

Josephson radiation sensors via temperature-to-phase conversion

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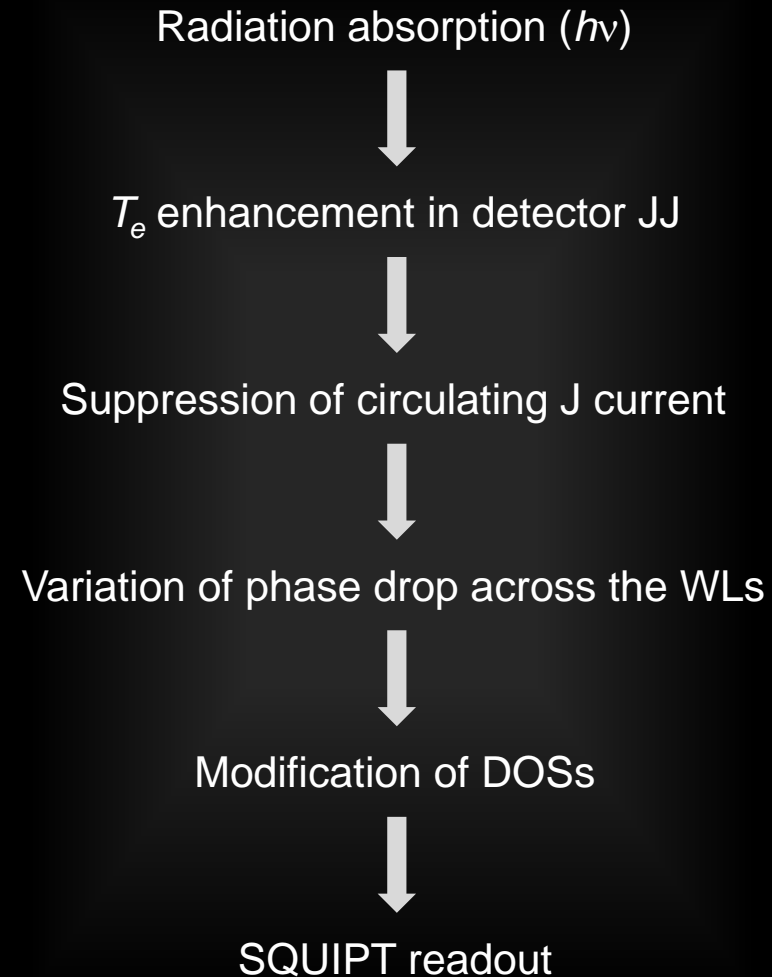
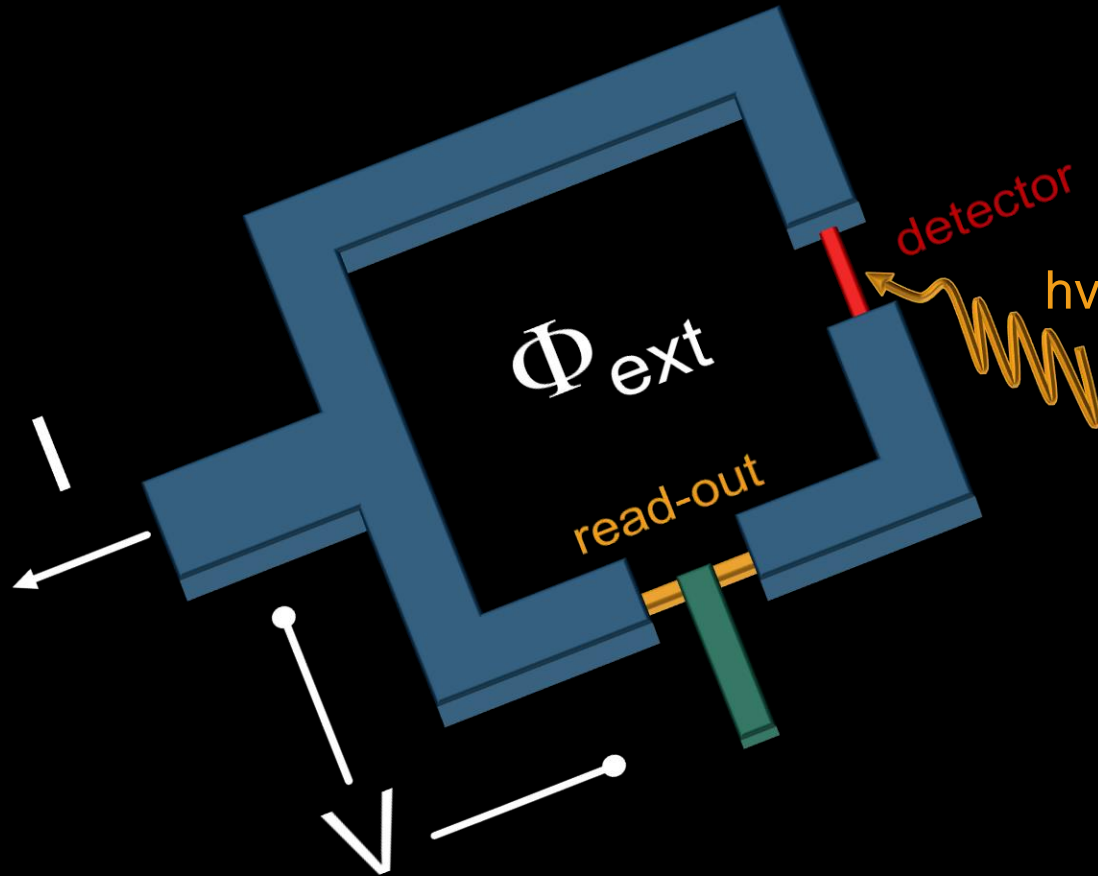
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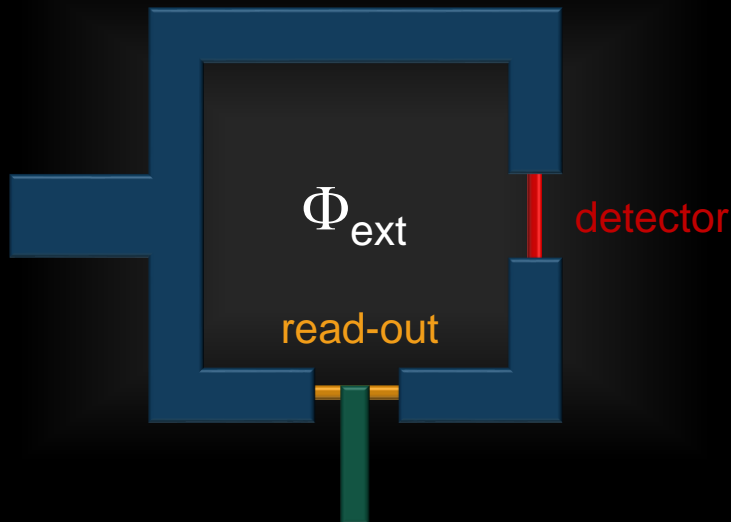
Outline

- Principle of photonic sensing based on T-to- ϕ transduction
- Theoretical model
- Superconducting Quantum Interference Proximity Transistor (SQIPT)
- Thermometer
- Nanocalorimeter operation
 - Thermal model, S/N ration, resolving power
- Bolometer operation
 - Thermal model, NEP
- Performances

Temperature-to-phase conversion: principle



Temperature-to-phase conversion: principle

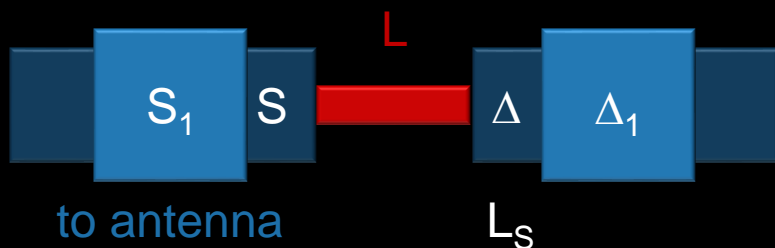


Assumptions:

- **Short** ($\Delta \leq \hbar D / L^2$) SNS JJ junctions:
 - Drastic reduction of detector volume V
 - Known analytical solutions for I and DOS
 - Enhanced response in the readout JJ
- Fluxoid quantization

$$\varphi_d + \varphi_r = 2\pi \frac{\Phi_{ext}}{\Phi_0} + 2k\pi$$

Andreev Mirror: $\Delta_1 \gg \Delta$
 Heat is confined in the detector region
 Appl. Phys. Lett. **63**, 3075 (1993)



$$L_S \gg \xi_0$$

Proximity only due to S

Good coupling with radiation when:

$$v > \frac{2\varepsilon_{gap}}{\hbar}$$

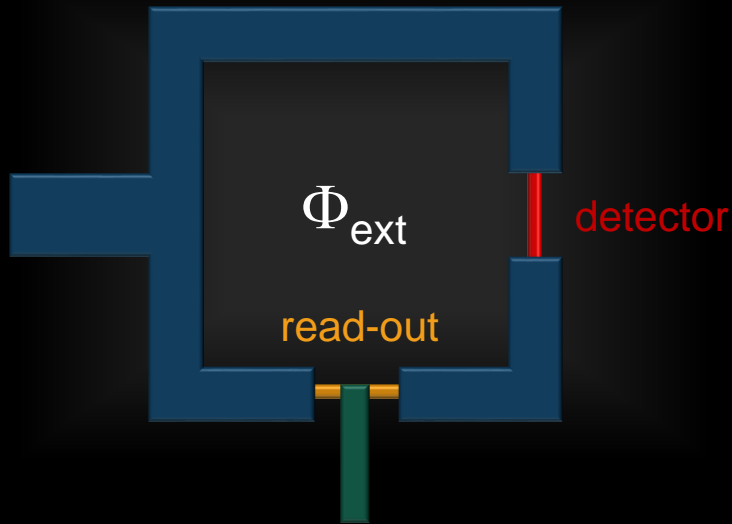
where the wire energy gap is

$$\varepsilon_{gap}^{d,r} = \Delta(T) \cos\left(\frac{\varphi_{d,r}}{2}\right)$$

For $\varepsilon_{gap} = 50\mu eV$
 $v > 15GHz$

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Temperature-to-phase conversion: model



$$I_C^{d,r}(T, \varphi_{d,r}) = \frac{\pi \Delta(T)}{e R_N^{d,r}} \cos\left(\frac{\varphi_{d,r}}{2}\right) \int_{\Delta(T) \cos(\varphi_{d,r}/2)}^{\Delta(T)} d\varepsilon \times \frac{1}{\sqrt{\varepsilon^2 + \Delta^2(T) \cos^2\left(\frac{\varphi_{d,r}}{2}\right)}} \tanh\left(\frac{\varepsilon}{2k_B T}\right)$$

T_e -induced suppression of supercurrent for $\Phi_{ext} \neq 0$

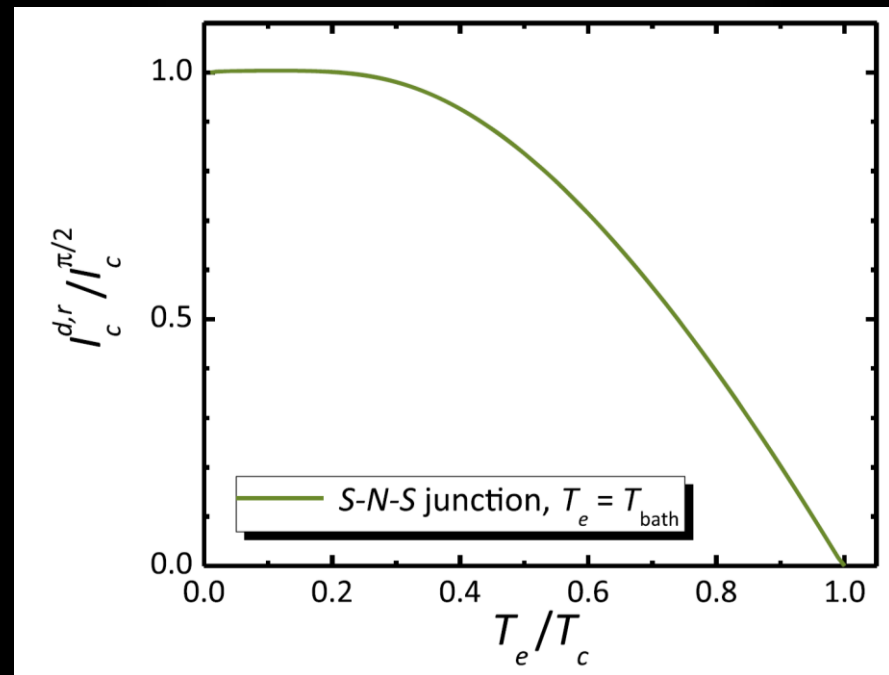


Variation finite phase drop across the WLs

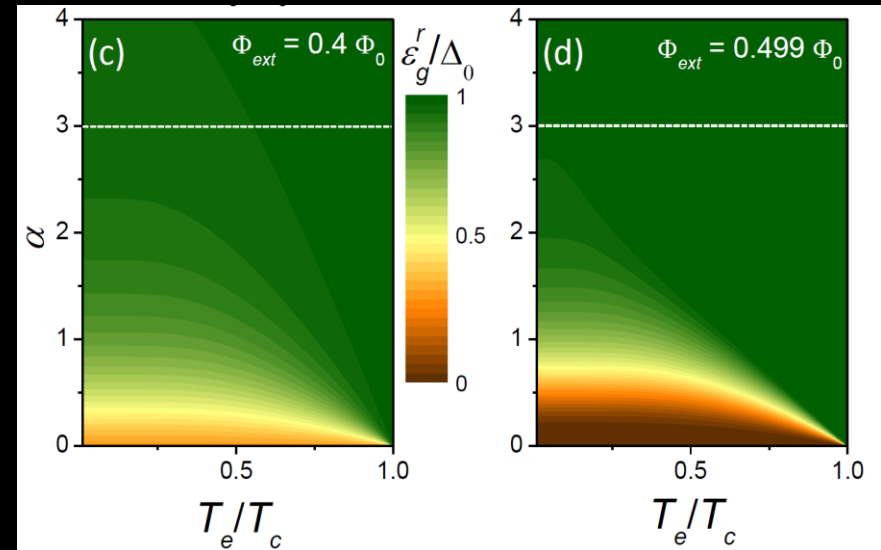
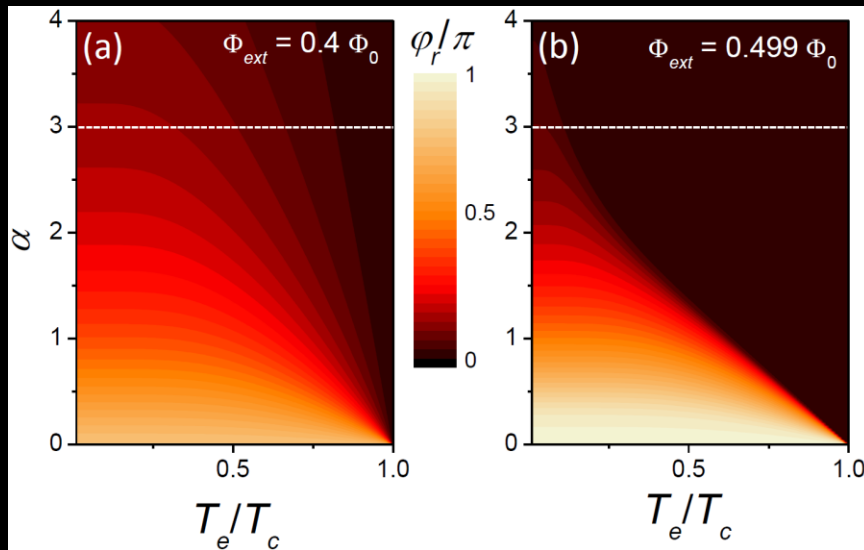


Conservation of circulating supercurrent

$$I_C^d(T_e, \varphi_d) = I_C^r(T_{bath}, \varphi_r)$$



Temperature-to-phase conversion: model



$$\alpha = \frac{R_N^d}{R_N^r} = \frac{L_J^d}{L_J^r}$$

$$\varepsilon_{gap}^{d,r} = \Delta(T) \cos\left(\frac{\varphi_{d,r}}{2}\right)$$

- For $\alpha \gg 1$, phase drop will occur predominantly across read-out JJ
- For $\alpha \ll 1$, phase drop will occur predominantly across detector JJ

Gapped metal ($\varphi = 0$)

Gapless metal ($\varphi = \pi$)

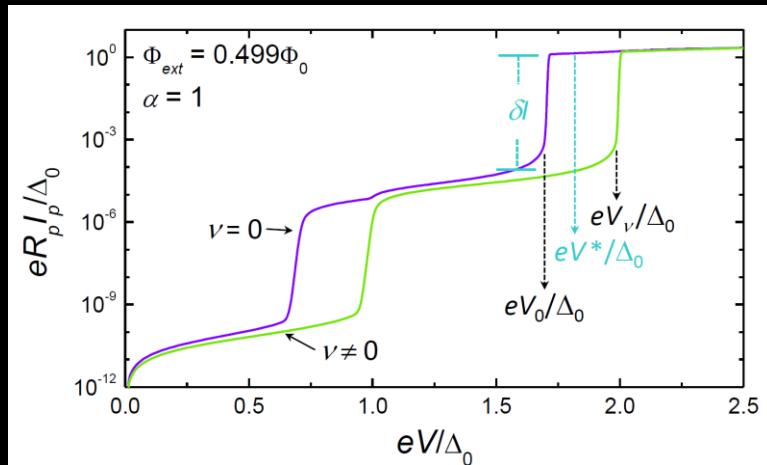
Temperature-to-phase conversion: SQUIPT response

$$I_P(V) = \frac{1}{ewR_N} \int_{\frac{L-w}{2}}^{\frac{L+w}{2}} dx \int_{-\infty}^{\infty} d\varepsilon \mathcal{N}_N^r(x, \varepsilon, \varphi_r) \mathcal{N}_S^p(\tilde{\varepsilon}) F(\varepsilon, \tilde{\varepsilon})$$

$$\mathcal{N}_S^p(\varepsilon, T_e) = \left| \text{Re} \left[\frac{(\varepsilon + i\Gamma)}{\sqrt{(\varepsilon + i\Gamma)^2 - \Delta^2(T_e)}} \right] \right|$$

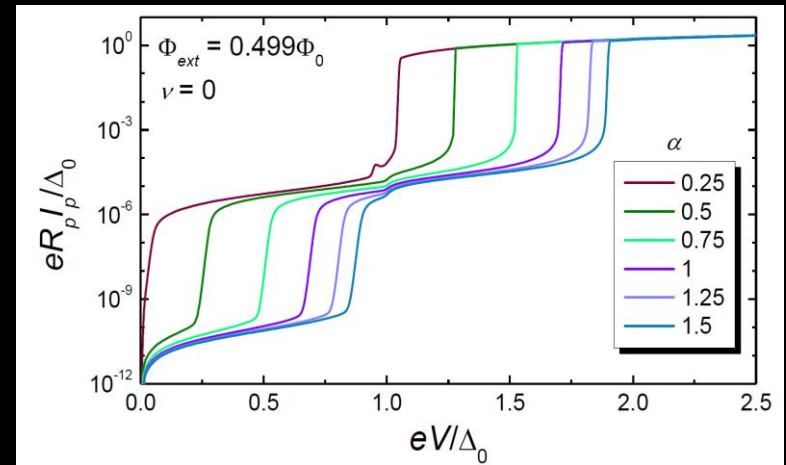
$$F(\varepsilon, \tilde{\varepsilon}) = [f_0(\tilde{\varepsilon}) - f_0(\varepsilon)]$$

Nature Phys. (2010), PRB (2011), PRB (2011), PRAppl (2014), APL (2015)



Absorption of a photon:

reduction of tunneling current



Different α :

selection of operation voltage

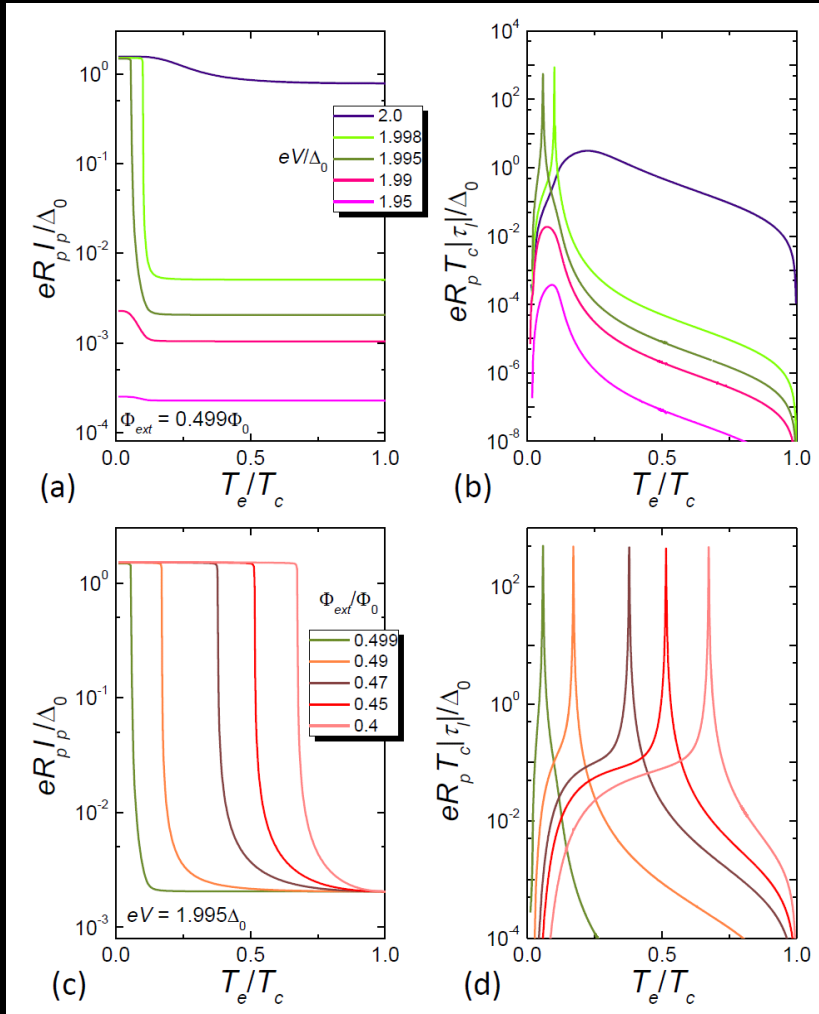
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Temperature-to-phase conversion: thermometer

Current

Transfer function

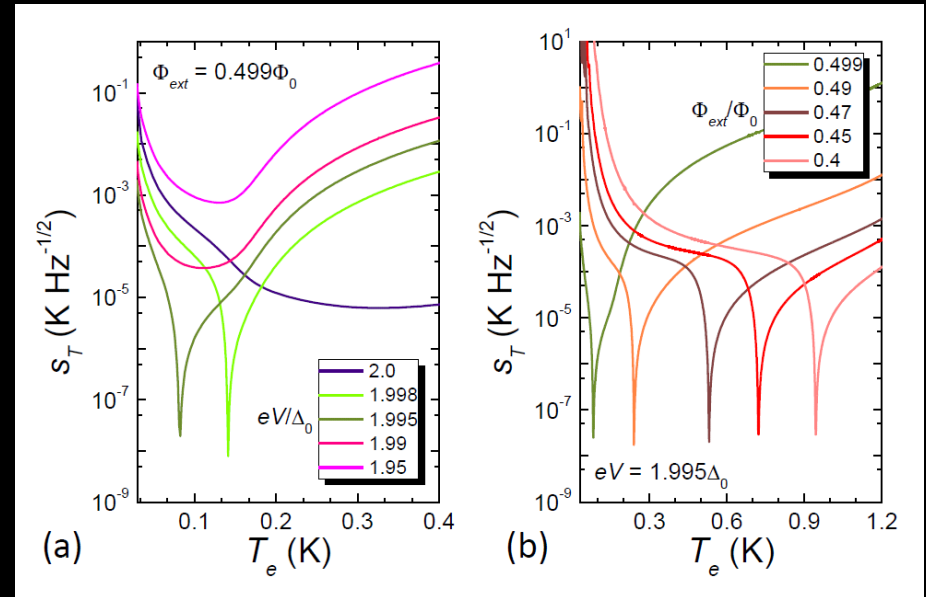
Temperature noise per unit of bandwidth



$$s_T = \frac{\sqrt{S_I}}{|\tau_I|}$$

Low frequency current noise

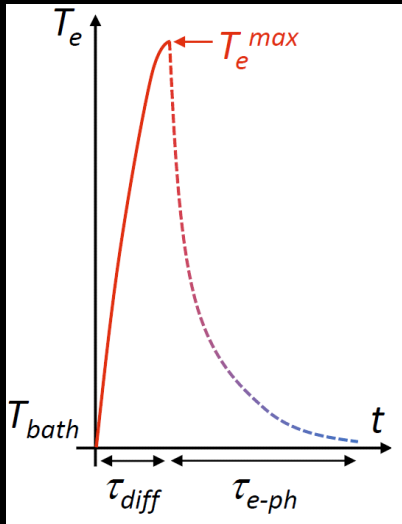
$$S_I(V) = 2eI_P(V) \coth\left(\frac{eV}{2k_B T_{bath}}\right)$$



$$s_T \sim tens \text{ nK}/\sqrt{Hz}$$

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Temperature-to-phase conversion: calorimeter



Energy required to change T_e in the detection junction:

$$\delta U(T_e, T_{bath}) = \delta Q^d(T_e, T_{bath}) + \delta W^d(T_e, T_{bath}) + \delta W^r(T_{bath})$$

δQ : heat of quasiparticles

δW : work to change phase difference

Total electron heat capacity

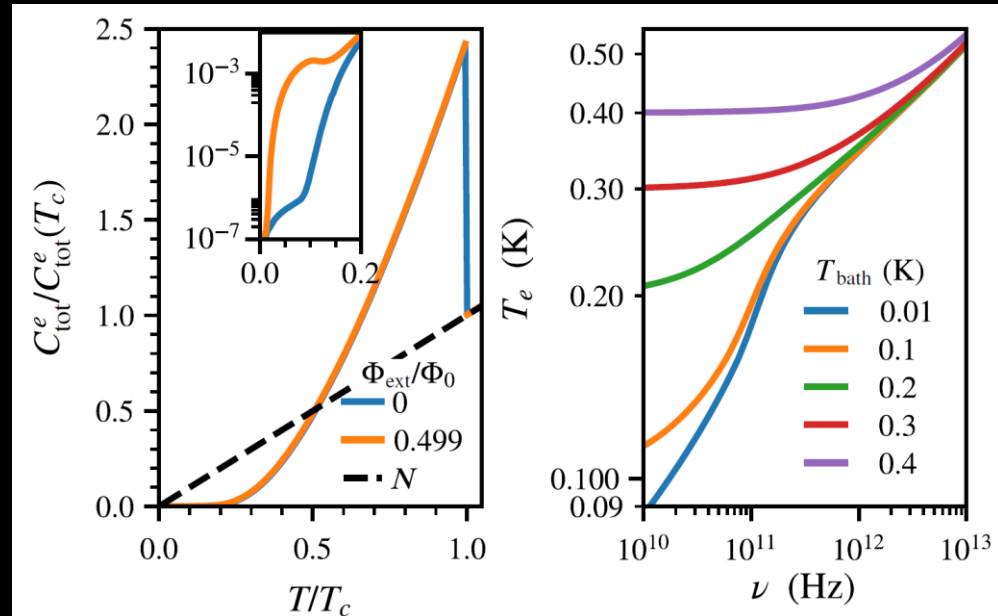
$$C_{tot}^e(T_e, T_{bath}) = \frac{d}{dT_e} \delta U(T_e, T_{bath})$$

$$\tau_{diff} = \frac{L^2}{D} \sim 10^{-12} \text{ s}$$

$$\tau_{e-ph} \approx \frac{k_B^2 v_F}{0.34 \Sigma} T_{bath}^{-3} \sim [10^{-4} - 10^{-7}] \text{ s}$$

Photon energy

$$\delta U = h\nu$$



Temperature-to-phase conversion: calorimeter

e-ph relaxation haltime

$$\tau_{1/2}(\nu, T_{\text{bath}}) = \int_{[T_e^{\text{max}}(\nu) + T_{\text{bath}}]/2}^{T_e^{\text{max}}(\nu)} dT_e \frac{C_{\text{tot}}^e(T_e)}{\dot{Q}_{e-ph}^{\text{tot}}(T_e, T_{\text{bath}})}$$

$$\dot{Q}_{e-ph} = -\frac{\Sigma_d V_d}{96\zeta(5)k_B^5} \int_{-\infty}^{\infty} dEE \int_{-\infty}^{\infty} d\varepsilon M_{E,E+\varepsilon}(\Delta_0(T_e), \Gamma_0)$$

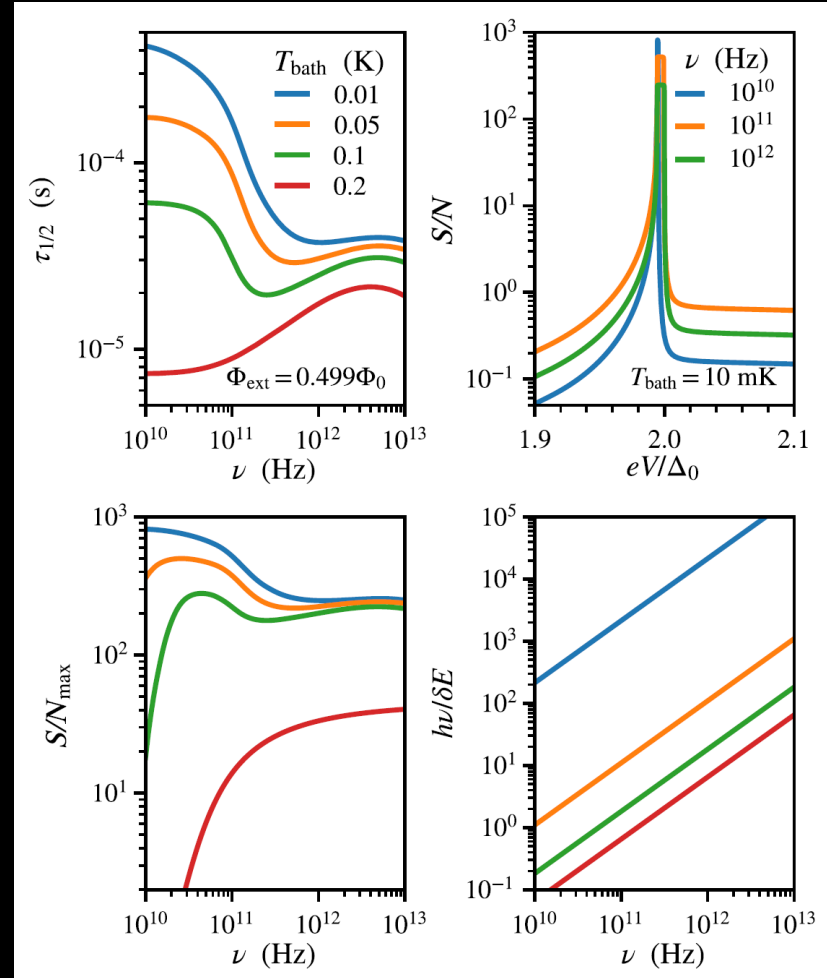
$$\times \varepsilon |\varepsilon| \left[\coth\left(\frac{\varepsilon}{2k_B T_{\text{bath}}}\right) (f_E - f_{E+\varepsilon}) - f_E f_{E+\varepsilon} + 1 \right]$$

signal to noise ratio

$$\frac{S}{N}(T_e(h\nu)) = \frac{I_P(V, T_{\text{bath}}) - I_P(V, T_e(h\nu))}{\sqrt{S_I(V, T_{\text{bath}})}\sqrt{\omega}}$$

resolving power

$$\frac{h\nu}{\delta E}(\nu, T_{\text{bath}}) = \frac{h\nu}{4\sqrt{2\ln 2 k_B T_{\text{bath}}^2 C_{\text{tot}}^e(\varphi_d, T_{\text{bath}})}}$$



Resolving power > 100 for frequencies above 100GHz at 10mK

Suited for microwave and Far InfraRed single-photon detection

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Temperature-to-phase conversion: bolometer

energy balance equation

$$P_{opt} + \dot{Q}_{e-ph}^{tot}(T_e, T_{bath}) = 0$$

$$\dot{Q}_{e-ph} = -\frac{\Sigma_d V_d}{96\zeta(5)k_B^5} \int_{-\infty}^{\infty} dEE \int_{-\infty}^{\infty} d\varepsilon M_{E,E+\varepsilon}(\Delta_0(T_e), \Gamma_0) \times \varepsilon|\varepsilon| \left[\coth\left(\frac{\varepsilon}{2k_B T_{bath}}\right) (f_E - f_{E+\varepsilon}) - f_E f_{E+\varepsilon} + 1 \right]$$

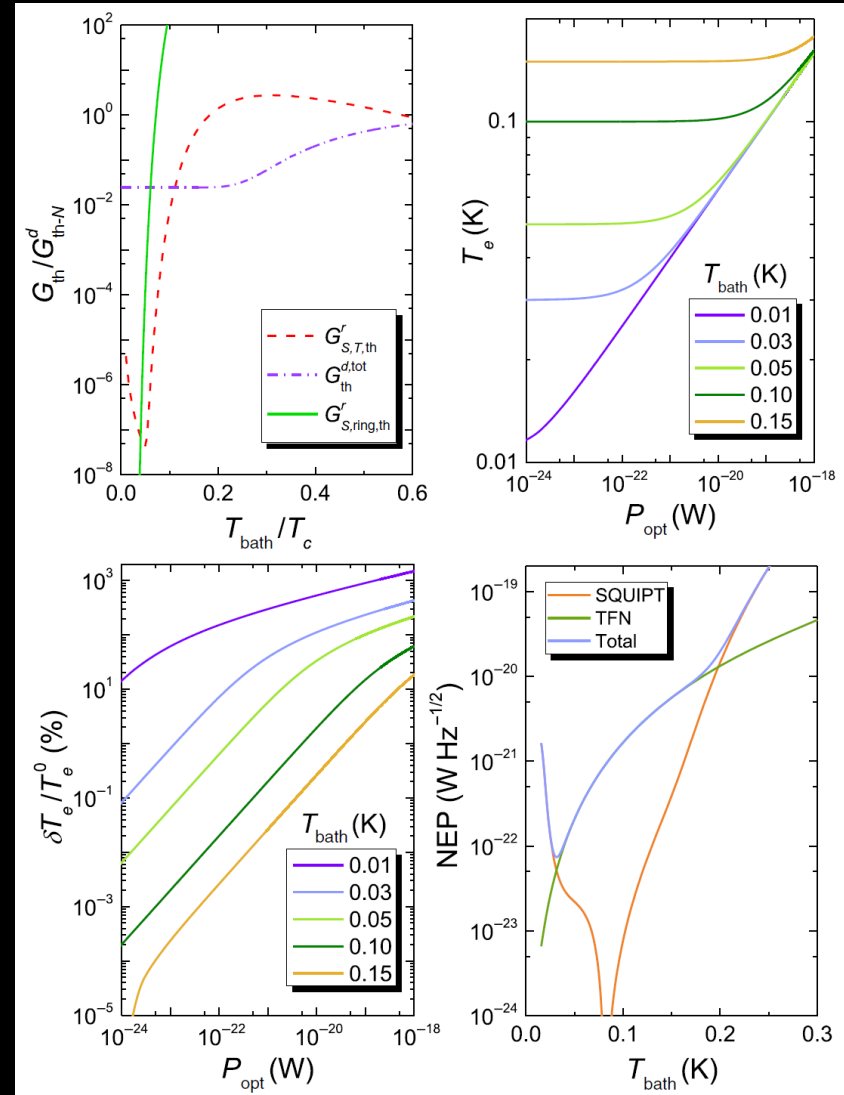
noise equivalent power

$$NEP = \sqrt{4k_B T_{bath}^2 G_{th}^{tot}}$$

e-ph thermal conductance

$$G_{th}^{tot} = \frac{\partial \dot{Q}_{e-ph}}{\partial T_e}$$

NEP below 10⁻²²WHz^{-1/2} at 50mK



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Temperature-to-phase conversion: performances

- T -to- ϕ conversion due to kinetic inductance changes & SQUIPT spectroscopy
- Thermal sensitivity amplified by **small thermal capacity** (superconductivity & small volume)
- As electron thermometer: sensitivity up to **tens of nK/Hz^{1/2}** tunable with Φ in the 10mK...1K temperature range
- Calorimetric mode: resolving power $h\nu/\delta E > 10^2$ for 50GHz...10THz photons
- Bolometric mode: NEP $\sim 10^{-22}$ W/Hz^{1/2} below 50 mK
- Attractive as MW and FIR radiation detector for open problems in **astrophysics & quantum electronic circuits**

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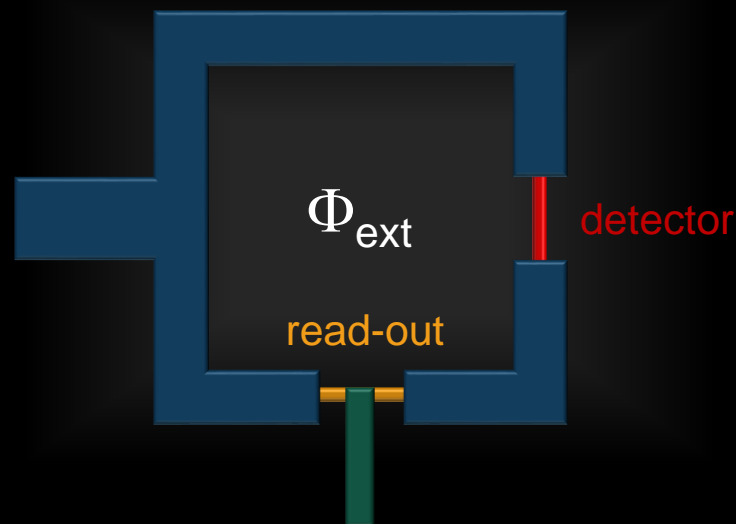
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