H γγ BEYOND THE STANDARD MODEL

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G.CACCIAPAGLIA, A.DEANDREA & J.L.P, ARXIV: 0901.0927 JHEP: 0906:054, 2009

OUTLINE

- Introduction of SM Higgs Physics
- Observables and New Physics
- Model-independent parametrization
- Use for data analysis and models survey

WHERE IS THE HIGGS BOSON?

- LEP and Tevatron results
 Higgs at low mass ?
- $115 \text{ GeV} < M_H < 150 \text{ GeV}$
- SM Higgs Radiative instability of Higgs mass
- New Physics ????
 Link with EWSB



J. Ellis, J.R. Espinosa , G.F. Giudice, A. Hoecker, A. Riotto, e-Print: arXiv:0906.0954 [hep-ph]

LIGHT HIGGS AT LHC



- Production → Gluons Fusion
- Detection Channel Decay into two photons

LIGHT HIGGS AT ILC



- Detection Channel ->> Decay into photons & gluons

LOOP INDUCED PROCESSES AND NEW PHYSICS



- Couplings sensitive to NP running into the loops
- Links between New Physics and EWSB (even beyond direct detection threshold)



SM → Masses proportional to Higgs VEV: v_{SM}
 → No decoupling

$$A_F(0) = \frac{4}{3}, \quad A_W(0) = -7, \quad A_S(0) = \frac{1}{3}, \quad A_{F,W,S}(\infty) \longrightarrow 0$$



• New Physics \rightarrow Charged and Colored Particles \rightarrow Couplings not necessarily proportional to Higgs VEV $A_{NP} = \frac{v_{SM}}{m_{NP}} \frac{\partial m_{NP}}{\partial v} A_{F,W,S}$.



• New Physics \rightarrow Charged and Colored Particles \rightarrow Couplings not necessarily proportional to Higgs VEV $A_{NP} = \underbrace{v_{SM}}_{m_{NP}} \frac{\partial m_{NP}}{\partial v} A_{F,W,S}$. Coupling normalized by the SM one



New Physics → Charged and Colored Particles
 → Couplings not necessarily proportional to Higgs VEV
 A_{NP} = ^v_{SM} ∂m_{NP} ∂m_{NP} A_{F,W,S}.
 Only depends on spin and mass of NP

A MODEL INDEPENDENT PARAMETRIZATION

$$\Gamma_{\gamma\gamma} = \frac{G_F \alpha^2 m_H^3}{128\sqrt{2}\pi^3} \left| A_W(\tau_W) + 3\left(\frac{2}{3}\right)^2 A_t(\tau_t) \left[1 + \kappa_{\gamma\gamma}\right] + \dots \right|^2$$

$$\Gamma_{gg} = \frac{G_F \alpha_s^2 m_H^3}{16\sqrt{2}\pi^3} \left| \frac{1}{2} A_t(\tau_t) \left[1 + \kappa_{gg}\right] + \dots \right|^2 ,$$

where
$$\kappa_{\gamma\gamma} = \sum_{NP} \frac{3}{4} N_{c,NP} Q_{NP}^2 \frac{v_{SM}}{m_{NP}} \frac{\partial m_{NP}}{\partial v} \frac{A_{F,W,S}(m_{NP})}{A_t}$$
,
 $\kappa_{gg} = \sum_{NP} 2C(r_{NP}) \frac{v_{SM}}{m_{NP}} \frac{\partial m_{NP}}{\partial v} \frac{A_{F,W,S}(m_{NP})}{A_t}$

- Normalization with the top contribution
- Introduction of 2 new parameters
- Only loop effects & no tree level modifications

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2

where
$$\kappa_{\gamma\gamma} = \sum_{NP} \frac{3}{4} N_{c,NP} Q_{NP}^2 \frac{v_{SM}}{m_{NP}} \frac{\partial m_{NP}}{\partial v} \begin{bmatrix} A_{F,W,S}(m_{NP}) \\ A_t \end{bmatrix}$$

 $\kappa_{gg} = \sum_{NP} 2C(r_{NP}) \frac{v_{SM}}{m_{NP}} \frac{\partial m_{NP}}{\partial v} \begin{bmatrix} A_{F,W,S}(m_{NP}) \\ A_t \end{bmatrix}$

• If $2 M_{NP} >> M_{H}$, only spin dependent ratios

$$\frac{A_{NP}}{A_t} = \begin{cases} 1 & \text{for fermions} \\ -\frac{21}{4} & \text{for vectors} \\ \frac{1}{4} & \text{for scalars} \end{cases}$$

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 $\kappa_{gg} = \sum_{NP} 2C(r_{NP}) \underbrace{\frac{v_{SM}}{m_{NP}} \frac{\partial m_{NP}}{\partial v}}_{M_{NP}} \frac{A_{F,W,S}(m_{NP})}{A_t}$

• If there is decoupling

$$\frac{v_{SM}}{m_{NP}}\frac{\partial m_{NP}}{\partial v} \sim \frac{v_{SM}^2}{m_{NP}^2}$$

CORRECTIONS TO STANDARD MODEL

$$\Gamma_{\gamma\gamma} = \frac{G_F \alpha^2 m_H^3}{128\sqrt{2}\pi^3} \left| A_W(\tau_W) + 3\left(\frac{2}{3}\right)^2 A_t(\tau_t) \left[1 + \kappa_{\gamma\gamma}\right] + \dots \right|^2$$

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 New Physics can modified Standard Model W and top couplings:

$$\begin{aligned} \kappa_{\gamma\gamma}(top) &= \kappa_{gg}(top) = \frac{v_{SM}}{m_t} \frac{\partial m_t}{\partial v} - 1 \,. \\ \kappa_{\gamma\gamma}(W) &= \frac{3}{4} \left(\frac{v_{SM}}{m_W} \frac{\partial m_W}{\partial v} - 1 \right) \, \frac{A_W(\tau_W)}{A_F(\tau_{top})} \,, \\ \kappa_{gg}(W) &= 0 \,. \end{aligned}$$

PRINCIPAL INTERESTS OF THIS PARAMETRIZATION

- Top is charged and colored
- New Physics probably linked to the top physics
- For a top partner: $\kappa_{\gamma\gamma} = \kappa_{gg}$
- Same sign for one new particle
- Positive kappas enhance gluon and decrease photon widths.
- Easy to extend to non minimal Higgs sector (Multiple Higgs, mixing with other scalars,...)

NEW PARAMETRIZATION AND LHC/ILC OBSERVABLES

$$\bar{\sigma}(H) = \left(\frac{\sigma_{gg}^{NP} + \sigma_{VBF}^{SM} + \sigma_{VH,\bar{t}tH}^{SM}}{\sigma_{gg}^{SM} + \sigma_{VBF}^{SM} + \sigma_{VH,\bar{t}tH}^{SM}}\right) \simeq \left(\frac{(1 + \kappa_{gg})^2 \sigma_{gg}^{SM} + \sigma_{VBF}^{SM} + \sigma_{VH,\bar{t}tH}^{SM}}{\sigma_{gg}^{SM} + \sigma_{VBF}^{SM} + \sigma_{VH,\bar{t}tH}^{SM}}\right)$$
Inclusive cross section normalized to SM

- Corrections to Hbb and HWW couplings neglected
- If correction to Hbb too important
 Parametrization non valid (SUSY in general)
- But an extended parametrization is possible using new parameters

OTHER INTERESTING OBSERVABLES

$$\overline{BR}(H \to \gamma\gamma) = \frac{\Gamma_{\gamma\gamma}^{NP}}{\Gamma_{\gamma\gamma}^{SM}} \frac{\Gamma_{\text{tot}}^{SM}}{\Gamma_{gg}^{NP} + \Gamma_{\gamma\gamma}^{NP} + \Gamma_{\text{others}}^{SM}}$$

$$\simeq \left(1 + \frac{\kappa_{\gamma\gamma}}{\frac{9}{16}A_W(\tau_W) + 1}\right)^2 \frac{\Gamma_{\text{tot}}^{SM}}{(1 + \kappa_{gg})^2 \Gamma_{gg}^{SM} + (\Gamma_{\text{tot}}^{SM} - \Gamma_{gg}^{SM})}$$

$$\overline{Photon Branching Ratio normalized to SM}$$

$$\overline{BR}(H \to gg) = \frac{\Gamma_{gg}^{NP}}{\Gamma_{gg}^{SM}} \frac{\Gamma_{\text{tot}}^{SM}}{\Gamma_{gg}^{NP} + \Gamma_{\gamma\gamma}^{NP} + \Gamma_{\text{others}}^{SM}}$$

$$\simeq \frac{(1 + \kappa_{gg})^2 \Gamma_{\text{tot}}^{SM}}{(1 + \kappa_{gg})^2 \Gamma_{gg}^{SM} + (\Gamma_{\text{tot}}^{SM} - \Gamma_{gg}^{SM})}.$$

Gluon Branching Ration normalized to SM

FROM SPECTRA OF GIVEN MODELS TO NP PARAMETERS

- Easy to compute these contributions
- For instance: 4^{th} chiral generation $190 \text{ GeV} < m_Q < 2 \text{ TeV}$ $100 \text{ GeV} < m_L < 2 \text{ TeV}$ arXiv:0706.3718 [hep-ph]

$$\kappa_{gg} = 2$$

$$\kappa_{\gamma\gamma} = \frac{3}{4} \left[3 \left(\frac{2}{3} \right)^2 + 3 \left(-\frac{1}{3} \right)^2 + 1 \right] = 2$$

----> Equivalent to 2 extra tops!

SURVEY OF MODELS

٠	Fourth Generation
.	SuperSymmetry in MSSM Golden Region
	Simplest Little Higgs (W' @ 2TeV)
*	Littlest Higgs (<i>T</i> - <i>Parity</i> : <i>f</i> =500 GeV, without <i>T</i> -parity: <i>f</i> =5 TeV)
	Color Octet Model
	Lee Wick Standard Model (LW Higgs mass @ 1 TeV)
\otimes	Universal Extra Dimension Model (M _{KK} @ 500 GeV)
*	Model of Gauge Higgs Unification in Flat Space (W ⁽¹⁾ @ 2 TeV)
•	Minimal Composite Higgs (GHU in Warped Space) (W ⁽¹⁾ @ 2.4 TeV)
▼	Brane Higgs Model with Flavour in Flat Space (W ⁽¹⁾ @ 2 TeV)
٨	Brane Higgs Model with Flavour in Warped Space (W ⁽¹⁾ @ 2 TeV)

MH=120 GEV @ ILC



A-BR(H $\rightarrow\gamma\gamma$) B-BR(H \rightarrow GG)

•	4 th	*	Littlest Higgs	•	Warped GHU Space	\star	Flat GHU
•	SUSY		Color Octet		Flat BH with Flavour	\otimes	UED Model
	SLH		Lee Wick SM		Warped BH with Flavour		

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	SLH		Lee Wick SM	•	Warped BH with Flavour		

MH=120 GEV @ LHC



MH=150 GEV @ LHC



A-INCLUSIVE $H \rightarrow \gamma \gamma$ B- VBF $\rightarrow H \rightarrow VV$

•	4 th	*	Littlest Higgs	•	Warped GHU Space	\star	Flat GHU
•	SUSY		Color Octet		Flat BH with Flavour	\otimes	UED Model
	SLH		Lee Wick SM		Warped BH with Flavour		

CONCLUSION

- ✓ **Usefull tool** to study EWSB complementary to direct detection
- ✓ Sizeable effects and possible detection of new particles even beyond direct detection threshold → Good for LC
- ✓ **Model independent parametrization** with few parameters
- Models pointing in well-defined directions in parameter space
 Good discrimination power !
- Now, need to be implemented in MC for more detailed analysis.
- ♦ Possible analysis with different observables ($BR(H \rightarrow VV),...$)
- Possible extension for non minimal Higgs sector

THE END

EASY IMPLEMENTATION OF THIS PARAMETRIZATION

- If m_H<<m_{NP} ==> NP only modifies Higgs loops ==> Same kinematics
- For instance, in Madgraph with HEFT model: Higgs loops described with effective vertices
 => Just tunable coefficients

c Higgs effective couplings (couplings of Higgs's directly to gluons)

c Coupling to gluons c Higgs coupling:

```
tau = hmass**2/(4d0*tmass**2)
series_t = 1d0 + tau*7d0/30d0 +
tau**2*2d0/21d0 + tau**3*26d0/525d0
series_p = 1d0 + tau/3d0 +
tau**2*8d0/45d0 + tau**3*4d0/35d0
```

c kglu is the effective coupling modification to H-> gluon gluon c such that (1+kglu) multiplies the top triangle contribution

kglu = 0d0

ga(1) = dcmplx(Zero, Zero) ga(2) = dcmplx(Zero, Zero)

MH=120 GEV @ LHC



A-INCLUSIVE $H \rightarrow \gamma \gamma$ B- VBF $\rightarrow H \rightarrow \gamma \gamma$

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MH=120 GEV @ LHC





GENERALIZATION

- Previous parametrization for a SM-like Higgs sector
- Extension for a non minimal scalar sector
 - Multiple higgses:

$$\phi_i = \frac{1}{\sqrt{2}} (v_i + c_i h + \dots)$$

$$\frac{v_{SM}}{m} \frac{\partial m(v)}{\partial v} \rightarrow \frac{v_{SM}}{m} \sum_i \frac{\partial m}{\partial v_i} c_i$$

• Mixing with no vev scalars:

$$S_{j} = s_{j}h + \dots$$

$$\frac{v_{SM}}{m} \frac{\partial m(v)}{\partial v} \rightarrow \frac{v_{SM}}{m} \left(\sum_{i} \frac{\partial m}{\partial v_{i}} c_{i} + \sum_{j} g_{Sj} s_{j} \right)$$

SUSY: MSSM IN THE GOLDEN REGION (ARXIV: HEP-PH 0702038)

- Motivated by naturalness, minimal fine tuning, precision tests, ...
- Large tan β compensated by small mixing angle α in Higgs sector:
 - Sparticles: large mixing in the stop sector
 - Charginos: mostly higgsinos
 - Heavy Higgses masses above 1TeV

$$\frac{g_{W^+W^-h}}{g_{SM}} = \sin(\beta - \alpha), \qquad \frac{g_{\bar{t}th}}{g_{SM}} = \frac{\cos\alpha}{\sin\beta}, \qquad \frac{g_{\bar{b}bh}}{g_{SM}} = -\frac{\sin\alpha}{\cos\beta}$$

MODELS WITH FLAVOR IN EXTRA DIMENSIONS



- O(1) Yukawa on the brane and O(1) mass in the bulk

FLAVOR IN EXTRA-DIM GAUGE BOSONS

Negligible contribution of W gauge boson

FERMIONS

• In general, different bulk masses of right handed and left handed fermions $M_L \neq M_R$

$$\sum_{n} \frac{\beta}{m_n} \frac{\partial m_n}{\partial \beta} \simeq -\pi^2 \beta^2$$

 β is linked to the VEV and the Yukawa of the field

All fermions contribute to the kappas

$$\kappa_{\gamma\gamma} = \kappa_{gg} \simeq 6(-\pi^2 \beta^2) - \frac{\pi^2 \alpha^2}{6} \sim -0.45 \left(\frac{2\text{TeV}}{m_{KK}}\right)^2,$$