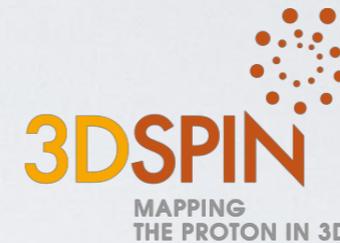


# TMDs at JLab : present and future

Pavia, 19-20 December 2018



## TRANSVERSITY and TENSOR CHARGE



Marco Radici  
INFN - Pavia

# why transversity (PDF / TMD) ?

1<sup>st</sup> Mellin moment of transversity  $\Rightarrow$  tensor “charge”

$$\delta q \equiv g_T^q = \int_0^1 dx [h_1^q(x, Q^2) - h_1^{\bar{q}}(x, Q^2)]$$

tensor charge connected to tensor operator

$$\begin{aligned} \langle p, S_p | \bar{q} \sigma^{\mu\nu} q | p, S_p \rangle &= (P^\mu S_p^\nu - P^\nu S_p^\mu) g_T^q(Q^2) \\ &= (P^\mu S_p^\nu - P^\nu S_p^\mu) \int dx h_1^{q-\bar{q}}(x, Q^2) \end{aligned}$$

tensor operator not accessible in tree-level Standard Model  
low-energy footprint of new physics at higher scales ?

# why di-hadron mechanism ?

## collinear framework

- simple product of PDF and IFF

Ex.: SIDIS

$$A_{\text{SIDIS}}^{\sin(\phi_R + \phi_S)}(x, z, M_h^2) \sim - \frac{\sum_q e_q^2 h_1^q(x) \frac{|\mathbf{R}_T|}{M_h} H_{1,q}^{\triangleleft}(z, M_h^2)}{\sum_q e_q^2 f_1^q(x) D_{1,q}(z, M_h^2)}$$

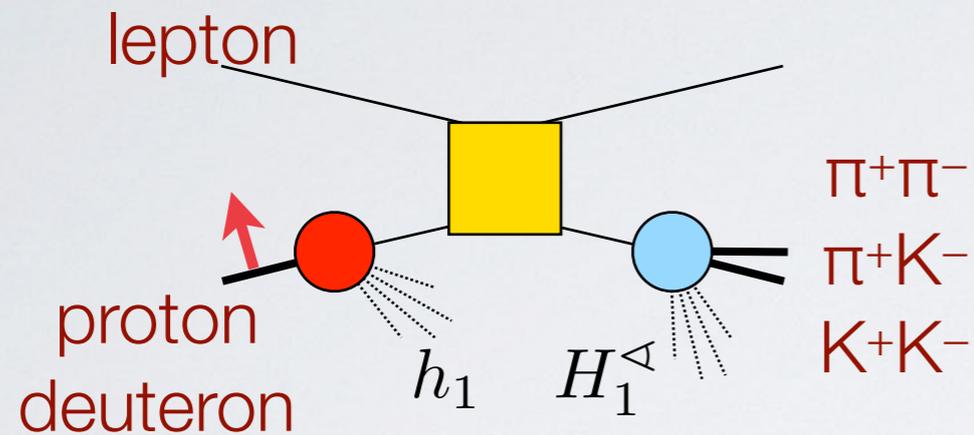
x-dependence of  $A_{\text{SIDIS}}$  all in PDF

- factorization theorems for all hard processes  
→ universality of  $h_1 H_1^{\triangleleft}$  mechanism

# available experimental data

## factorization theorems for all hard processes

SIDIS



## data used in the global fit

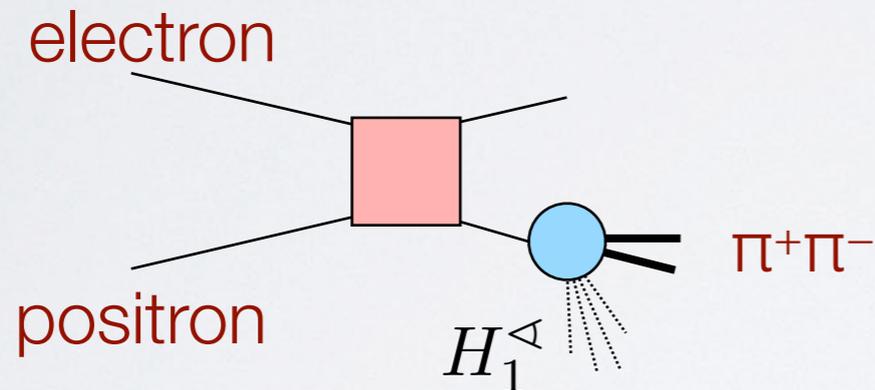


Airapetian et al.,  
*JHEP* **0806** (08) 017



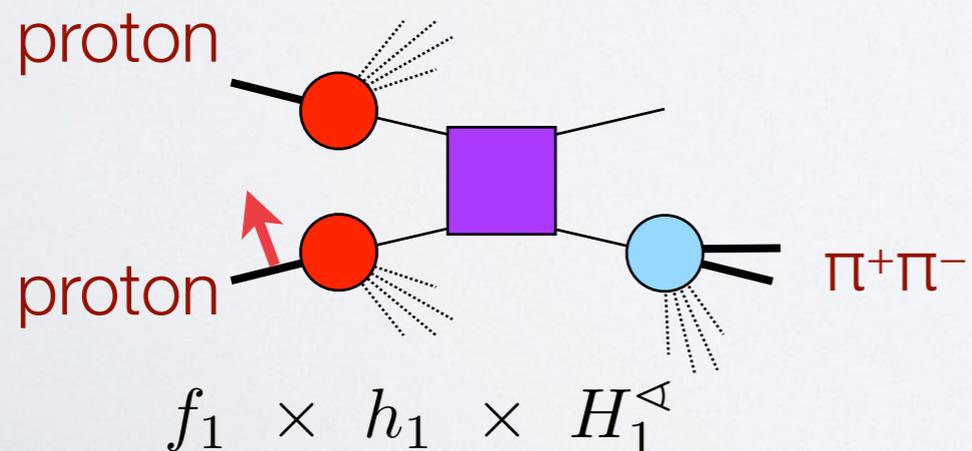
Adolph et al., *P.L.* **B713** (12)  
Braun et al., *E.P.J. Web Conf.* **85** (15)

$e^+e^-$



Vossen et al., *P.R.L.* **107** (11) 072004

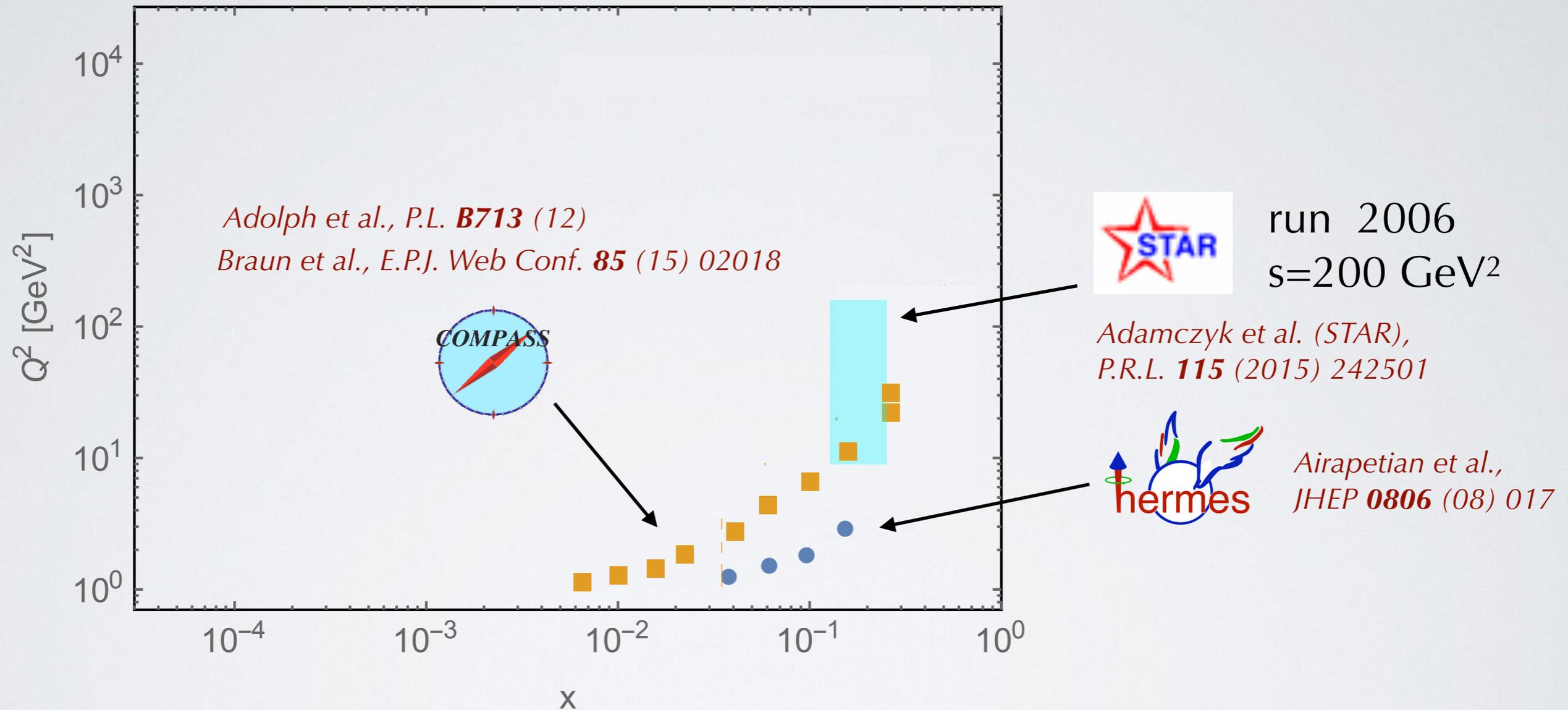
$p p^\uparrow$



run 2006 ( $s=200$ )

Adamczyk et al. (STAR),  
*P.R.L.* **115** (2015) 242501

# the phase space



- mostly medium/high  $x \rightarrow$  not enough for sea quark explorations
- guess low- $x$  behavior (relevant for calculation of tensor charge)

# choice of functional form

functional form whose Mellin transform can be computed analytically  
and complying with Soffer Bound at any  $x$  and scale  $Q^2$

$$h_1^{qv}(x; Q_0^2) = F^{qv}(x) \left[ \text{SB}^q(x) + \overline{\text{SB}}^{\bar{q}}(x) \right]$$



**Soffer Bound**

$$2|h_1^q(x, Q^2)| \leq 2 \text{SB}^q(x, Q^2) = |f_1^q(x, Q^2) + g_1^q(x, Q^2)|$$

MSTW08

DSSV

# choice of functional form

functional form whose Mellin transform can be computed analytically and complying with Soffer Bound at any  $x$  and scale  $Q^2$

$$h_1^{q_v}(x; Q_0^2) = F^{q_v}(x) \left[ \text{SB}^q(x) + \overline{\text{SB}}^{\bar{q}}(x) \right]$$

Soffer Bound

$$2|h_1^q(x, Q^2)| \leq 2 \text{SB}^q(x, Q^2) = |f_1^q(x, Q^2) + g_1^q(x, Q^2)|$$

MSTW08

DSSV

$$F^{q_v}(x) = \frac{N_{q_v}}{\max_x [|F^{q_v}(x)|]} x^{A_{q_v}} [1 + B_{q_v} \text{Ceb}_1(x) + C_{q_v} \text{Ceb}_2(x) + D_{q_v} \text{Ceb}_3(x)]$$

Ceb<sub>n</sub>(x) Chebyshev polynomial

10 fitting parameters

constrain parameters

$$|N_{q_v}| \leq 1 \Rightarrow |F^{q_v}(x)| \leq 1 \quad \text{Soffer Bound ok at any } Q^2$$

# low-x behavior

$$\lim_{x \rightarrow 0} xSB^q(x) \propto x^{a_q} \left. \vphantom{\lim_{x \rightarrow 0} xSB^q(x)} \right\}$$

$$\lim_{x \rightarrow 0} F^{qv}(x) \propto x^{A_q}$$

$$h_1^q(x) \stackrel{x \rightarrow 0}{\approx} x^{A_q + a_q - 1}$$

tensor charge  $\delta q(Q^2) = \int_{x_{\min}}^1 dx h_1^{q-\bar{q}}(x, Q^2)$

constrain parameters

# low-x behavior

$$\left. \begin{aligned} \lim_{x \rightarrow 0} xSB^q(x) &\propto x^{a_q} \\ \lim_{x \rightarrow 0} F^{qv}(x) &\propto x^{A_q} \end{aligned} \right\}$$

$$h_1^q(x) \stackrel{x \rightarrow 0}{\approx} x^{A_q + a_q - 1}$$

$$\text{tensor charge } \delta q(Q^2) = \int_{x_{\min}}^1 dx h_1^{q-\bar{q}}(x, Q^2)$$

constrain parameters

1)  $\delta q$  finite  $\Rightarrow A_q + a_q > 0$

2) “massive” jet in DIS  $\rightarrow h_1$  at twist 3  
violation of Burkardt-Cottingham s.r.

$$\int_0^1 dx g_2(x) \propto \int_0^1 dx \frac{h_1(x)}{x} \longrightarrow A_q + a_q > 1$$

*Accardi and Bacchetta, P.L. **B773** (17) 632*

3) small-x dipole picture  $\Rightarrow h_1^{qv}(x) \stackrel{x \rightarrow 0}{\approx} x^{1-2\sqrt{\frac{\alpha_s(Q^2)N_c}{2\pi}}} \longrightarrow \text{at } Q_0 \quad A_q + a_q \sim 1$

*Kovchegov & Sievert, arXiv:1808.10354*

# low-x behavior

$$\left. \begin{aligned} \lim_{x \rightarrow 0} xSB^q(x) &\propto x^{a_q} \\ \lim_{x \rightarrow 0} F^{qv}(x) &\propto x^{A_q} \end{aligned} \right\}$$

$$h_1^q(x) \stackrel{x \rightarrow 0}{\approx} x^{A_q + a_q - 1}$$

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constrain parameters

low-x behavior important

1)  $\delta q$  finite  $\Rightarrow A_q + a_q > 0$

2) “massive” jet in DIS  $\rightarrow h_1$  at twist 3 violation of Burkardt-Cottingham s.r.

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*Accardi and Bacchetta, P.L. **B773** (17) 632*

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*Kovchegov & Sievert, arXiv:1808.10354*

our choice

$$A_q + a_q > \frac{1}{3}$$

$$\left| \int_0^{x_{\min}} dx \right| \sim 1\% \text{ of } \left| \int_{x_{\min}}^1 dx \right|$$

for  $x_{\min}=10^{-6}$  from MSTW08

# theoretical uncertainties

## unpolarized Di-hadron Fragmentation Function $D_1$

- **quark**  $D_1^q$  is **well** constrained by  $e^+e^- \rightarrow (\pi^+\pi^-) X$  (Montecarlo)
- **gluon**  $D_1^g$  is **not** constrained by  $e^+e^- \rightarrow (\pi^+\pi^-) X$  (currently, LO analysis)
- **no data** available yet for  $p p \rightarrow (\pi^+\pi^-) X$

we don't know anything about the gluon  $D_1^g$

our choice:

set  $D_1^g(Q_0) =$

$$\begin{cases} 0 \\ D_1^u(Q_0) / 4 \\ D_1^u(Q_0) \end{cases}$$

←  $\sim$  1-hadron  $D_1^g(Q_0)$

deteriorates our  $e^+e^-$  fit as  $\chi^2/\text{dof} =$

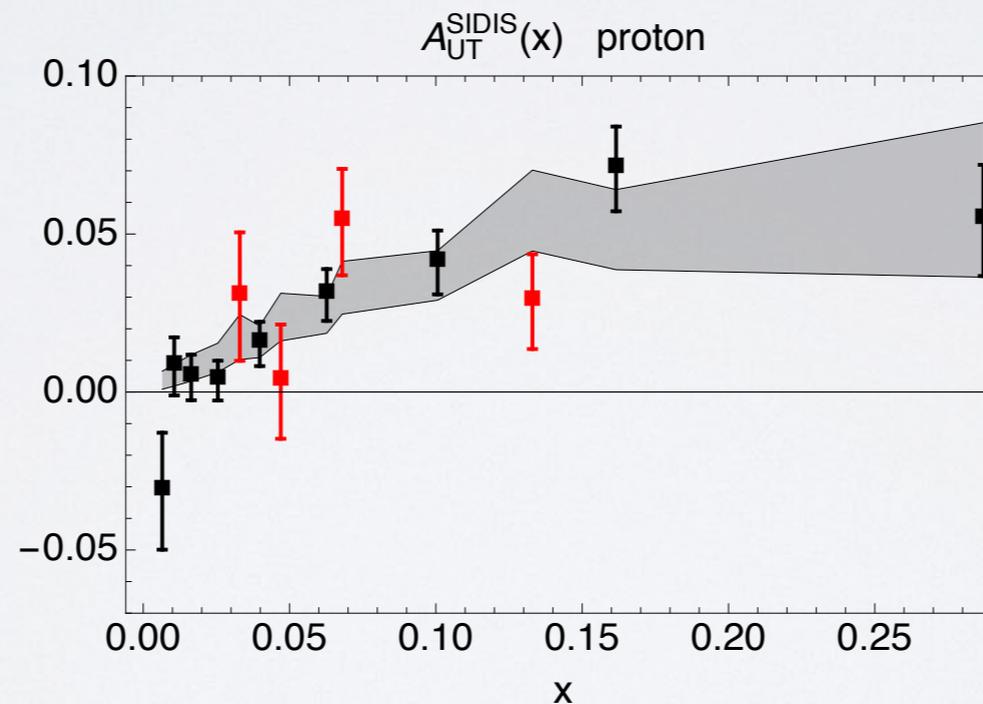
$$\begin{cases} 1.69 & 1.28 \\ 1.81 & 1.37 \\ 2.96 & 2.01 \end{cases}$$

background  $\rho$  channels

# statistical uncertainty

## the bootstrap method

- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50, 100, 200 sets... until average and standard deviation reproduce original exp. points (here,  $200 \times 3 = 600$ )
- exclude largest and smallest 5% => 90% band



automatically accounts for correlations

# results

global fit published in

*Radici and Bacchetta, P.R.L. 120 (18) 192001*

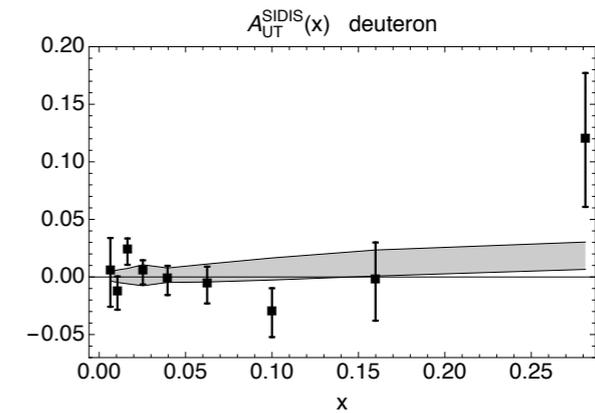
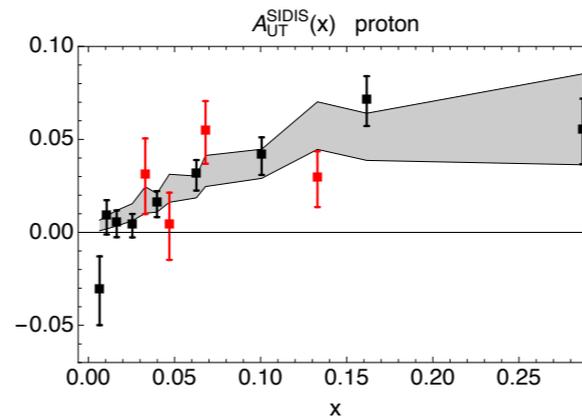
## SIDIS



*Adolph et al., P.L. B713 (12)*



*Airapetian et al.,  
JHEP 0806 (08) 017*

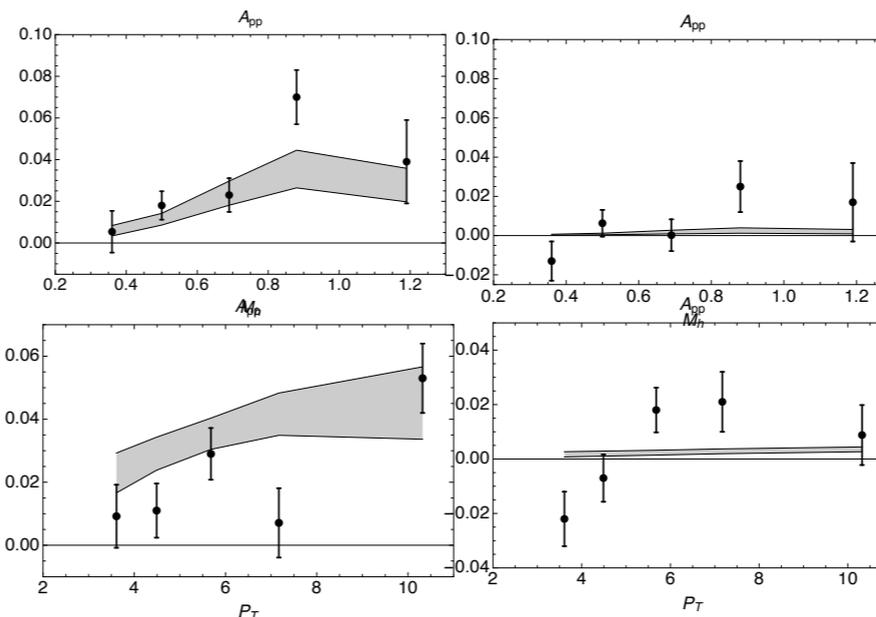


## pp collisions

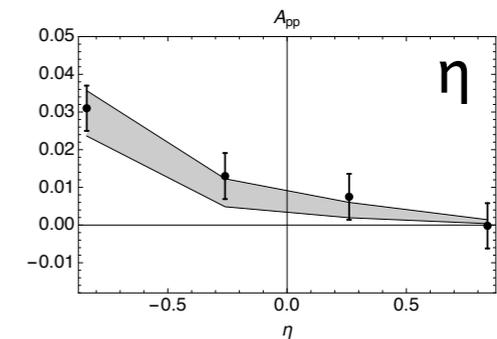


*Adamczyk et al.,  
P.R.L. 115 (2015) 242501*

$M_h, \eta < 0$



$M_h, \eta > 0$



$p_T, \eta < 0$

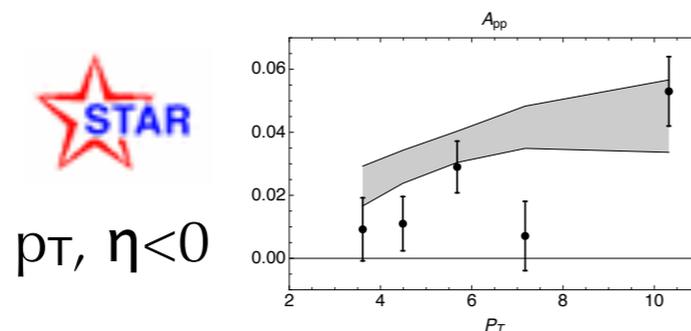
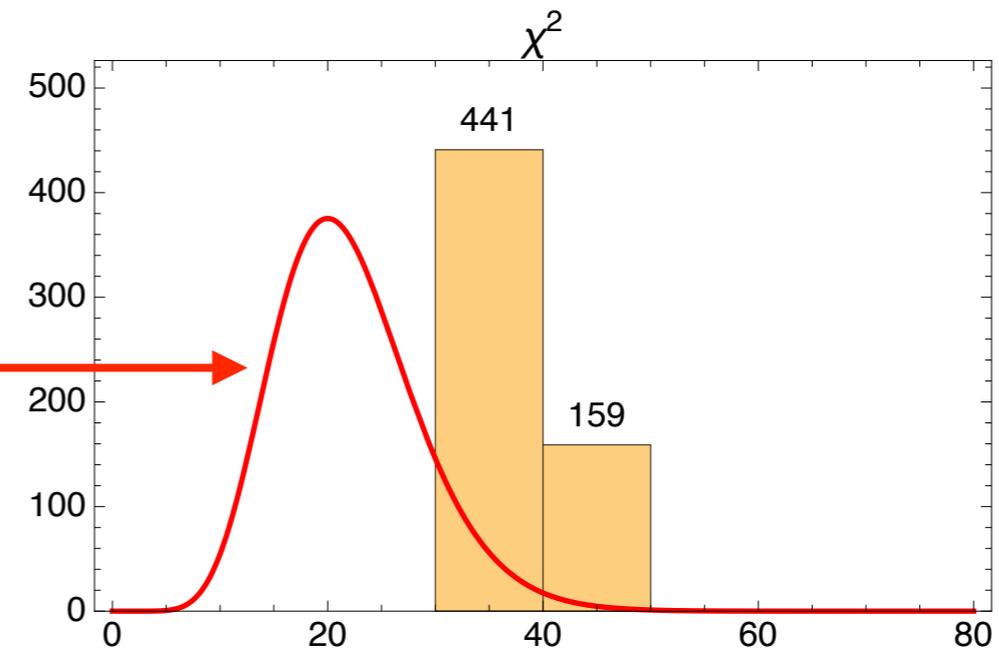
$p_T, \eta > 0$

# $\chi^2$ of the fit

|   |                  |   |  |   |
|---|------------------|---|--|---|
| proton SIDIS  | 13 data points = | 4  | + 9                                    |  |
| deuteron SIDIS  | 9 data points =  |   | + 9                                    |   |
|  | 24 data points   | (4 $\eta$ ) $\times$ $\frac{4}{24}$   | + (10 $M_h$ ) $\times$ $\frac{10}{24}$ | + (10 $p_T$ ) $\times$ $\frac{10}{24}$  |
| global fit  | 10 parameters    |   |  |   |
| d.o.f.  | 22               |   |  |   |

probability density function of  $\chi^2$  distribution for 22 d.o.f.

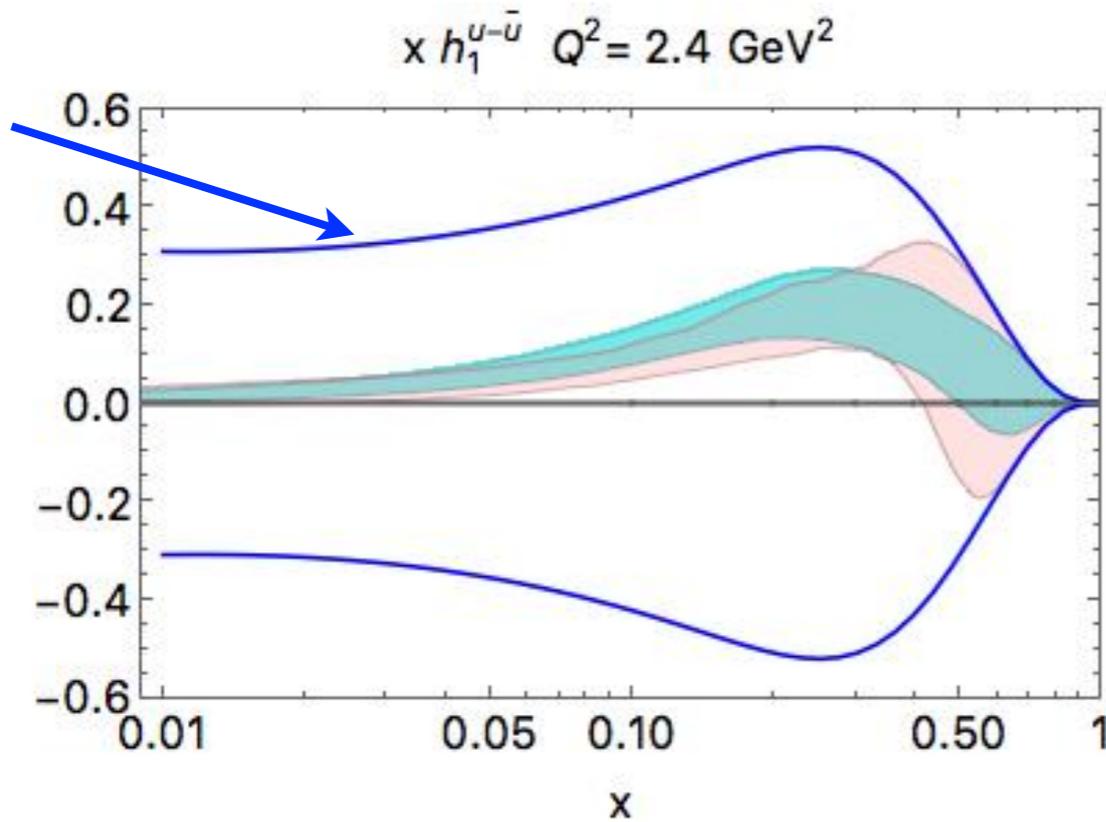
for  $\chi^2/\text{dof} = 1$  perfect overlap



$$\chi^2/\text{dof} = 1.76 \pm 0.11$$

# comparison with previous fit

Soffer bound



*Radici & Bacchetta,*  
*P.R.L. **120** (18) 192001*

global fit

up

higher  
precision

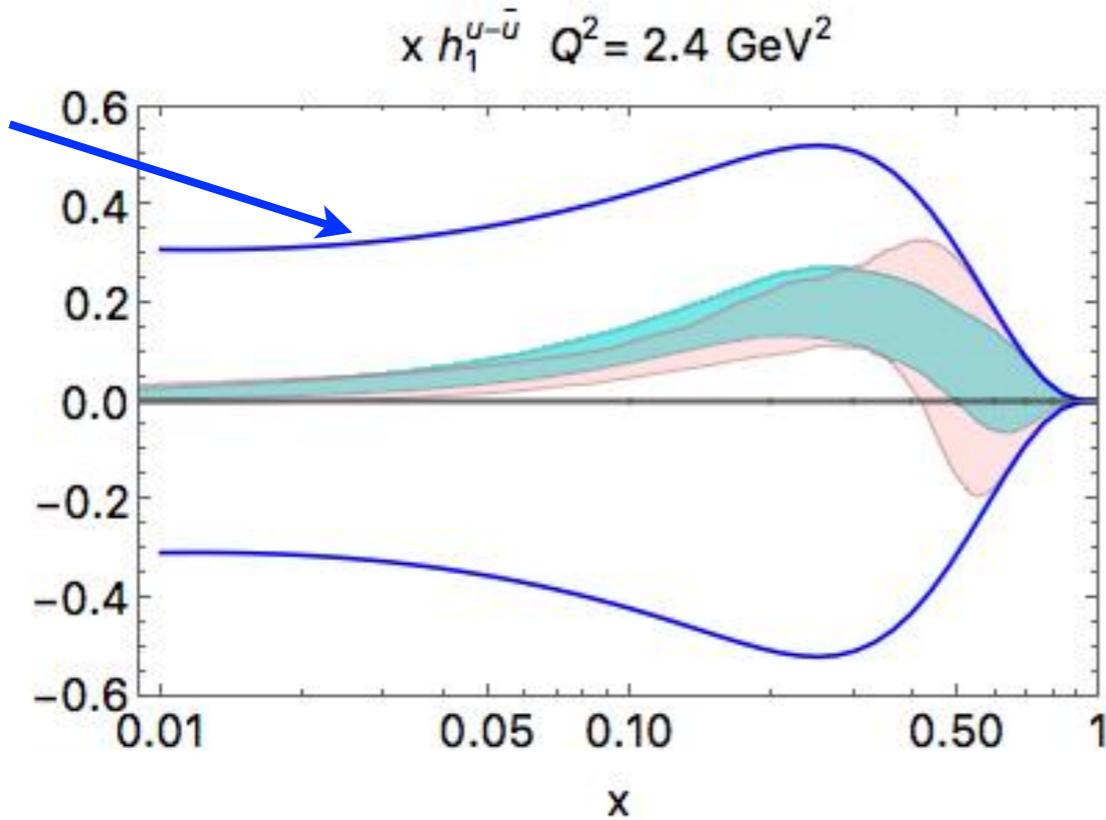
old fit (only SIDIS data)

*Radici et al.,*  
*JHEP **1505** (15) 123*

equivalent to  
Collins extraction

# comparison with previous fit

Soffer bound



*Radici & Bacchetta,*  
*P.R.L. 120 (18) 192001*

global fit

up

higher precision

old fit (only SIDIS data)

*Radici et al.,*  
*JHEP 1505 (15) 123*

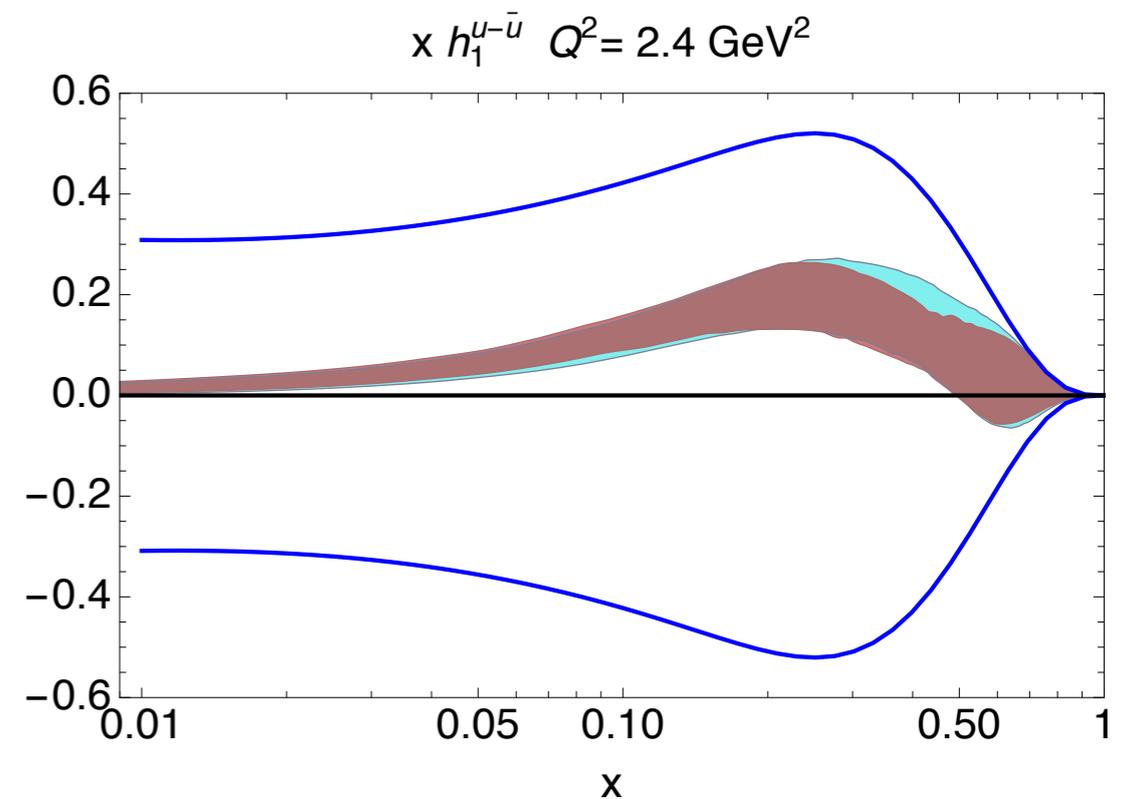
equivalent to  
Collins extraction

up

insensitive to  
uncertainty on  
gluon  $D_1$

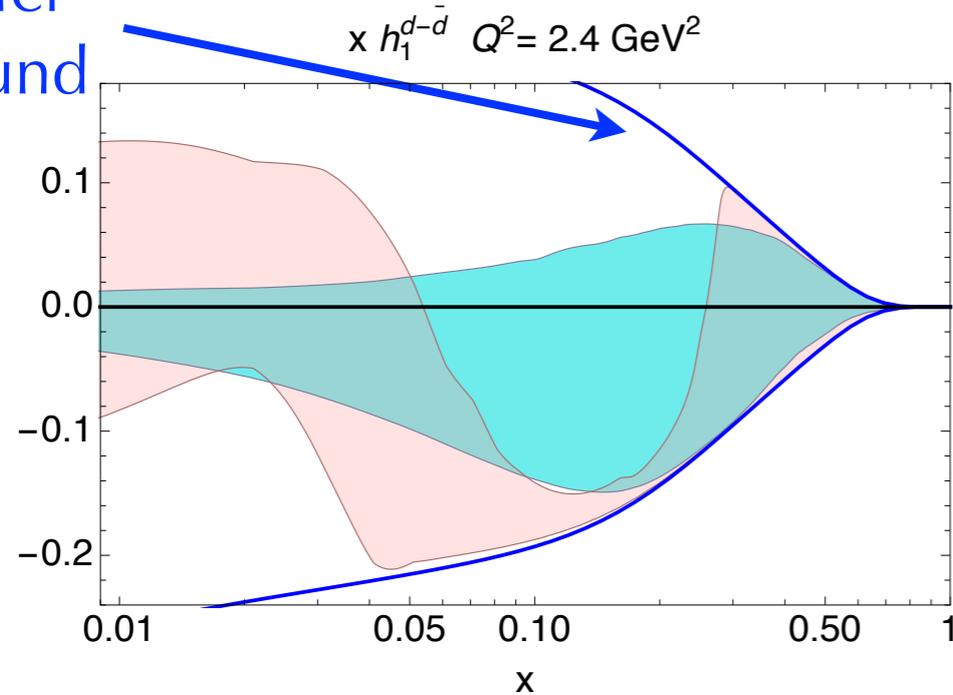
$$D_{1g}(Q_0) = 0$$

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_1^u/4 \\ D_1^u \end{cases}$$



# comparison with previous fit

Soffer bound



*Radici & Bacchetta,*  
*P.R.L. **120** (18) 192001*

global fit

old fit

*Radici et al.,*  
*JHEP **1505** (15) 123*

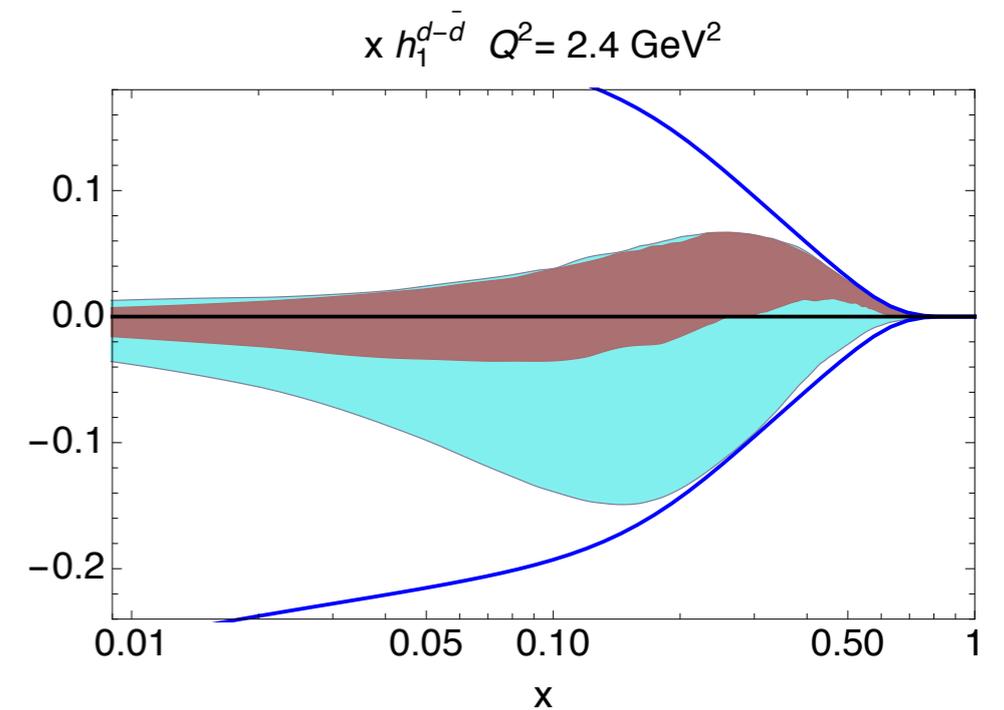
down

down

sensitive to  
uncertainty on  
gluon  $D_1$

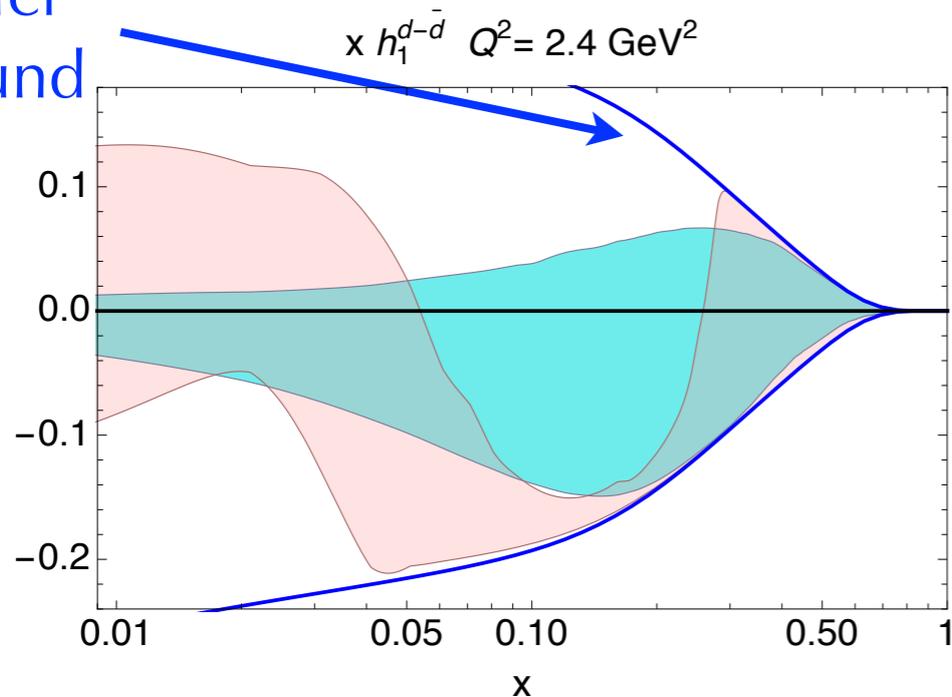
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# comparison with previous fit

Soffer bound



Radici & Bacchetta,  
*P.R.L.* **120** (18) 192001

global fit

old fit

Radici et al.,  
*JHEP* **1505** (15)

down

p-p: up~down, gluon @LO  
but  
SIDIS: up~(8x)down, gluon @NLO

need data from target more sensitive  
to down: deuteron, <sup>3</sup>He ...

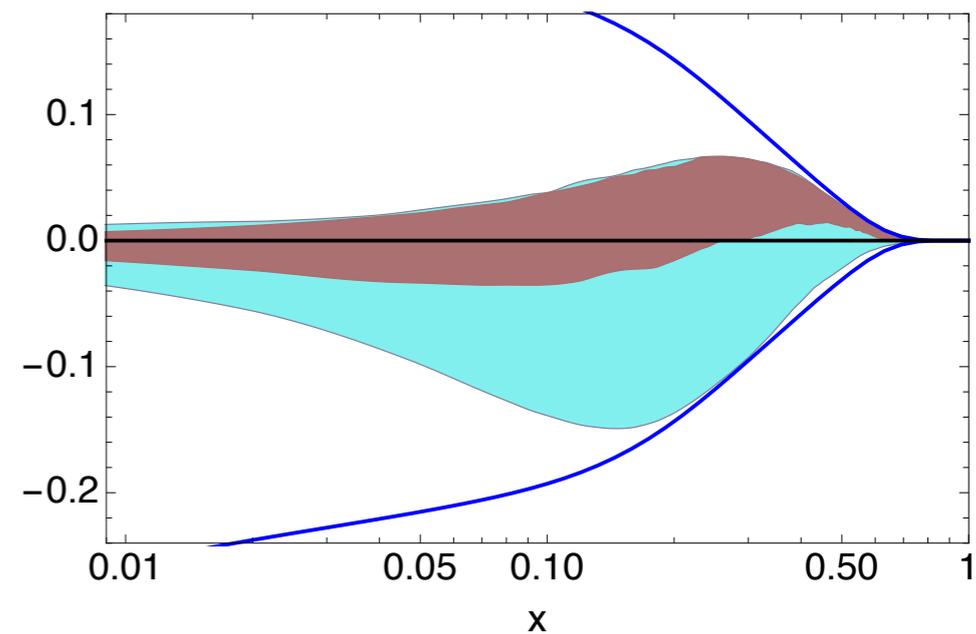
down

sensitive to  
uncertainty on  
gluon  $D_1$

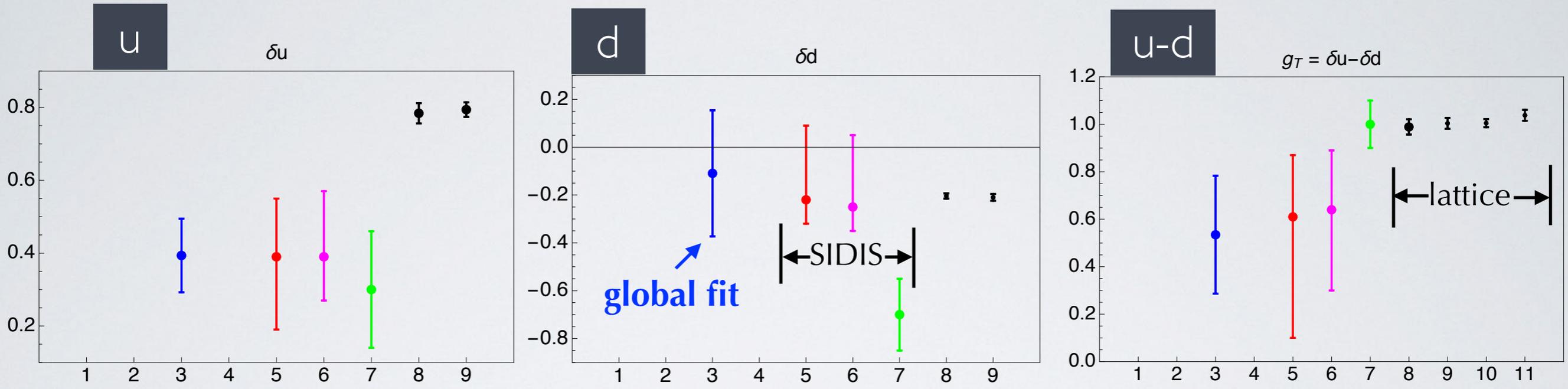
$$D_{1g}(Q_0) = 0$$

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_1^u / 4 \\ D_1^u \end{cases}$$

$x h_1^{d-\bar{d}} \quad Q^2 = 2.4 \text{ GeV}^2$



# tensor charge



$Q^2=4 \text{ GeV}^2$  \*

**JAM** includes  
"lattice data"

Radici & Bacchetta,  
*P.R.L.* 120 (18) 192001

3) **global fit '17**

Kang et al., *P.R. D*93 (16) 014009

5) **"TMD fit" \*  $Q^2=10$**

Anselmino et al., *P.R. D*87 (13) 094019

6) **Torino fit \*  $Q^2=1$**

Lin et al., *P.R.L.* 120 (18) 152502

7) **JAM fit '17 \*  $Q_0^2=2$**

8) **PNDME '18**

*Gupta et al., P.R. D*98 (18) 034503

9) **ETMC '17**

*Alexandrou et al., P.R. D*95 (17) 114514;  
*E P.R. D*96 (17) 099906

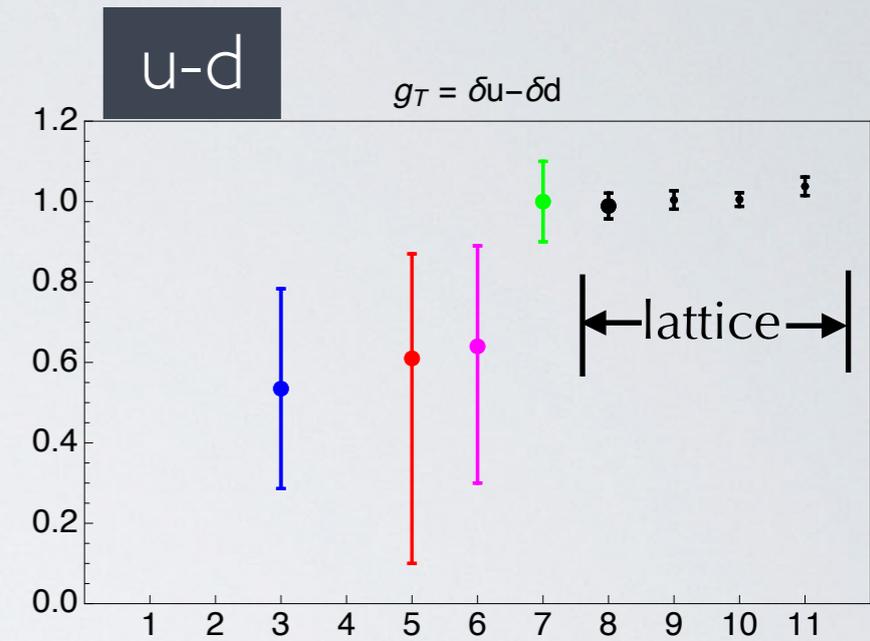
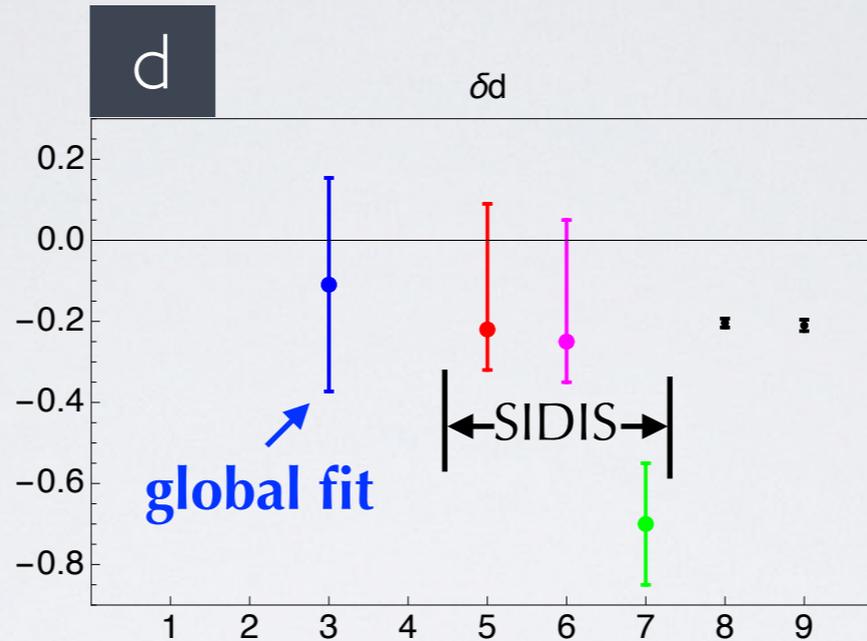
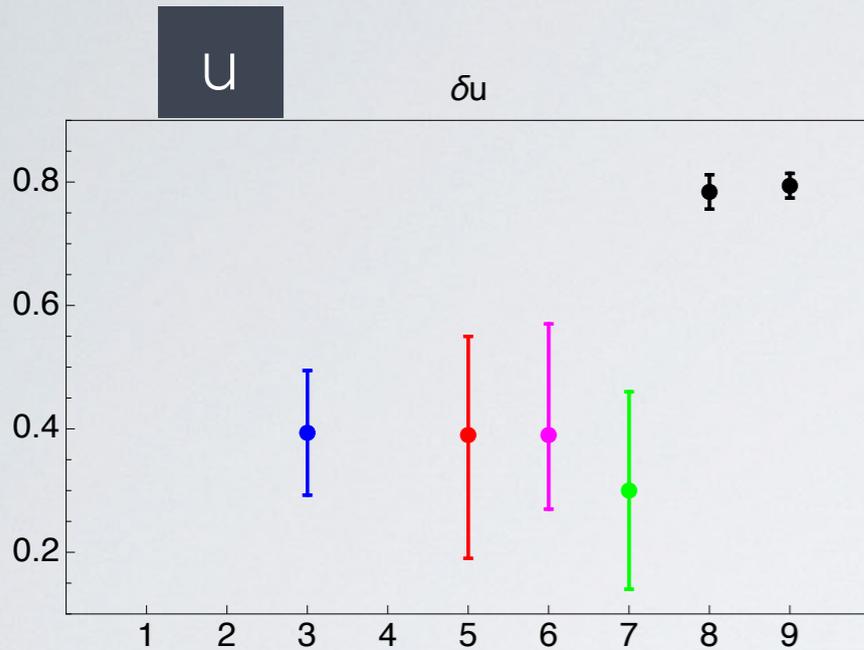
10) **RQCD '14**

*Bali et al., P.R. D*91 (15)

11) **LHPC '12**

*Green et al., P.R. D*86 (12)

# tensor charge



no simultaneous compatibility between lattice and phenomenology

$Q^2=4 \text{ GeV}^2$  \*

JAM includes "lattice data"

Radici & Bacchetta, P.R.L. 120 (18) 192001

3) global fit '17

Kang et al., P.R. D93 (16) 014009

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Alexandrou et al., P.R. D95 (17) 114514;  
E P.R. D96 (17) 099906

10) RQCD '14

Bali et al., P.R. D91 (15)

11) LHPC '12

Green et al., P.R. D86 (12)

# results

global fit published in

*Radici and Bacchetta, P.R.L. 120 (18) 192001*

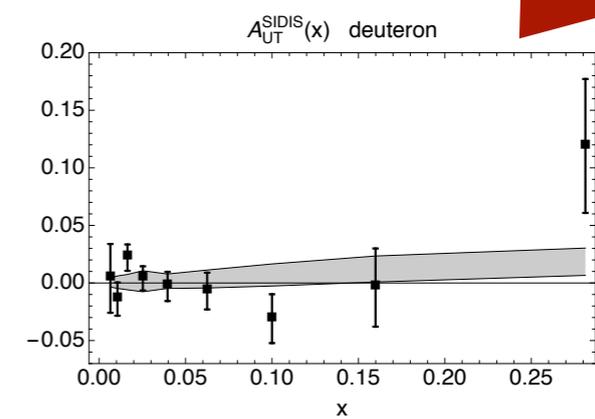
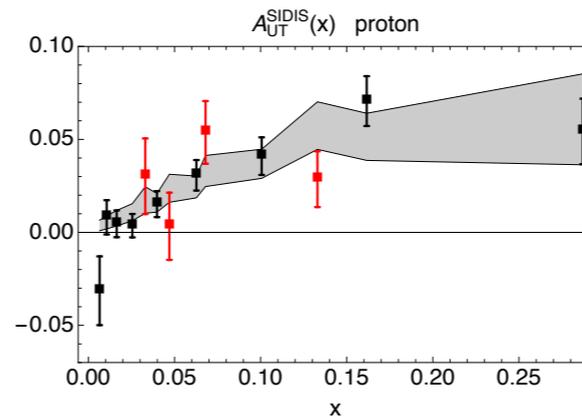
## SIDIS



*Adolph et al., P.L. B713 (12)*



*Airapetian et al.,  
JHEP 0806 (08) 017*

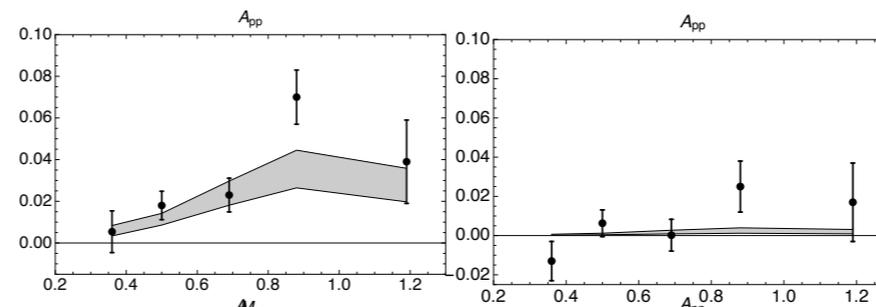


## pp collisions

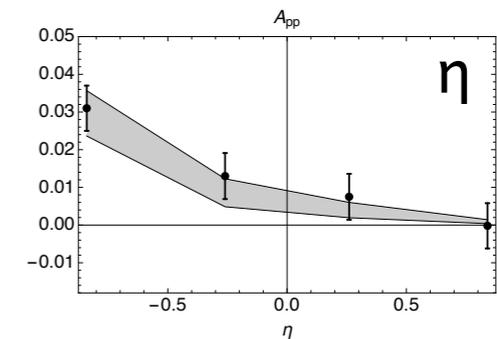


*Adamczyk et al.,  
P.R.L. 115 (2015) 242501*

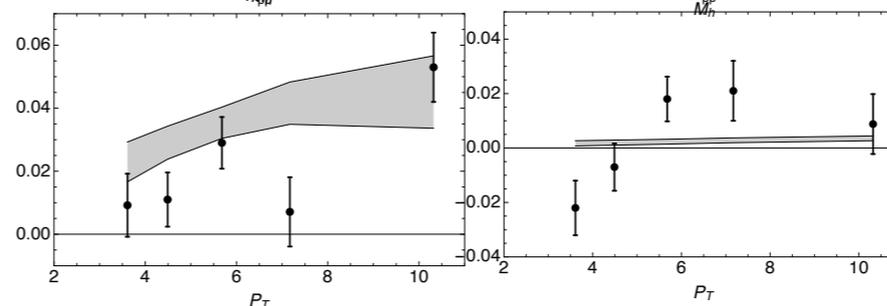
$M_h, \eta < 0$



$M_h, \eta > 0$



$p_T, \eta < 0$



$p_T, \eta > 0$

# Compass pseudo-data

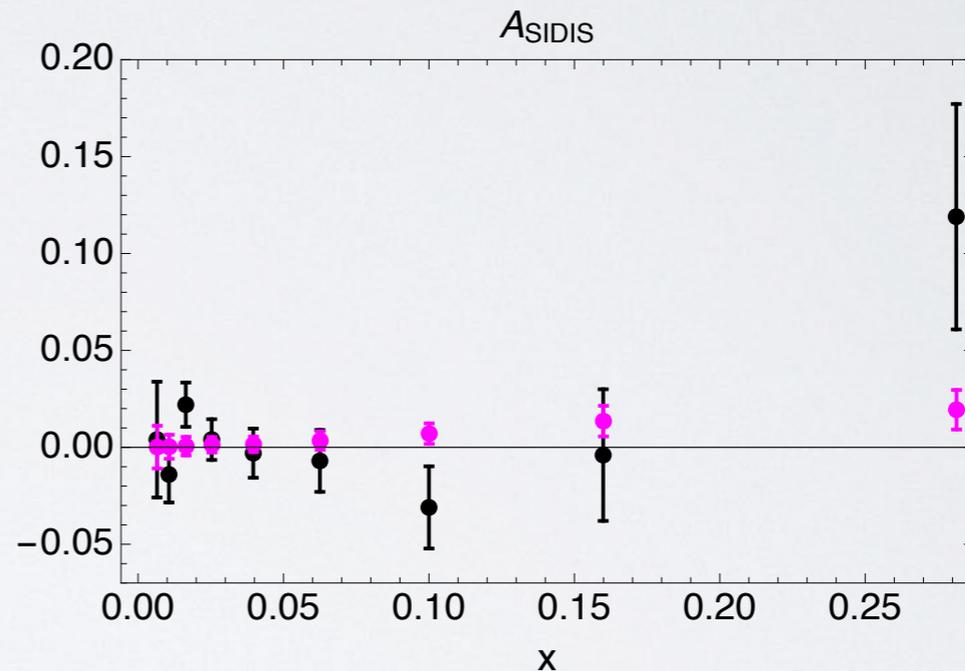
add to previous set of data  
a new set of SIDIS pseudo-data for **deuteron** target



*Adolph et al., P.L. B713 (12)*



pseudodata



statistical error  $\sim 0.6 \times$  [ error in 2010 proton data ]  
<A> = average value of replicas in previous global fit

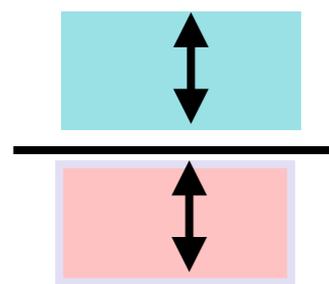
study impact on precision of previous global fit

# impact of pseudo-data for deuteron

global fit + pseudodata

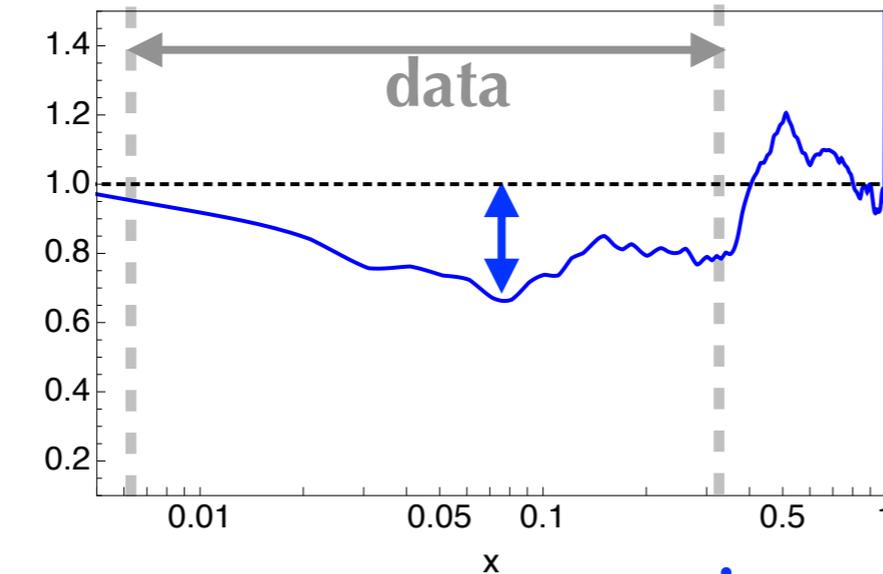
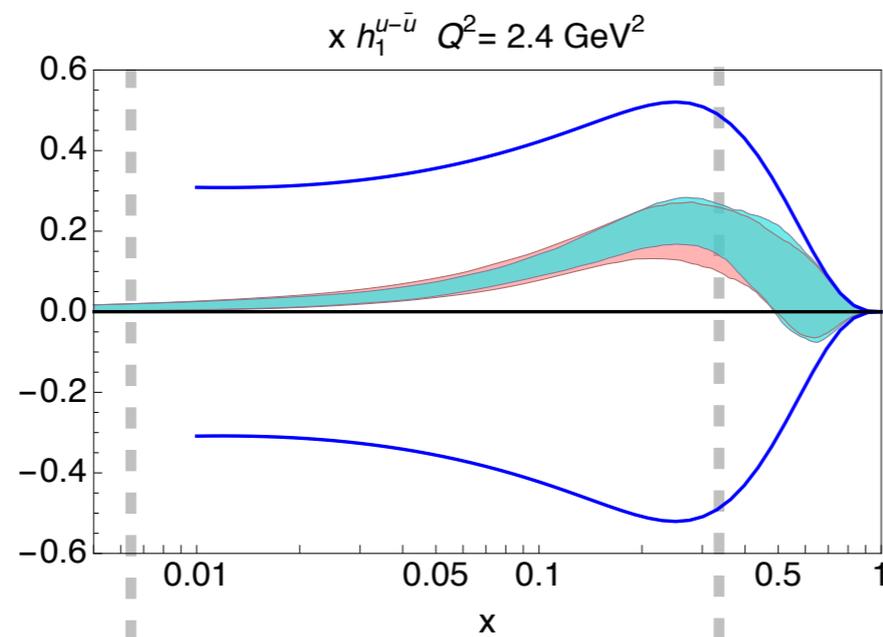
global fit

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_1^u/4 \\ D_1^u \end{cases}$$

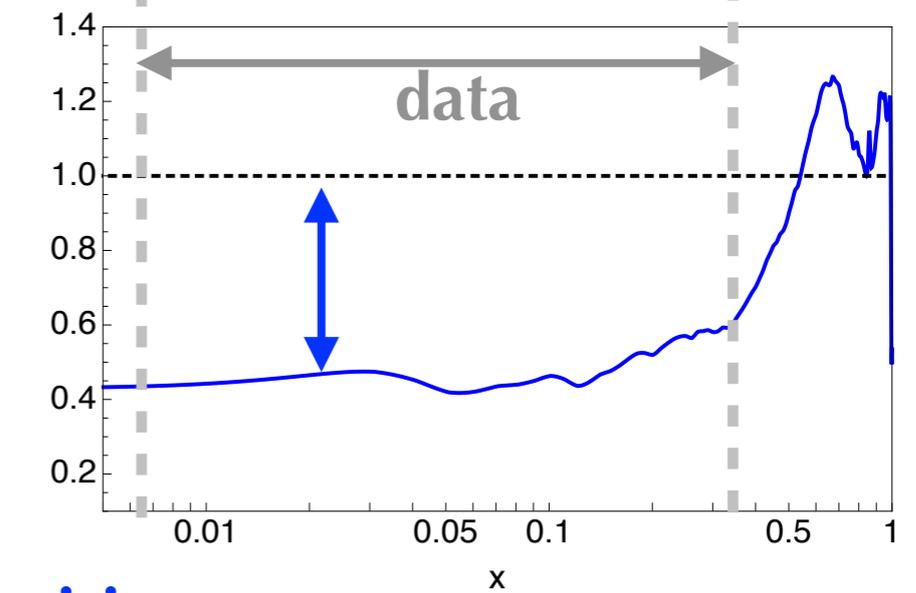
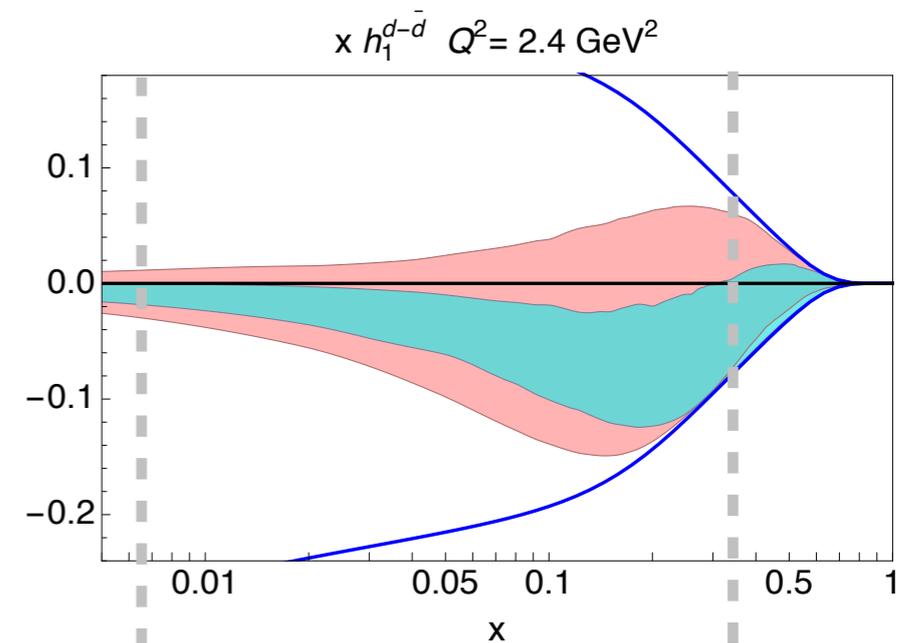


ratio of widths

up

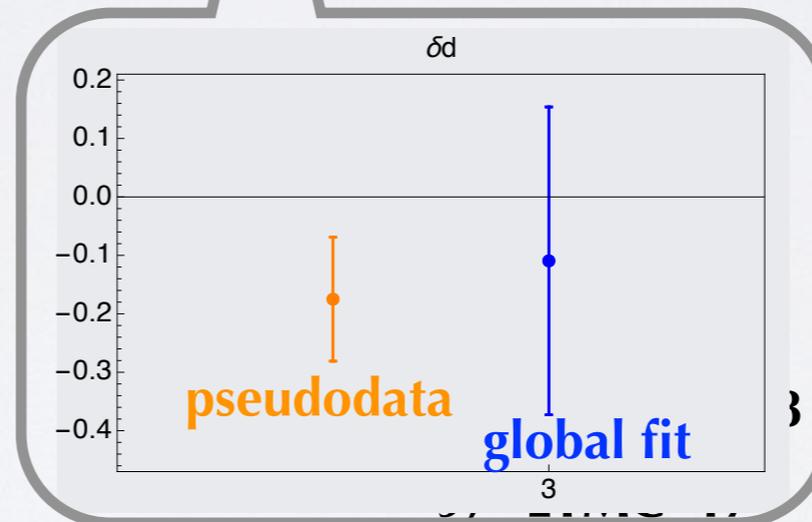
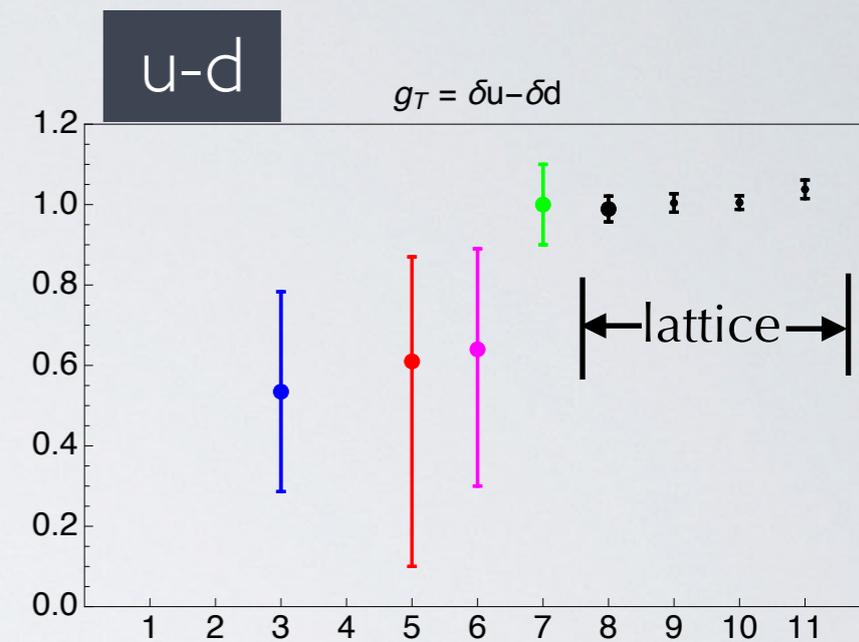
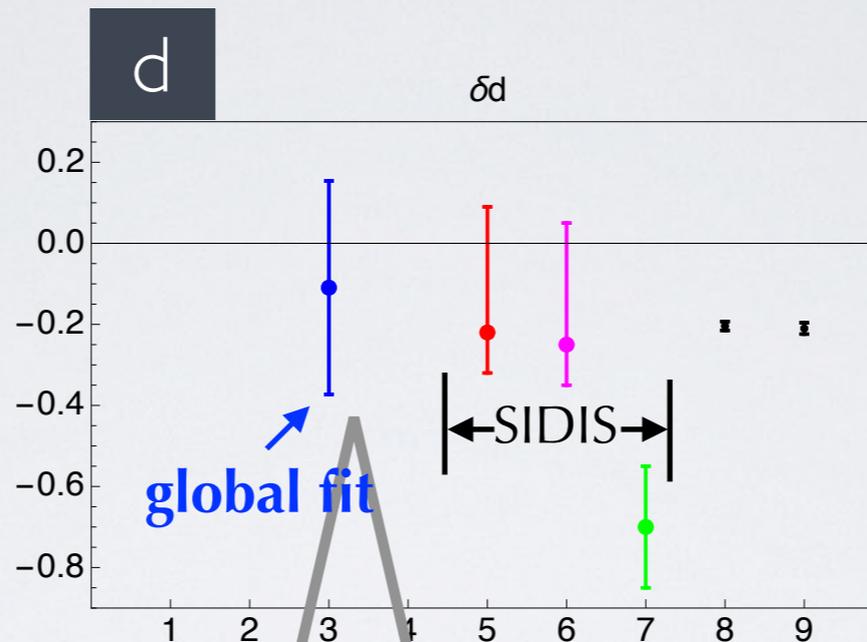
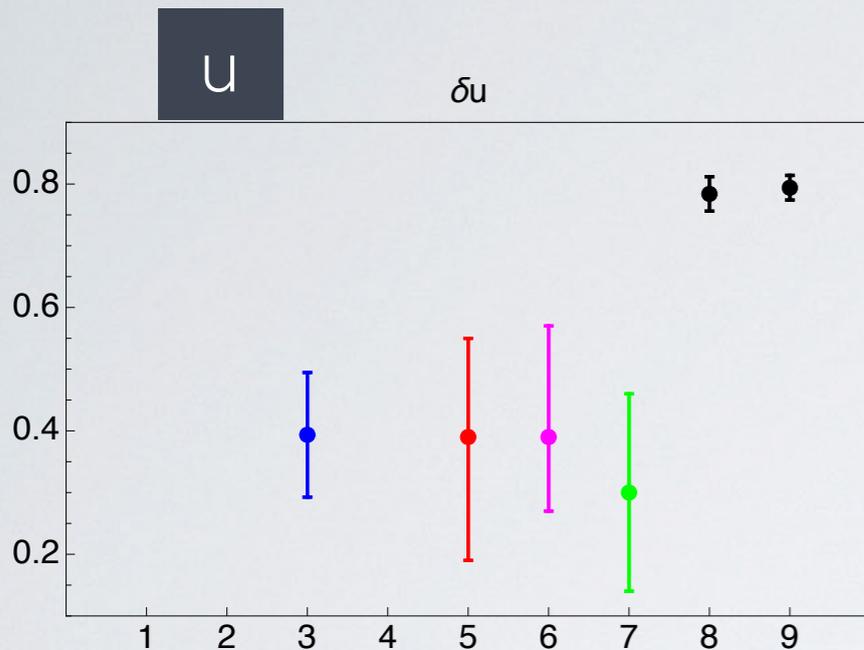


down



increase precision

# tensor charge



**JAM** includes  
"lattice data"

Radici & Bacchetta,  
*P.R.L.* 120 (18) 192001

3) **global fit '17**

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Gupta et al., *P.R.* D98 (18) 034503

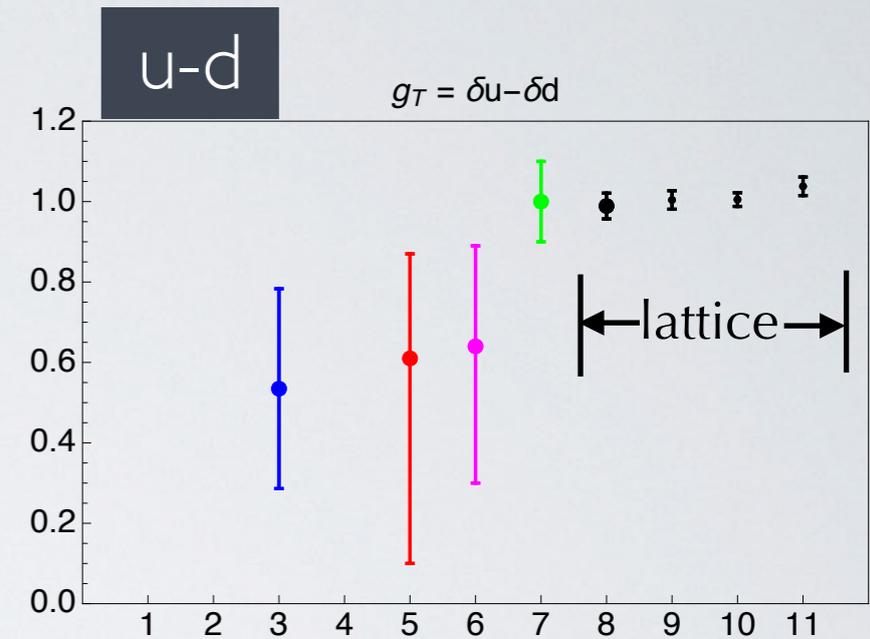
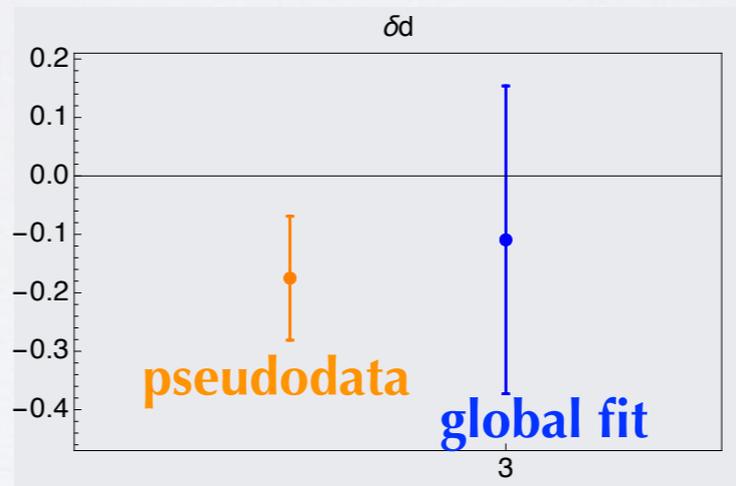
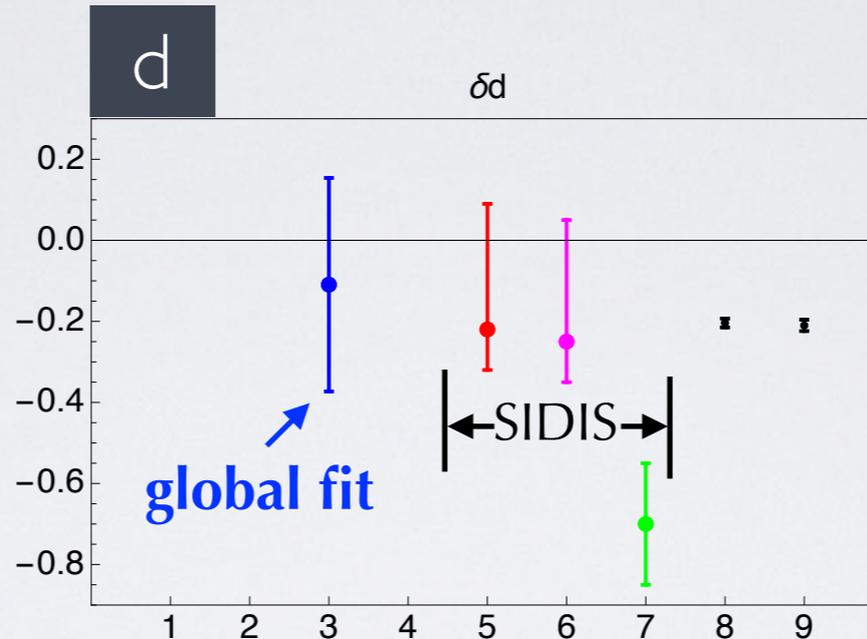
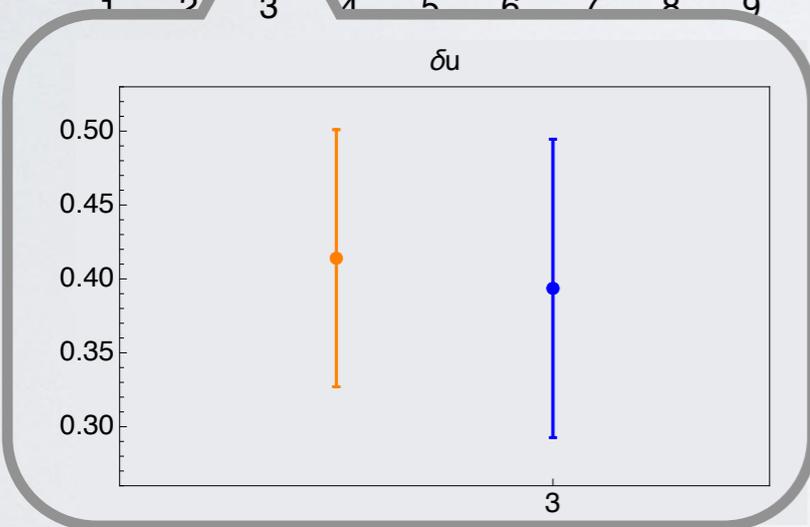
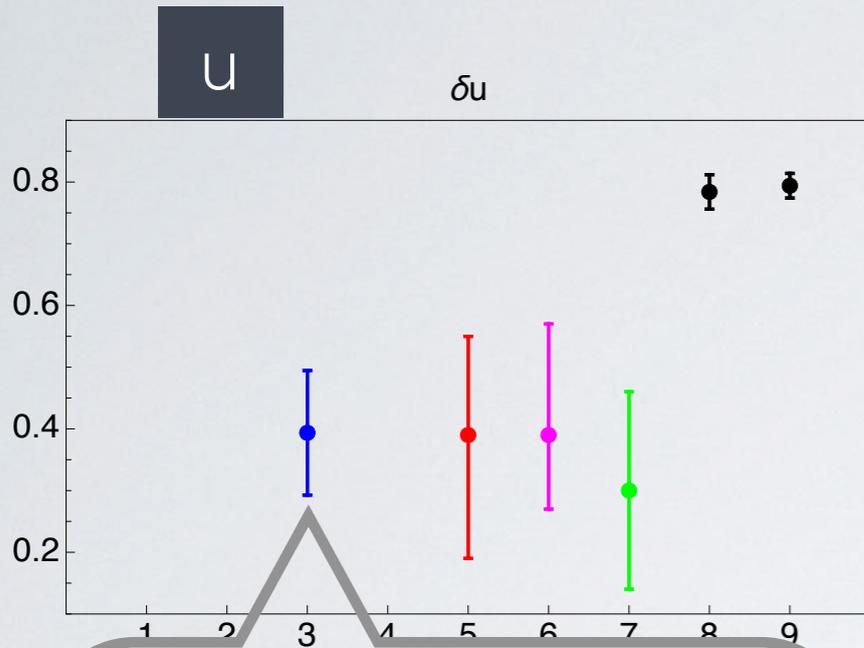
Alexandrou et al., *P.R.* D95 (17) 114514;

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# tensor charge



**JAM** includes  
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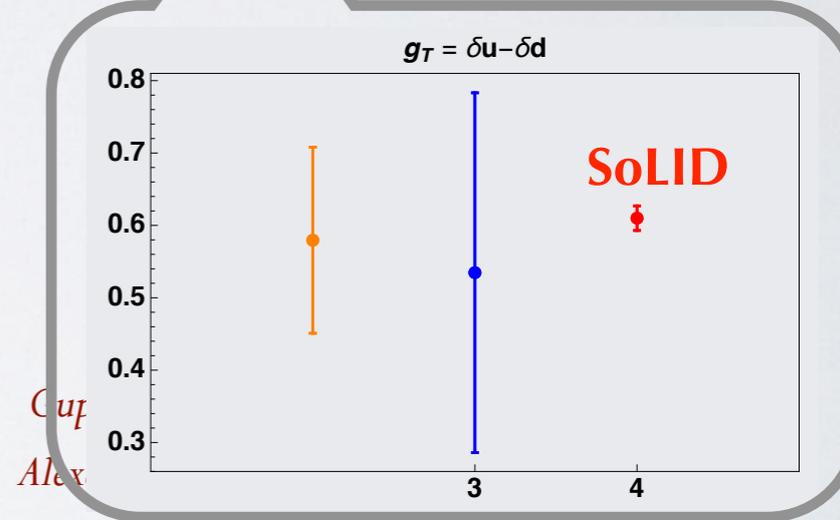
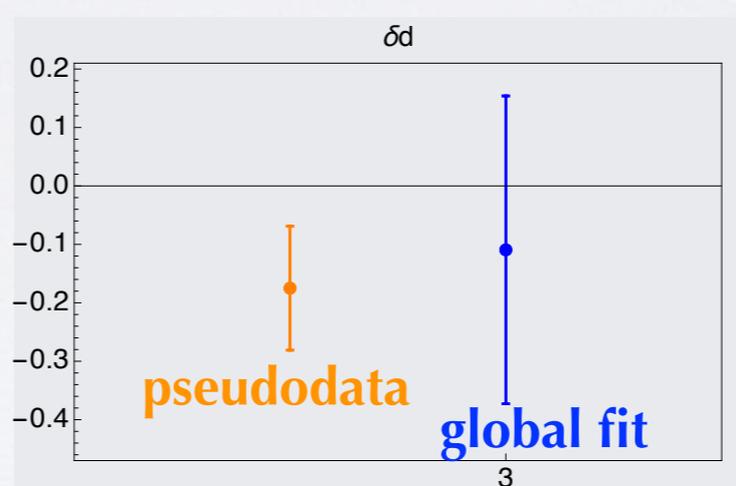
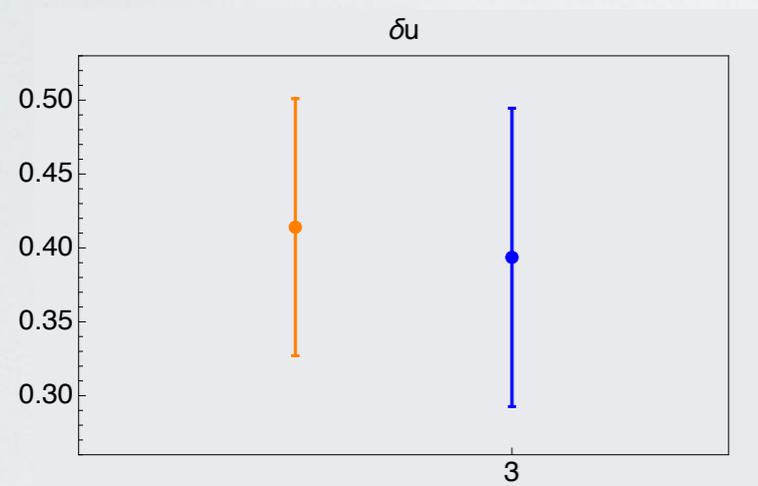
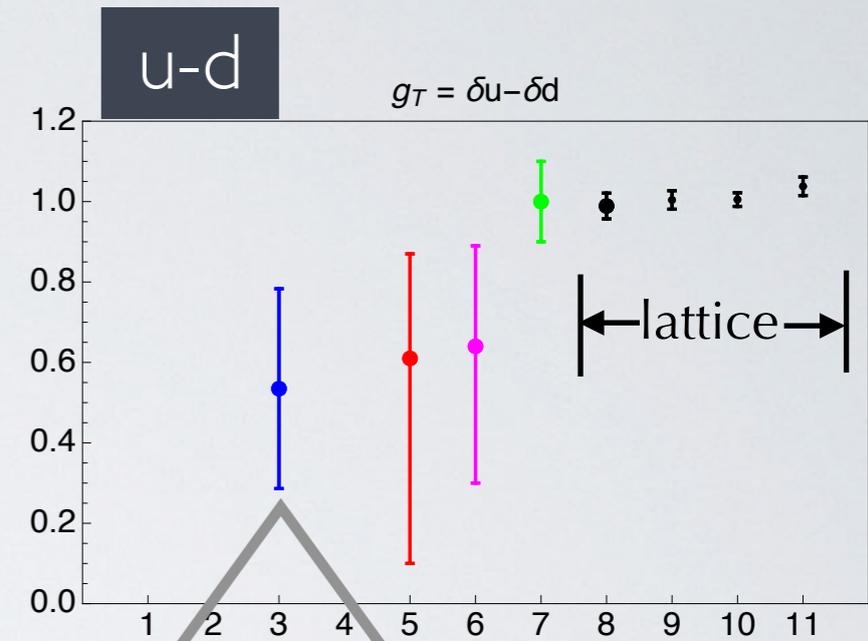
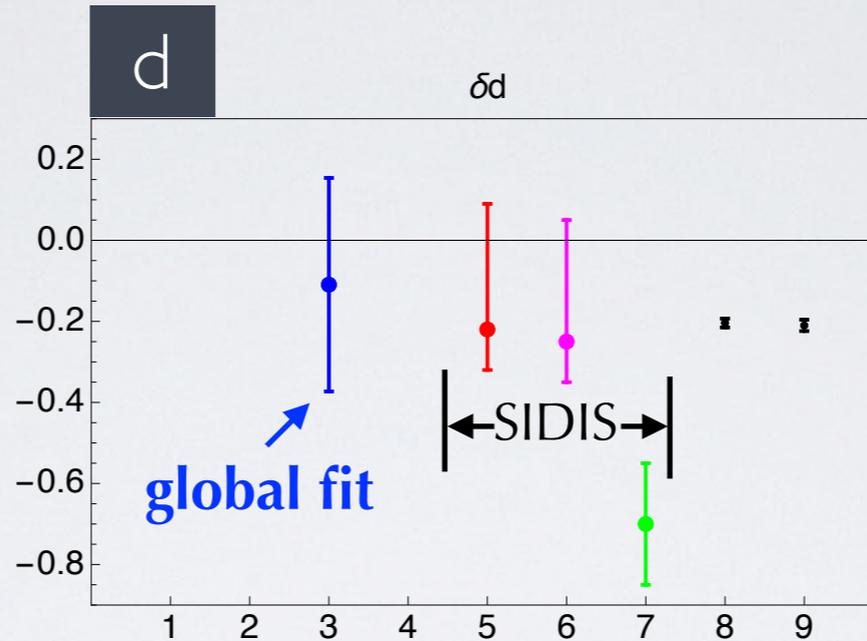
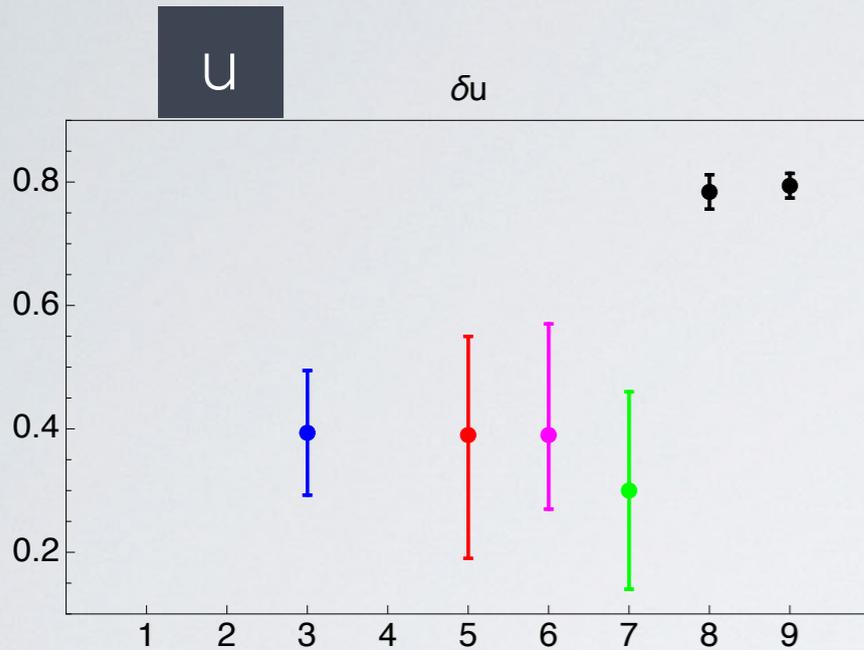
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Cu $\bar{u}$   
Alex

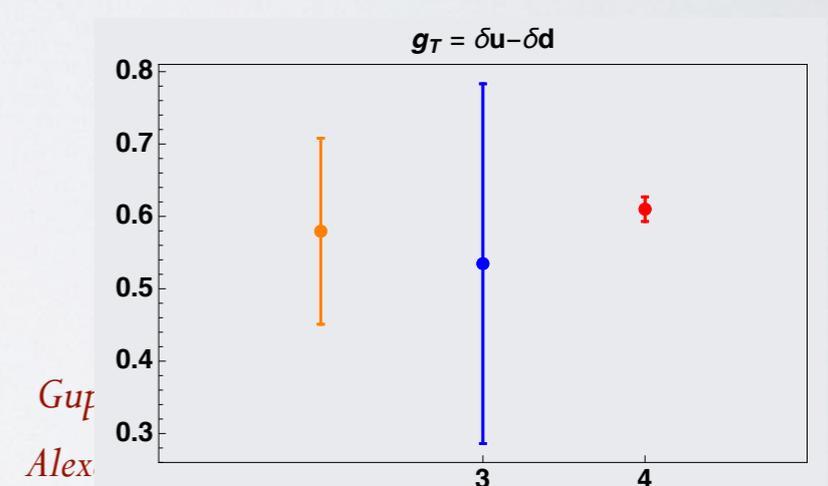
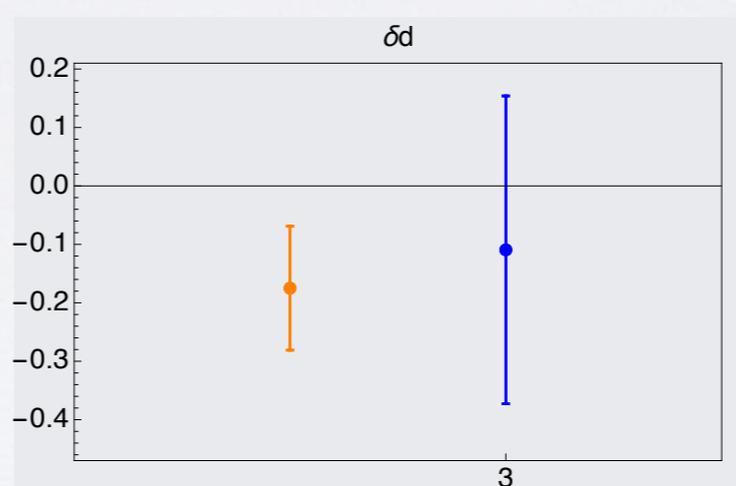
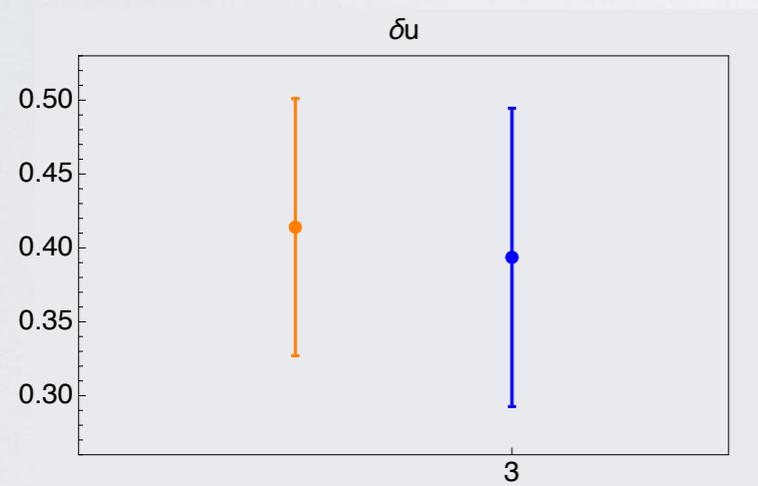
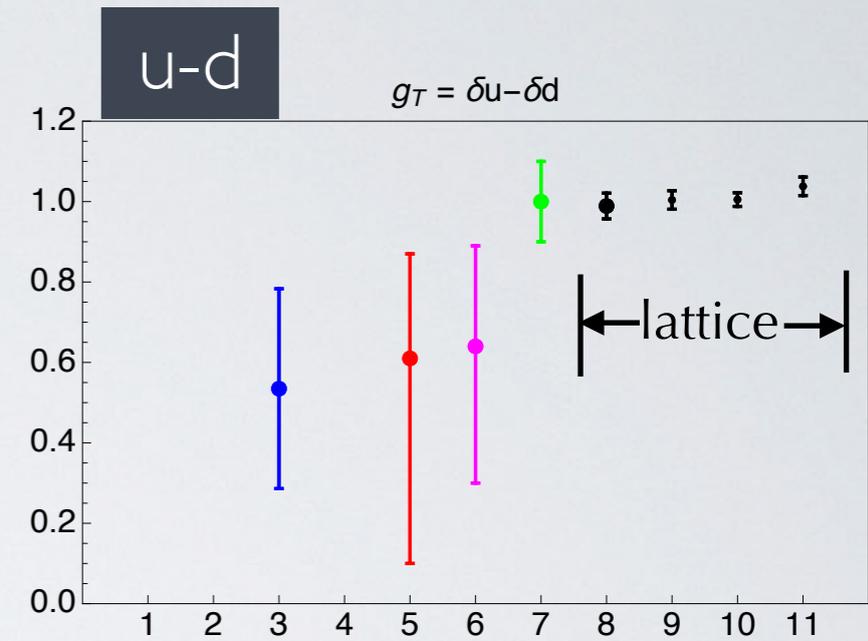
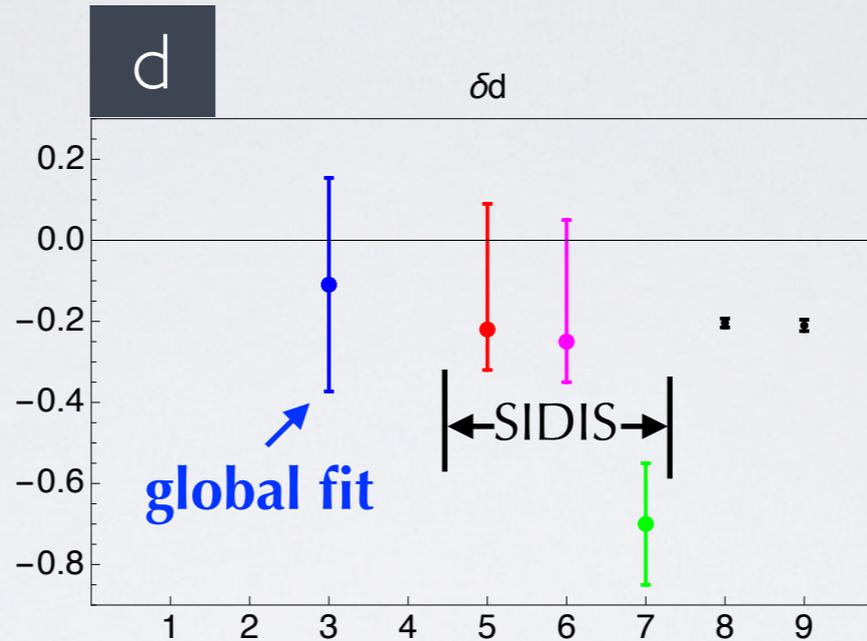
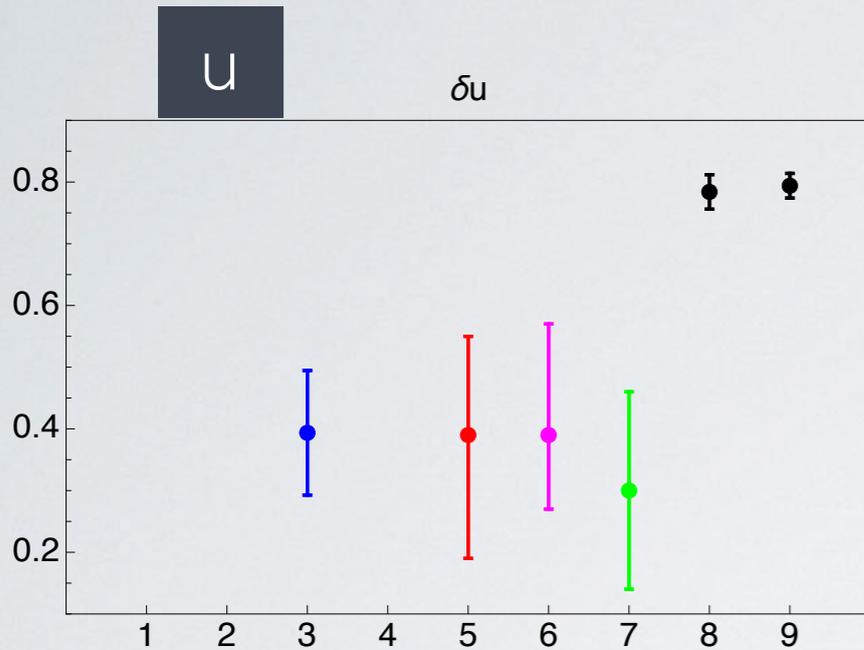
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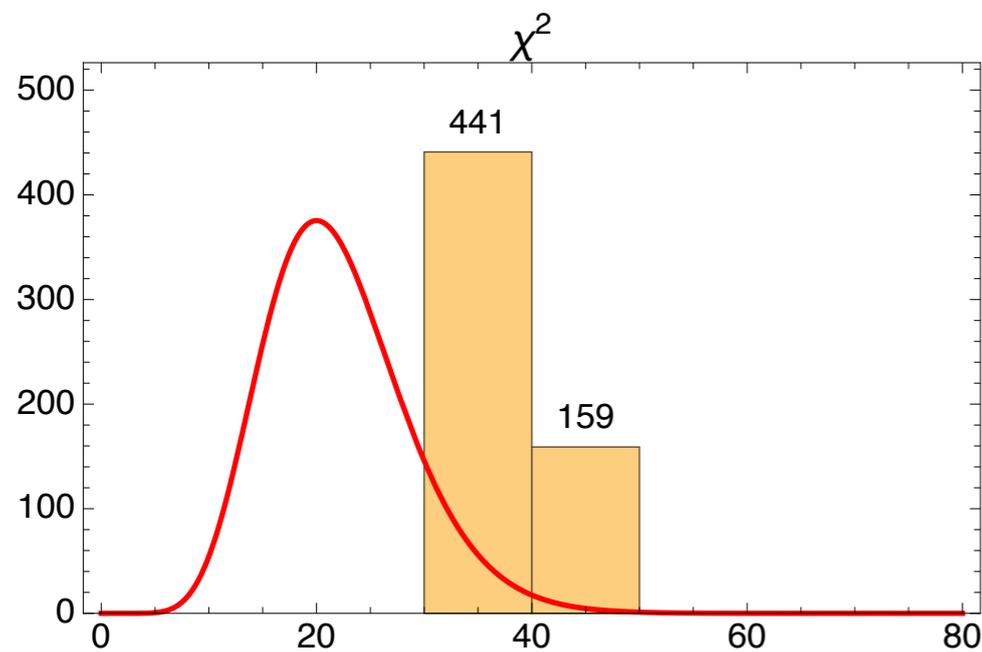
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better precision

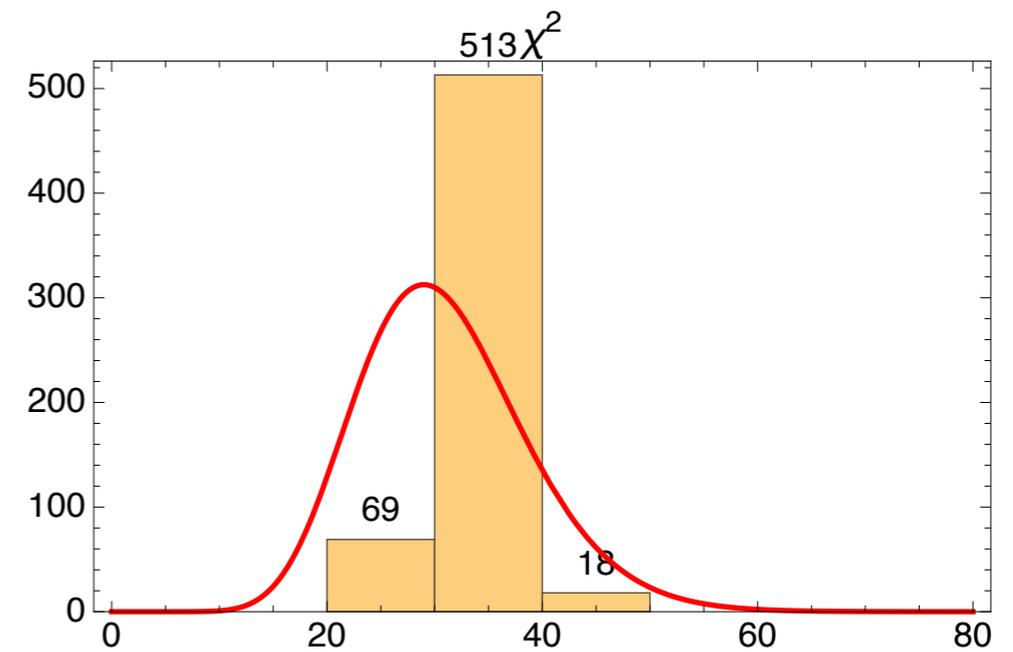
but tension with lattice confirmed

# better $\chi^2$

$$\chi^2/\text{dof} = 1.76 \pm 0.11$$



$$\chi^2/\text{dof} = 1.12 \pm 0.09$$



probability density function of  
 $\chi^2$  distribution for  
22 d.o.f.      31 d.o.f.

but central value of pseudodata not known  
→ only spreading is meaningful

# results

global fit published in

*Radici and Bacchetta, P.R.L. 120 (18) 192001*



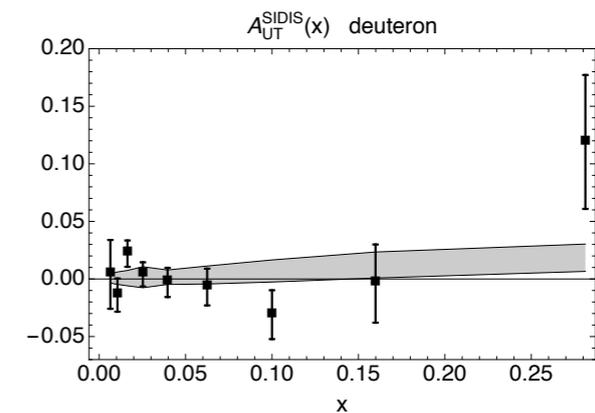
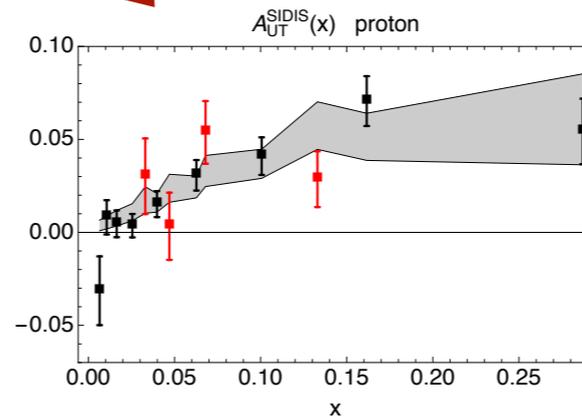
**SIDIS**



*Adolph et al., P.L. B713 (12)*



*Airapetian et al.,  
JHEP 0806 (08) 017*

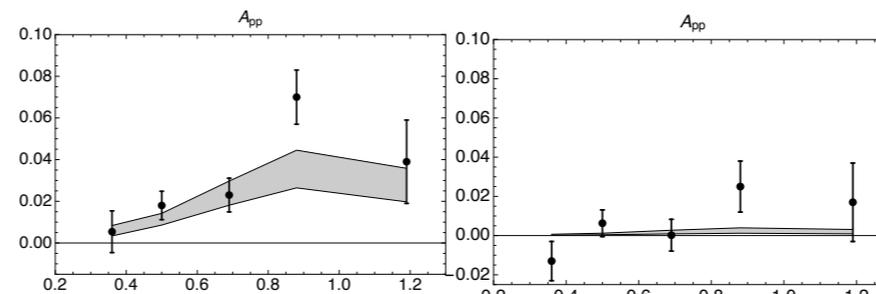


**pp collisions**

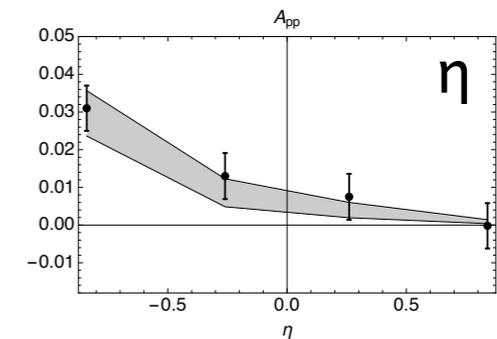


*Adamczyk et al.,  
P.R.L. 115 (2015) 242501*

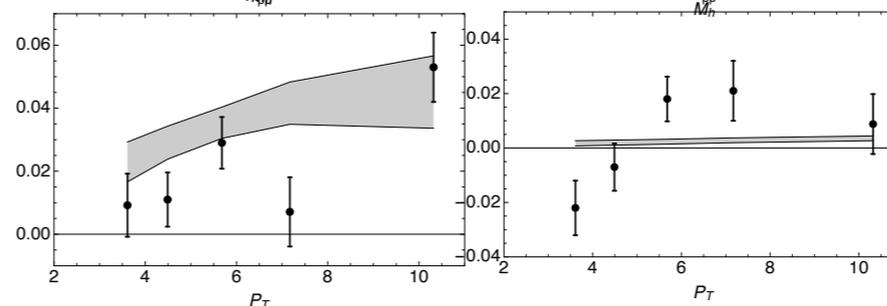
$M_h, \eta < 0$



$M_h, \eta > 0$



$p_T, \eta < 0$



$p_T, \eta > 0$

# CLAS12 pseudo-data

add to previous set of data  
 a new set of SIDIS pseudo-data for **proton** target



*Adolph et al., P.L. B713 (12)*

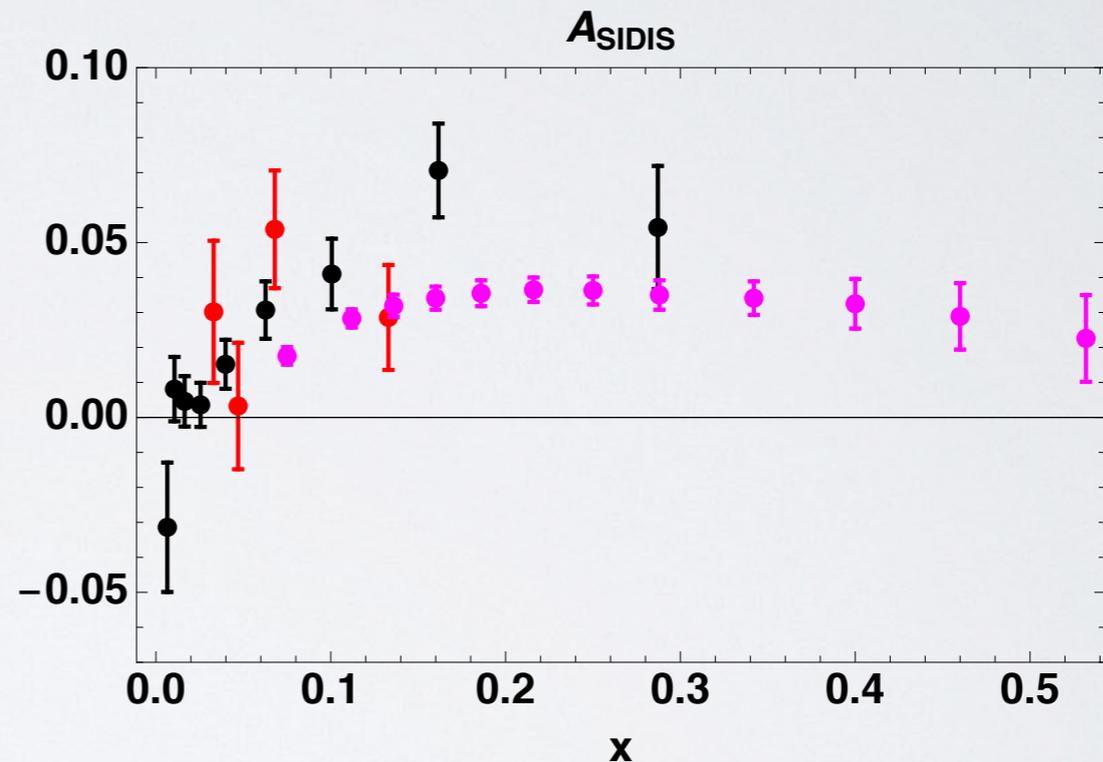


*Airapetian et al., JHEP 0806 (08) 017*

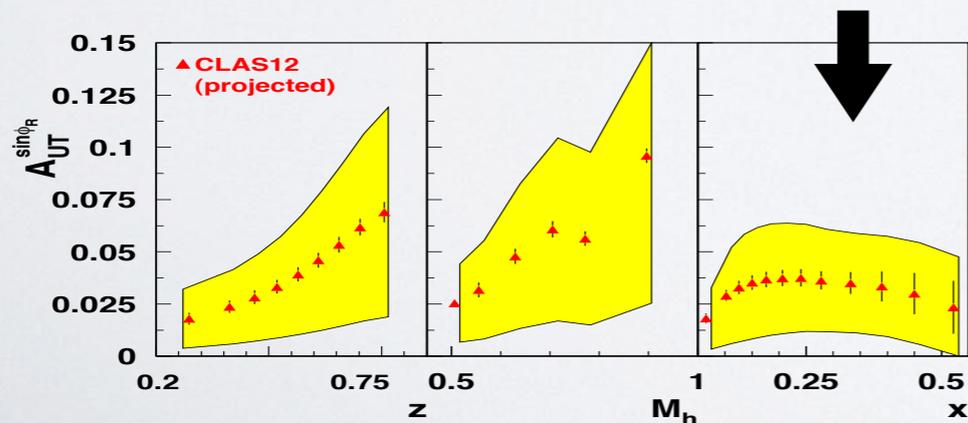


pseudodata C12-12-009

A 12 GeV Research Proposal to Jefferson Lab (PAC 39)



Measurement of transversity with dihadron production in SIDIS with transversely polarized target



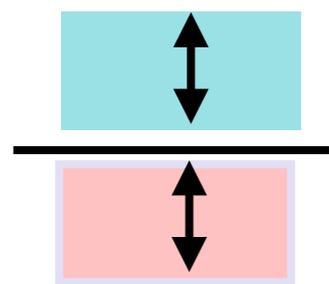
study impact on precision of published global fit

# impact of pseudo-data for proton

global fit + pseudodata

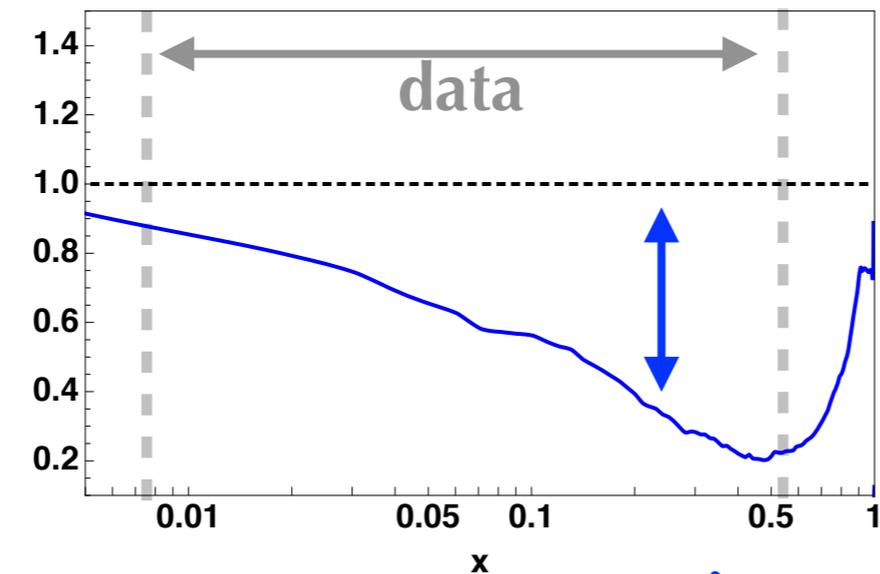
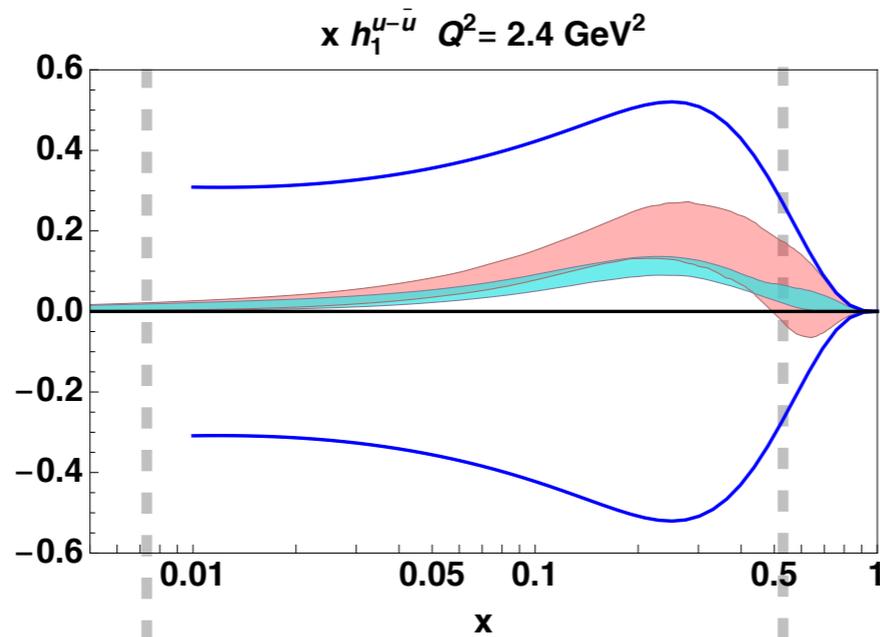
global fit

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_{1^u}^u/4 \\ D_{1^u}^u \end{cases}$$

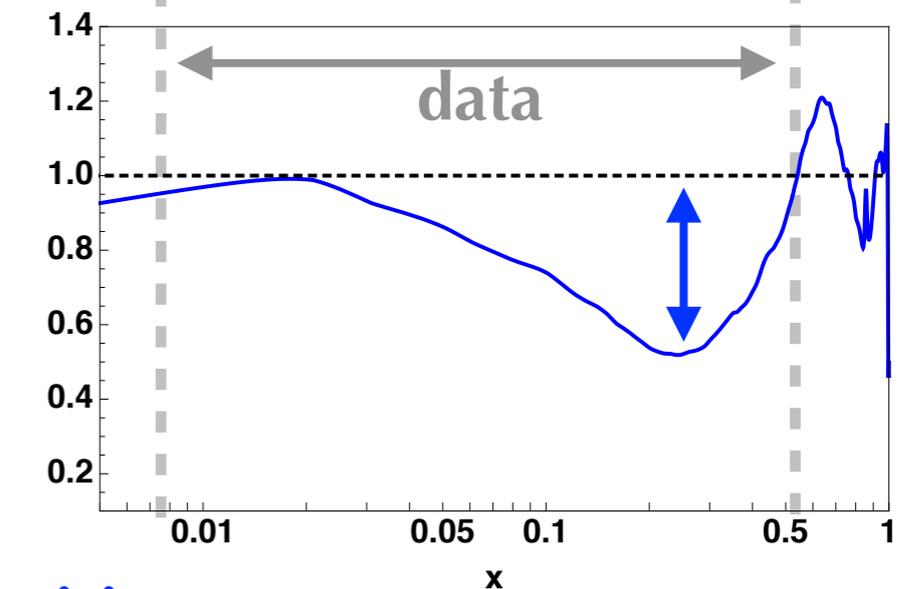
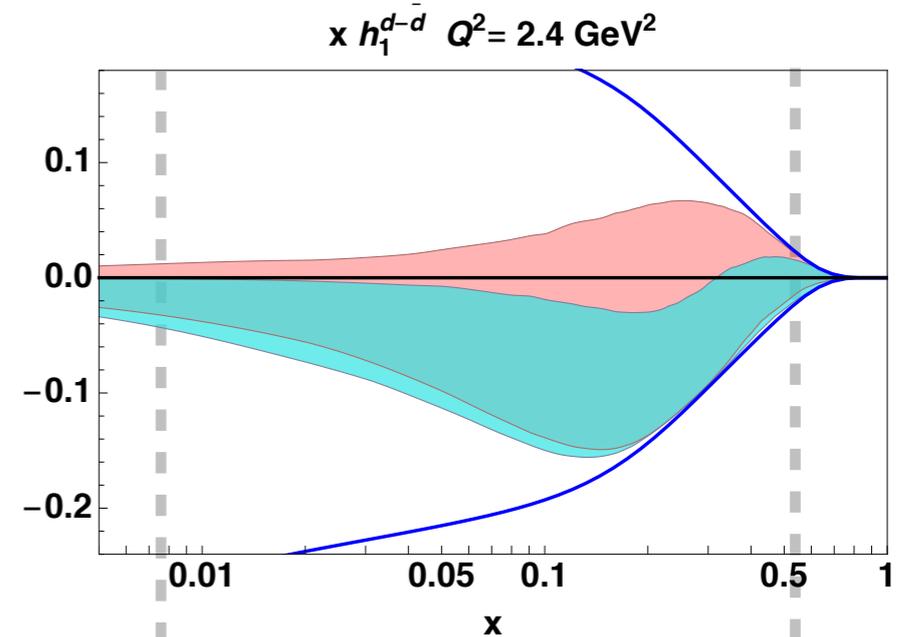


ratio of widths

up



down



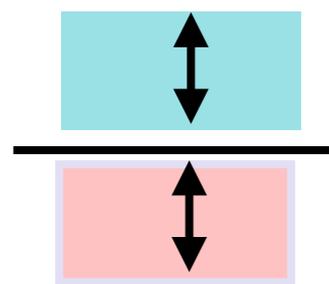
increase precision

# linear scale

global fit + pseudodata

global fit

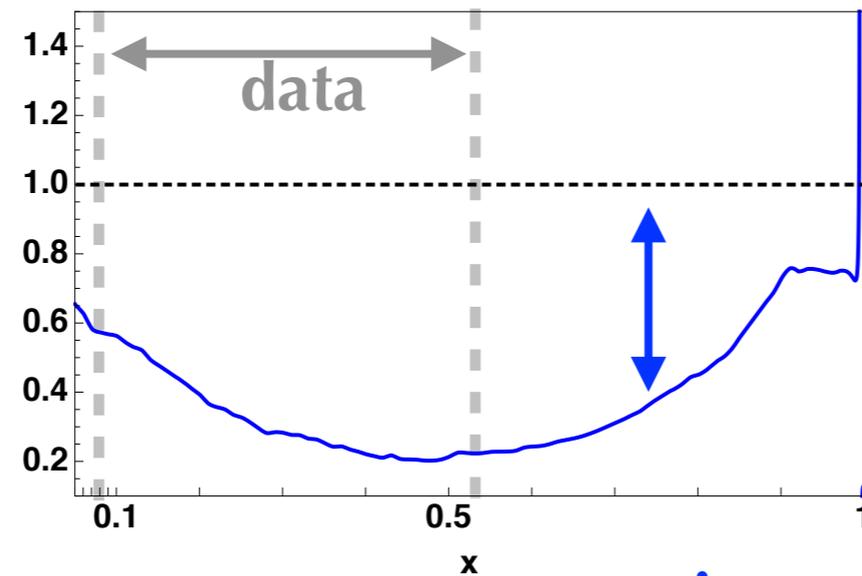
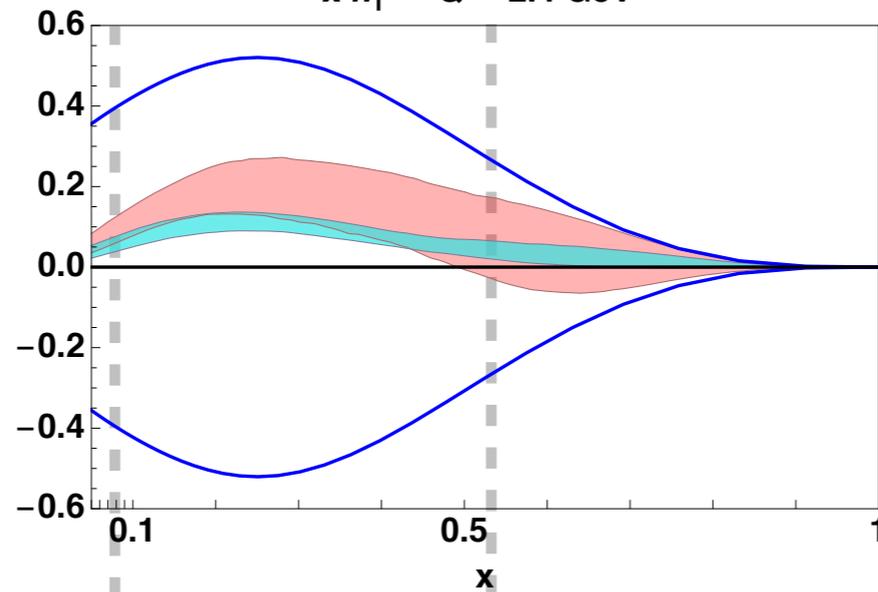
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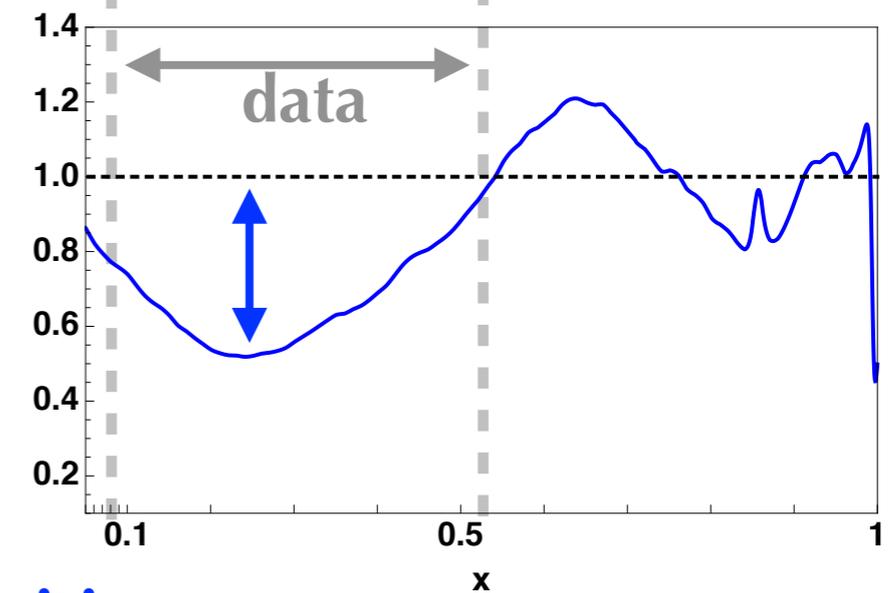
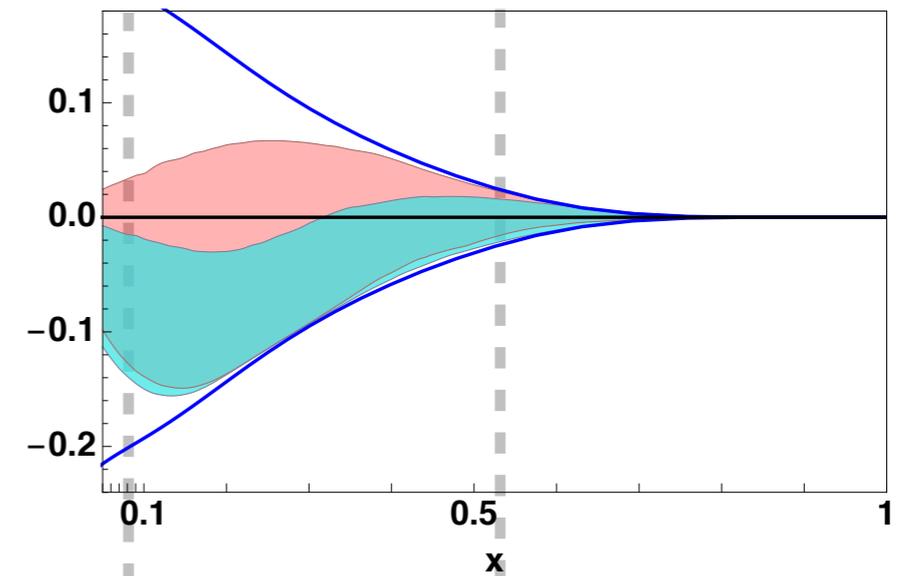
up

$x h_1^{u-\bar{u}} Q^2 = 2.4 \text{ GeV}^2$



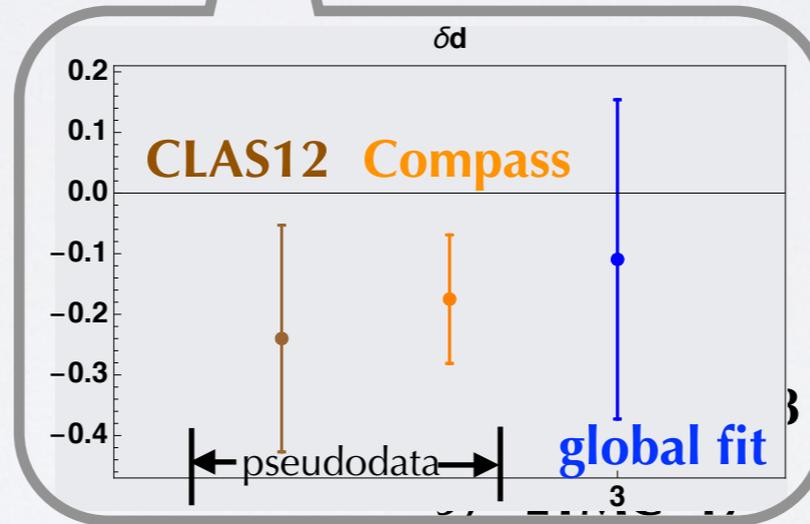
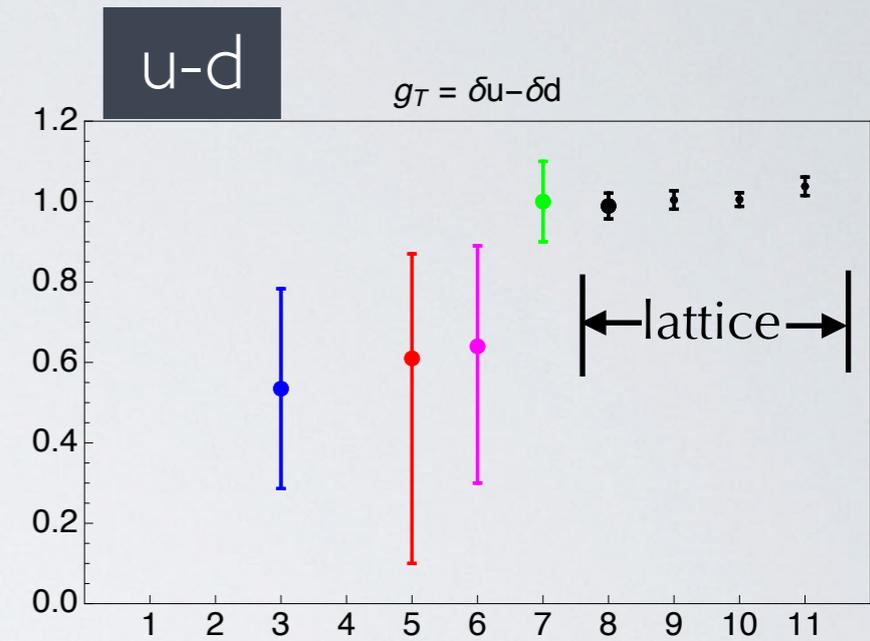
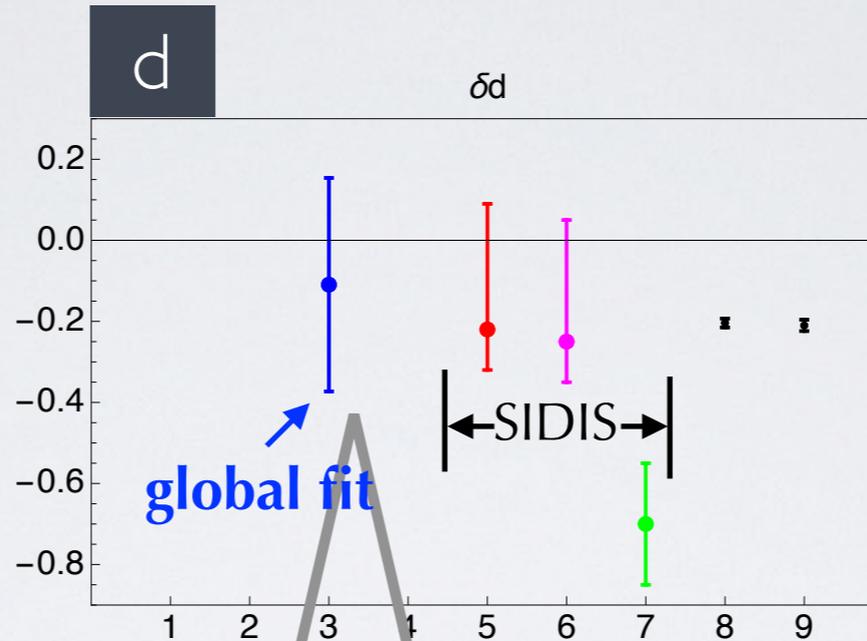
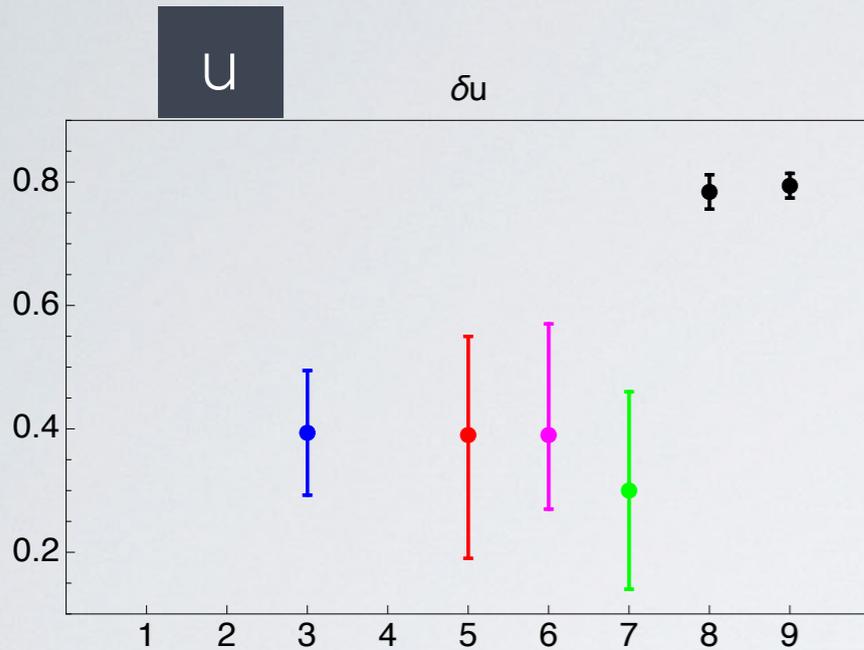
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increase precision

# tensor charge



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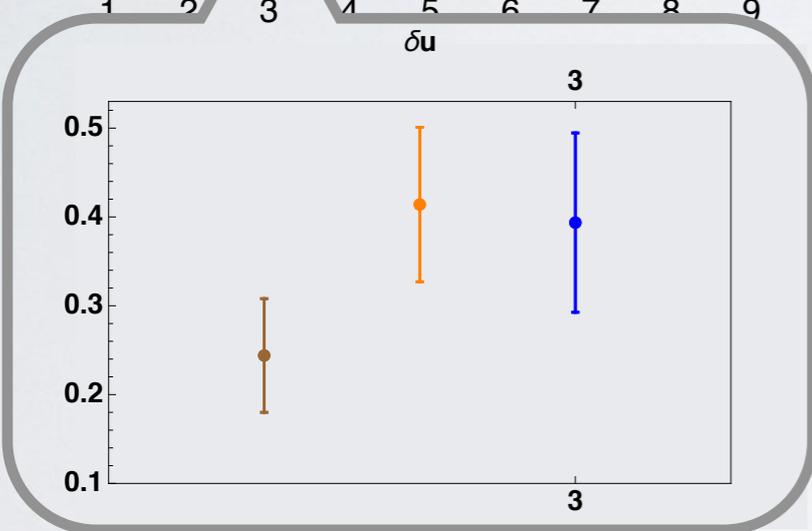
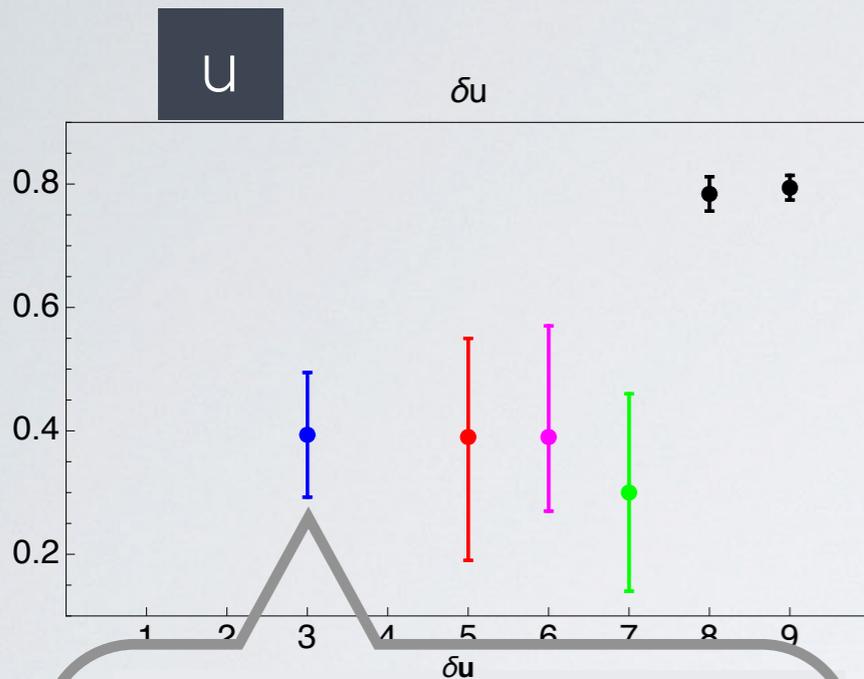
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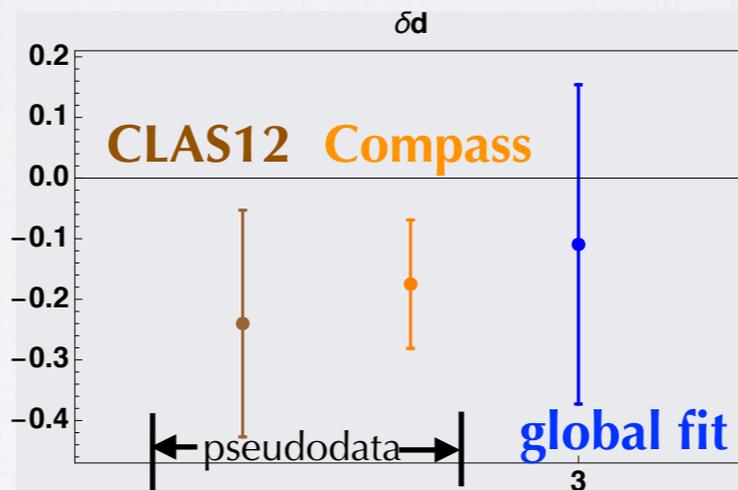
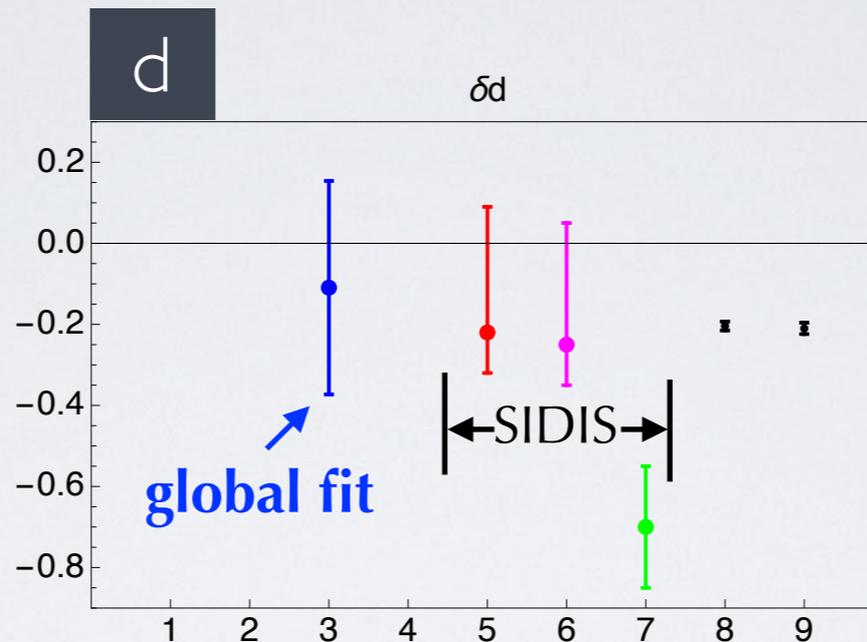
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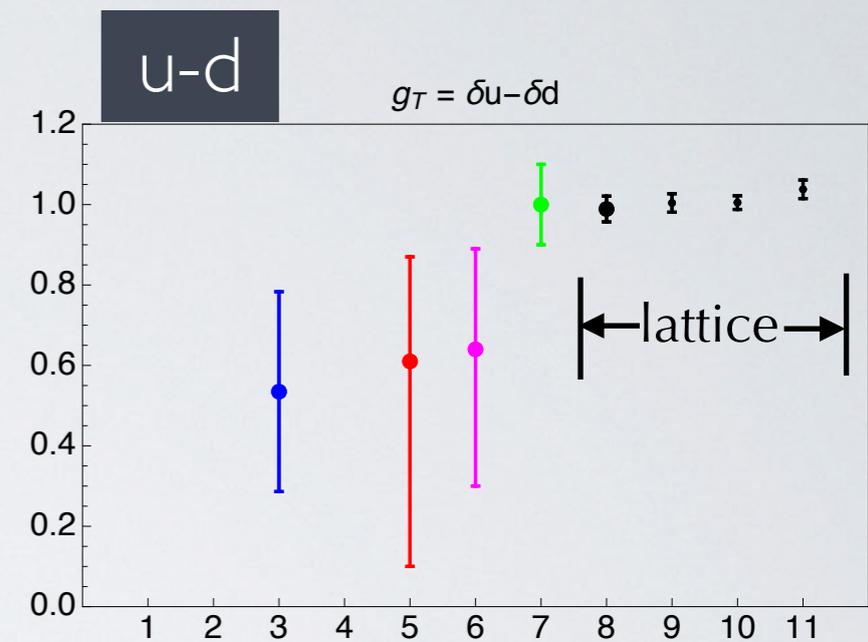
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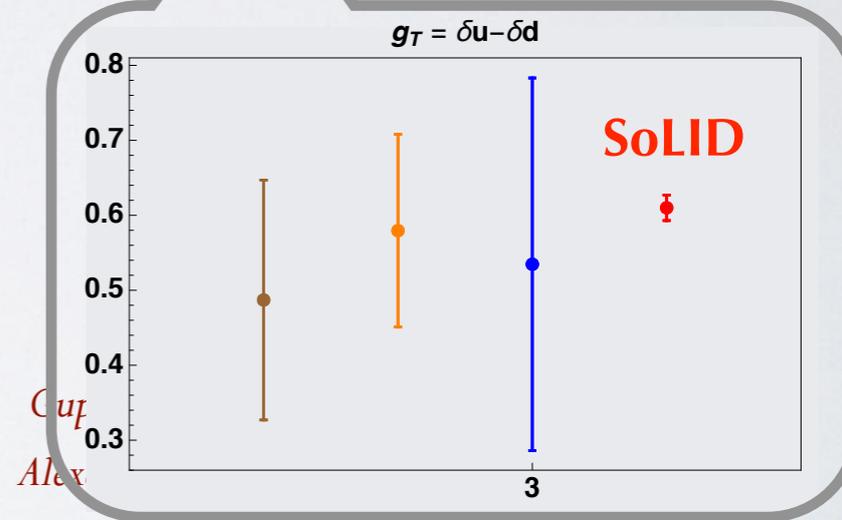
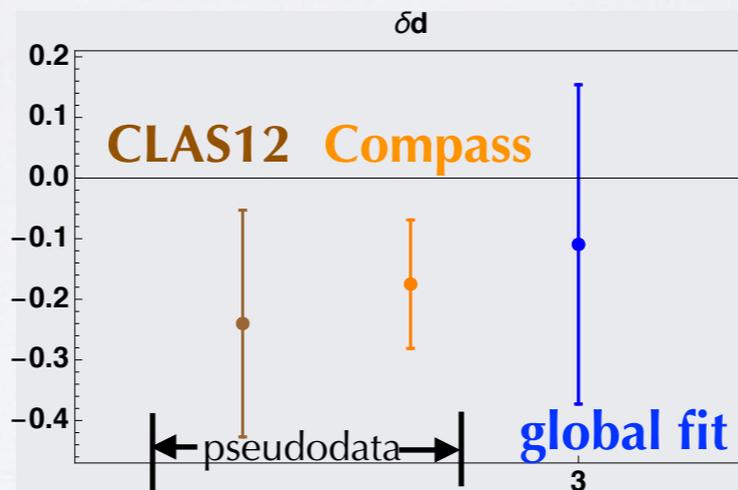
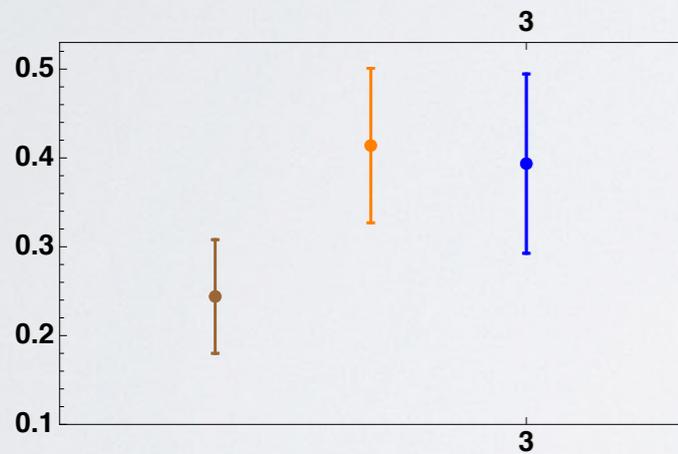
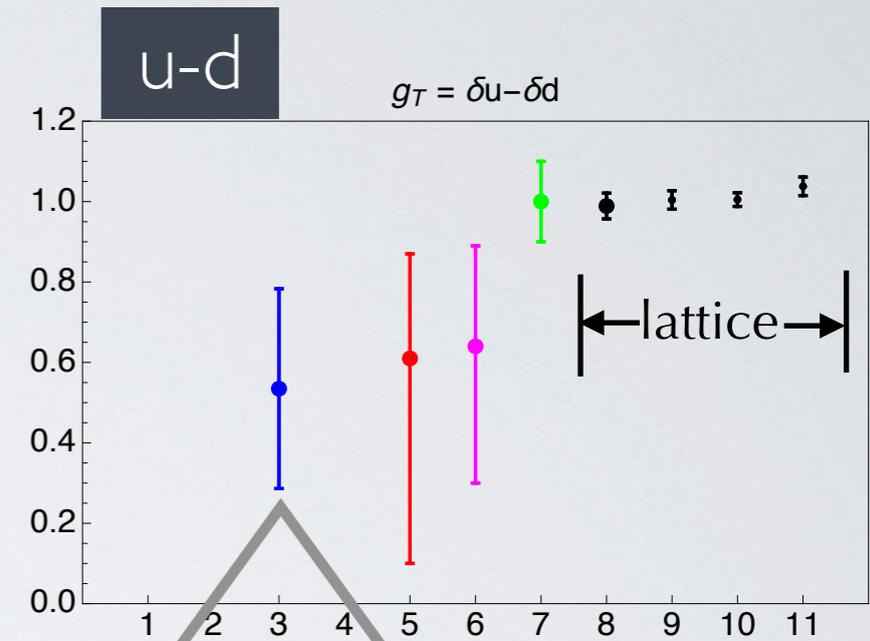
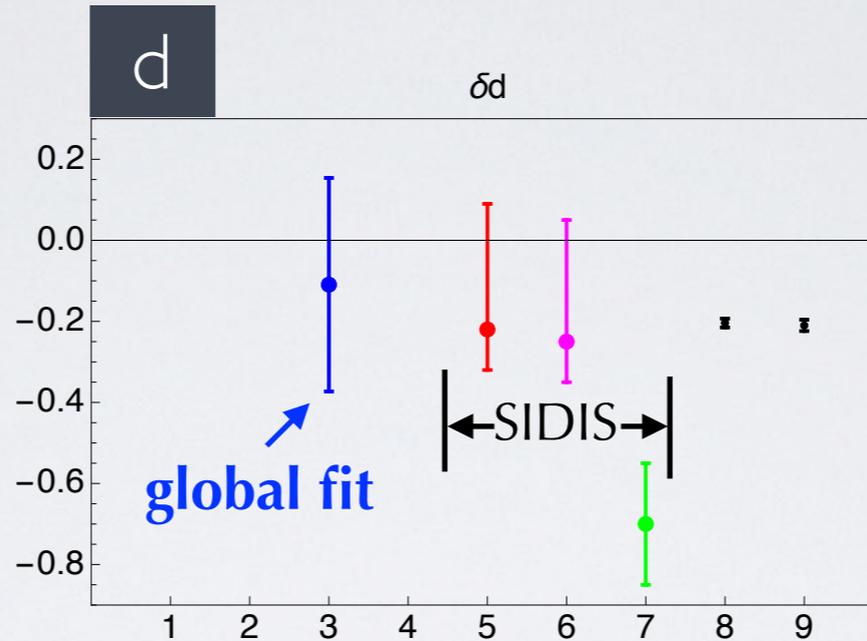
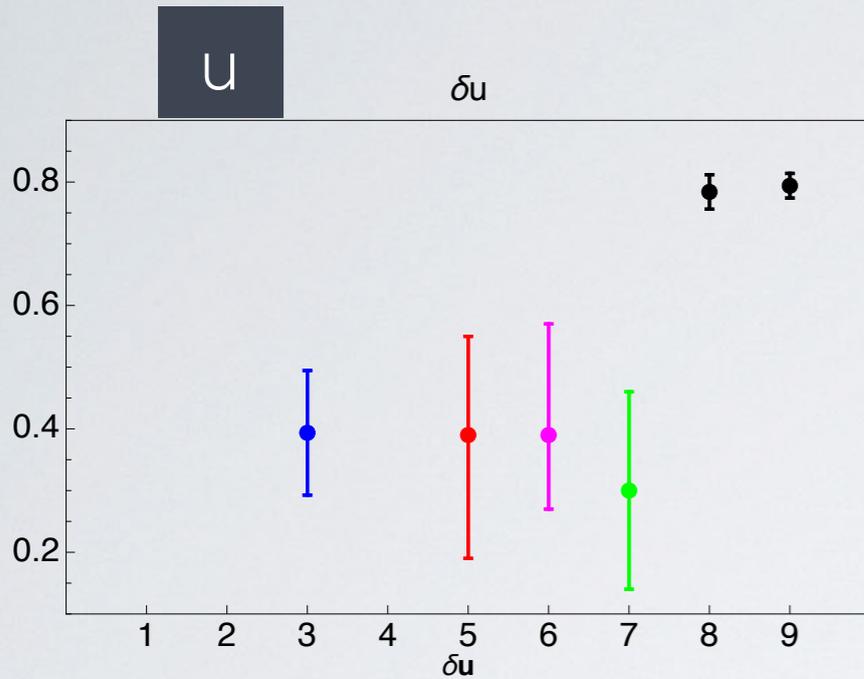
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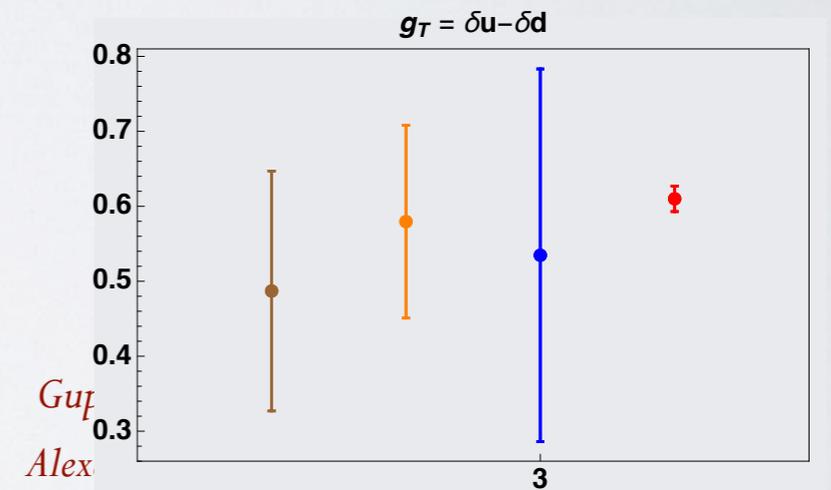
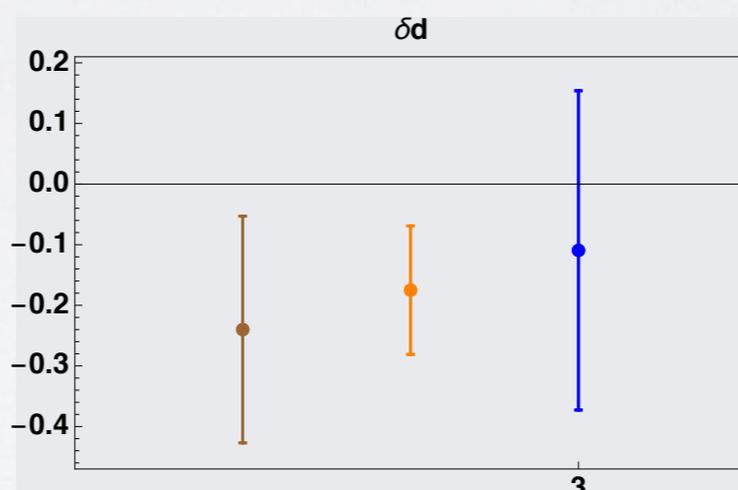
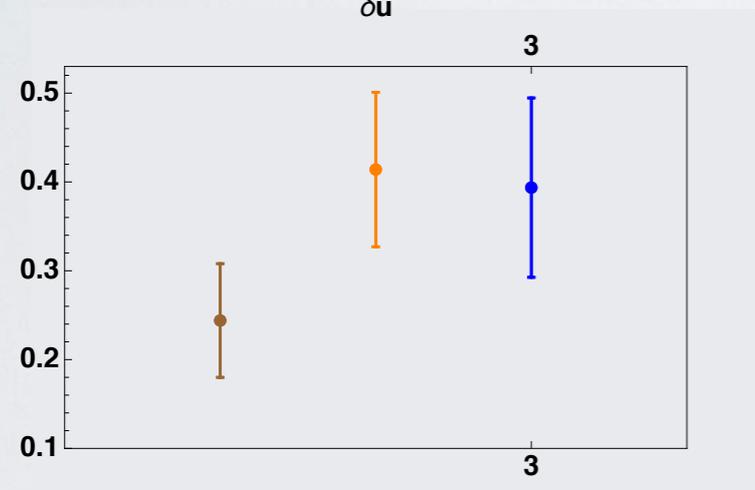
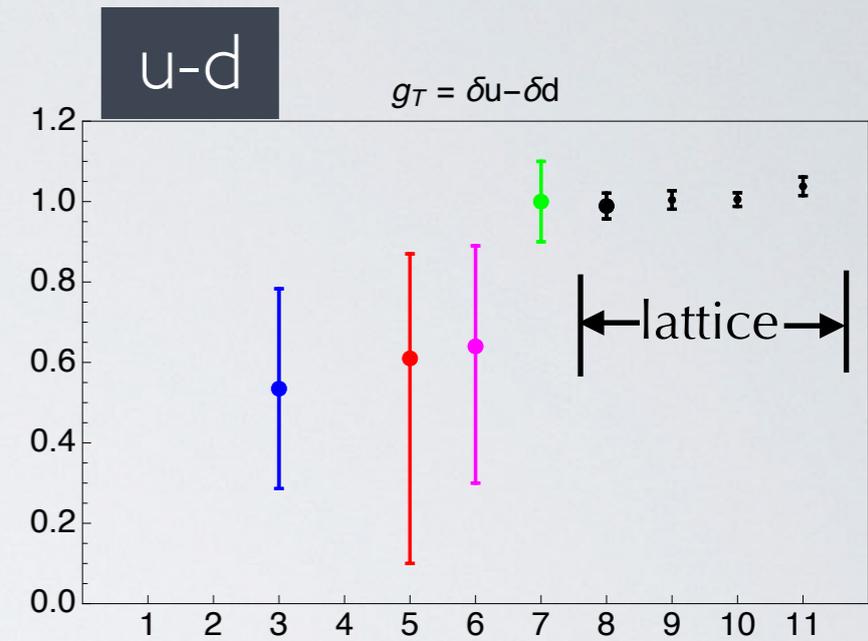
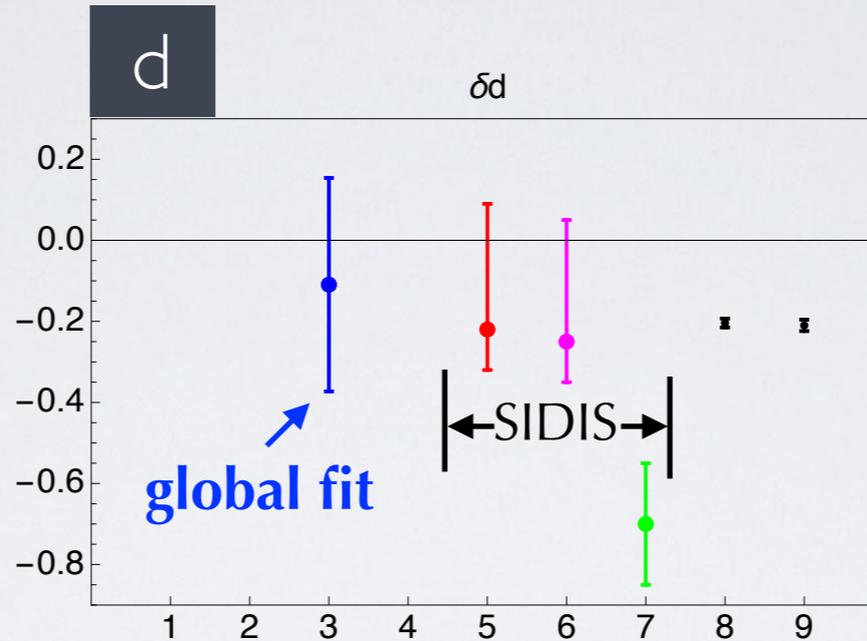
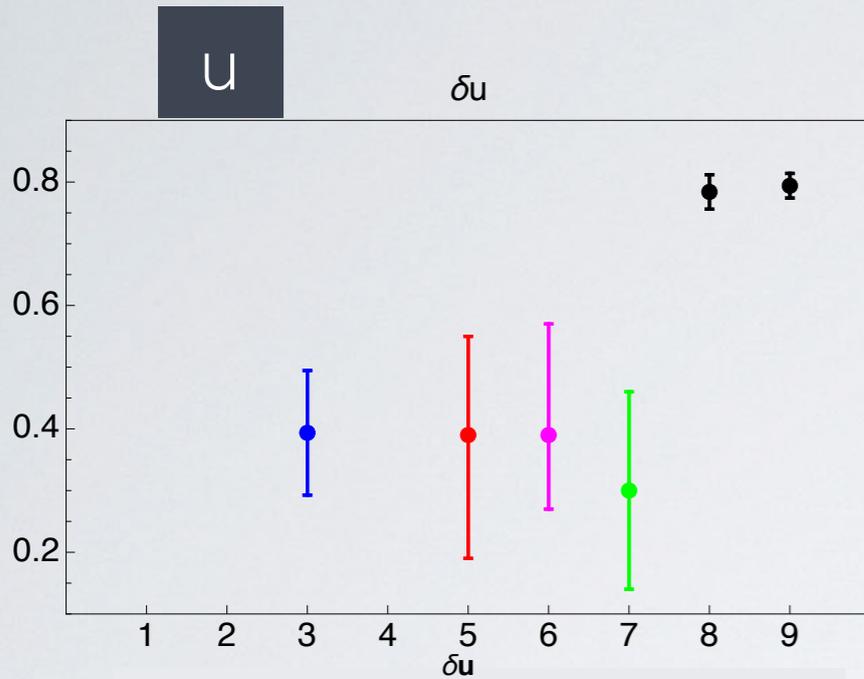
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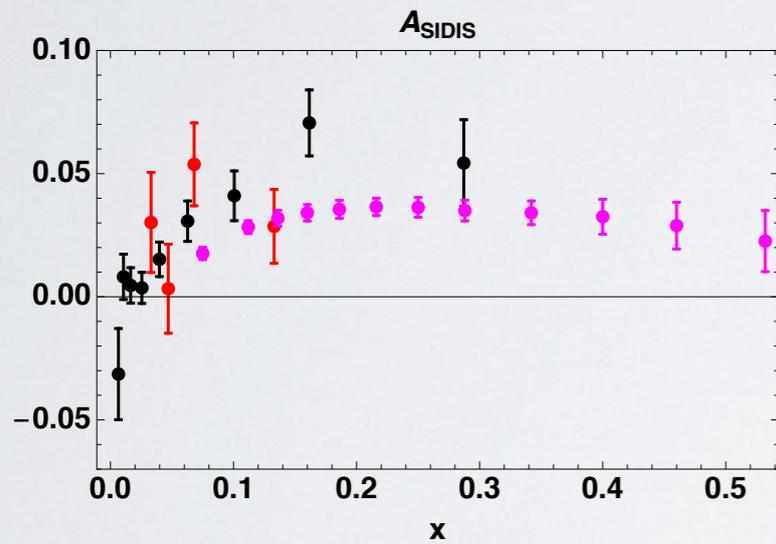
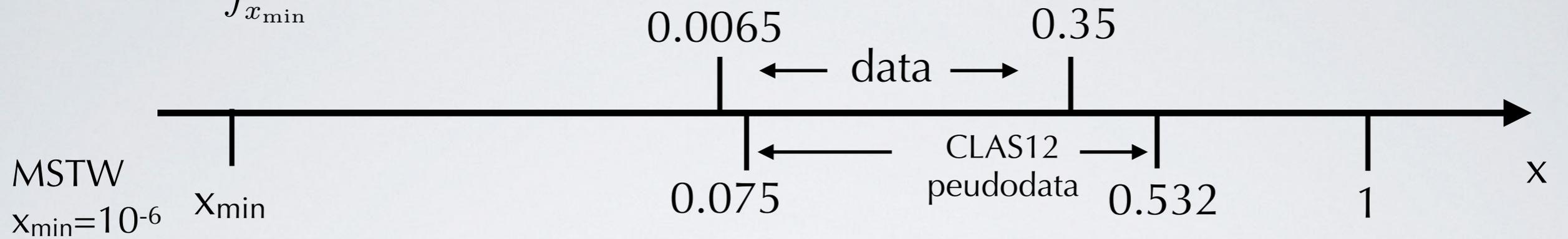
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but tension with lattice confirmed

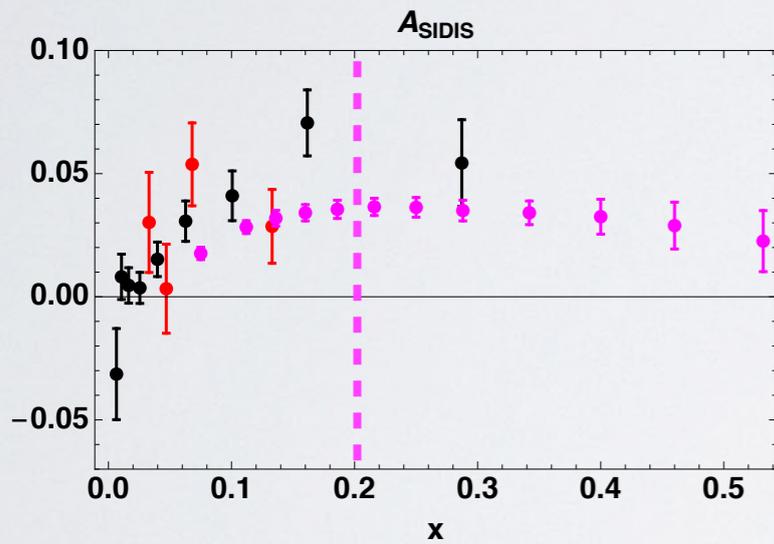
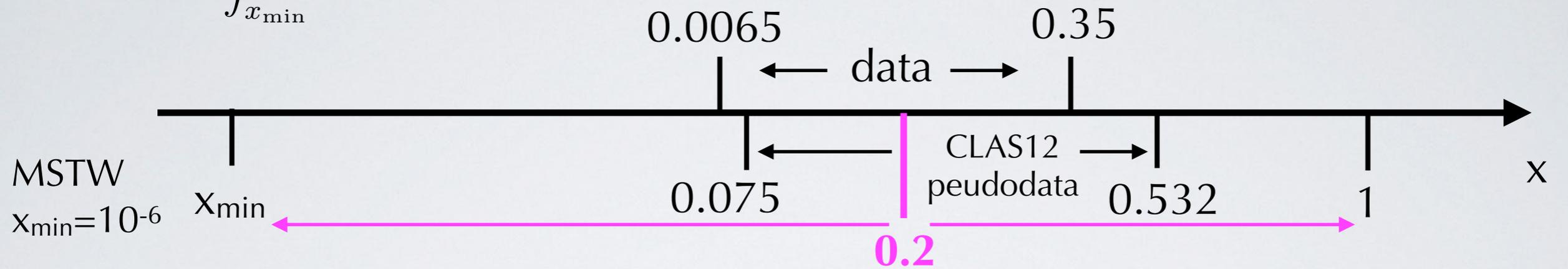
# break down of Mellin moment

$$\delta q(Q^2) = \int_{x_{\min}}^1 dx h_1^{q-\bar{q}}(x, Q^2)$$



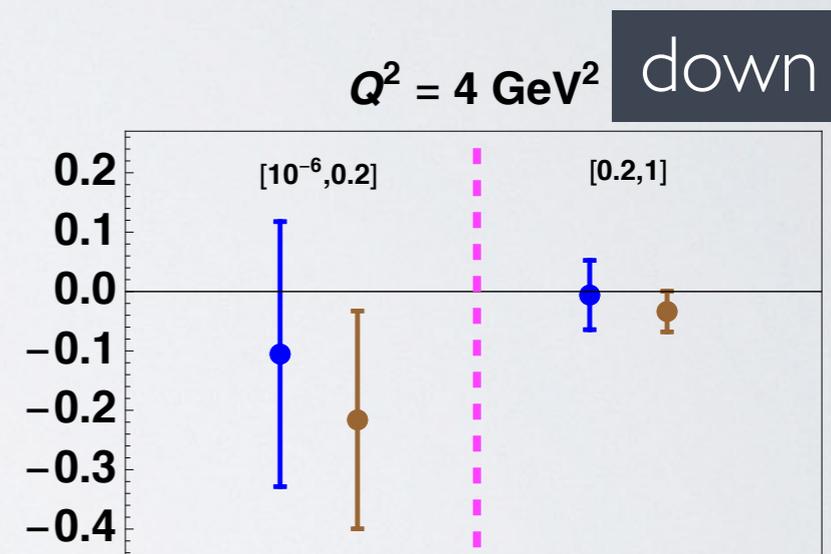
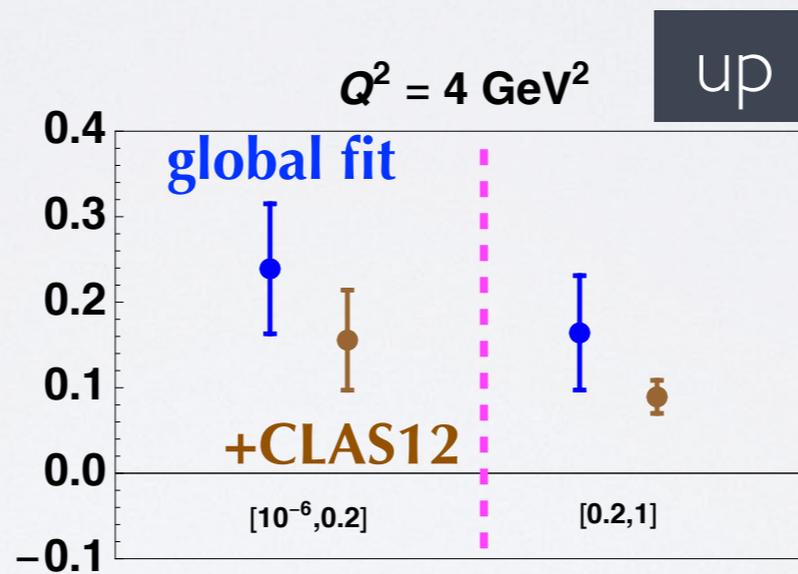
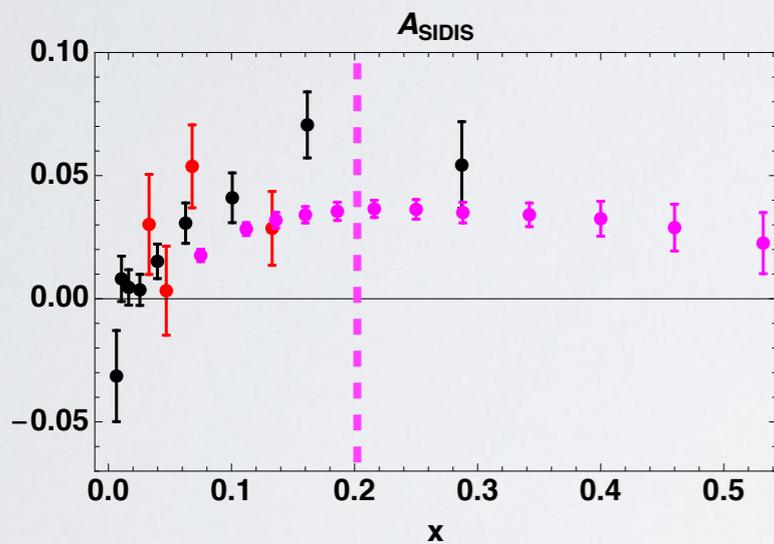
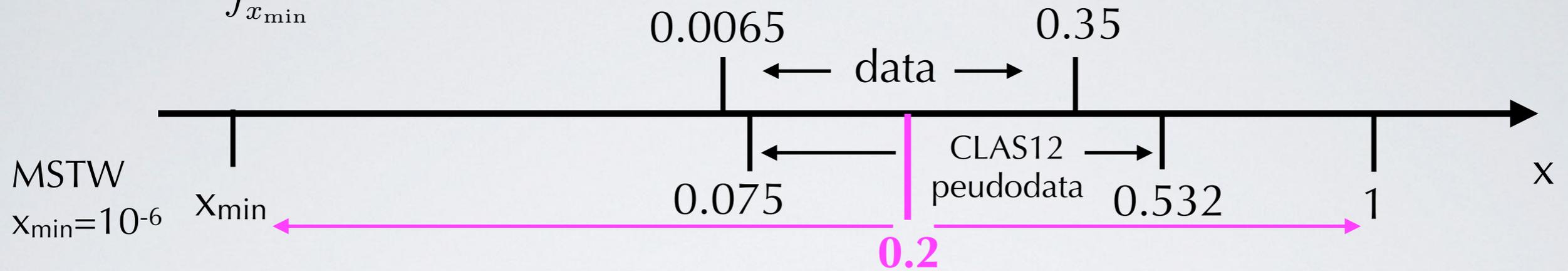
# break down of Mellin moment

$$\delta q(Q^2) = \int_{x_{\min}}^1 dx h_1^{q-\bar{q}}(x, Q^2)$$



# break down of Mellin moment

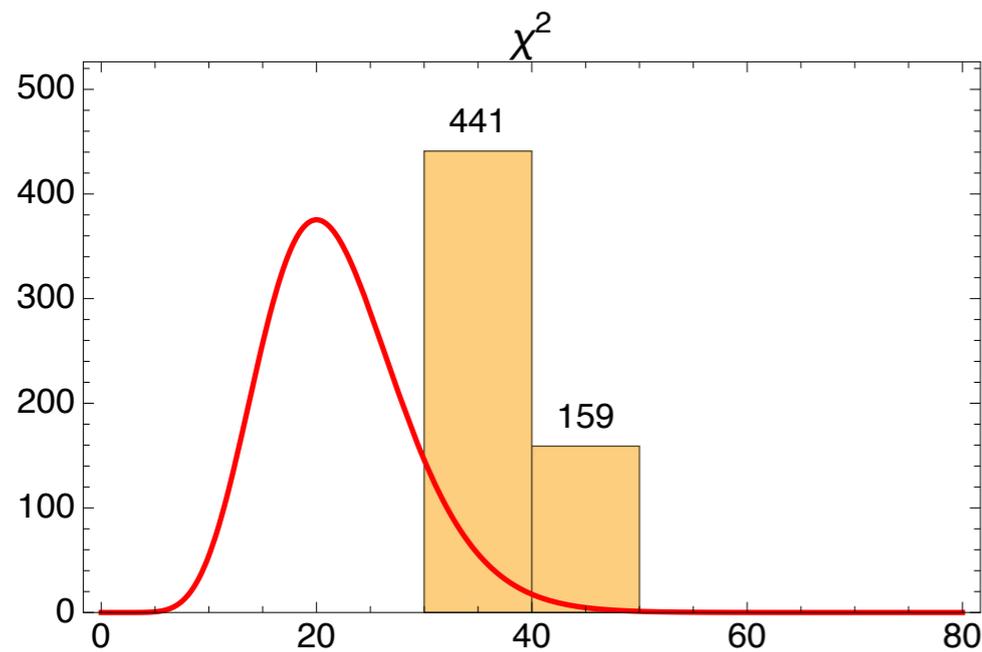
$$\delta q(Q^2) = \int_{x_{\min}}^1 dx h_1^{q-\bar{q}}(x, Q^2)$$



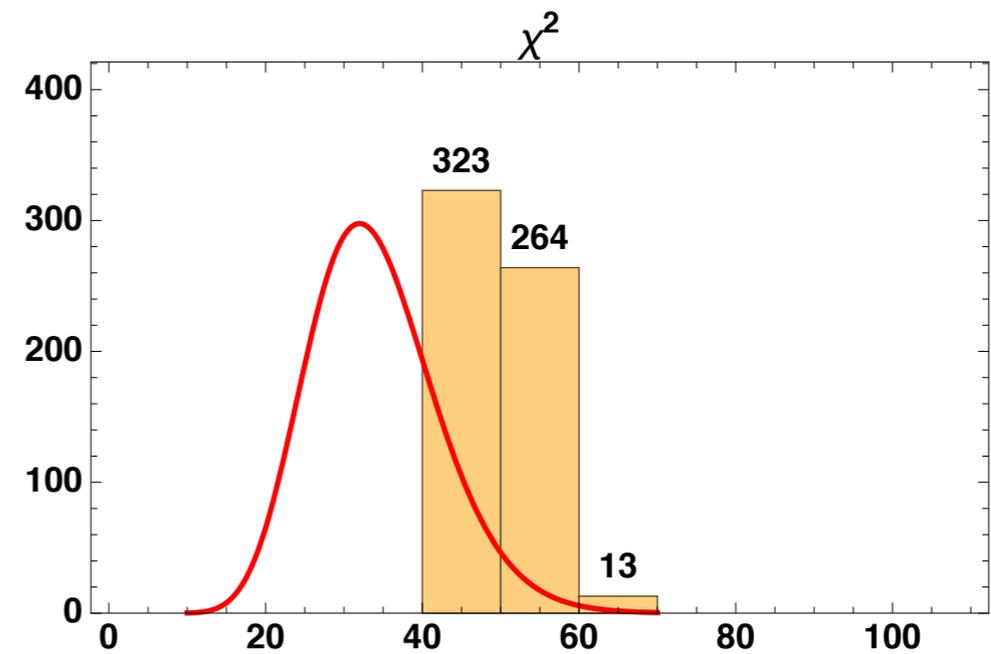
impact of CLAS12 pseudodata at large  $x$  ( $>0.2$ )  
 gives  $\sim 50\%$  of up tensor charge  
 relative error  $\Delta g_T/g_T$  from 82%  $\rightarrow$  43%

# better $\chi^2$

$$\chi^2/\text{dof} = 1.76 \pm 0.11$$



$$\chi^2/\text{dof} = 1.48 \pm 0.10$$

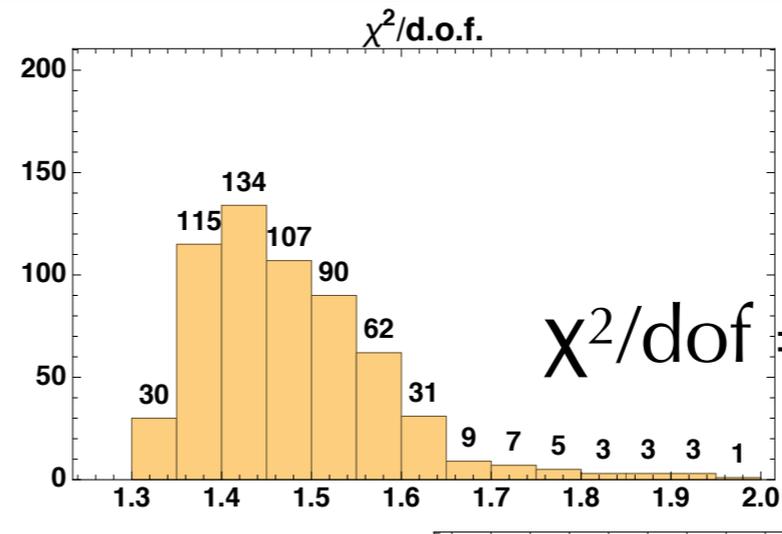
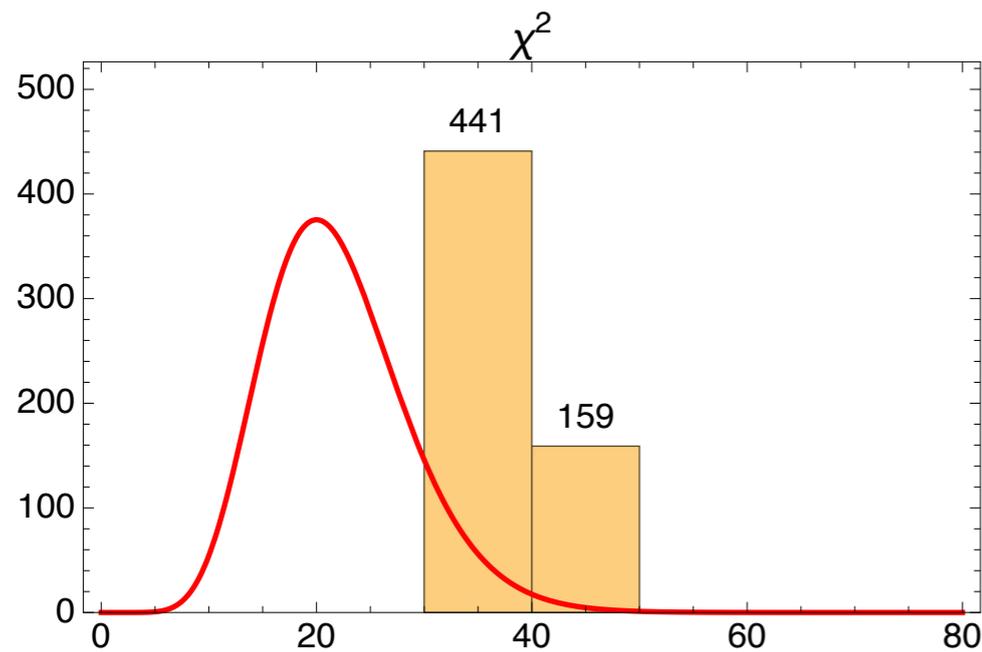


probability density function of  
 $\chi^2$  distribution for  
22 d.o.f.      34 d.o.f.

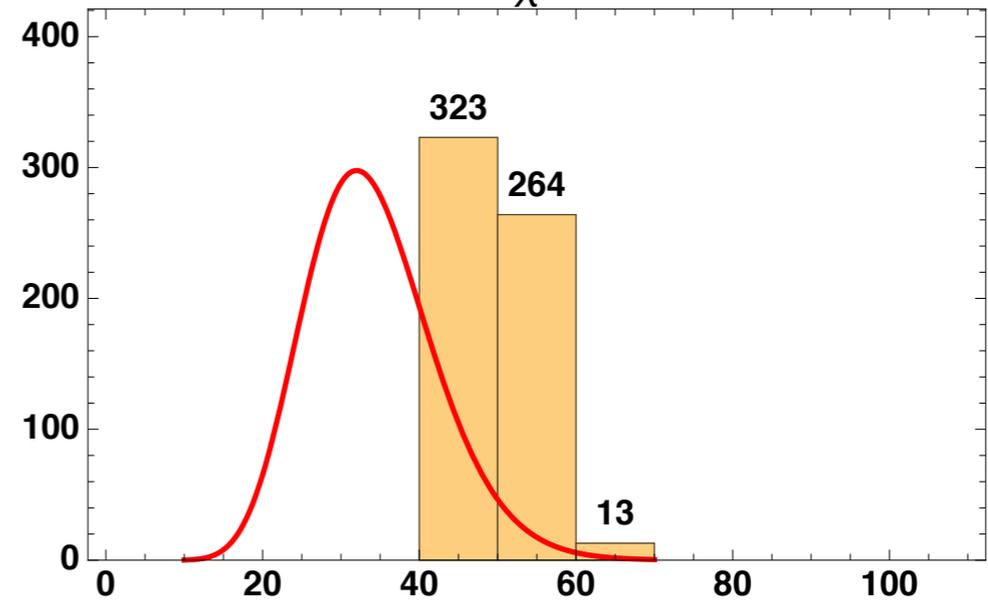
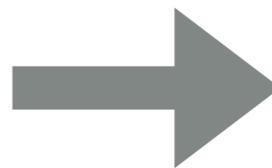
but central value of pseudodata not known  
→ only spreading is meaningful

# better $\chi^2$

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probability density function of  
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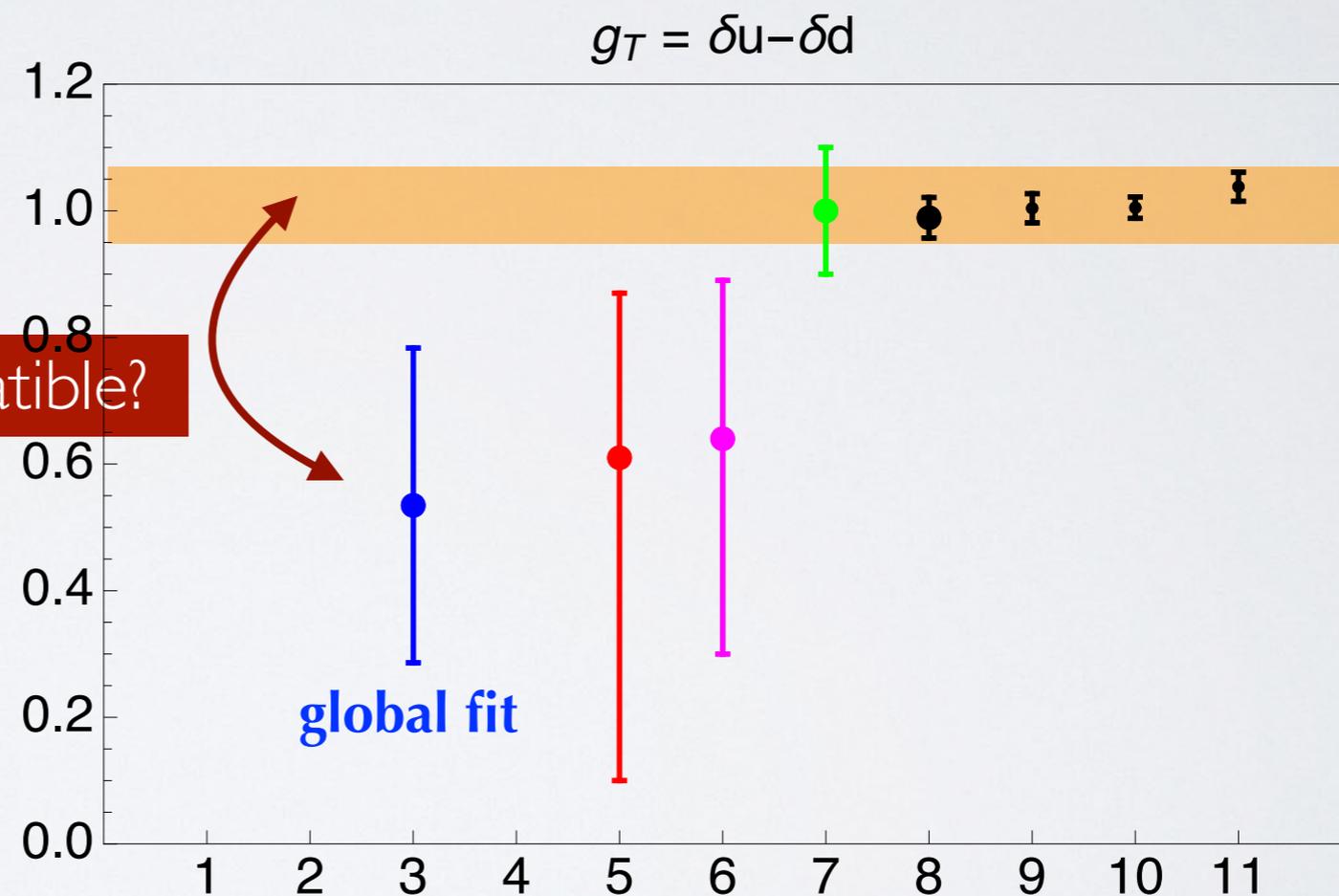
but central value of pseudodata not known  
→ only spreading is meaningful

# compatibility with lattice

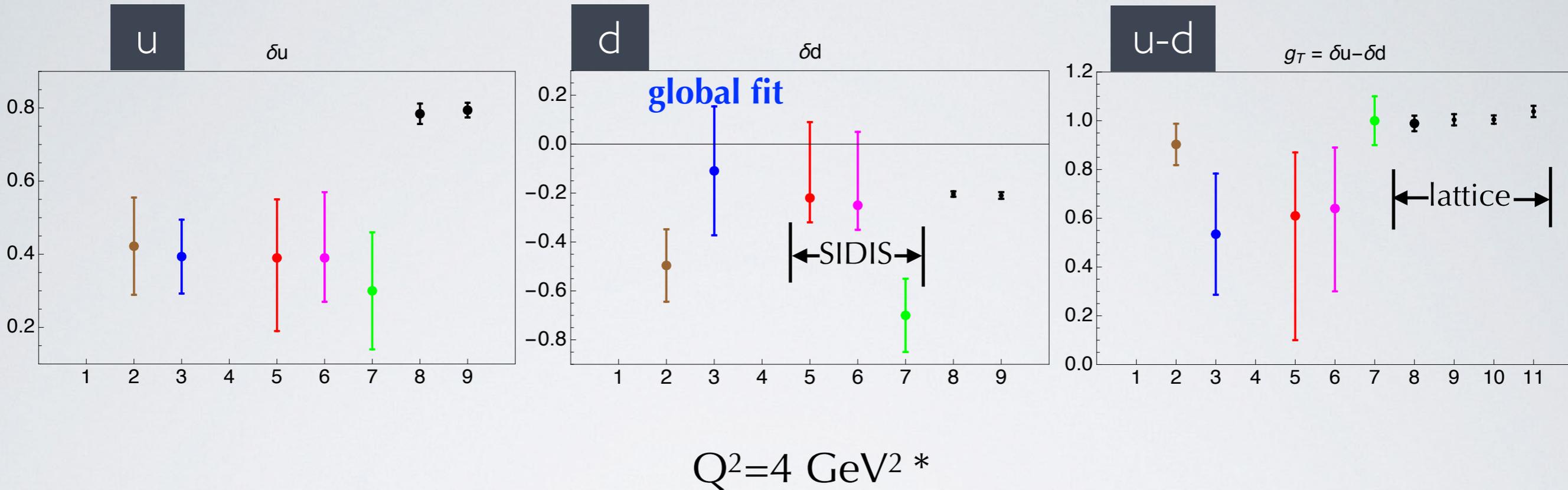
add to SIDIS+pp data  
constraint to reproduce  $g_T$  from lattice

$$\overline{g_T^{\text{latt}}} = 1.004 \pm 0.057$$

are they compatible?



# tensor charge



## 2) global fit + constrain $g_T$

Radici & Bacchetta,  
*P.R.L. 120 (18) 192001*

3) **global fit '17**

Kang et al., *P.R. D93 (16) 014009*

5) **"TMD fit" \*  $Q^2=10$**

Anselmino et al., *P.R. D87 (13) 094019*

6) **Torino fit \*  $Q^2=1$**

Lin et al., *P.R.L. 120 (18) 152502*

7) **JAM fit '17 \*  $Q_0^2=2$**

8) **PNDME '18**

*Gupta et al., P.R. D98 (18) 034503*

9) **ETMC '17**

*Alexandrou et al., P.R. D95 (17) 114514;  
E P.R. D96 (17) 099906*

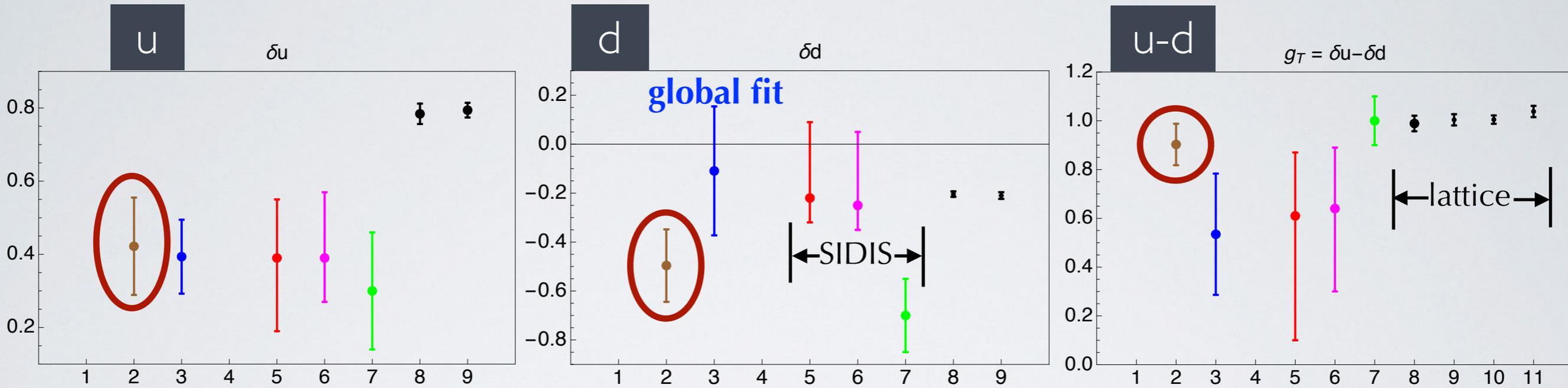
10) **RQCD '14**

*Bali et al., P.R. D91 (15)*

11) **LHPC '12**

*Green et al., P.R. D86 (12)*

# tensor charge



$Q^2=4 \text{ GeV}^2$  \*

not yet full compatibility

## 2) global fit + constrain $g_T$

Radici & Bacchetta,  
*P.R.L.* 120 (18) 192001

3) **global fit '17**

Kang et al., *P.R.* D93 (16) 014009

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*Bali et al., P.R. D91 (15)*

11) **LHPC '12**

*Green et al., P.R. D86 (12)*

# impact of lattice $g_T$ constraint

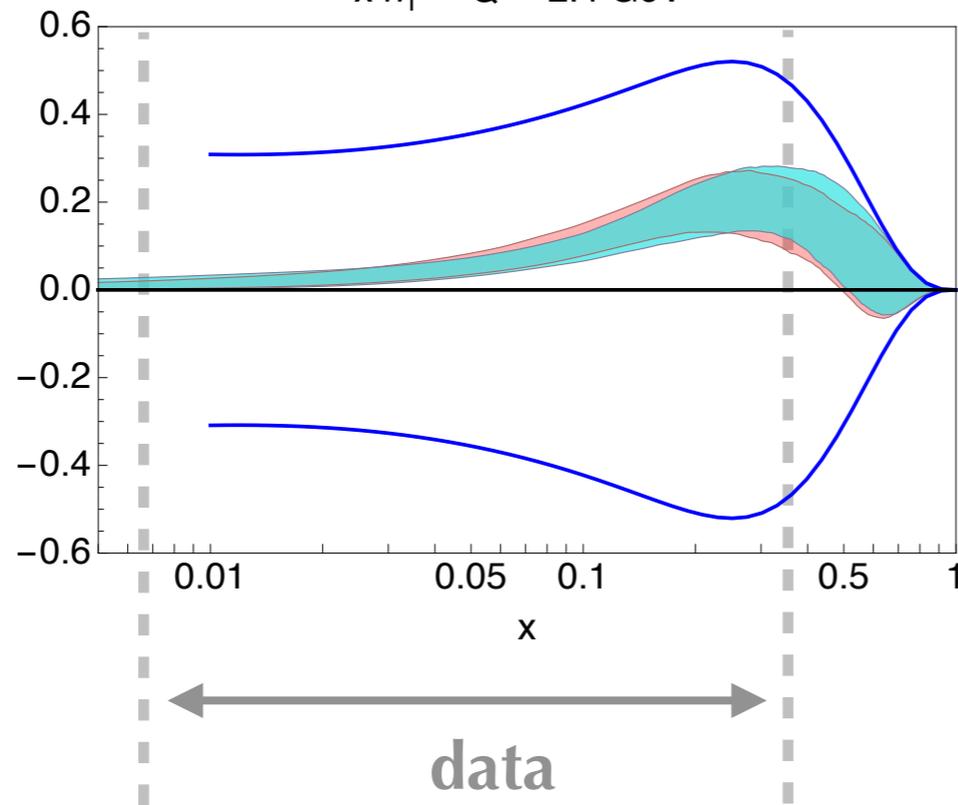
global fit + lattice  $g_T$  constraint

global fit

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_{1^u}^u/4 \\ D_{1^u}^u \end{cases}$$

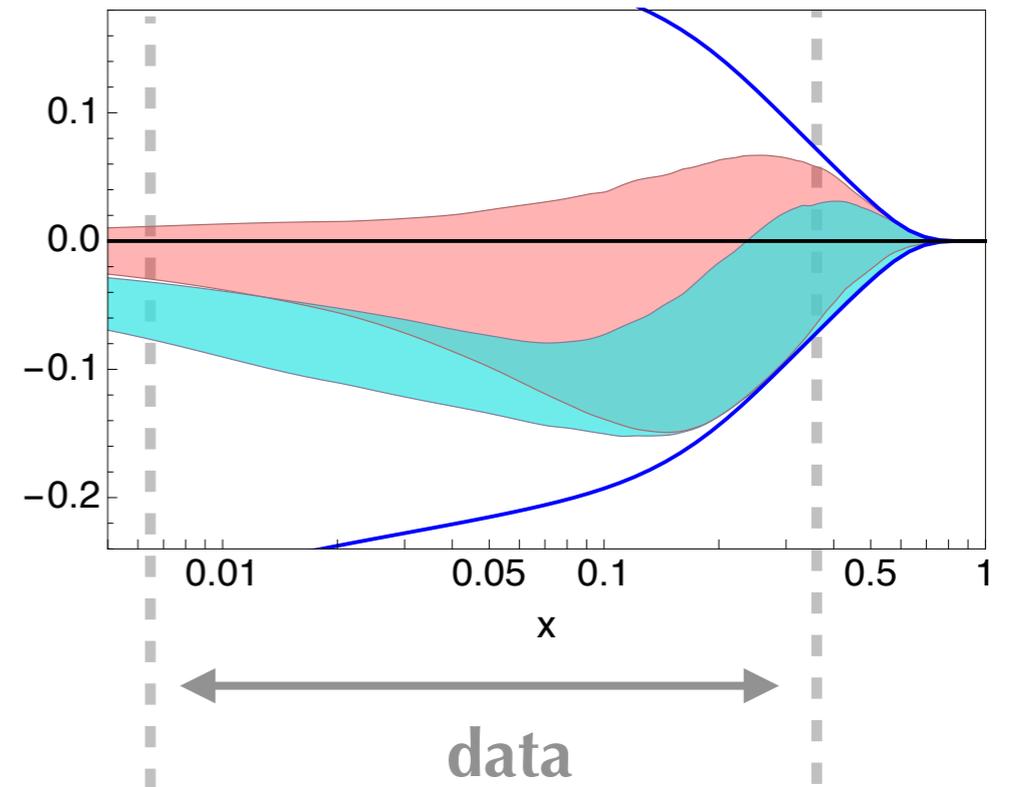
up

$x h_1^{u-\bar{u}} Q^2 = 2.4 \text{ GeV}^2$



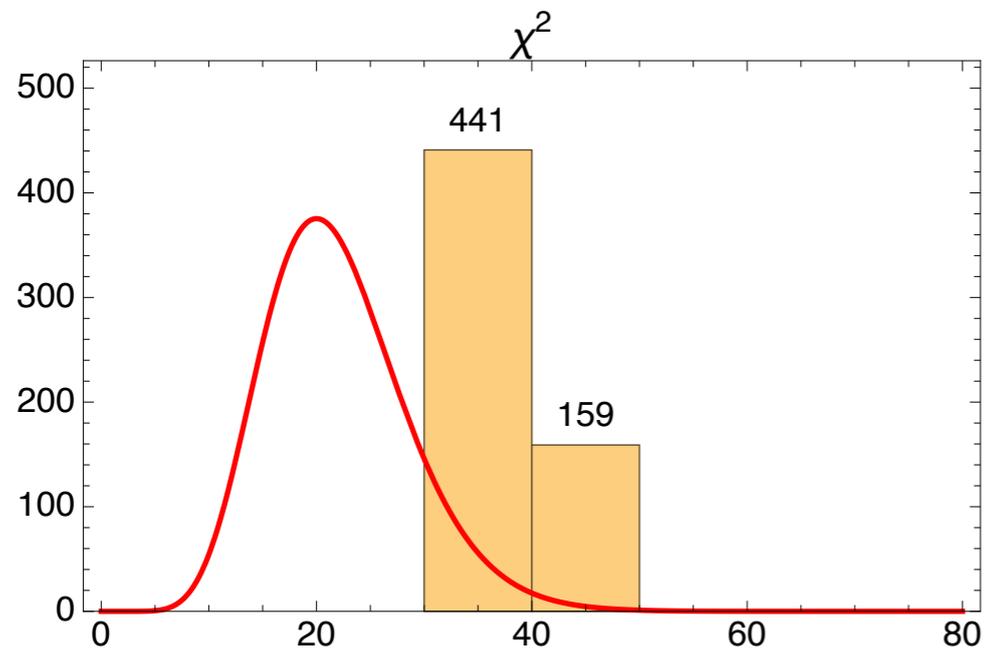
down

$x h_1^{d-\bar{d}} Q^2 = 2.4 \text{ GeV}^2$

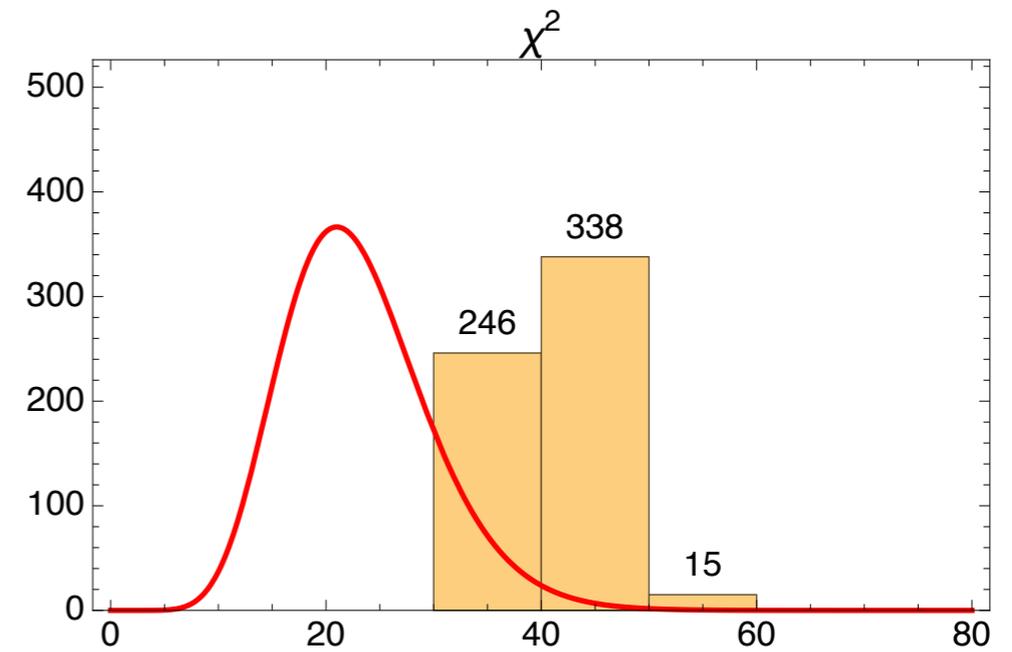


# $\chi^2$

$$\chi^2/\text{dof} = 1.76 \pm 0.11$$



$$\chi^2/\text{dof} = 1.82 \pm 0.25$$



probability density function of  
 $\chi^2$  distribution for  
22 d.o.f.      23 d.o.f.

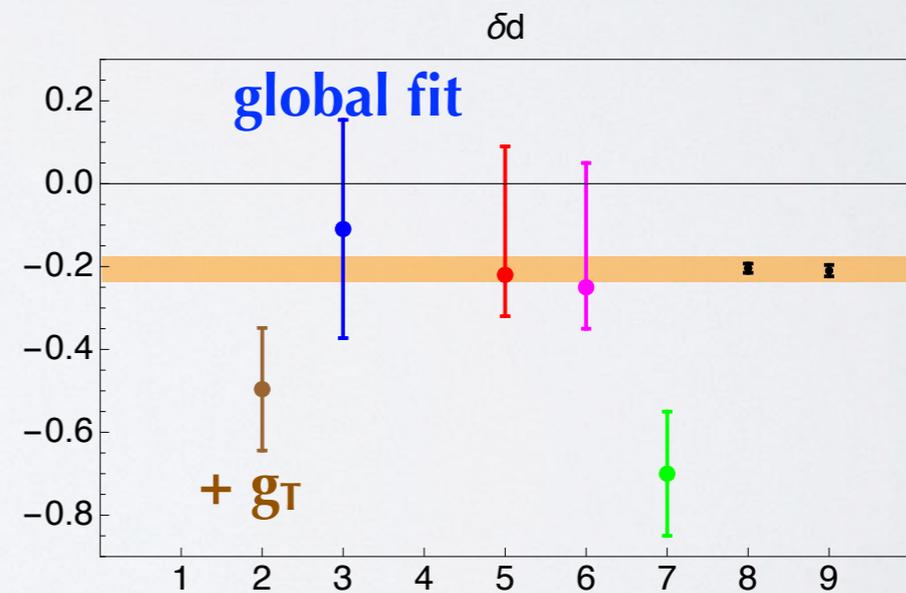
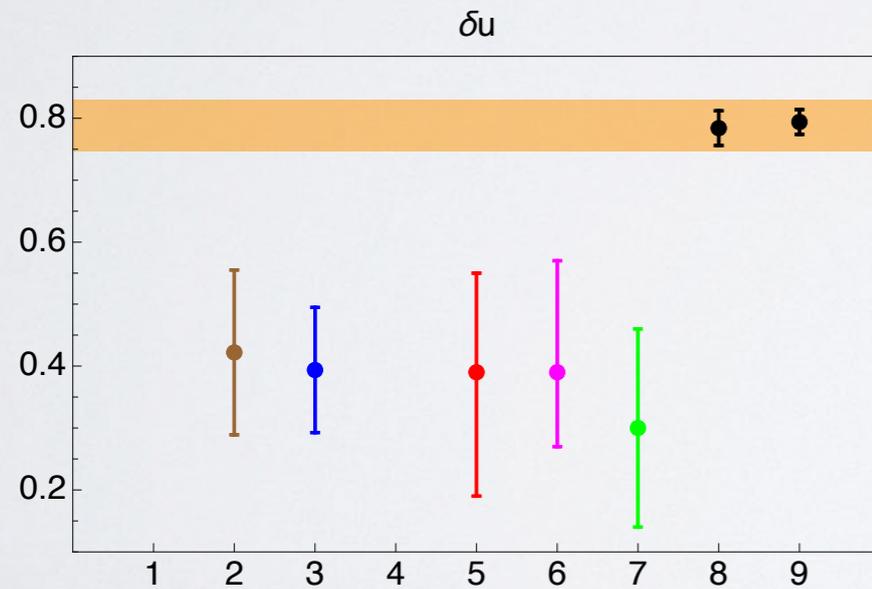
# compatibility with lattice

add to SIDIS+pp data  
constraint to reproduce from lattice  
 $g_T, \delta u, \delta d$

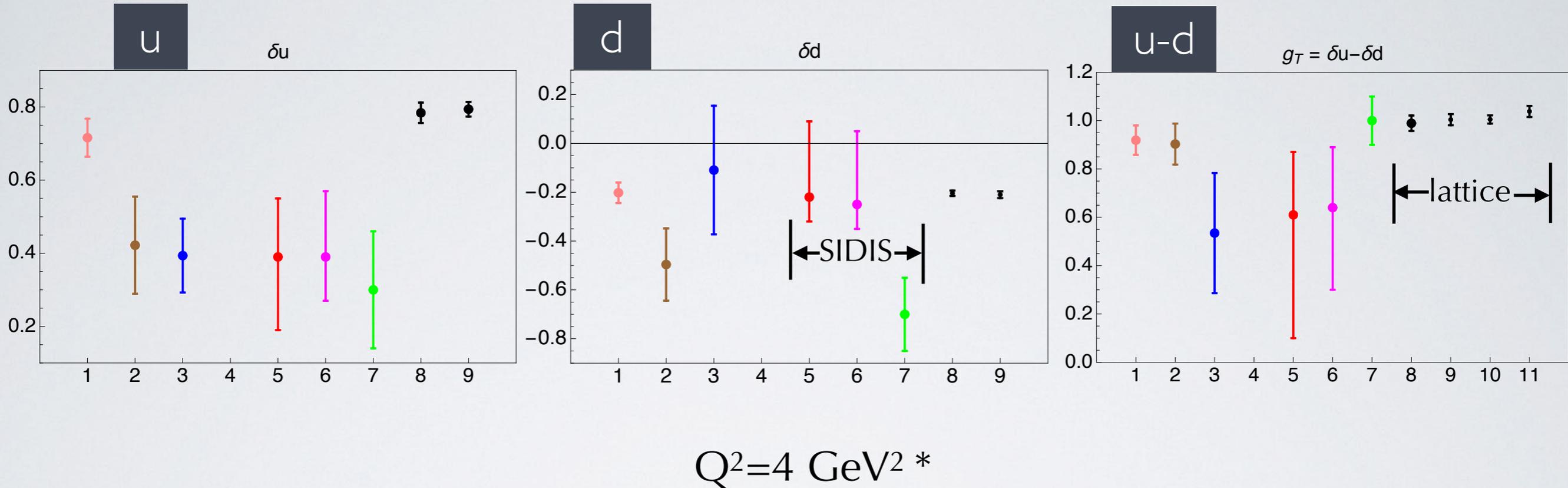
$$\overline{g_T}^{\text{latt}} = 1.004 \pm 0.057$$

$$\overline{\delta u}^{\text{latt}} = 0.782 \pm 0.031$$

$$\overline{\delta d}^{\text{latt}} = -0.218 \pm 0.026$$



# tensor charge



$Q^2=4 \text{ GeV}^2$  \*

1) global fit + constrain  $g_T$ ,  $\delta u$ ,  $\delta d$

2) global fit + constrain  $g_T$

Radici & Bacchetta,  
*P.R.L.* 120 (18) 192001

3) global fit '17

Kang et al., *P.R. D93* (16) 014009

5) "TMD fit" \*  $Q^2=10$

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8) PNDME '18

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9) ETMC '17

Alexandrou et al., *P.R. D95* (17) 114514;  
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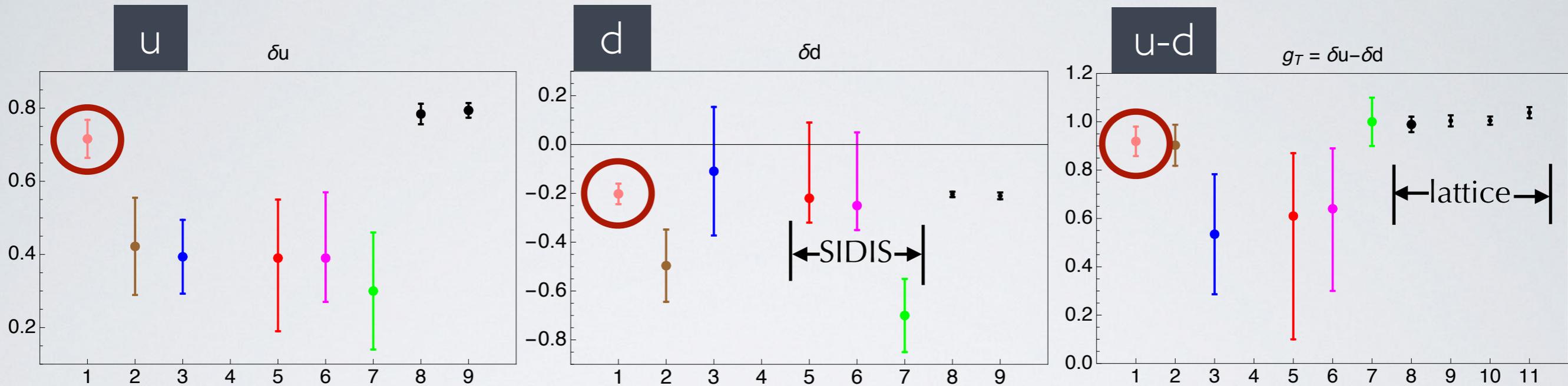
10) RQCD '14

Bali et al., *P.R. D91* (15)

11) LHPC '12

Green et al., *P.R. D86* (12)

# tensor charge



$Q^2=4 \text{ GeV}^2$  \*

compatible, but...

1) global fit + constrain  $g_T$ ,  $\delta u$ ,  $\delta d$

2) global fit + constrain  $g_T$

Radici & Bacchetta,  
*P.R.L.* 120 (18) 192001

3) global fit '17

Kang et al., *P.R. D*93 (16) 014009

5) "TMD fit" \*  $Q^2=10$

Anselmino et al., *P.R. D*87 (13) 094019

6) Torino fit \*  $Q^2=1$

Lin et al., *P.R.L.* 120 (18) 152502

7) JAM fit '17 \*  $Q_0^2=2$

8) PNDME '18

*Gupta et al., P.R. D*98 (18) 034503

9) ETMC '17

*Alexandrou et al., P.R. D*95 (17) 114514;  
*E P.R. D*96 (17) 099906

10) RQCD '14

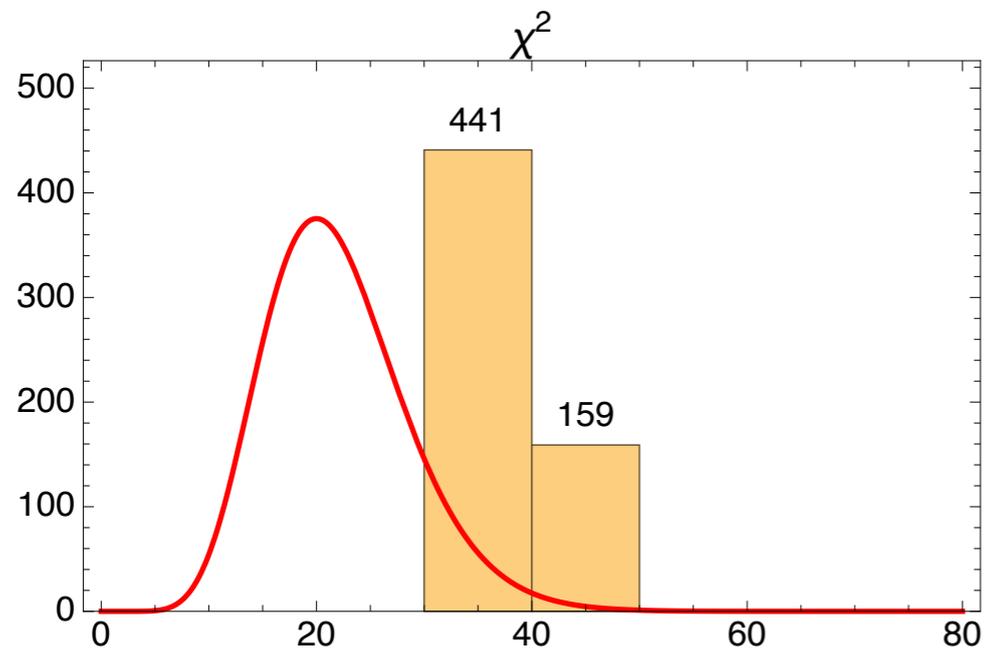
*Bali et al., P.R. D*91 (15)

11) LHPC '12

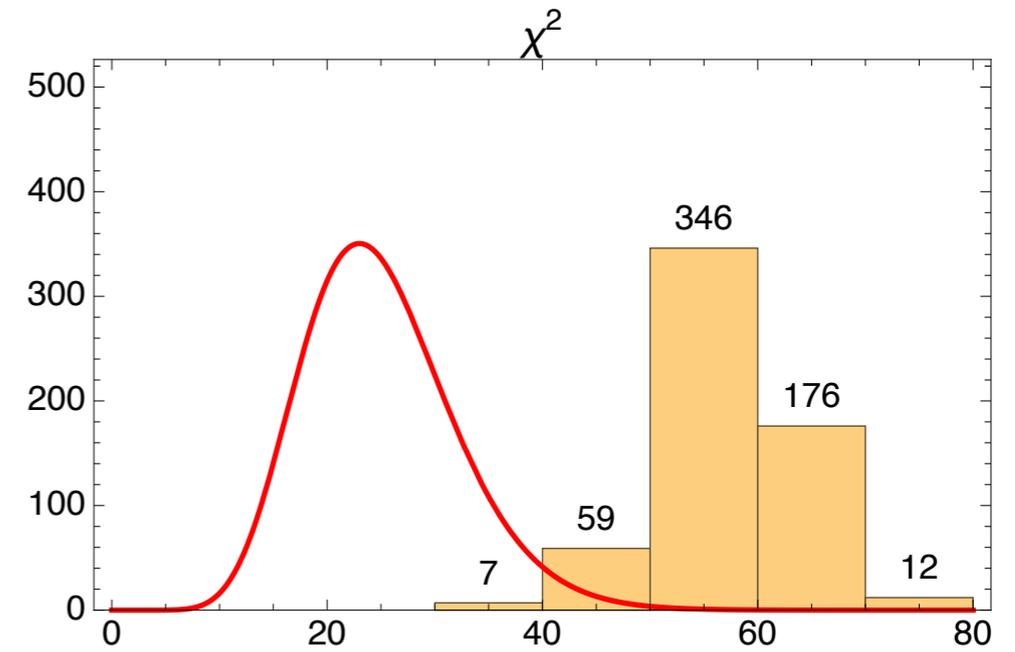
*Green et al., P.R. D*86 (12)

# $\chi^2$

$$\chi^2/\text{dof} = 1.76 \pm 0.11$$



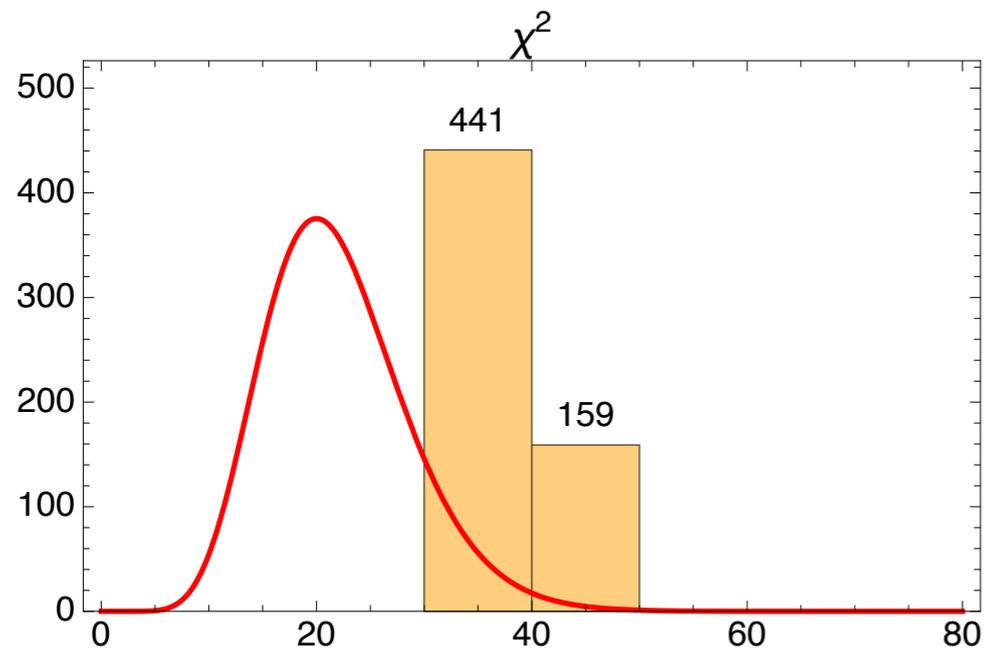
$$\chi^2/\text{dof} = 2.29 \pm 0.25$$



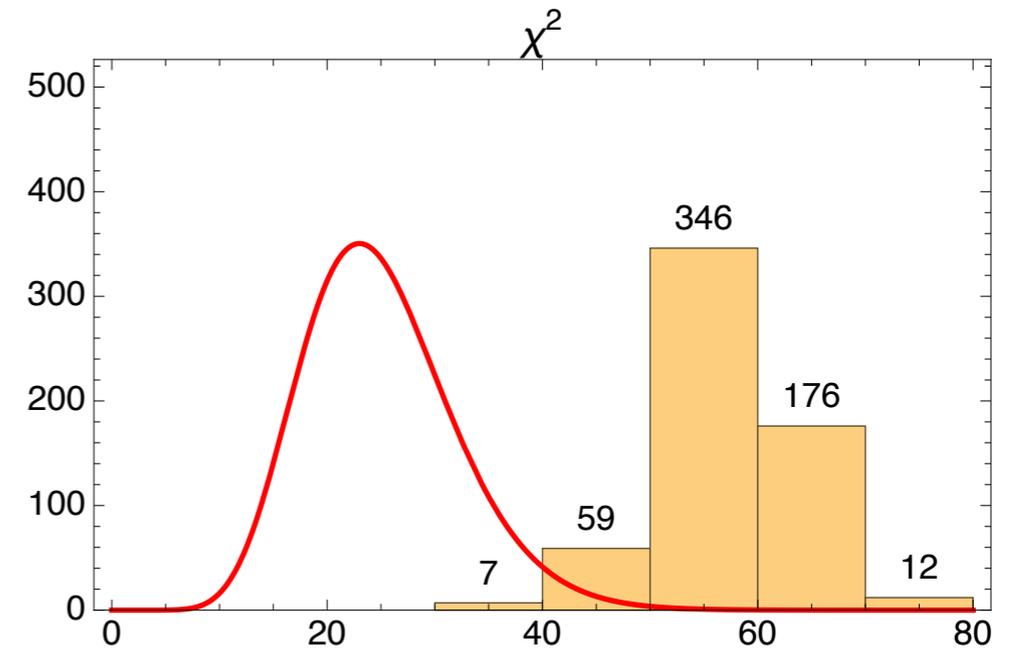
probability density function of  
 $\chi^2$  distribution for  
22 d.o.f.      25 d.o.f.

# $\chi^2$

$$\chi^2/\text{dof} = 1.76 \pm 0.11$$



$$\chi^2/\text{dof} = 2.29 \pm 0.25$$



probability density function of  
 $\chi^2$  distribution for  
22 d.o.f.      25 d.o.f.

compatible, but... statistically very unlikely !

# More (existing) data ...

## Di-hadron

→ refit di-hadron fragmentation functions using new data:

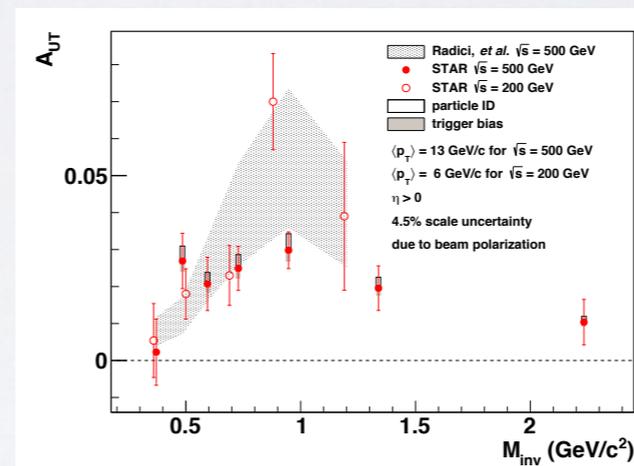
$e^+e^- \rightarrow (\pi\pi) X$  constrains  $D_{1^q}$   
(currently only by Montecarlo)



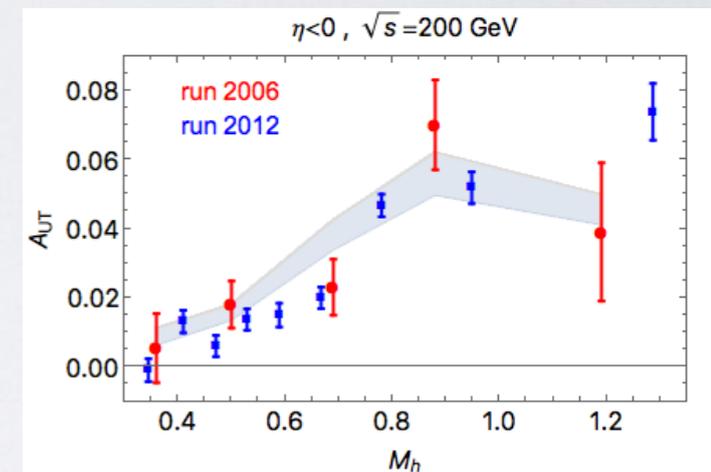
Seidl et al.,  
*P.R. D96 (17) 032005*

## p-p<sup>↑</sup> collisions

→ use also other (multi-dimensional) data from STAR run 2011 (s=500) and (later) run 2012 (s=200)



Adamczyk et al. (STAR),  
*P.L. B780 (18) 332*



Radici et al.,  
*P.R. D94 (16) 034012*

## SIDIS

→ use COMPASS data on  $\pi K$  and  $KK$  channels, and from  $\Lambda^\uparrow$  fragmentation: constrain strange contribution ?

# Conclusions / Open Problems

- first global fit of di-hadron inclusive data leading to extraction of **transversity as a PDF** in collinear framework
- inclusion of STAR p-p<sup>†</sup> data increases precision of up channel; **large uncertainty on down** due to unconstrained gluon unpolarized di-hadron fragmentation function → **need more/better “neutron target” data**
- **NO** apparent simultaneous **compatibility with lattice** for tensor charge in up, down, and isovector channels
- adding **Compass SIDIS pseudo-data for deuteron** increases precision of down, but leaves this scenario unaltered
- adding **CLAS12 SIDIS pseudo-data for proton** affects large x (error of up tensor charge reduced by ~2x), but tension with lattice even increased
- it is possible to force replicas to be **compatible with data and lattice** but situation is **statistically very unlikely**

enlarging the covered x-range is crucial → (target)<sup>†</sup> program at JLab I 2 !

**THANK YOU**

# Back-up

# the leading-twist PDF / TMD map

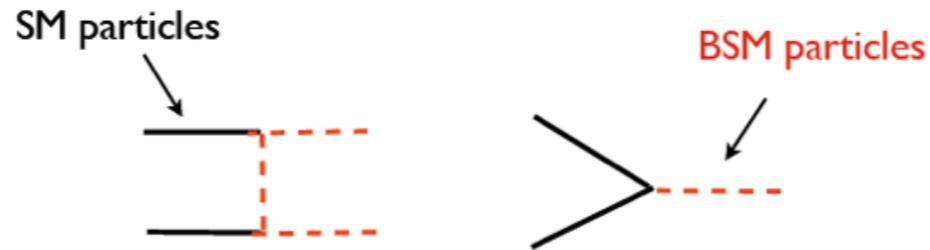
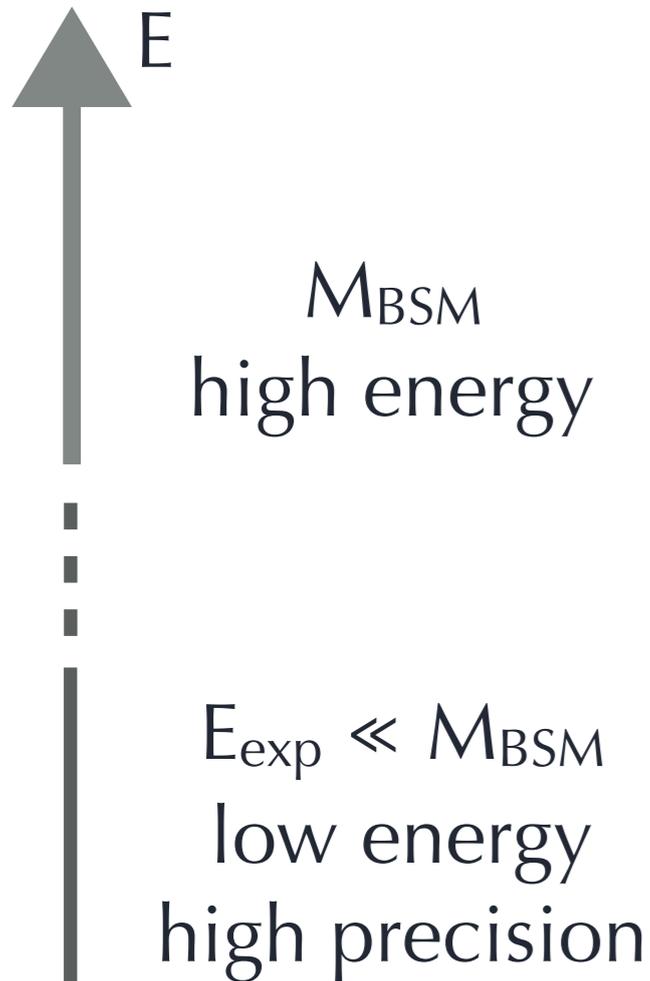
quark polarization

|                      |   |                                   |  |
|----------------------|---|-----------------------------------|--|
|                      | U | L                                 | <b>T</b>   |
| nucleon polarization | U | <b>f<sub>1</sub></b>              | <b>h<sub>1</sub><sup>⊥</sup></b>                       |
|                      | L |                                   | <b>h<sub>1L</sub><sup>⊥</sup></b>                      |
|                      | T | <b>f<sub>1T</sub><sup>⊥</sup></b> | <b>g<sub>1T</sub></b>                                  |
|                      |   |                                   | <b>h<sub>1</sub></b> <b>h<sub>1T</sub><sup>⊥</sup></b> |

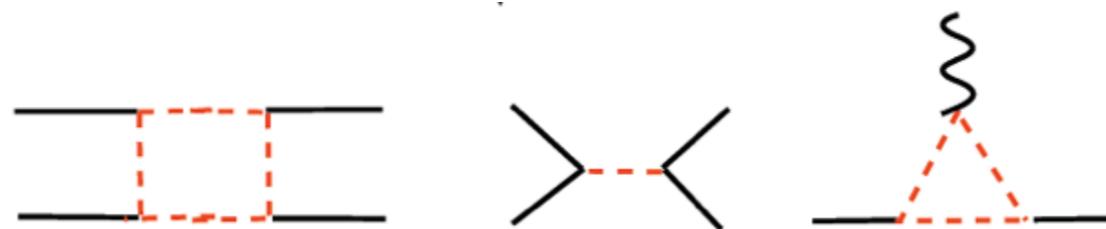
- 1- **h<sub>1</sub>** needed as the 3rd basic quark **PDF** for spin-1/2 objects
- 2- address novel QCD dynamics in the chiral-odd sector, also as **TMD**

# potential for BSM discovery ?

at least, two ways of searching :



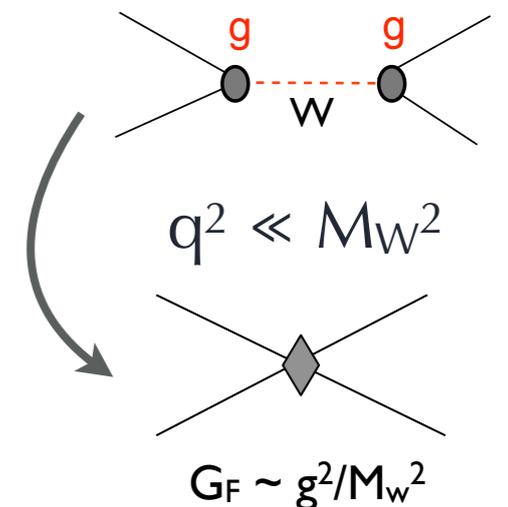
1- direct access to new particles



2- indirect access virtual effects



Example: weak CC interaction



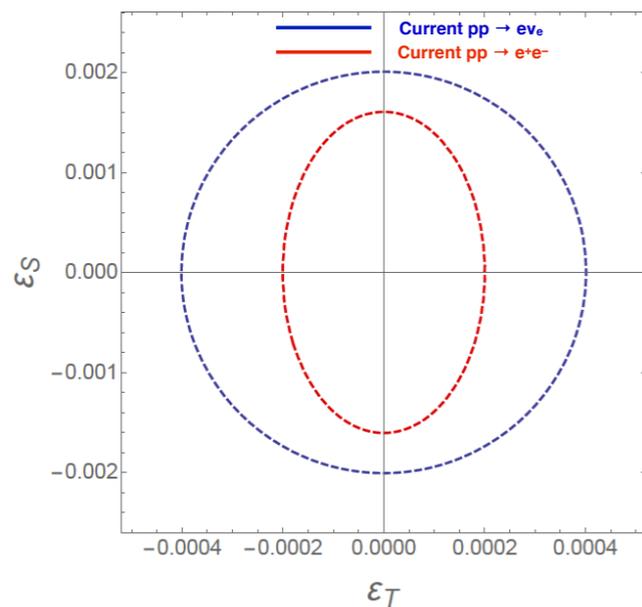
footprint:  
new local  
operators

# Examples of direct access

-  $pp \rightarrow e^- \nu + X$  search for  $W' \rightarrow e^- \nu$  with  $W'$  heavy partner of  $W$

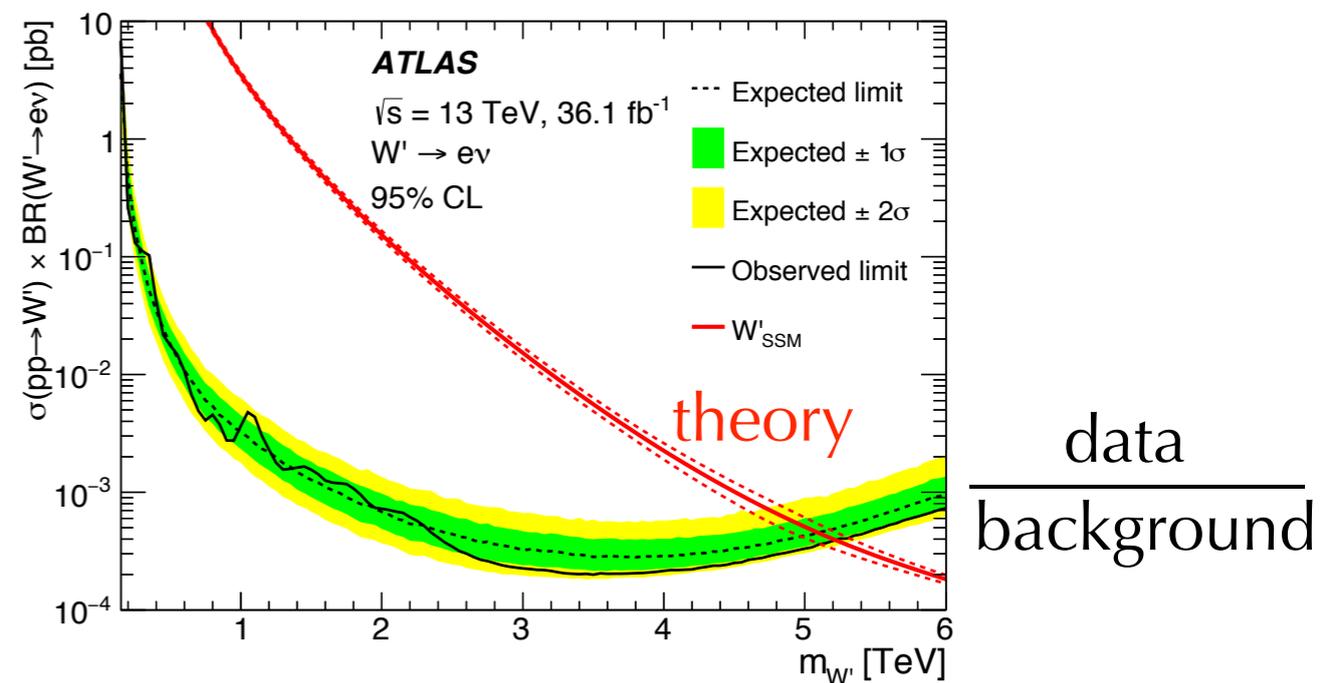
$M_{W'} > 5.1-5.2$  TeV at 95% C.L.

puts constraints on BSM operators including scalar ( $\epsilon_S$ ) & tensor ( $\epsilon_T$ )



*Gupta et al. (PNDME), P.R. D98 (18) 034503*

limits on cross section

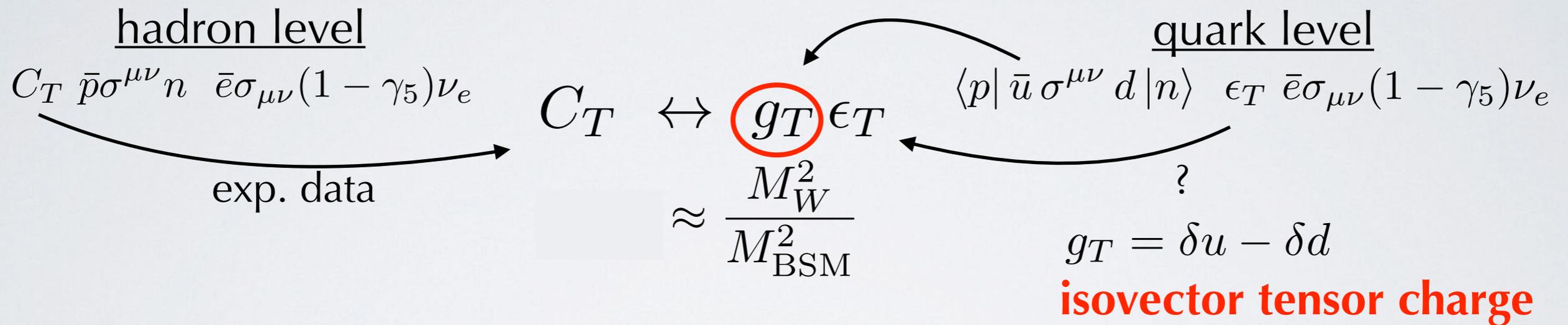


*Aaboud et al. (ATLAS), E.P.J. C78 (18) 401*

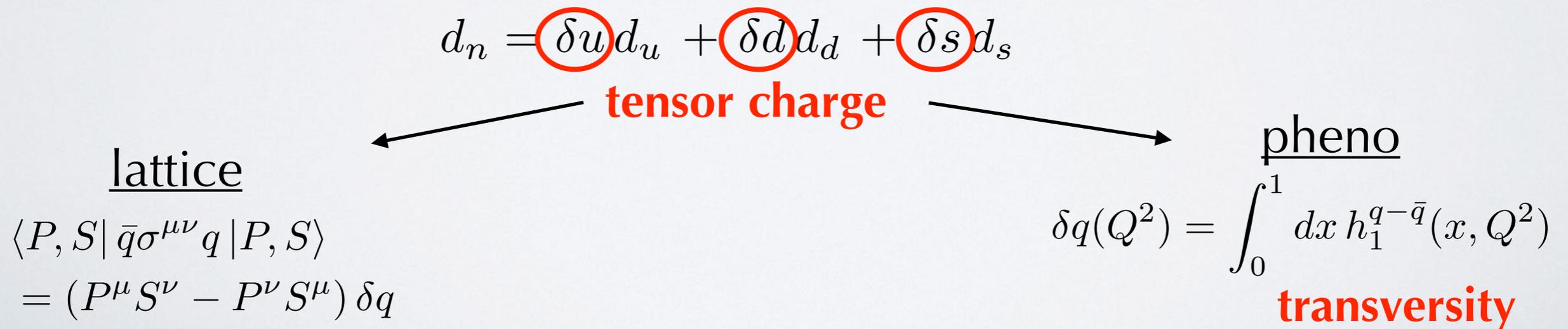
constraints reinforced including  
 $pp \rightarrow Z' \rightarrow e^- e^+ + X$

# Examples of indirect access

- **nuclear  $\beta$ -decay**: effective field theory including operators not in SM Lagrangian; for example, **tensor operator**



- **neutron EDM**: estimate CPV induced by quark chromo-EDM  $d_q$

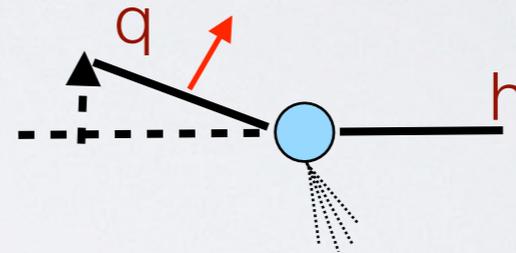


# extraction of transversity

transversity is chiral-odd  $\rightarrow$  need a chiral-odd partner

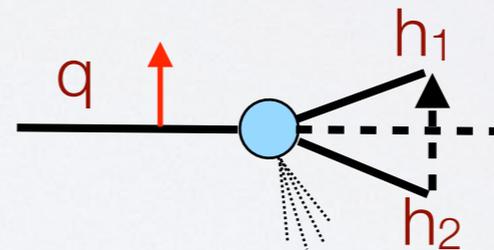
- **itself** : fully polarized Drell-Yan ✗

- **Collins function** : the Collins effect



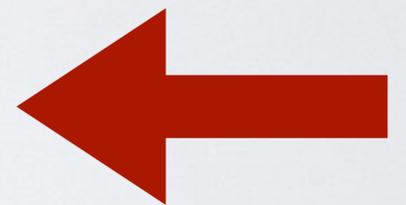
TMD framework  **$h_1$  as TMD**

- **IFF** : the di-hadron mechanism



collinear framework  **$h_1$  as PDF**

- **hadron-in-jet mechanism** : mixed framework  **$h_1$  as PDF**



# 2-hadron-inclusive production

framework  
collinear  
factorization

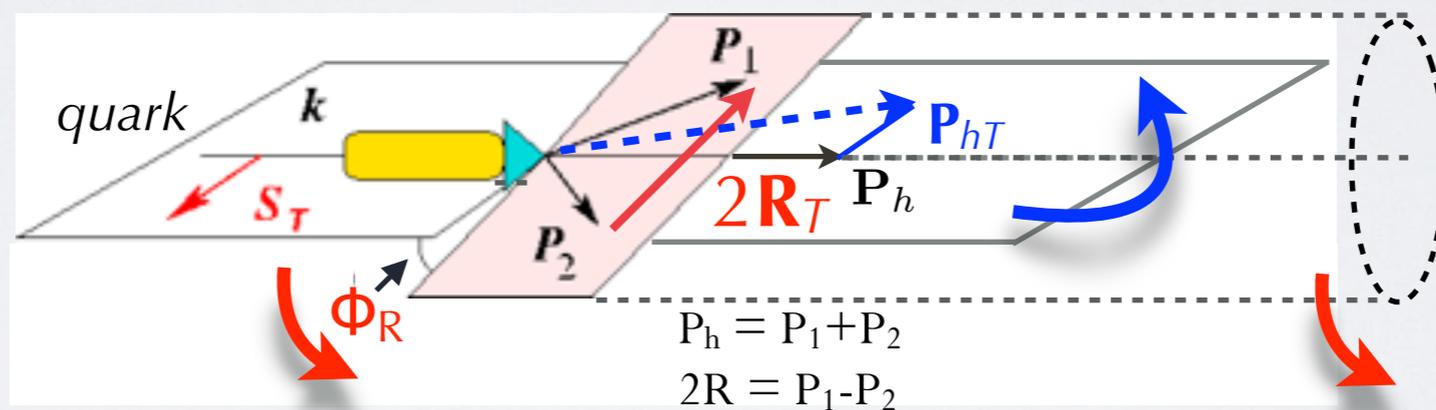
Collins, Heppelman, Ladinsky,  
N.P. **B420** (94)

$$R_T \ll Q$$

$$H_1^{\triangleleft}$$

$$M_h$$

invariant mass



correlation  $S_T$  and  $R_T \rightarrow$  azimuthal asymmetry

survives to  
polar  
symmetry  
(  $\int dP_{hT}$  )

# advantages of di-hadron mechanism

## collinear framework

- simple product of PDF and IFF

Ex.: SIDIS

$$A_{\text{SIDIS}}^{\sin(\phi_R+\phi_S)}(x, z, M_h^2) \sim -\frac{\sum_q e_q^2 h_1^q(x) \frac{|\mathbf{R}_T|}{M_h} H_{1,q}^{\triangleleft}(z, M_h^2)}{\sum_q e_q^2 f_1^q(x) D_{1,q}(z, M_h^2)}$$

x-dependence of  $A_{\text{SIDIS}}$  all in PDF

- flavor sum simplified by symmetries of IFF
  - + data on proton and deuteron targets
  - separate valence up and down
- factorization theorems for all hard processes
  - universality of  $h_1 H_1^{\triangleleft}$  mechanism

# advantages of di-hadron mechanism

$$A_{\text{SIDIS}}^{\sin(\phi_R + \phi_S)}(x, z, M_h^2) \sim - \frac{\sum_q e_q^2 h_1^q(x) \frac{|\mathbf{R}_T|}{M_h} H_{1,q}^{\triangleleft}(z, M_h^2)}{\sum_q e_q^2 f_1^q(x) D_{1,q}(z, M_h^2)}$$

$\pi^+\pi^-$   
tree level

$$H_1^{\triangleleft u} = -H_1^{\triangleleft d} \quad \text{isospin symmetry}$$

$$H_1^{\triangleleft q} = -H_1^{\triangleleft \bar{q}}$$

$$D_1^q = D_1^{\bar{q}}$$

} charge conjugation

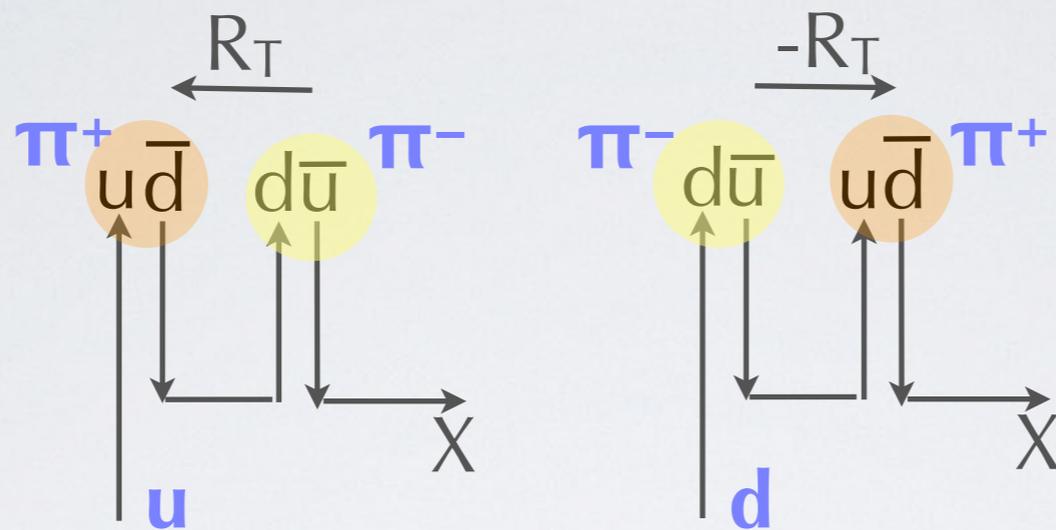
+ data on proton  
and deuteron targets

$$\text{proton} \quad xh_1^{u-\bar{u}} - \frac{1}{4}xh_1^{d-\bar{d}} = F [A_{\text{SIDIS}}^p \text{ data}, H_1^{\triangleleft u}, f_1^q D_1^q]$$

$$\text{deuteron} \quad xh_1^{u-\bar{u}} + xh_1^{d-\bar{d}} = \tilde{F} [A_{\text{SIDIS}}^D \text{ data}, H_1^{\triangleleft u}, f_1^q D_1^q]$$

separate valence up and down

# IFF symmetries



$$\begin{aligned}
 H_1^{\triangleleft u} &= -H_1^{\triangleleft d} && \text{isospin symmetry} \\
 H_1^{\triangleleft q} &= -H_1^{\triangleleft \bar{q}} && \left. \vphantom{H_1^{\triangleleft q}} \right\} \text{charge conjugation} \\
 D_1^q &= D_1^{\bar{q}}
 \end{aligned}$$

valid only for ( $\pi^+\pi^-$ ) pairs and at tree level

# hadronic collisions in Mellin space

$d\sigma(\eta, M_h, P_T)$  typical cross section for  $a+b \rightarrow c+d$  process

$$\frac{d\sigma_{UT}}{d\eta} \propto \int d|\mathbf{P}_T| dM_h \sum_{a,b,c,d} \int \frac{dx_a dx_b}{8\pi^2 \bar{z}} f_1^a(x_a) h_1^b(x_b) \frac{d\hat{\sigma}_{ab \rightarrow c+d}}{d\hat{t}} H_1^{\triangleleft c}(\bar{z}, M_h)$$

to be computed thousands times... usual trick: use **Mellin anti-transform**

$$h_1(x, Q^2) = \int_{C_N} dN x^{-N} h_1^N(Q^2) \quad N \in \mathbb{C}$$

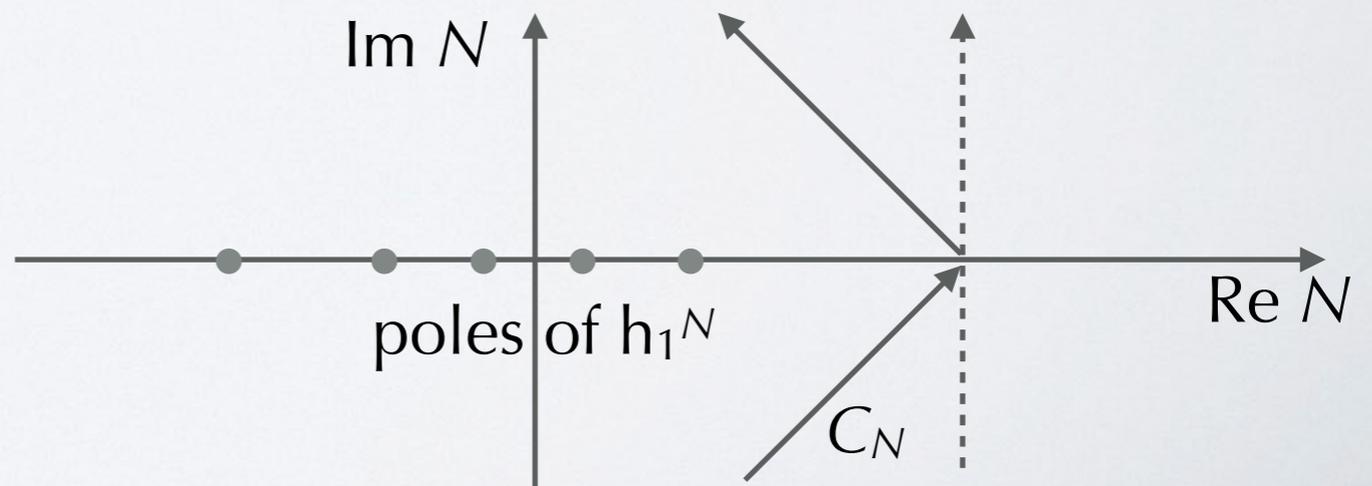
*Stratmann & Vogelsang,  
P.R. D64 (01) 114007*

$$\frac{d\sigma_{UT}}{d\eta} \propto \sum_b \int_{C_N} dN \int d|\mathbf{P}_T| h_{1b}^N(P_T^2) \int dM_h \sum_{a,c,d} \int \frac{dx_a dx_b}{8\pi^2 \bar{z}} f_1^a(x_a) x_b^{-N} \frac{d\hat{\sigma}_{ab \rightarrow c+d}}{d\hat{t}} H_1^{\triangleleft c}(\bar{z}, M_h)$$

$F_b(N, \eta, |\mathbf{P}_T|, M_h)$

pre-compute  $F_b$  only one time  
on contour  $C_N$

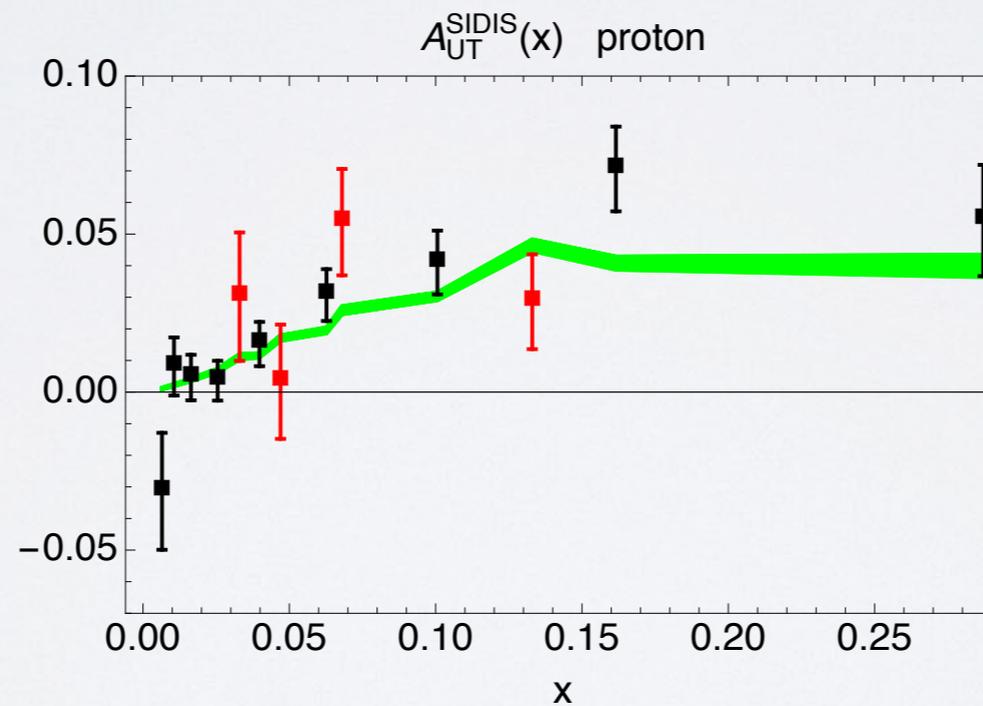
this **speeds up** convergence  
and facilitates  $\int dN$ , provided  
that  **$h_1^N$  is known analytically**



# statistical uncertainty

## the bootstrap method

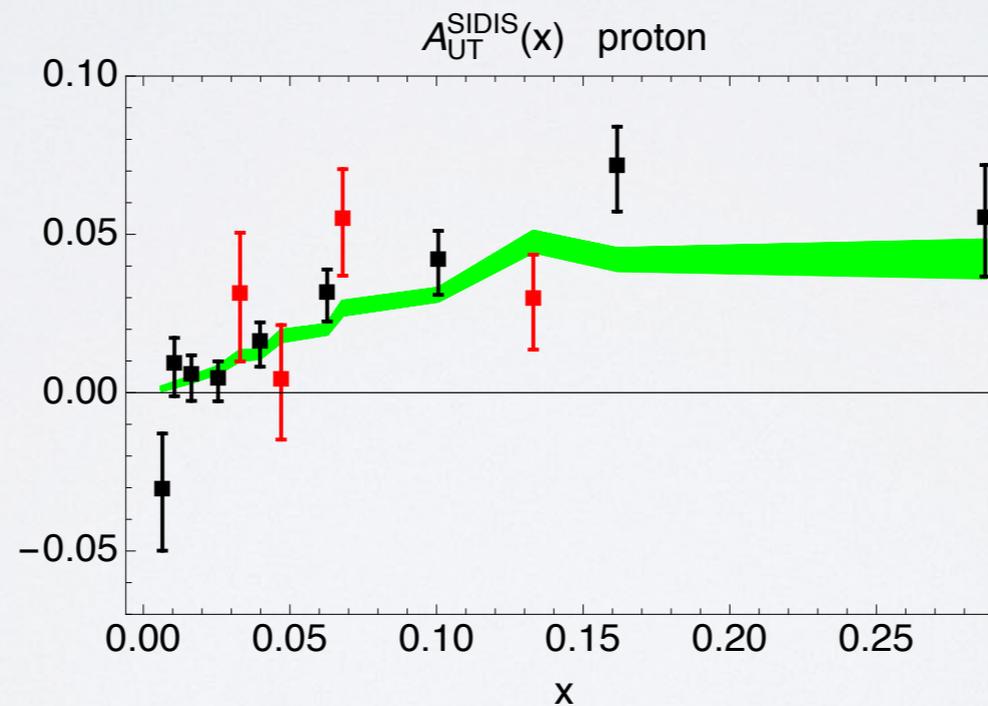
- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50



# statistical uncertainty

## the bootstrap method

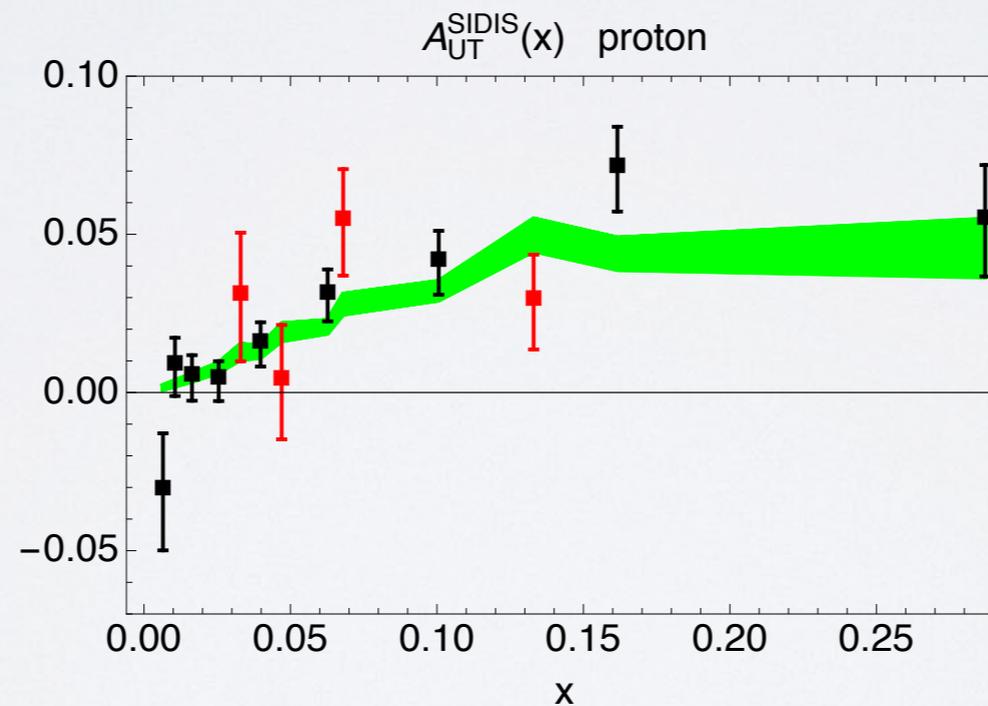
- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50, 100



# statistical uncertainty

## the bootstrap method

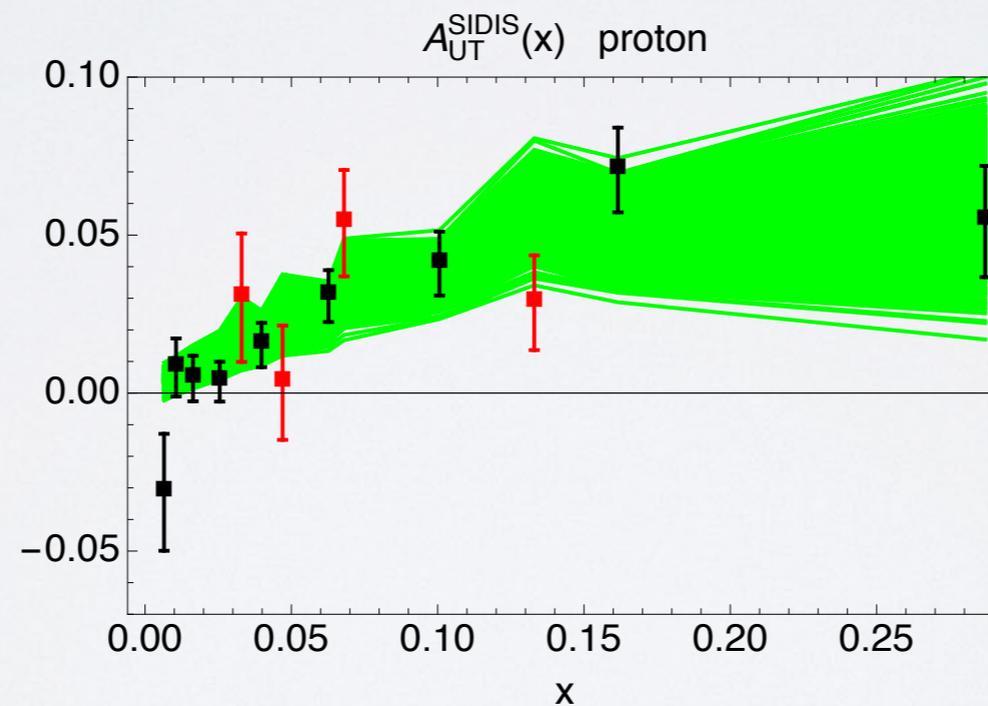
- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50, 100, 200 sets...



# statistical uncertainty

## the bootstrap method

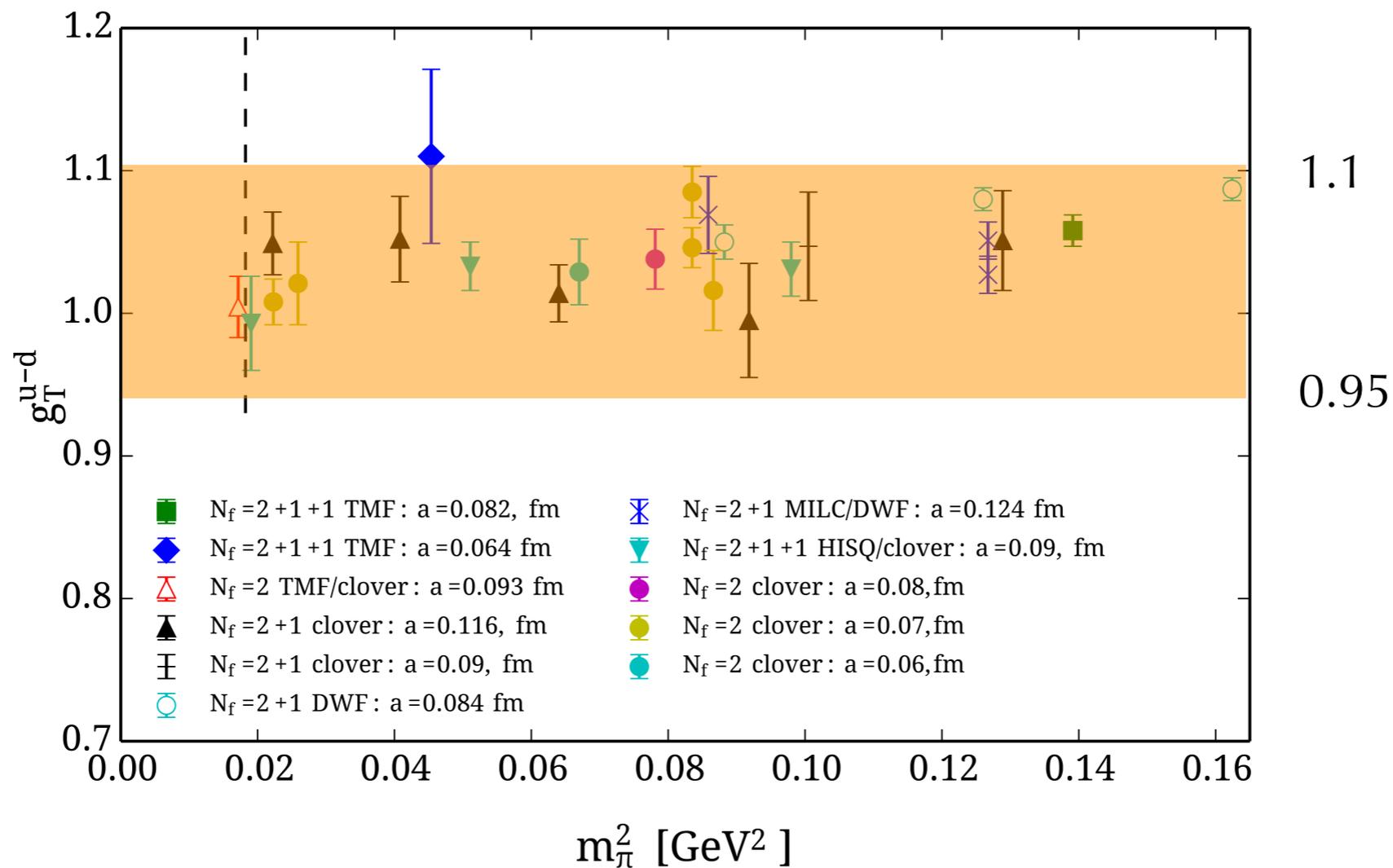
- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50, 100, 200 sets... until average and standard deviation reproduce original exp. points (here,  $200 \times 3 = 600$ )



# isovector tensor charge $g_T = \delta_u - \delta_d$

lattice results  
with different  
discretization schemes, lattice spacings, volumes

$Q^2=4 \text{ GeV}^2$



lattice quasi-PDF

see also arXiv:1803.04393 (LP<sup>3</sup>)

Alexandrou, arXiv:1612.04644

# impact of “full” lattice constraint

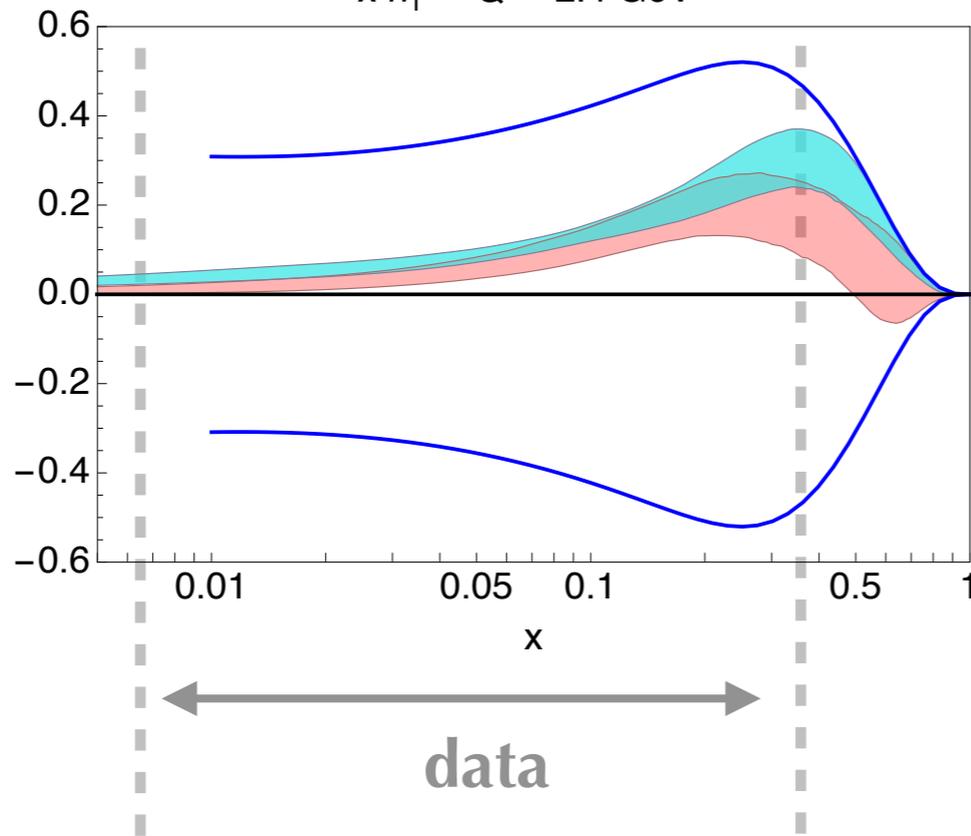
global fit + lattice ( $g_T, \delta u, \delta d$ ) constraint

global fit

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_{1^u} / 4 \\ D_{1^u} \end{cases}$$

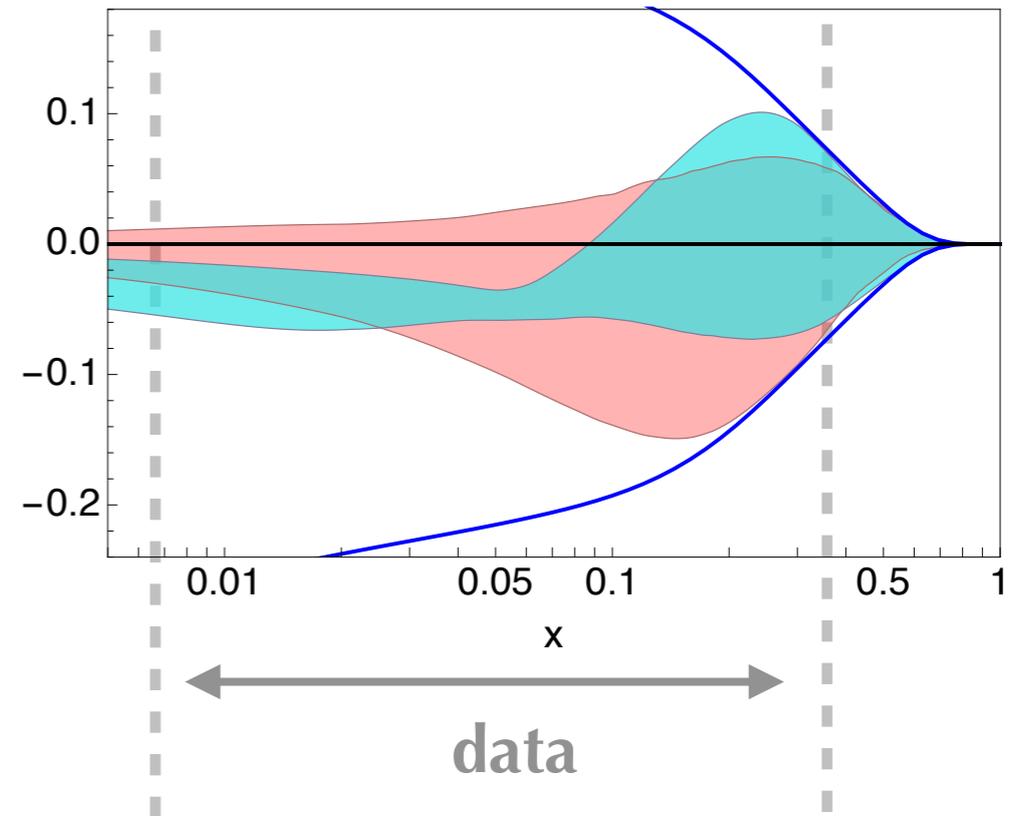
up

$x h_1^{u-\bar{u}} Q^2 = 2.4 \text{ GeV}^2$



down

$x h_1^{d-\bar{d}} Q^2 = 2.4 \text{ GeV}^2$



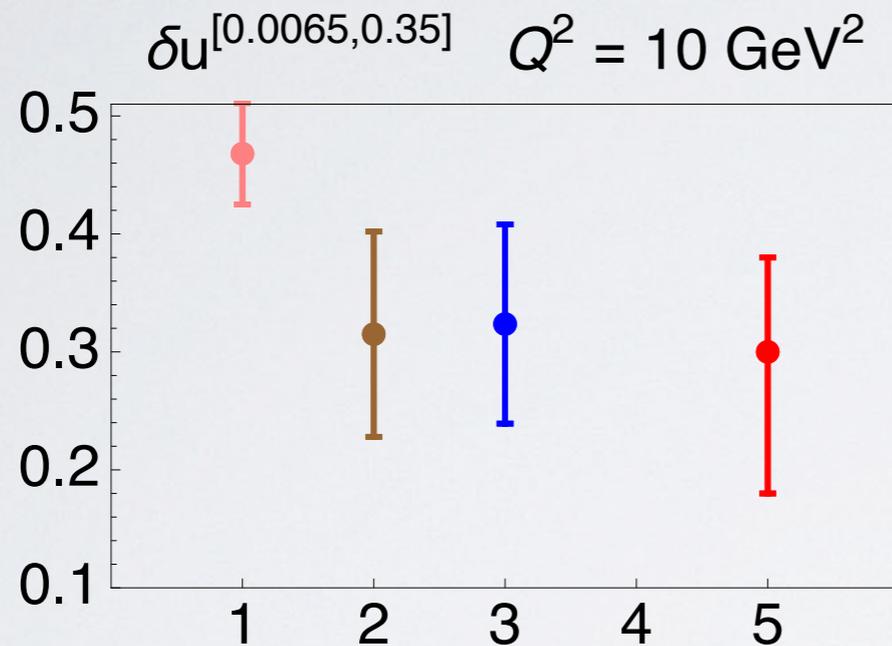
higher up, even within  
x-range of data

# truncated tensor charge

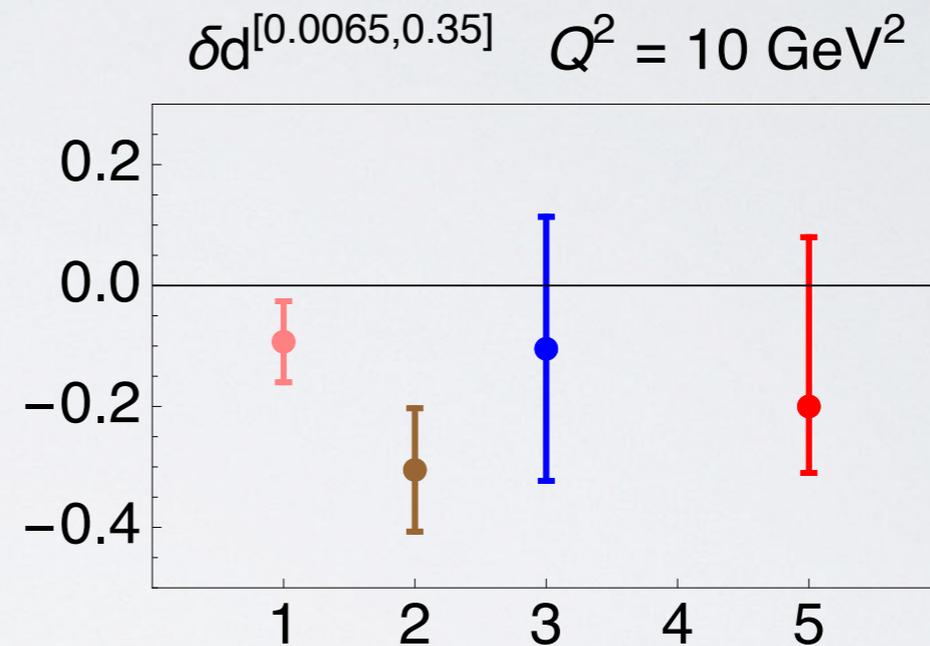
truncated

$$\delta q^{[0.0065,0.35]} \quad Q^2 = 10$$

up



down



1) **global fit + constrain  $g_T, \delta u, \delta d$**

2) **global fit + constrain  $g_T$**

3) **global fit '17** *Radici & Bacchetta, P.R.L. 120 (18) 192001*

5) **"TMD fit"** *Kang et al., P.R. D93 (16) 014009*