Interactions of UHE Cosmic Rays What do we learn from Nuclear Physics aspects?

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What do we observe? Cosmic Ray Flux

- It extends over many orders of magnitude, different techniques needed for measurements
- Power-law → similar mechanisms of production at different energies ?

shock wave

<u>Acceleration in shocks</u> → Astrophysical objects with jet formation: Gamma-Ray Bursts, Active Galactic Nuclei, ...

B field



What do we observe? Cosmic Ray Flux

- It extends over many orders of magnitude, different techniques needed for measurements
- Ultra-High Energies (mostly extragalactic sources)
 - \rightarrow Pierre Auger Observatory
 - \rightarrow Telescope Array
 - Measurement of energy deposit in atmosphere
 - Collection of particles at ground





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What do we observe? UHECR Composition



Cosmic Ray accelerators



 $E_{max} \sim q B R$





E_{max} > **10⁸ TeV**

Main difference: size of accelerator !

 \rightarrow AGN, GRBs: R ~ 100,000 – 10,000,000,000 km

UHECRs from sources to detection



UHECRs from sources to detection

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UHECR interactions in the extragalactic space

UHECR interactions

Example: protons interacting with CMB photons

- CR proton in the extragalactic space with 10²⁰ eV (Lorentz factor Γ~ 10¹¹)
- CMB photons ~ 10⁻³ eV (Lab frame)
- Energy scale of the processes: energy of the photon in the proton rest frame, $\epsilon' \simeq \Gamma \epsilon$, for this example order of ~ 100 MeV $\rightarrow -pion$ mass

- Photo-meson production is triggered → consequences:
 - Energy loss of protons
 - Production of secondary neutrinos and photons

UHECR interactions

- Highest energies: suppression of the spectrum \rightarrow **<u>GZK effect</u>** (maybe)
- Intermediate energies: dip → energy losses due to pair production (more robust feature wrt suppression)
- CAVEAT: this is valid for pure proton composition at the sources !

UHECR interactions

Example: Fe nuclei interacting with CMB photons

- CR iron in the extragalactic space with 10²⁰ eV → simple treatment: Fe-56 isotope is a superposition of 56 nucleons (Lorentz factor Γ ~ 2 x 10⁹)
- CMB photons ~ 10^{-3} eV (Lab frame)
- Not enough for pion production !
- Other relevant processes at UHE (photodisintegration):
 - Giant Dipole Resonance ~ 10 MeV in the NRF
 - <u>Quasi-Deuteron</u> → above ~ 30 MeV in the NRF, up to the threshold for photo-pion production
- Consequences:
 - Suppression of the flux → if nuclei, due to photodisintegration
 - Production of neutrinos and photons connected to GZK effect → if nuclei, less efficient !

Extragalactic Cosmic Ray Propagation

- In order to infer source properties from observations, we need to take into account:
 - adiabatic energy losses (expansion of the Universe)
 - interactions with extragalactic background photons
 - deflection of magnetic fields
- Simulation codes for extragalactic propagation
 - SimProp, Aloisio, Boncioli, di Matteo, Grillo, Petrera, Salamida, JCAP 2017
 - **CRPropa 3**, Alves Batista, Dundovic, Erdmann, Kampert, Kuempel, Müller, Sigl, van Vliet, Walz, Winchen, JCAP 2016
 - TransportCR, Kalashev & Kido, J.Exp.Theor.Phys 2016
 - **HERMES**, De Domenico, J.Exp.Theor.Phys 2013
 - **PRINCE**, Heinze & Fedynitch, in preparation

Extragalactic Cosmic Ray Propagation

- In order to infer source properties from observations, we need to take into account:
 - adiabatic energy losses (expansion of the Universe)
 - interactions with extragalactic background photons

 \rightarrow subject to uncertainties related to interaction cross sections and background photon fields

- \rightarrow similar uncertainties concern also interactions in the source
- \rightarrow interaction rate proportional to density of photons and cross section

UHECR interactions in a candidate source

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Example: Gamma-Ray Burst

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Nuclear Cascade in a GRB

[Biehl, Boncioli, Fedynitch, Winter, A&A 2017]

- Example calculated in a GRB shell, but a similar calculation can be performed with the photon fields of the extragalactic space
- Main control parameters for the development of the cascade
 - typical radius of the collisions
 - luminosity of the source

Classification

- luminosity of the source

Classification

Nuclear Physics uncertainties

Photo-disintegration – current status

- EXFOR contains 14 absorption cross sections < Fe
- 47 measurements where at least one inclusive cross section available
- Located mostly on main diagonal (stable elements)
- All other isotopes need model prediction → not always well reproducing data

Photo-disintegration – develompment of the cascade

- GRB photon field, Fe injected
- TALYS vs PSB: disintegration is more efficient, more channels are included
- If only measured cross section are included, cascade cannot fully develop (as in the PSB case)
- Precise knowledge of cross sections can influence the expectation on composition!

Nuclear Physics and the interpretation of UHECR data

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Interpretation of UHECR data - flux

- <u>Final aim: infer properties of UHECR sources</u> <u>from data</u>
- Exercise: fix spectral index and composition at the source → how the spectra after propagation change if different disintegration models are used?
- If disintegration is more efficient, the propagated spectra are more depleted, especially at the highest energies → a softer spectral index at the injection is required to obtain the same spectra as with the less efficient disintegration

[Alves Batista, Boncioli, di Matteo, van Vliet, Walz, JCAP 2015] [Alves Batista, Boncioli, di Matteo, van Vliet, in prep.]

Interpretation of UHECR data - composition

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Other uncertainties affecting UHECR physics

- UHECR source distribution
- UHECR propagation:
- EBL spectrum and evolution: factor-of-2 uncertainty in the far IR, even in modern models → influences the photodisintegration (Alves Batista, Boncioli, di Matteo, van Vliet, Walz JCAP 2015, Auger Collaboration JCAP 2017)

- Magnetic fields: stronger IGMF \rightarrow softer injection spectrum required (Auger Collaboration, ICRC 2017)

- UHECR composition estimates rely on extrapolations of hadronic interactions at energies not accessible at LHC → large differences
- Experimental measurements → differences among spectrum and (interpretation of) composition results by different experiments

Conclusions

- The study of UHECR interactions is needed for the understanding of the features of the UHECR spectrum and the UHECR composition
- Complementary information provided by secondary messengers (neutrinos and photons produced in the propagation) helps for having a deeper knowledge of sources → cosmological evolution of UHECR sources cannot be investigated only with UHECRs ("horizon" problem)
- Scarce knowledge about relevant aspects of nuclear physics in cosmic ray astrophyisics may introduce problems in interpretation of UHECR data

What we aim for the future

- UHECR experiments: better understanding of systematics and determination of UHECR composition

- Gamma-ray experiments: constraints on EBL spectrum and evolution
- Neutrino experiments: help in constraining UHECR source evolution
 - Measurements of total absorption cross sections
 - \rightarrow needed for interaction rate calculations
 - → measurements are sparse
 - Especially, clarify **isobar situation** (unstable elements)
 - Development of the cascade depends on branching ratios/multiplicities for different channels → data on residual cross sections will ensure that we

don't have systematic offsets in models \rightarrow We propose systematic measurements to improve the predictability of unmeasured cross sections \rightarrow measuring the total absorption cross section for two or more different isobars

Backup slides

UHECR interactions: extragalactic photons

Relevant backgrounds(ϵ = photon energy in lab frame) $\epsilon \leq 3 \text{ meV}$ (MW): cosmic microwave background (CMB) $1 \text{ meV} \leq \epsilon \leq 10 \text{ eV}$ (IR to UV): extragalactic background light (EBL)

M.G. Hauser and E. Dwek, Ann. Rev. Astron. Astrop. 39 (2001) 249

 $\varepsilon' \simeq \Gamma \varepsilon$

- EBL + intermediate energy CRs
- CMB → higher density wrt EBL

UHECR interactions: protons

Main processes (ϵ' = photon energy in nucleus rest frame) $\epsilon' \gtrsim 1$ MeV: pair production, $N + \gamma \rightarrow N + e^+ + e^-$ →A.M. Hillas, Phys. Lett. 24A 677 (1967) →G.R. Blumenthal, Phys. Rev. D Vol 1 1596 (1970)

 <u>Robustness of dip feature</u>: protons are collected from a large volume → feature not sensitive to local over-density or deficit of sources

• Its observation has been considered as <u>evidence for proton composition</u>, for example Berezinzky, Gazizov, Grigorieva Phys.Rev. D74 (2006) 043005

UHECR interactions: protons

Main processes $(\epsilon' = \text{photon energy in nucleus rest frame})$

 $\epsilon' \gtrsim 1$ MeV: pair production, $N + \gamma \rightarrow N + e^+ + e^-$

 $\epsilon'\gtrsim$ 150 MeV: pion production, e.g. $\mathbf{p}+\gamma\rightarrow\mathbf{n}+\pi^+$

→K. Greisen, PRL 16 748 (1966),
→G.T. Zatsepin and V.A. Kuzmin, Sov. Phys. JETP Lett. 4 78

<u>Suppression</u>: can be mimicked by acceleration cutoff and/or deficit of sources

UHECR interactions and effects in the observables

• Suppression due to effect of the minimum distance of the sources

[Aloisio & Boncioli, Astrop. Phys. 2011]

 Suppression due to maximum energy of acceleration of injected protons

- Even in the simple case of a pure proton composition, the suppression can be due to different aspects or to a combination of them.
- In general, degenerate scenarios looking at the spectrum → can be distinguished with secondary particles !

UHECR interactions: nuclei

Main processes $(\epsilon' = \text{photon energy in nucleus rest frame})$ $\epsilon' \gtrsim 1$ MeV: pair production, $N + \gamma \rightarrow N + e^+ + e^ \epsilon' \gtrsim 8$ MeV: disintegration, e.g. $Z^A + \gamma \rightarrow Z^{A-1} + n$ $\epsilon' \gtrsim 150$ MeV: pion production, e.g. $p + \gamma \rightarrow n + \pi^+$

[Alves Batista, Boncioli, di Matteo, van Vliet,

UHECR interactions and effects in the observables

 \rightarrow Secondary nucleons produced in the photo-disintegration chain have energies not larger than E(Fe) /A \Rightarrow in the case of cut-off=20.5 the secondary protons are confined at low energies wrt the case of cut-off=22

 \rightarrow this affects the composition observables

■ Suppression of the flux → established, but not its origin: Maximum energy or Photo-disintegration scenario?

UHECR interactions and effects in the observables

cut-off: 20.5

Degenerate scenarios looking at the spectrum → can be distinguished with composition, thanks to propagation!

 $\langle X_{max} \rangle = \langle X_{max} \rangle_p + f_E \langle \ln A \rangle$ $\sigma^2(X_{max}) = \langle \sigma_{sh}^2 \rangle + f_E^2 \sigma_{\ln A}^2$

UHECRs and cosmogenic neutrinos

- Normalization of the CR flux, injection spectral index and maximum energy
- Composition
- Source evolution history → not possible to constrain with only CR measurements !
- Details of EBL and cross section models

[van Vliet, Horandel, Alves Batista, ICRC 2017]

10⁰ L 10⁰

 10^{9}

UHECRs and cosmogenic neutrinos

[Aloisio, Boncioli, di Matteo, Grillo, Petrera, Salamida, JCAP 2015]

- Source evolution history \rightarrow not possible to constrain with only CR measurements ! \rightarrow see also Heinze, Boncioli, Bustamante, Winter, ApJ 2016
- Neutrino flux also dependent on EBL models !

Propagation models and secondary messengers

- Cosmogenic neutrinos and photons corresponding to <u>proton scenario</u> and <u>mixed</u> <u>composition scenario</u> (hard spectral index and low energy cutoff)
- Complete fit of composition, shape of initial spectrum <u>and source evolution</u>, including <u>neutrinos and photons</u> can help to have a global picture and constrain/exclude scenarios (photons can give more stringent constraints than neutrinos in the local universe, see for example Supanitski 2016)

Extragalactic magnetic fields

- Effect of magnetic fields on the propagation
 - It is higher for heavier masses
 - It limits the horizon of the universe contributing to the UHECRs
 - The average mass at Earth is lighter

 Global effect: more interactions → softer spectra (at the injection) can describe the data, because the magnetic fields naturally suppress the heavy nuclei at low energies

(as found for example in Mollerach & Roulet 2013 and Wittkowski for the Auger Coll, ICRC 2017)

CRPropa 2, Astropart. Phys. 2013]

Interpretation of UHECR data

[Auger Coll, JCAP 2017]

Interpretation of UHECR data

Fit CR data: spectrum and composition results

[Auger	Coll.	JCAP	20171
L	,		

	MC code	$\sigma_{\rm photodisint.}$	EBL model
SPG	SimProp	PSB	Gilmore 2012
STG	SimProp	TALYS	Gilmore 2012
SPD	SimProp	PSB	Domínguez 2011
CTG	CRPropa	TALYS	Gilmore 2012
CTD	CRPropa	TALYS	Domínguez 2011
CGD	CRPropa	Geant4	Domínguez 2011

Uncertainties of cross section models and EBL in CR sources and propagation

- Effect on CR observables:
 - Alves Batista, Boncioli, di Matteo, van Vliet, Walz, JCAP 2015
 - Boncioli, Fedynitch, Winter, Sci. Reports 2017
- Effect on cosmogenic neutrinos and photons
 - Aloisio, Boncioli, di Matteo, Grillo, Petrera, Salamida, JCAP 2015
 - Alves Batista, Boncioli, di Matteo, van Vliet, in prep.

Uncertainties: EBL models

Extragalactic background light at z=0

- High energies \rightarrow interactions happen mainly with CMB photons
- Low energies and light masses → effect of enhanced interactions with EBL and accumulation of lighter masses
- Difference is increased in the "hard" injection scenario → higher than Auger statistical errors in the spectrum !

[Alves Batista, Boncioli, di Matteo, van Vliet, Walz , JCAP 2015] [Alves Batista, Boncioli, di Matteo, van Vliet, in prep.]

Uncertainties: photo-disintegration cross sections

- One-path for photo-disintegration (**PSB model**, adapted from Puget, Stecker, Bredekamp ApJ 1976)
- More sophisticated models (like TALYS, http://www.talys.eu/) compute cross sections for all exclusive photo-disintegration channels (ejection of protons, neutrons, deuterons, tritiums, helium-3 and helium-4 nuclei, and any combinations thereof)

• Problems:

Z=8

- Models do not (always) agree with data !
- Data is not totally available !

[Alves Batista, Boncioli, di Matteo, van Vliet, Walz, JCAP 2015] [Boncioli, Fedynitch, Winter, Sci.Rep. 2017] [Alves Batista, Boncioli, di Matteo, van Vliet, in prep.]

Uncertainties: photo-disintegration cross sections

- Spectra at Earth \rightarrow differences more visible in hard injection
- Composition:
 - Hard injection: more efficient disintegration gives smoother transition from light to heavy (also found in GRB sources, see Boncioli, Fedynitch, Winter, Sci. Reports 2017)
 - Soft injection: the composition at Earth is more mixed in the whole energy range

[Alves Batista, Boncioli, di Matteo, van Vliet, Walz, JCAP 2015] [Alves Batista, Boncioli, di Matteo, van Vliet, in prep.] We are interested in the total <u>photoabsorption cross section</u> and in the <u>inclusive cross sections</u>

Interaction framework and terminology

Interactions of cosmic rays in the source environment or in the propagation can be rigorously followed with a system of differential equations describing the evolution of the differential particle density wrt time, taking into account all interactions that can modify their number and energy.

$$\frac{\partial N_i}{\partial t} = \frac{\partial}{\partial E} \left(-b(E)N_i(E) \right) - \frac{N_i(E)}{t_{\rm esc}} + \tilde{Q}_{ji}(E)$$

$$Q_{ji}(E_i) = \int dE_j N_j(E_j) \Gamma_j(E_j) \frac{dn_{j \to i}}{dE_i}(E_j, E_i)$$

$$\tilde{Q}_{ji}(E) = Q_i(E) + Q_{ji}(E)$$

- Production rate of particles of species *i* and energy *Ei* from the interactions or decay of the parent *j*
- After considering isotropy of the photon distribution, and calculating the quantities in the shock rest frame:

$$\begin{split} \Gamma_j(E_j) &= \int d\varepsilon \, n_\gamma(\varepsilon) \, f_j \, (y) \\ y &\equiv (E_j \varepsilon) / m_A \end{split} \begin{array}{l} \text{Escape rate of} \\ \text{the primary} \\ \text{particle} \end{split}$$

$$f_j(y) \equiv \frac{1}{2y^2} \int_{0}^{2y} d\epsilon_r \, \epsilon_r \, \sigma_j(\epsilon_r)$$

> All integrations need to be performed only once if the target photon density is constant over time → the interaction rate is only a function of energy

Data set used in the current work

Volume 17, number 1	PHYSICS LETTERS	15 June 1965	The nuclear gamma-ray absorption cross sec- tion of ⁴⁰ Ca has been measured with the 260 MeV electron synchrotron of the Levedev Physical In-	
NUCLEAR γ-R.	AY ABSORPTION CROSS SECTION OF ⁴⁰ The giant resonance region	Ca IN	9-channel pair magnetic spectrometer as a de- tector. The resolution of the spectrometer for	
B.S. DOLBILKIN, V.I. KORIN, L.E. LAZAREVA and F.A. NIKOLAEV P.N. Lebedev Physical Institute, Moscow, USSR 40Ca), 70.84 g/c1		γ -quanta of energy $E_{\gamma} = 20$ MeV was approxima- tely 220 keV. A block of natural calcium (96.97% 40 Ca), 70.84 g/cm ² thick, was used as absorber		
Yad.Fiz. 33, 581 (19	81)		Spectra of photoprotons from the nucleus Na-23 are	
B.S.Ishkhanov, I.M.I	Kapitonov, V.I.Shvedunov, A.I.Gutii,	A.M.Parlag	measured in the bremsstrahlung beam. Cross section of the reaction Na-23(gamma,p)Ne-22 with production the final nucleus in various states are obtained f the photoproton spectra.	

Investigation of the Reaction ${}^{23}Na(\gamma, p){}^{22}Ne$ with Production of the Final Nucleus in Various States

Total Photonuclear Cross Sections for Low Atomic Number Elements

J. M. Wyckoff, B. Ziegler, H. W. Koch, and R. Uhlig Phys. Rev. 137, B576 - Published 8 February 1965

Total photonuclear cross sections have been measured in an attenuation experiment using a scintillation pair spectrometer and an x-ray spectrum with a fixed maximum energy of 90 MeV. The cross sections as a function of x-ray photon energy for beryllium, carbon, oxygen, sodium, magnesium, aluminum, silicon, sulfur, calcium, nickel, cobalt, copper, and silver show detailed structure in many cases at x-ray energies of 15-30 MeV and display a consistent trend in shapes and magnitudes. The integrated cross sections up to 35 MeV relative to the classical dipole sum rule show a monotonic increase with atomic weight. Other analyses of the total photonuclear cross sections in terms of mean energies and of the ratios of the total cross sections to photoneutron cross sections are also presented.

Situation on experimental data and theoretical models

> We use the EXFOR database https://www-nds.iaea.org/exfor/exfor.htm

> No measurements of <u>absorption cross section</u> for the same isobar

Our current model:

- TALYS 1.8 is used with the strenght function strenght 1, based on a Kopecky-Uhl generalized Lorentzian model, as in Khan et al. paper
- TALYS is not recommended for A<12. For these nuclei we use a collection from CRPropa2 (Khampert et al, Astropart.Phys. 42 (2013) 41-51), based partially on data

What is TALYS? www.talys.eu

TALYS is software for the simulation of nuclear reactions. Many state-of-the-art nuclear models are included to cover all main reaction mechanisms encountered in light particle-induced nuclear reactions. **TALYS** provides a complete description of all reaction channels and observables, and is user-friendly.

Model predictions and parametrizations

 \rightarrow use of interpolated or fitted absorption cross sections where available, as done in PEANUT, ENDF-B-VII.1, JENDL/PD-2004

 \rightarrow use of parametrizations if cross sections are totally unknown

Other models:

> PSB model is obtained from Puget, Stecker and Bredekamp, Astrophys. J. 205, 638 (1976). Use of one nucleus for each mass; cross section for one and two nucleon emissions is approximated by a Gaussian in the low energy range and by a constant above 30 MeV. Threshold for reactions taken from Stecker and Salamon, Astrophys. J. 512 (1999). The list of nuclei has been slightly modified to be used in the current code for photodisintegration

> Box approximation is used in Murase and Beacom, Phys Rev. D81 2010

 $\sigma_{\rm GDR} \approx 1.45 \times 10^{-27} A \ {\rm cm}^2$

ENDF-B-VII.1 18 is an evaluated nuclear data library based on calculations using the GNASH code system. Its photo-nuclear part contains absorption cross-sections a sometimes inclusive emission spectra of neutrons and protons, but no residual cross-sections. Comparisons with data reveal a very good agreement with the measurements.

JENDL/PD-2004 [19] is another evaluated library, based on Lorentz fits at GDR energies and quasideuteron emission above. Elements without $\sigma_{\rm abs}$ measurements are evaluated through branching ratios from pre-quilibrium and evaporation models, together with photo-neutron data. The description of $\sigma_{\rm abs}$ is good for all measured elements.

Photo-disintegration cross sections: current situation

- Ca-40: double magic nucleus
- TALYS predictions not dependent on the element
- PEANUT predictions are different in the same isobar; if data available, at least the central GDR peak is reproduced
- Box approximation, used for example in Murase and Beacom, Phys Rev. D81 2010, underestimates data and models for A=40

[Boncioli, Fedynitch, Winter, Sci.Rep. 2017]

Photo-disintegration cross sections: current situation

[Boncioli, Fedynitch, Winter, Sci.Rep. 2017]

Photo-disintegration cross sections: current situation

Interpretation of UHECR data \rightarrow propagation model MODEL

- *SimProp* propagation
- PSB cross sections
- Gilmore EBL
- EPOS-LHC air interactions

	parameters
Rcut	18.68
gamma	0.96
Н	0.0
He	67.3
N	28.1
Si	4.6
Dmin	174.4/119

Interpretation of UHECR data \rightarrow propagation model MODEL

- SimProp propagation
- PSB cross sections
- Dominguez EBL
- EPOS-LHC air interactions

	parameters
Rcut	18.19
gamma	-1.02
Н	62.6
Не	36.8
N	0.6
Si	0.03
Dmin	187.0/119

Interpretation of UHECR data → propagation model MODEL

SimProp propagation

- TALYS cross sections
- Gilmore EBL

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• EPOS-LHC air interactions

	parameters
Rcut	18.62
gamma	0.77
Н	0.0
Не	7.0
N	85.2
Si	7.7
Dmin	175.9/119

[Auger Coll, JCAP 2017]

Interpretation of UHECR data \rightarrow air interaction models

[Auger Coll, JCAP 2017]

Source class I: Empty Cascade

- Low luminosity / large collision radii
- Only a few isotopes in the cascade populated relative to injected energy in primaries
- Maximum energy determined by adiabatic cooling, i.e. rigidity-dependent / Peters cycle
- Optically thin to photo-hadronic interactions of all
- Nuclei stay mostly intact and escape as CR

Source class II: Populated Cascade

- Intermediate luminosity / collision radii
- Cascade broadly populated along the main diagonal relative to injected energy
- Maximum energy determined by photo- hadronic processes, no Peters cycle!
- Optically thick to photo-hadronic interactions of heavy nuclei, still opt. thin to light nuclei
- Nuclei disintegrate partially

Source class III: Optically Thick Case

- High luminosity / small collision radii
- Cascade populated but more narrow, most of the energy is dumped into nucleons
- Maximum energy determined by photo- hadronic processes, no Peters cycle!
- Optically thick to photo-hadronic interactions of all
- Nuclei disintegrate very efficiently

- > Usual assumption:
 - GRB jet initially consists of free nucleons
 - nucleons can recombine into deuterium or alpha as the jet expands and cool but nuclei heavier than carbon cannot be produced
- ➤ The mechanisms to launch a relativistic jet is not yet completely understood → initial composition is not necessarily free nucleons
 - some observations of GRBs suggest that their relativistic jets were initially dominated by the magnetic field energy flux, being able to involve also heavy nuclei (A=56)
- In Shibata & Tominaga, arXiv:1503.03662 [astro-ph.HE], the jet is assumed to be due to falling matter during a relativistic jet-induced explosion. They adopt Wolf-Rayet stars as progenitors
- Other possibility: nucleosynthesis in from free nucleons in the central engine of GRBs (Metzger et al, 2011)