

# Unpolarized TMD distributions and their evolution from DY and SIDIS data

Alexey Vladimirov

Regensburg University



Universität Regensburg

# Introduction

I am presenting a (global) fit of the DY and SIDIS data and extraction  
non-perturbative functions using TMD-factorization approach

## Theory details

- ▶ "CSS-like": same equations, but different solution –  $\zeta$ -prescription
- ▶ Optimal TMD distribution as the boundary
- ▶ NNLO (or N<sup>3</sup>LO) evolution+coefficient functions.
- ▶ NNLO matching to collinear PDF at  $b \rightarrow 0$

artemide

- ▶ Complete freedom in definition of NP-input
- ▶ Variety of implementation of TMD evolution (“classic” CSS,  $\gamma$ -improved, resummed,...)
- ▶ Tools for theoretical studies (variation bands, cuts, ...)
- ▶ Modular structure
- ▶ FORTRAN 95 + python interface (`harpy`)
- ▶ "Mostly" documented

[github.com/VladimirovAlexey/artemide-public](https://github.com/VladimirovAlexey/artemide-public)

## TMD factorization formula in $\zeta$ -prescription

$$\frac{d\sigma}{dy dQ^2 d^2 \mathbf{q}_T} = \sigma_0 \sum_{ff'} \int d^2 b e^{i(\mathbf{b} \cdot \mathbf{q}_T)} H_{ff'} \left( \frac{Q}{\mu} \right) \left( \frac{Q^2}{\zeta_Q} \right)^{-2\mathcal{D}(Q,b)} F_{f \leftarrow h}(x_1, b) F_{f' \leftarrow h}(x_2, b)$$



Universität Regensburg

TMD factorization formula in  $\zeta$ -prescription

Rapidity  
anomalous dimension  
(universal)

N	q	U	L	T
U	$f_i$			$h_i^\perp$
L			$g_i$	$h_{iL}^\perp$
T	$h_i^\perp$		$g_{iT}$	$h_i h_{iT}^\perp$

$$\frac{d\sigma}{dy dQ^2 d^2 \mathbf{q}_T} = \sigma_0 \sum_{ff'} \int d^2 b e^{i(\mathbf{b} \cdot \mathbf{q}_T)} H_{ff'} \left( \frac{Q}{\mu} \right) \left( \frac{Q^2}{\zeta_Q} \right)^{-2\mathcal{D}(Q,b)} F_{f \leftarrow h}(x_1, b) F_{f' \leftarrow h}(x_2, b)$$

- ▶ TMD distributions and evolution are independent
  - ▶  $\zeta_Q$  is the (known) function of  $\mathcal{D}$
  - ▶ TMD distributions could be (or not) matched to collinear PDFs – it is a part of the modeling
  - ▶  $Q$ -dynamics split from  $q_T$ -dynamics

# TMD factorization formula in $\zeta$ -prescription

Rapidity  
anomalous dimension  
(universal)

$$\frac{d\sigma}{dy dQ^2 d^2 \mathbf{q}_T} = \sigma_0 \sum_{ff'} \int d^2 b e^{i(\mathbf{b} \cdot \mathbf{q}_T)} H_{ff'} \left( \frac{Q}{\mu} \right) \left( \frac{Q^2}{\zeta_Q} \right)^{-2\mathcal{D}(Q,b)}$$

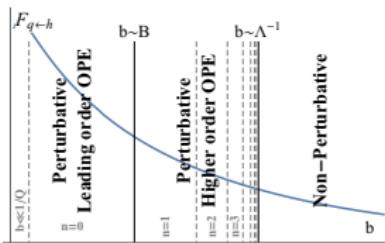
$N^3\text{LO}$   
(universal)

small-b  
 $N^3\text{LO}$

$\mathbf{q}$	$\mathbf{U}$	$\mathbf{L}$	$\mathbf{T}$
$\mathbf{N}$	$f_1$		$h_1^\perp$
$\mathbf{U}$		$g_1$	$h_{1L}^\perp$
$\mathbf{L}$			
$\mathbf{T}$	$f_{1T}^\perp$	$g_{1T}$	$h_1 h_{1T}^\perp$

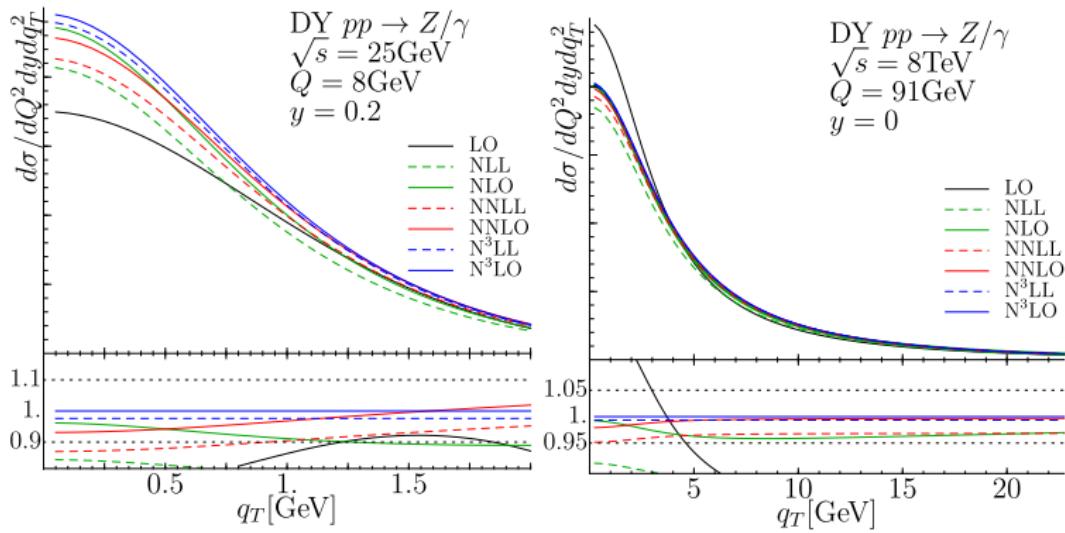
$$F_{f \leftarrow h}(x_1, b) F_{f' \leftarrow h}(x_2, b)$$

small-b  
 $\text{NNLO, NLO, LO}$   
depending on TMD



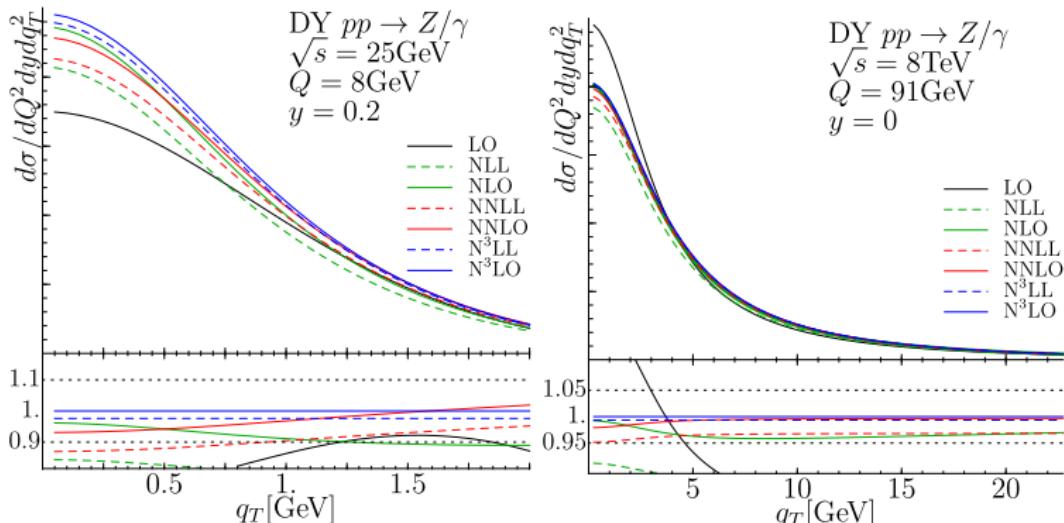
Small-b  $\leftrightarrow$  large  $q_T$   
At small-b everything is perturbative  
turns to resummation

All scales are at  $\mu \sim Q$   
 Perturbative series is saturated



Difference between NNLO and N<sup>3</sup>LO is marginal

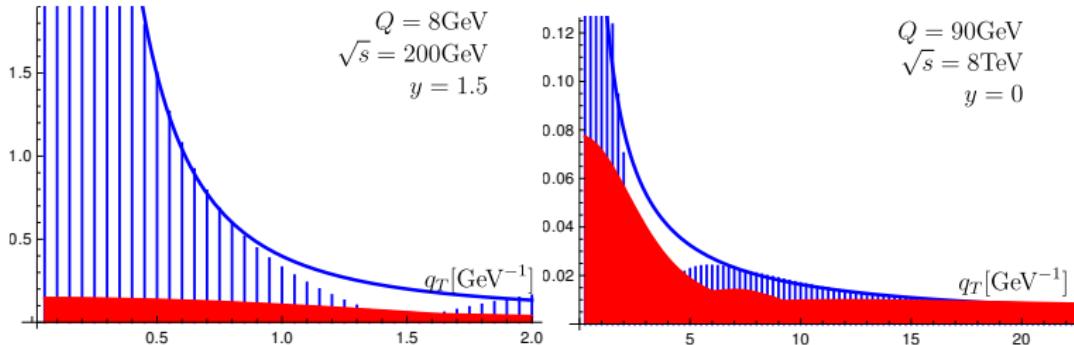
All scales are at  $\mu \sim Q$   
 Perturbative series is saturated



All that follows is done in  
 NNLO or  $N^3\text{LO}$  of TMD factorization  
 + NNLO matching of TMD model to collinear PDF

How large is non-perturbative component?

DY(resummation)/DY(TMD)

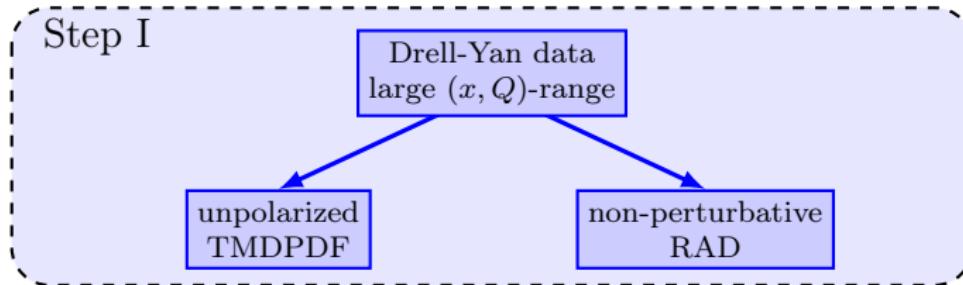


At  $q_T \lesssim 10 \text{ GeV}$  non-perturbative corrections are essential  
and are to be extracted from the experimental data

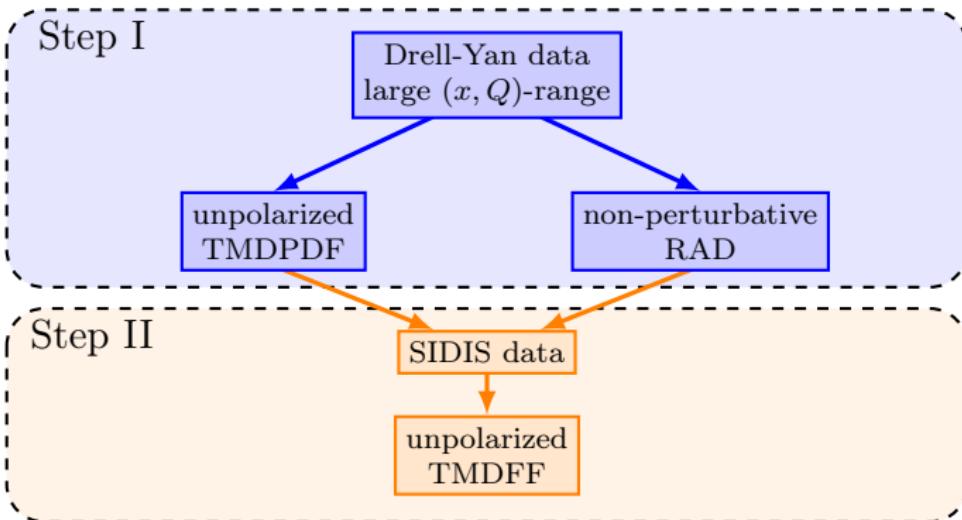


Universität Regensburg

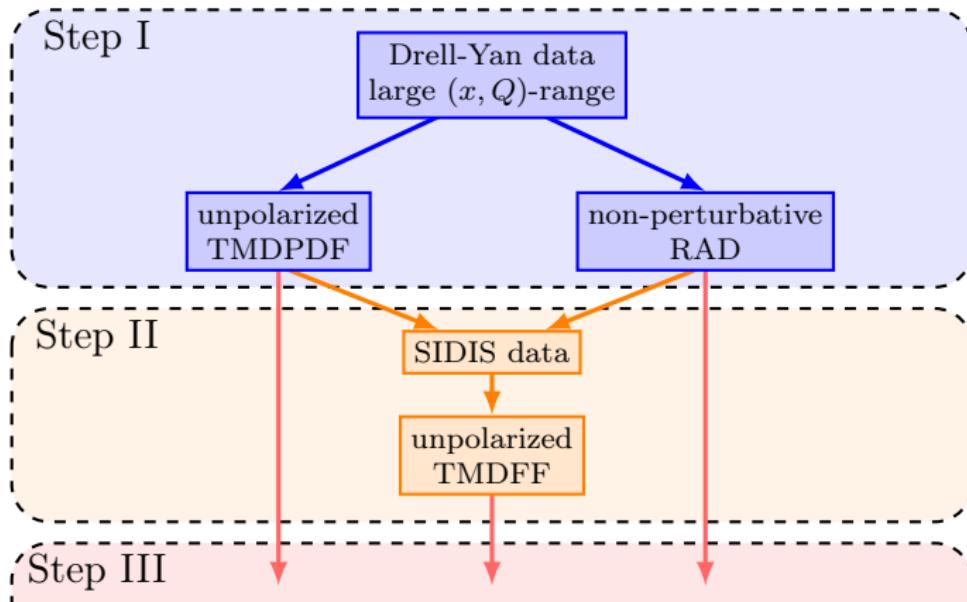
## Fit strategy and the test of universality



## Fit strategy and the test of universality

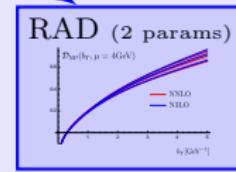
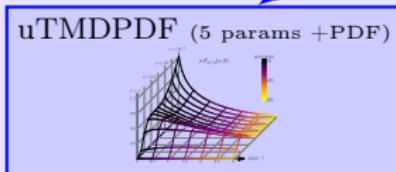
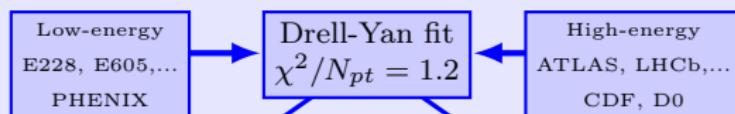


## Fit strategy and the test of universality



## Check of universality

Step I



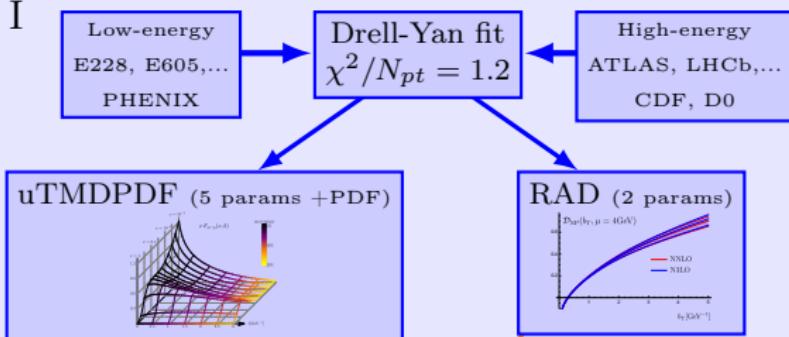
[V.Bertone,I.Scimemi,AV;1902.08474]



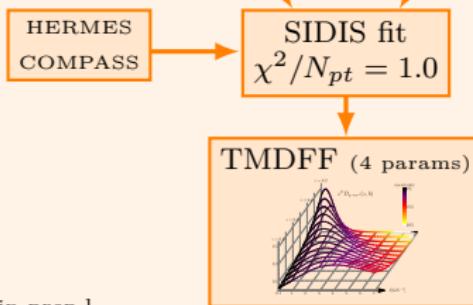
Universität Regensburg

## Check of universality

### Step I

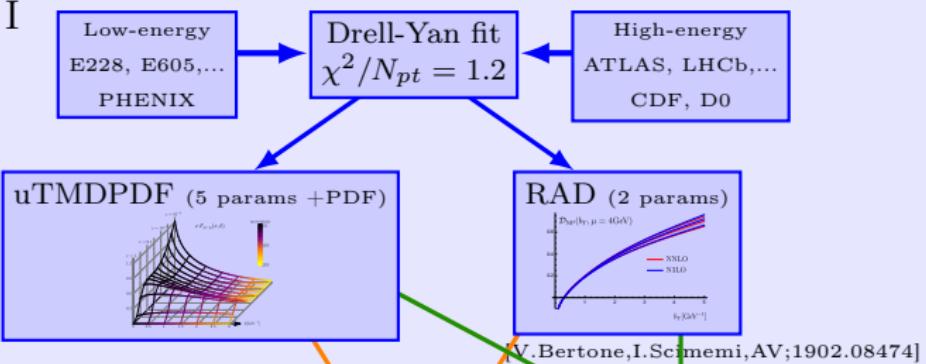


### Step II

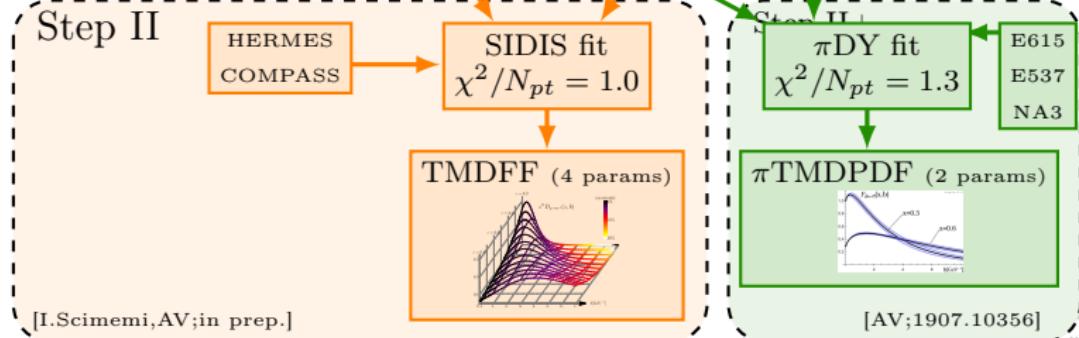


## Check of universality

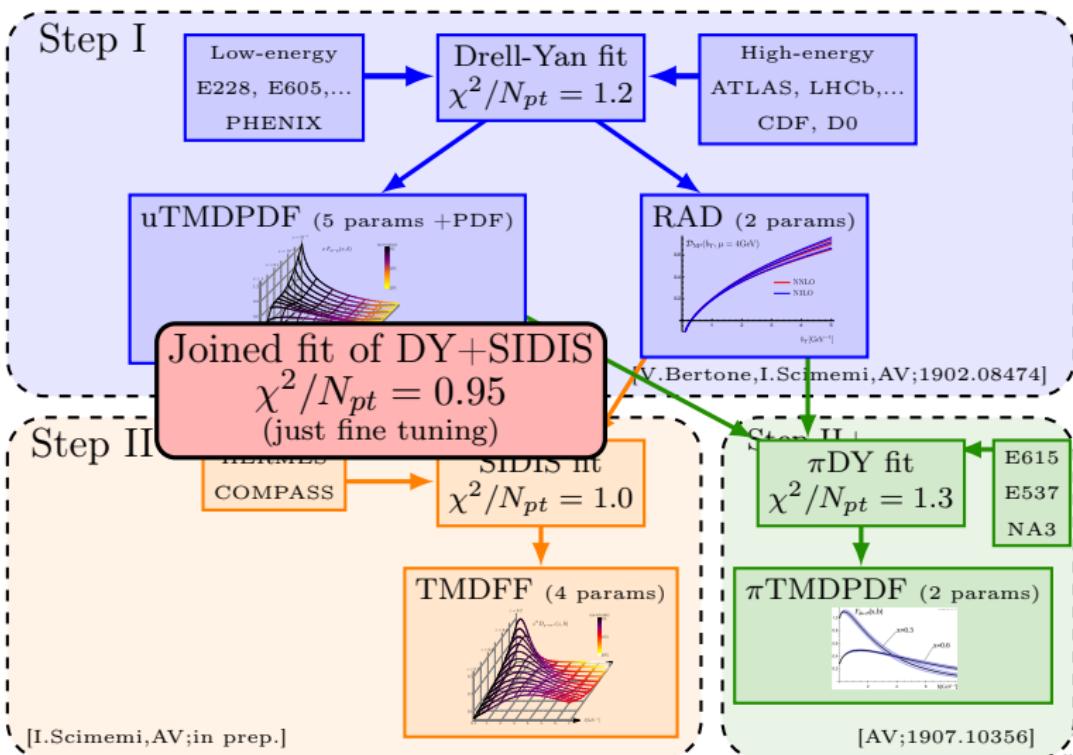
### Step I



### Step II



## Check of universality



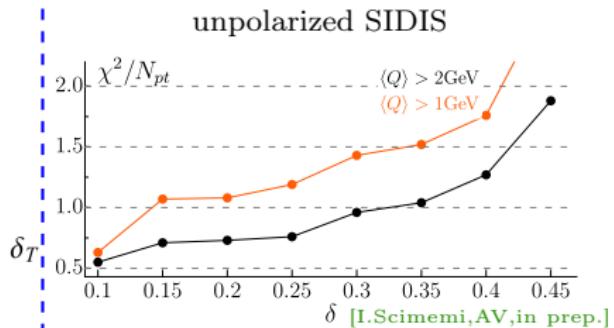
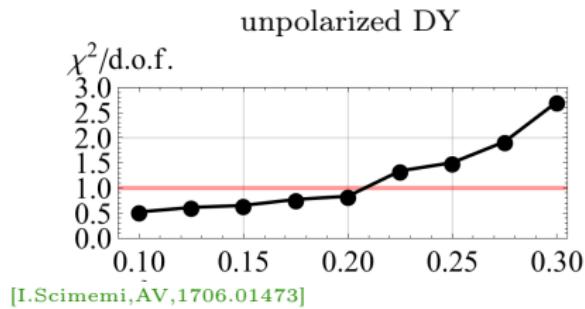
There are plenty of details/questions

- ▶ How to cut the data?
- ▶ Where is the limit of TMD factorization?
- ▶ Power corrections:
  - ▶ Induced (ATLAS, LHCb) (linear in  $q_T$ )
  - ▶ Kinematic
  - ▶ In the definition of collinear frame
  - ▶ Target mass and produced mass
- ▶ Universality and correlations
- ▶ ... many others ...
- ▶ **What do we learn from it?**



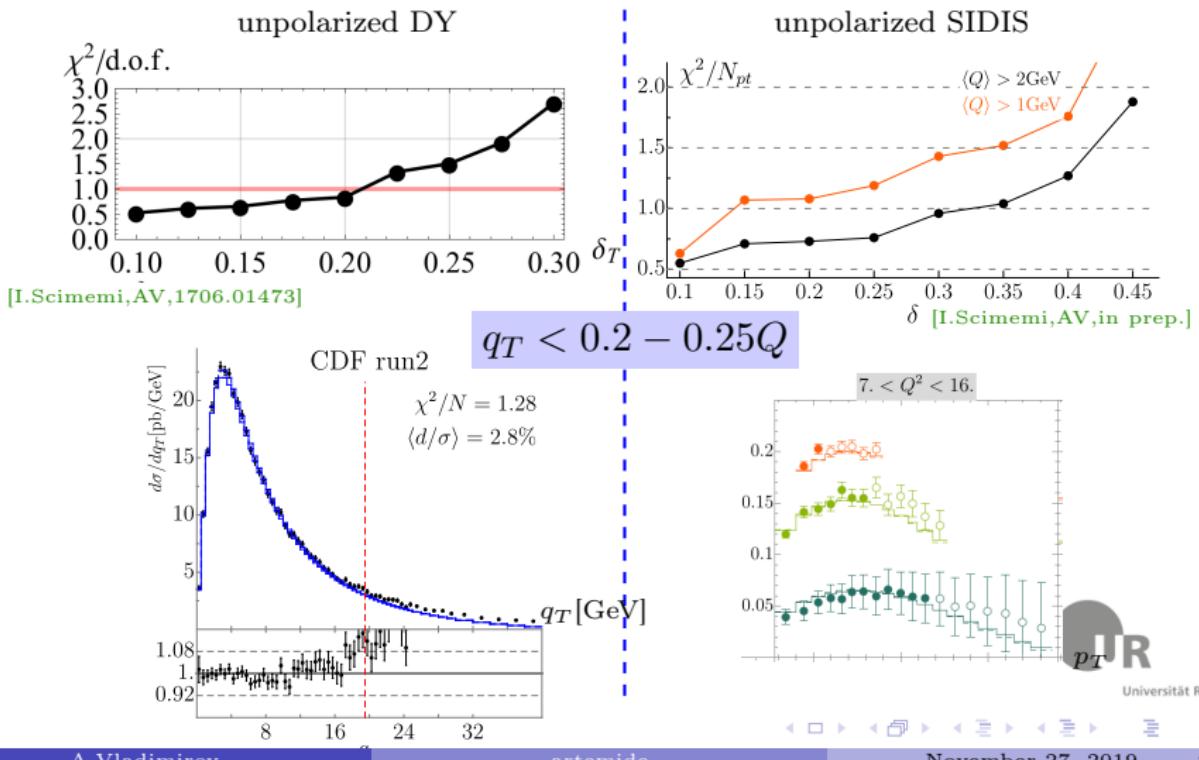
Universität Regensburg

Test by inclusion of the data with  
 $q_T < \delta \cdot Q$



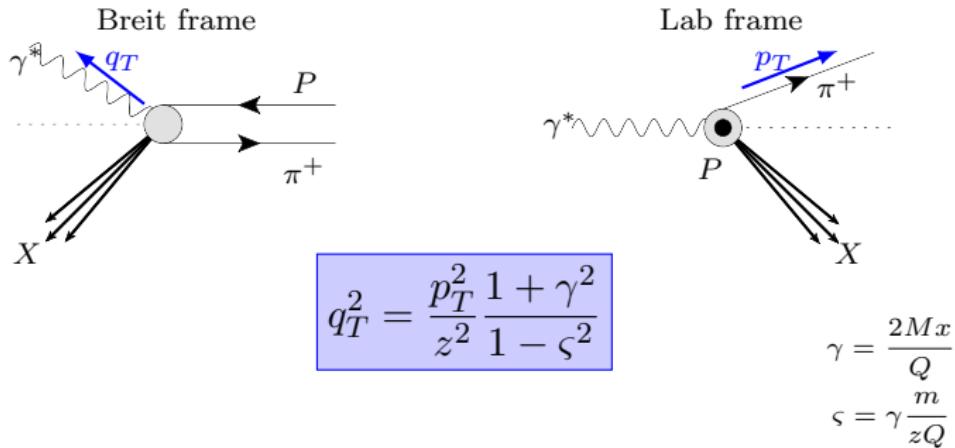
Limits of TMD factorization

Test by inclusion of the data with  
 $q_T < \delta \cdot Q$



## TMD factorization for SIDIS

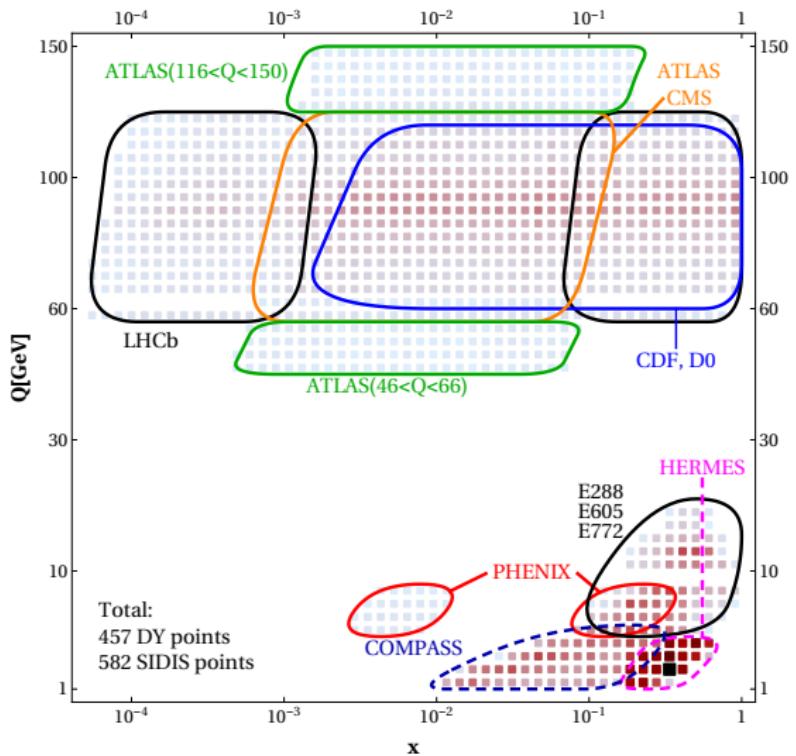
Proof of factorization is done in the Breit frame



In practice:  $q_T < 0.25Q$

- ▶ Most part of data is not TMD factorisable.
- ▶ Low  $z$ 's are not accessible
- ▶ H1, ZEUS data have no TMD points, too low  $z$ .

Data survived after the cut



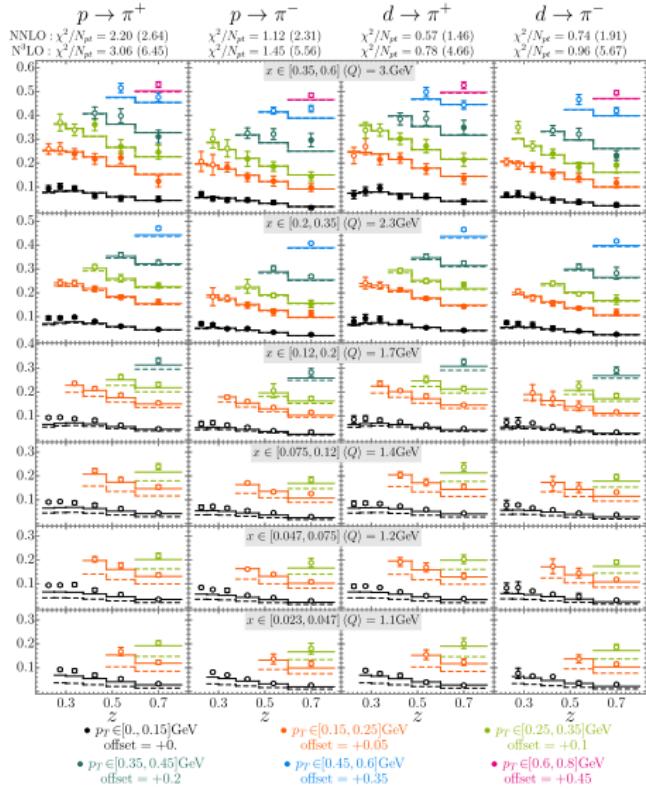
**High energy DY:** 194 points  
**Low energy DY:** 263 points  
**(Low energy) SIDIS:** 582 points

NNLO:  $\chi^2_{global}/N_{pt} = 0.95$

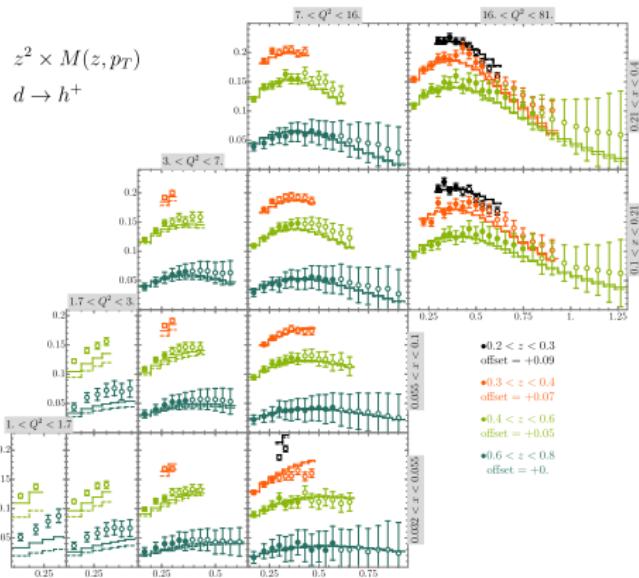
N<sup>3</sup>LO:  $\chi^2_{global}/N_{pt} = 1.05$

artemide v.2.02



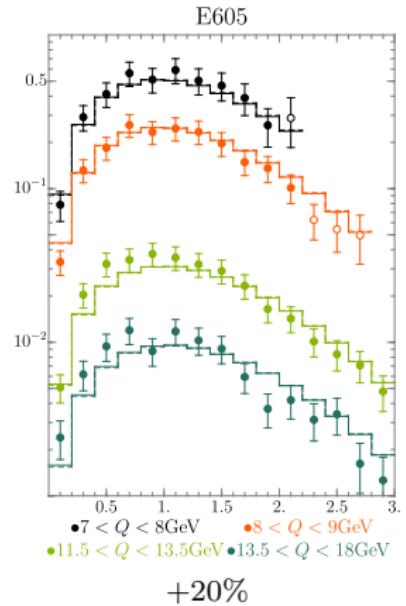
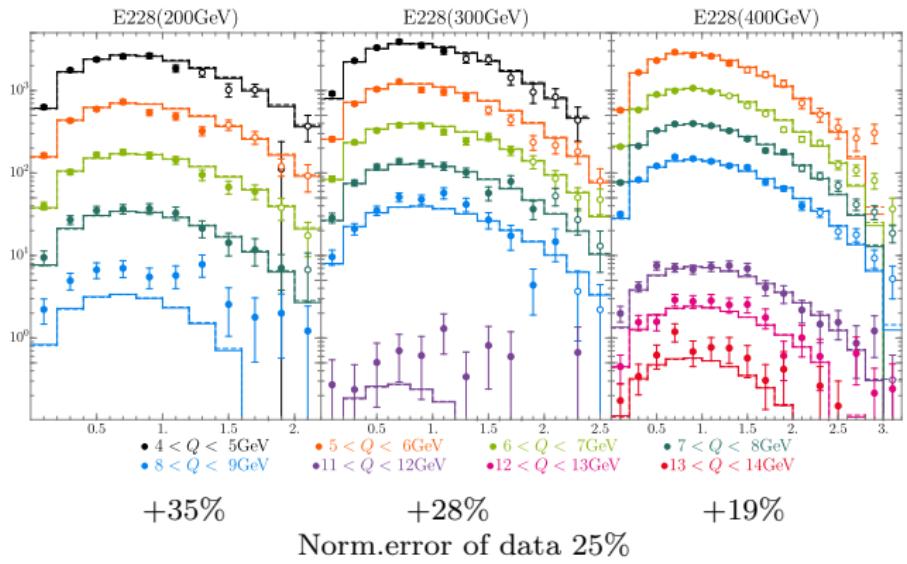


## Example of SIDIS data



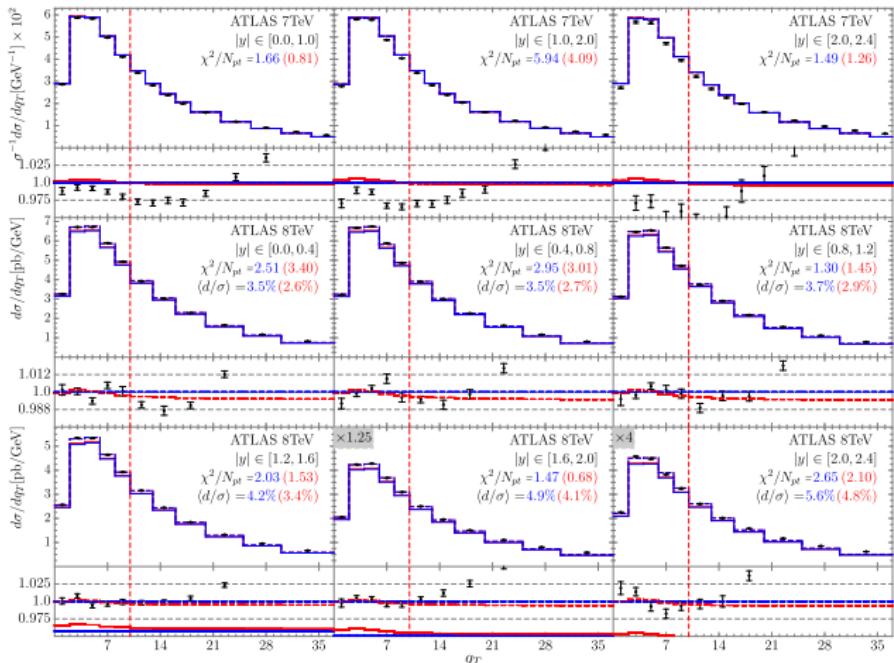
Universität Regensburg

## Example of low-energy DY



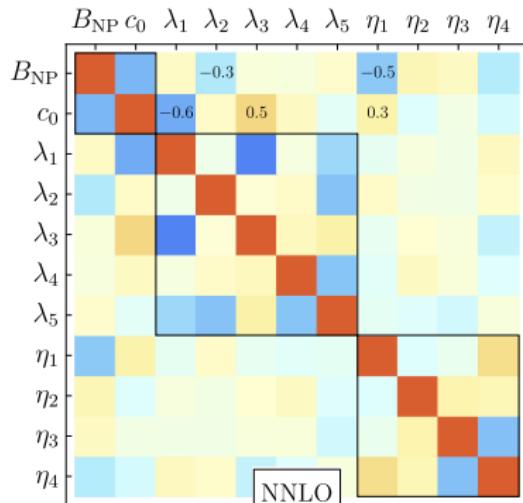
Universität Regensburg

## Example of Z-boson data



Universität Regensburg

Evolution : 2 parameters  
 TMDPDF : 5 parameters + PDF  
 TMDF : 4 parameters + NNFF

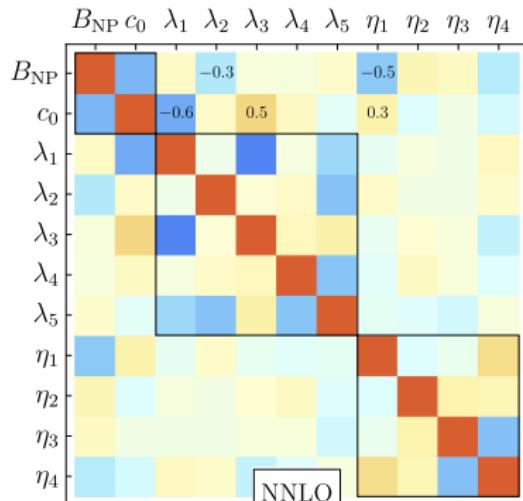


Different NP functions are almost decorrelated



Universität Regensburg

Evolution : 2 parameters  
 TMDPDF : 5 parameters + PDF  
 TMDFF : 4 parameters + NNFF



Different NP functions are almost  
decorrelated

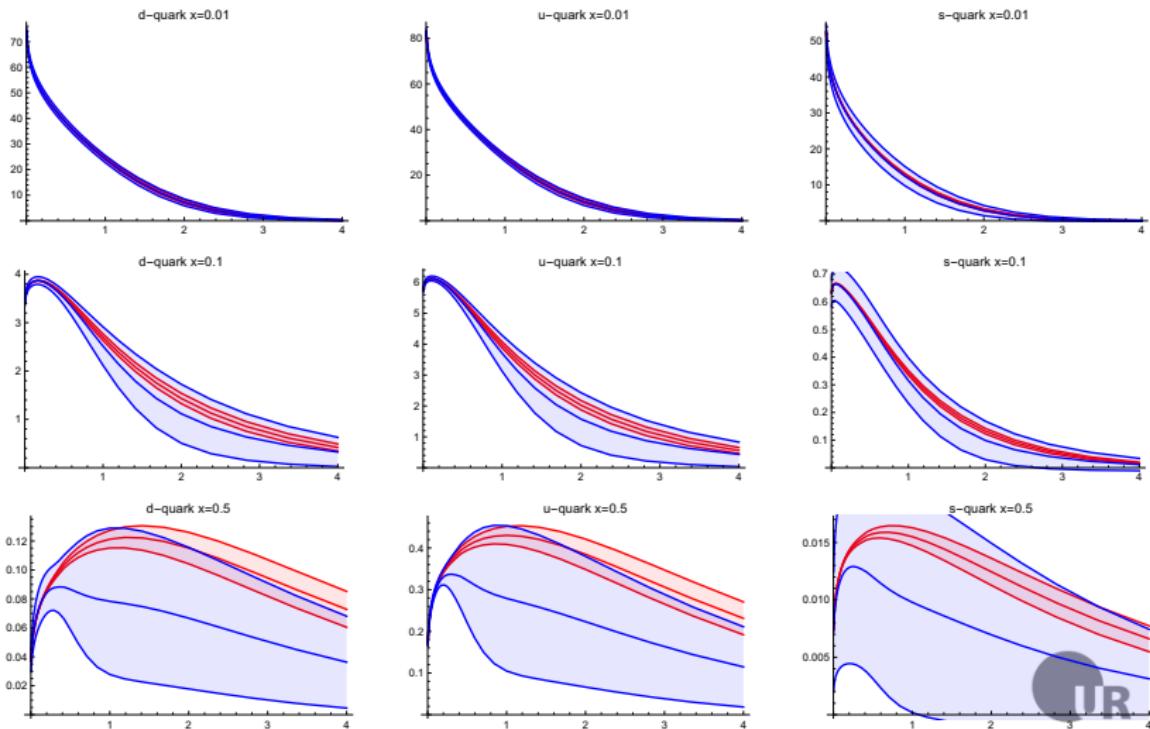
Fit quality essentially depends on the collinear input.

Vary NNPDF within the  $1\sigma$  band

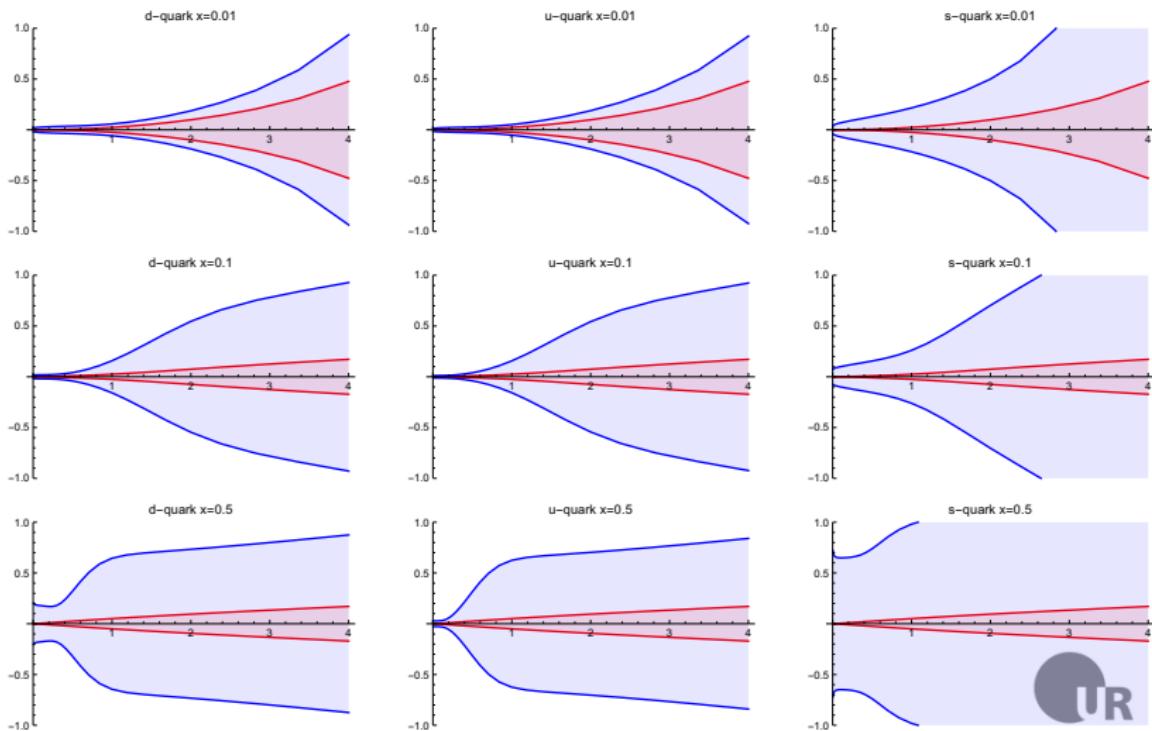
$$\chi^2/N_{pt} \in [0.8, 6.]$$

We cannot estimate accurately the PDF uncertainty.

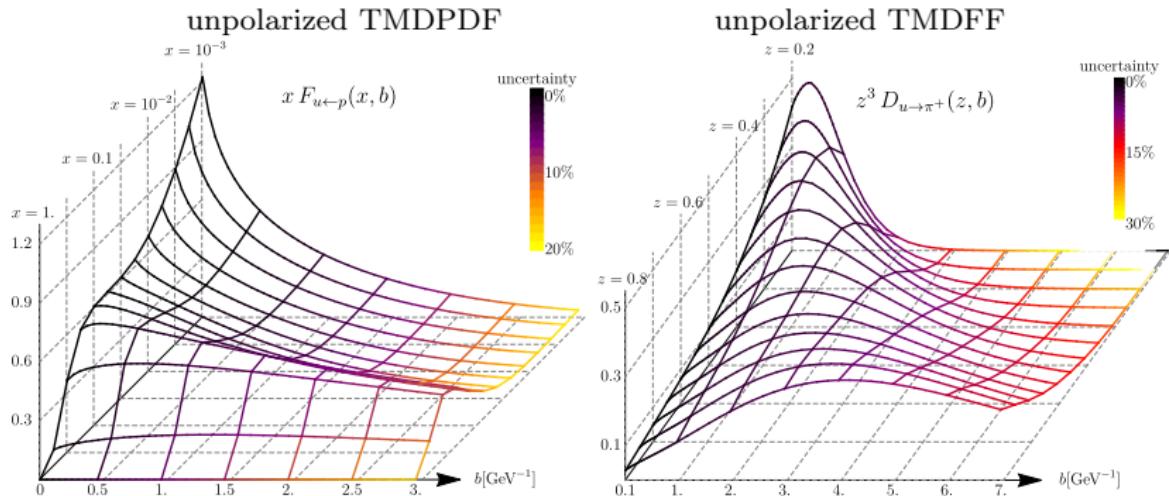
$$F(x, b) = C(x, b, \mu_{OPE}) \otimes f_1(x, \mu_{OPE}) f_{NP}(x, b)$$



$$F(x, b) = C(x, b, \mu_{OPE}) \otimes f_1(x, \mu_{OPE}) f_{NP}(x, b)$$

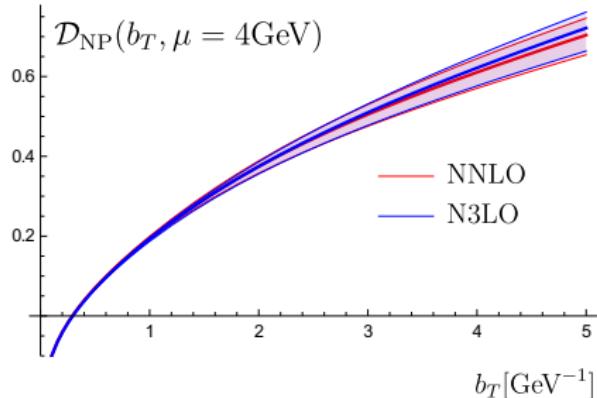


unpolarized TMD-distributions



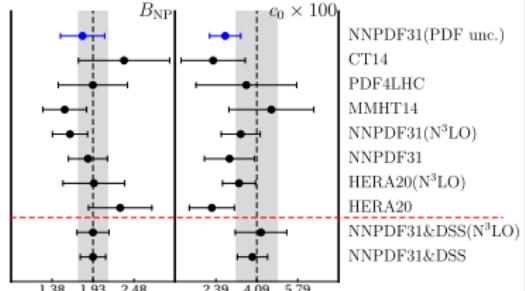
... systematic error 10-20%  
... evolution does not have this systematics

## Universal TMD evolution kernel



$$\mathcal{D}_{\text{NP}}(b, \mu) = \mathcal{D}_{\text{perp}}(b^*, \mu) + c_0 b b^*, \quad b^* = b / \sqrt{1 + b^2 / B_{NP2}}$$

- ▶ Linear asymptotic at  $b \rightarrow \infty$  (ed.assumption)
- ▶ RAD is independent on PDF set



# What could be learned about QCD from TMD physics?



Universität Regensburg

# What could be learned about QCD from TMD physics?

## TMD distributions

- ▶ Distribution of spin, momentum inside of hadron
- ▶ Hadronization properties



Universität Regensburg

# What could be learned about QCD from TMD physics?

## TMD distributions

- ▶ Distribution of spin, momentum inside of hadron
- ▶ Hadronization properties

## Non-perturbative TMD evolution kernel (RAD, CSS-kernel, non-perturbative Sudakov)

- ▶ ???



Universität Regensburg

# What could be learned about QCD from TMD physics?

## TMD distributions

- ▶ Distribution of spin, momentum inside of hadron
- ▶ Hadronization properties

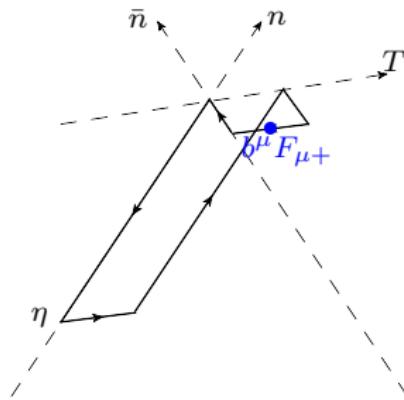
## Non-perturbative TMD evolution kernel (RAD, CSS-kernel, non-perturbative Sudakov)

- ▶ Properties of QCD vacuum



Universität Regensburg

Rapidity anomalous dimension can be defined as a primary object.



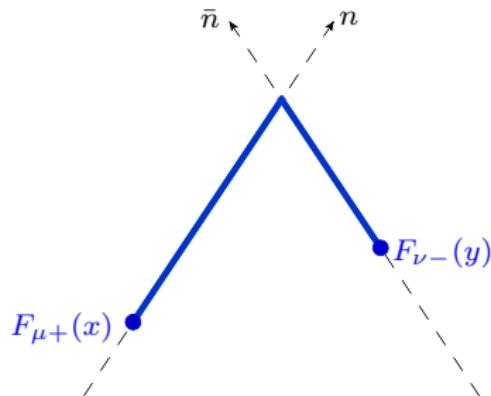
$$\mathcal{D}(b) = \frac{1}{2} \lim_{\eta \rightarrow \infty} \frac{S'(b; \eta \lambda)}{S(b)}$$

$$S'(b) = ig \int_0^b dr \frac{\text{Tr}}{N_c} \langle 0 | F_{b+}(-\lambda n + r) P \exp \left( -ig \int_{C'} dx^\mu A_\mu(x) \right) | 0 \rangle$$

- Route to non-perturbative calculation and modeling.

## Power correction to $\mathcal{D}$

$$\mathcal{D}(\mathbf{b}) = \underbrace{\mathcal{D}_{\text{pert}}(\ln(\mu^2 \mathbf{b}^2))}_{\text{known at N}^3\text{LO}} + \mathbf{b}^2 \mathcal{D}_1(\ln(\mu^2 \mathbf{b}^2)) + \mathbf{b}^4 \mathcal{D}_2(\ln(\mu^2 \mathbf{b}^2)) + \dots$$



$\mathcal{D}_1$  is expressed via 2-point correlators connected by “a minimal distance link”

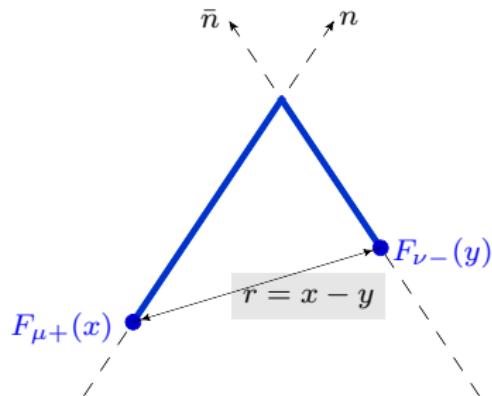
$$g^2 \frac{\text{Tr}}{N_c} \langle 0 | F_{\mu x}(x) [x, 0] [0, y] F_{\nu y}(y) | 0 \rangle = \\ \left( g^{\mu\nu} - \frac{y^\mu x^\nu}{(xy)} \right) \varphi_1(x, y) + (\dots)^{\mu\nu} \varphi_2$$

At LO  $x^2 = y^2 = 0$   
 $\varphi_1(x, y) = \varphi_1(r^2)$   
 $\varphi_2(x, y) = 0$

$$\mathcal{D}_1(\mathbf{b}) = \frac{1}{2} \int_0^\infty \frac{d\mathbf{r}^2}{\mathbf{r}^2} \varphi_1(\mathbf{r}^2)$$

## Power correction to $\mathcal{D}$

$$\mathcal{D}(\mathbf{b}) = \underbrace{\mathcal{D}_{\text{pert}}(\ln(\mu^2 \mathbf{b}^2))}_{\text{known at N}^3\text{LO}} + \mathbf{b}^2 \mathcal{D}_1(\ln(\mu^2 \mathbf{b}^2)) + \mathbf{b}^4 \mathcal{D}_2(\ln(\mu^2 \mathbf{b}^2)) + \dots$$



$\mathcal{D}_1$  is expressed via 2-point correlators connected by “a minimal distance link”

$$g^2 \frac{\text{Tr}}{N_c} \langle 0 | F_{\mu x}(x) [x, 0] [0, y] F_{\nu y}(y) | 0 \rangle = \\ \left( g^{\mu\nu} - \frac{y^\mu x^\nu}{(xy)} \right) \varphi_1(x, y) + (\dots)^{\mu\nu} \varphi_2$$

At LO  $x^2 = y^2 = 0$   
 $\varphi_1(x, y) = \varphi_1(r^2)$   
 $\varphi_2(x, y) = 0$

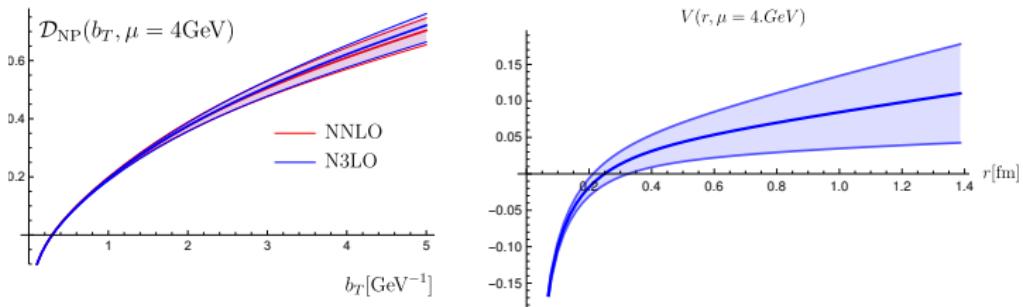
$$\mathcal{D}_1(\mathbf{b}) = \frac{1}{2} \int_0^\infty \frac{d\mathbf{r}^2}{\mathbf{r}^2} \varphi_1(\mathbf{r}^2) \lesssim \frac{\pi^2}{72} \frac{G_2}{\Lambda_{\text{QCD}}^2} \approx (0.01 - 0.05) \text{GeV}^2$$

Extracted value:  $\mathcal{D}_1 = 0.022 \pm 0.009 \text{GeV}^2$

To get interpretation – apply a model

## Example

In stochastic vacuum model



$$V(r) = r \frac{\pi}{4} \mathcal{D}''(0) + \frac{\mathcal{D}'(0)}{2} + \frac{r^2}{2} \int_r^\infty \frac{dx}{x^2} \frac{\mathcal{D}'(x)}{\sqrt{x^2 - r^2}} + \dots$$

QCD string tension:  $\sigma = 0.014 \pm 0.01 \text{GeV}^2$  vs.  $\sim 0.1 - 0.2 \text{GeV}^2$



Universität Regensburg

# Conclusion

TMD factorization is consistent and universal approach

I have demonstrated

- ▶ Large bulk of data (DY+SIDIS) supports it
- ▶ Extracted NP-functions are (almost) uncorrelated
- ▶ Previous estimation of Limits for TMD factorization confirmed  $q_T/Q < 0.25$
- ▶ Perfect perturbative stability
- ▶ Target mass corrections and proper definition of the kinematic variables helps

Plenty of points to improve

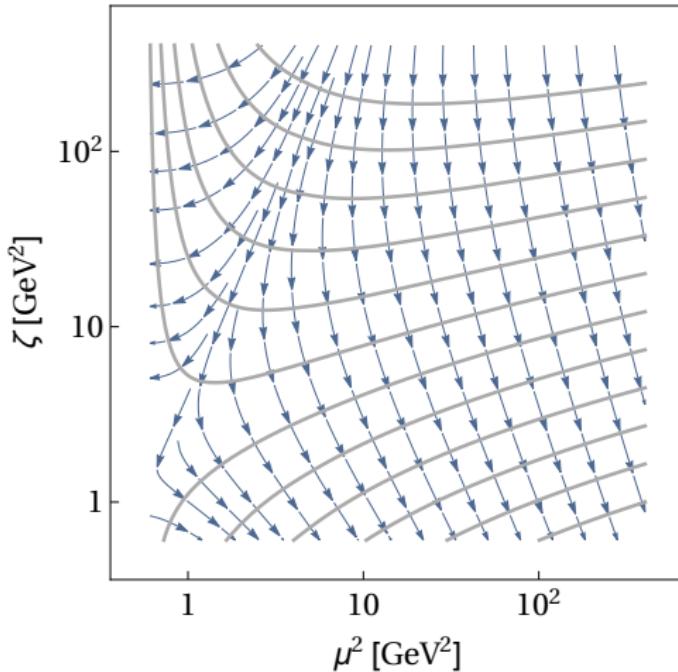
- ▶ Power corrections and higher  $q_T$
- ▶ Realistic uncertainties and mode-bias
- ▶ More processes
- ▶ ...

# Backup slides



Universität Regensburg

TMD distribution is not defined by a scale  $(\mu, \zeta)$   
It is defined by an equipotential line.

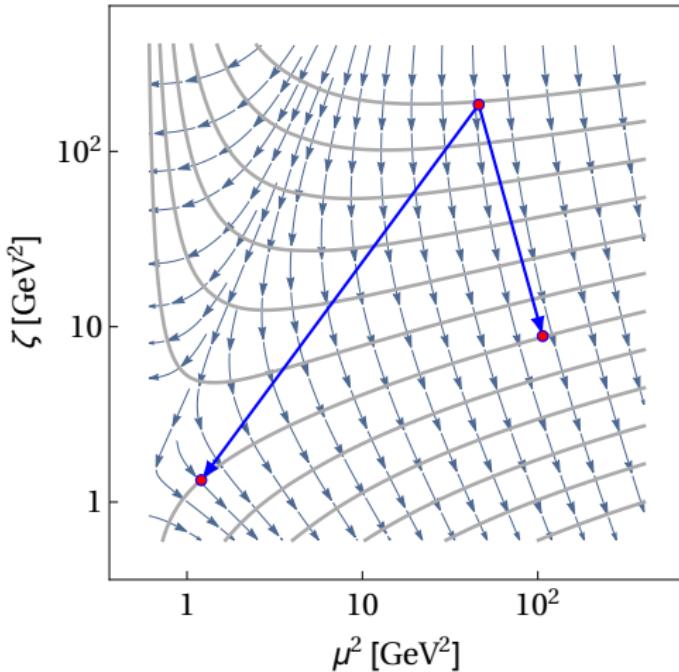


The scaling is defined by  
**a difference between scales**  
a difference between potentials



Universität Regensburg

TMD distribution is not defined by a scale  $(\mu, \zeta)$   
It is defined by an equipotential line.



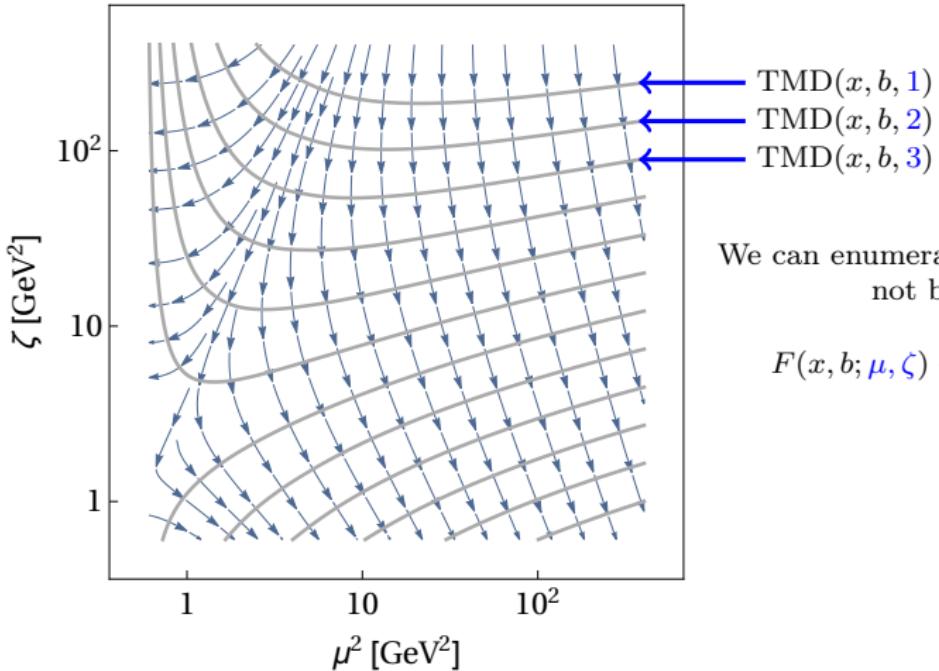
The scaling is defined by  
**a difference between scales**  
a difference between potentials

Evolution factor to both points  
is the same  
although the scales are  
different by  $10^2 \text{ GeV}^2$



Universität Regensburg

TMD distributions on the same equipotential line are equivalent.



We can enumerate them by a lines  
not by  $(\mu, \zeta)$

$$F(x, b; \mu, \zeta) \rightarrow F(z, b; \text{line})$$



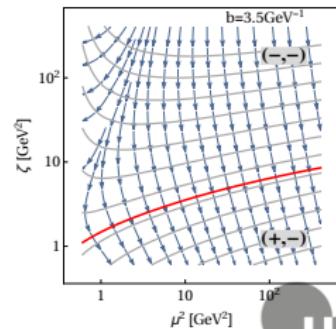
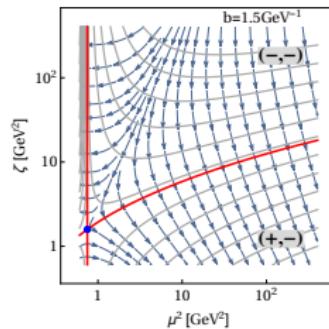
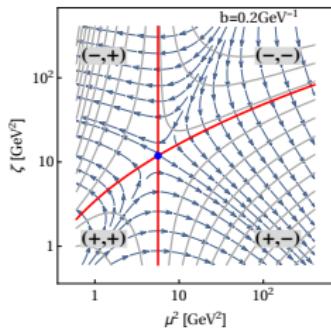
Universität Regensburg

There is a unique line which passes through all  $\mu$ 's

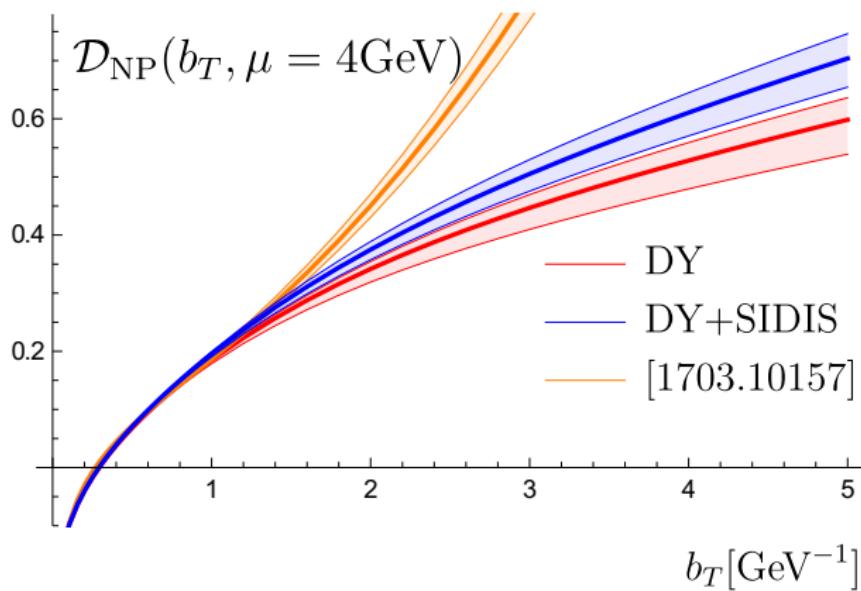
## The optimal TMD distribution

$$F(x, b) = F(x, b; \mu, \zeta_\mu)$$

where  $\zeta_\mu$  is the special line.



## Comparison with other evolutions



Universität Regensburg