LHC

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Perchè la luminosità

Per ogni processo fisico con sezione d'urto σ , il rate di eventi R è proporzionale alla luminosità istantanea \mathcal{L} :

- μ = numero di collisioni pp inelastiche per bunch crossing
- n_b = numero di coppie di bunch collidenti
- f_r = frequenza di rivoluzione di LHC (11245 Hz)
- $\sigma_{\rm inel}$ = sezione d'urto totale pp (~80 mb a 13 TeV)
- \mathcal{E} = accettanza e efficienza del rivelatore
- $\mu_{\rm vis}~=$ numero medio di collisioni per bunch crossing rivelato
- $\sigma_{\rm vis}~=$ sezione d'urto "visible" = costante di calibrazione

Misure di luminosità precise:

- *L* istantanea → monitor online LHC/trigger: performance e operatitivà (lumi levelling, monitoring del fascio...). Precisione richiesta: <3-5%
 </p>
- L integrata → analisi fisiche: misure di sezione d'urto, test SM, nuova fisica (vincolare le incertezze PDF che limitano le teorie). Precisione richiesta: <1-2%

Monitor di luminosità di ATLAS per il Run 2



Monitor di luminosità di CMS per il Run 2



Calibrazione assoluta: vdM scan

μ_{vis} in funzione delle separazione del fascio δ_{x(y)}
 Luminosità tramite parametri del fascio



Misure di corrente (da LHC)

Larghezza dei fasci (da fit sulla curva di scan)

Sezione d'urto visibile calibrata per ogni rivelatore/algoritmo

$$\sigma_{vis} = \frac{R}{L} = \frac{\mu f_r n_b}{L} = \frac{\mu m_{vis}^{peak}}{n_{p1} n_{p2}}$$

Assunzione: funzione densità di probabilità protoni fattorizzabile

$$\mathscr{L}\left(\delta_{x},\delta_{y}\right) = f_{x}\left(\delta_{x}\right)f_{y}\left(\delta_{y}\right)$$





Incertezze dominanti sulla calibrazione vdM

	ATLAS	CMS	ATLAS	CMS
	2017	2017	2018 (preliminary)	2018
Non fattorizzabilità	0.2 ± 0.2%	0.8 ± 0.8%	0.3 ± 0.5%	ongoing
Riproducibilità scan per scan	± 1.2 %	± 0.9 %	±0.6%	ongoing
Calibration transfer	±1.3%	±1.5%	±1.3%	ongoing

CMS-PAS-LUM-17-004



http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PLOTS/LUMI-2017-001/

Stabilità durante l'anno



- ATLAS: confronto fra luminometri attivi durante la presa dati, incluso Z-counting (non usato per valutare l'incertezza) e Track-counting 2017: ±1.3%; 2018: ±0.8% (preliminary)
 CAAS: DAAS del represente LIEET (DCC (00, 4% della basis esità nel 2017)
- CMS: RMS del rapporto HFET/PCC (99.4% della luminosità nel 2017) 2017: ±0.5%; 2018: ongoing

Incertezza sistematica totale

	ATLAS	CMS	ATLAS	CMS	ATLAS	CMS	ATLAS	CMS	ATLAS	CMS
Running	2012	2012	2015	2015	2016	2016	2017	2017	2018	2018
period	рр	рр	рр	рр	рр	рр	рр	рр	рр	рр
√s [TeV]	8	8	13	13	13	13	13	13	13	13
σ _L /L [%]	1.9	2.6	2.1	2.3	2.2	2.5	2.4	2.3	2.0 pret	ongoing

ATLAS: incertezza sistematica finale combinata per il Run 2:

- Incertezze non correlate: analisi vdM, stabilità a lungo termine....
- Incertezze correlate: calibration transfer...

Dataset	Int. luminosity (fb ⁻¹)	Rel. uncertainty
2015+2016	36.21 ± 0.77	2.1%
2017	44.31 ± 1.04	2.4%
2018	58.45 ± 1.16	2.0%
2015-2018	139.0 ± 2.4	1.7% preli

ATLAS Ref: https://twiki.cern.ch/twiki/bin/viewauth/Atlas/LuminosityForPhysics

CMS Ref: CMS-PAS-LUM-17-004/17-004/15-001/13-001

HL-LHC: ATLAS upgrade

- Misura direttamente proporzionale alla luminosità: LUCID
- PMT saturati ad alto μ sostituiti con fibre di quarzo, readout lontano da zona radiattiva
- Misura luminosità bunch-per-bunch: BCM'
- stazioni di diamanti con 3 sensori indipendenti per luminosità, fondo e sistema di aborto del fascio
- Timing: High Granularity Timing Detector
- separazione spaziale dei vertici ad alto μ tramite informazione temporale: σ_t = 30 ps.
- Low-Gain Avalanche Detector (LGAD)
- Linearità fra numero hit nel detector e numero di interazioni per bunch crossing







HL-LHC: CMS upgrade

Approccio multi detector

(difficoltà nel dividere effetti di non linearità con solo due rivelatori)

Inner Tracker endcap (TEPX):

indipendente dal readout di CMS

- Drift Tube (DT): bassa occupanza → misure bunch per bunch
- VeloPix: Monitor di luminosità dedicato indipendente altamente resistente a danni da radiazioni
- Timing Detector: separazione spaziale dei vertici ad alto μ tramite precisa informazione temporale: σ_t = 30-50 ps/traccia.





Conclusioni

- Luminosità parametro chiave per analisi di fisica ma misura complicata negli acceleratori adronici
- Ridondanza di luminometri cruciale per studio delle prestazioni e valutazione delle sistematiche: tipica incertezza sistematica totale ~2%!
- Misura della luminosità a LHC durante Run 1 e Run 2 ottenuta con precisione finale tale da permettere test di precisione sullo Standard Model e ricerche di nuova fisica.

Prospettive future (LHC Fase II): studi in corso su upgrade dei luminometri per nuove condizioni di LHC.

Back up

- Brevi scan vdM all'inizio e alla fine di ogni fill di LHC
 - Fasci separati in X e Y per 7/9 steps di 10s/punto
 - Livello di precisione inferiore rispetto vdM:
 - intervallo di scan limitato (non sensibile alle code)
 - possibli effetti di non fattorizzabilità (modo di produzione del fascio differente)
 - effetti di dinamica del fascio
 - Utili per misure di luminosità relativa
- Controllo online effetti di linearità e stabilità:
 - LHC: diagnostica e controllo delle prestazioni;
 - CMS (2017) e ATLAS (2018): controllo online dei luminometri.



 $\Sigma_{r}^{2} = 2$



At zero crossing angle

Luminosity algorithms

- Event- (or zero-) counting algorithms:
 - Based on Poisson statistics: count of events with at least one hit

$$P_{OR} = \frac{N_{OR}}{N_{orbits}} = 1 - e^{-\mu\varepsilon_{OR}} \Longrightarrow \mu = -\ln(1 - \frac{N_{OR}}{N_{orbits}})$$

- If μ too large \rightarrow "zero starvation" or "saturation"
- Hit-counting algorithms: Now: ATLAS: # LUCID hits. CMS: # pixel clusters.
 - Count of total hits in a given BX
 - $-\,$ based on Poisson statistics but saturation at higher μ
- Track- (& vertex-) counting algorithms:
 - conceptually similar to hit-counting. Examples: ATLAS.
- Particle-counting algorithms (summed over all bunches)
 - Examples in ATLAS: current in hadronic-calorimeter photomultipliers or charge measurements (LUCID).

Lumi from	Detector type	Data flow	Name	Lumi Algo
ATLAS	P-CVD diamond pads	Bunch-by-bunch (bbb)	ВСМ	Event counting
	Quartz Cherenkov tubes	bbb	LUCID	Event counting Hit counting
	Si strip + pixel tracker: #vertices	bbb	"Vtx"	Vtx counting
	Si strip + pixel tracker: #tracks	bbb	"Trks"	Trks counting
	Fwd LAr / E.M. EndCap calo: gap currents	Bunch-averaged (ba)	FCal	Particle flux
	TILE calorimeter	ba	TILE	Particle flux
	Pixelated radiation monitor	ba	ТРХ	Hit counting
CMS	Pixel trk: #clusters	bbb	PCC	Hit counting
	Fwd Fe/quartz calo	bbb	HFET	E_T flow (analog)
	Fwd Fe/quartz calo.	bbb	HFOC	Hit counting
	Pixel telescope	bbb	PLT	Hit counting (3-fold coinc)
	Fast Beam Conditions Monitor	bbb	BCM1f	Trk segment counting
	Muon drift tube	ba	DT	Hit counting

ATLAS/CMS luminosity ratio





- Largest contribution: emittance_x > emittance_y, coupled with horizontal (x) crossing in CMS vs. vertical (y) crossing in ATLAS
- Analysis complicated by residual µ- or time-dependence of reported L, that could be different in the two experiments

>most trusted offline algorithms: track-cntg (ATLAS), pixel-cluster cntg (CMS)

ightarrow dedicated experiment: crossing-angle scan

Vdm Scan calibration: difficulties

Central role of beam dynamics -> two beam-beam effects:

- beam-beam deflection: if bunches not exactly centred → angular kick due to e.m repulsion;
- 2) dynamic β : mutual (de)focusing of the two colliding bunches;

Effect: < 0.5% PbPb, 1 - 2% for 7/8/13 TeV pp and around 4% for 5 TeV pp.

Scan curve distorted by interactions of the two beams during a scan.



beam separation larger than nominal separation



Beams focus/defocus each other by an amount that is a function of separation

Dynamic-β

Non-factorisation correction procedure

$$\mathscr{L}(\delta_x, \delta_y) = f_x(\delta_x) f_y(\delta_y) ?$$

 Single beam profiles are parameterised by fitting the beam-separation dependence of the luminosity & of the beamspot displacement and width during a vdM scan.

This allows to:

- → estimate the true luminosity (i.e. unbiased by non-factorisation effects)
- → estimate correction for non-factorisation, *R*, with an associated uncertainty





 $R = rac{\mathscr{L} ext{ not assuming factorisation}}{\mathscr{L} ext{ assuming factorisation}}$

 The [ATLAS/ALICE] procedure above is closely related to the "beam-beam imaging" scans [pioneered by LHCb & now established method in CMS] in which one beam is scanned transversely as a probe across the other.



Beam separation (x-scan)

Non-factorization correction: beam-beam imaging

- Principle: use one beam (\sim wire) to probe the other
 - keep witness beam (B1) stationary; scan probe beam (B2) across it in x, then in y; repeat with B1 $\leftarrow \rightarrow$ B2
 - measure 2-d distribution of reco'd evt vertices at each step: N_{vtx}(x, y) ={r_{witness}(x,y) x r_{probe}(x,y)} (X) R_{vtx position}(x,y) (see ArXiv_1603.0356 [hep-ex])
 - extract single-beam parameters of B1 & B2 from fit to 2-d vertex distributions in the 4 scans (B1/B2, x/y)
 - closely related to the ATLAS & ALICE luminous-region evolution method (but uses only transverse info, not L/z)
 - common key issue: vertex-position resolution $R_{vtx position}$
 - pros & cons of the 2 approaches to be clarified

Non-factorization correction: beam-beam imaging





Example of pull distributions of the fitted single-beam model of the single-gaussian (factorizable, left) and double-gaussian (non-factorizable, right) type to the vertex distribution accumulated during scan Y3 of bunch pair1631. (Caption adapted from Fig. 11 of CMS-PAS-LUM-2015-001)

Incertezze dominanti sulla calibrazione vdM



Correzioni per effetti di non-fattorizzabilità

- dipendenza della luminosità verticale dalla separazione orizzontale e viceversa
- Correzione analisi vdM scan sotto assunzione di fattorizzabilità
 - $2017: 0.8 \pm 0.8\%$ 2018: ongoing
 - ATLAS 2017 0.2 \pm 0.2% $2018\ 0.3\pm 0.5\%$

1ttps://twiki.cern.ch/twiki/bin/viewauth/Atlas/LuminosityForPhysics CMS-PAS-LUM-17-004

Riproducibilità calibrazione nei diversi scans:

2017: ± 0.9 % 2018: ongoing ATLAS 2017 ±1.2%,

NB: differenti strategie di valutazione delle sistematiche tra ATLAS e CMS \rightarrow difficile confronto diretto

Calibration transfer: dal vdM alla fisica

Non linearità nella risposta da vdM scan (bassa \mathcal{L} , basso μ , pochi bunch distanziati) a fisica (alta \mathcal{L} , alto μ , più di 2000 bunches ogni 25 ns)

ATLAS: $\mu_{\text{Tracking}}/\mu_{\text{Algo}}$ 1.05 LUCID HitOR Correzione non linearità tramite TILE tracciatore LUCID @ $\mu = 50$ 0.95 2017: - 9.1%; 2018: - 8.9% Incertezza sistematica dal confronto con 0.9⊢ ATLAS Preliminary calorimetri √s=13 TeV 2017: ±1.3%; 2018: ±1.3% LHC Fill 6259, Sep. 30, 2017 0.85 50 Interactions per Bunch Crossing ($\mu_{_{Alac}}$

> CMS:

- Correzione non linearità tramite analisi degli scan in emittanza
 - HFET $@ \mu = 50$ nel 2017: 1.5 %
- Incertezza sistematica dal metodo dei residui sul confronto fra luminometri 2017: ±1.5%; 2018: ongoing

ATLAS Ref.: https://twiki.cern.ch/twiki/bin/viewa uth/Atlas/LuminosityForPhysics CMS Ref.: CMS-PAS-LUM-17-004

GROUPS,

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OTS/

ATLAS: Non factorization evidence



Figure 3a:

Non-factorization correction factor R ($\sigma_{vis}^{corr}=\sigma_{vis}/R$) for several colliding-bunch pairs and scan sets (I-V), extracted from fits to the beam-separation dependence, during van der Meer (vdM) scans, of only the luminosity L. The beam-separation dependence of the luminosity is modeled by a two-dimensional (2-D) Gaussian function multiplied by a sixth-order polynomial (g.p6). The error bars are statistical only. The horizontal red lines represent the weighted average over all colliding-bunch pairs and scan sets, with the shaded bands indicating the RMS spread of the individual R values associated with each colliding-bunch pair.

ATLAS: Calibration transfer



http://atlas.web.cern.ch/Atlas/GRO UPS/PHYSICS/PLOTS/LUMI-2017-001/

Figure 8:

Ratio of the luminosity measured by the E3 and E4 Tile scintillators (averaged over the A and C sides of ATLAS) to that from track counting, in the 2017 vdM fill and a closely following high-luminosity physics fill. The ratios are normalised to unity in the vdM fill. Each point corresponds to the average over 30 luminosity blocks (approximately 30 minutes). The luminosity block numbers in the two runs have been offset so the physics fill begins at 1000.

CMS: Non factorization evidence



Figure 6: Pull distributions using the Super Double Gaussian fit model. The pull is defined as the difference between the number of measured vertices and the number of vertices predicted by the fit, divided by the statistical uncertainty of the measurement. These plots show the results from the scan constraining beam 1 in *x*. Left: 2-D pull distribution as a function of *x* and *y* position. Right: 1-D projections of the 2-D pull distribution, in slices of constant radius (top) and constant azimuthal angle (bottom).

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CMS: summary of 2017 uncertainties

Table 4: Summary of the systematic uncertainties entering the CMS luminosity measurement for $\sqrt{s} = 13$ TeV pp collisions. When applicable, the percentage correction is shown.

	Systematic	Correction (%)	Uncertainty (%)
	Length scale	-0.9	0.3
	Orbit drift	_	0.2
	<i>x-y</i> correlations	+0.8	0.8
	Beam-beam deflection	+1.6	0.4
Normalization	Dynamic-β*	/_/	0.5
	Beam current calibration	$\langle -$	0.3
	Ghosts and satellites		0.1
	Scan to scan variation	\rightarrow	0.9
	Bunch to bunch variation	\setminus	0.1
	Cross-detector consistency	0.4-0.6	0.6
	Afterglow (HF)	\rightarrow $/ / /$	0.2⊕0.3
Integration	Cross-detector stability	\checkmark \checkmark \checkmark	0.5
integration	Linearity	\searrow	1.5
	CMS deadtime		0.5
	Total		2.3

ATLAS reference luminometer: LUCID-

2



ATLAS Luminosity performance summary

Y	ear	c.m. energy (TeV)	Mu max	L_max (10 ³³ cm ⁻² s ⁻¹)	L_int (fb ⁻¹)	NBCID	Dt (ns)	Tot Sys Unc. (%)	vdM Sys. Unc. (%)	Reference Detector (online & offline)
2	010	7	5	0.2	0.047	348	150	3.5	3.4	LUCID-I
2	011	7	20	3.6	5.5	1331	50	1.8	1.5	BCM
2	012	8	40	7.7	22.7	1368	50	1.9	1.2	BCM
2	015	13	28	5	4.2	2232	25	2.1	1.7	LUCID-2
2	016	13	45	14	38.5	2208	25	2.2	1.2	LUCID-2
2	017	13	80	20	40	2544	25	2.4	1.5	LUCID-2
	Year	c.m. energy (TeV)		β* (m)	L_Ins (10 ³⁰ c ⁻²	Η L _ s ⁻¹)	_Int (μb⁻¹)	Tot Sys Unc. (%)	vdM Sys. Unc.(%)	Reference detector
	2011	7		90	5*10 ⁻	3	80	2.3	1.5	BCM
	2012	8		90	5*10 ⁻	2	500	1.5	1.2	BCM
	2012	8		1000	0.8*10	-3	22	1.4	1.2	LUCID

CMS: PLT

- Uses same pixel sensors and readout chips as phase-0 pixel detector
- 48 silicon sensor planes arranged in 16 "telescopes" (8 on either side of CMS) outside the pixel endcap (|η|~ 4.2)
- Use special "fast-or" readout mode of chip to look for events where all three planes in a telescope register a hit ("threefold coincidence") to measure luminosity
- Provide online bunch-by-bunch measurements to LHC and CMS with a statistical precision of 1% every 1.5s to allow for fast feedback (e.g., for beam optimizations)

CMS: BCM1F

- 24 sensors located on face of PLT/BCM1F carriage with fast readout (6.25 ns) to distinguish luminosity from machine background
- 2015-2016 all sensors were diamond, but severe problems with efficiency loss in 2016

EYETS 16/17 In EYETS sensors were replaced and upgraded to a mix of polycrystalline diamond, single crystal diamond, and silicon



CMS: PCC

- Pixel cluster counting uses the raw rates of pixel clusters in the main CMS pixel detector
 - Primary offline measurement in 2015 and 2016 (Phase 0 pixel detector)
 - Limited by CMS DAQ and trigger for online practicality
- Two major corrections necessary:
 - "Type 1" affect the next BX after a colliding BX for signal spillover
 - "Type 2" affect several BXes after for material activation



CMS: HF

36 PMT boxes each end

 Uses existing HF calorimeter with dedicated readout for luminosity information



- Two algorithms used:
 - HFOC: uses raw occupancy rate in HF (fraction of towers with hit energy above noise threshold).
 Standard in 2015-16 but some nonlinearities at higher pileup.
 - HFET: uses sum of E_T deposited in all HF towers. Commissioned during 2016 and is now the primary algorithm for 2017-2018 running.

Upgrade per HL-LHC

Prestazioni attese HL-LHC:

- Energia nel centro di massa: $\sqrt{s} = 14$ TeV
- L istantanea = $5.0 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
- L integrata = 3000 fb⁻¹
- Numero medio interazioni per bunch crossing: $<\mu>=200$



Difficoltà tecniche:

- Alto flusso di particelle, alto trigger rate
 - nuovo sistema di trigger
- Alta occupanza rivelatori
 - Limitazioni readout
 - Complessità di ricostruzione evento
- Aumento della fluenza vicino alla beam pipe fino a $10^{16} n_{eq} / \text{cm}^{-2}$
 - danni da radiazione
 - attivazione dei materiali



https://twiki.cern.ch/twiki/bin/view/AtlasPublic/ UpgradeEventDisplays

HL-LHC: ATLAS



HL-LHC: ATLAS

ATLAS *L*-upgrade: (one of the applications of) the HGTD [High-Granularity Timing Detector]

- O Original (& primary) motivation
 - Build Not the set of th
 - Spatially overlapping vertices can be resolved in the time dimension using accurate vertex timing measurements
- O HGTD in a nushell
 - ◎ two endcap disks at z = ± 3.5 m
 - O Active area: 120 mm < R < 640 mm
 ⇒ 2.4 < |η| < 4.0
 - Si-based Low-Gain Avalanche Detector (LGAD) technology
 ⇒ σ_t = 30 ps/track over the lifetime of HL-LHC
 - 2 Si layers per disk
 - - ⇒ 1.3 mm x 1.3 mm pixels

W. Kozanecki



FCAL Collaboration meeting, Krakow, 10 May 2018

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HL-LHC: ATLAS from CMS experience



- O Use pads 2.8 <|η| < 3.1</p>
 - nigh granularity → no saturation
 - 10% occupancy easily handled by Poisson formalism
 - good statistical power over full



Deadtime-less, bbb readout

- Hit count per ASIC (2 cm x 2 cm area) at 40 MHz (every BX on every turn) for
 - · central time window, and, separately, for
 - sideband(s) for afterglow subtraction
 - take advantage of good time resolution

O Excellent linearity...



... in simulation!

This effort would greatly benefit from acquiring real-life experience with PCC-based *L* determination using the forward-pixel disks in the present ATLAS detector CMS: special read out of Inner Tracker

W. Kozanecki

Sara Valentinetti

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Luminosity measurement with HGTD

- Exploit correlation between the number of hits in the detector and the number of interactions per bunch crossing
- □ The number of multi-hits in a single pixel is expected to be negligible (especially in the outer part of the detector $2.4 < |\eta| < 3.1$)
- The after-glow can be suppressed comparing the number of hits during a bunch crossing with the number of hits after a bunch crossing





- HGTD will allow unbiased high-statistics per-BCID luminosity measurements integrating all the hits in a period of 1s.
- □ A statistical precision of 0.14% for a run with an average number of simultaneous p-p interactions <µ>=1 and down to 0.01% for a run with<µ>=200 is possible.
- Better precision could also be achieved through a longer averaging time.

HL-LHC: ATLAS BCM'

Reasons:

- handle on the background activity in the tracker, with the aim of aborting LHC beams when danger levels approached.
- online bunch-by-bunch luminosity measurement at 1% precision.

Run 2 problematics:

- Current BCM provided by two systems:
 - fast BCM: both as a bunch-by-bunch abort and luminosity device.
 - Beam Loss Monitoring (BLM): diamond sensor integrating the charge over a minimum of 40μs.
- two functionalities provided by a resistive splitting of the very fast signal -> non-optimal performance on both tasks:
 - For luminosity: the speed results in degraded S/N, so signal amplitude variations induce efficiency drift.
 - For abort beam: luminosity requires single MIP sensitivity but this limits the level of the abort threshold settings, making them hard to justify, especially with only a few circulating bunches in the machine ,when the slow BLM's fail to show any activity

HL-LHC: ATLAS BCM'

 stations with 3 fully separate devices (pCVD diamond pad sensors): luminosity, fast and slow abort. The 1 cm2 diamond sensors of luminosity and abort devices are divided into7 pads of decreasing sizes from 32 to 1 mm2to effectively cover the foreseen dynamic range. The time-of-flight flagging of beam induced background used in the current BCM is kept. Background-induced showers originating upstream provoke signals early by2z/c≈12.5 ns (for the 1.9 m "golden" z location) in 4 stations on the upstream side with respect to the signal



Inner Tracker Endcap

- The larger tracker for HL-LHC will take up the space currently occupied by PLT & BCM1F
- However the endcap disks (TEPX) are in a perfect position to do a similar luminosity measurement



https://indico.cern.ch/event/697164/contributions/2987411/a **CMSLuminosity** 30/ 26337 ttachments/1647549 FCALWorkshop.pdf

Standalone Lumi Detector

- It is also desirable to have a completely independent lumi system
- Promising approach using VeloPix: very radiation-hard pixel chip being developed for LHCb VELO upgrade
- Could squeeze a single layer in the space outside the tracker



DT Lumi

- As currently, it looks like the most advantageous approach to using the muon chambers is to take advantage of tracks reconstructed at trigger level
- Should be possible in principle to make bunch-by-bunch measurements available (may even be possible, at least as a prototype, in Run3 after LS2)
- Overall occupancy should remain quite low even at HL-LHC levels

HLT TRACK RECONSTRUCTION



2. Level-2 (standalone) Reconstruction Reconstruction of the track inside the muon system

 Level-3 (global) Reconstruction Reconstruction of the track combining the information from tracker and muon system

Lumi with Timing Detectors

- Timing layer outside tracker to provide highprecision time measurements to improve PF reconstruction
- High time resolution (30-50ps) makes this promising as a luminosity measurement also
- Still just an idea at this point...will need much more development



ATLAS Z-counting



The invariant mass distribution of the muon pairs of the 240,000 Z ->mumu boson events selecting two muons with pT>27 GeV, pseudorapidity < 2.4 and 66 < m(mumu) < 116 GeV. The statistical errors are smaller than the symbol size.

ATLAS Z-counting



Figure 2: *Top:* The instantaneous luminosity determined from the $Z \rightarrow \mu\mu$ counting rate, $L_{ZCounting}$ (full circles), selecting two muons with $p_T^{\mu} > 27$ GeV, $|\eta^{\mu}| < 2.4$ and $66 < m_{\mu\mu} < 116$ GeV, and the ATLAS-preferred luminosity measurement L_{ATLAS} (red line) based on the LUCID online luminometer, both averaged over 20 Luminosity Blocks (LB). The LHC fill 6283 was taken at $\sqrt{s} = 13$ TeV on October 8, 2017. The Z counting rate is corrected for in situ data-driven trigger and reconstruction efficiencies including the residual Monte Carlo correction, and is normalised to the integrated ATLAS luminosity for this fill. The x-axis represents the elapsed time in units of Luminosity Blocks with a typical length of one minute per LB. Error bars are the statistical uncertainties of the $L_{ZCounting}$ determination.

ATLAS Z-counting



Figure 3: Fractional difference between the run-integrated luminosity determined from the $Z \rightarrow \mu\mu$ counting rate, L_{Z Counting}, selecting two muons with $p_T^{\mu} > 27$ GeV, $|\eta^{\mu}| < 2.4$ and $66 < m_{\mu\mu} < 116$ GeV, and the run-integrated, ATLAS-preferred luminosity measurement L_{ATLAS} (based on the LUCID online luminometer) per LHC fill taken at $\sqrt{s} = 13$ TeV in 2017. The Z counting rate is corrected for in situ data-driven trigger and reconstruction efficiencies including the residual Monte Carlo correction, and is normalised to the integrated ATLAS luminosity of the whole 2017 data taking period. The x-axis represents the date when the fill was recorded; only runs with at least 10,000 Z counts and a minimum length of about 40 min are included. Error bars reflect the statistical uncertainties of the L_{Z Counting} measurements only. The dashed line indicates zero.

ATLAS/CMS Z-counting ratio



ATLAS: sistematica finale Run 2

- Integrated luminosity in each year L_i and its uncertainty σ_i
 - Covariance matrix of the absolute integrated luminosities in each year V_L
 - This encodes the correlations between the year-by-year luminosities
- Then $L_{
 m tot}=\Sigma_i L_i$ and variance $\,\,\sigma_{L_{
 m tot}}^2={f GV_L} ilde{f G}$
 - Derivatives vector: $\mathbf{G} = \left(\frac{\mathrm{d}L_{\mathrm{tot}}}{\mathrm{d}L_1}, \frac{\mathrm{d}L_{\mathrm{tot}}}{\mathrm{d}L_2}, \frac{\mathrm{d}L_{\mathrm{tot}}}{\mathrm{d}L_3}, \ldots\right) = (1, 1, 1, \ldots)$
 - Covariance matrix V_L constructed by summing individual sources with uncertainties σ_i in each year (many separate correlated/uncorrelated sources):

$$V_L = \begin{pmatrix} \sigma_1^2 & 0 & 0 \\ 0 & \sigma_2^2 & 0 \\ 0 & 0 & \sigma_3^2 \end{pmatrix} + \begin{pmatrix} \sigma_1^2 & \sigma_1\sigma_2 & \sigma_1\sigma_3 \\ \sigma_1\sigma_2 & \sigma_2^2 & \sigma_2\sigma_3 \\ \sigma_1\sigma_3 & \sigma_2\sigma_3 & \sigma_3^2 \end{pmatrix} + \dots$$

An uncorrelated source A correlated source

- Some sources may not contribute to all years, so some σ_i=0
 - Sources with both correlated and uncorrelated parts are handled by being broken into two separate contributions to V₁

https://twiki.cern.ch/twiki/bin/view/Atlas/LuminosityForPhysics