

Diffractive leptonproduction of ρ and ϕ light vector mesons via small-x unintegrated gluon density

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Outline

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 - Motivation
 - Unintegrated Gluon Distribution (UGD)
 - Leptoproduction of light vector mesons
- 2 **Theoretical framework**
 - Helicity Amplitudes in κ_T -factorization
 - UGD models
 - Cross section and $\mathbf{b}(Q^2)$ -slope parametrization
- 3 **Results**
 - Numerical results
- 4 **Conclusions and Outlook**
 - Conclusions
 - Outlook

Motivation

- ▶ **Parton densities** are relevant to the search for **new Physics**

They describe the internal structure of the nucleon in terms of its elementary components (quarks and gluons)

⇒ enter the expression for cross sections

⇒ nonperturbative objects

⇒ can be extracted from experiments through global fits

- ▶ Several types of distributions...

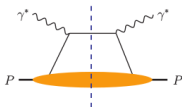
- exhibit particular **universality properties**
- obey distinct **evolution equations**
- respect different types of **factorization theorems**

...A brief overview

Integrated parton densities:

► PDF (or collinear) factorization

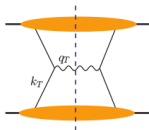
- inclusive processes
- $\kappa_T \sim$ hardest scale



Unintegrated parton densities:

► TMD factorization

- inclusive processes
- $\kappa_T \ll$ hardest scale



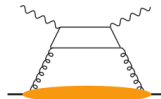
► GPD factorization

- exclusive processes
- skewness effects



► κ_T -factorization (or small-x factorization)

- inclusive or exclusive processes
- small-x, large κ_T
- **Unintegrated gluon distribution**



What is the UGD?

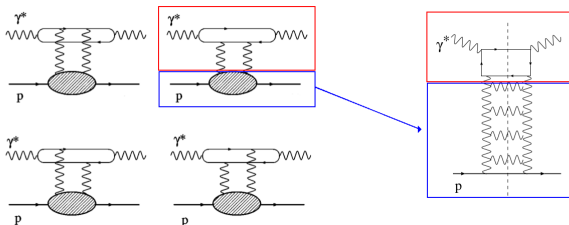
- DIS: conventionally described in terms of PDFs
- less inclusive processes: need to use distributions unintegrated over the parton κ_T

- example: **virtual photoabsorption** in κ_T -**factorization**

$$\sigma_{\text{tot}}(\gamma^* p \rightarrow X) = \text{Im}_s \{ \mathcal{A}(\gamma^* p \rightarrow \gamma^* p) \} \equiv \Phi_{\gamma^* \rightarrow \gamma^*} \circledast \mathcal{F}(x, \kappa^2)$$

- $\mathcal{F}(x, \kappa^2)$ is the **unintegrated gluon distribution (UGD)** in the proton

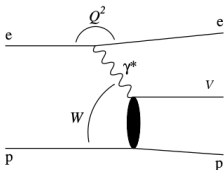
- small- x limit: $\text{UGD} = [\text{BFKL gluon ladder}] \circledast [\text{proton impact factor}]$



Leptonproduction of ρ and ϕ mesons at HERA

$e - p$ collisions provide

$$\gamma^* + \text{proton} \longrightarrow V + \text{proton} \quad \dots \text{exclusive process!}$$



- $V: \rho, \phi$
- High-energy regime:
 $s \equiv W^2 \gg Q^2 \gg \Lambda_{\text{QCD}}^2 \implies \text{small } x = \frac{Q^2}{W^2}$
- photon virtuality Q is the **hard scale** of the process

- **Process solved in helicity** \implies probe cross sections in the HERA energy range:

$$\text{H1: } 2.5 \text{ GeV}^2 < Q^2 < 60 \text{ GeV}^2 \\ 35 \text{ GeV} < W < 180 \text{ GeV}$$

$$\text{ZEUS: } 2 \text{ GeV}^2 < Q^2 < 60 \text{ GeV}^2 \\ 32 \text{ GeV} < W < 180 \text{ GeV}$$

- HERA data available for $\sigma_L, \sigma_T, \sigma_{TOT}$

[H1 collaboration: F.D. Aaron et al., *JHEP* 05 032 (2010)]

[ZEUS collaboration: S. Chekanov et al, *Nucl. Phys. B* 718 (2005)]

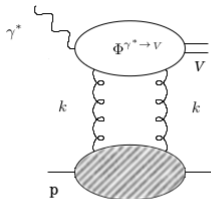
Helicity Amplitudes in κ_T -factorization

- Leading **helicity amplitudes** are known

Assumption:

- $\text{Im}_s \{ \mathcal{A}(\gamma^* p \rightarrow V p) \}$
- same W - and t -dependence for T_{11} and T_{00}
 - same physical mechanism, scattering of small transverse size of dipole on the proton target, at work \implies **κ_T -factorization**

$$T_{\lambda_V \lambda_\gamma}(s; Q^2) = is \int \frac{d^2 \kappa}{(\kappa^2)^2} \Phi^{\gamma^*(\lambda_\gamma) \rightarrow V(\lambda_V)}(\kappa^2, Q^2) \mathcal{F}(x, \kappa^2), \quad x = \frac{Q^2}{s}$$



Interesting transitions:

- $\gamma_L^* \rightarrow V_L \xrightarrow{\text{encoded by}} \Phi^{\gamma_L^* \rightarrow V_L}$
- $\gamma_T^* \rightarrow V_T \xrightarrow{\text{encoded by}} \Phi^{\gamma_T^* \rightarrow V_T}$

$V = \rho, \phi$ via **distribution amplitudes (DAs)**: $\varphi(y) = \varphi^{\text{WW}}(y) + \varphi^{\text{gen}}(y)$

\implies **DAs** enter $\Phi^{\gamma^* \rightarrow V} = [\text{H}_{\text{LO}}] \circledast [\text{DA}]$

T_{11} and T_{00}

Assumption:

- Wandzura-Wilczek (WW) approximation \rightarrow genuine terms neglected

$$T_{11} = is \frac{2 B C}{Q^2} \int \frac{d^2 \kappa}{(\kappa^2)^2} \mathcal{F}(x, \kappa^2) \int_0^1 \frac{dy}{(y\bar{y} + \tau)} \varphi_+^{\text{WW}}(y, \mu^2) \frac{\alpha(\alpha + 2y\bar{y} + 2\tau)}{(\alpha + y\bar{y} + \tau)^2} + o(\tau^2)$$

$$T_{00} = is \frac{4 B C}{Q} \int \frac{d^2 \kappa}{(\kappa^2)^2} \mathcal{F}(x, \kappa^2) \int_0^1 dy \frac{y\bar{y}}{(y\bar{y} + \tau)} \left(\frac{\alpha}{\alpha + y\bar{y} + \tau} \right) \varphi_1^{\text{as}}(y, \mu^2)$$

where $B = \frac{\pi \alpha_s f_V e_V}{N_c}$, $C = \sqrt{4\pi \alpha_{\text{em}}}$, $\tau = m_q^2 / Q^2$, $\alpha = \kappa^2 / Q^2$.

- Generalized massive formula: $\tau = 0 \rightarrow$ no quark mass $\Rightarrow \rho$ -production
 $\tau \neq 0 \rightarrow$ with quark mass $\Rightarrow \phi$ -production

\Rightarrow Vector meson-DAs employed:

- asymptotic $\varphi_1^{\text{as}}(y) \xrightarrow{\text{fixing}} a_2(\mu^2) = 0$
- $\varphi_+^{\text{WW}}(y, \mu^2) = (2y - 1) \varphi_{1T}^{\text{WW}}(y, \mu^2) + \varphi_{AT}^{\text{WW}}(y, \mu^2)$

UGD models

$\Rightarrow \mathcal{F}(x, \kappa^2)$ has to be modeled!

Following UGD models are selected:

- **Ivanov-Nikolaev:**

soft and hard components

$\xrightarrow{\text{to probe}}$ different regions of κ

[I. P. Ivanov and N. N. Nikolaev, *Phys. Rev. D* 65 (2002)]

- **GBW:**

FT of dipole cross section

\Rightarrow evolution saturation scale
is not needed



Standard GBW model

[K.J. Golec-Biernat, M. Wüsthoff, *Phys. Rev. D* 59 (1998) 014017]

Cross section and $b(Q^2)$ -slope parametrization

► Polarized cross sections:

$$\sigma_L(\gamma^* p \rightarrow V p) = \frac{1}{16\pi b(Q^2)} \frac{|\mathbf{T}_{00}(\mathbf{s}; Q^2)|^2}{W^2}$$

$$\sigma_T(\gamma^* p \rightarrow V p) = \frac{1}{16\pi b(Q^2)} \frac{|\mathbf{T}_{11}(\mathbf{s}; Q^2)|^2}{W^2}$$

► $b(Q^2)$ -slope for light vector mesons:

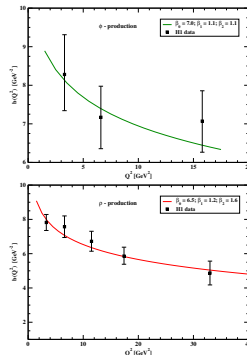
$$b(Q^2) \approx \beta_0 - \beta_1 \log \left[\frac{Q^2 + m_V^2}{m_{J/\psi}^2} \right] + \frac{\beta_2}{Q^2 + m_V^2},$$

● for ϕ -meson:

$$\beta_0 = 7.0 \text{ GeV}^{-2}, \beta_1 = 1.1 \text{ GeV}^{-2}, \\ \beta_2 = 1.1;$$

● for ρ -meson:

$$\beta_0 = 6.5 \text{ GeV}^{-2}, \beta_1 = 1.2 \text{ GeV}^{-2}, \\ \beta_2 = 1.6.$$

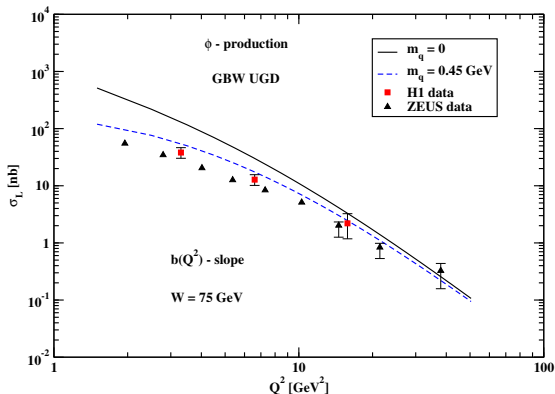


Numerical results

ϕ -production

ϕ -production - Quark mass m_q effect on σ_L

- ▶ WW approximation
- ▶ **GBW UGD** employed

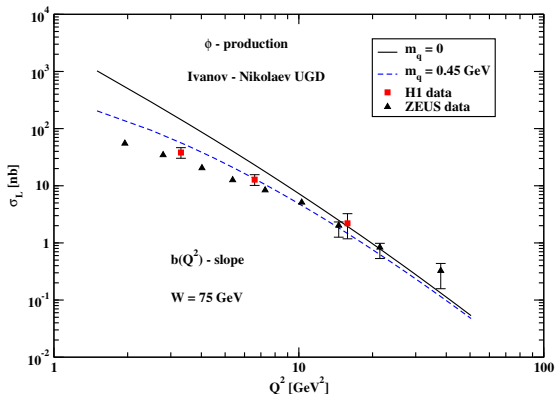


[A.D. Bolognino, W. Schäfer, A. Szczurek (in progress)]

- Crucial effect of quark mass m_q in order to catch data

ϕ -production - Quark mass m_q effect on σ_L

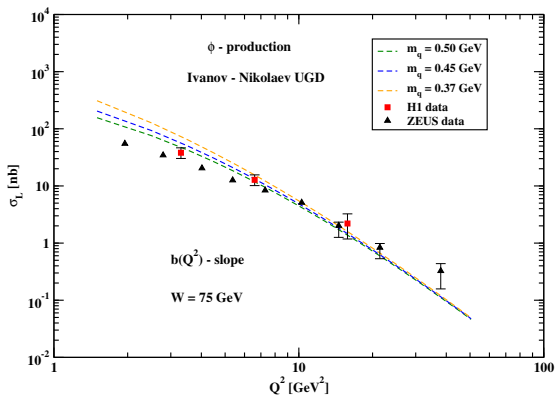
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[A.D. Bolognino, W. Schäfer, A. Szczurek (in progress)]

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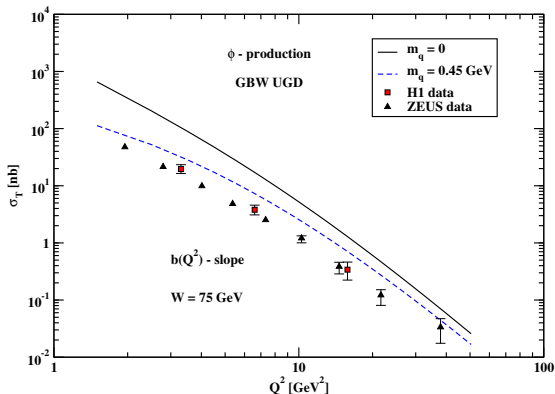
ϕ -production - Stability of σ_L on m_q values



[A.D. Bolognino, W. Schäfer, A. Szczurek (in progress)]

ϕ -production - Quark mass m_q effect on σ_T

- ▶ WW approximation
- ▶ **GBW UGD** employed

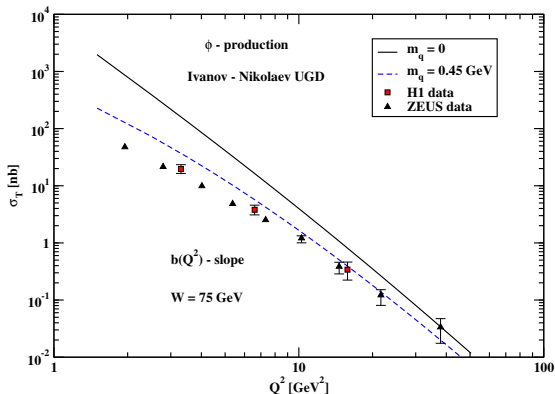


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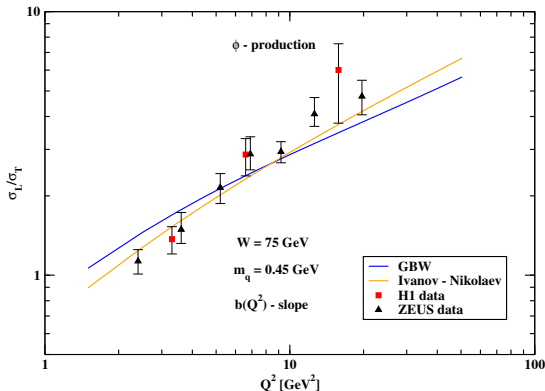


[A.D. Bolognino, W. Schäfer, A. Szczurek (in progress)]

- Crucial effect of quark mass m_q in order to catch data

ϕ -production - Cross section ratio σ_L/σ_T

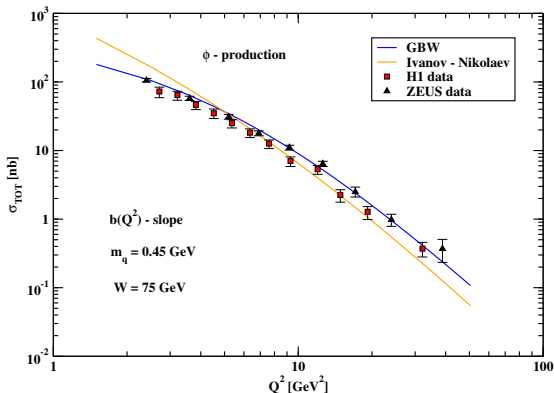
► GBW and Ivanov-Nikolaev models compared



[A.D. Bolognino, W. Schäfer, A. Szczurek (in progress)]

ϕ -production - Total cross section σ_{TOT}

- **GBW** and **Ivanov-Nikolaev** models compared
- $\sigma_{TOT} = \sigma_L + \epsilon \sigma_T$, with $\epsilon \approx 1$ in HERA kinematics



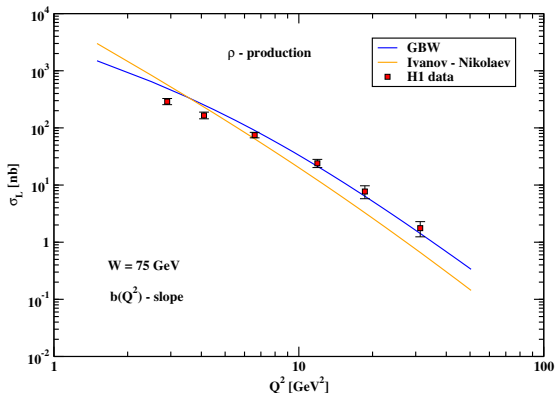
[A.D. Bolognino, W. Schäfer, A. Szczurek (in progress)]

Numerical results

ρ -production

ρ -production - Longitudinal cross section σ_L

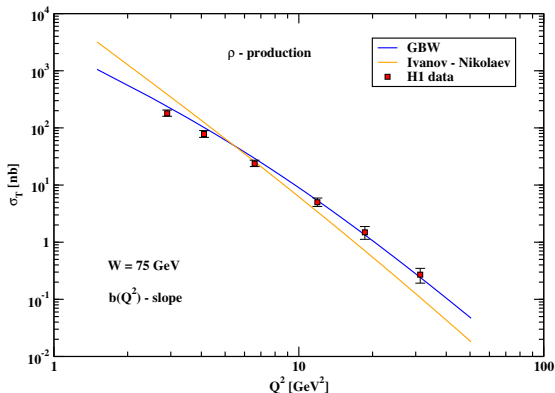
- WW approximation
- GBW and Ivanov-Nikolaev models compared



[A.D. Bolognino, F.G. Celiberto, D.Yu. Ivanov, A. Papa, W. Schäfer, A. Szczurek (in progress)]

ρ -production - Transverse cross section σ_T

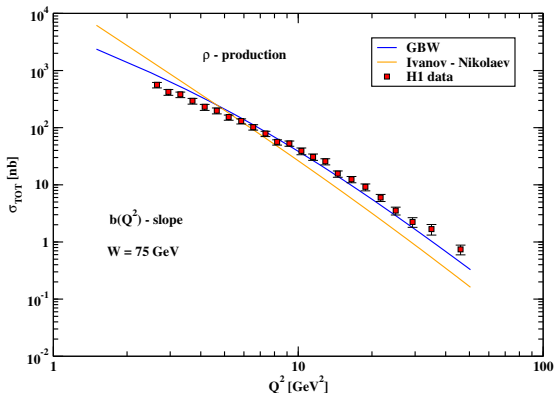
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[A.D. Bolognino, F.G. Celiberto, D.Yu. Ivanov, A. Papa, W. Schäfer, A. Szczurek (in progress)]

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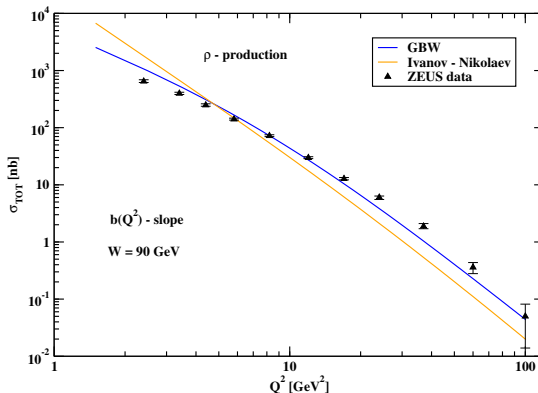
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[A.D. Bolognino, F.G. Celiberto, D.Yu. Ivanov, A. Papa, W. Schäfer, A. Szczurek (in progress)]

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[A.D. Bolognino, F.G. Celiberto, D.Yu. Ivanov, A. Papa, W. Schäfer, A. Szczurek (in progress)]

Conclusions...

Exclusive leptonproduction of **polarized light vector mesons** to describe **cross sections**

- ▶ **Exclusive** final state + **small-x** limit \Rightarrow **κ_T -factorization** allowed
- ▶ t -dependence of σ parametrized via $b(Q^2)$ -slope:
 - \Rightarrow Improved predictions for **ϕ -production** via quark mass m_q ✓
 - \Rightarrow Standard GBW model able to catch HERA data in **ρ -production** ✓

...Outlook

For ρ -production:

- ▶ NLO impact factor in ρ -meson leptonproduction $\Rightarrow \Phi\gamma^*\rightarrow\rho = [H_{\text{NLO}}] \circledast [DA]$
- ▶ Test different UGD models $\xrightarrow{\text{in order to}}$ calculate cross section

For ϕ -production:

- ▶ Cross section encoded by non-forward helicity amplitude and spin-flip
- ▶ Quark mass m_q in genuine contribution

...And in general

- ▶ UGD extraction from different channels:
 - ◇ small- x gluon TMD \rightarrow BFKL evolution + TMD input
[A.D. Bolognino, F.G. Celiberto, ... (in progress)]
- ▶ Consider other processes as testfield for UGD:
 - ◇ Heavy quark and heavy-meson production

Thanks for your
attention!!

BACKUP slides

UGD models

Ivanov and Nikolaev' (IN) UGD: a soft-hard model

$$\mathcal{F}(x, \kappa^2) = \mathcal{F}_{\text{soft}}^{(B)}(x, \kappa^2) \frac{\kappa_s^2}{\kappa^2 + \kappa_s^2} + \mathcal{F}_{\text{hard}}(x, \kappa^2) \frac{\kappa^2}{\kappa^2 + \kappa_h^2},$$

The soft term:

$$\diamond \mathcal{F}_{\text{soft}}^{(B)}(x, \kappa^2) = a_{\text{soft}} C_F N_c \frac{\alpha_s(\kappa^2)}{\pi} \left(\frac{\kappa^2}{\kappa^2 + \mu_{\text{soft}}^2} \right)^2 V_N(\kappa)$$

- $\mu_{\text{soft}}^2 \rightarrow$ soft parameter
- $a_{\text{soft}} \rightarrow$ weight of soft term compared to the hard one

The hard term:

$$\diamond \mathcal{F}_{\text{hard}}(x, \kappa^2) = \mathcal{F}_{\text{pt}}^{(B)}(\kappa^2) \frac{\mathcal{F}_{\text{pt}}(x, Q_c^2)}{\mathcal{F}_{\text{pt}}^{(B)}(Q_c^2)} \theta(Q_c^2 - \kappa^2) + \mathcal{F}_{\text{pt}}(x, \kappa^2) \theta(\kappa^2 - Q_c^2)$$

- $\mathcal{F}_{\text{pt}}(x, \kappa^2) = \frac{\partial x g(x, \kappa^2)}{\partial \ln \kappa^2}$
- $\mathcal{F}_{\text{pt}}^{(B)}(x, \kappa^2) = C_F N_c \frac{\alpha_s(\kappa^2)}{\pi} \left(\frac{\kappa^2}{\kappa^2 + \mu_{\text{pt}}^2} \right)^2 V_N(\kappa)$

The coupling constant:

- $\alpha_s \leq 0.82$ (frozen)

[I. P. Ivanov and N. N. Nikolaev, *Phys. Rev. D* 65 (2002)]

UGD models

Golec-Biernat–Wüsthoff' (GBW) UGD

$$\mathcal{F}(x, \kappa^2) = \kappa^4 \sigma_0 \frac{R_0^2(x)}{2\pi} e^{-\kappa^2 R_0^2(x)}$$

- ◇ derives from the effective dipole cross section $\hat{\sigma}(x, r)$ for the scattering of a $q\bar{q}$ pair off a nucleon $\xrightarrow{\text{through}}$ a reverse Fourier transform of

$$\sigma_0 \left\{ 1 - \exp\left(-\frac{r^2}{4R_0^2(x)}\right) \right\} = \int \frac{d^2\kappa}{\kappa^4} \mathcal{F}(x, \kappa^2) (1 - \exp(i\vec{\kappa} \cdot \vec{r})) (1 - \exp(-i\vec{\kappa} \cdot \vec{r}))$$

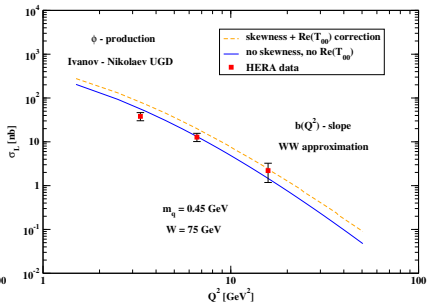
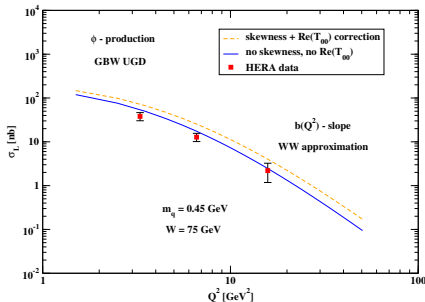
- ◇ $R_0^2(x) = \frac{1}{\text{GeV}^2} \left(\frac{x}{x_0}\right)^{\lambda_p}$
- ◇ The normalization σ_0 and the parameters x_0 and $\lambda_p > 0$ of $R_0^2(x)$ have been determined by a global fit to $F_2(x)$:

$$\sigma_0 = 23.03 \text{ mb}, \quad \lambda_p = 0.288, \quad x_0 = 3.04 \cdot 10^{-4}.$$

[K.J. Golec-Biernat, M. Wüsthoff, *Phys. Rev. D* 59 (1998) 014017]

Skewness effects - ϕ -production

Skewness effects on σ_L



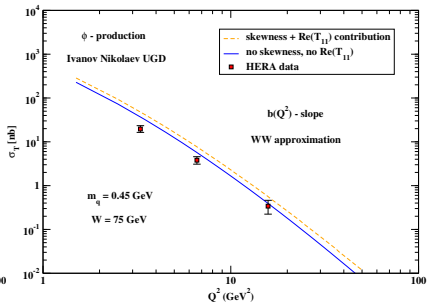
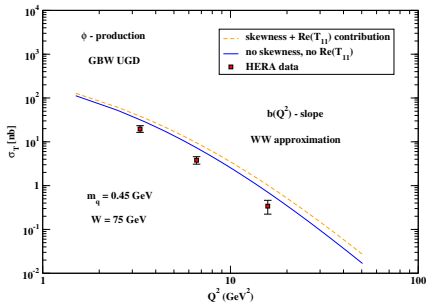
$$\sigma_L = Rg^2(1 + \rho^2) \tilde{\sigma}_L$$

where

- $Rg = \frac{2^{2\lambda+3}}{\sqrt{\pi}} \frac{\Gamma(\lambda+5/2)}{\Gamma(\lambda+4)}$ with $\lambda = \frac{\delta \log(1/s \Im T_{00})}{\delta \log(1/x)}$
- $\rho = \tan\left(\frac{\pi\lambda}{2}\right) = \frac{\Re T_{00}}{\Im T_{00}}$

Skewness effects - ϕ -production

Skewness effects on σ_T



$$\sigma_T = Rg^2(1 + \rho^2) \tilde{\sigma}_T$$

where

- $Rg = \frac{2^{2\lambda+3}}{\sqrt{\pi}} \frac{\Gamma(\lambda+5/2)}{\Gamma(\lambda+4)}$ with $\lambda = \frac{\delta \log(1/s \Im T_{11})}{\delta \log(1/x)}$
- $\rho = \tan\left(\frac{\pi\lambda}{2}\right) = \frac{\Re T_{11}}{\Im T_{11}}$