

WORKSHOP ON THE MICROBUNCHING INSTABILITY III



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Coherent Optical Radiation: Operational Experience @ FLASH

M. Castellano, <u>E. Chiadroni</u> (INFN – LNF)

A. Cianchi (University and INFN – Roma "Tor Vergata")

K. Honkavaara, G. Kube (DESY – HH)

Non-Intercepting Devices

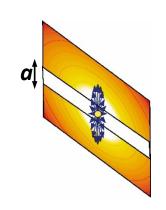
Emittance measurements are often performed imaging a beam on a metallic screen

BUT

- Traditional diagnostic based on OTR cannot be used for high power beam
- All other intercepting devices are easily damaged or destroyed from these type of beams
- It is fundamental to develop non intercepting alternatives for emittance measurements
 - Even more desirable after recent observations of COTR in linacdriven FELs

Diffraction Radiation

DR is produced by the interaction between the EM fields of the traveling charge and the conducting screen



> The radiation intensity is

$$I \propto \frac{-\frac{\pi \alpha}{\gamma \lambda}}{I}$$

DR impact parameter is $\frac{\gamma\lambda}{2\pi}$ if a

$$\frac{\gamma\lambda}{2\pi}$$
 \rightarrow if a

$$>> \frac{\gamma \lambda}{2\pi}$$
 No radiation
$$\cong \frac{\gamma \lambda}{2\pi}$$
 DR
$$<< \frac{\gamma \lambda}{2\pi}$$
 TR

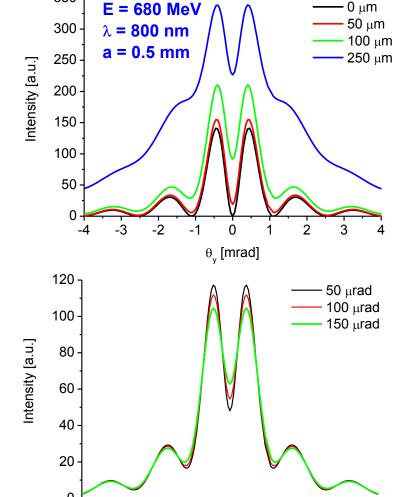
Excellent candidate to measure beam parameters parasitically

ODR as Transverse Beam Diagnostics

350

The visibility of the interference fringes can be used to determine the transverse size of a bunch of electrons crossing the slit:

The beam angular divergence too gives rise to a reduced fringes visibility, opening the way to a possible single shot emittance measurement.



-3

-2

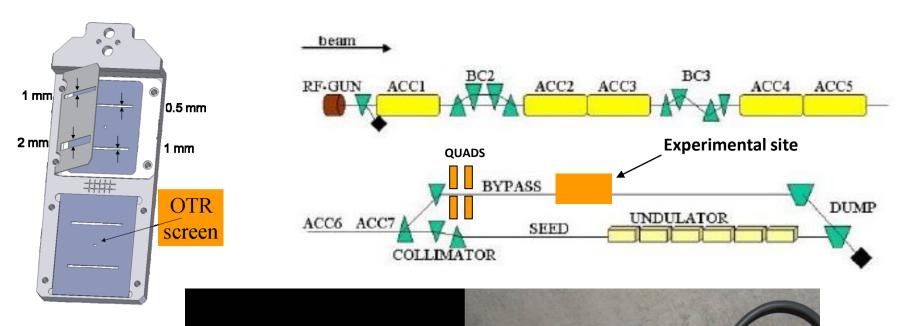
 $\theta_{..}$ [mrad]

2

3

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ODR Experiment @ FLASH



Double convex coating with f = DR angular dist

Achromatic doub with f = 250 mm for beam imagin tsu CCD Model 8-LGLAG2)

Hamamatsu Camera

- Back illuminated, more than 80% of quantum efficiency at both 550 nm and 800 nm
- Peltier cooled @ -55 °C
- No dark noise signal
- Integration time up to hours
- > 30 bunches of 1 nC @ 5 Hz integrated for 1 second => 150 nC!

Optical <u>Diffraction Radiation</u> Interferometry (ODRI)

1 mm

2 mm

1111

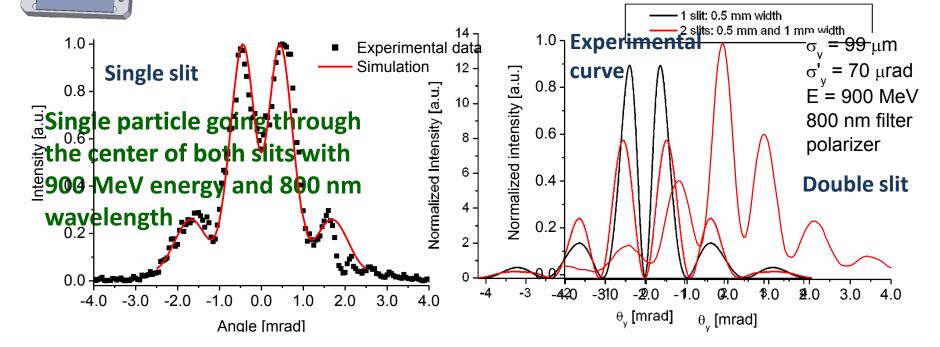
0.5 mm

1 mm

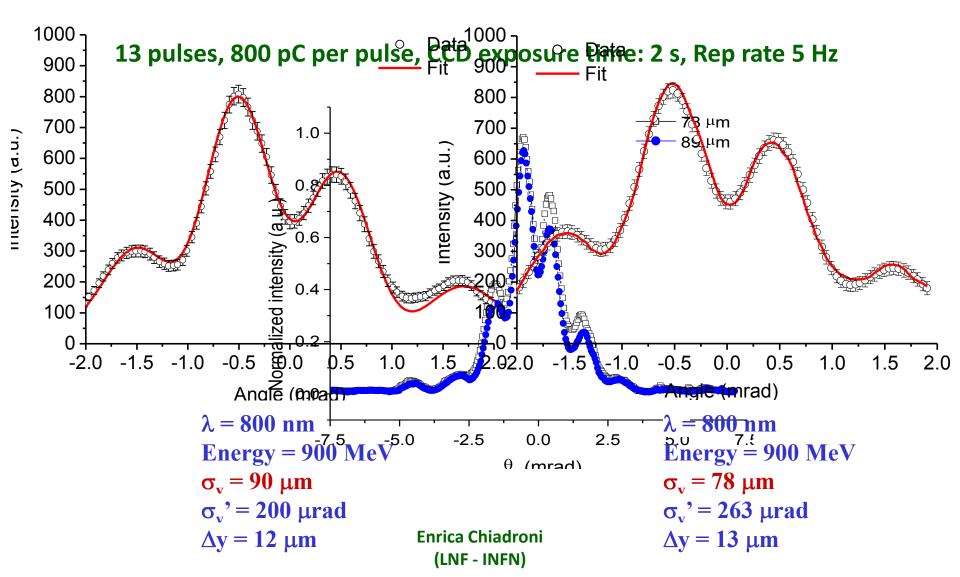
OTR

In case of 800 nm wavelength and 1 GeV beam energy the 1 mm cut is not large enough to prevent the production of ODR in the forward direction, reflected by the screen and interfering with the backward ODR produced by the screen itself.

The two interfering amplitudes have different intensity and angular distribution.



ODRI Angular Distribution for Two Different Beam Sizes



Coherent Radiation

The total radiation intensity emitted by a bunch of electrons is given by

$$I_{tot}(\lambda = I_{sp} N + N(N -)F_{||}(\lambda F_{\perp}'\lambda 9)$$

in which I_{sp} is the intensity emitted by a single particle and $F_{//}(\lambda)$ and $F_{\tau}(\lambda, 9)$ the bunch longitudinal and transverse form factors, respectively

$$F_{\parallel}(\lambda) = \int_{-\infty}^{\infty} S(z)e^{i\frac{2\pi}{\lambda}z}dz \bigg|^{2}, \quad F_{\perp}'\lambda \vartheta = \int_{-\infty}^{\infty} T(\rho)e^{i\frac{2\pi}{\lambda}\rho\sin(\vartheta)}d\rho \bigg|^{2}$$

with S(z) the longitudinal density and $T(\rho)$ the transverse distribution of the bunch.

If part of the bunch emits coherently (no compression mode), then

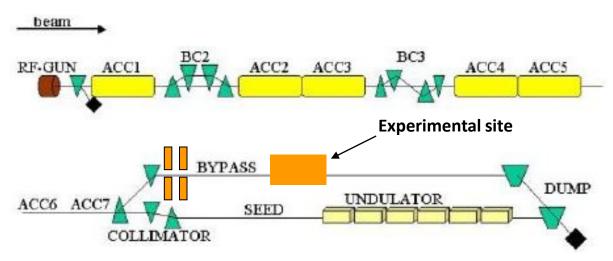
$$I_{coh} = N_{coh}^2 F_{||}(\lambda \ F_{\perp}'\lambda \ \mathcal{G} \ I_{sp} \quad \Longrightarrow \quad I_{tot} \cong I_{sp} \ \mathbb{V} + N_{coh}^2 F_{||}(\lambda \ F_{\perp}'\lambda \ \mathcal{G})$$

We assume the # of coherent particles remains the same at 800 nm and 550 nm.

We expect a different behavior at 800 nm and 550 nm w.r.t. the OTR incoherent emission.

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FLASH Linac Settings

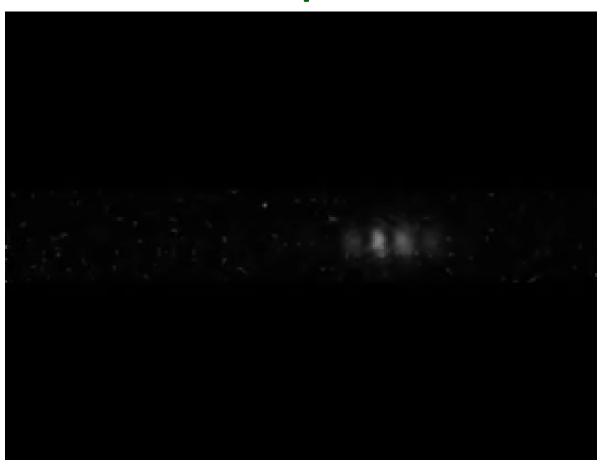


- > No compression mode => on crest in both bunch compressors
- No coherent emission detected by any diagnostics on the straight line BUT
- Emphasized by the fact that the beam travels through a dispersion section, i.e. the by-pass
- > Experimental evidences of COTR/CODRI emission due to u-bunching instabilities at ODRI station strongly dependent on the trajectory
 - 0.1% energy variation as knob to turn coherent emission on or off

- Fluctuation shot by shot more than 50% of intensity
 - ➤ fluctuation of longitudinal modulation or of the number of particles which emits coherently
- Charge fluctuation of 2%

Coherent ODRI (CODRI) Angular Distribution

20 shots: 1 pulse 0.2 s

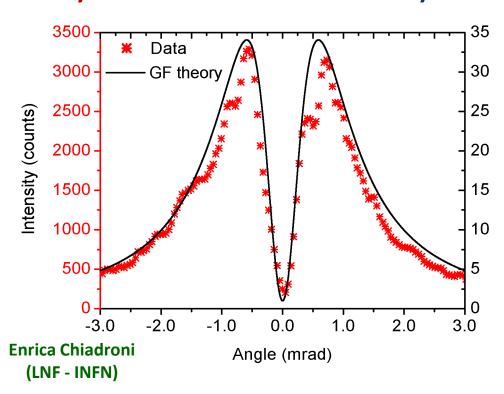


- This phenomenon has been observed in the no-compression mode
- Fluctuation shot by shot more than 50% of intensity
 - ➢ fluctuation of longitudinal modulation or of the number of particles which emits coherently
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A great enhancement of total intensity that could not be described by the

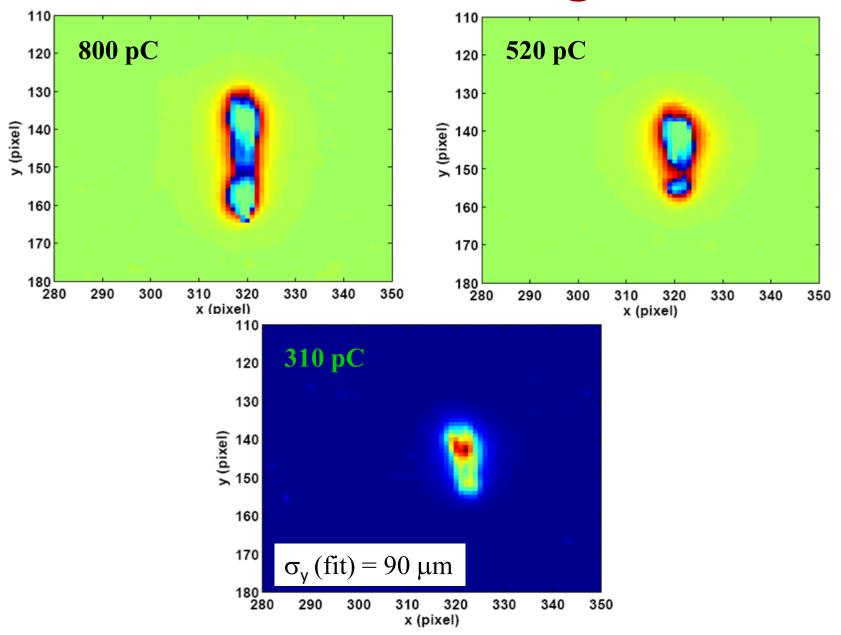
incoherent emission theory

200 ms exposure time 1 micropulse 800 pC



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- Charge fluctuation of 2%
- ➤ A great enhancement of total intensity that could not be described by the incoherent emission theory
- CCD often saturated not allowing always the measurement of beam size
 - ➤ Visible structures in the OTR beam image: we cannot evaluate because of the saturation and because of lack of resolution

OTR beam images



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- Big differences between 550 nm and 800 nm
 - > unlike the OTR incoherent emission which has white spectrum

Coherent OTR vs Incoherent OTR

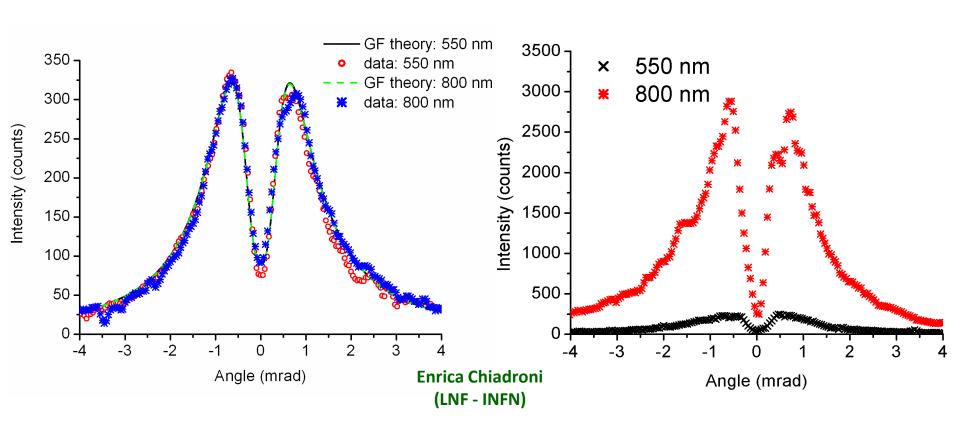
200 ms exposure time, 1 micro-pulse, 800 pC

Incoherent emission

→ OTR angular distribution does not depend on the wavelength

Coherent emission

→ OTR angular distribution shows dependence on the wavelength



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- ➤ A great enhancement of total intensity that could not be described by the incoherent emission theory
- CCD often saturated not allowing always the measurement of beam size
 - ➤ Visible structures in the OTR beam image: we cannot evaluate because of the saturation and because of lack of resolution
- ➢ Big differences between 550 nm and 800 nm
 - > unlike the OTR incoherent emission which has white spectrum
- Angular distribution with single pulse even down to 0.3 nC (while more than 100 nC, integrated, in standard operation)!!!
- No relevant change with charge

Evaluation of μ -bunch length and # of coherent particles

With the assumption of a longitudinal Gaussian-distributed beam and from the measured intensity at different wavelengths

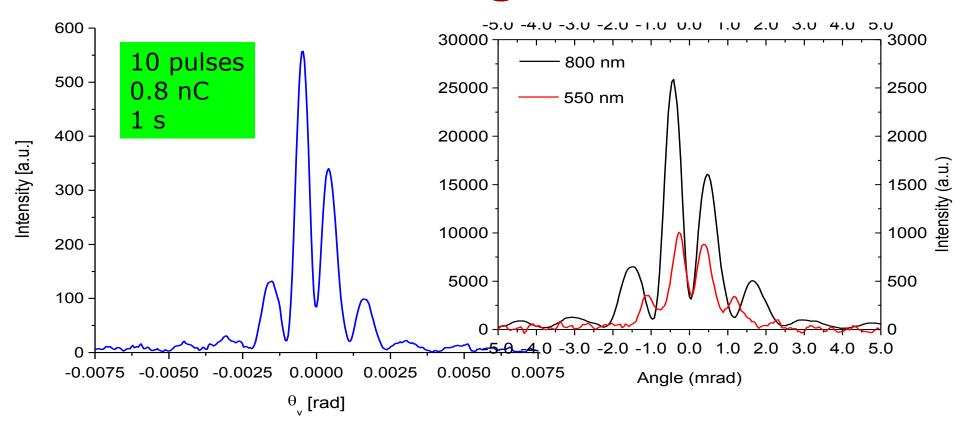
$$\frac{I_{tot}(\lambda_{-}) - I_{incoh}(\lambda_{-})}{I_{tot}(\lambda_{-}) - I_{incoh}(\lambda_{-})} \cdot \frac{I_{incoh}(\lambda_{-})}{I_{incoh}(\lambda_{-})} = \frac{f(\lambda_{-})}{f(\lambda_{-})}$$
thus
$$\sigma_{-} = \sqrt{\frac{\lambda_{-} \lambda_{-}^{2} \ln\left(\frac{I_{tot}(\lambda_{-})}{I_{tot}(\lambda_{-})}\right)}{(\lambda_{-} \lambda_{-}^{2})}}$$

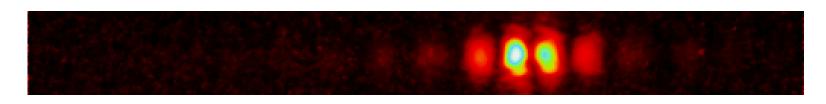
Once the form factor is evaluated, the number of particles which contributes to coherent emission can be determined if the incoherent emission is known

$$N_{coh} = \sqrt{N_e \frac{I_{coh}}{I_{incoh}} \frac{1}{f(\lambda)}}$$

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Coherent ODRI Angular Distribution





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Coherent OTR Angular Distribution

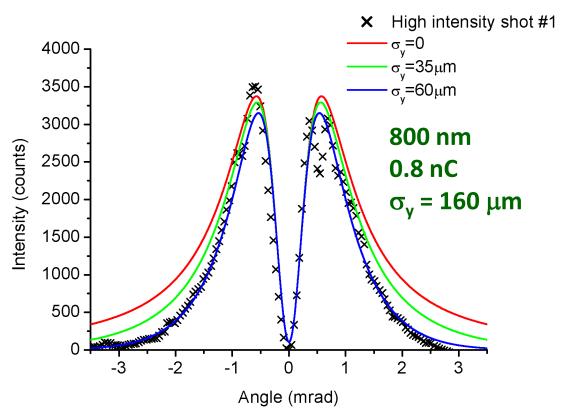
Reminding that the coherent term of the spectral angular distribution depends on the transverse form factor as

$$I_{coh} = N_{coh}^2 F_{||}(\lambda F_{\perp}'\lambda \mathcal{G} I_{sp}) \implies I_{tot} \cong I_{sp} V + N_{coh}^2 F_{||}(\lambda F_{\perp}'\lambda \mathcal{G})$$

where, assuming a Gaussian distributed beam

$$F_{\perp} (\lambda \theta) = e^{-\left(\frac{2\pi}{\lambda}\sigma_{y}\sin\theta\right)}$$

and the effect of the $\sin \vartheta$ dependence is evident at larger angles.



Conclusions

- > Evidences of coherence effects in the optical wavelength range have been observed
- ➤ Both COTR and CODRI angular distributions have been detected resulting in a promising tool to measure the transverse beam size
- The phenomenon has been observed in the no compression mode
 - > it is compression independent and no effects have been seen on any screen of the straight line
 - > the coherent optical emission we observed can be attributed to a microbunching in the visible-NIR range ($\approx 1~\mu m$ modulation) due to longitudinal density fluctuations
- \triangleright A preliminary analysis allowed us to quantify both the # of coherent particles (≈ 10⁻⁴ N_e) and the longitudinal part of the bunch which contributes to the Coherent Optical Emission